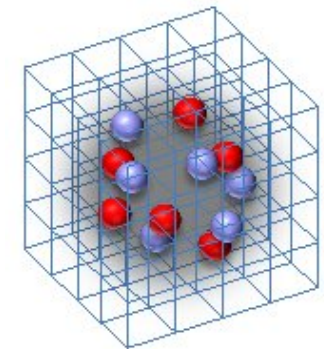


Nuclear Physics from Lattice Simulations

Ulf-G. Meißner, Univ. Bonn & FZ Jülich



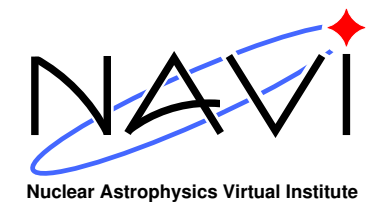
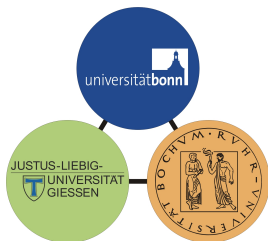
NLEFT

Supported by DFG, SFB/TR-16

and by EU, I3HP EPOS

and by BMBF 06BN9006

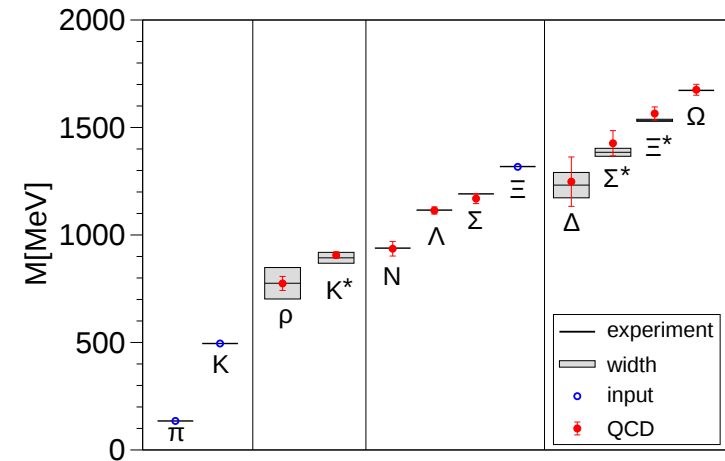
and by HGF VIQCD VH-VI-417



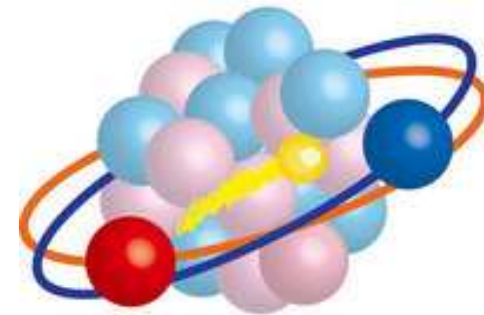
Introduction

FACETS of STRONG QCD

- quarks and gluons form hadrons
 - ⇒ **lattice QCD**
 - ⇒ **exploring the strong color force**
- nucleons and mesons form nuclei
 - ⇒ **nuclear physics**
 - ⇒ **exploring the residual color force**



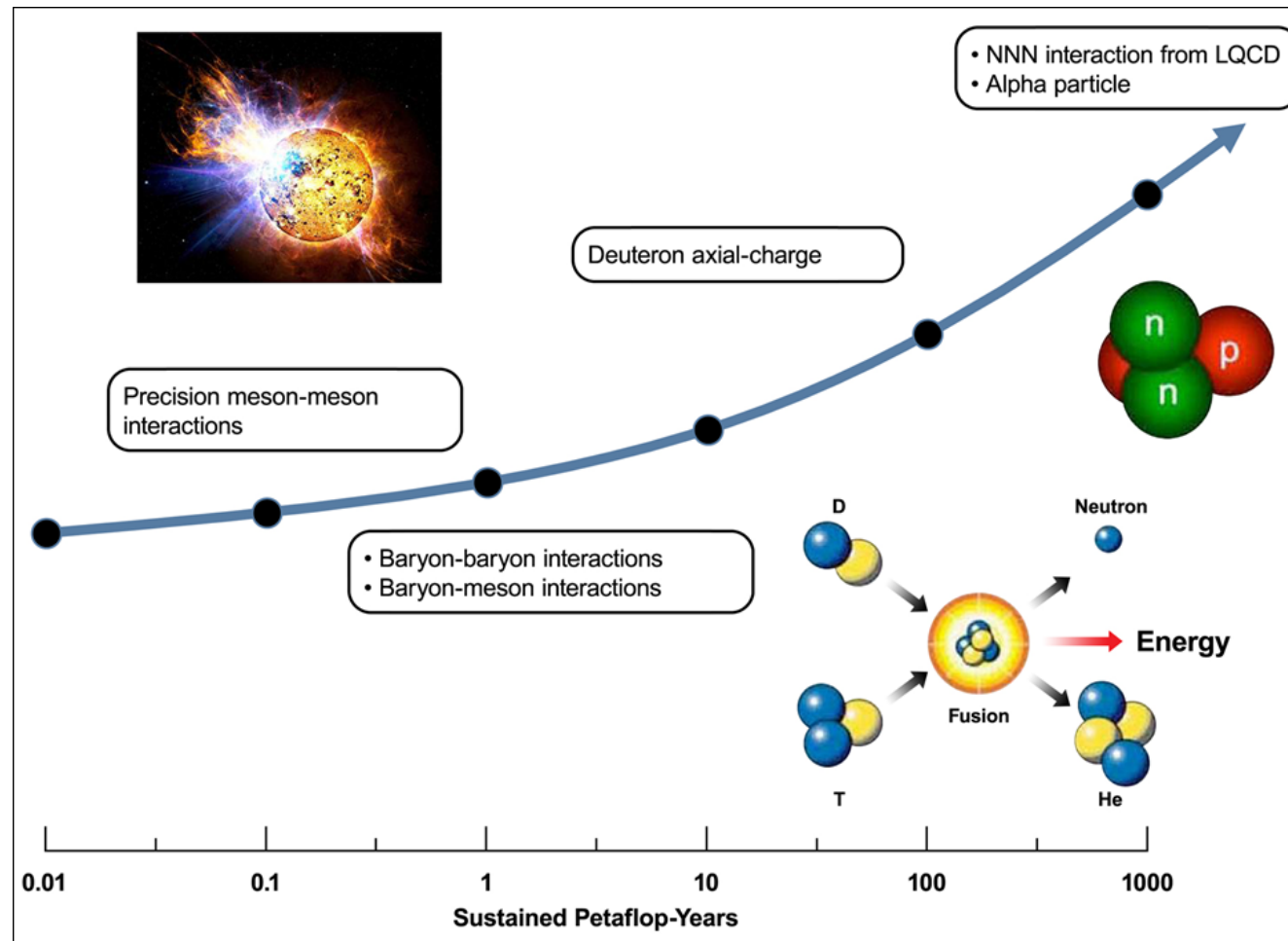
BMW collaboration (2008)



NUCLEAR PHYSICS from LATTICE QCD

- LQCD requires hundreds of sustained petaflop-yrs — not quite there yet

USDOE, Computing at the Extreme Scale, 2009



THE NUCLEAR LANDSCAPE: AIMS & METHODS

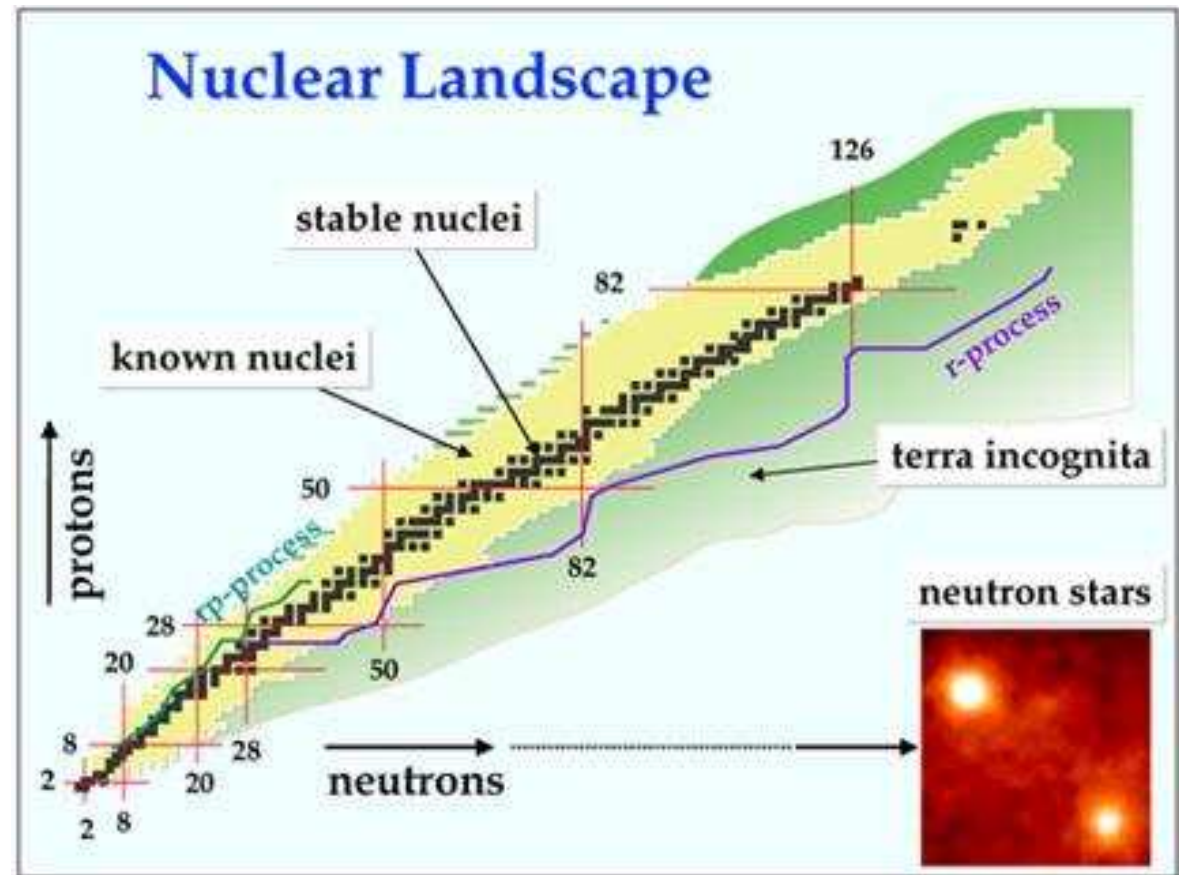
7

- Theoretical methods:

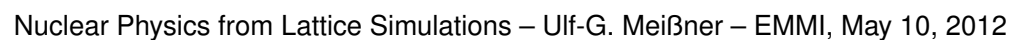
- Lattice QCD: $A = 0, 1, 2, \dots$
- NCSM, Faddeev-Yakubowsky, GFMC, ... :
 $A = 3 - 16$
- coupled cluster, ... : $A = 16 - 100$
- density functional theory, ... : $A \geq 100$

- Chiral EFT:

- provides accurate NN and 3N forces
- successfully applied in light nuclei
with $A = 2, 3, 4$
- *combine with simulations to go to larger A*



⇒ Nuclear Lattice Simulations



Effective Field Theory for Nuclear Physics

only a brief reminder → details in

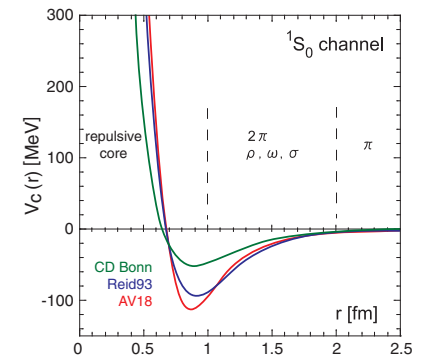
E. Epelbaum, H.-W. Hammer, UGM, Rev. Mod. Phys. **81** (2009) 1773
[arXiv:0811.1338 [nucl-th]]

Gasser, Leutwyler, **Weinberg**, van Kolck, Epelbaum, Bernard, Kaiser, UGM, . . .

- Scales in nuclear physics:

Natural: $\lambda_\pi = 1/M_\pi \simeq 1.5 \text{ fm}$ (Yukawa 1935)

Unnatural: $|a_{np}(^1S_0)| = 23.8 \text{ fm}$, $a_{np}(^3S_1) = 5.4 \text{ fm} \gg 1/M_\pi$



- this can be analyzed in a suitable EFT based on

$$\mathcal{L}_{\text{QCD}} \rightarrow \mathcal{L}_{\text{EFF}} = \mathcal{L}_{\pi\pi} + \mathcal{L}_{\pi N} + \mathcal{L}_{NN} + \dots$$

- pion and pion-nucleon sectors are perturbative in $Q/\Lambda_\chi \rightarrow$ **chiral perturbation th'y**
- $\mathcal{L}_{NN}$ collects short-distance contact terms, to be fitted
- NN interaction requires non-perturbative resummation

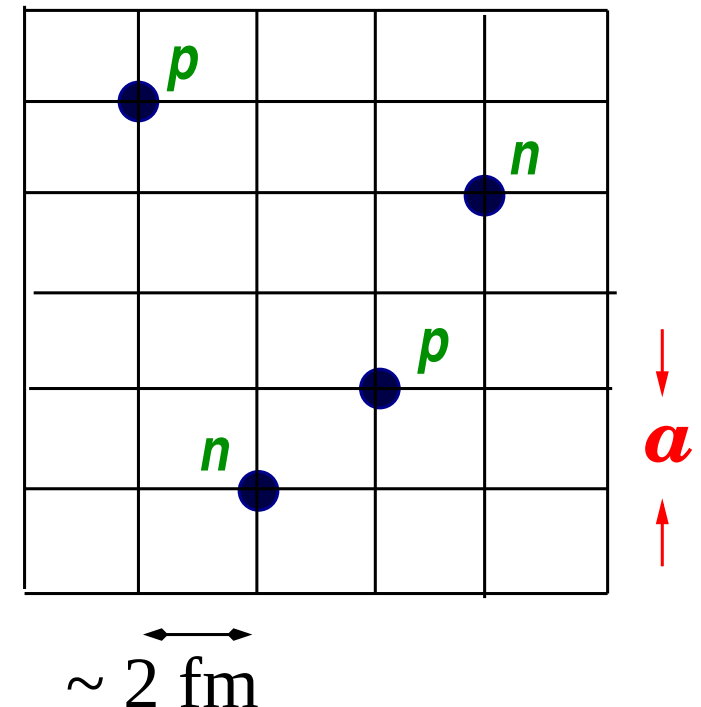
→ chirally expand $V_{\text{NN}(\text{N})}$, use in regularized LS/FY equation

Nuclear lattice simulations – Formalism –

Frank, Brockmann (1992), Koonin, Müller, Seki, van Kolck (2000), Lee, Schäfer (2004), . . .
Borasoy, Krebs, Lee, UGM, Nucl. Phys. **A768** (2006) 179; Borasoy, Epelbaum, Krebs, Lee, UGM, Eur. Phys. J. **A31** (2007) 105

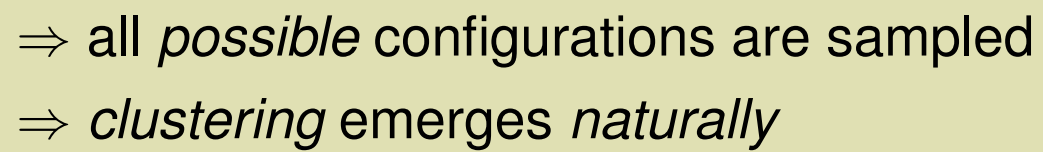
- *new method* to tackle the nuclear many-body problem
- discretize space-time $V = L_s \times L_s \times L_s \times L_t$:
nucleons are point-like fields on the sites
- discretized chiral potential w/ pion exchanges
and contact interactions
- typical lattice parameters

$$\Lambda = \frac{\pi}{a} \simeq 300 \text{ MeV [UV cutoff]}$$



- strong suppression of sign oscillations due to approximate Wigner SU(4) symmetry
- hybrid Monte Carlo & transfer matrix (similar to LQCD)

J. W. Chen, D. Lee and T. Schäfer, Phys. Rev. Lett. **93** (2004) 242302



TRANSFER MATRIX METHOD

- Correlation-function for A nucleons: $Z_A(t) = \langle \Psi_A | \exp(-tH) | \Psi_A \rangle$
with Ψ_A a Slater determinant for A free nucleons

Euclidean time

- Ground state energy from the time derivative of the correlator

$$E_A(t) = -\frac{d}{dt} \ln Z_A(t)$$

→ ground state filtered out at large times: $E_A^0 = \lim_{t \rightarrow \infty} E_A(t)$

- Expectation value of any normal-ordered operator $\mathcal{O}$

$$Z_A^{\mathcal{O}} = \langle \Psi_A | \exp(-tH/2) \mathcal{O} \exp(-tH/2) | \Psi_A \rangle$$

$$\lim_{t \rightarrow \infty} \frac{Z_A^{\mathcal{O}}(t)}{Z_A(t)} = \langle \Psi_A | \mathcal{O} | \Psi_A \rangle$$

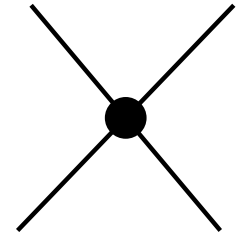
- $$\langle \Psi_A | \mathcal{O} | \Psi_A \rangle = \lim_{t \rightarrow \infty} \frac{\langle \Psi_A | \exp(-tH/2) \mathcal{O} \exp(-tH/2) | \Psi_A \rangle}{\langle \Psi_A | \exp(-tH) | \Psi_A \rangle}$$

-
- operator insertion for expectation value
- $\langle \Psi_{Z,N}^{\text{free}} |$ $SU(4) \pi$ full LO O full LO $SU(4) \pi$ $|\Psi_{Z,N}^{\text{free}} \rangle$
- $2L_{t_o} + L_{t_i}$ $L_{t_o} + L_{t_i}$ $L_{t_o} + L_{t_i}/2$ L_{t_o} 0
- inexpensive filter
- inexpensive filter

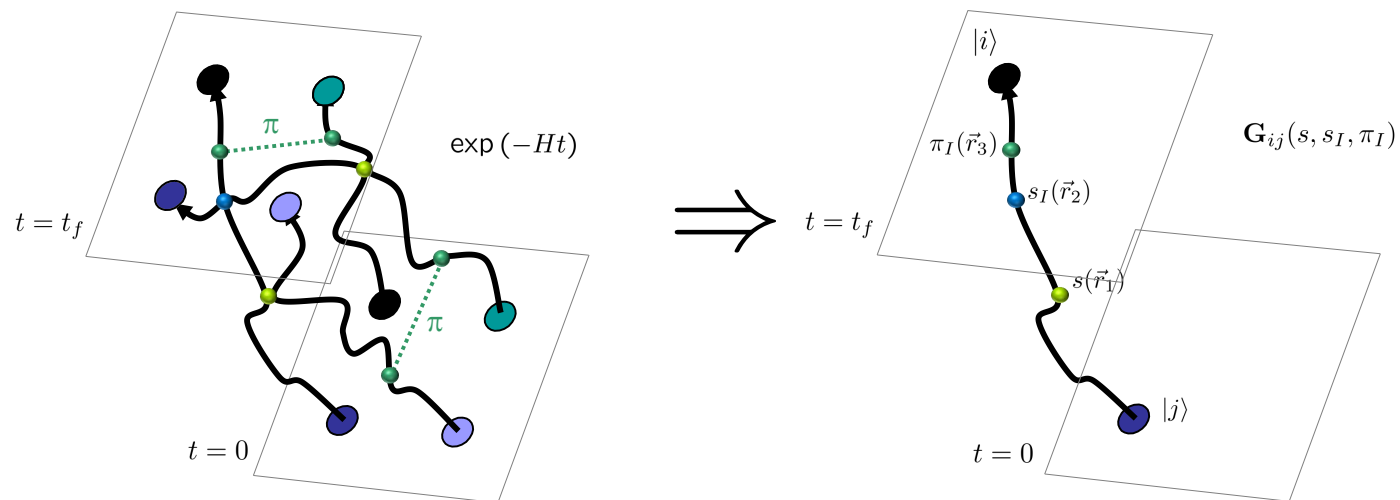
MONTE CARLO with AUXILIARY FIELDS

- Contact interactions represented by auxiliary fields s, s_I

$$\exp(\rho^2/2) \propto \int_{-\infty}^{+\infty} ds \exp(-s^2/2 - s\rho), \quad \rho \sim N^\dagger N$$



- Correlation function = path-integral over pions & auxiliary fields



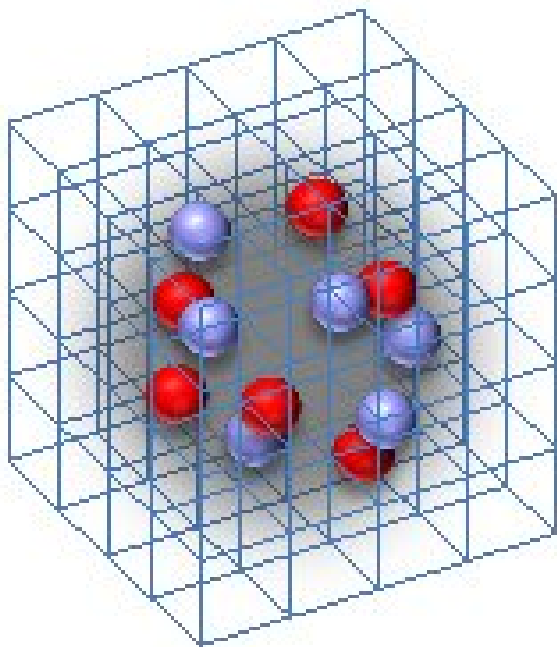
-
- A wooden dice cup, made of light-colored wood with a visible grain, is tipped over on its side. Six white dice with black pips are scattered on the surface in front of the cup's opening. The surface is covered with a light blue and green patterned cloth. The background is a plain, light-colored wall.



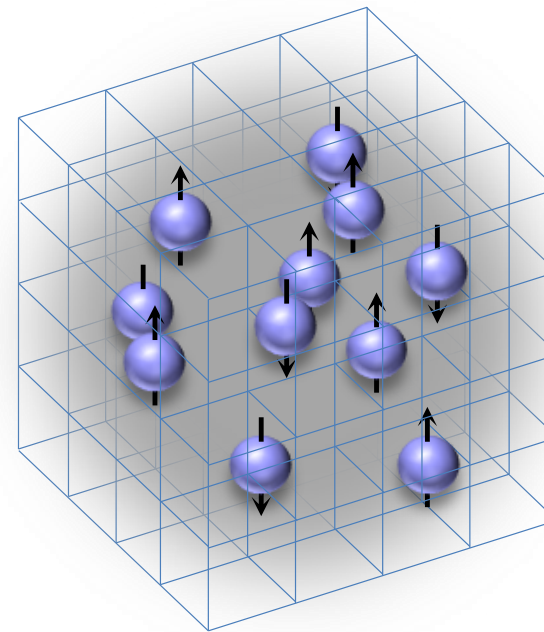
Nuclear lattice simulations

– Results –

nuclei



neutron matter

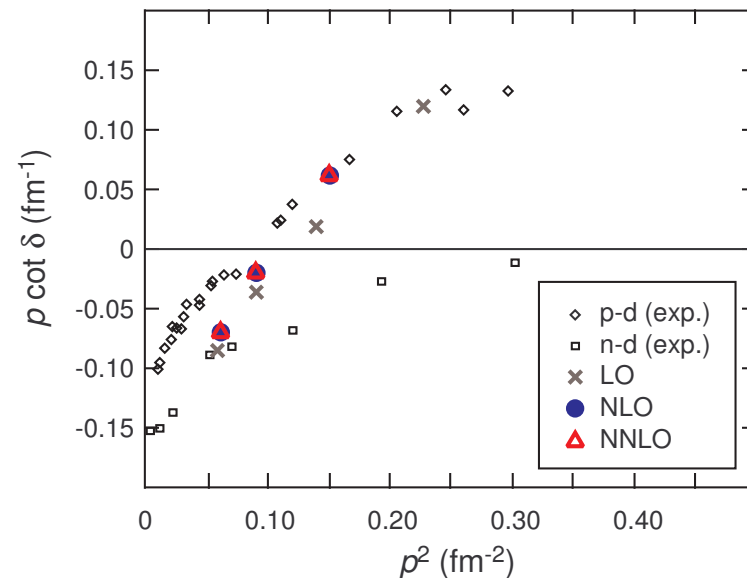
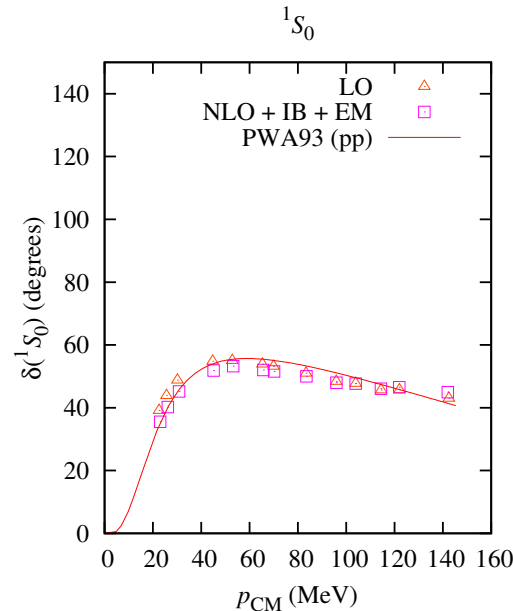
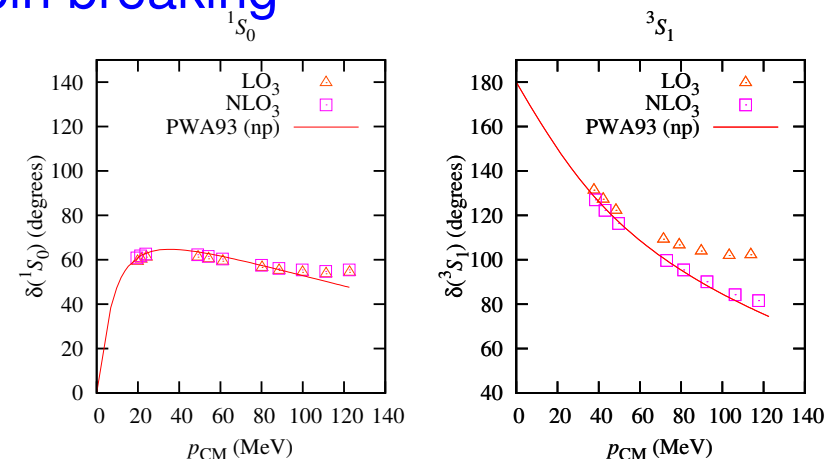


FIXING PARAMETERS & FIRST PREDICTIONS

- work at NNLO including strong and em isospin breaking
- 9 NN LECs from np scattering and Q_d
- 2 LECs for isospin-breaking (np , pp , nn)
- 2 LECs D , E related to the leading 3NF

⇒ make predictions

- pp vs np scattering
- nd spin-3/2 quartet channel
- . . .



Ground states

PREDICTIONS: TRITON & HELIUM-3

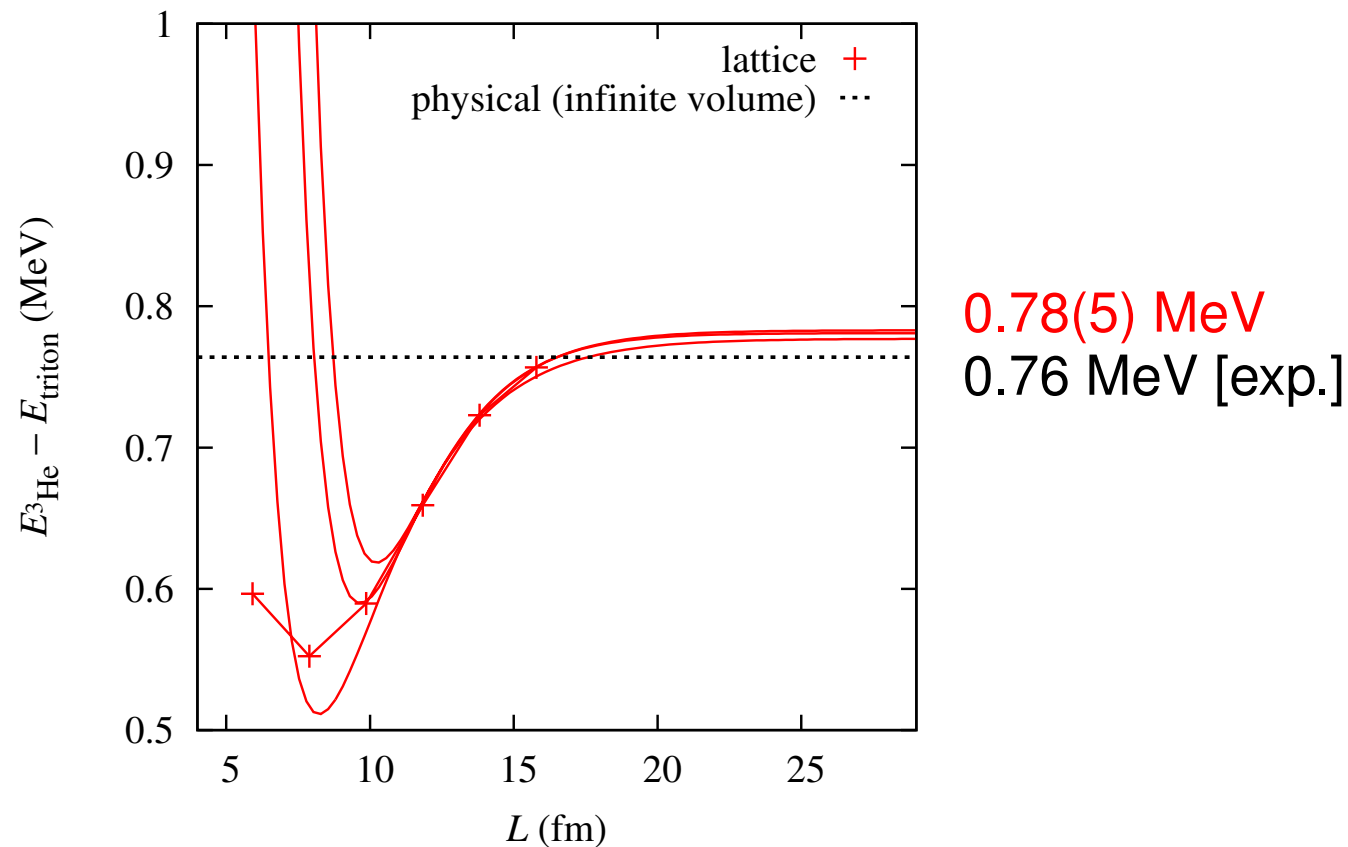
22

Epelbaum, Krebs, Lee, UGM, Phys. Rev. Lett. **104** (2010) 142501; Eur. Phys. J. **A 45** (2010) 335

- binding energies of 3N systems: $E(L) = \text{B.E.} - \frac{a}{L} \exp(-bL)$

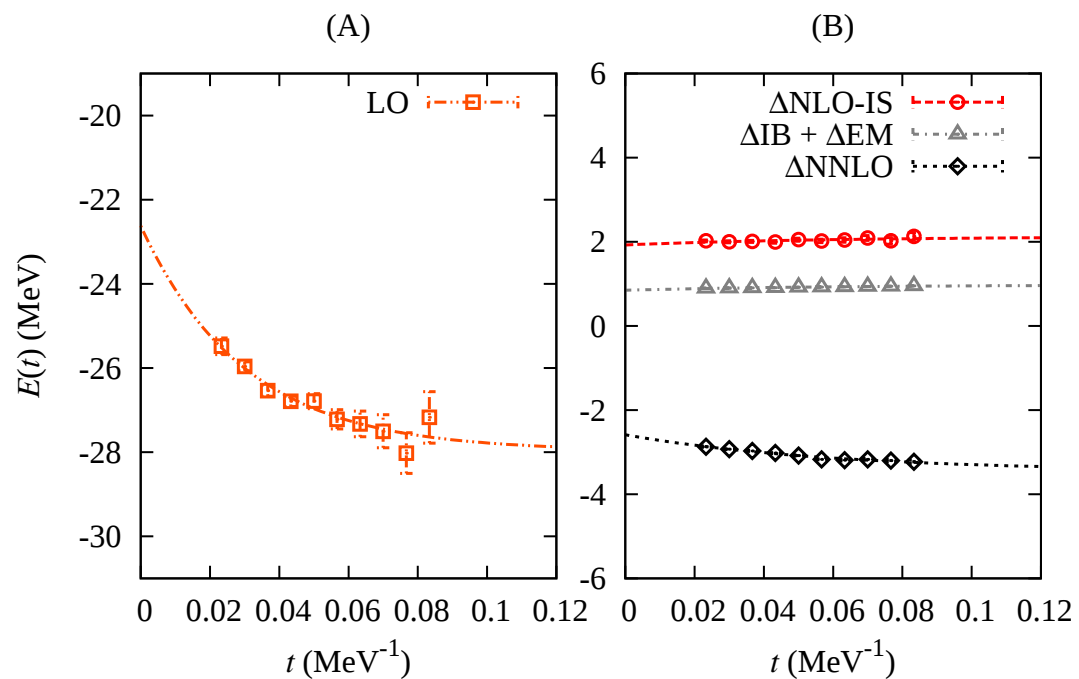
see also Hammer, Kreuzer (2011)

⇒ predict the energy difference $E(^3\text{He}) - E(^3\text{H})$



Ground state of ${}^4\text{He}$

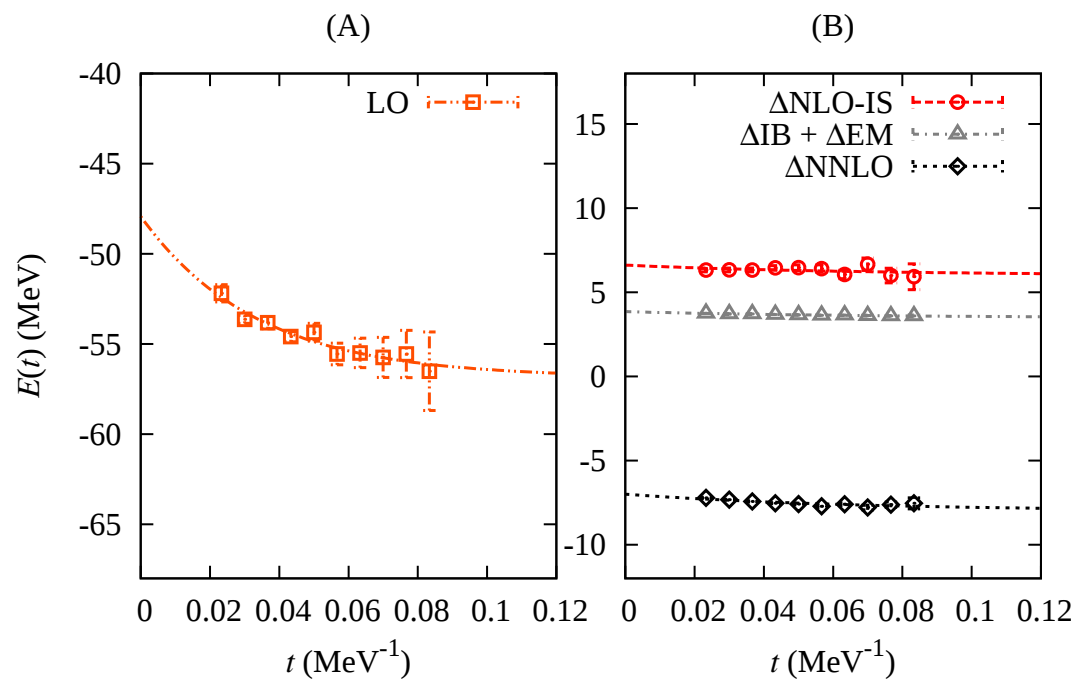
$L = 11.8 \text{ fm}$



LO ($\mathcal{O}(Q^0)$)	$-28.0(3) \text{ MeV}$
NLO ($\mathcal{O}(Q^2)$)	$-24.9(5) \text{ MeV}$
NNLO ($\mathcal{O}(Q^3)$)	$-28.3(6) \text{ MeV}$
Exp.	-28.3 MeV

Ground state of ^8Be

$L = 11.8 \text{ fm}$

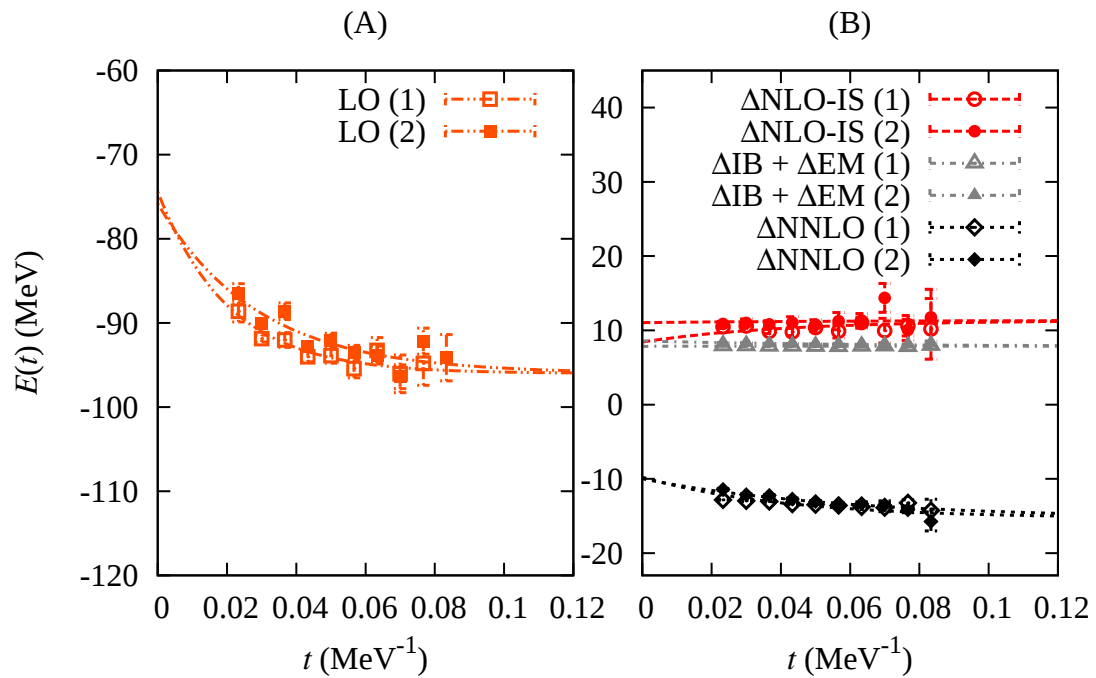


to be published

LO ($\mathcal{O}(Q^0)$)	−57(2) MeV
NLO ($\mathcal{O}(Q^2)$)	−47(2) MeV
NNLO ($\mathcal{O}(Q^3)$)	−55(2) MeV
Exp.	−56.5 MeV

Ground state of ^{12}C

$L = 11.8 \text{ fm}$

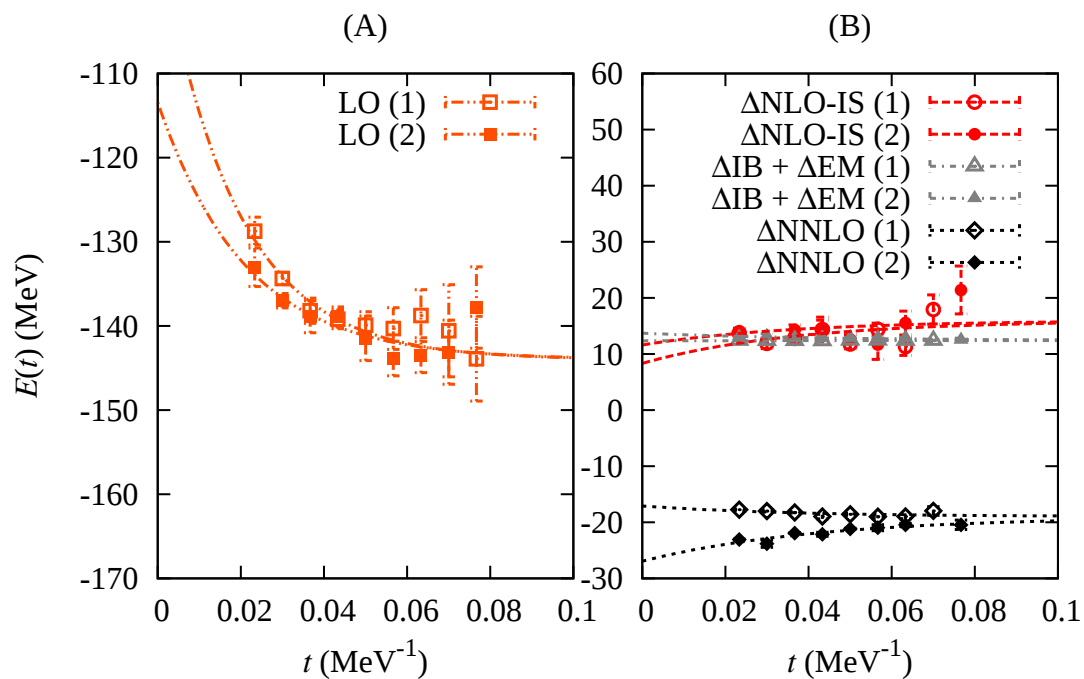


to be published

LO ($\mathcal{O}(Q^0)$)	−96(2) MeV
NLO ($\mathcal{O}(Q^2)$)	−77(3) MeV
NNLO ($\mathcal{O}(Q^3)$)	−92(3) MeV
Exp.	−92.2 MeV

Ground state of ^{16}O

$L = 11.8 \text{ fm}$



to be published

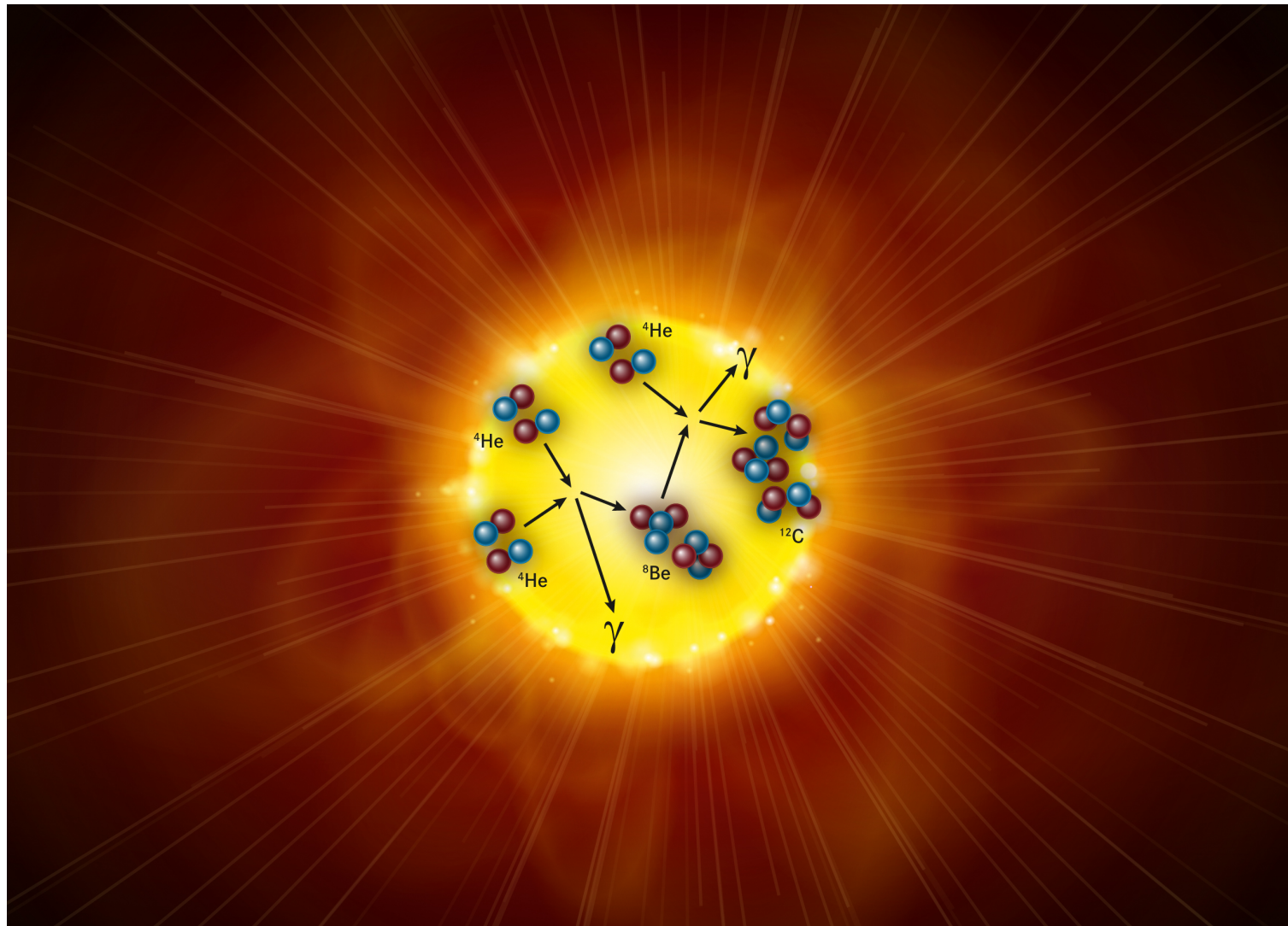
LO ($\mathcal{O}(Q^0)$)	-144(4) MeV
NLO ($\mathcal{O}(Q^2)$)	-116(6) MeV
NNLO ($\mathcal{O}(Q^3)$)	-135(6) MeV
Exp.	-127.6 MeV

SPECTRUM OF ^{12}C & the HOYLE STATE

27

Epelbaum, Krebs, Lee, UGM, Phys. Rev. Lett. **106** (2011) 192501

Viewpoint: Hjorth-Jensen, Physics **4** (2011) 38



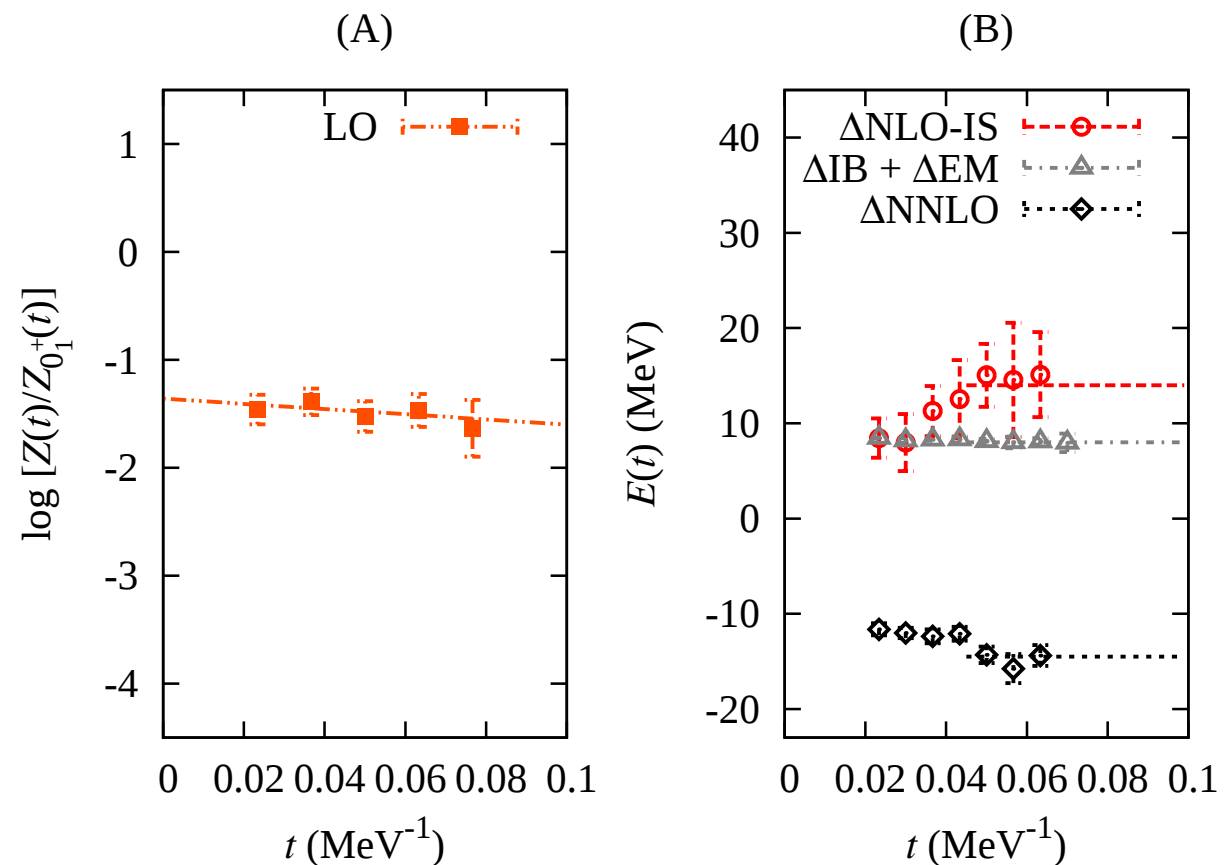
EXCITED STATES of ^{12}C

28

- Multi-channel projection MC: start with many different initial standing waves
- Extract four orthogonal energy levels with even parity and zero total momentum
- Two states have $J_z = 0$, two has $J_z = 2 \rightarrow$ clean signals!
- 2_1^+ state:

$$E(2_1^+) = -86(3) \text{ MeV} \\ [-87.7 \text{ MeV}]$$

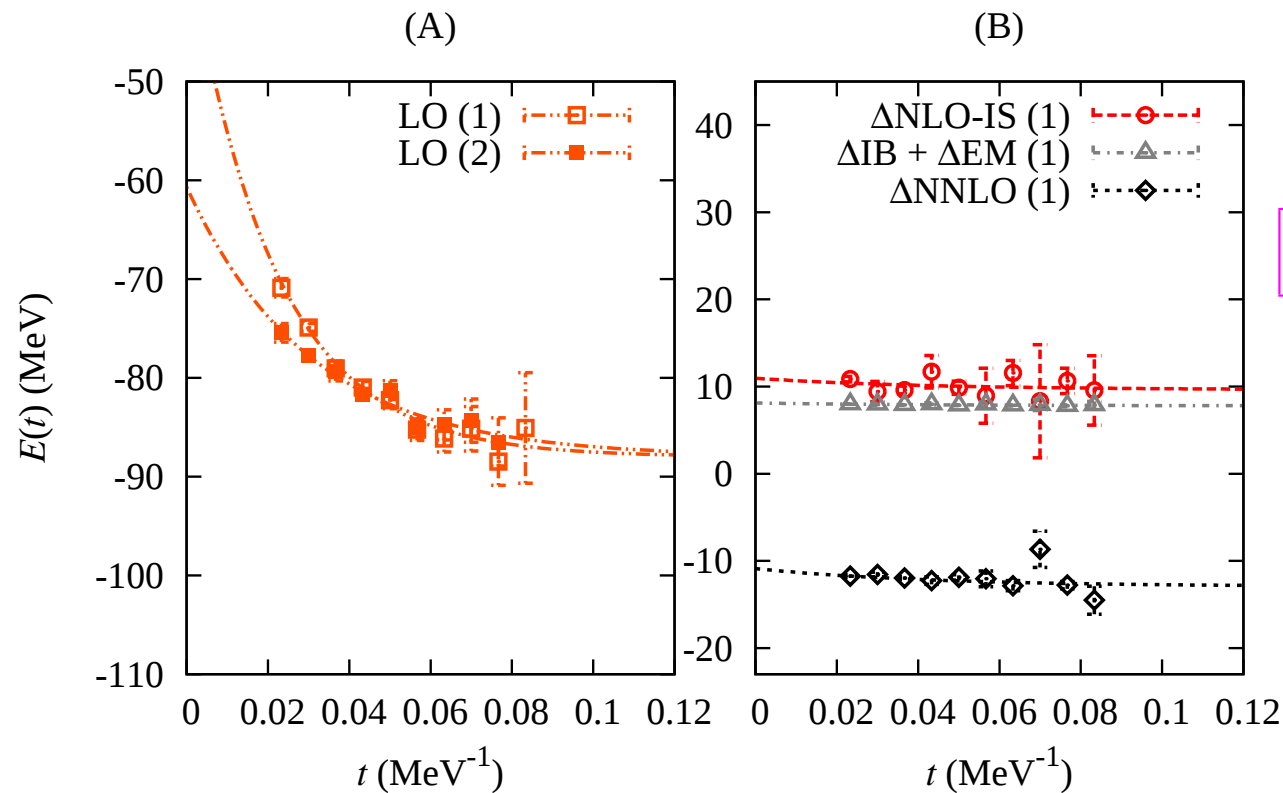
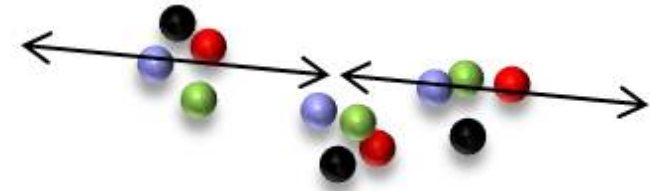
to be published



THE HOYLE STATE (0_2^+)

29

- energy: $E(0_2^+) = -84(3)$ MeV
- close to $E(^4\text{He}) + E(^8\text{Be}) = -83.8(2.0)$ MeV
- structure: “bent” alpha-chain like (not “BEC”)



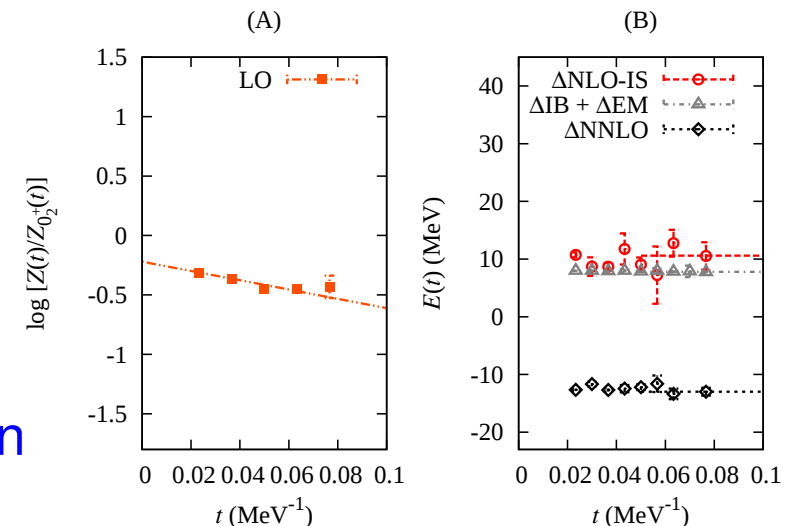
- a 2^+ state 4 MeV above the Hoyle state
- interpretation:
 - a rotational band of the Hoyle state
 - generated from excitations of the alpha-chain

a 2^+ state 3.51 MeV above the Hoyle state seen in $^{11}\text{B}(d, n)^{12}\text{C}$

Ajzenberg-Selove, Nucl. Phys. A506 (1990) 1

Bency John et al., Phys. Rev. C 68 (2003) 014305

⇒ ab initio prediction requires experimental confirmation



SPECTRUM OF ^{12}C

31

- Summarizing the results for carbon-12:

	0_1^+	2_1^+	0_2^+	2_2^+
LO	$-96(2)$ MeV	$-94(2)$ MeV	$-88(2)$ MeV	$-84(2)$ MeV
NLO	$-77(3)$ MeV	$-72(3)$ MeV	$-71(3)$ MeV	$-66(3)$ MeV
NNLO	$-92(3)$ MeV	$-86(3)$ MeV	$-84(3)$ MeV	$-79(3)$ MeV
Exp.	-92.2 MeV	-87.7 MeV	-84.5 MeV	$-80.7(4)$ MeV

- importance of consistent 2N & 3N forces
- good agreement w/ experiment, can be improved
- test of the *Anthropic Principle* possible

to be published

- consider a moving finite-size particle in a box
 $\Rightarrow$ the wavefunction picks up a *topological phase*
 i.e. the particle touches all walls $\rightarrow$ energy shift
- e..g a moving dimer with momentum $\vec{k}$
 (b.s. of two particles with masses m_1, m_2)

$$\Rightarrow \boxed{\frac{\Delta E_{\vec{k}}(L)}{\Delta E_{\vec{0}}(L)} = \sum_{i=1,2,3} \cos(2\pi\alpha k_i) \equiv \tau(\vec{k}, \alpha)}$$

$$\alpha = \frac{m_1}{m_1 + m_2}$$



- ↪ result is **universal**
- ↪ easily generalized to other waves and more constituents
- ↪ easily generalized to the scattering of composite particles

FERMION-DIMER SCATTERING

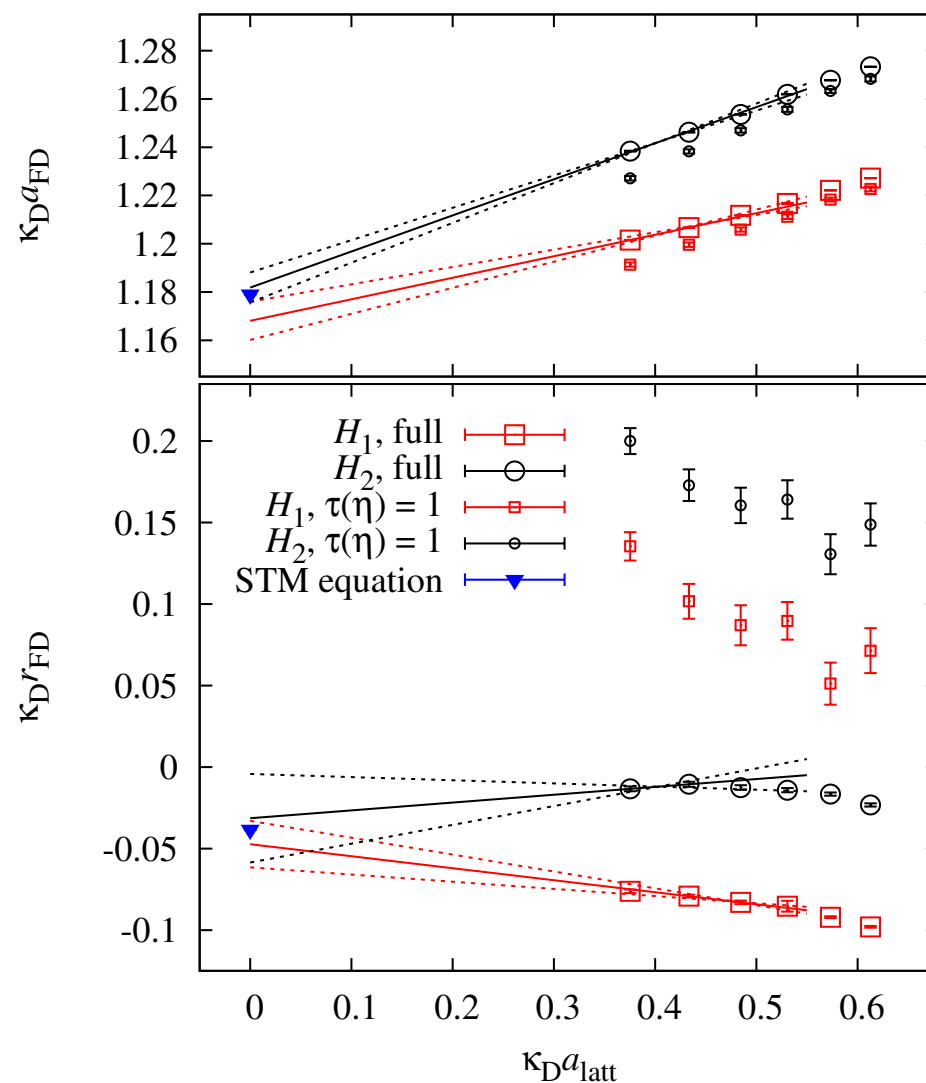
33

Bour, Hammer, König, Lee, UGM, Phys. Rev. D **84** (2011) 091503(R)

- consider fermion-dimer scattering on the lattice
- consider two lattice Hamiltonians $\mathcal{H}_{1,2}$ that reproduce the same continuum limit
Bour et al., Phys.Rev. A **83** (2011) 063619
- compare with exact solution of the STM eq.
Skorniakov, Ter-Martirosian (1957)

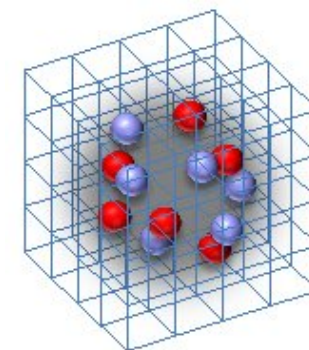
⇒ topological volume factor τ
must be included for obtaining
the correct continuum result !

- extension to bound nuclear systems from lattice QCD
Davoudi, Savage, Phys. Rev. D **84** (2011) 114502



SUMMARY & OUTLOOK

- Nuclear lattice simulations as a new quantum many-body approach
- Formulate continuum EFT on space-time lattice $V = L_s \times L_s \times L_s \times L_t$
- New method to extract phase shifts & mixing angles
- Fix parameters in few-nucleon systems $\rightarrow$ predictions
- Promising results for $A = 2, 3, 4, 8, 12, 16$ at NNLO
- ^{12}C spectrum at NNLO $\rightarrow$ **Hoyle state** & 2^+ excitation
- First ever ab initio MC calculation of ^{16}O
- Reaction theory is in reach $\rightarrow$ topological phases
- Computational scaling $\rightarrow$ medium nuclei in reach



$\Rightarrow$ larger A and higher precision

