

# Review Saclay approach

Miguel Ángel Escobedo

Instituto Galego de Física de Altas Enerxías  
Universidade de Santiago de Compostela

December 12, 2022



Work done in collaboration with Jean-Paul Blaizot

# Review

- Discussed in Blaizot and Escobedo, 2021.

# Review

- Discussed in Blaizot and Escobedo, 2021.
- Valid for the case  $E \gg \Gamma$ .

# Review

- Discussed in Blaizot and Escobedo, 2021.
- Valid for the case  $E \gg \Gamma$ .
- We consider the dissipative part of the interaction as a perturbation.

# Review

- Discussed in Blaizot and Escobedo, 2021.
- Valid for the case  $E \gg \Gamma$ .
- We consider the dissipative part of the interaction as a perturbation.
- We take into account the effects of the energy gap between singlets and octets.

# Review

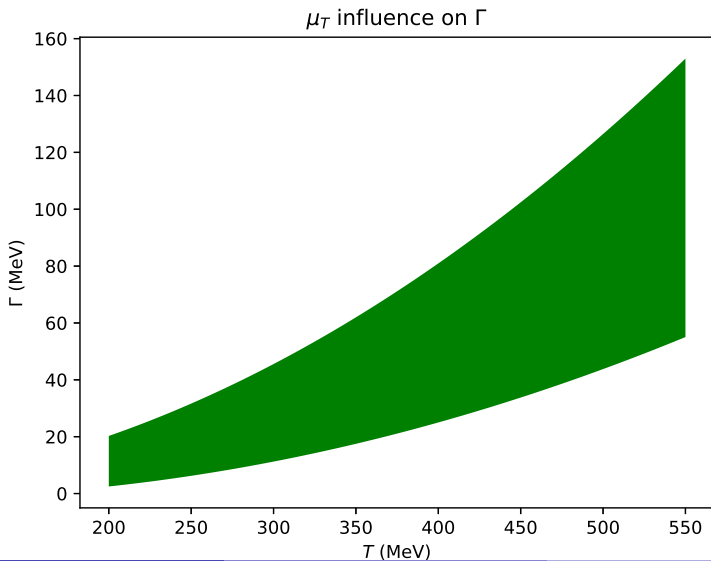
- Discussed in Blaizot and Escobedo, 2021.
- Valid for the case  $E \gg \Gamma$ .
- We consider the dissipative part of the interaction as a perturbation.
- We take into account the effects of the energy gap between singlets and octets.
- For the real part of the potential we consider two scenarios. HTL and a lattice inspired scenario.

# Review

- Discussed in Blaizot and Escobedo, 2021.
- Valid for the case  $E \gg \Gamma$ .
- We consider the dissipative part of the interaction as a perturbation.
- We take into account the effects of the energy gap between singlets and octets.
- For the real part of the potential we consider two scenarios. HTL and a lattice inspired scenario.
- The decay width comes from a HTL computation using as input the binding energy and wave function of the singlet.

# Decay vs T

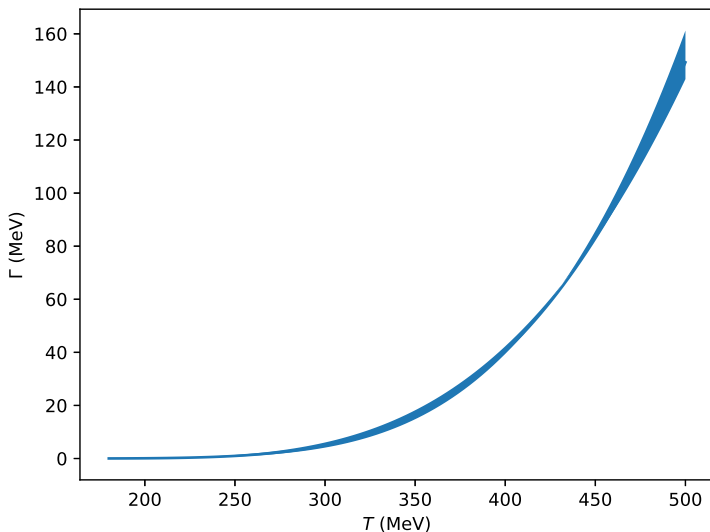
Perturbative case





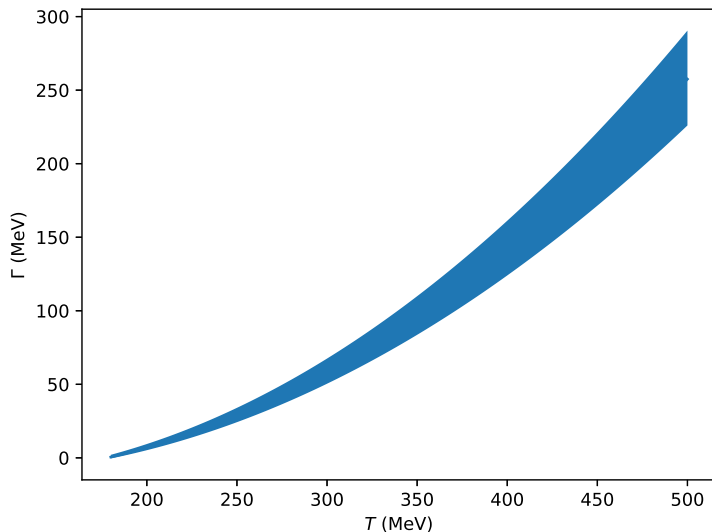
# Decay vs T

Lattice inspired scenario.  $\Upsilon(1S)$



# Decay vs T

Lattice inspired scenario.  $\Upsilon(2S)$



# Masses and binding energies

Perturbative case

- We use the 1S mass at tree level,  $M_b = \frac{M_{\Upsilon(1S)}}{2}$ .

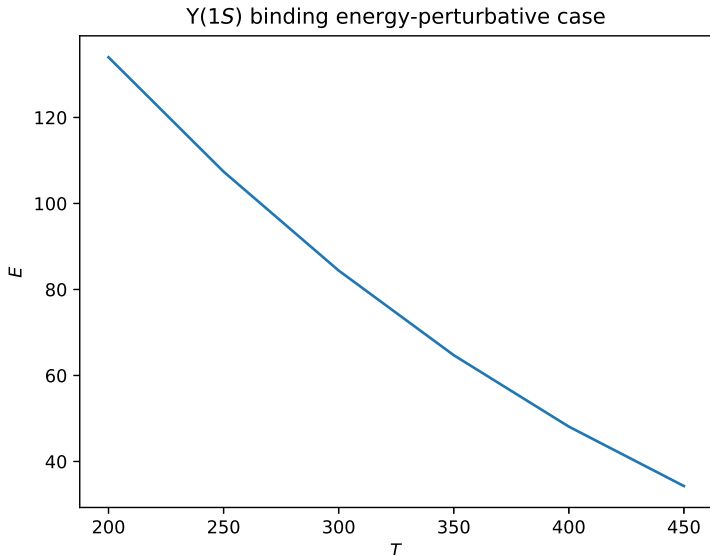
# Masses and binding energies

## Perturbative case

- We use the 1S mass at tree level,  $M_b = \frac{M_{\Upsilon(1S)}}{2}$ .
- We solve the Schrödinger equation with the real part of the HTL potential and from there we obtain the binding energy.

# Masses and binding energies

Perturbative case



# Masses and binding energies

## Lattice inspired scenario

- We use  $M_b = 4882 \text{ MeV}$ .

# Masses and binding energies

## Lattice inspired scenario

- We use  $M_b = 4882 \text{ MeV}$ .
- This is the mass used in the paper from which we get the static potential data (Rothkopf and Lafferty, 2020) to reproduce bottomonium spectroscopy.

# Masses and binding energies

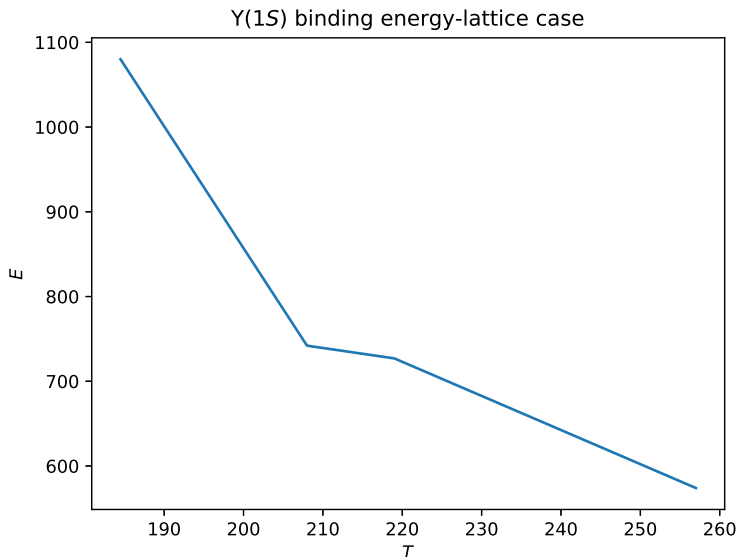
## Lattice inspired scenario

- We use  $M_b = 4882 \text{ MeV}$ .
- This is the mass used in the paper from which we get the static potential data (Rothkopf and Lafferty, 2020) to reproduce bottomonium spectroscopy.
- As real potential, we use a parametrization that was shown in (Rothkopf and Lafferty, 2018) to reproduce the static potential within errors.



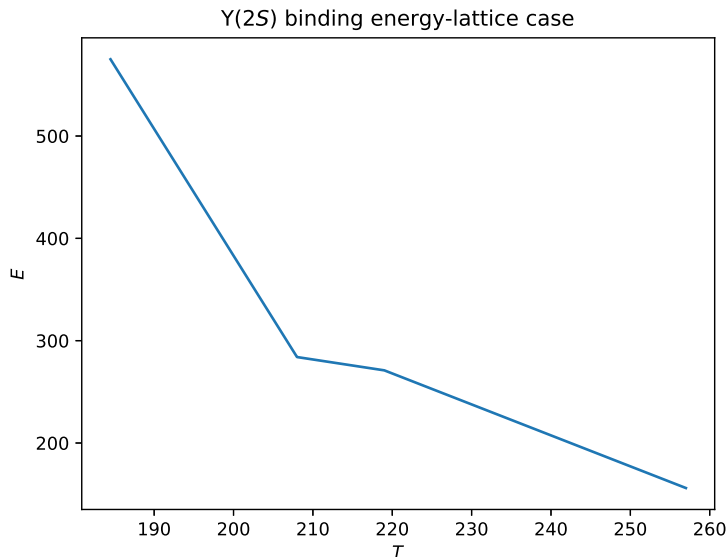
# Masses and binding energies

Lattice inspired scenario.  $\Upsilon(1S)$



# Masses and binding energies

Lattice inspired scenario.  $\Upsilon(2S)$



## Gamma vs $p$

We do not include  $p$  dependence in our approach. Hence,  $\Gamma$  depends only on temperature.

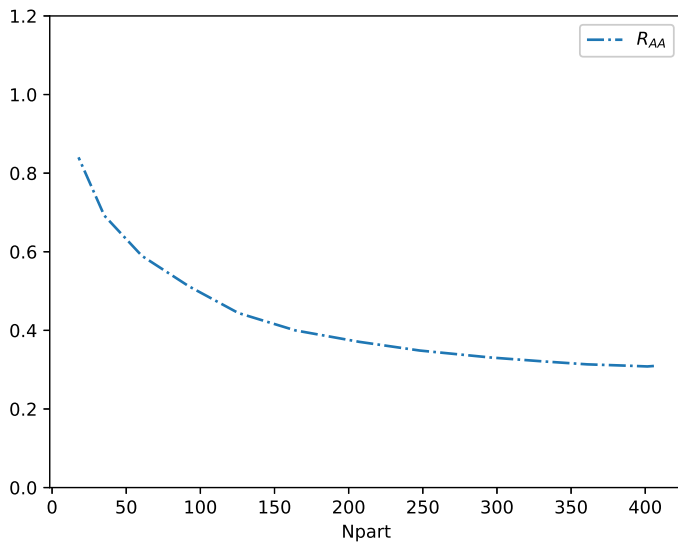
| T   | $\Upsilon(1S)$ perturbative | $\Upsilon(1S)$ lattice | $\Upsilon(2S)$ lattice |
|-----|-----------------------------|------------------------|------------------------|
| 200 | 6.52                        | 0.0889                 | 7                      |
| 300 | 19.54                       | 4.96                   | 58.5                   |
| 400 | 40.45                       | 40.4                   | 142                    |

## $R_{AA}$ for fixed $\Gamma$

$$\Gamma = \begin{cases} 0 & T < 200 \text{ MeV} \\ \frac{T}{2} - 100 \text{ MeV} & T > 200 \text{ MeV} \end{cases}$$

- Bjorken evolution.
- Glauber model.
- Initial temperature scales with the number of participants.

# $R_{AA}$ for fixed $\Gamma$



## $R_{AA}$ vs $p$

We do not include  $p$  dependence in our model. However, we can compute  $R_{AA}$  in the given centrality window

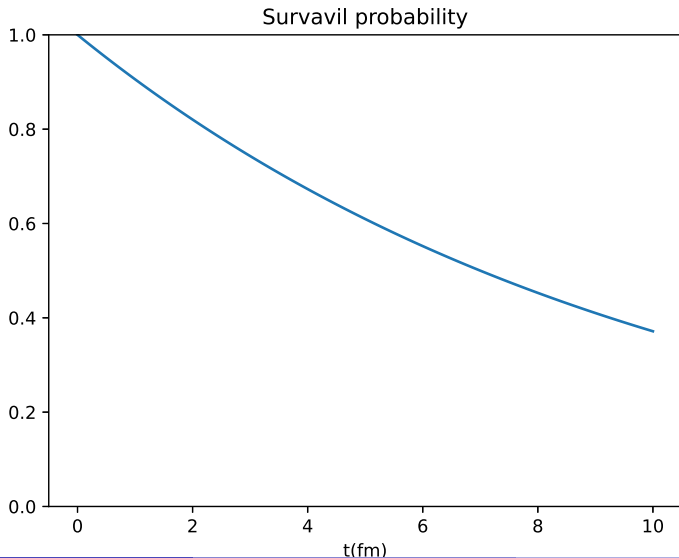
$$R_{AA}|_{0-10 \text{ centrality}} = 0.32$$

# Survival probability

- $T = 300 \text{ MeV}$
- Initial state is a medium  $\Upsilon(1S)$  state.

# Survival probability

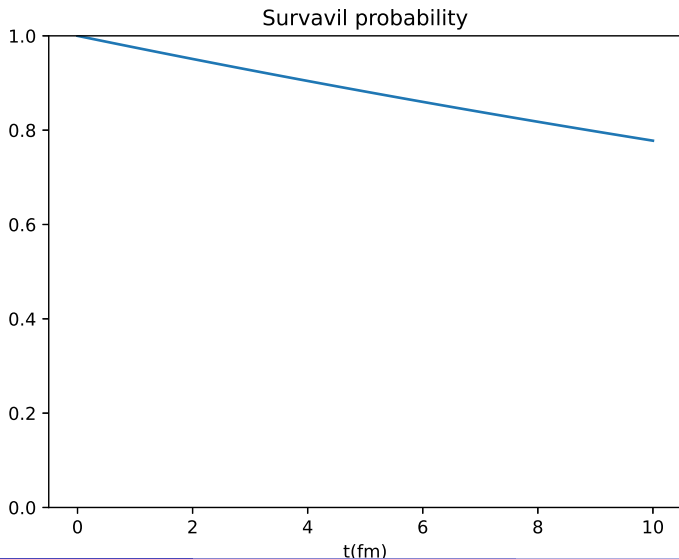
Perturbative case





# Survival probability

Lattice inspired scenario



# Vacuum state transition

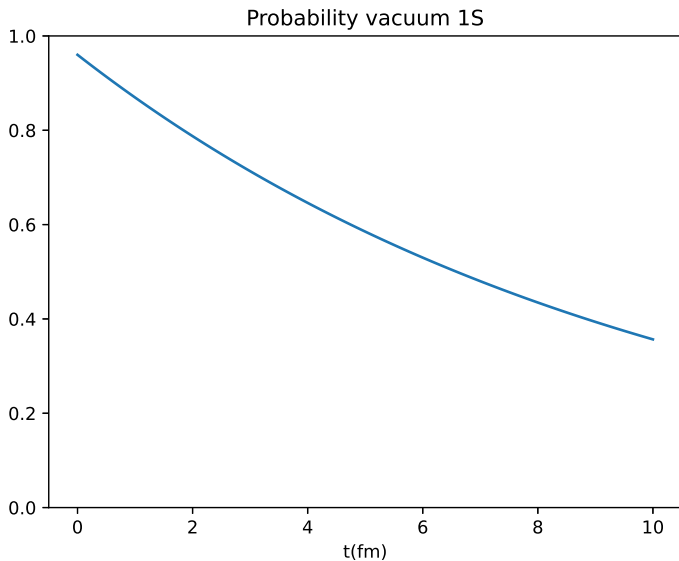
## Perturbative case

We computed numerically the overlap between Yukawa potential eigenvectors and Coulomb ones.

## Lattice inspired scenario

- We could not directly compare with the  $T = 0$  eigenvectors in an easy way. Our code is set to work with potentials that go zero at infinity, which is not the case of the Cornell potential.
- Then we compare with a very small temperature in which medium effects are very mild.
- We see that medium and vacuum eigenvectors are almost identical at the given temperature.

# Vacuum state transition



# Vacuum state transition

