

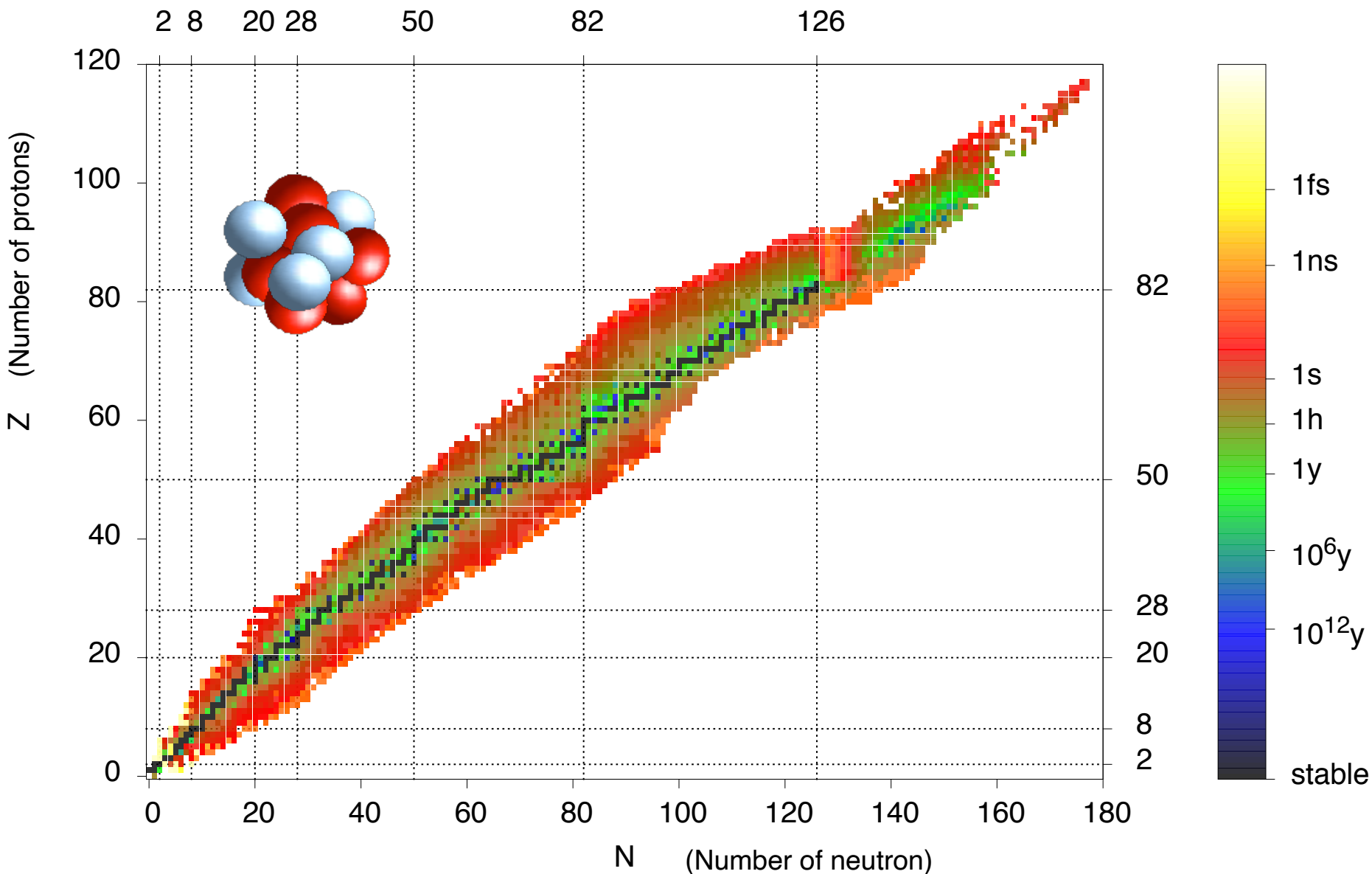
Time dependent dynamics of nuclear many-body states

Alexander Volya
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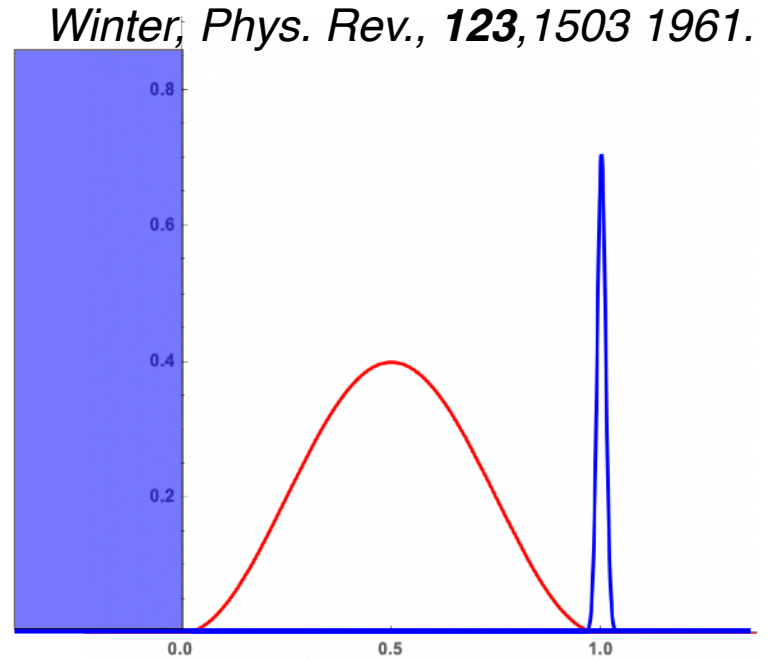
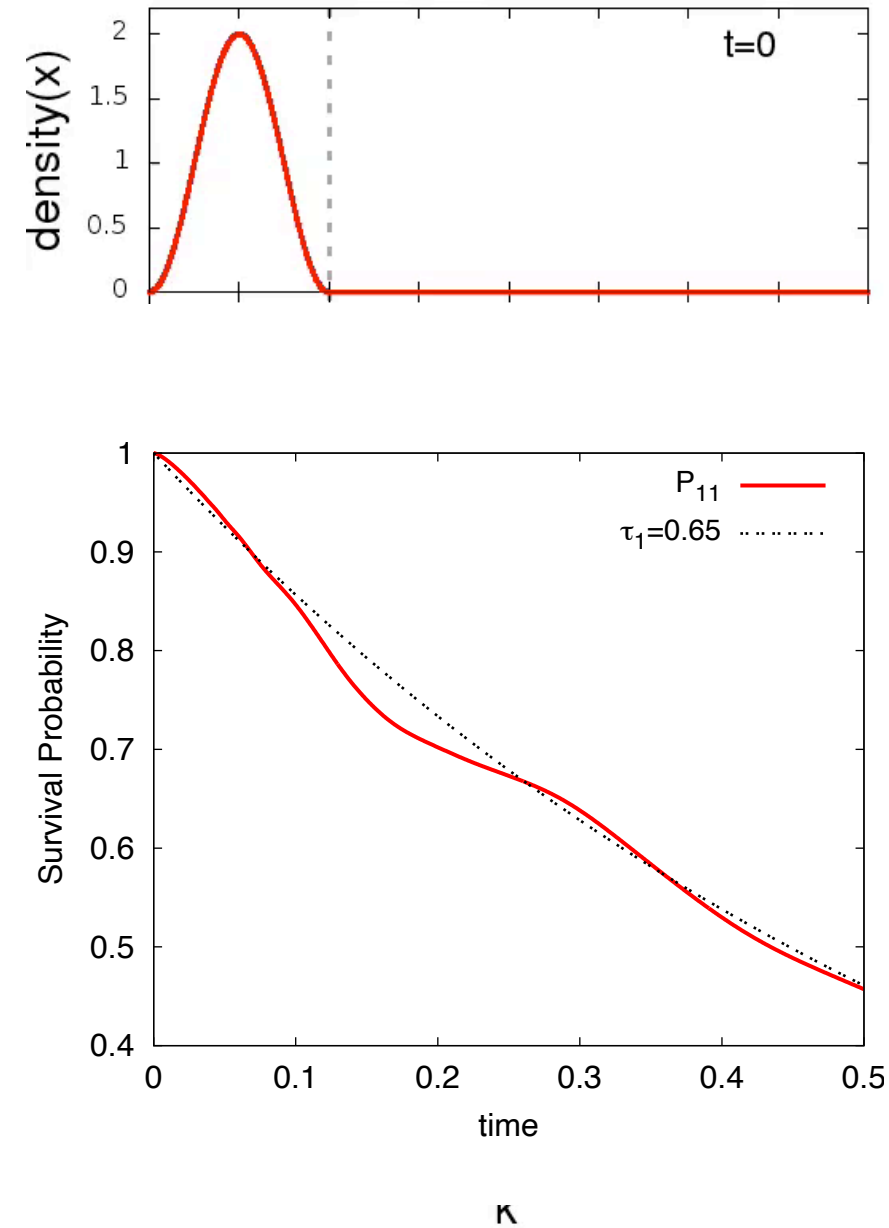
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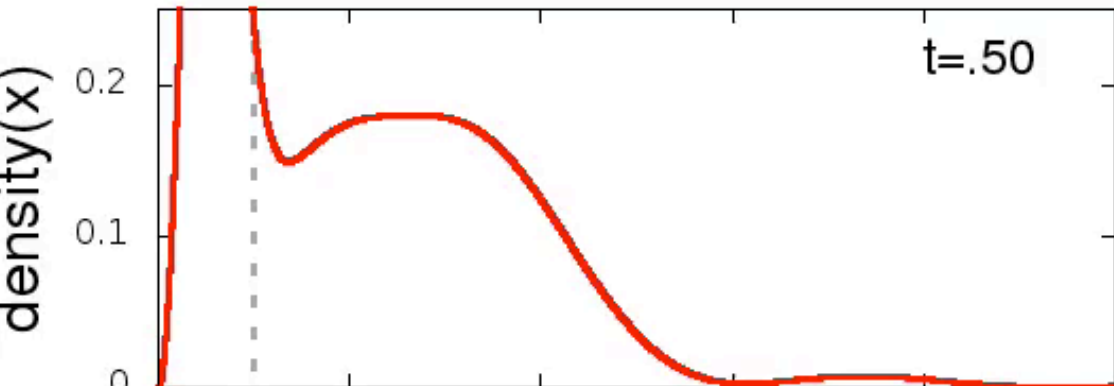


Atomic nucleus and open quantum system

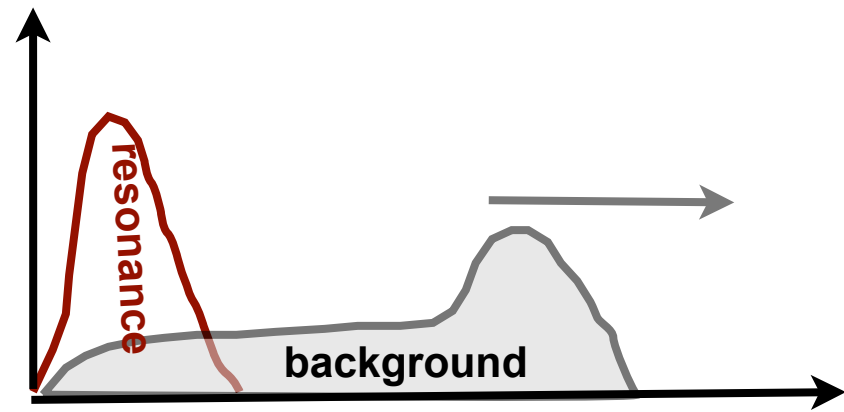
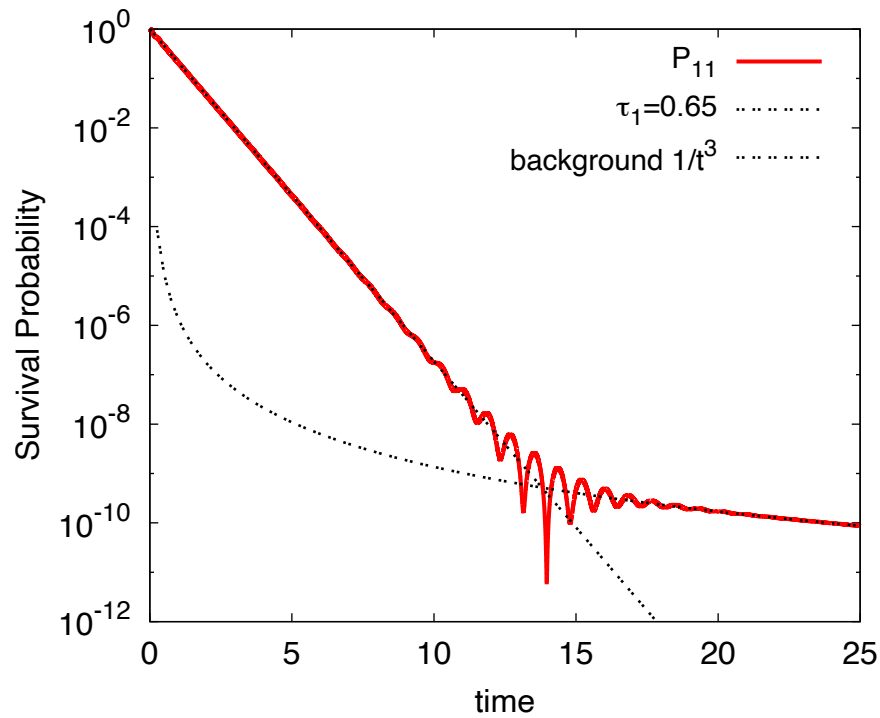


Time dependence of decay





Winter's model: Dynamics at remote times



Internal dynamics in decaying system

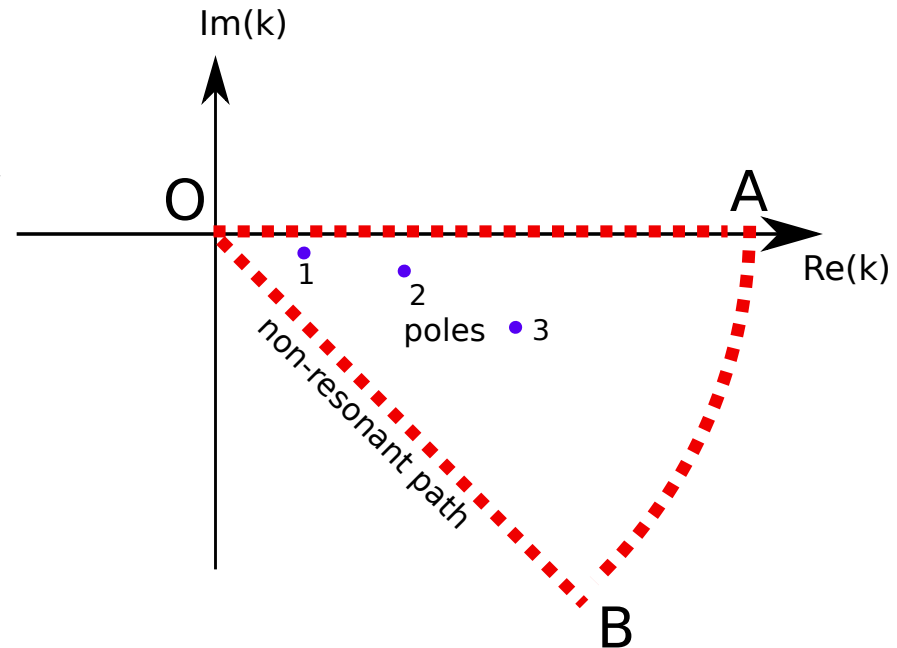
Winter's model

$$M_{nn'}(t) = \langle n | e^{-iHt} | n' \rangle = \int_0^\infty e^{-ik^2 t} \langle n | k \rangle \langle k | n' \rangle dk.$$

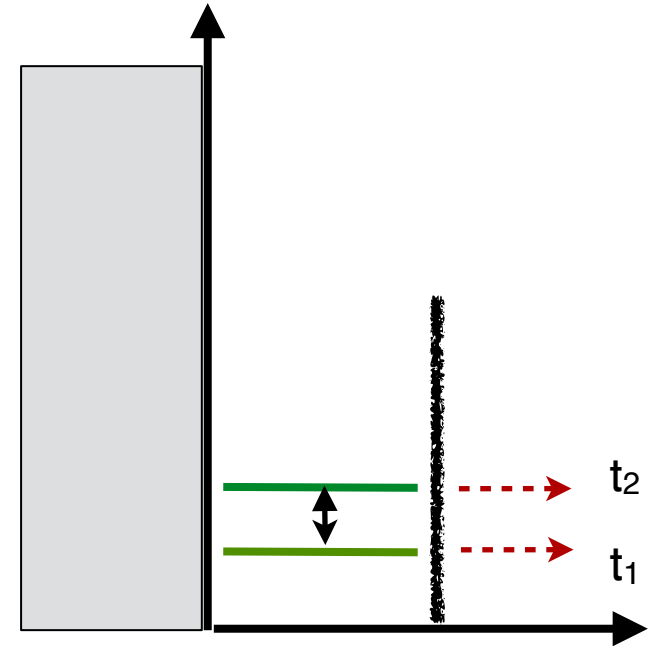
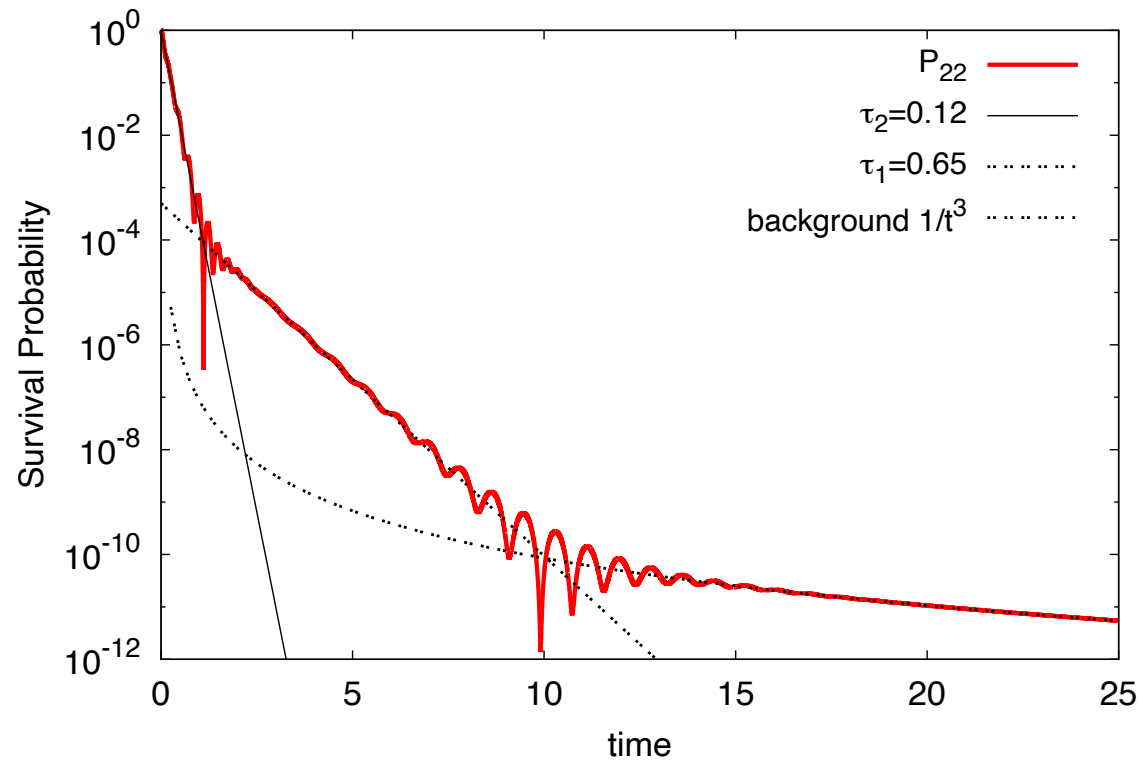
$$M_{nn'}(t) = \sum_{k_r} M_{nn'}^{(r)} e^{-ik_r^2 t} + M_{nn'}^{(NR)}(t)$$

$$M_{nn'}^{(r)}(t) = -2\pi i \text{Res}(k_r)$$

$$M_{nn'}^{(NR)}(t) = \frac{1+i}{\pi^{5/2} \sqrt{2}} (1+G)^2 \frac{1}{nn'} \frac{1}{t^{3/2}}.$$

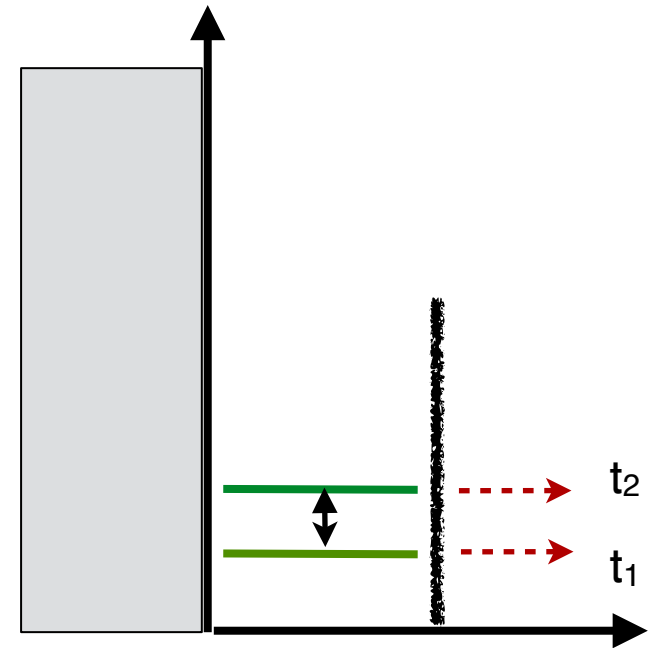
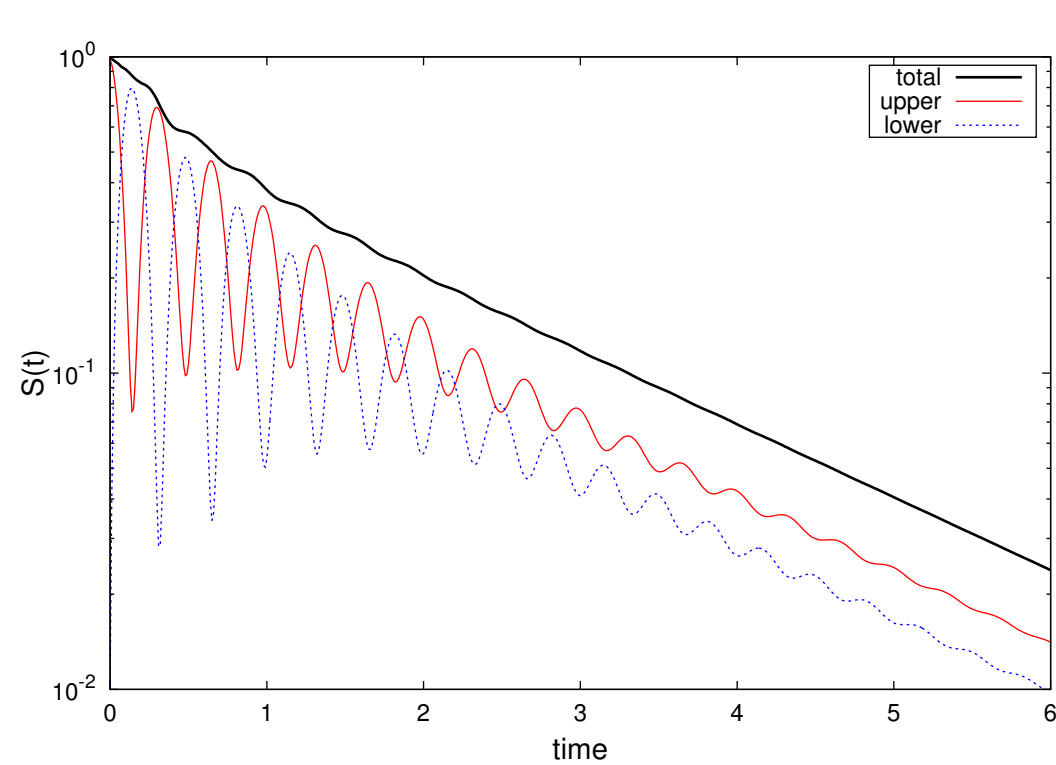


Internal dynamics in decaying system Winter's model



M. Peshkin, A. Volya, and V. Zelevinsky, “Non-exponential and oscillatory decays in quantum mechanics,” *EPL*, vol. 107, no. 4, p. 40001, 2014.

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Time-dependent Continuum Shell Model Approach

- Expand Using evolution operator in Chebyshev polynomials

$$\exp(-iHt) = \sum_{n=0}^{\infty} (-i)^n (2 - \delta_{n0}) J_n(t) T_n(H)$$

- Chebyshev polynomial $T_n[\cos(\theta)] = \cos(n\theta)$
- Use iterative relation and matrix-vector multiplication to generate

$$|\lambda_n\rangle = T_n(H)|\lambda\rangle$$

$$|\lambda_0\rangle = |\lambda\rangle, \quad |\lambda_1\rangle = H|\lambda\rangle \quad |\lambda_{n+1}\rangle = 2H|\lambda_n\rangle - |\lambda_{n-1}\rangle$$

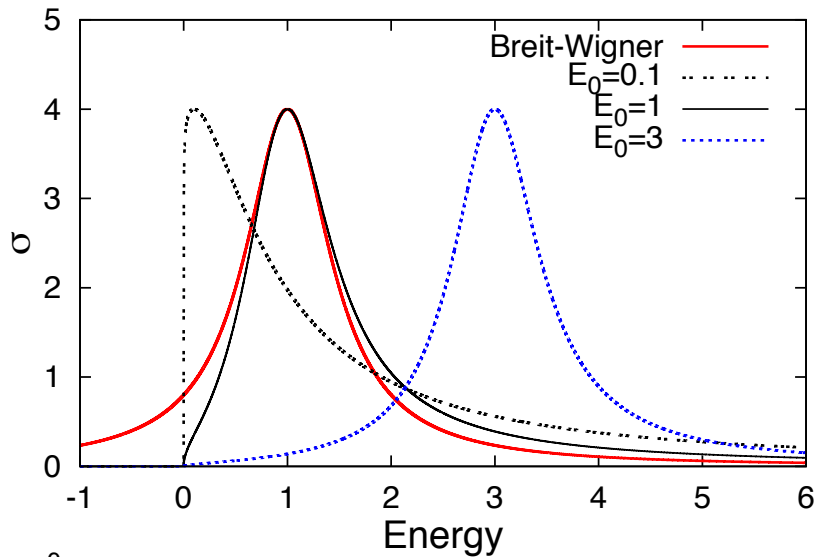
$$\langle\lambda'|T_{n+m}(H)|\lambda\rangle = 2\langle\lambda'_m|\lambda_n\rangle - \langle\lambda'|\lambda_{n-m}\rangle, \quad n \geq m$$

- Use FFT to find return to energy representation

Probing the Non-exponential Decay Regime in Open Quantum Systems

- Broad threshold resonance (^9N , ^9He)
 - Pronounced non-exponentiality
 - Very short half-life
- Three-body decay (^6Be , ^{13}Li , ^{16}Be)
 - Nucleon-nucleon correlations
 - Energy dependence
- Overlapping resonances (^{13}C , ^{13}N)
 - Interference, pronounced non-exponentiality
 - Superradiance

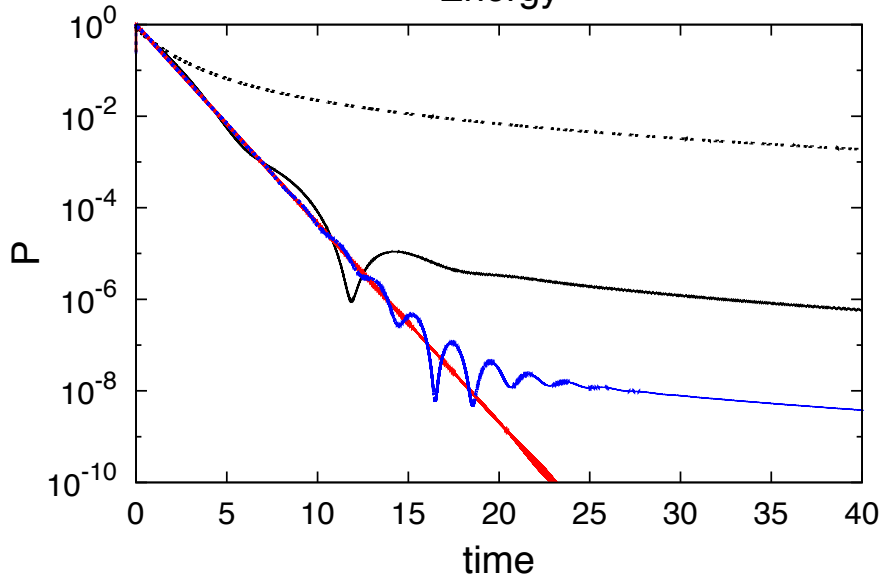
Time-dependent picture



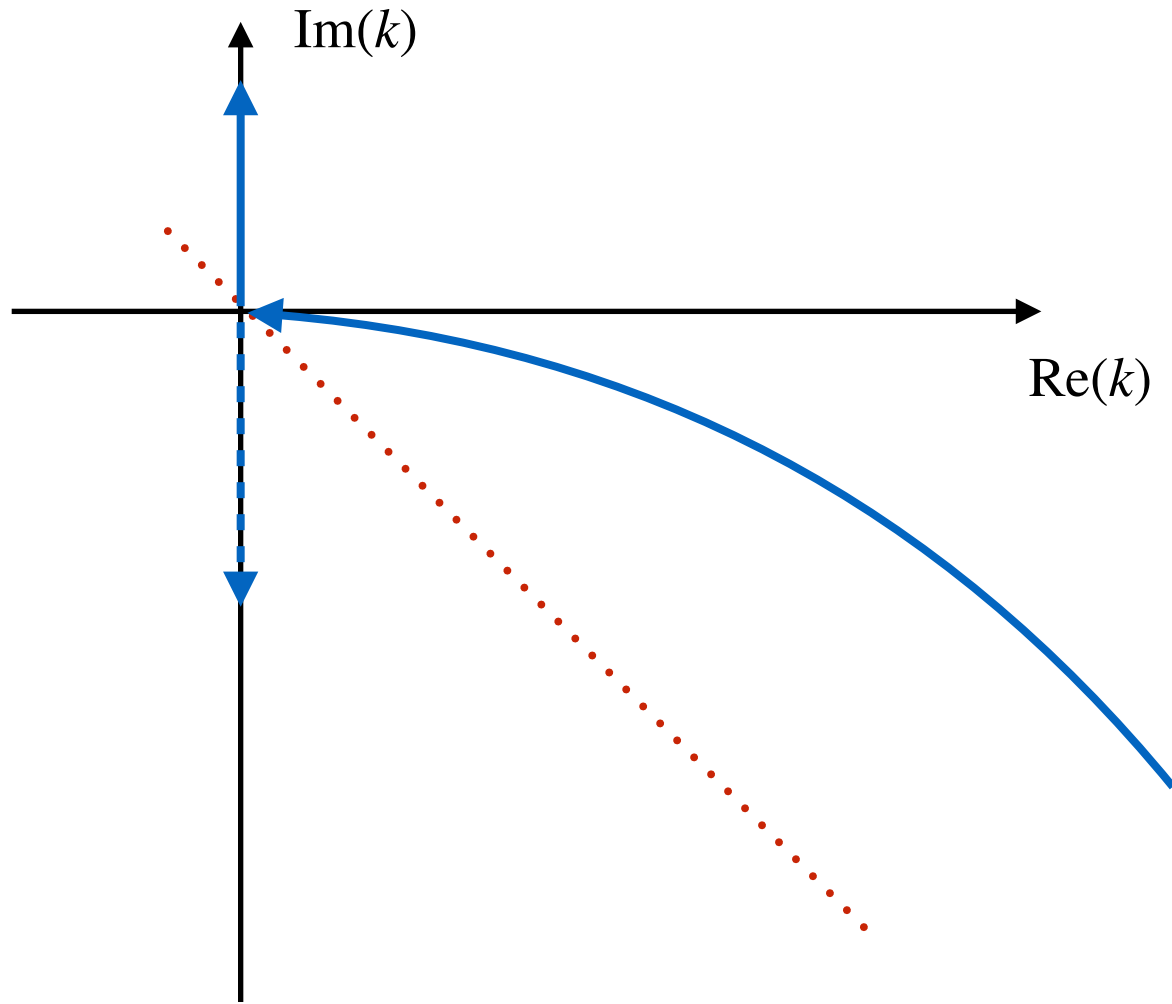
$$\mathcal{G} = \frac{1}{E - E_0 + i/2\Gamma(E)}$$

$$\Gamma(E) \propto \sqrt{E}$$

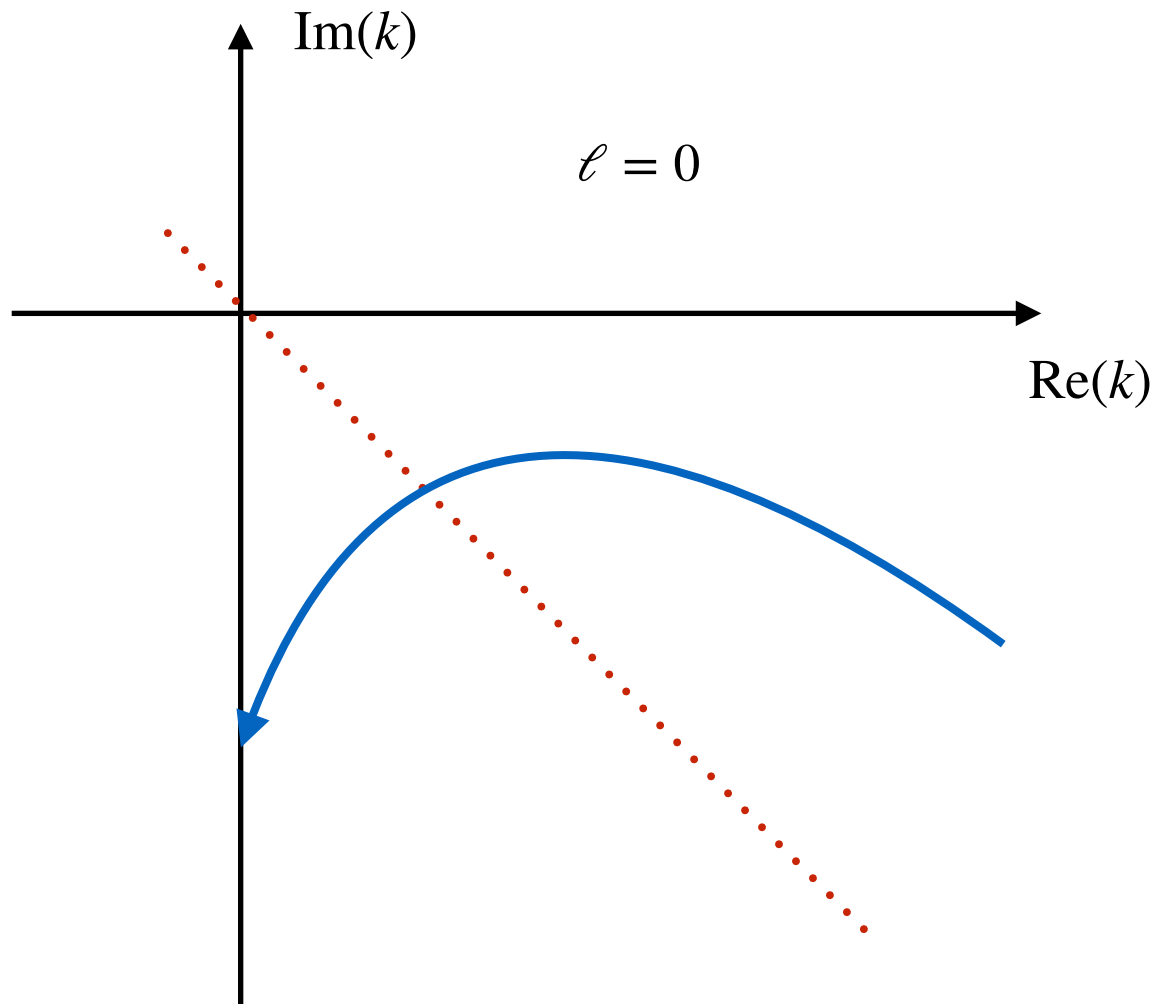
Power-law remote decay rate!



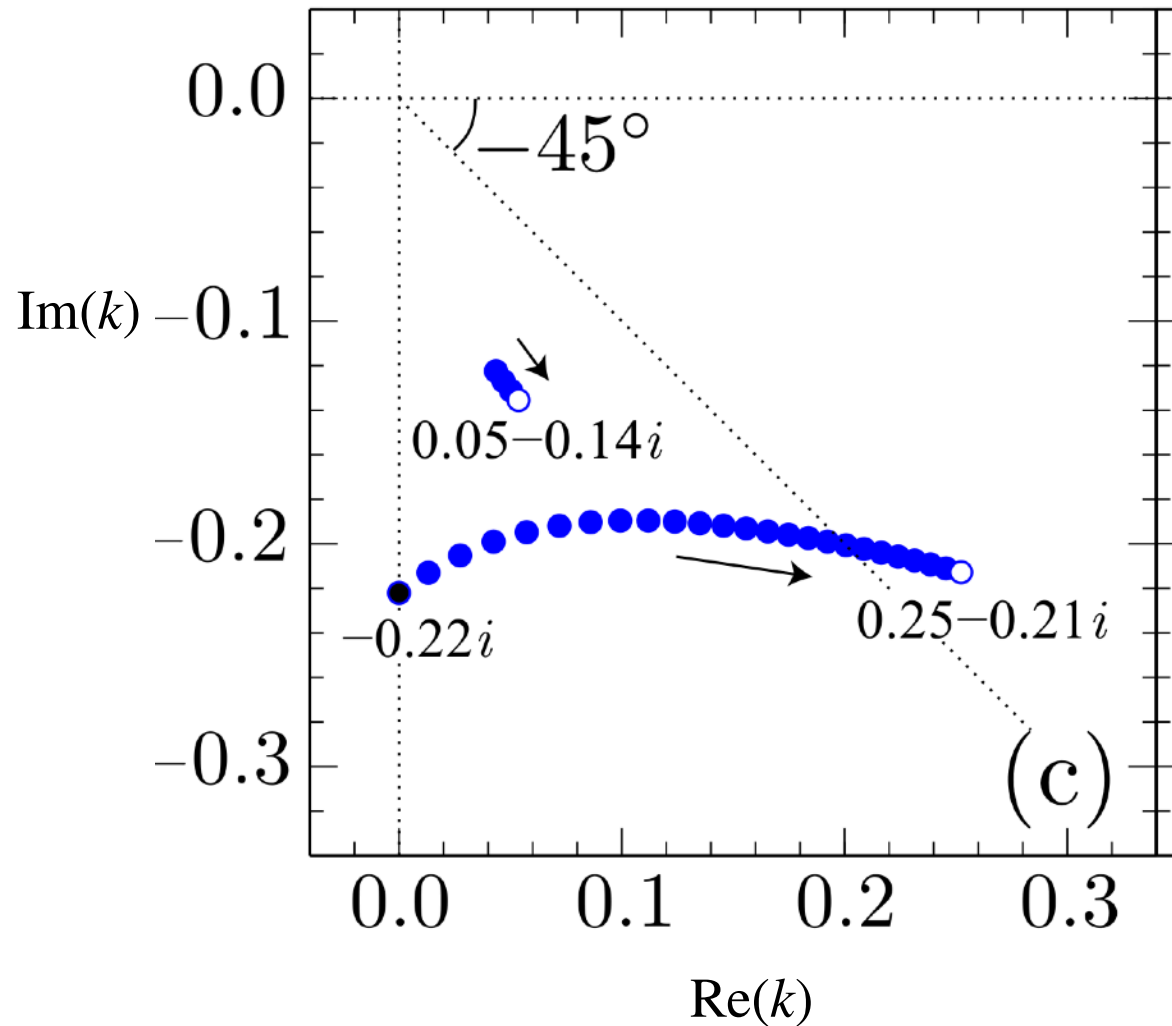
Resonance to bound state pole dynamics



Resonance to bound state pole dynamics



Resonance to bound state pole dynamics



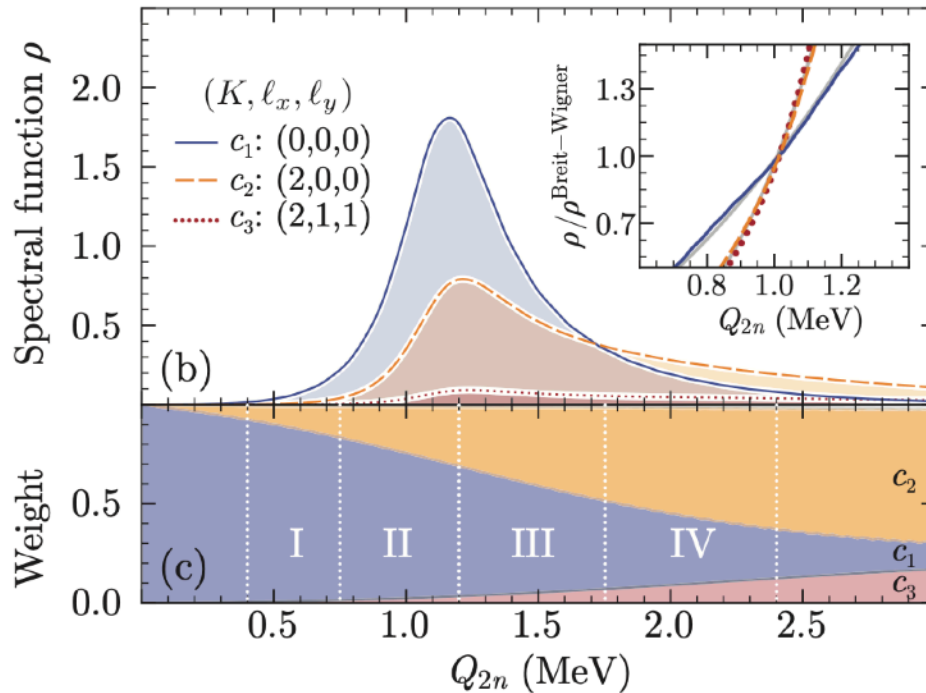
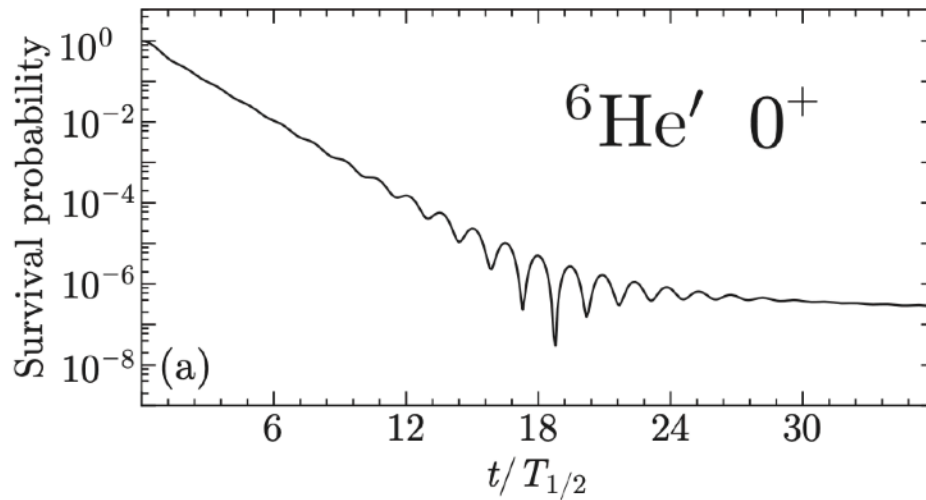
$n+{}^9\text{Li}$ $p+{}^9\text{C}$

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Time-dependent picture

Two-neutron decay

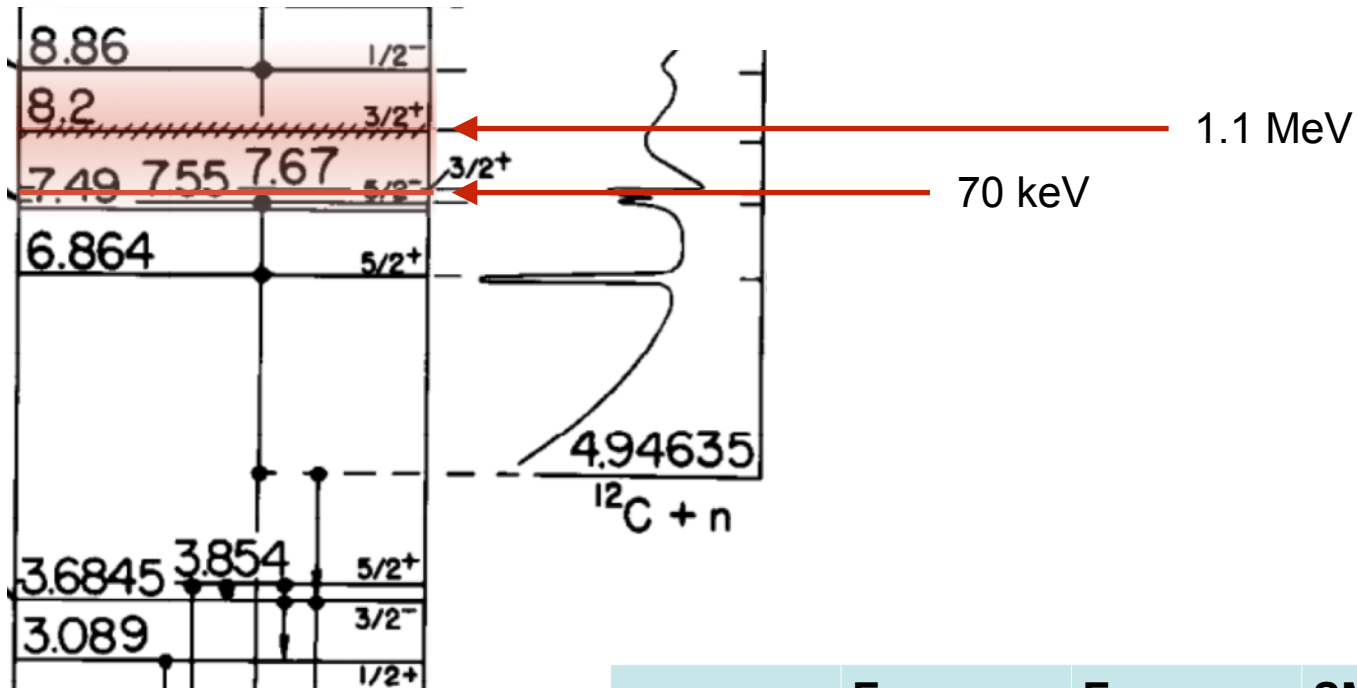


S. M. Wang, W. Nazarewicz, A. Volya, and Y. G. Ma,
“Probing the Non-exponential Decay Regime in Open
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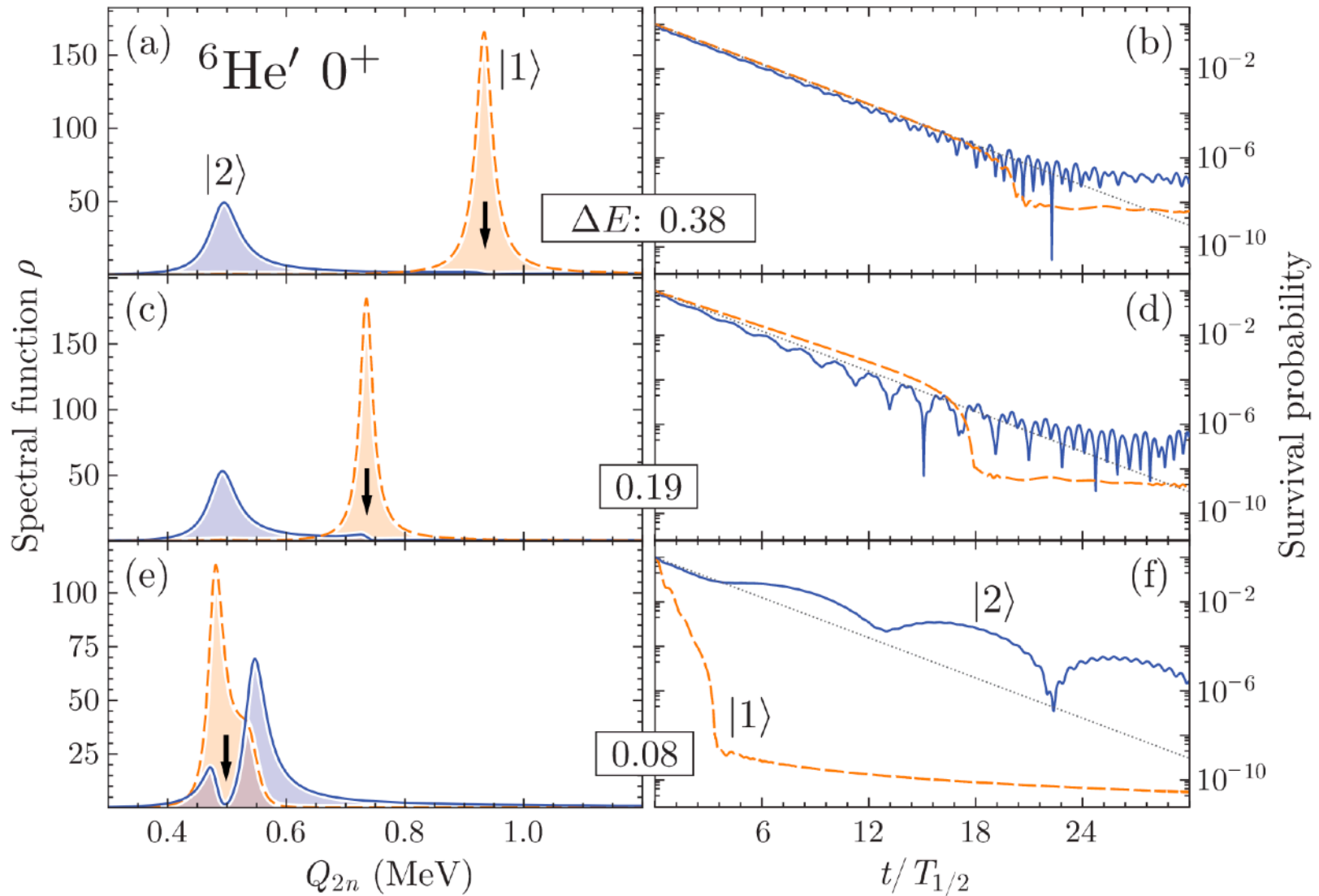
Superradiance in ^{13}C



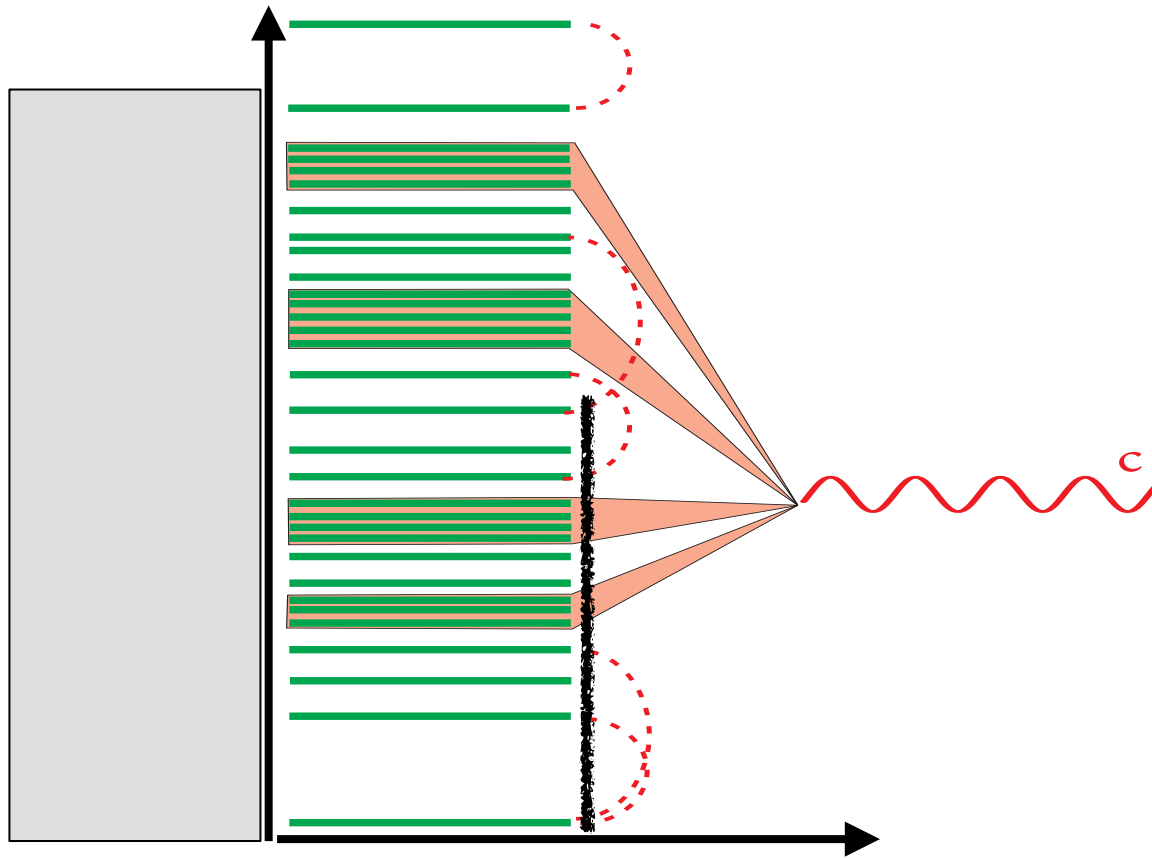
^{13}C

	Exp. Energy	Exp. Width	SM. Width	CSM Width
3/2+ (1)	7.686(6)	0.070(5)	0.858	0.098
3/2+ (2)	8.2(1)	1.1(3)	0.342	1.031

Example of interference between resonances



Coupling via continuum



Universal law of relaxation

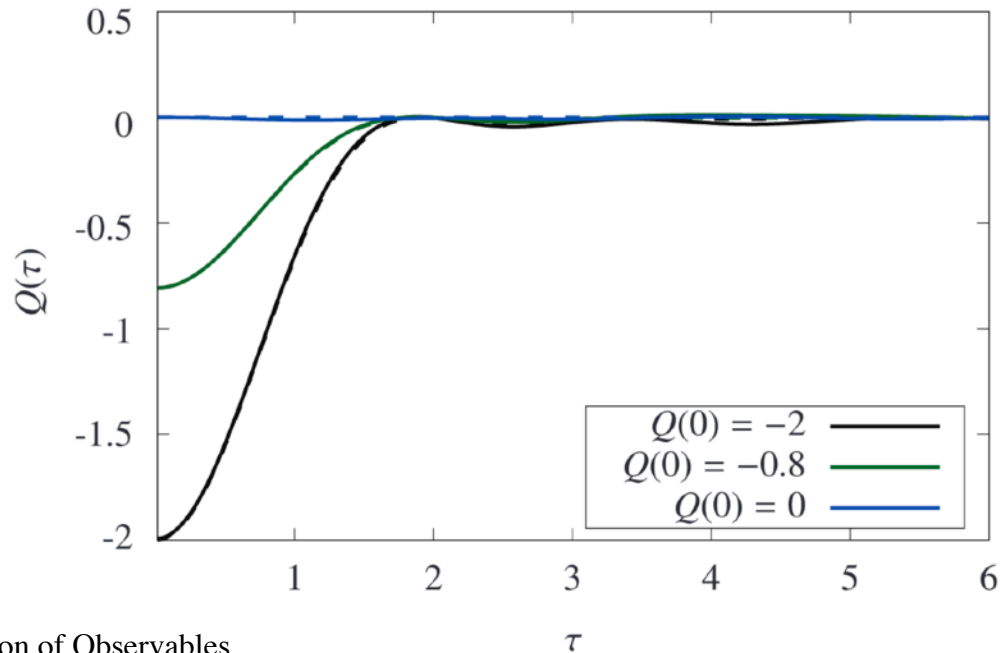
$$Q_a(t) = \langle \Psi_a(t) | \hat{Q} | \Psi_a(t) \rangle, \quad f_{a'a}(t) = \langle a' | e^{-i\hat{H}t} | a \rangle,$$

$$f_{aa'}(0) = \delta_{aa'}, \quad f_{aa'}(t) = f_{a'a}(-t),$$

$$\overline{f_{aa'}(t)} = \sum_n e^{-iE_n t} \overline{\langle a | n \rangle \langle n | a' \rangle} = \delta_{aa'} f(t), \quad Q(t) = Q(0) |f(t)|^2$$

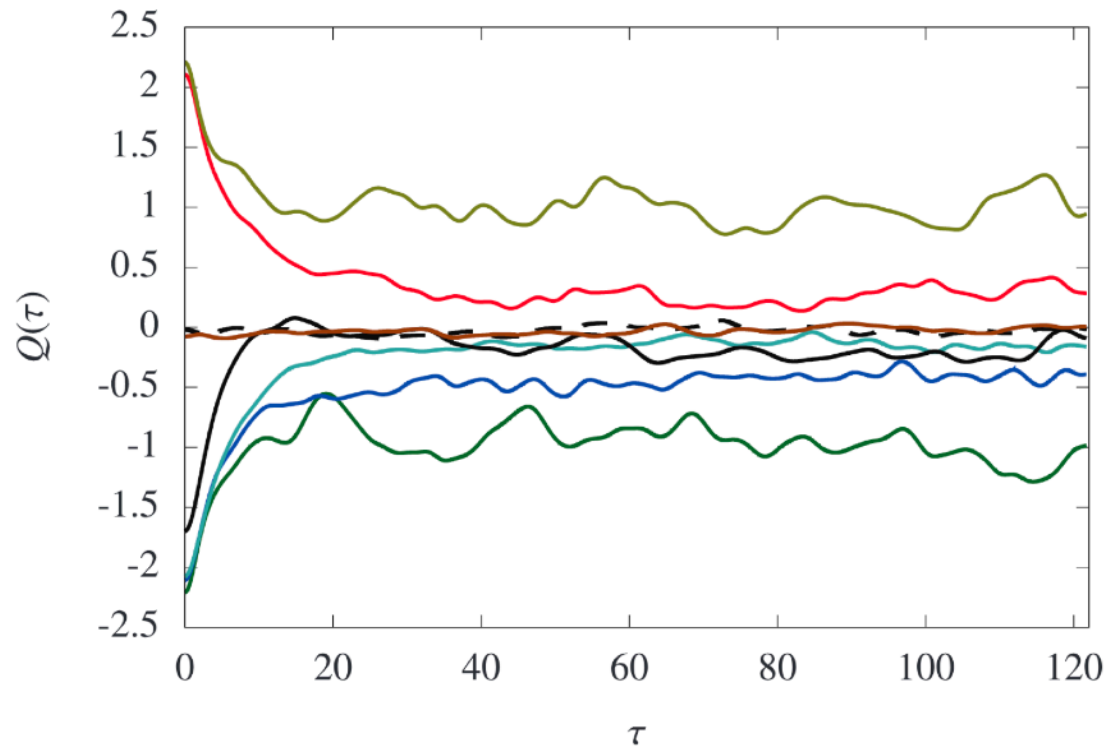
Under the above assumptions of full randomness, the function $f(t)$ is universal

$$f(\tau) = \frac{1}{\tau} J_1(2\tau),$$



Realistic shell model

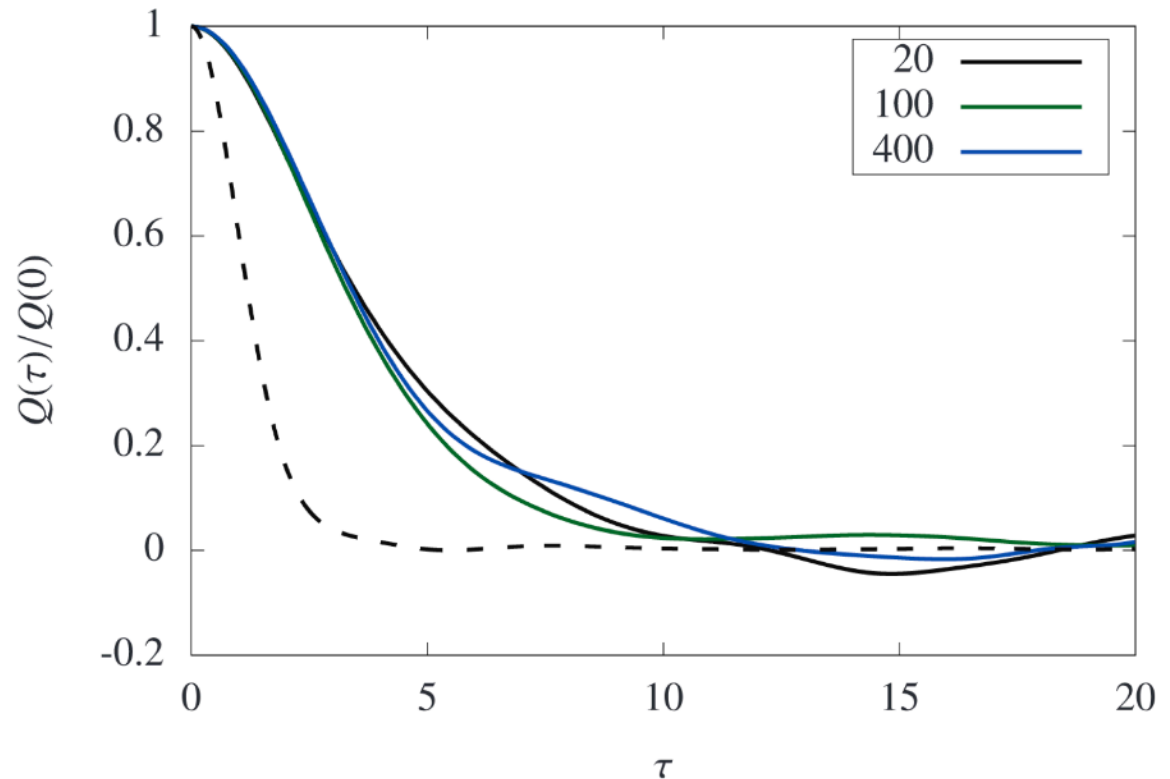
Realistic sd shell model
USDB Hamiltonian
Realistic quadrupole moment
Spreading width ~ 10 MeV



- Most states thermalize at zero
- Extreme states retain their average quadrupole moment
- Giant resonances
- Scars in chaotic dynamics

Relaxation in realistic systems

Realistic sd shell model
USDB Hamiltonian
Realistic quadrupole moment
Spreading width ~ 10 MeV



Relaxation times in realistic systems
is significantly longer.
Extended lifetime of perturbation

A. Volya and V. Zelevinsky, Time-Dependent Relaxation of Observables
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K. Ma, A. Volya, and K. Yang, *Eigenstate Thermalization and
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Conclusions

- Discussion of dynamics becomes central to exotic nuclear systems on the verge of stability.
- Exponential decay is not exact and is a result of delicate approximations.
- Poles of the scattering matrix differentiate the types of dynamics.
- The initial state (of real energy) matters, as do the reaction mechanism and memory effect.
- Experimental studies of correlations can be used to probe the state and its time evolution.
- There is strong competition between internal relaxation and decay.

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