The Nantes approach

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The model

- ► Work by Blaizot & Escobedo: SU(3): J-P. Blaizot, M. Escobedo (2018)
 - Heavy quarks-Plasma interaction described using NRQCD
 - Derivation of quantum master equations in the quantum brownian motion regime (high temperature) to describe the evolution of the density operator (not Lindblad equations)

Our work:

- Extension to preserve positivity ⇒ Lindblad equations
- Direct resolution in 1D and application to charmonium system
- Study of the validity of a semi-classical treatment (not covered in this talk)
- New potential developped specifically for 1D studies R. Katz, S.D., P-B. Gossiaux (2022)

Quantum Master Equation



$$\mathcal{L} = \mathcal{L}_0 + \mathcal{L}_1 + \mathcal{L}_2 + \mathcal{L}_3 + \mathcal{L}_4$$

 \mathcal{L}_0 : Kinetic terms

expected to be subleading \mathcal{L}_1 : Static screening (V)

Higher-order terms,

 \mathcal{L}_2 : Fluctuations (W) Dynamical processes $\mathcal{L}_3/\mathcal{L}_4$: Dissipation (W'/W"/W"')

> Transition between color states and dissipation effects

Quantum Master Equation

$$\mathcal{L}_{0}\mathcal{D} = -i[H_{Q}, \mathcal{D}] \qquad \qquad \blacktriangleright n$$

$$\mathcal{L}_{1}\mathcal{D} = -\frac{i}{2} \int_{xx'} V(x - x') [n_{x}^{a} n_{x'}^{a}, \mathcal{D}] \qquad \qquad \square n$$

$$\mathcal{L}_{2}\mathcal{D} = \frac{1}{2} \int_{xx'} W(x - x') \left(\{n_{x}^{a} n_{x'}^{a}, \mathcal{D}\} - 2n_{x}^{a} \mathcal{D} n_{x'}^{a} \right)$$

$$\mathcal{L}_{3}\mathcal{D} = -\frac{i}{4T} \int_{xx'} W(x - x') \left(\dot{n}_{x}^{a} \mathcal{D} n_{x'}^{a} - n_{x}^{a} \mathcal{D} \dot{n}_{x'}^{a} + \frac{1}{2} \{\mathcal{D}, [\dot{n}_{x}^{a}, n_{x'}^{a}]\}$$

•
$$n_x^a$$
: color charge density
 $n_x^a = \delta(x - r) t^a \otimes \mathbb{I} - \mathbb{I} \otimes \delta(x - r) \tilde{t}^a$

▶ Can recover \mathcal{L}_3 from \mathcal{L}_2 by performing:

$$\left(\left\{n_{x}^{a}n_{x'}^{a},\mathcal{D}\right\}-2n_{x}^{a}\mathcal{D}n_{x'}^{a}\right) \longrightarrow \left\{\left(n_{x}^{a}-\frac{i}{4T}\dot{n}_{x}^{a}\right)\left(n_{x'}^{a}+\frac{i}{4T}\dot{n}_{x'}^{a}\right),\mathcal{D}\right\}-2\left(n_{x}^{a}+\frac{i}{4T}\dot{n}_{x}^{a}\right)\mathcal{D}\left(n_{x'}^{a}-\frac{i}{4T}\dot{n}_{x'}^{a}\right)$$

• Additionnal terms
$$\Rightarrow \mathcal{L}_4$$

1D Potential



- Based on a 3D potential inspired from Lattice results D. Lafferty, A. Rothkopf (2020)
- Real part: parametrization to reproduce 3D mass spectra
- Imaginary part: separated in a coulombic and string part, aims at reproducing 3D decay widths

1D Potential



- Very good agreement for the mass spectra
- Good agreement for the decay widths, differences due to the large distance behaviour of the imaginary part

Reaction rates



- Increase with temperature and momentum
- Stronger increase for J/Ψ

Charmonium gaussian singlet initial state T = 300 MeV



- Initial gaussian singlet state at T = 300 MeV (σ = 0.1 fm)
- Octet populated as a dipole
- Delocalization of initial state along s = s' axis
- Remaining central correlation

Charmonium gaussian singlet initial state T = 300 MeV



- Instantaneous projections on vacuum eigenstates defined as P_Φ(t) =< Φ|D_s(t)|Φ >
- Equilibration phase with transitions between states
 - *\chi_c* populated later due to different transitions
- Decay phase afterwards, with same decay rate for all states

Asymptotic Wigner distribution



- Large distance
- Distribution progressively becomes Gaussian

Asymptotic Wigner distribution



- $\sqrt{< p^2 >}$ does not scale as $\sqrt{\frac{MT}{2}}$ (dotted lines)
- ▶ Equilibrium limit modified by \mathcal{L}_4
- ► At large distances, scaling as $\sqrt{\frac{1}{1+\frac{\gamma}{2}}\frac{MT}{2}}$ with $\gamma = \frac{\tilde{W}^{(4)}(0)}{16MT\tilde{W}''(0)}$ (dashed lines)

Charmonium gaussian octet initial state in a cooling medium



 Cooling medium with gaussian octet initial state (σ = 0.1 fm)

$$T(t) = T_0 \left(\frac{1}{1+t} \right)^{1/3}$$

- Delocalization of initial state along s = s' axis
- Seems to reach same kind of limit

Charmonium gaussian octet initial state in a cooling medium



- Formation of bound states at early times
- Helped by the initial proximity of the two quarks
- Global evolution similar to the fixed temperature case