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- Stronger suppression where the comover densities (multiplicities) are larger

- Rate equation governing the quarkonium density:

$$\tau \frac{d\rho^Q}{d\tau}(b, s, y) = -\sigma^{co-Q} \rho^{co}(b, s, y) \rho^Q(b, s, y)$$

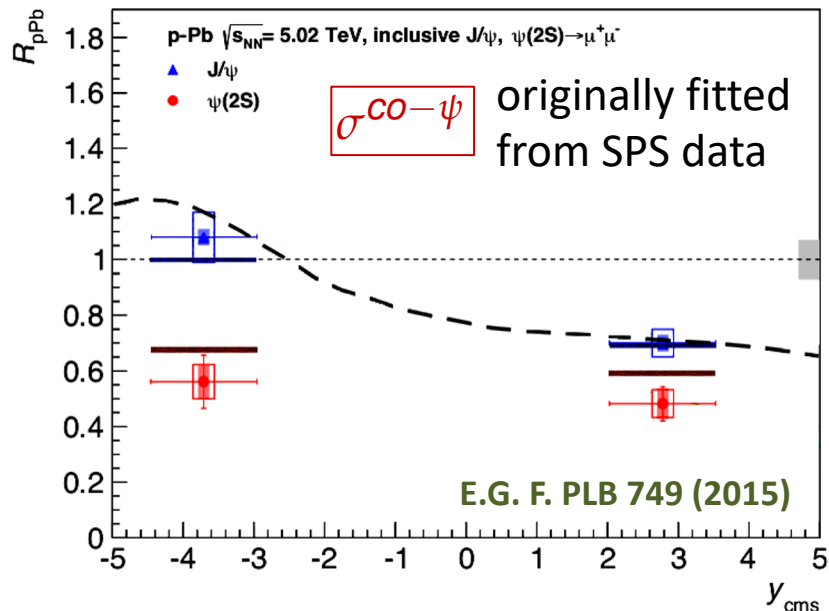
σ^{co-Q} cross section of quarkonium dissociation due to interactions with comoving medium

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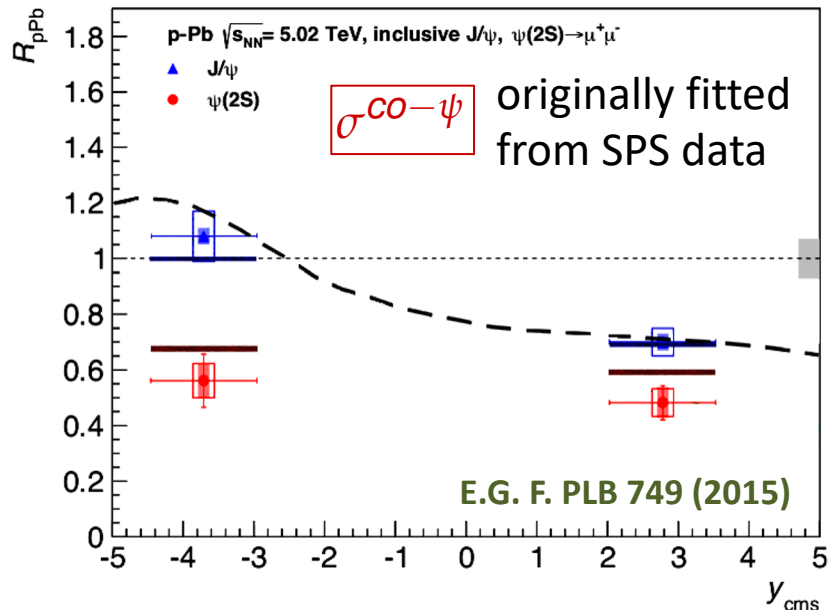


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CIM ψ results @LHC:



Our aim:

- to investigate if the relative suppression for Y states can be explained by CIM
- to apply CIM to pPb and PbPb
- to investigate the nature of the comovers

Upsilon CMS suppression in pPb

At the time of the CMS Υ PbPb analysis, **no nuclear effects** were expected **to apply differently to different states**, in particular nuclear breakup

PRL **109**, 222301 (2012)

 Selected for a **Viewpoint** in *Physics*
PHYSICAL REVIEW LETTERS

week ending
30 NOVEMBER 2012



Observation of Sequential Υ Suppression in PbPb Collisions

S. Chatrchyan *et al.**
(CMS Collaboration)

*In addition to QGP formation, differences between quarkonium production yields in PbPb and pp collisions can also arise from **cold-nuclear-matter effects** [21]. However, such effects should have a **small impact on the double ratios** reported here. Initial-state nuclear effects are expected to affect similarly each of the three Υ states, thereby canceling out in the ratio. Final-state “nuclear absorption” becomes weaker with increasing energy [22] and is expected to be negligible at the LHC [23].*

$\frac{[Y(nS)/Y(1S)]_{ij}}{[Y(nS)/Y(1S)]_{pp}}$	2S	3S
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pPb	0.83 ± 0.05 (stat.) ± 0.05 (syst.)	0.71 ± 0.08 (stat.) ± 0.09 (syst.)

- CMS assumption **contradicted** by their **pPb data** **CMS JHEP04(2014)103**
- If this relative suppression can be attributed to comover effects, how does that **translate to PbPb collisions?** [comover suppression is related to the multiplicity]

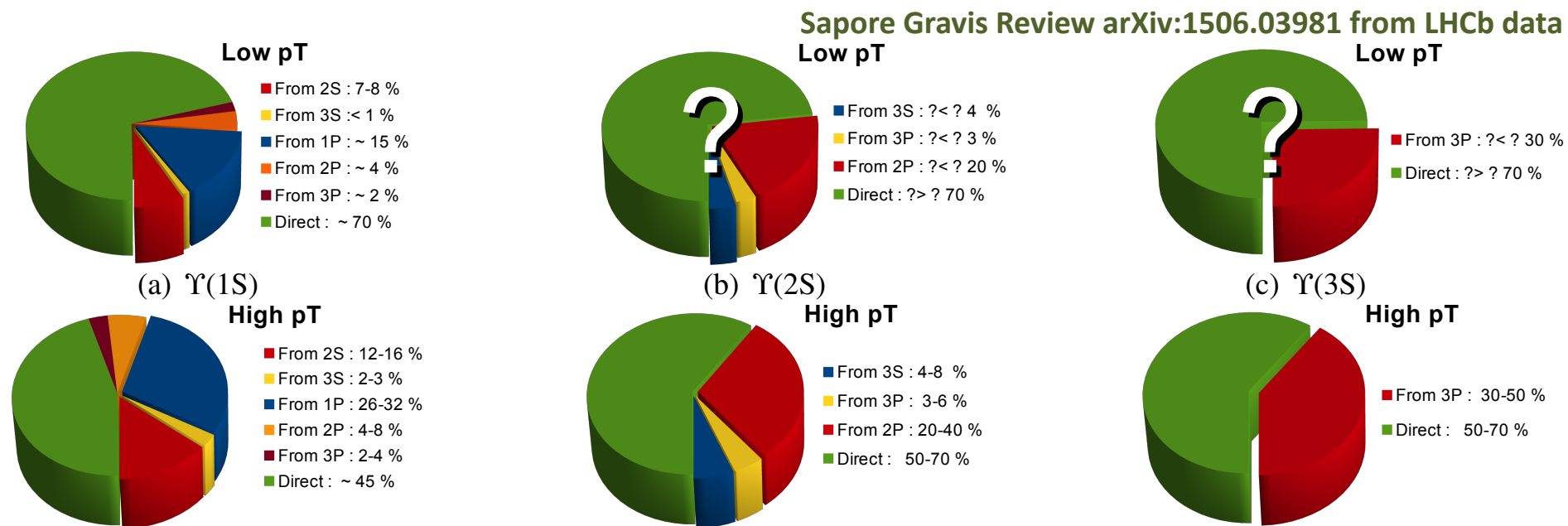
A closer look into the Y family and its feed down structure

- The bottomonium family is much richer than the charmonium one
- χ_b'' first particle discovered at the LHC ATLAS PRL 108 (2012) 152001
- It allows for a much finer studies with 3 Y states (decaying into dimuons)
- It comprises excited states which are not too fragile [as opposed to e.g. the ψ']

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Feed-down structure at **low** p_T is quite different than CDF measurement at $p_T > 8\text{GeV}$



- $\Upsilon(3S)$ is **far from being 100% direct**
- In the region of the Υ PbPb and pPb data, the $\Upsilon(1S)$ is **not 50% direct**

Setting the scene for the bottomonium family

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We take:

$$\sigma^{co-Q_{b\bar{b}}} = \sigma_{\text{geom}} \left(1 - \frac{E_{\text{Binding}}}{E_{co}}\right)^n$$

E. G. F., J.P. Lansberg, arXiv:1804.04474

$$\sigma_{\text{geom}} \equiv \pi r_{Q_{b\bar{b}}}^2$$

$E_{\text{Binding}} \equiv 2M_B - M_{Q_{b\bar{b}}}$, i.e. the threshold energy to break the bound state

$E^{co} = \sqrt{p^2 + m_{co}^2}$ the average energy of the comovers in the quarkonium rest frame

We average over B-E phase space distribution of the comovers $1/(e^{E^{co}/T_{eff}} - 1)$

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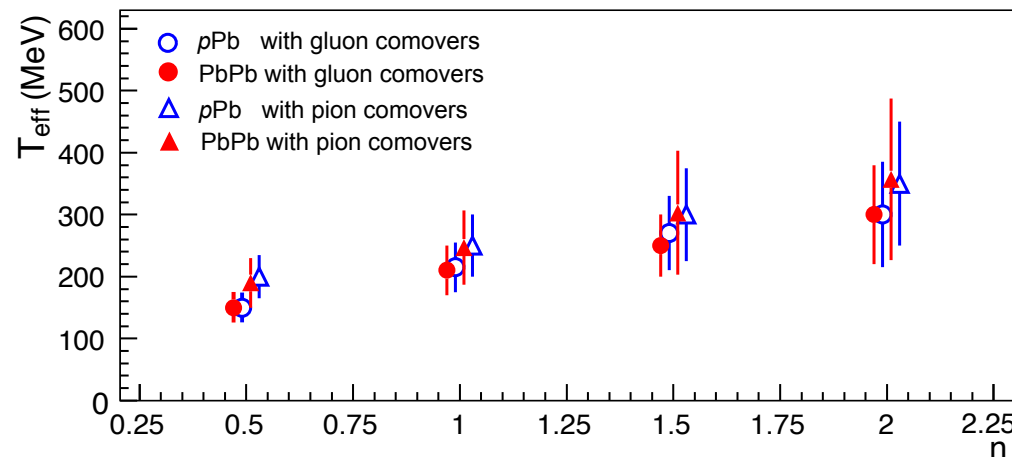
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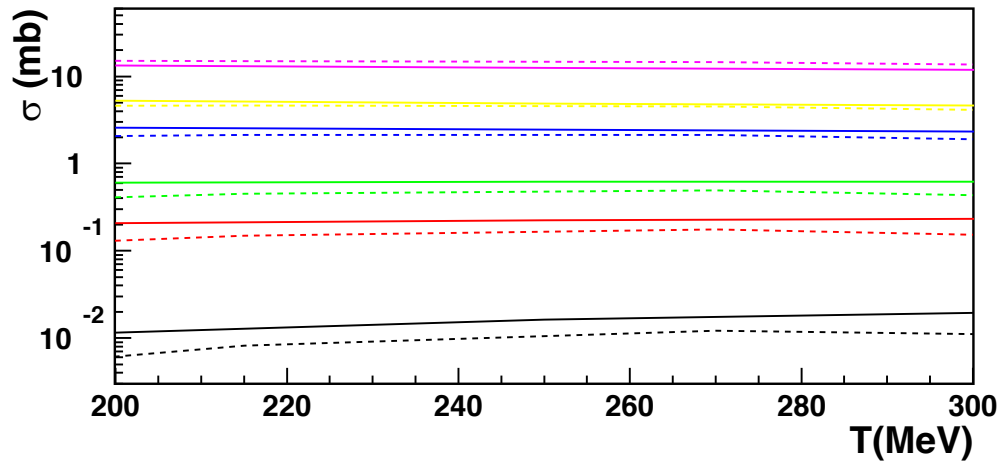


Using pPb CMS and ATLAS data at 5.02 TeV we fit T_{eff} and n . Also with PbPb CMS data

By varying n between 0.5 and 2, we obtain T_{eff} in the range from 200 to 300 MeV both for **partons** or **hadrons**

Setting the scene for the bottomonium family

- High stability in the mentioned temperature range with running n



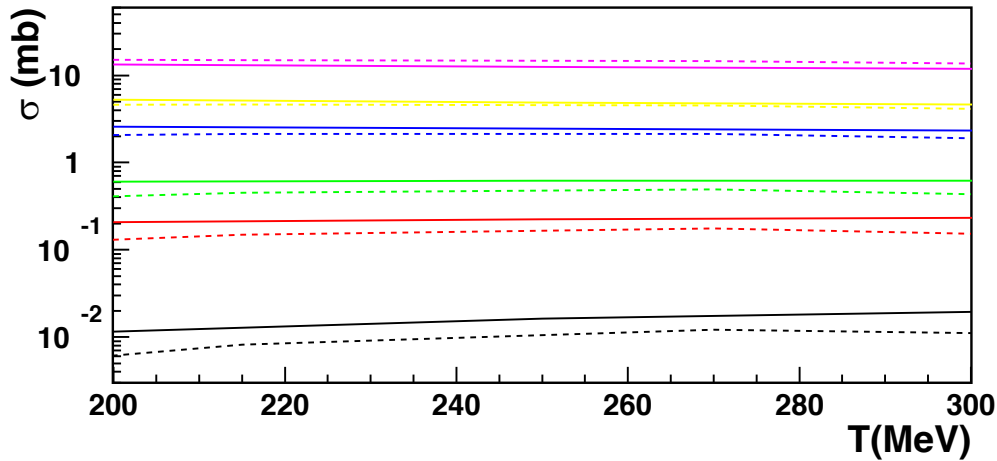
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The mean values for the dissociation cross-sections for the bottomonium family in a comover medium made of pions (continuous line) or gluons (discontinuous line).

From down to up: 1S, 1P, 2S, 2P, 3S, 3P

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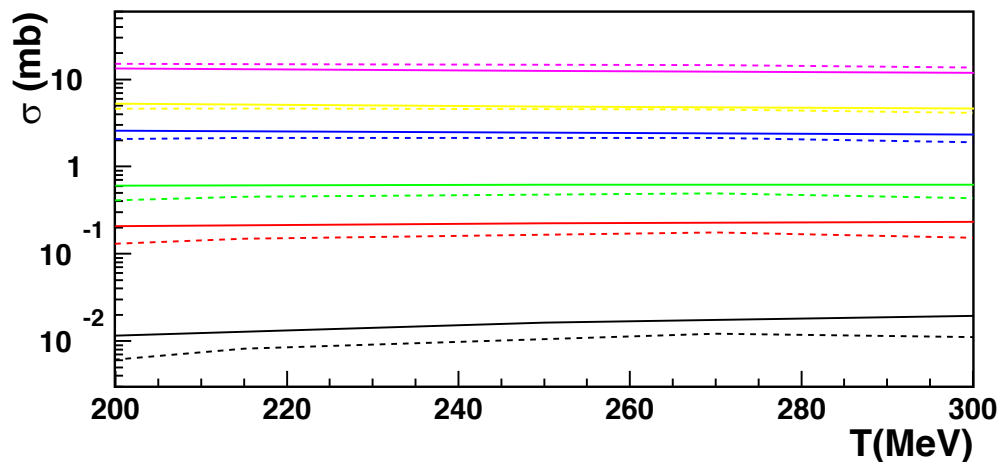
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- The feed-downs discussed above were used:

low P_T	direct	from χ_b	from Y'	from χ'_b	from Y''	from χ''_b
Y	$\sim 70\%$	$\sim 15\%$	$\simeq 8\%$	$\sim 5\%$	$\simeq 1\%$	$\sim 1\%$
Y'	$\sim 63\%$	–	–	$\sim 30\%$	$\simeq 4\%$	$\sim 3\%$
Y''	$\sim 60\%$	–	–	–	–	$\sim 40\%$

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- Varying the feed-down fractions for 2 limiting cases does not change the results
80% of direct 1S and 50% of direct 3S or 60% of direct 1S and 70% of direct 3S

Double ratio $\Upsilon(nS)/\Upsilon(1S)$ in pPb & PbPb @ 2.76 & 5.02 TeV

For $n=1$ and $T=250 \pm 50$ MeV:

	$E_{Binding}$	$r_{Q_{b\bar{b}}}$	$\sigma^{co-Q_{b\bar{b}}}$
$\Upsilon(1S)$	1100 MeV	0.14 fm	$0.02^{+0.020}_{-0.010}$ mb
χ_{B1}	670 MeV	0.22 fm	$0.23^{0.14}_{-0.12}$ mb
$\Upsilon(2S)$	540 MeV	0.28 fm	$0.61^{+0.33}_{-0.28}$ mb
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Υ pPb at 5.02 TeV

	CIM	Exp
	$-1.93 < y < 1.93$	CMS data
$\Upsilon(2S)/\Upsilon(1S)$	0.91 ± 0.03	0.83 ± 0.05 (stat.) ± 0.05 (syst.)
$\Upsilon(3S)/\Upsilon(1S)$	0.72 ± 0.02	0.71 ± 0.08 (stat.) ± 0.09 (syst.)
	$-2.0 < y < 1.5$	ATLAS data
$\Upsilon(2S)/\Upsilon(1S)$	0.90 ± 0.03	0.76 ± 0.07 (stat.) ± 0.05 (syst.)
$\Upsilon(3S)/\Upsilon(1S)$	0.71 ± 0.02	0.64 ± 0.14 (stat.) ± 0.06 (syst.)

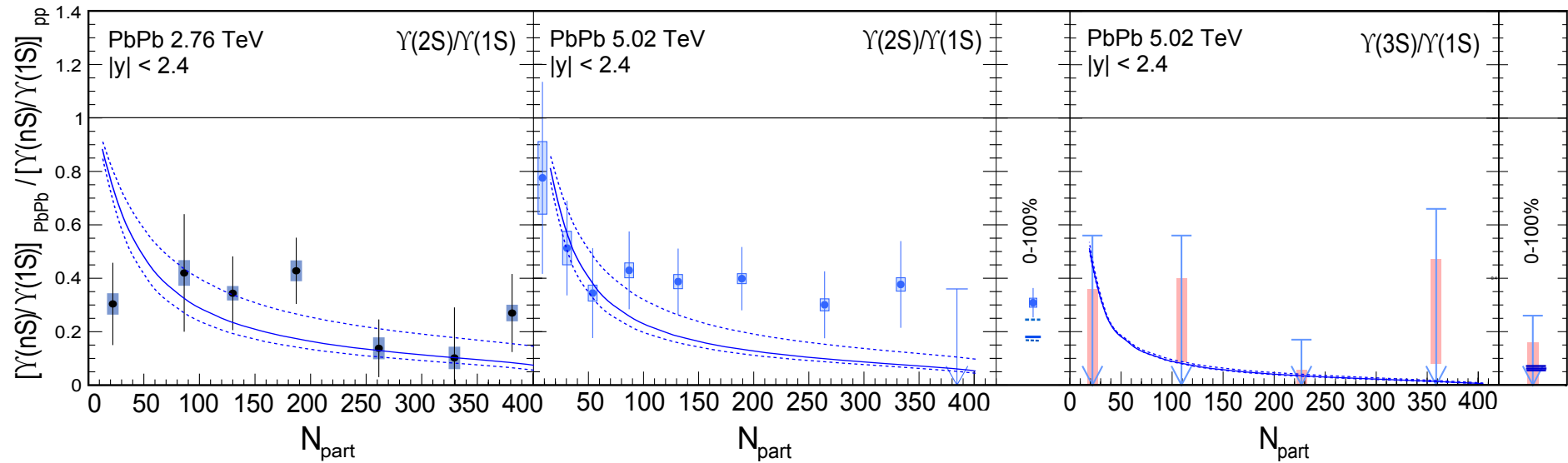
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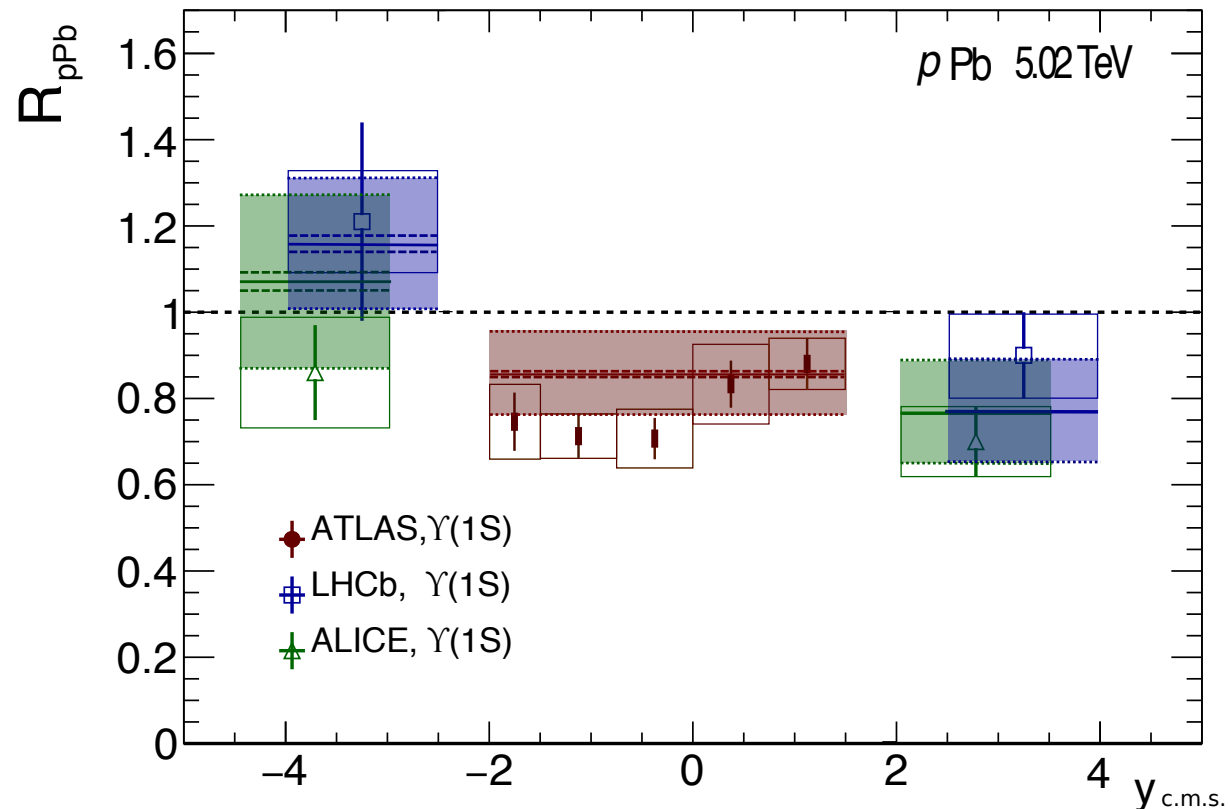


$Y(nS)/Y(1S)$ well reproduced in PbPb collisions without any other phenomena needed

Consistency check: $Y(1S)$ nuclear modification factor in pPb

- Now that the $\sigma^{c\bar{c}-Q_{b\bar{b}}}$ are fixed, we need to check the consistency with the absolute suppression of $Y(1S)$
- Other nuclear effects which cancel in the double ratio, **do not cancel** anymore, i.e. shadowing
- We take into account **nCTEQ15**
- Comovers damp down the antishadowing peak
=> **better agreement with ALICE**

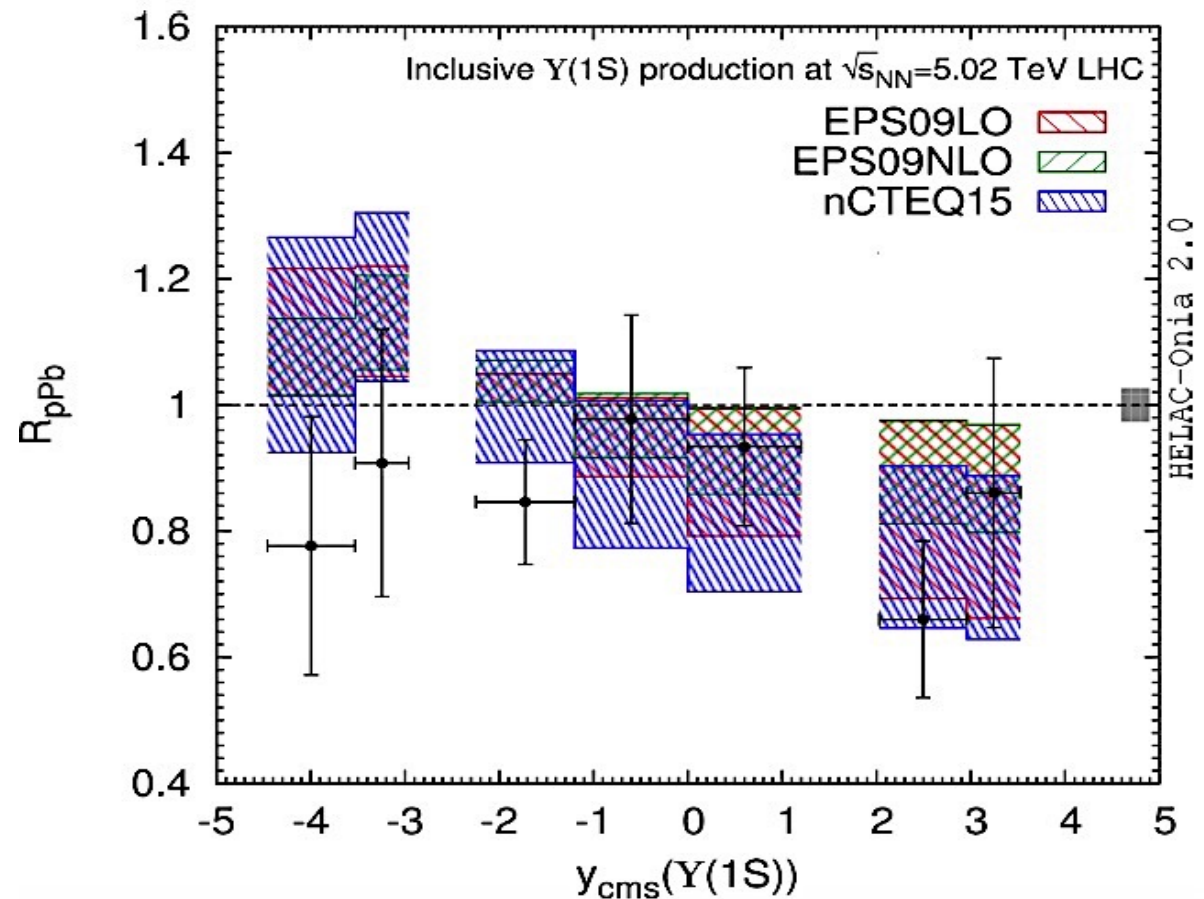
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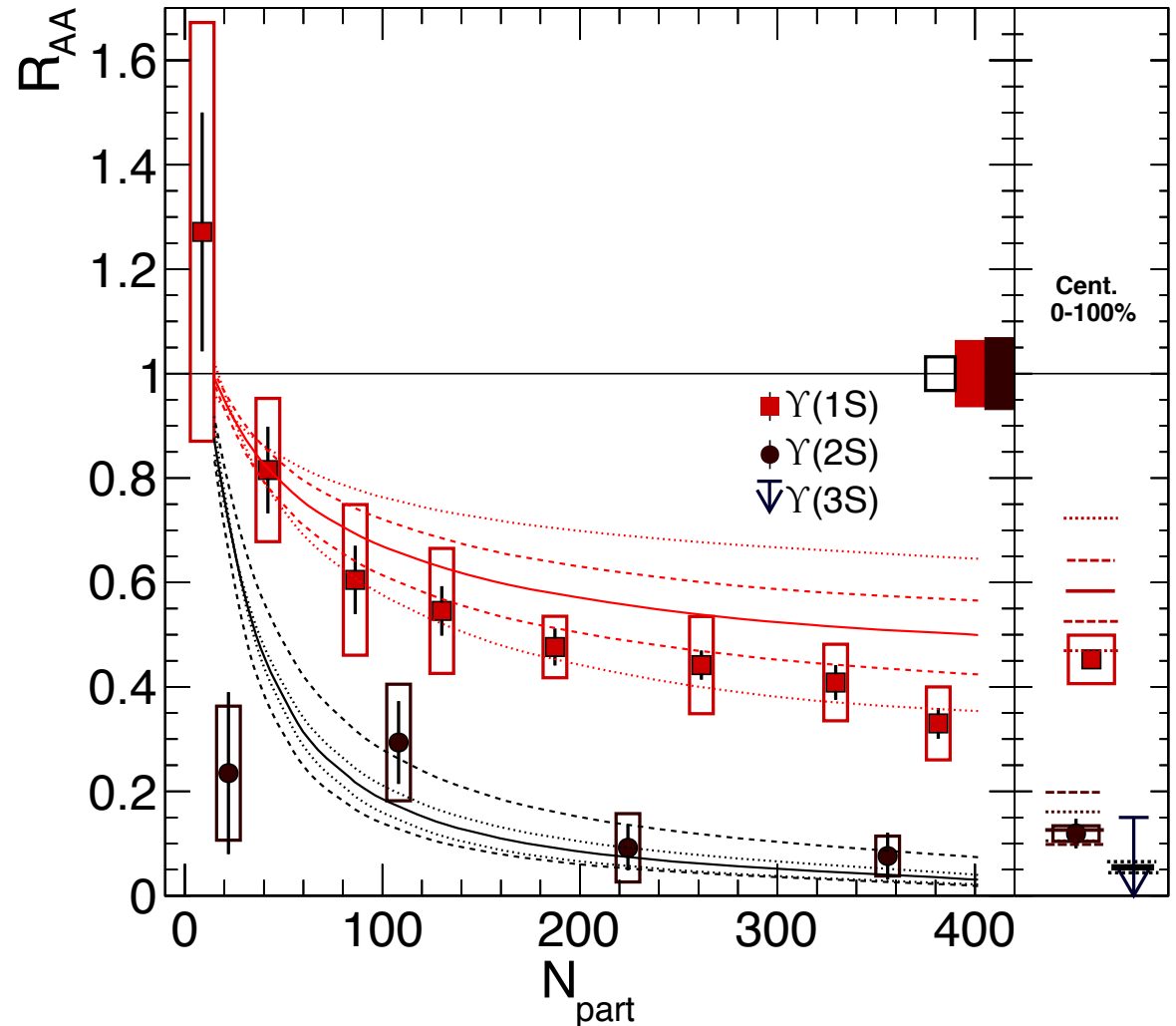
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Consistency check: R_{pPb} for $Y(1S)$ and $Y(2S)$ @ 2.76 TeV

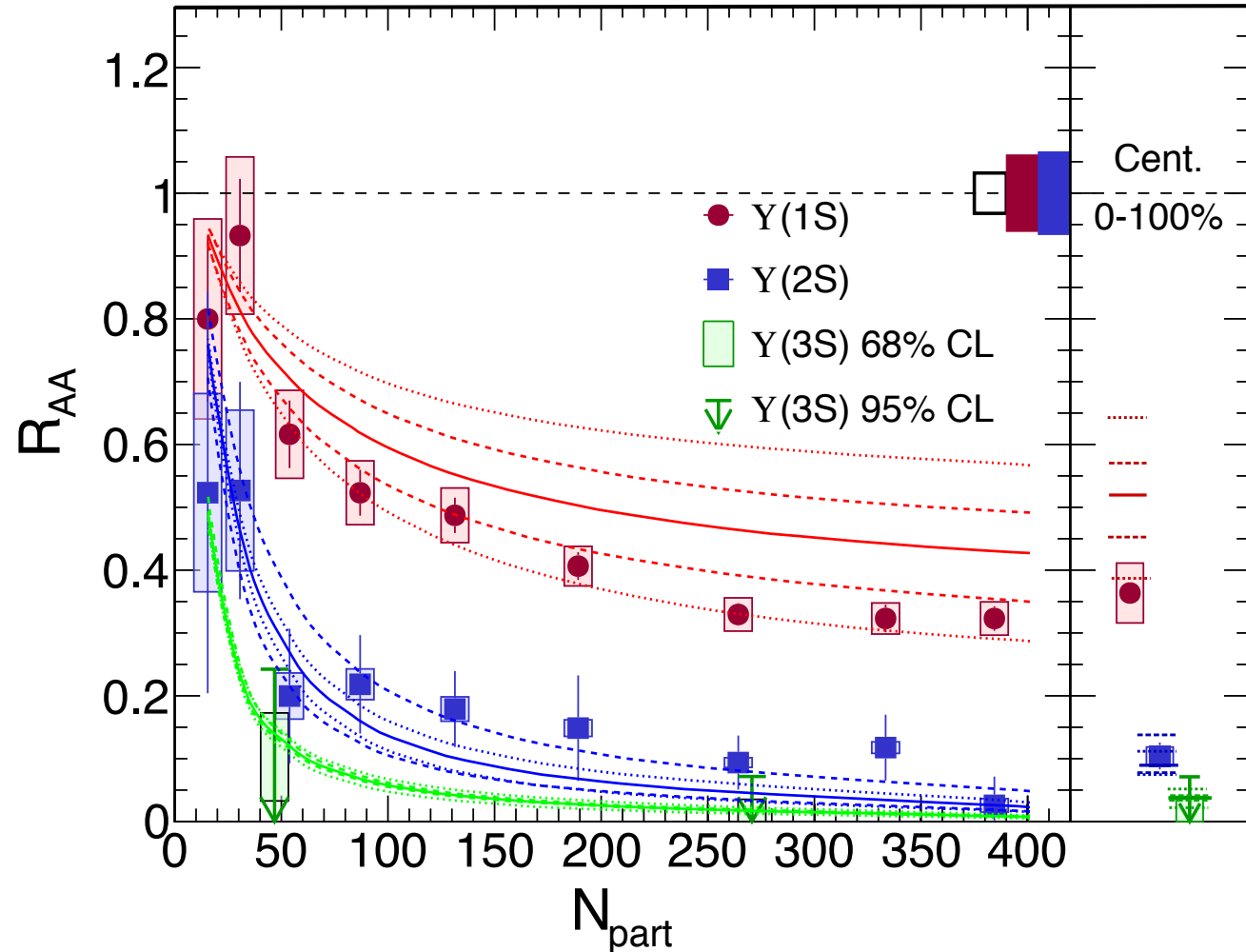
- We take into account **nCTEQ15** (as for R_{pPb})
- We do show the **significant uncertainty** of the barely known gluon nPDFs



The magnitude of suppression -taking into account nCTEQ15- is well reproduced without the need to invoke any other phenomena

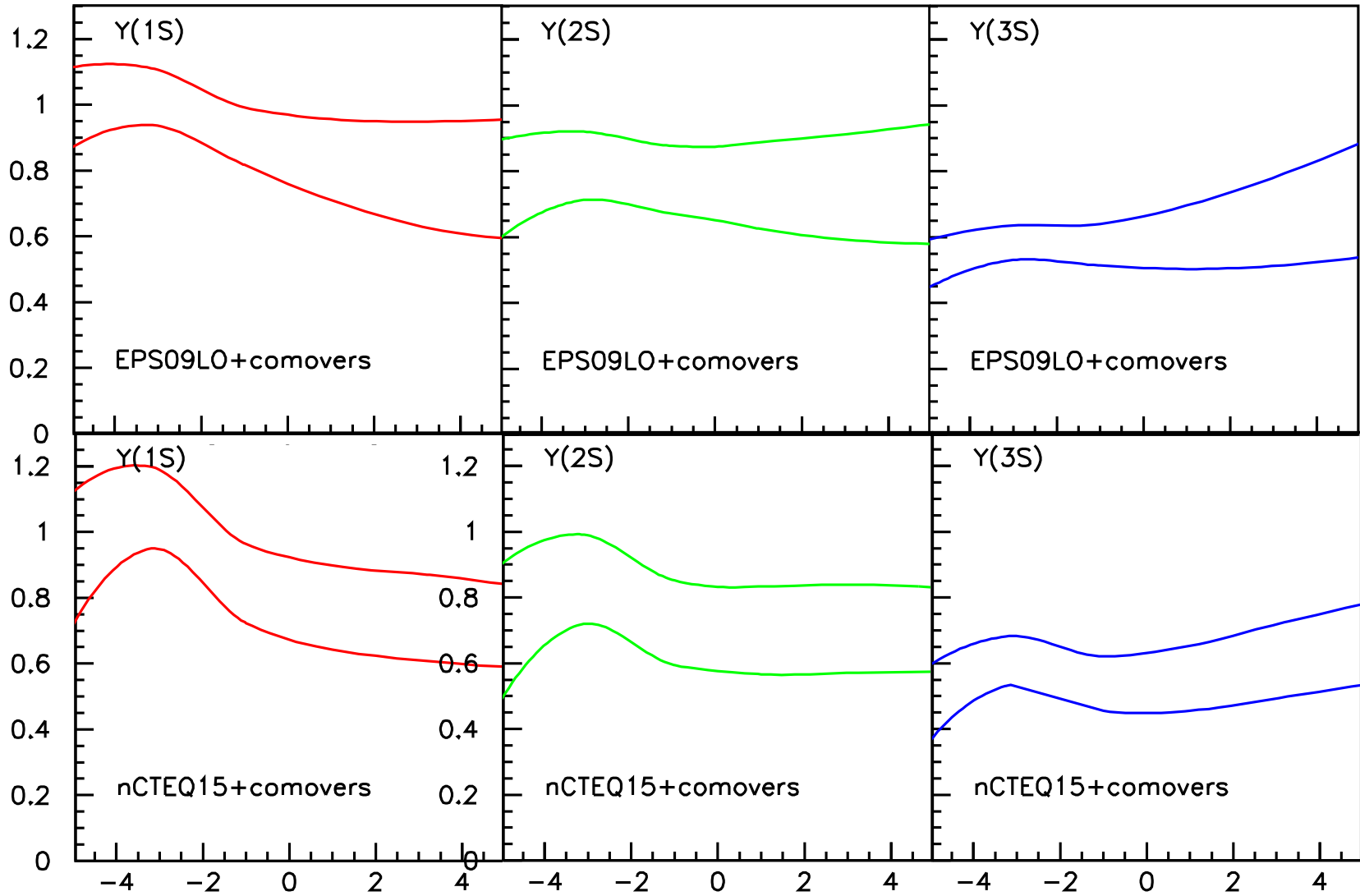
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Predictions: R_{pPb} for $Y(1S)$, $Y(2S)$ and $Y(3S)$ @ 8.16 TeV



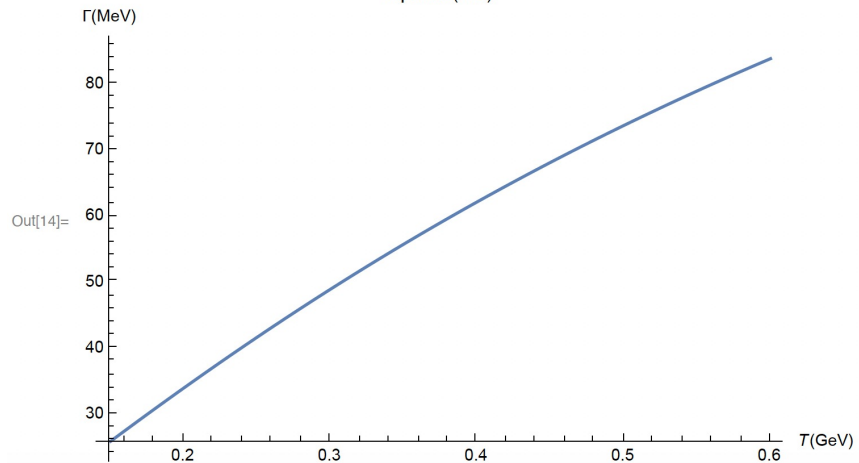
Conclusions

- We have updated our understanding of the **feed-down pattern** within the bottomium family close to $\langle p_T \rangle$ where it matters for heavy-ion studies
- In the absence of any other explanation for the **relative suppression of excited quarkonia in pA collisions** (and its rapidity dependence), we have assumed that the **reinteraction with comovers** explains it all
- This allowed us to **fit all the comover-bottomonium-interaction cross sections** from the **CMS and ATLAS pPb double ratios** in a coherent way
- A fit to the **CMS PbPb double ratios** at 2.76 and 5.02 TeV gives similar results
- We have checked that it yields a consistent magnitude for the **Y suppression** as measured by ATLAS, ALICE and LHCb **in pPb collisions** when combined with nCTEQ15 **shadowing** (which does not affect the double ratio)
- Both the **double ratios** of $Y(2S)/Y(1S)$ & $Y(3S)/Y(1S)$ (insensitive to shadowing) and the **magnitude of the suppression** (with nCTEQ15) of $Y(1S)$ & $Y(2S)$ are **well reproduced in PbPb collisions** without the need to invoke any other phenomena

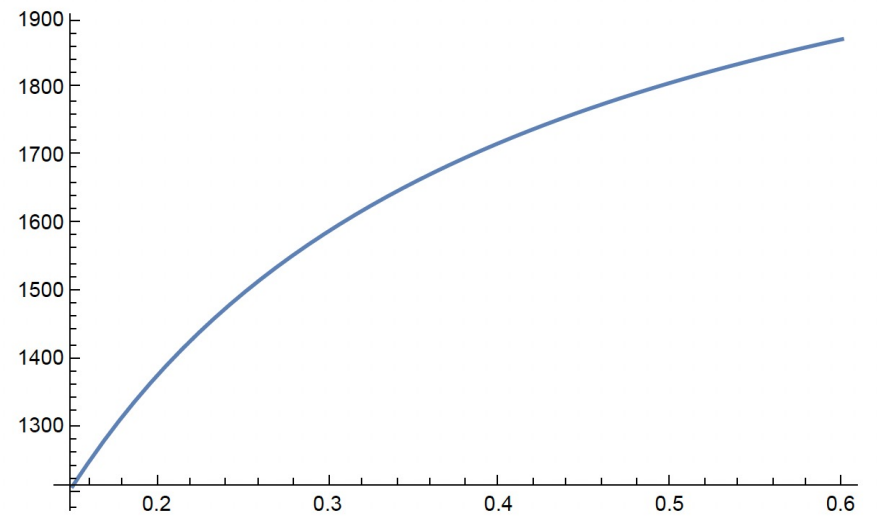
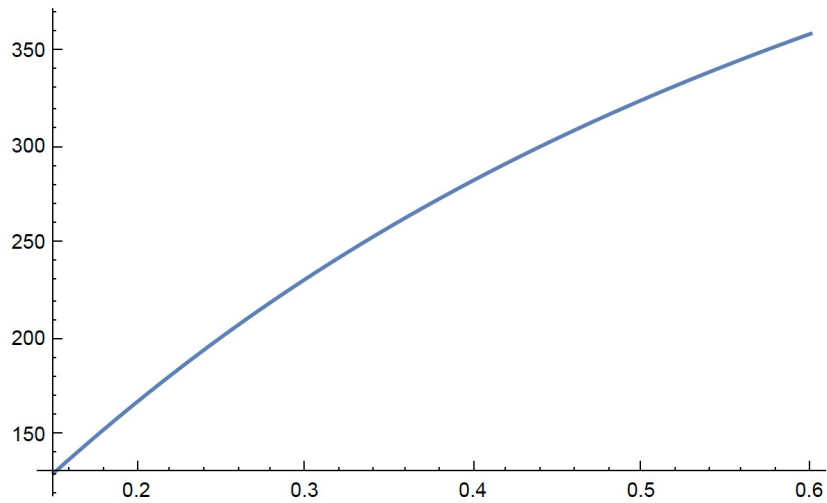
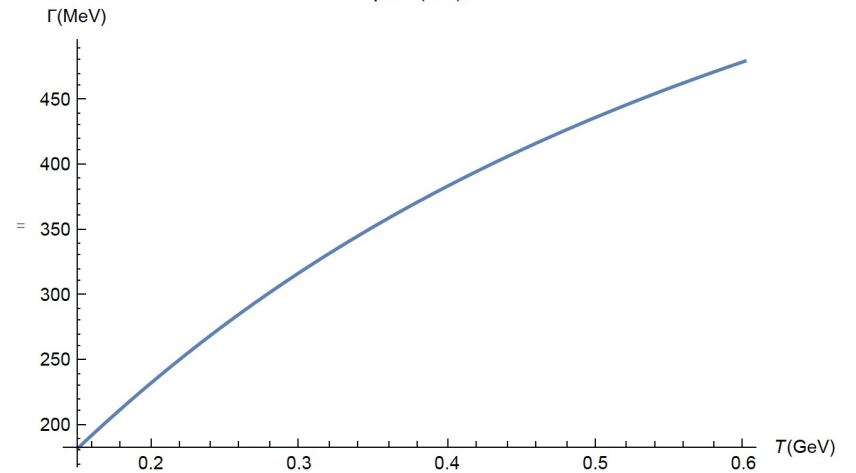
Physical interpretation: what the nature of the comovers is

- **Case I:** The medium is **hadronic in pPb** collisions, while it is **gluonic in PbPb**
 - The most common expectation: The relevant d.o.f. are hadrons in pPb collisions where the QGP is not produced whereas the gluons become relevant in the hotter PbPb environment with the presence of QGP
- **Case II:** Both in **pPb and PbPb** collisions, the medium is made of **hadrons**, i.e. the comovers can be identified with pions
 - Both in pA and AA collisions, Υ not affected by the hot (deconfined) medium
 - Possible interpretation: melting temperature of the $\Upsilon(1S)$ and $\Upsilon(2S)$ is too high to be observed and the $\Upsilon(3S)$ is fragile enough to be entirely broken by hadrons. Bottomonia unaffected by the presence of a possible QGP
- **Case III:** Both in **pPb and PbPb** collisions, the medium is made of **partons**, i.e. the comovers can be identified with gluons
 - Comovers are to be considered as partons in a (deconfined) medium
 - A QGP-like medium is formed following pPb collisions at LHC energies
 - CIM: **effective modelling** of bottomonium dissociation in the **QGP**

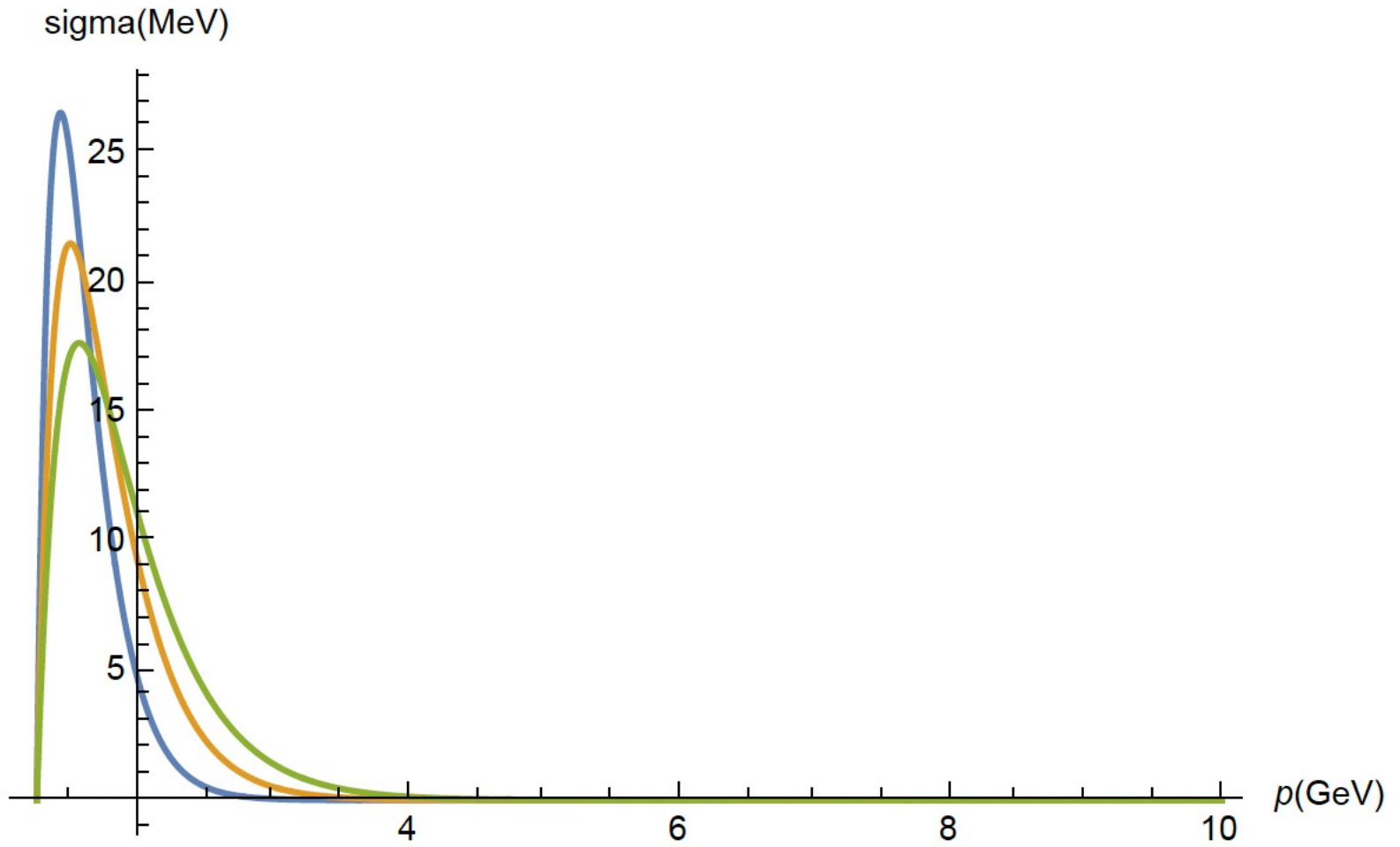
Upsilon(1 S)



Upsilon(2 S)



Upsilon(1S)



Upsilon(1S)

