

# Isospin Symmetric Island of Inversion at the N=Z line

Duy Duc Dao, Frédéric Nowacki

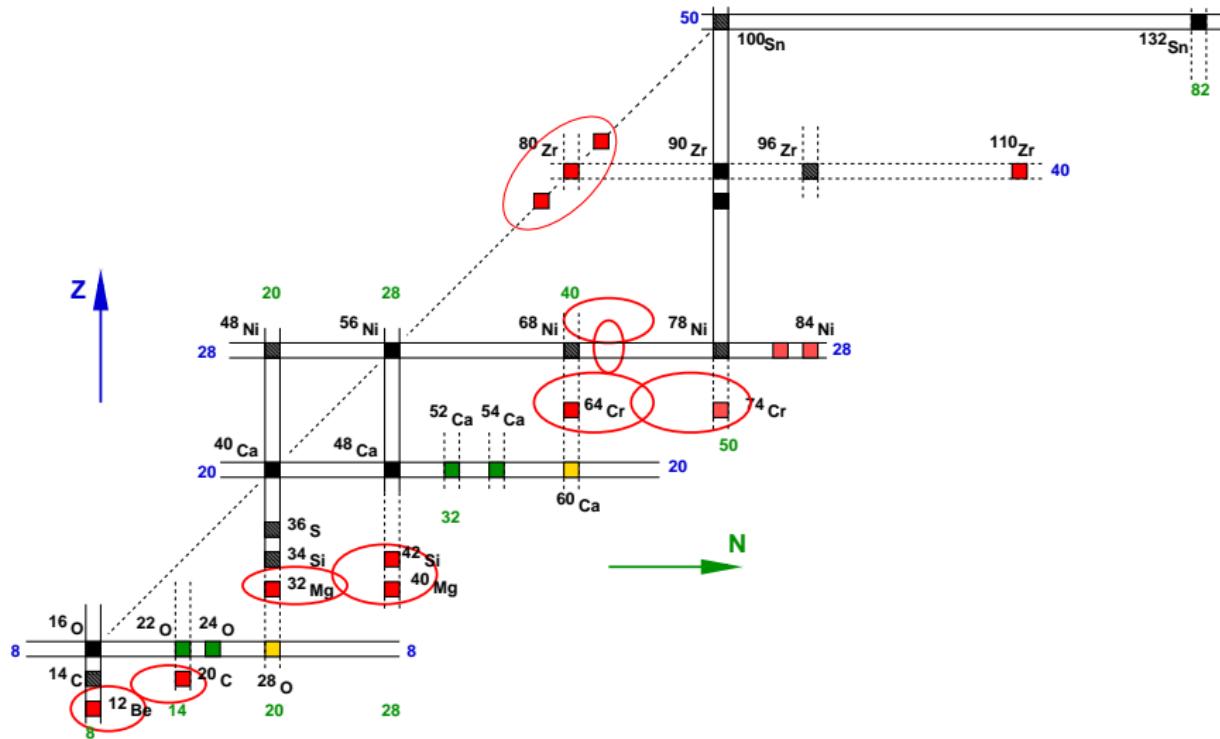


The 12th International Conference on Direct Reactions with  
Exotic Beams DREB2024

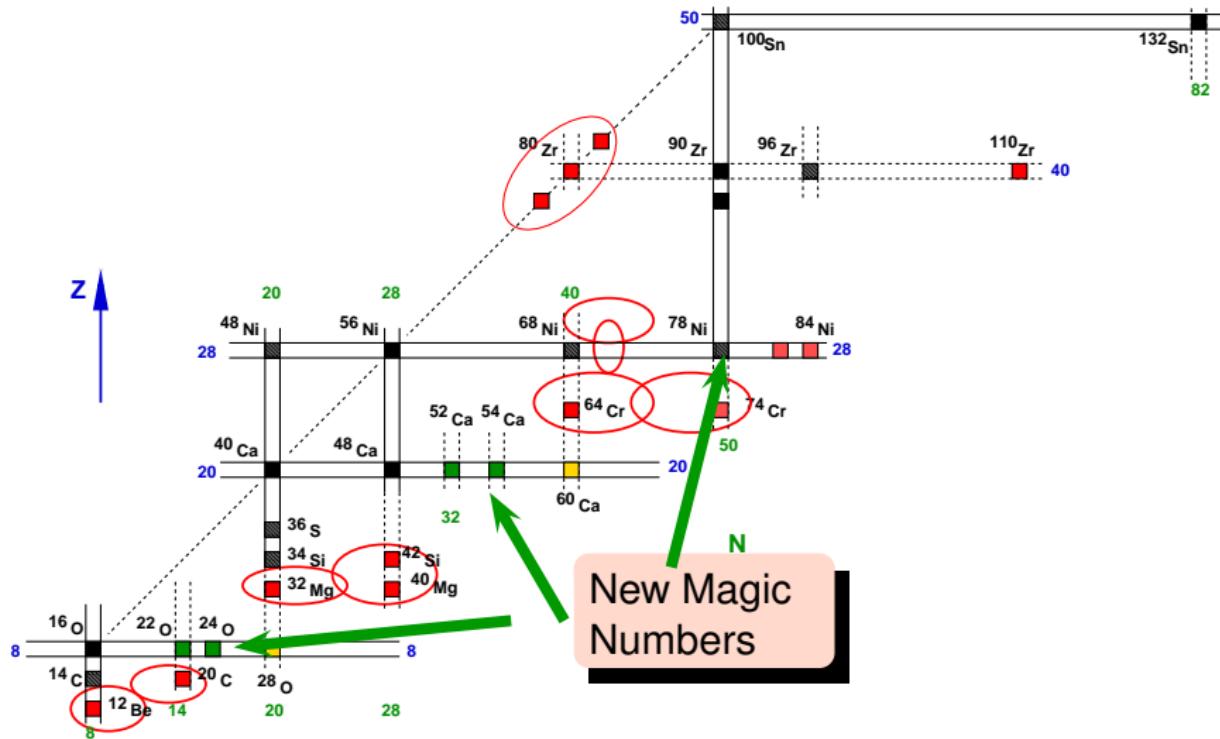
Wiesbaden, Germany, June 24th – 28th, 2024



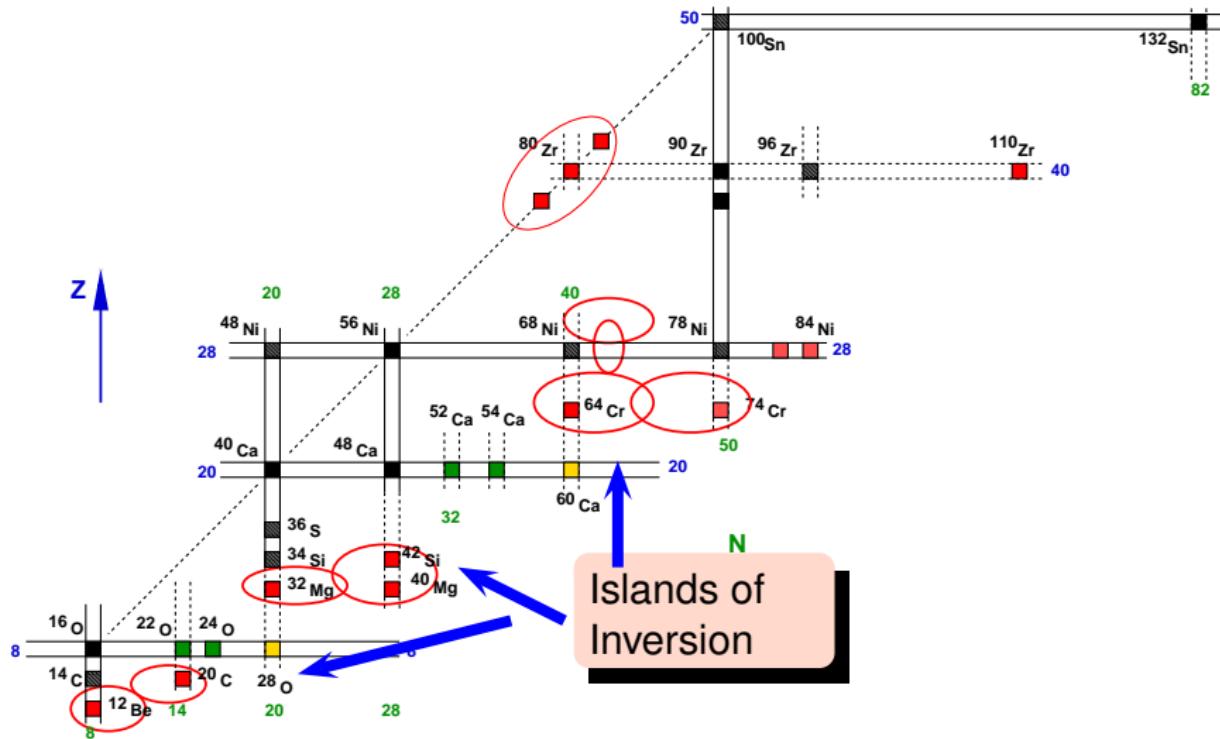
# Landscape of medium mass nuclei



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## UNDERSTANDING REGULARITIES for both SPHERICAL and DEFORMED systems

132Sn  
82

- New Magic Numbers:  $^{24}\text{O}$ ,  $^{48}\text{Ni}$ ,  $^{54}\text{Ca}$ ,  $^{78}\text{Ni}$ ,  $^{100}\text{Sn}$
- Vanishing of shell closures:  $^{12}\text{Be}$ ,  $^{32}\text{Mg}$ ,  $^{42}\text{Si}$ ,  $^{64}\text{Cr}$ ,  $^{80}\text{Zr}$  ...
- Island of deformation around  $A \sim 32$ ,  $A \sim 64$
- Low-lying dipole excitations in Ne, Ni isotopes

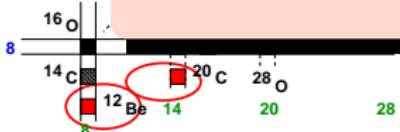
$Z$

- Variety of phenomena dictated by shell structure
- Close connection between collective behaviour and underlying shell structure
- 

$$\mathcal{H} = \mathcal{H}_m + \mathcal{H}_{\mathcal{M}}$$

Interplay between

- Monopole field (spherical mean field)
- Multipole correlations (pairing, Q.Q, ...)



# Effective Hamiltonian

## Monopole and multipole

Multipole expansion:

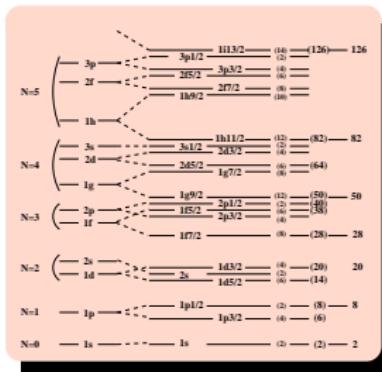
$$H = H_{\text{monopole}} + H_{\text{multipole}}$$

- Spherical mean-field

$H_{\text{monopole}}$ :

- Evolution of the spherical single particle levels

A. Poves and A. Zuker (Phys. Report 70, 235 (1981))



$H_{\text{multipole}}$ :

- Correlations
- Energy gains

- Pairing ( $SU_2$ ) semi-magic (n-n) (p-p)
- Quadrupole ( $SU_3/p$ - $SU_3/q$ - $SU_3$ ) p-n in H.O. or  $\Delta j = 2$

# Effective Hamiltonian

## Monopole and multipole

Multipole expansion:

- **Pairing regime: spherical nuclei**

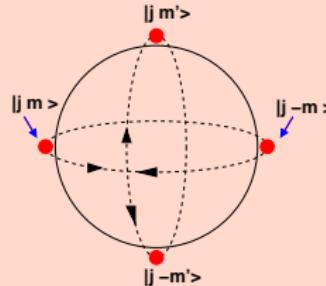
ground state = pairs of like-particles coupled at  $J=0$  (seniority  $v=0$ )

$2^+$  state (break of pair;  $v=2$ ) at high energy

$H_{monopole}$ :

superfluid nucleus:

A. Poves an



Typical example: **semi-magic Tin isotopes**

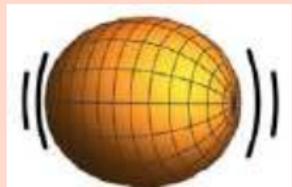
$H_{multipole}$ :

- **Quadrupole regime: deformed nuclei**

• *Pai*

prolate nucleus:

• *Qu*



M. Dufour a

Typical example: **open shell N=Z nuclei**

# Effective Hamiltonian

## Monopole and multipole

Multipole expansion:

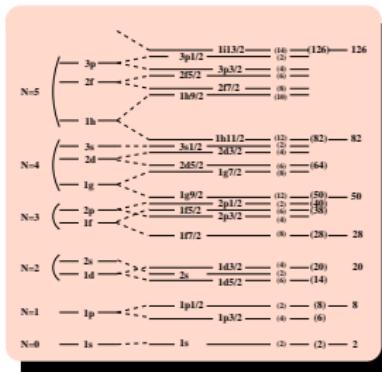
$$H = H_{\text{monopole}} + H_{PP} + H_{QQ}$$

- Spherical mean-field

$H_{\text{monopole}}$ :

- Evolution of the spherical single particle levels

A. Poves and A. Zuker (Phys. Report 70, 235 (1981))



$H_{\text{multipole}}$ :

- Correlations
- Energy gains

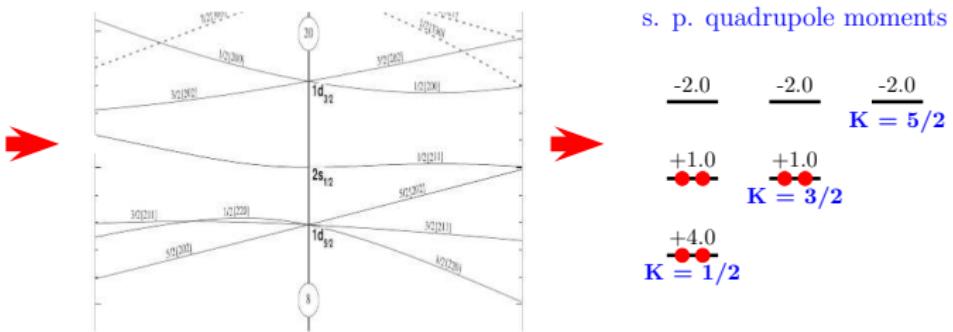
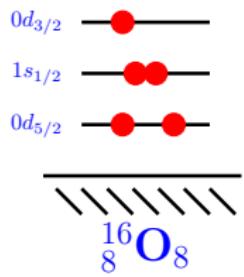
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- Quadrupole ( $SU3/p$ - $SU3/q$ - $SU3$ ) p-n in H.O. or  $\Delta j = 2$

# Nilsson-SU3 estimates

PHYSICAL REVIEW C **92**, 024320 (2015)

## Nilsson-SU3 self-consistency in heavy $N = Z$ nuclei

A. P. Zuker,<sup>1</sup> A. Poves,<sup>2,3</sup> F. Nowacki,<sup>1</sup> and S. M. Lenzi<sup>4</sup>



s. p. quadrupole moments

-2.0      -2.0      -2.0  
+1.0      +1.0      +1.0  
 **$K = 5/2$**

**$K = 3/2$**

+4.0

$$Q_0 = 2q^{20} = (2n_z - n_x - n - y)$$

# Island of Inversion at the N=Z line

## Strongly deformed states at $N = Z$ :

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- Most deformed cases for  $^{76}\text{Sr}$ ,  $^{80}\text{Zr}$
- New spectroscopy for  $^{84}\text{Mo}$  and  $^{86}\text{Mo}$

NSCL/GRETINA Experiment

R.D.O. Llewellyn *et al.*, PRL 124, 152501 (2020)

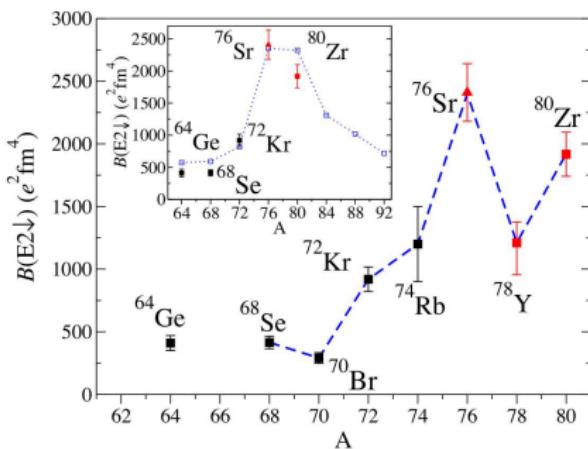
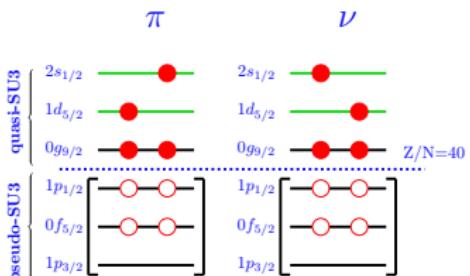


FIG. 3. Schematics of the  $B(E2\downarrow)$  values for the  $N = Z$  nuclei



- ZBM3 valence space:  
extension of JUN45  
to pseudo-SU3 + Quasi-SU3
- New effective interactions:
  - Realistic TBME + Monopole “3N” constraints
  - ab-initio N3LO (2N) interaction
  - ongoing ab-initio N3LO (2N) + 3N ( $\ln l$ ) interaction
- SM + DNO-SM for most deformed cases

# Discrete Non-Orthogonal Shell Model

**Generator Coordinate Method:**  $|\Psi_{\text{eff}}\rangle = \sum_i f_i |\Phi_i\rangle$

1) Deformed Hartree-Fock (HF) Slater determinants

2) Restoration of rotational symmetry

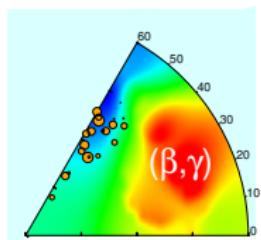
3) Mixing of shapes:

$$|\Psi_{\text{eff}}\rangle = \text{shape}_1 + \text{shape}_2 + \text{shape}_3 + \dots$$

## Intrinsic/Laboratory Description

- **Deformation structure of nuclear states:**  $\{J_\alpha^\pi\}$ ,  $q = (\beta, \gamma)$

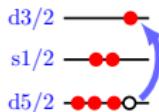
$$M_\alpha^{(J)}(q, K) = \sum_{q', K'} [\hat{N}^{1/2}]_{K' K}^{(J)}(q', q) f_\alpha^{(J)}(q', K')$$



- ◊ Probability of a configuration  $(\beta, \gamma)$ :

$$P_\alpha^{(J)}(q) = \sum_K |M_\alpha^{(J)}(q, K)|^2$$

- **particle-hole interpretation:**



M-scheme

- **K-quantum numbers:**

$$P_\alpha^{(J)}(K) = \sum_q |M_\alpha^{(J)}(q, K)|^2$$

# Discrete Non-Orthogonal Shell Model

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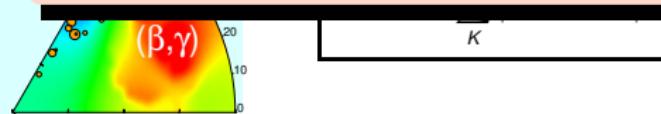
PHYSICAL REVIEW C 105, 054314 (2022)

- Nuclear structure within a discrete nonorthogonal shell model approach: New frontiers

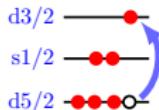
D. D. Dao and F. Nowacki

Université de Strasbourg, CNRS, IPHC UMR7178, 23 rue du Loess, F-67000 Strasbourg, France

(Received 8 March 2022; accepted 6 May 2022; published 23 May 2022)



- **particle-hole interpretation:**



M-scheme

- **K-quantum numbers:**

$$P_{\alpha}^{(J)}(K) = \sum_q |M_{\alpha}^{(J)}(q, K)|^2$$

# Recent developments of the DNO shell model

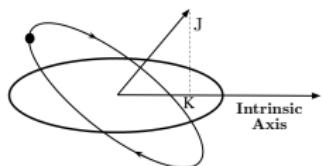
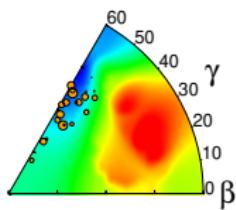
## ◇ Variation-After-Projection DNO-SM approach

$$\mathcal{H}_{\text{eff}}|\Psi_{\alpha}^{JM}\rangle = E_{\alpha}^{(J)}|\Psi_{\alpha}^{JM}\rangle \implies \delta \frac{\langle \Psi_{\alpha}^{JM} | \mathcal{H}_{\text{eff}} | \Psi_{\alpha}^{JM} \rangle}{\langle \Psi_{\alpha}^{JM} | \Psi_{\alpha}^{JM} \rangle} = 0, \quad |\Psi_{\alpha}^{JM}\rangle = \sum_{q,K} f_{\alpha}^{(J)}(q, K) \mathcal{P}_{MK}^J |\Phi(q)\rangle$$

Double variation AFTER Angular Momentum Projection: Mixing coefficient

Slater state

## ◇ DNO-SM( $\beta, \gamma$ )

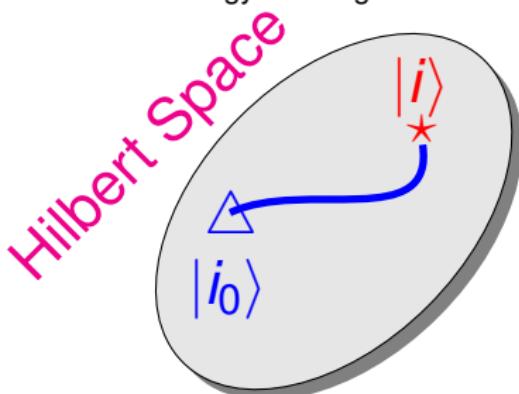


$$|\Psi_{\text{eff}}\rangle = \bigoplus_{q=1}^{\infty} \bigoplus_{J_{\alpha}^{\pi}=0_1^+, \dots} \bigoplus_{K} |\Psi_{\alpha}^{JM}\rangle$$

(D.D. Dao and F. Nowacki, PRC 105, 054314 (2022))

## ◇ DNO-SM(VAP)

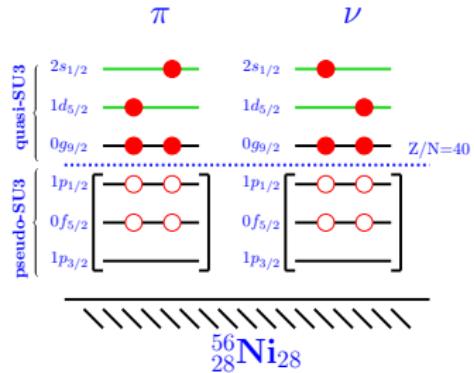
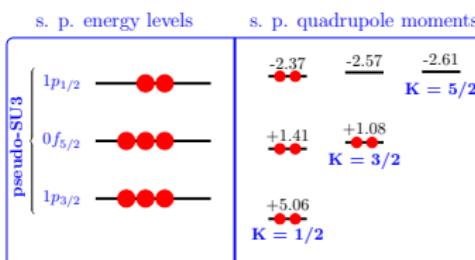
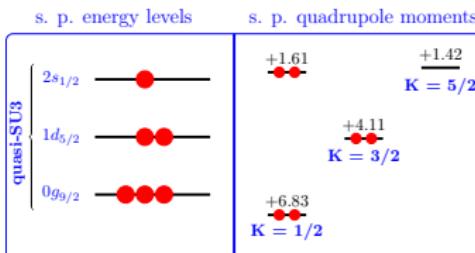
- $q = 1, 2, 3, \dots$
- $J_{\alpha}^{\pi} = 0_1^+, \dots$
- Best energy-favoring Slater states



# Island of Inversion at the N=Z line

## Strongly deformed states at $N = Z$

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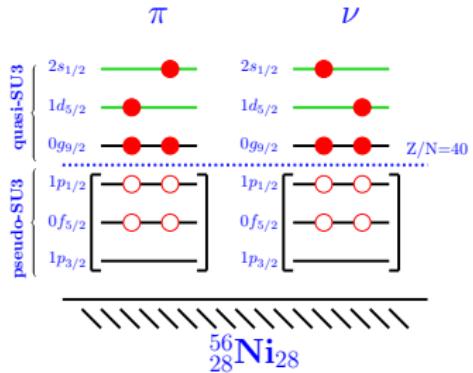
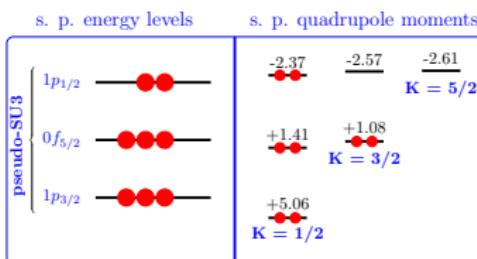
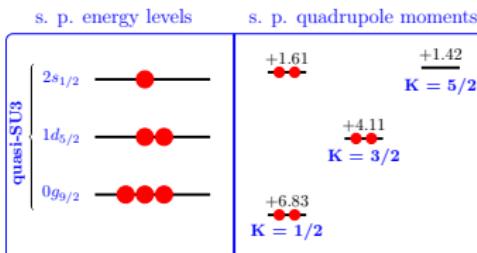


nucleus	Np-Nh*	ZRP	PHF	B(E2)(e <sup>2</sup> .fm <sup>4</sup> )	Exp.	DNO-SM*	SM
$^{84}\text{Mo}$	4p-4h	1104	1193	$1740^{+580}_{-430}$	$1740^{+580}_{-430}$	1765	-
	8p-8h	1891	1732				
$^{86}\text{Mo}$	0p-0h	542	196				
	2p-2h	1030	871				
	4p-4h	1416	1179				
	6p-6h	1858	1655				

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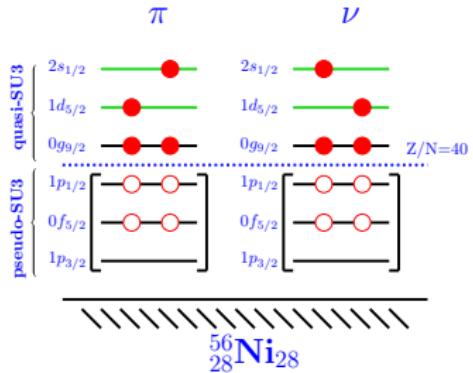
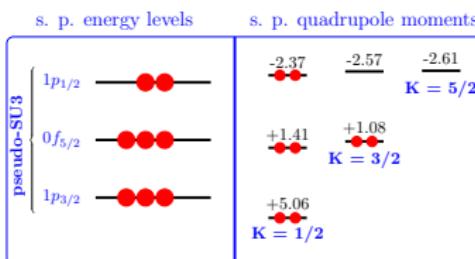
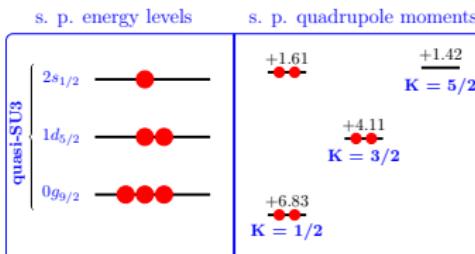


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						<b>707(71)</b>	980 731

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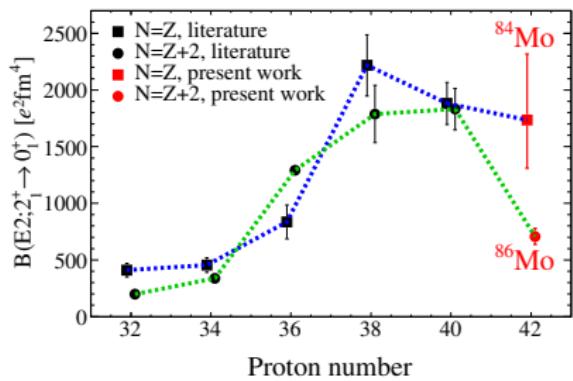
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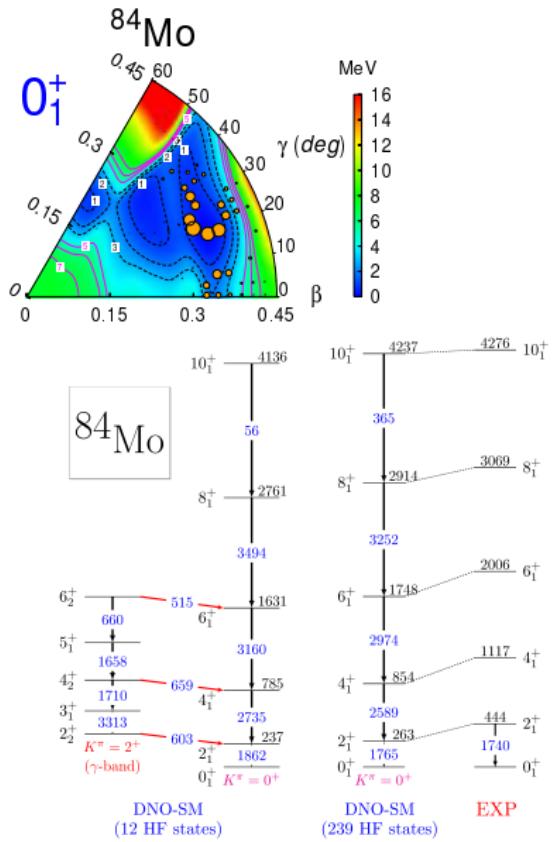
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NSCL/GRETINA Experiment



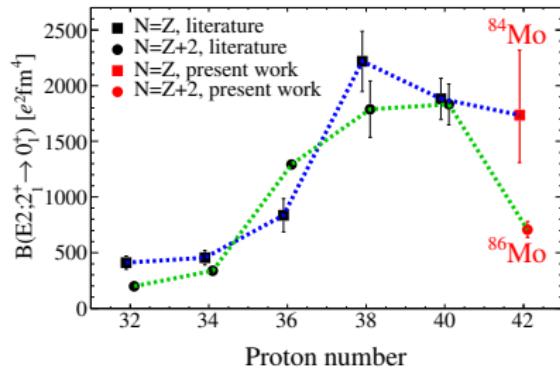
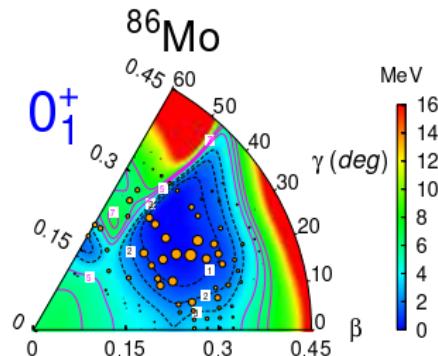
J. Ha, F. Recchia et al., submitted to NATURE



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## Three-Body Forces and the Limit of Oxygen Isotopes

Takaharu Otsuka,<sup>1,2,3</sup> Toshio Suzuki,<sup>4</sup> Jason D. Holt,<sup>5</sup> Achim Schwenk,<sup>5</sup> and Yoshinori Akaishi<sup>6</sup><sup>1</sup>*Department of Physics, University of Tokyo, Hongo, Tokyo 113-0033, Japan*<sup>2</sup>*Center for Nuclear Study, University of Tokyo, Hongo, Tokyo 113-0033, Japan*<sup>3</sup>*National Superconducting Cyclotron Laboratory, Michigan State University, East Lansing, Michigan, 48824, USA*<sup>4</sup>*Department of Physics, College of Humanities and Sciences, Nihon University, Sakurajosui 3, Tokyo 156-8550, Japan*<sup>5</sup>*TRIUMF, 4004 Wesbrook Mall, Vancouver, BC, V6T 2A3, Canada*<sup>6</sup>*RIKEN Nishina Center, Hirosawa, Wako-shi, Saitama 351-0198, Japan*

(Received 17 August 2009; published 13 July 2010)

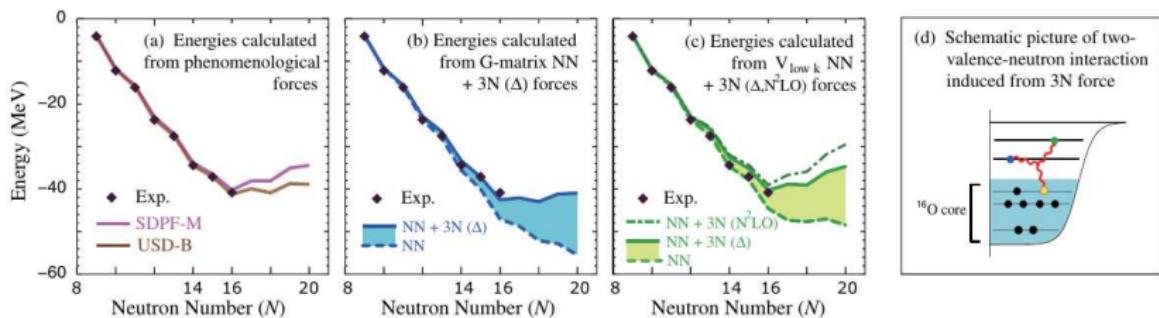


FIG. 4 (color online). Ground-state energies of oxygen isotopes measured from  $^{16}\text{O}$ , including experimental values of the bound 16–24 O. Energies obtained from (a) phenomenological forces SDPF-M [13] and USD-B [14], (b) a  $G$  matrix and including FM 3N forces due to  $\Delta$  excitations, and (c) from low-momentum interactions  $V_{\text{low } k}$  and including chiral EFT 3N interactions at  $N^2\text{LO}$  as well as only due to  $\Delta$  excitations [25]. The changes due to 3N forces based on  $\Delta$  excitations are highlighted by the shaded areas. (d) Schematic illustration of a two-valence-neutron interaction generated by 3N forces with a nucleon in the  $^{16}\text{O}$  core.

## Evolution of Shell Structure in Neutron-Rich Calcium Isotopes

G. Hagen,<sup>1,2</sup> M. Hjorth-Jensen,<sup>3,4</sup> G. R. Jansen,<sup>3</sup> R. Machleidt,<sup>5</sup> and T. Papenbrock<sup>1,2</sup>

<sup>1</sup>*Physics Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA*

<sup>2</sup>*Department of Physics and Astronomy, University of Tennessee, Knoxville, Tennessee 37996, USA*

<sup>3</sup>*Department of Physics and Center of Mathematics for Applications, University of Oslo, N-0316 Oslo, Norway*

<sup>4</sup>*National Superconducting Cyclotron Laboratory and Department of Physics and Astronomy,*

*Michigan State University, East Lansing, Michigan 48824, USA*

<sup>5</sup>*Department of Physics, University of Idaho, Moscow, Idaho 83844, USA*

(Received 16 April 2012; published 17 July 2012)

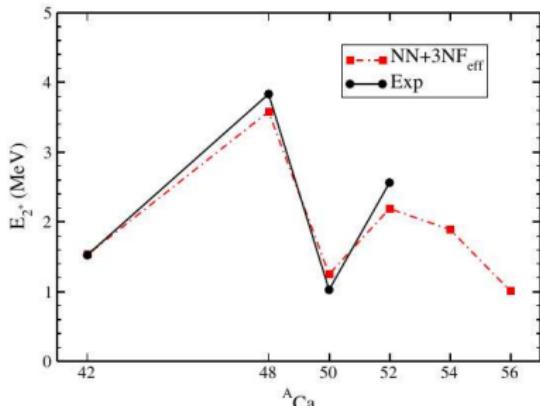


FIG. 2 (color online). (Excitation energies of  $J^\pi = 2^+$  states in the isotopes  $^{42,48,50,52,54,56}\text{Ca}$  (experiment: black circles, theory: red squares))

# Shell closures and 2N forces only

PHYSICAL REVIEW C 74, 061302(R) (2006)

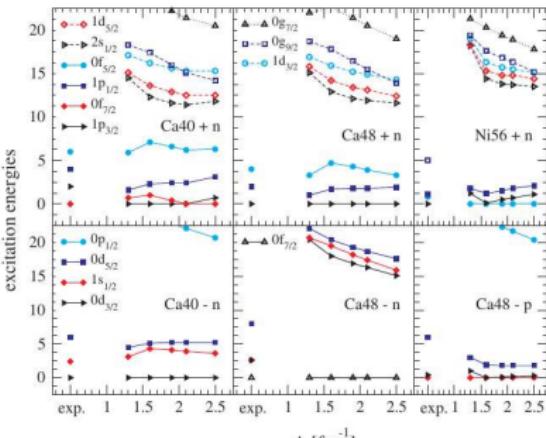
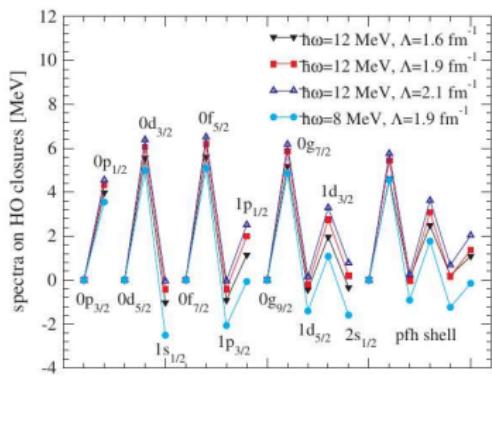
## Shell-model phenomenology of low-momentum interactions

Achim Schwenk<sup>1,\*</sup> and Andrés P. Zuker<sup>2,†</sup>

<sup>1</sup>Nuclear Theory Center, Indiana University, 2401 Milo B. Sampson Lane, Bloomington, Indiana 47408, USA

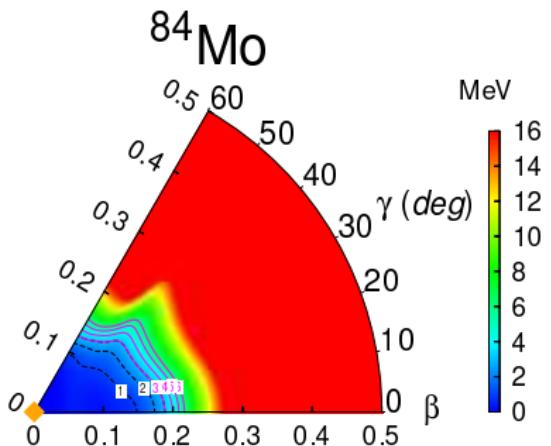
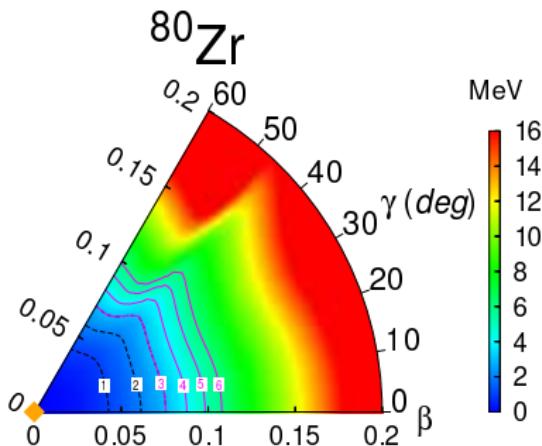
<sup>2</sup>Institut de Recherches Subatomiques, IN2P3-CNRS, Université Louis Pasteur, F-67037 Strasbourg, France

(Received 14 January 2005; revised manuscript received 20 September 2006; published 12 December 2006)



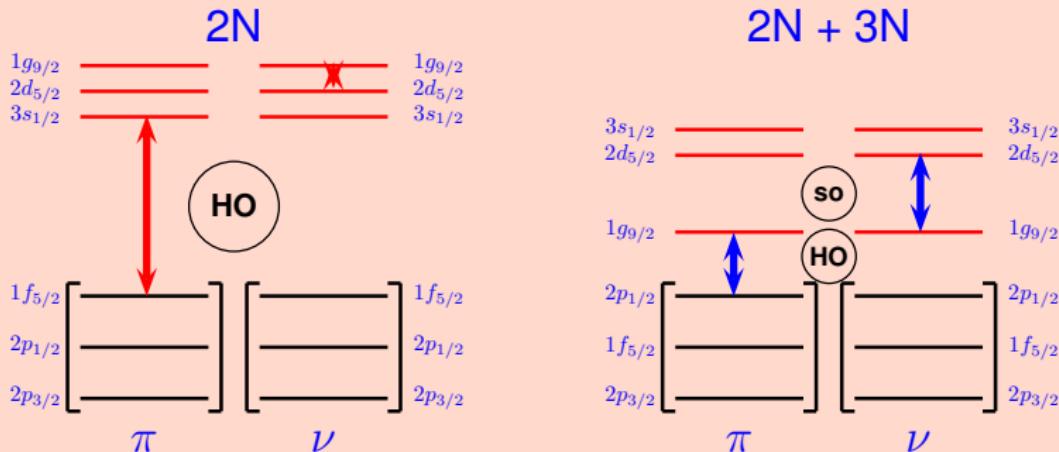
- no Spin-orbit shell closures in  $^{12}\text{C}$ ,  $^{22}\text{O}$ ,  $^{48}\text{Ca}$ ,  $^{56}\text{Ni}$
- too strong H. O. shell closures  $^{16}\text{O}$ ,  $^{40}\text{Ca}$ , ... and  $^{80}\text{Zr}$  !!!

# N3LO NN calculations



nucleus	NpNh*	B(E2)(e <sup>2</sup> .fm <sup>4</sup> )				N3LO
		ZRP	PHF	Exp.	DNO-SM	
$^{80}\text{Zr}$	4p-4h	587	637			
	8p-8h	1713	1509	<b>1910(180)</b>	2325	
	12p-12h	2663	2396			0.03
$^{84}\text{Mo}$	4p-4h	1104	1193	<b>1740<sup>+580</sup><sub>-430</sub></b>	1740	174
	8p-8h	1891	1732			

# N3LO NN calculations



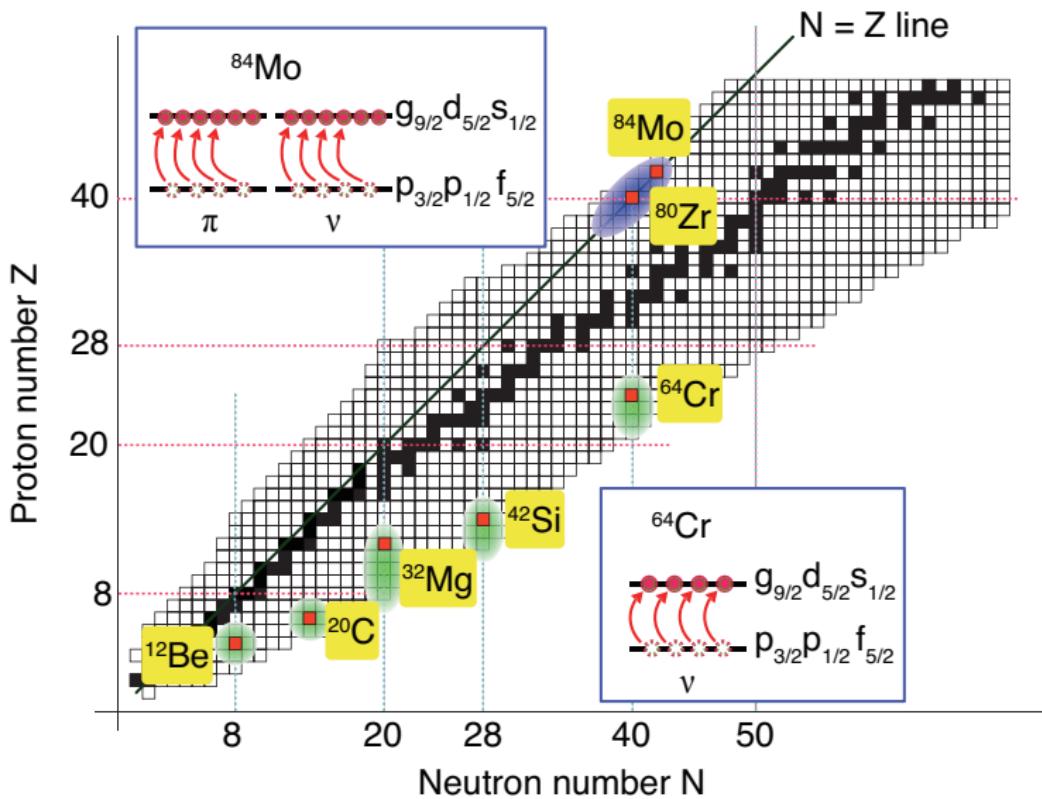
Three body forces and persistence of spin-orbit shell gaps in medium-mass nuclei: Towards the doubly magic  $^{78}\text{Ni}$ ,

K. Sieja, F. Nowacki

Phys. Rev. C85, 051301(R) (2012)

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	8p-8h	1891	1732				

# Isospin Symmetric Island of Inversion



# Summary

- Monopole drift develops in all regions but the Interplay between correlations (pairing + quadrupole) and spherical mean-field (monopole field) determines the physics.
- New “island of inversion” or “island of deformation” present for neutron-rich systems show up also at N=Z line with very deformed rotors dominated by Many-particles-Many-holes configurations.
- New spectroscopy for  $^{84}\text{Mo}$  and  $^{86}\text{Mo}$  and first fingerprint of 3N forces in deformed systems
- Around A~ 80, an “island of enhanced collectivity” show very deformed rotors dominated by Many-particles-Many-holes configurations.
- Ongoing NN + 3N( $\text{lnl}$ ) ab-initio calculations

# Summary

Special thanks to:

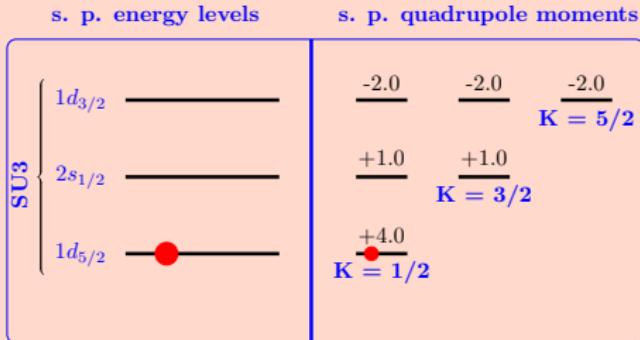
- D. D. Dao
- G. Martinez-Pinedo, A. Poves, S. Lenzi
- A. Gade, O. Sorlin, A. Obertelli

# Nilsson-SU3 estimates

PHYSICAL REVIEW C **92**, 024320 (2015)

## Nilsson-SU3 self-consistency in heavy $N = Z$ nuclei

A. P. Zuker,<sup>1</sup> A. Poves,<sup>2,3</sup> F. Nowacki,<sup>1</sup> and S. M. Lenzi<sup>4</sup>



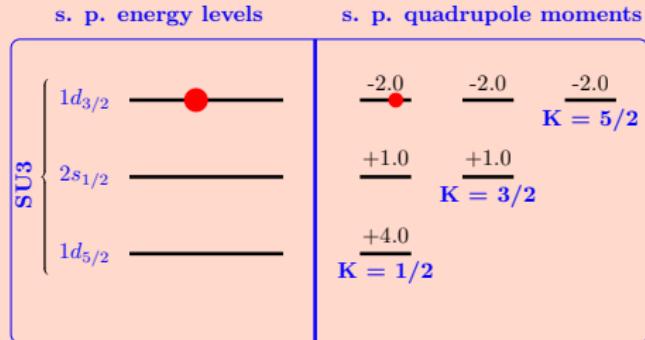
$$Q_0 = 2q^{20} = (2n_z - n_x - n - y)$$

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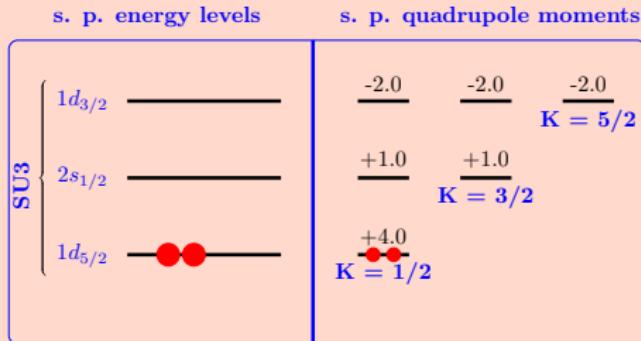
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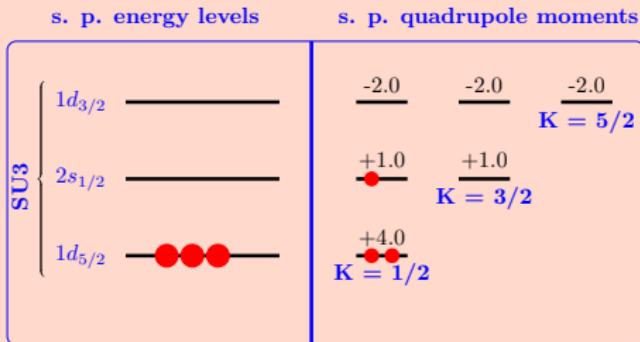
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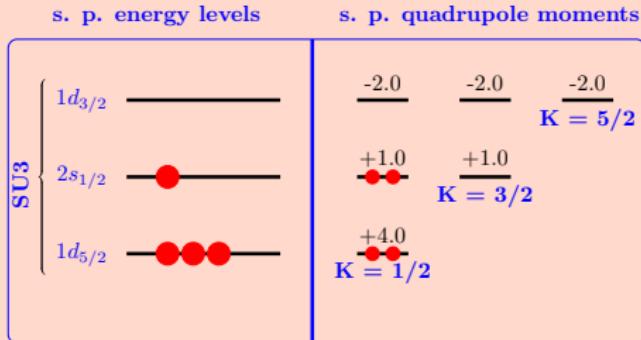
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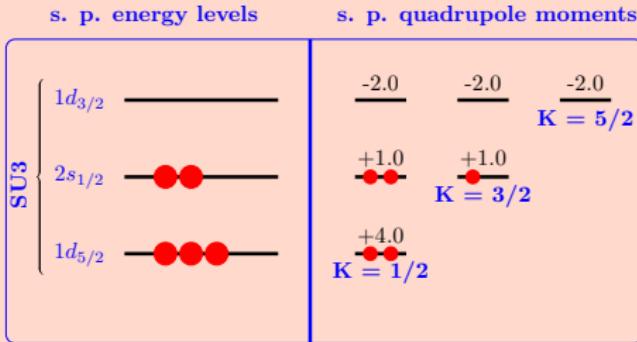
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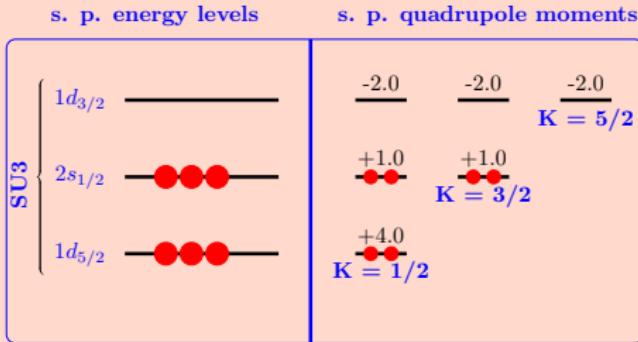
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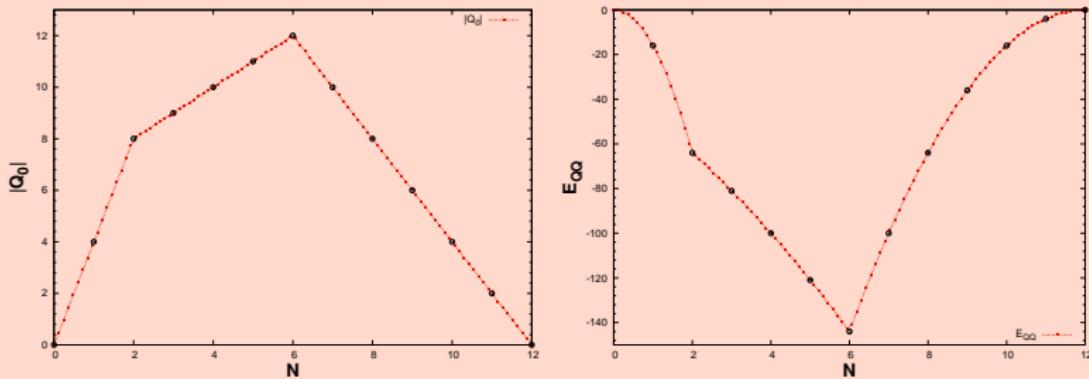
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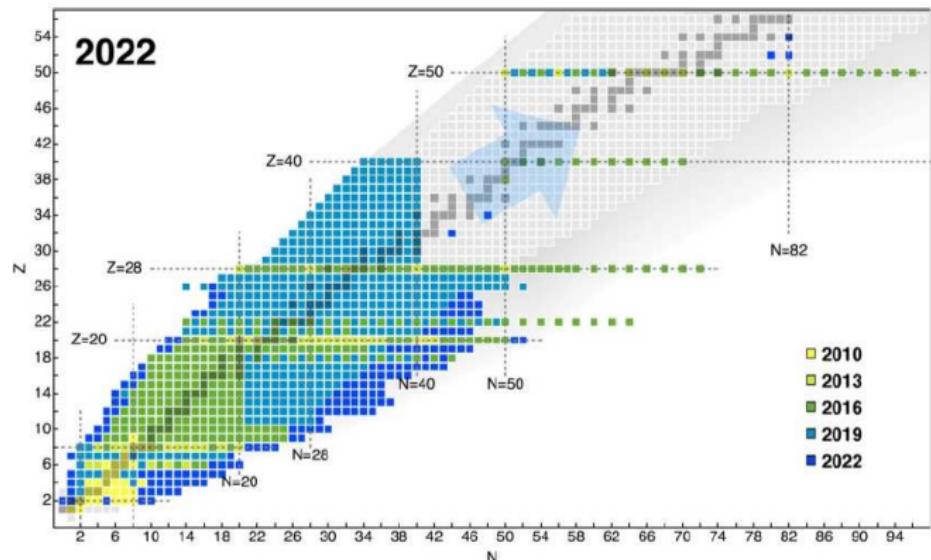
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# Ab-initio predictions ?



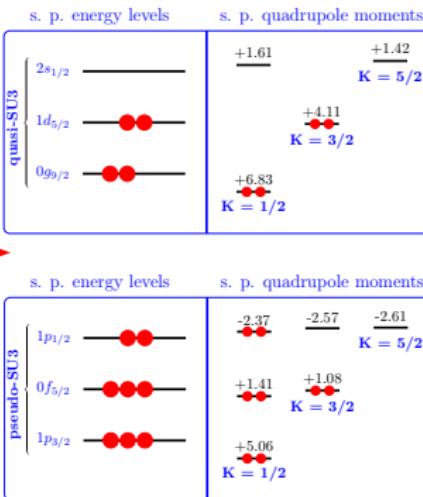
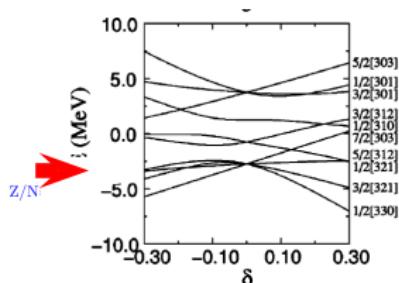
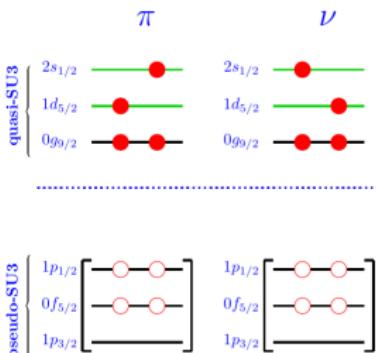
## Ab Initio Progress: How Heavy Can We Go?

Tremendous progress in ab initio reach, largely due to polynomially scaling methods!



# Nilsson-SU3 estimates

single particle energy levels

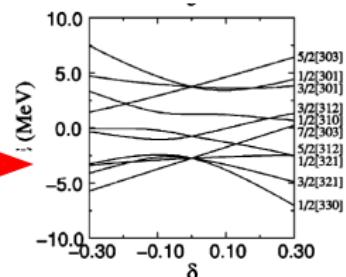
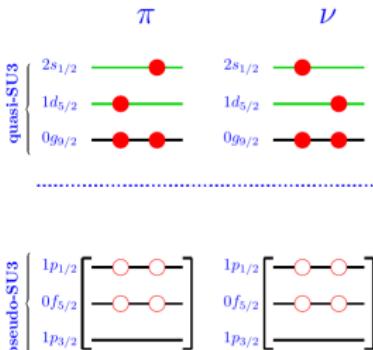


56Ni<sub>28</sub>

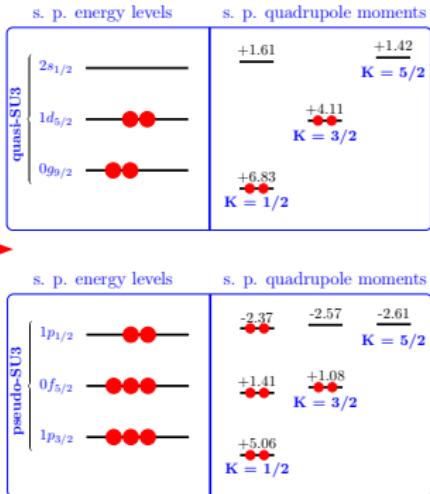
nucleus	NpNh*	ZRP	PHF	B(E2)(e <sup>2</sup> .fm <sup>4</sup> )	Exp.	DNO-SM
<sup>76</sup> Sr	4p-4h	924	806			
	8p-8h	2189	2101		<b>2390(240)</b>	1847
	12p-12h	2316	2300			
<sup>80</sup> Zr	4p-4h	587	637		<b>1910(180)</b>	2325
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# Nilsson-SU3 estimates

single particle energy levels



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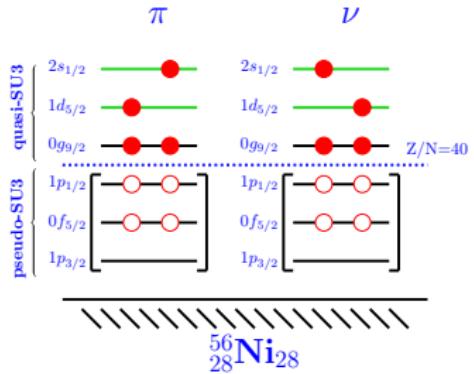
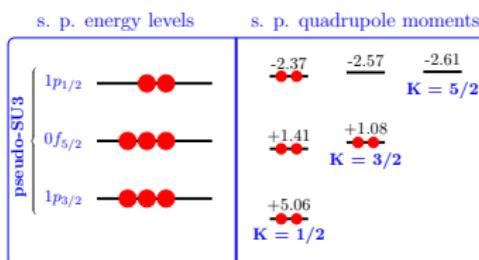
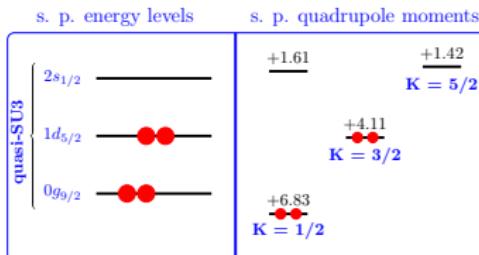


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# Island of Inversion at the N=Z line

## Strongly deformed states at $N = Z$

- Configuration mixing in  $^{72}\text{Kr}$
  - Most deformed cases for  $^{76}\text{Sr}$ ,  $^{80}\text{Zr}$
  - New spectroscopy for  $^{84}\text{Mo}$  and  $^{86}\text{Mo}$
- NSCL/GRETINA Experiment**



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**NSCL/GRETINA Experiment**

R.D.O. Llewellyn *et al.*, Phys. Rev. Lett. **124**, 152501 (2020)

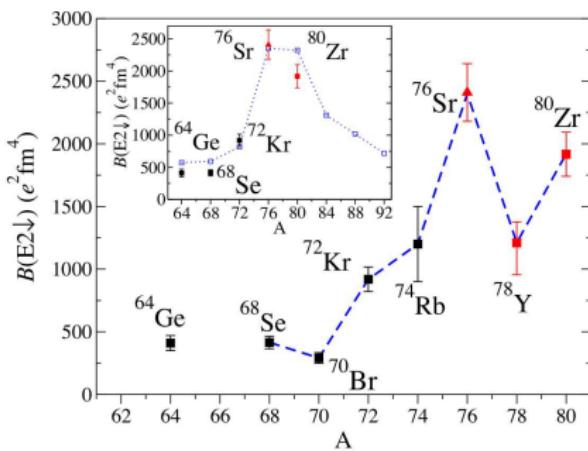
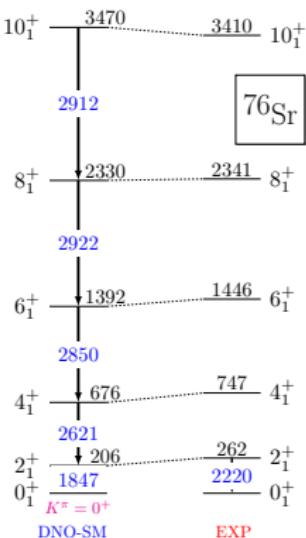
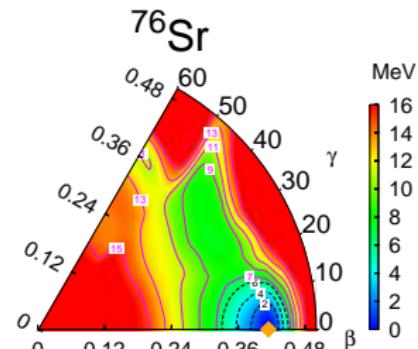


FIG. 3. Schematics of the  $B(E2\downarrow)$  values for the  $N = Z$  nuclei



(20 KE states)

EXP

# Development of deformation at N=8,20,40,70

F. Nowacki, A. Obertelli and A. Poves

Progress in Particle and Nuclear Physics 120 (2021) 103866

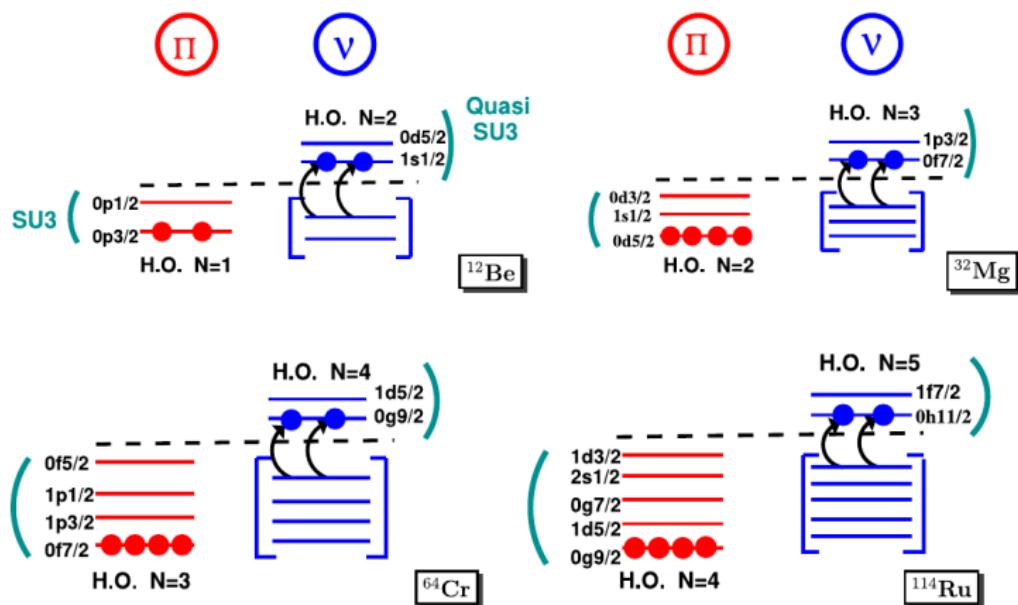
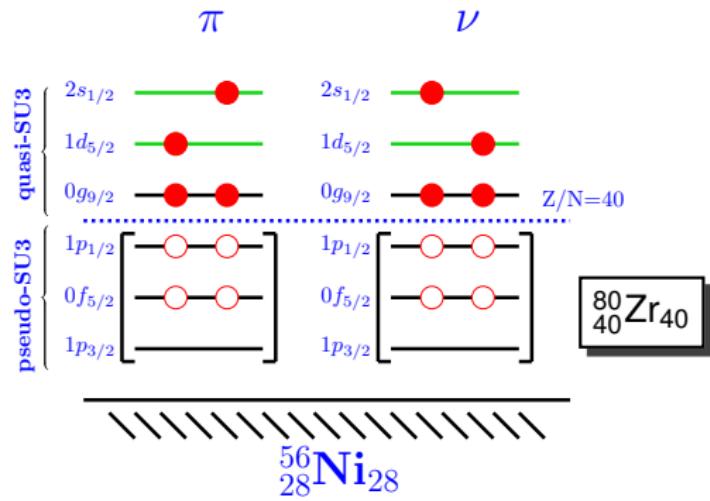


Fig. 40. Schematic view of the valence spaces at N = 8, 20, 40 and 70. The intruder configurations that develop quadrupole collectivity are highlighted.

# N=40 at N=Z



- p shell:  $^{16}\text{O}$   
spherical/doubly magic
- sd shell:  $^{40}\text{Ca}$   
spherical/doubly magic
- pf shell:  $^{80}\text{Zr}$   
deformed nucleus

- Low-lying states in H.O. N=Z=8: CS , 4p4h, 8p8h
- Low-lying states in H.O. N=Z=20: CS , 4p4h, 8p8h
- Low-lying states in H.O. N=Z=40: 4p4h ? 8p8h ? 12p12h ?

# $^{88}\text{Ru}$ : boundary of the N=Z island of inversion

$^{88}\text{Ru}$ , DNP-ZBM3 effective interaction

dimension	$(\beta, \gamma)$	$(\beta, \gamma) + np-nh$	DNO-SM(VAP)	SM
	39	207	10	$\sim 2 \times 10^{12}$
$E(0_{gs}^+)$ (MeV)	-416.465	-418.029	-418.244	-419.780
$E^*(2_1^+)$ (MeV)	0.25	—	0.30	0.49
BE(2) ( $e^2 \cdot \text{fm}^4$ )	1322	—	979	635

(D. D. Dao, F. Nowacki, A. Poves, *Isospin-symmetric island of inversion at the N=Z line*)

