CHIRAL THERMODYNAMICS: from NUCLEAR MATTER to NEUTRON STARS

Wolfram Weise
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- QCD interface with nuclear physics:
 in-medium Chiral Effective Field Theory
- Nuclear Equation of State:

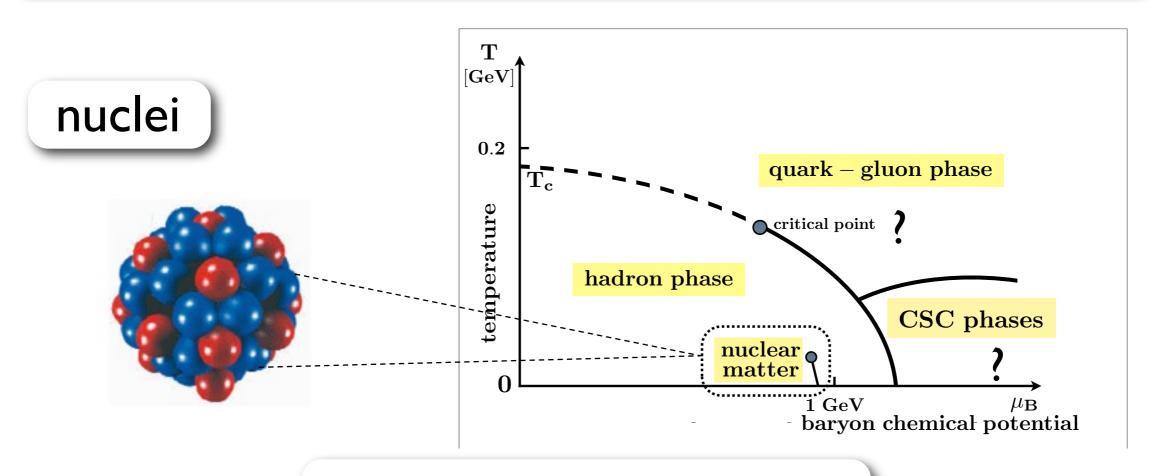
 liquid-gas phase transition and its N/Z evolution
- Density and temperature dependence of the Chiral (Quark) Condensate
- **Equation of State** of dense and cold matter:

 new constraints from **Neutron Stars**





NUCLEAR MATTER and QCD PHASES



Scales in nuclear matter:

Momentum scale:

Fermi momentum

- NN distance:
- Energy per nucleon:
- Compression modulus:

$$m k_F \simeq 1.4~fm^{-1} \sim 2m_\pi$$

$$m d_{NN} \simeq 1.8 ~fm \simeq 1.3 ~m_\pi^{-1}$$

$$\mathbf{E}/\mathbf{A} \simeq -\mathbf{16} \ \mathbf{MeV}$$

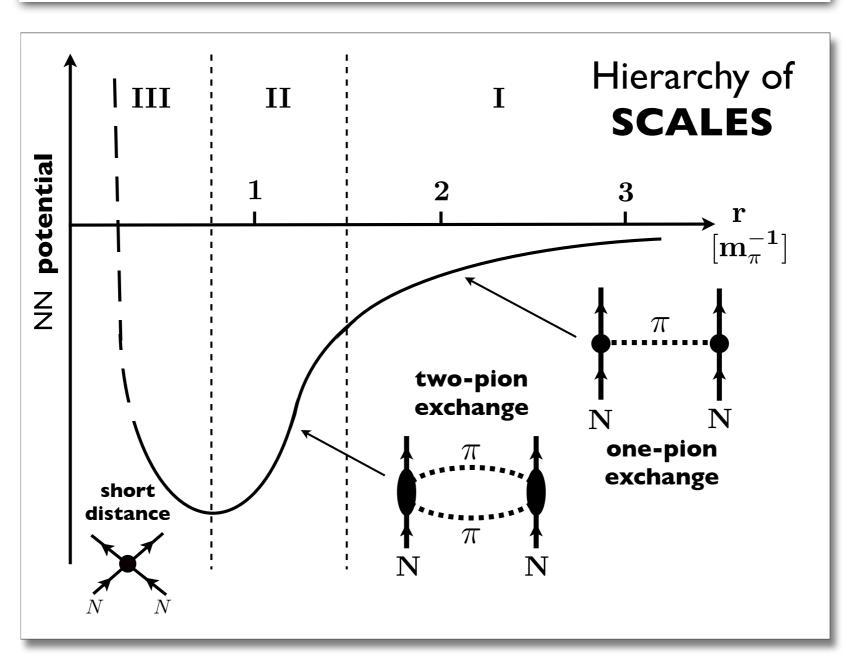
$$\mathbf{K} = (\mathbf{260} \pm \mathbf{30}) \; \mathbf{MeV} {\sim} \; \mathbf{2m}_{\pi}$$





Nuclear Forces

- Contemporary Developments -



Early history:

M. Taketani, S. Nakamura, M. Sasaki (1951)

Today's approach:

Chiral Effective Field Theory + Lattice QCD





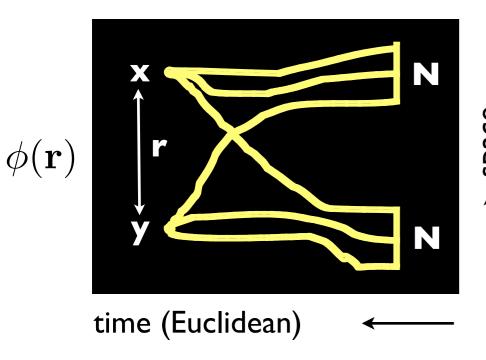
Short distance:

NN POTENTIAL from LATTICE QCD

N. Ishii, S. Aoki, T. Hatsuda: Phys. Rev. Lett. 99 (2007) 022001

K. Murano, N. Ishii, S. Aoki, T. Hatsuda: Prog. Theor. Phys. 125 (2011) 1225

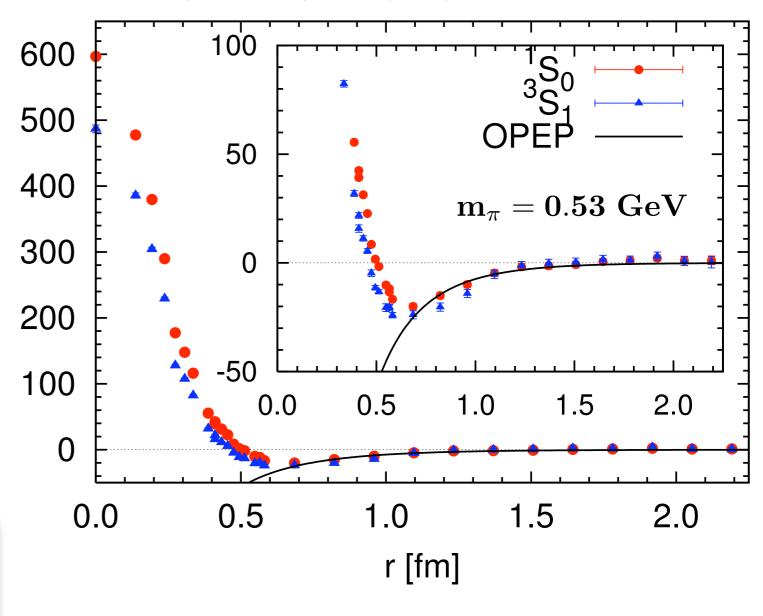
 $V_{\rm C}({\rm r})$ [MeV]



still: mostly quenched QCD "large" quark/pion masses

Reconstruct potential from wave function:

$$\mathbf{V_C}(\mathbf{r}) = \mathbf{E} + \frac{\nabla^2 \phi(\mathbf{r})}{2\mu \ \phi(\mathbf{r})}$$



 Repulsive core from Lattice QCD





Chiral Effective Field Theory and the Nuclear Many-Body Problem

PIONS and NUCLEI in the context of LOW-ENERGY QCD

- CONFINEMENT of quarks and gluons in hadrons
- Spontaneously broken CHIRAL SYMMETRY

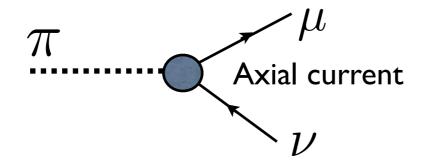
LOW-ENERGY / LOW-TEMPERATURE QCD:

Effective Field Theory of weakly interacting

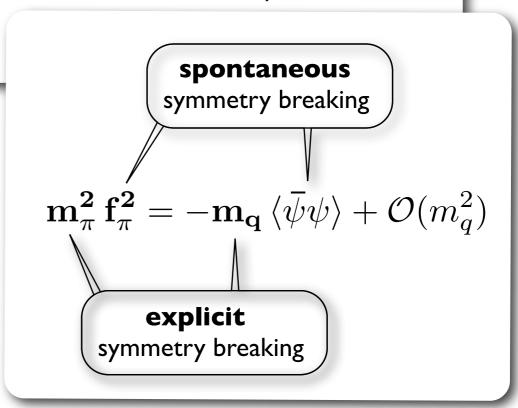
Nambu-Goldstone Bosons (PIONS)

representing QCD at (energy and momentum) scales

$$\mathbf{Q} << \mathbf{4}\pi\,\mathbf{f}_\pi \sim \,\mathbf{1}\,\mathbf{GeV}$$



$$\mathbf{f}_{\pi} = \mathbf{92.4\,MeV}$$





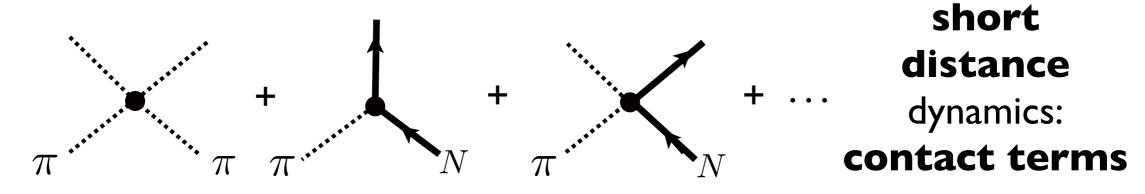


CHIRAL EFFECTIVE FIELD THEORY

- Systematic framework at interface of QCD and Nuclear Physics
- Interacting systems of PIONS (light / fast) and NUCLEONS (heavy / slow):

$$\mathcal{L}_{eff} = \mathcal{L}_{\pi}(U, \partial U) + \mathcal{L}_{N}(\Psi_{N}, U, ...)$$
$$U(x) = exp[i\tau_{a}\pi_{a}(x)/f_{\pi}]$$

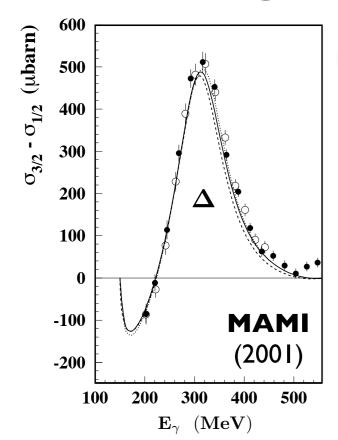
Construction of Effective Lagrangian: Symmetries





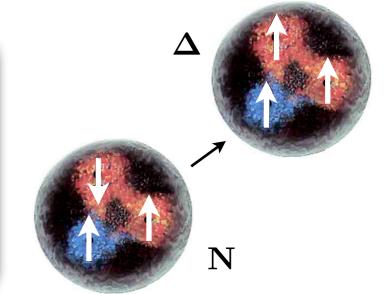
Explicit $\Delta(1230)$ DEGREES of FREEDOM

Large spin-isospin polarizabilty of the Nucleon



example: polarized Compton scattering

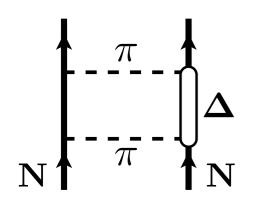
$$eta_\Delta = rac{\mathbf{g_A^2}}{\mathbf{f_\pi^2}(\mathbf{M_\Delta} - \mathbf{M_N})} \sim 5\,\mathrm{fm^3}$$
 $\mathbf{M_\Delta} - \mathbf{M_N} \simeq 2\,\,\mathbf{m_\pi} << 4\pi\,\mathbf{f_\pi}$ (small scale)



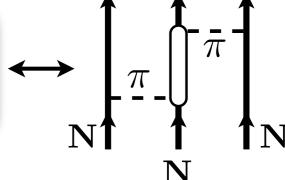
Pionic Van der Waals - type intermediate range central potential

N. Kaiser, S. Gerstendörfer, W.W., NPA637 (1998) 395

N. Kaiser, S. Fritsch, W.W., NPA750 (2005) 259



$$egin{aligned} \mathbf{\Delta} & \mathbf{V_c}(\mathbf{r}) = -rac{9\,\mathbf{g_A^2}}{32\pi^2\,\mathbf{f_\pi^2}}\,eta_\Delta\,rac{\mathbf{e^{-2m_\pi r}}}{\mathbf{r^6}}\,\mathbf{P}(\mathbf{m_\pi r}) \end{aligned}$$



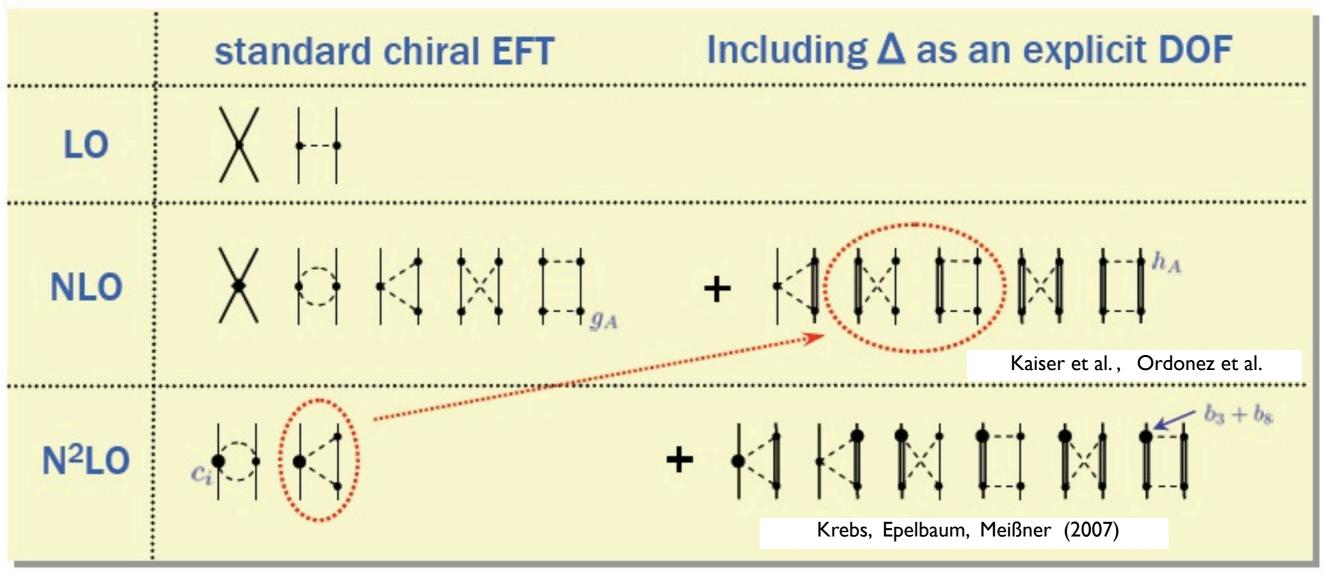
strong 3-body interaction

J. Fujita, H. Miyazawa (1957) Pieper, Pandharipande, Wiringa, Carlson (2001)





Explicit $\Delta(1230)$ DEGREES of FREEDOM (contd.)



(with thanks to Evgeny Epelbaum)

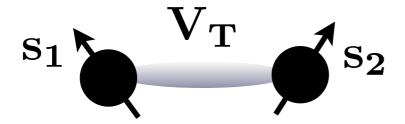
- ullet Important physics of $oldsymbol{\Delta}(1230)$ promoted to NLO
 - Improved convergence



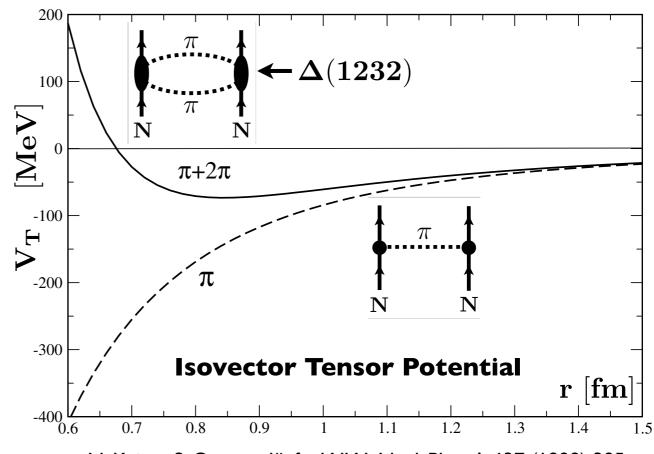


Important pieces of the CHIRAL NUCLEON-NUCLEON INTERACTION

ISOVECTORTENSOR FORCE

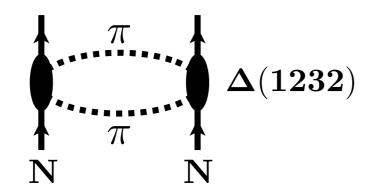


 \bullet note: **no** ρ meson



N. Kaiser, S. Gerstendörfer, W.W.: Nucl. Phys. A 637 (1998) 395

CENTRAL ATTRACTION from TWO-PION EXCHANGE



lacktriangle note: **no** σ boson

Van der WAALS - like force:

$$\mathbf{V_c}(\mathbf{r}) \propto -rac{\exp[-2\mathbf{m_\pi r}]}{\mathbf{r^6}}\mathbf{P}(\mathbf{m_\pi r})$$

... at intermediate and long distance





CHIRAL DYNAMICS and the **NUCLEAR MANY-BODY PROBLEM**

N. Kaiser, S. Fritsch, W.W. (2002 - 2005)

- **Small** scales:
- ${
 m k_F} \sim 2\,{
 m m_\pi} \sim {
 m M_\Delta} {
 m M_N} << 4\pi\,{
 m f_\pi}$
- PIONS (and DELTA isobars) as explicit degrees of freedom
- IN-MEDIUM CHIRAL PERTURBATION THEORY

pion exchange processes in presence of filled Fermi sea

2nd order TENSOR force + nucleon's SPIN-ISOSPIN polarizability

short-distance dynamics: N contact interactions (incl. resummations)



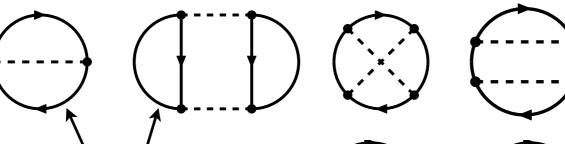


IN-MEDIUM CHIRAL PERTURBATION THEORY

Loop expansion of (In-Medium) Chiral Perturbation Theory
↑

Systematic expansion of **ENERGY DENSITY** $\mathcal{E}(\mathbf{k_F})$ in **powers** of **Fermi momentum** [modulo functions $\mathbf{f_n}(\mathbf{k_F}/\mathbf{m_\pi})$] (works for $\mathbf{k_F} << 4\pi\,\mathbf{f_\pi} \sim 1\,\mathrm{GeV}$)

Nuclear thermodynamics: compute free energy density



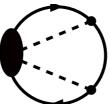
(3-loop order)

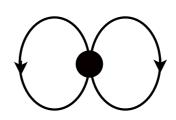
N. Kaiser, S. Fritsch, W.W. (2002-2005)

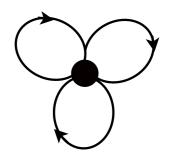
in-medium

nucleon propagators incl. Pauli blocking











FINITE NUCLEI

... including the nuclear surface :
Energy Density Functional

J.W. Holt, N. Kaiser, W.W.: Eur. Phys. J. A 47 (2011) 128

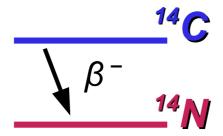
- many quantitatively successful applications throughout the nuclear chart e.g. P. Finelli et al.: Nucl. Phys. A 770 (2007) I
 - binding energies and charge radii
 - pairing and ground state deformations
- systematics through isotopic chains governed by isospin dependent forces from chiral pion dynamics



... further applications

Gamow-Teller beta decays

interesting case:



anomalously long lifetime (5739 y) enables radiocarbon dating

Theoretically **not** understood on the basis of **two-nucleon** interactions only

- Solution: chiral effective interaction including three-body force J.W. Holt, N. Kaiser, W.W.: Phys. Rev. C79 (2009) 054331, Phys. Rev. C81 (2010) 024002
- Spin-orbit interactions
 - Role of 2nd order tensor force from pion exchange and three-body interactions

N. Kaiser: Phys. Rev. C68 (2003) 054001; N. Kaiser and W.W.: Nucl. Phys. A804 (2008) 60

- In-medium Chiral SU(3) dynamics and hypernuclei
 - \triangleright Weak Λ -nuclear spin-orbit coupling

N. Kaiser, W.W.: Phys. Rev. C71 (2005) 015203

P. Finelli, N. Kaiser, D. Vretenar, W.W.: Phys. Lett. B658 (2007) 90; Nucl. Phys. A831 (2009) 163



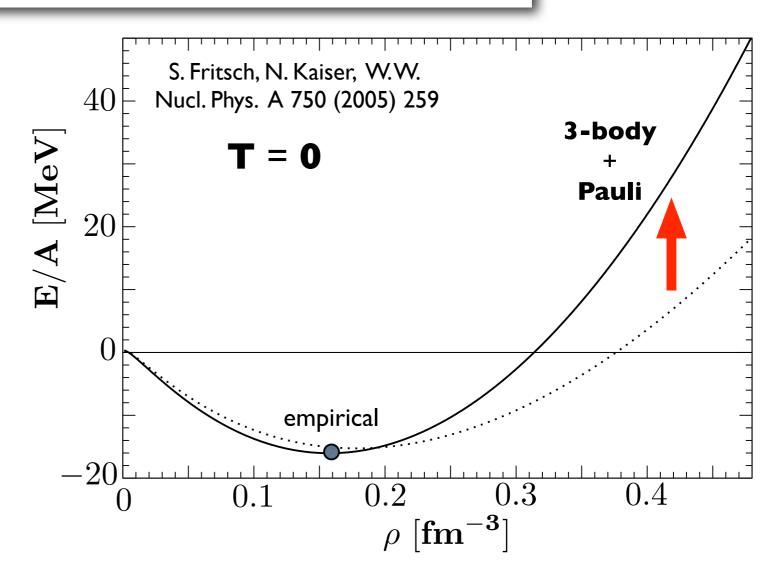
NUCLEAR MATTER

- In-medium ChPT 3-loop $(\pi, \mathbf{N}, \boldsymbol{\Delta})$
- **Input** parameters: two contact terms
- basically: analytic calculation
- **Output**:
 - Binding & saturation

$$E_0/A = -16 \, \mathrm{MeV} \; , \;
ho_0 = 0.16 \, \mathrm{fm}^{-3} \; , \; K = 290 \, \mathrm{MeV}$$



- Asymmetry energy $A(k_F^0) = 34 \, MeV$
- Quasiparticle interaction and Landau parameters



J.W. Holt, N. Kaiser, W.W. Nucl. Phys. A 870 (2011) 1, Nucl. Phys. A 876 (2012) 61



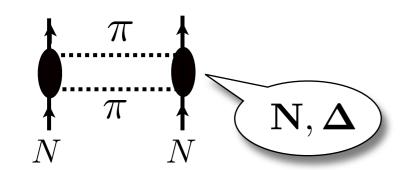


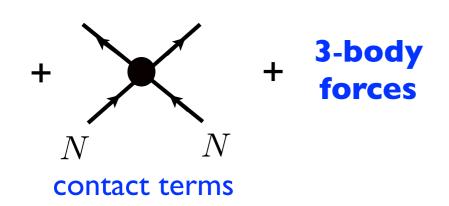
NUCLEAR THERMODYNAMICS

NUCLEAR CHIRAL (PION) DYNAMICS

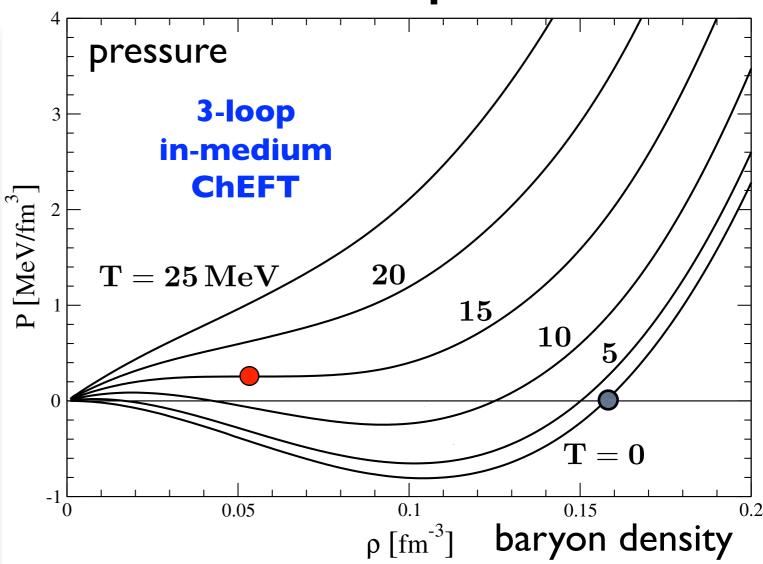
BINDING & SATURATION:

Van der Waals + Pauli





nuclear matter: equation of state



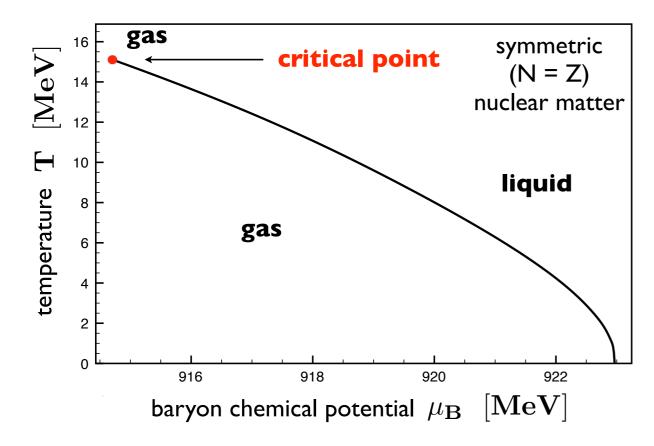
Liquid - Gas Transition at Critical Temperature $T_c = 15 \text{ MeV}$

(empirical: $T_c = 16 - 18 \text{ MeV}$)





PHASE DIAGRAM of NUCLEAR MATTER

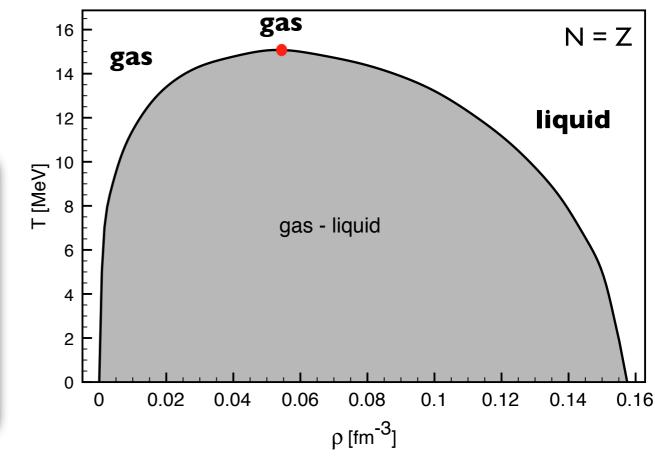


- Pion-nucleon dynamics incl. delta isobars
- Short-distanceNN contact terms
- Three-body forces

In-medium
 chiral effective field theory
 (3-loop calculation of free energy density)

S. Fritsch, N. Kaiser, W.W. Nucl. Phys. A 750 (2005) 259

S. Fiorilla, N. Kaiser, W.W. Nucl. Phys. A 880 (2012) 65

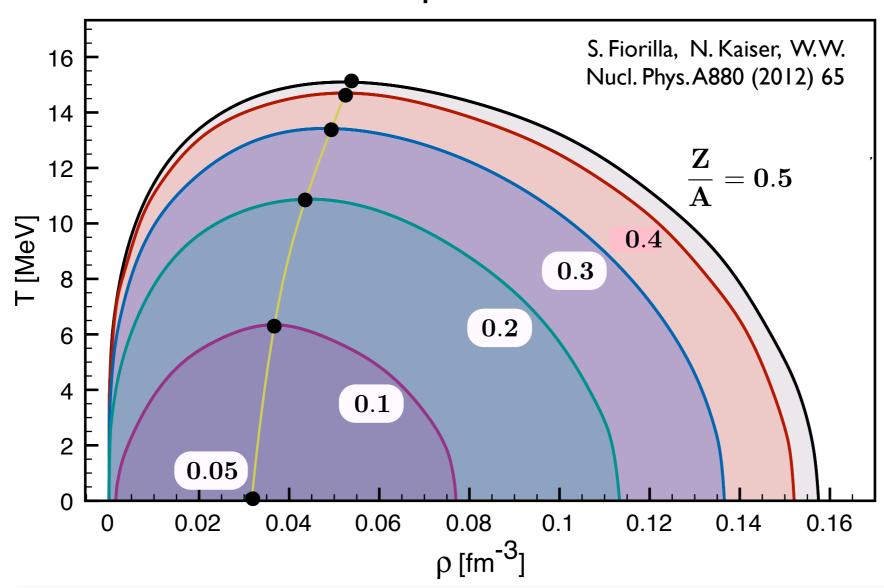






PHASE DIAGRAM of NUCLEAR MATTER

Trajectory of **CRITICAL POINT** for **asymmetric matter** as function of proton fraction Z/A



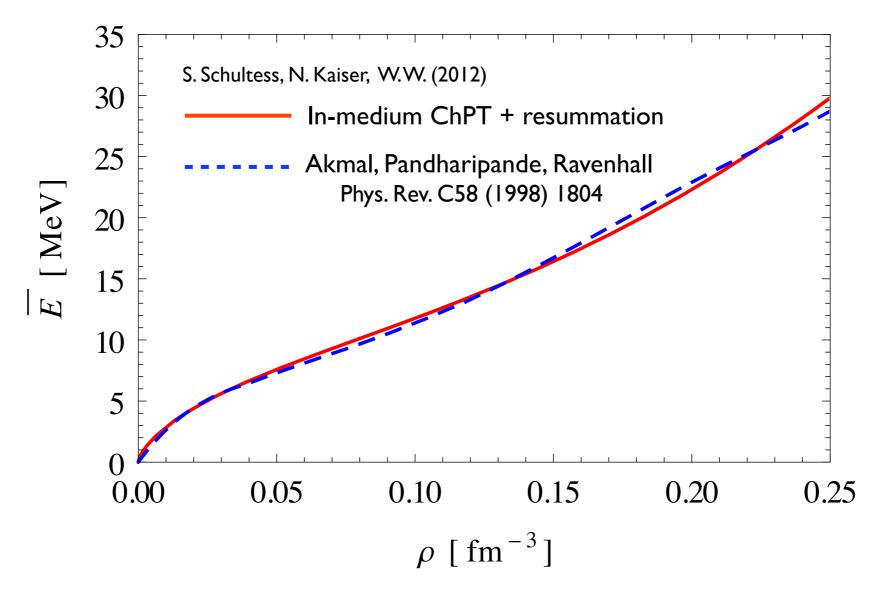
... determined almost entirely by isospin dependent (one- and two-) pion exchange dynamics



NEUTRON MATTER

In-medium chiral effective field theory (3-loop) with resummation of short distance contact terms (large nn scattering length, $a_{\rm s}=19~{\rm fm})$

N. Kaiser, Nucl. Phys. A 860 (2011) 370



perfect agreement with sophisticated many-body calculations
 (e.g. VCS (Urbana) or QMC methods (P. Armani et al., arXiv:1110.0993)



... short digression:

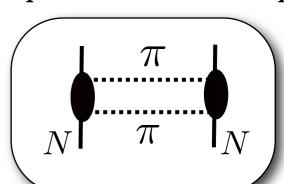
Nuclear Thermodynamics and the Chiral Condensate

CHIRAL CONDENSATE at finite BARYON DENSITY

- Chiral (quark) condensate $\langle \bar{q}q \rangle$: $m_\pi^2 \, f_\pi^2 = -2 \, m_q \langle \bar{q}q \rangle$ Order parameter of spontaneously broken chiral symmetry in QCD
- Hellmann Feynman theorem: $\langle \Psi | \bar{\mathbf{q}} \mathbf{q} | \Psi \rangle = \langle \Psi | \frac{\partial \mathcal{H}_{\mathbf{QCD}}}{\partial \mathbf{m_q}} | \Psi \rangle = \frac{\partial \mathcal{E}(\mathbf{m_q}; \rho)}{\partial \mathbf{m_q}}$

$$\mathbf{m_q} \frac{\partial \mathbf{M_N}}{\partial \mathbf{m_q}}$$

in-medium chiral effective field theory



$$rac{\langle ar{\mathbf{q}} \mathbf{q}
angle_{
ho}}{\langle ar{\mathbf{q}} \mathbf{q}
angle_{\mathbf{0}}} = \mathbf{1} - rac{
ho}{\mathbf{f_{\pi}^2}} \left[rac{ar{\sigma_{\mathbf{N}}}}{\mathbf{m_{\pi}^2}} \left(\mathbf{1} - rac{\mathbf{3} \, \mathbf{p_{F}^2}}{\mathbf{10} \, \mathbf{M_{\mathbf{N}}^2}} + \ldots
ight) + rac{\partial}{\partial \mathbf{m_{\pi}^2}} \left(rac{\mathbf{E_{int}}(\mathbf{p_F})}{\mathbf{A}}
ight)
ight]$$

(free) Fermi gas of nucleons

nuclear interactions (dependence on pion mass)





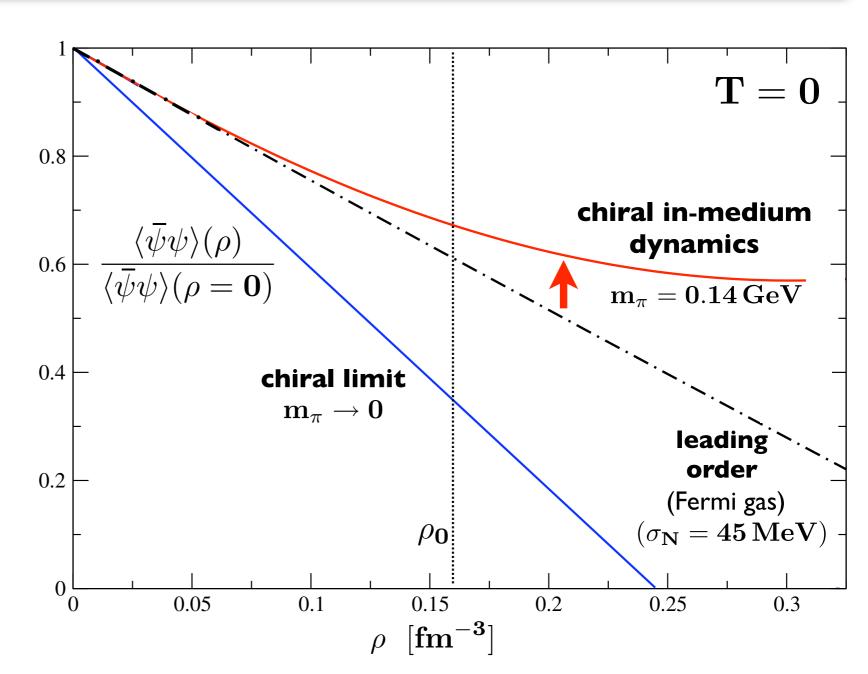
CHIRAL CONDENSATE: DENSITY DEPENDENCE

In-medium
Chiral
Effective
Field Theory

(NLO 3-loop)

constrained by realistic nuclear equation of state

N. Kaiser, Ph. de Homont, W.W. Phys. Rev. C 77 (2008) 025204



- Substantial change of symmetry breaking scenario between chiral limit $m_q=0$ and physical quark mass $m_q\sim 5\,MeV$
- Nuclear Physics would be very different in the chiral limit!





CHIRAL CONDENSATE:

DENSITY and TEMPERATURE DEPENDENCE

• Free energy density

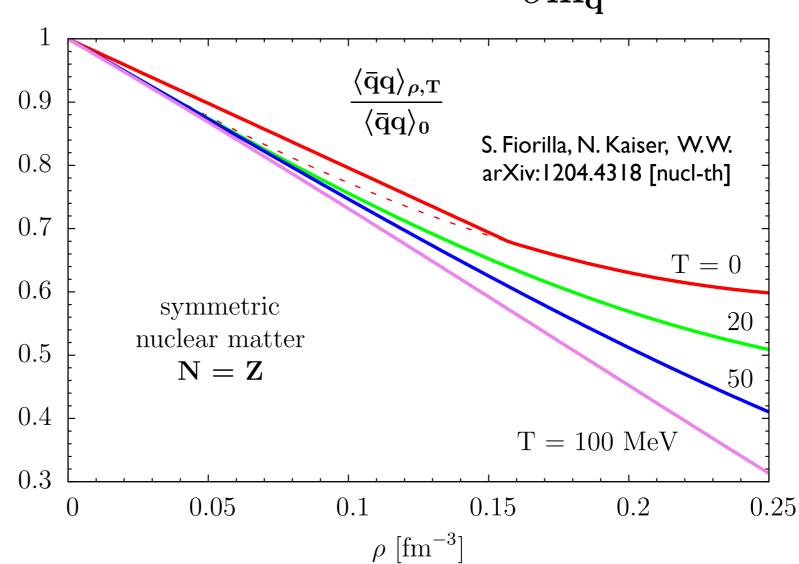
$$\mathcal{F}(\mathbf{m}_{\mathbf{q}}; \rho, \mathbf{T})$$

In-medium
Chiral
Effective
Field Theory

(NLO 3-loop)

constrained by realistic nuclear equation of state

$$\langle \mathbf{\Psi} | \mathbf{ar{q}} \mathbf{q} | \mathbf{\Psi}
angle_{
ho, \mathbf{T}} = rac{\partial \mathcal{F}(\mathbf{m_q};
ho, \mathbf{T})}{\partial \mathbf{m_q}}$$



No indication of first order chiral phase transition for

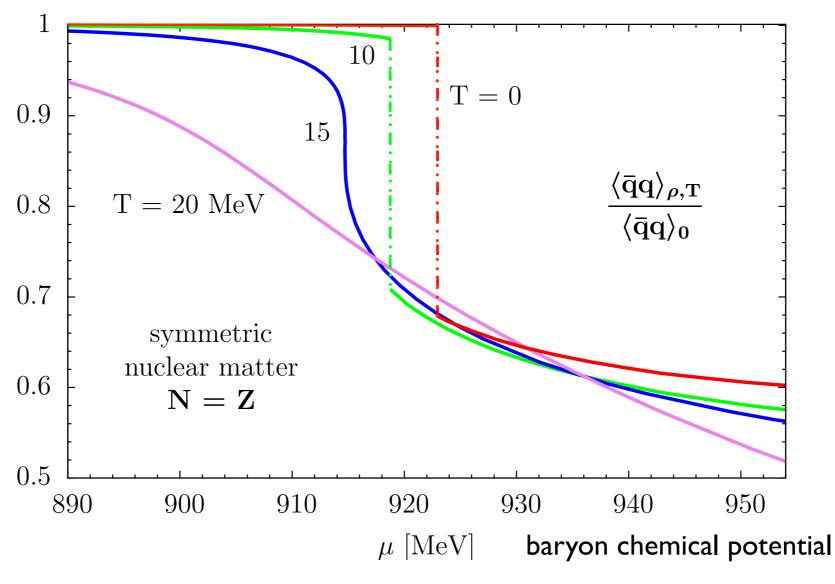
$$\rho \lesssim 2 \rho_0$$
, $T \lesssim 100 \,\mathrm{MeV}$





CHIRAL CONDENSATE:

Dependence on TEMPERATURE and BARYON CHEMICAL POTENTIAL



S. Fiorilla,
N. Kaiser,
W.W.
arXiv:1204.4318
[nucl-th]

- liquid-gas phase transition leaves its signature also in chiral condensate
- but: **no** tendency toward **chiral first order transition** in the range

$$\mu_{\mathbf{B}} \lesssim 1 \; \mathbf{GeV}$$

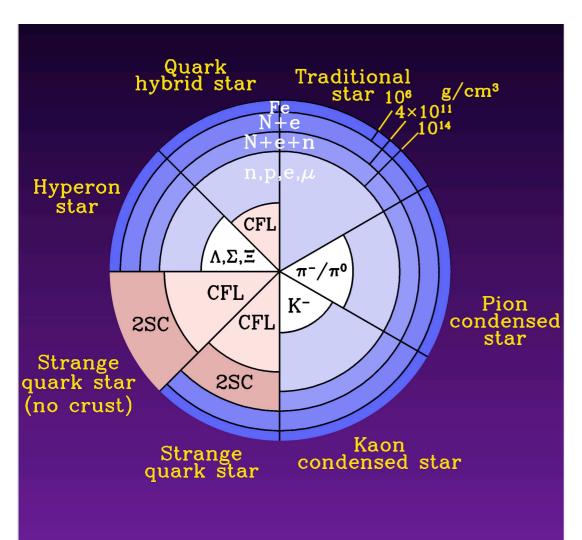


Outlooks:

New Constraints from NEUTRON STARS







Neutron Star Scenarios

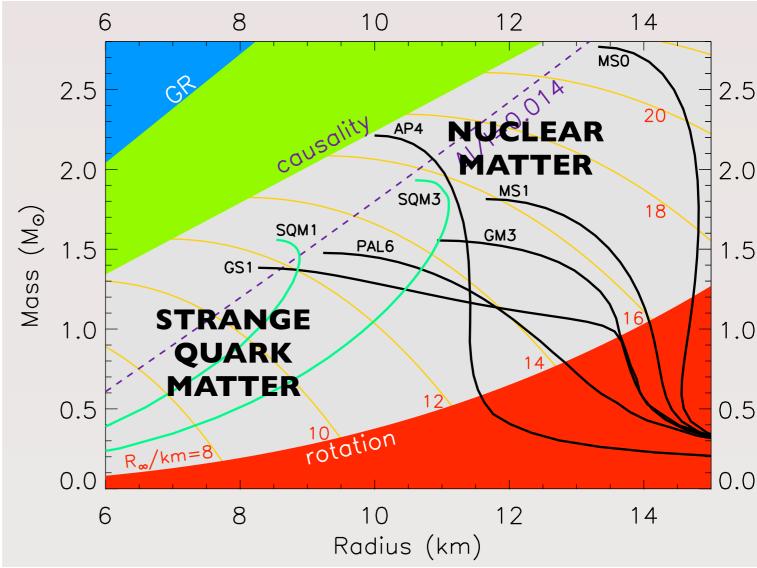
Tolman-Oppenheimer-Volkov equations

$$egin{split} rac{d\mathbf{P}}{d\mathbf{r}} &= -rac{\mathbf{G}}{\mathbf{c^2}} rac{(\mathbf{M} + 4\pi \mathbf{P}\mathbf{r^3})(\mathcal{E} + \mathbf{P})}{\mathbf{r}(\mathbf{r} - \mathbf{G}\mathbf{M}/\mathbf{c^2})} \ & rac{d\mathbf{M}}{d\mathbf{r}} = 4\pi \mathbf{r^2} rac{\mathcal{E}}{\mathbf{c^2}} \end{split}$$

NEUTRON STARS and the EQUATION OF STATE of DENSE BARYONIC MATTER

J. Lattimer, M. Prakash: Astrophys. J. 550 (2001) 426 Phys. Reports 442 (2007) 109

Mass-Radius Relation







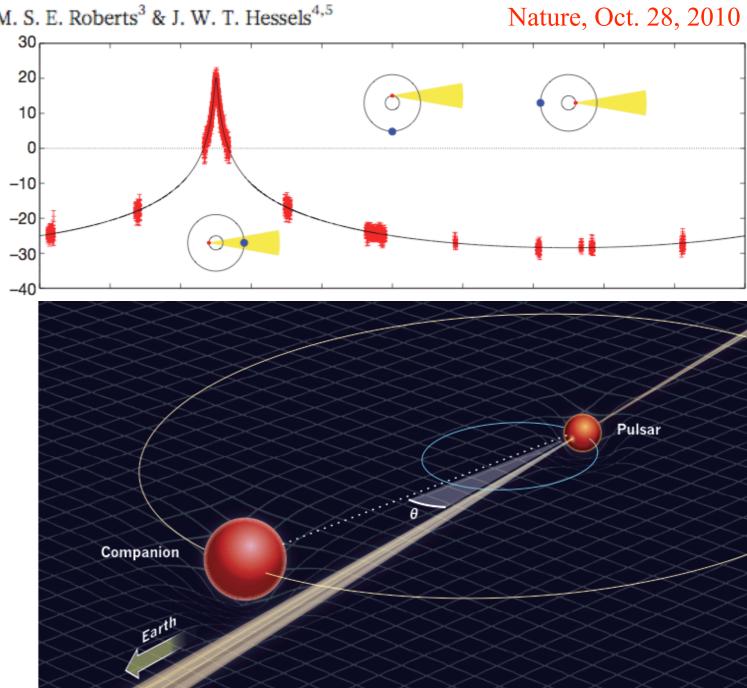
A two-solar-mass neutron star measured using Shapiro delay

P. B. Demorest¹, T. Pennucci², S. M. Ransom¹, M. S. E. Roberts³ & J. W. T. Hessels^{4,5}

direct measurement of neutron star mass from increase in travel time near companion

J1614-2230 most edge-on binary pulsar known (89.17°) + massive white dwarf companion (0.5 M_{sun})

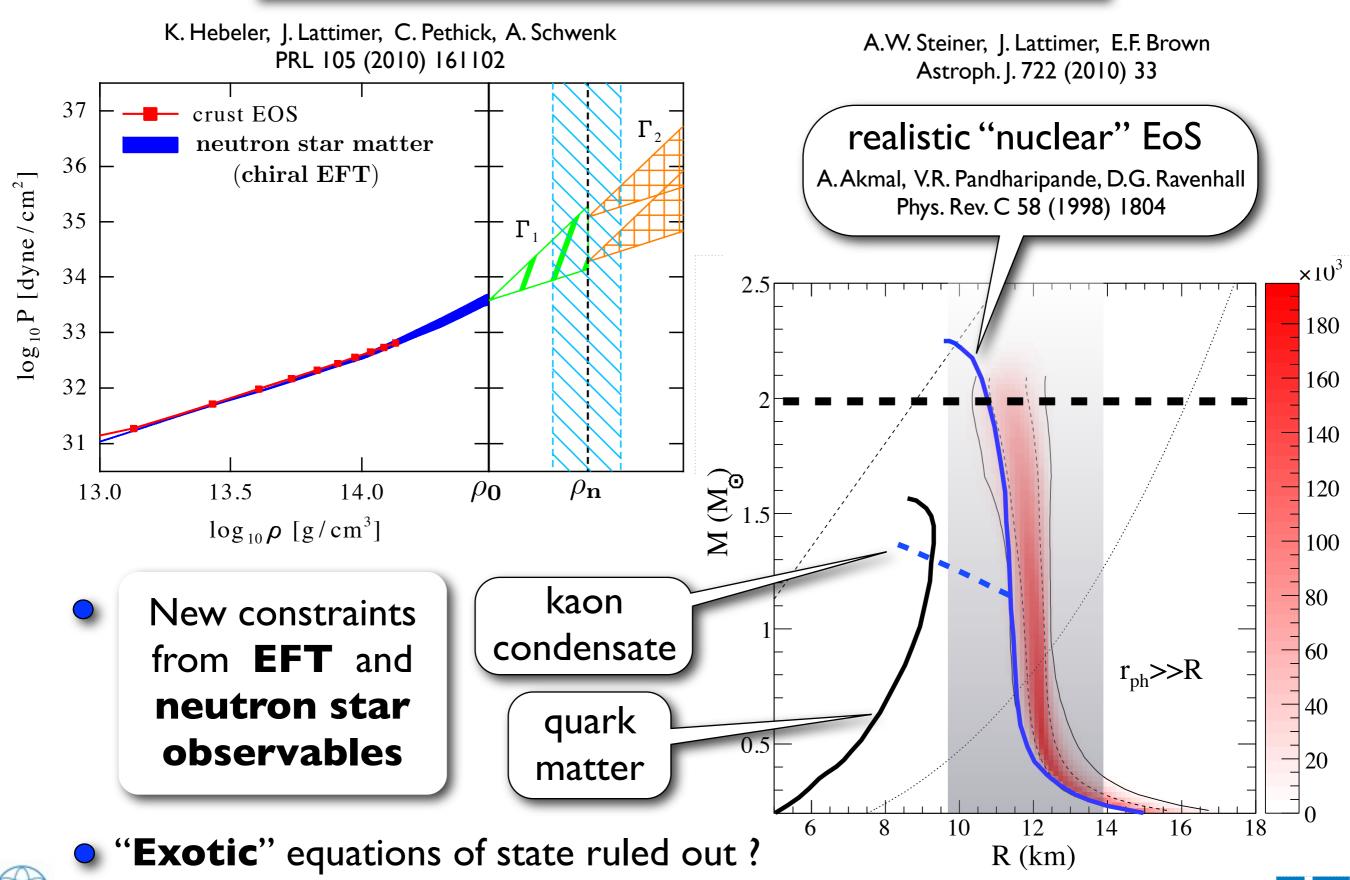
heaviest neutron star with 1.97±0.04 M_{sun}





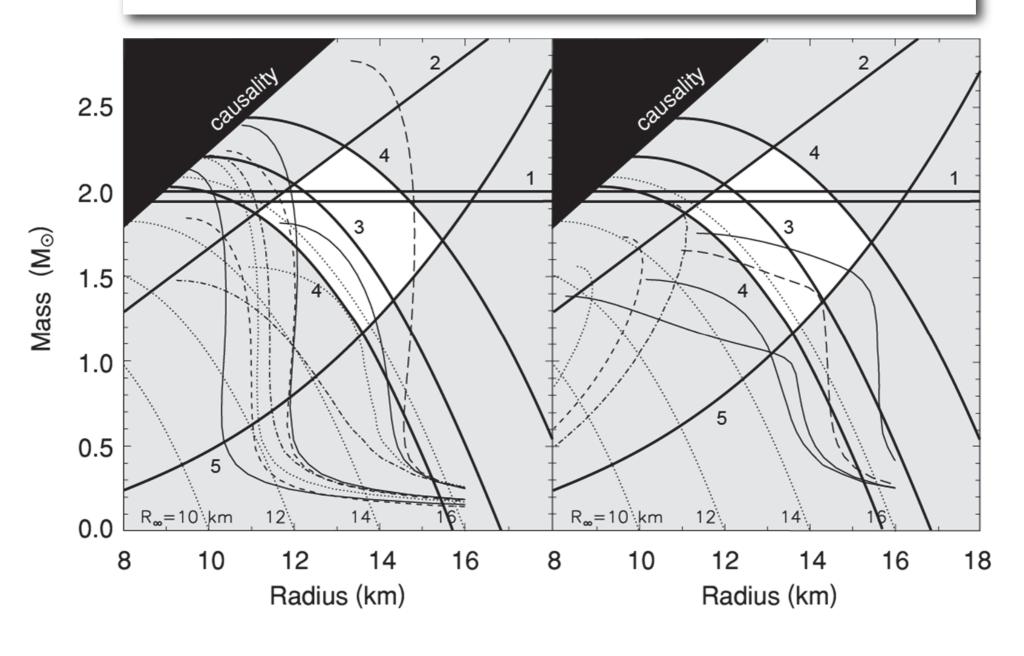


News from NEUTRON STARS





NEUTRON STARS: MASS and RADIUS constraints



from: J.Trümper

Irsee Symposium 2012

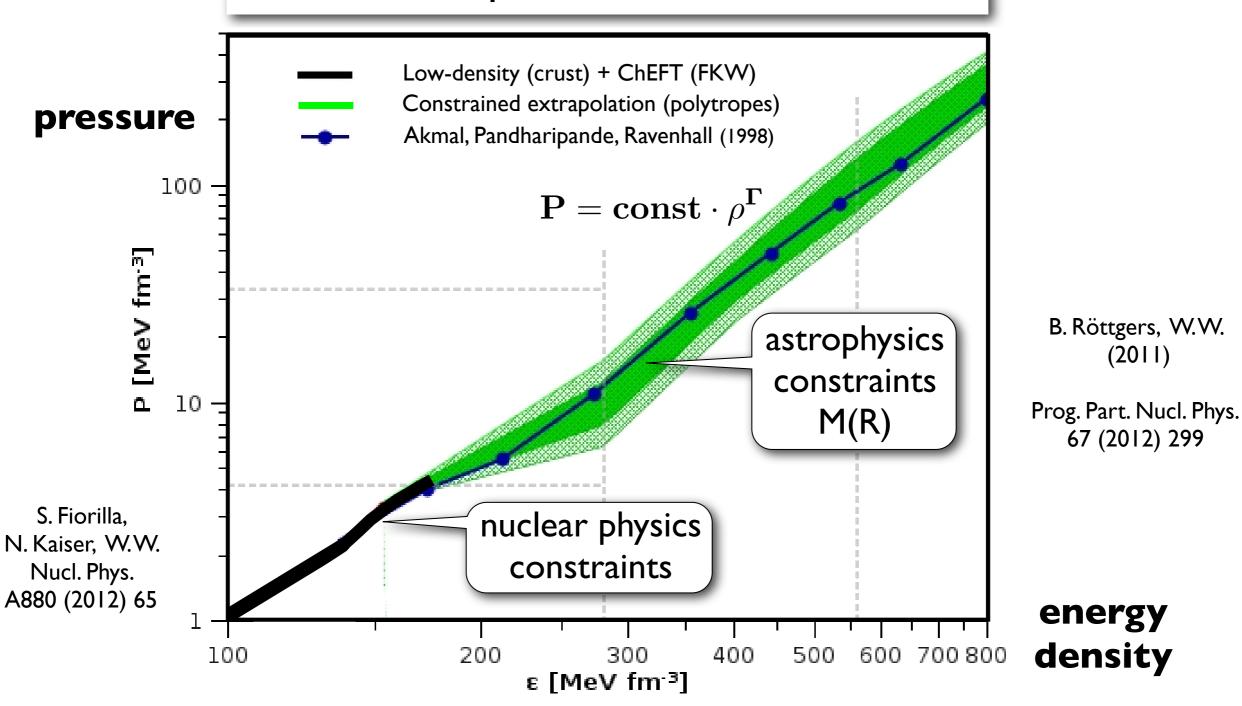
- 1 Largest mass J1614 2230 (Demorest et al. 2010)
- 2 Maximum gravity XTE 1814 338 (Bhattacharyya et al. 2005)
- 3 Minimum radius RXJ1856 3754 (Trümper et al. 2004)

- 4 Radius, 90% confidence limits LMXB 47 Tuc (Heinke et al. 2006)
- 5 Largest spin frequency J1748 2446 (Hessels et al. 2006)



NEUTRON STAR MATTER

Equation of State



Including new neutron star constraints plus Chiral Effective Field Theory at lower density

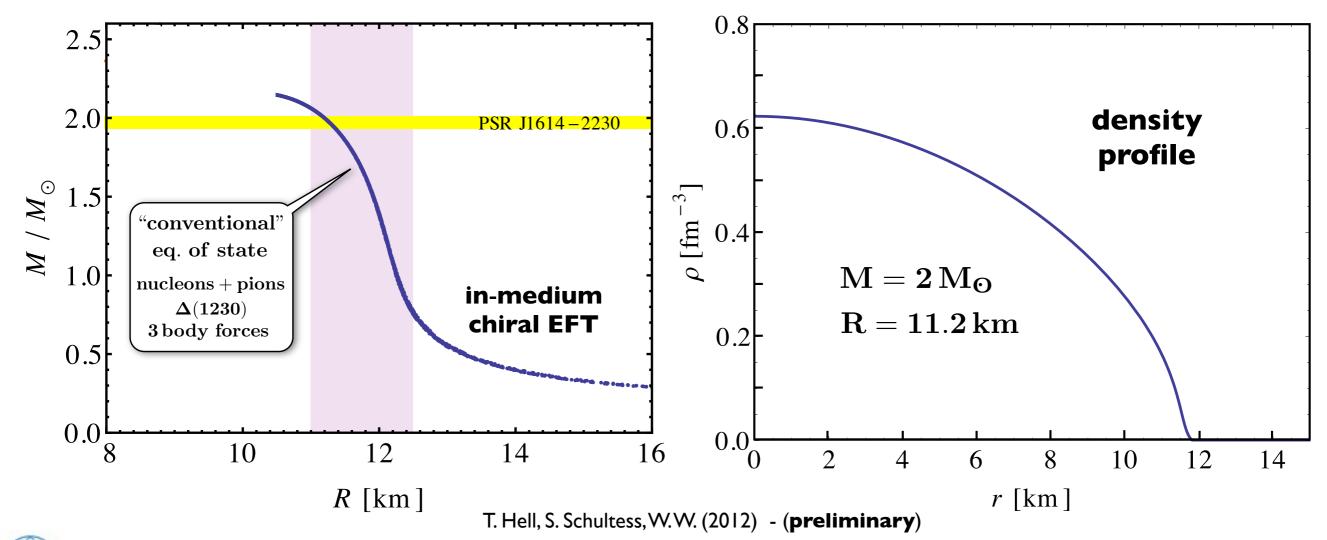




NEUTRON STAR MATTER

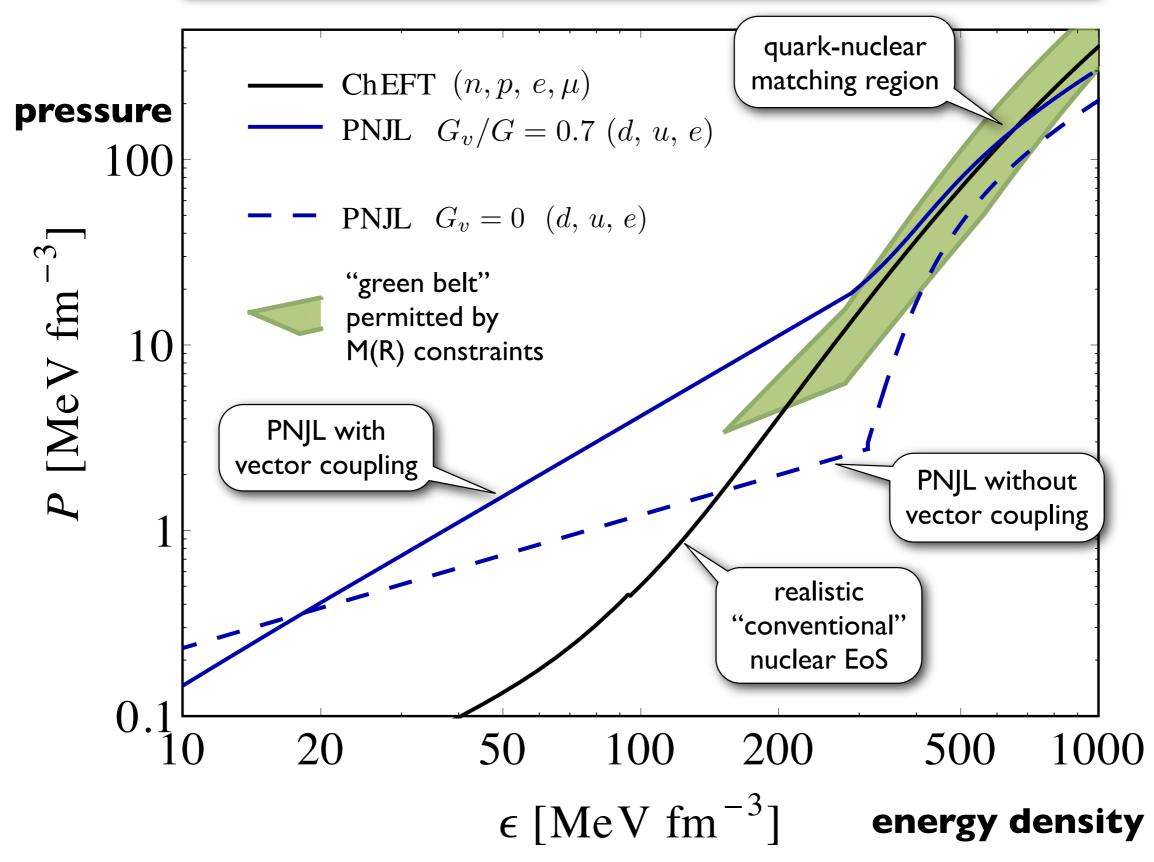
Mass - Radius relation

- Option I:
 Conventional hadronic (baryonic + mesonic) degrees of freedom
- In-medium Chiral Effective Field Theory up to 3 loops (reproducing thermodynamics of normal nuclear matter) Nucl. Phys. A 880 (2012) 65 including beta equilibrium condition ${\bf n}\leftrightarrow {\bf p}+e,\mu$





NEUTRON STAR Equation of State







NEUTRON STAR MATTER

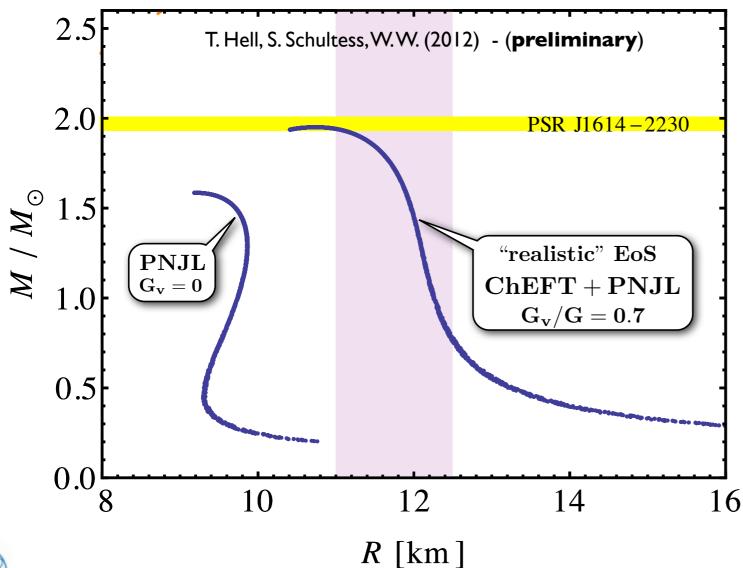
Mass - Radius relation (contd.)

Option II:

Polyakov - Nambu - Jona-Lasinio (PNJL) model (u-,d- and s-quarks as quasiparticles with dynamically generated constituent masses)

• ... features first order chiral phase transition

... produces too soft equation of state for neutron matter -> does not work!



Option III:
 Conventional
 hadronic EoS (ChEFT)
 matched at

$$ho \sim 4
ho_0$$

to **PNJL** EoS incl. **vector coupling**

soft crossover between hadronic and quark phases



SUMMARY

- Nuclear thermodynamics based on in-medium Chiral Effective Field Theory valid for $k_F < 4\pi\,f_\pi \sim 1\,GeV$
 - Proper treatment of in-medium two-pion exchange and three-body forces
 - ► Fermi liquid → interacting Fermi gas (1st order transition)
 - Realistic equations of state for nuclear and neutron matter
 - No indication of first order **chiral** transition in the density and temperature range $ho \lesssim 2\,
 ho_0\,,~T \lesssim 100\,{
 m MeV}$
 - Convergence issues: four-body correlations?
- New dense & cold matter constraints from neutron stars:
 - Mass-radius relation; observation of two-solar-mass neutron star
 - Stiff EoS required; "non-exotic" equation of state works best!



The End

thanks to:

Nino Bratovic Thomas Hell Matthias Drews
Jeremy Holt
Sebastian Schulteß

Salvatore Fiorilla Norbert Kaiser





arXiv:1202.1671 [nucl-th]

Chemical freeze-out in heavy ion collisions at large baryon densities

Stefan Floerchinger¹ and Christof Wetterich²

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We argue that the chemical freeze-out in heavy ion collisions at high baryon density is not associated to a phase transition or rapid crossover. We employ the linear nucleon-meson model with parameters fixed by the zero-temperature properties of nuclear matter close to the liquid-gas quantum phase transition. For the parameter region of interest this yields a reliable picture of the thermodynamic and chiral properties at non-zero temperature. The chemical freeze-out observed in low-energy experiments occurs when baryon densities fall below a critical value of about 15 percent of nuclear density. This region in the phase diagram is far away from any phase transition or rapid crossover.

$$\mathcal{L} = \bar{\psi}_{a} i \gamma^{\nu} (\partial_{\nu} - i g \omega_{\nu} - i \mu \delta_{0\nu}) \psi_{a}$$

$$+ \sqrt{2} h \left[\bar{\psi}_{a} \left(\frac{1 + \gamma_{5}}{2} \right) \phi_{ab} \psi_{b} + \bar{\psi}_{a} \left(\frac{1 - \gamma_{5}}{2} \right) (\phi^{\dagger})_{ab} \psi_{b} \right]$$

$$+ \frac{1}{2} \phi_{ab}^{*} (-\partial_{\mu} \partial^{\mu}) \phi_{ab} + U_{\text{mic}}(\rho, \sigma)$$

$$+ \frac{1}{4} (\partial_{\mu} \omega_{\nu} - \partial_{\nu} \omega_{\mu}) (\partial^{\mu} \omega^{\nu} - \partial^{\nu} \omega^{\mu}) + \frac{1}{2} m_{\omega}^{2} \omega_{\mu} \omega^{\mu}$$

$$\phi_{ab} = \begin{pmatrix} \frac{1}{\sqrt{2}}(\sigma + i\pi^0) & i\pi^- \\ i\pi^+ & \frac{1}{\sqrt{2}}(\sigma - i\pi^0) \end{pmatrix}$$

$$U_{\text{mic}}(\rho, \sigma) = \bar{U}(\rho) - m_{\pi}^2 f_{\pi} \sigma$$

$$\rho = \frac{1}{2}(\sigma^2 + \pi^2)$$

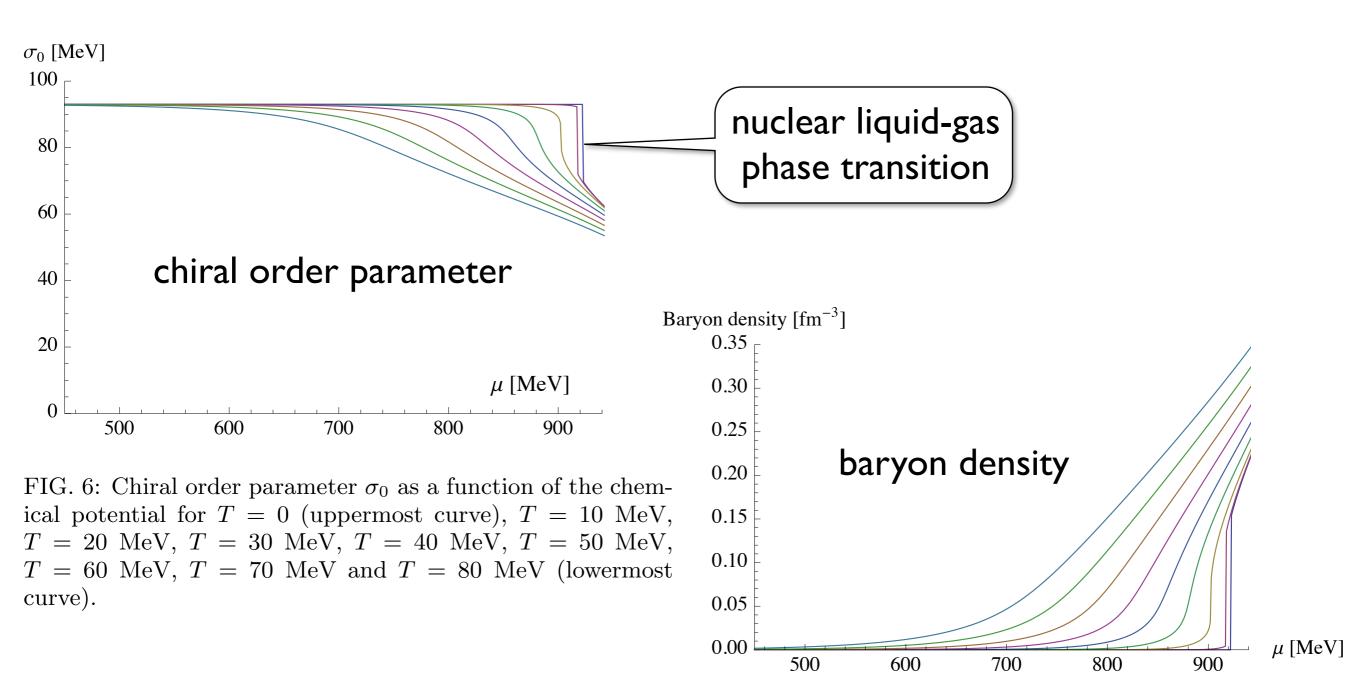


FIG. 7: Baryon number density as a function of the chemical potential for T=0 (lowermost curve), T=10 MeV, T=20 MeV, T=30 MeV, T=40 MeV, T=50 MeV, T=60 MeV, T=70 MeV and T=80 MeV (uppermost curve).

CHEMICAL FREEZE-OUT

