

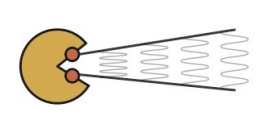
*Via Piella, Bo*



# Accessing the three-body dynamics with multi-particle correlations at LHC

V. Mantovani Sarti (TUM)

4th Workshop on Anti-Matter, Hyper-Matter and Exotica Production at the LHC  
15.02.2023



# TUM Many-body dynamics in nuclear physics

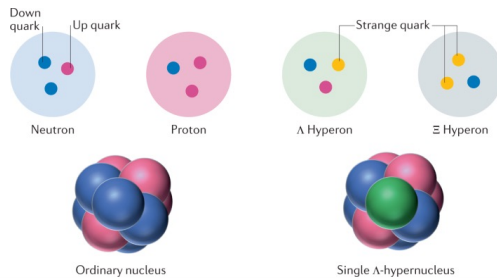
1. Properties of nuclei and hypernuclei cannot be described satisfactorily with two-body forces only.

*L.E. Marcucci et al., Front. Phys. 8:69 (2020)*

*Hypernuclei session Tue. 14.02*

2. Many-body scattering requires three-body calculations (e.g. neutron-deuteron).

*L. Girlanda et al., PRC 102, 064003 (2020)*



$\rho_0$

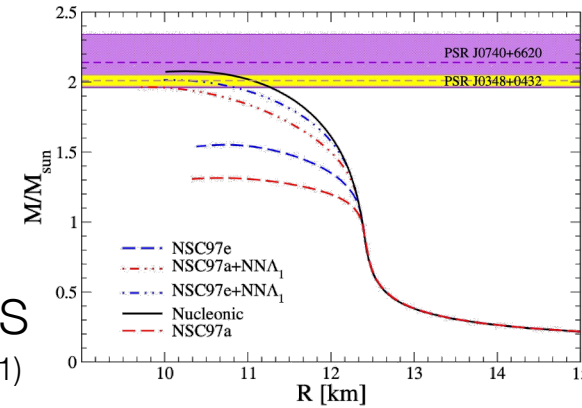
$\rho$  (fm<sup>-3</sup>)

# TUM Many-body dynamics in astrophysics

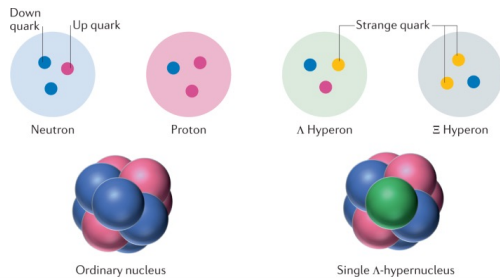
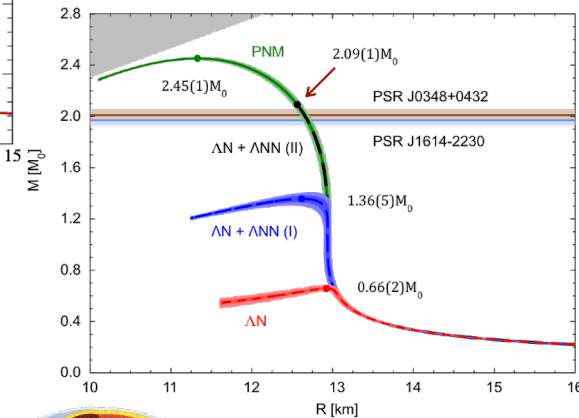
## 3. Equation of state of matter inside neutron stars

- Production of hyperons energetically favourable around  $2-3 \rho_0$
- Only two-body  $\Lambda N$ 
  - Too soft EoS, incompatible with measured heavy NS
  - Large improvement in 2-body  $YN$  with femtoscopy<sup>(1)</sup>
- Introduction of three-body  $\Lambda NN$  forces
  - Stiffens EoS, model-dependent<sup>(2)</sup>
  - Need for additional experimental constraints

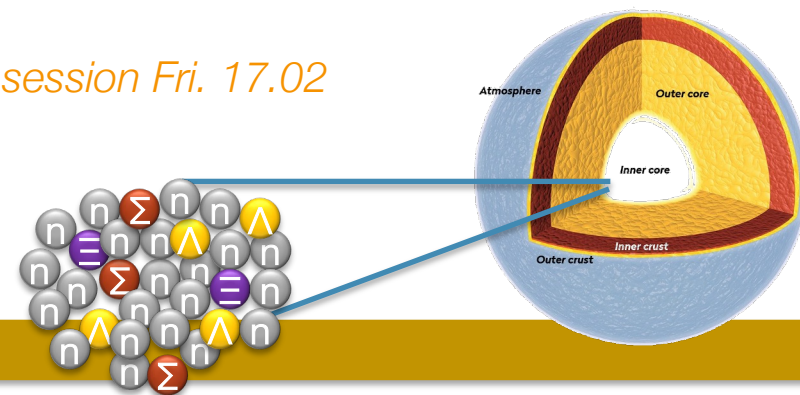
D. Logoteta et al., Eur.Phys.J.A 55 (2019)



D. Lonardoni et al., PRL 114 (2019)



NS session Fri. 17.02



$\rho_0$

(1) ALICE Coll. PLB 833 (2022)  
ALICE Coll. Nature 588 (2020)  
L. Fabbietti et al. Ann.Rev.Nucl.Part.Sci. 71 (2021)

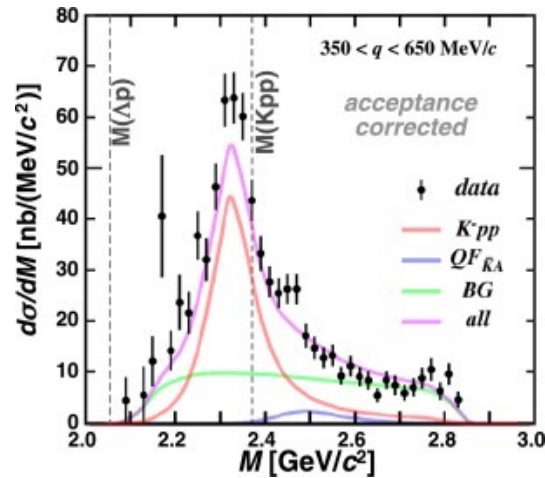
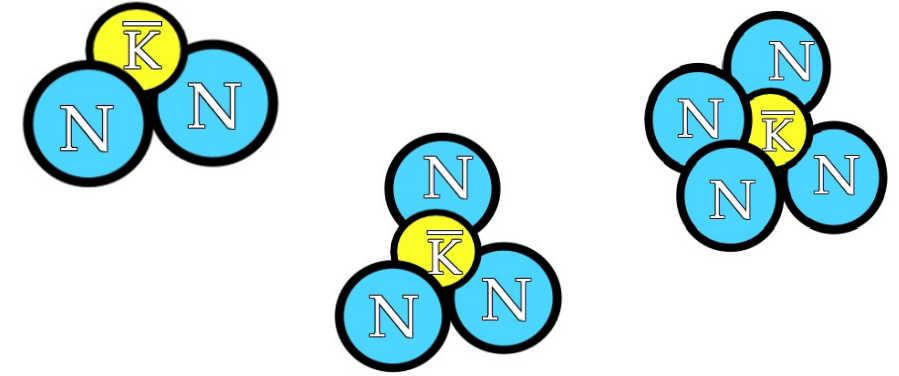
$\sim 3 \rho_0$

(2) D. Logoteta et al., Eur.Phys.J.A 55 (2019)  
D. Gerstung et al., Eur.Phys.J.A 56 (2020)  
D. Lonardoni et al., PRL 114 (2019)

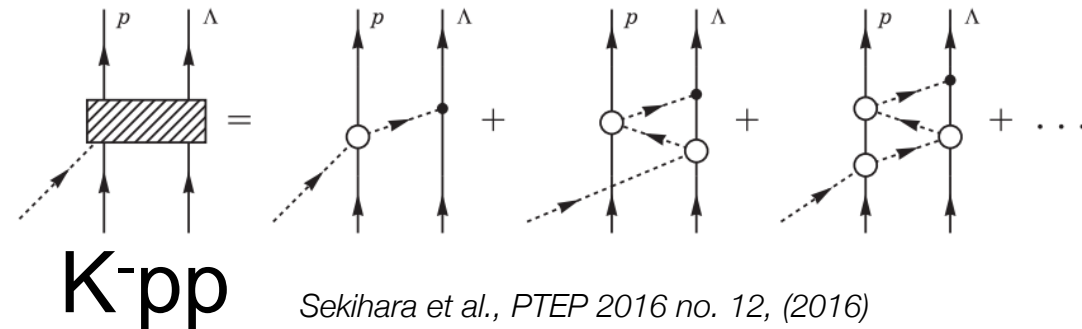
$\rho$  ( $\text{fm}^{-3}$ )

# TUM Many-body dynamics with mesons and bound states

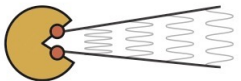
- Exotic bound states of antikaons with nucleons due to the strongly attractive KN interaction in  $l = 0$  channel  
*S. Wycech, NPA 450 (1986); Y. Akaishi, T. Yamazaki, PRC 65 (2002)*
- First experimental evidence of the p-p-K- bound state by the E15 Collaboration. *E15 Coll., PLB 789 (2019) 620*



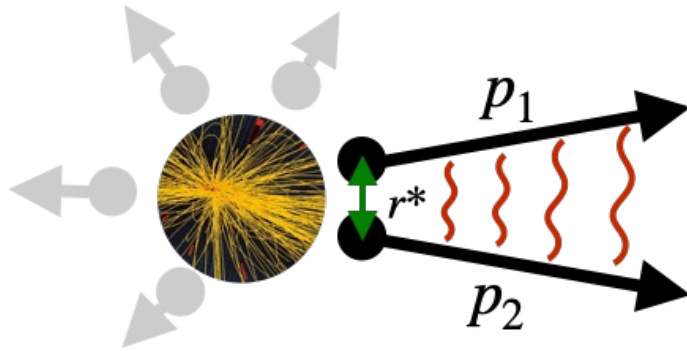
## Kaonic bound state formation mechanism



Next experimental challenge: genuine three-body interaction measurements



# TUM Two-body femtoscopy at LHC



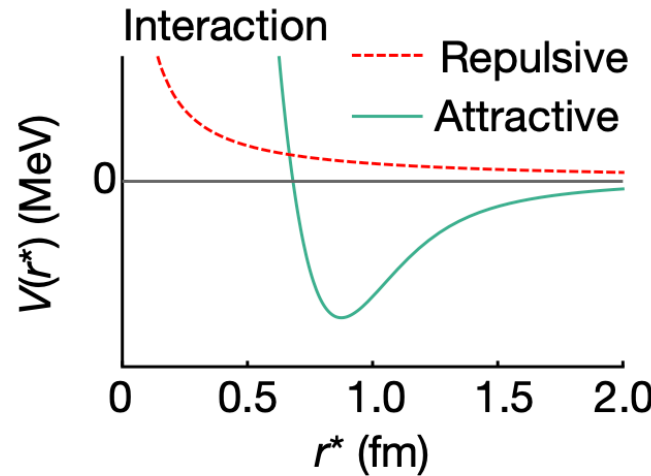
Emission source  $S(r^*)$   
ALICE Coll. PLB 811 (2020) 135849

ALICE Collaboration  
 $p\Lambda$ : PLB 833 (2022), 137272  
 $p\Sigma^0$ : PLB 805 (2020) 135419  
 $p\Xi/p\Omega$ : Nature 588 (2020) 232-238  
 $p\phi$ : PRL 127 (2021), 172301  
 $Kp$ : PRL 124 (2020) 09230,  
 PLB 822 (2021) 136708  
 arXiv: 2205.15176 [nucl-ex]  
 and many more:  
 PRC 99 (2019) 2, 024001  
 PLB 797 (2019) 134822  
 PRL 123 (2019) 112002  
 PRC 103 (2021) 5, 055201  
 PLB 829 (2022), 137060  
 arXiv: 2201.05352 [nucl-ex]  
 arXiv: 2204.10258 [nucl-ex]

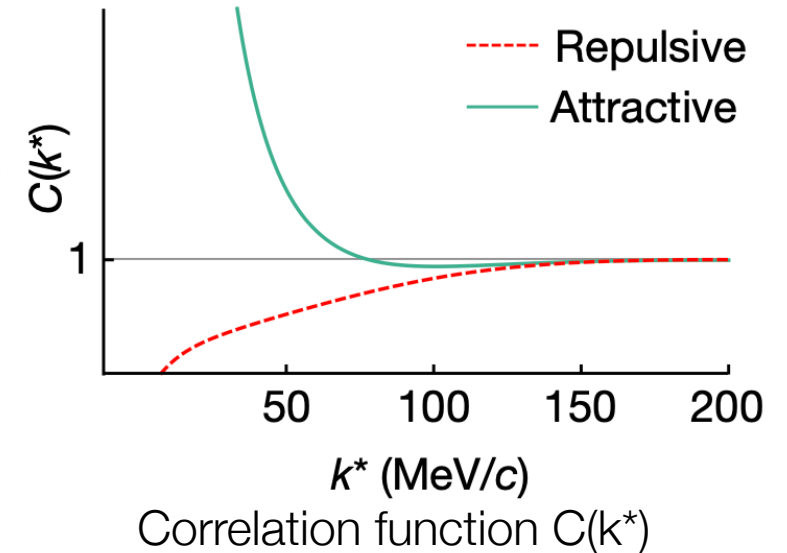
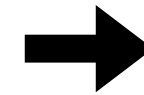
- Two-particle correlation function Today session

$$C(k^*) = N \frac{N_{\text{same}}(k^*)}{N_{\text{mixed}}(k^*)} = \int S(r^*) |\psi(k^*, r^*)|^2 d^3 r^*$$

- Measurements in small colliding systems ( $\sim 1-2$  fm)  
 → Access to the strong interaction and short-range dynamics

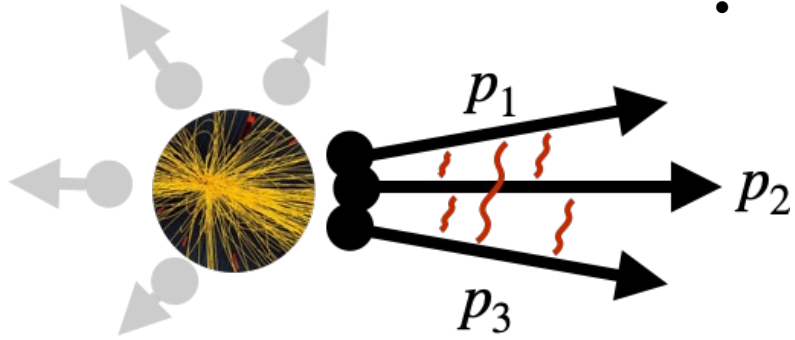


Schrödinger equation  
 Two-particle wave function  
 $|\psi(k^*, r^*)|$



Correlation function  $C(k^*)$

# TUM Three-body femtoscopy at LHC



- Three-particle correlation function

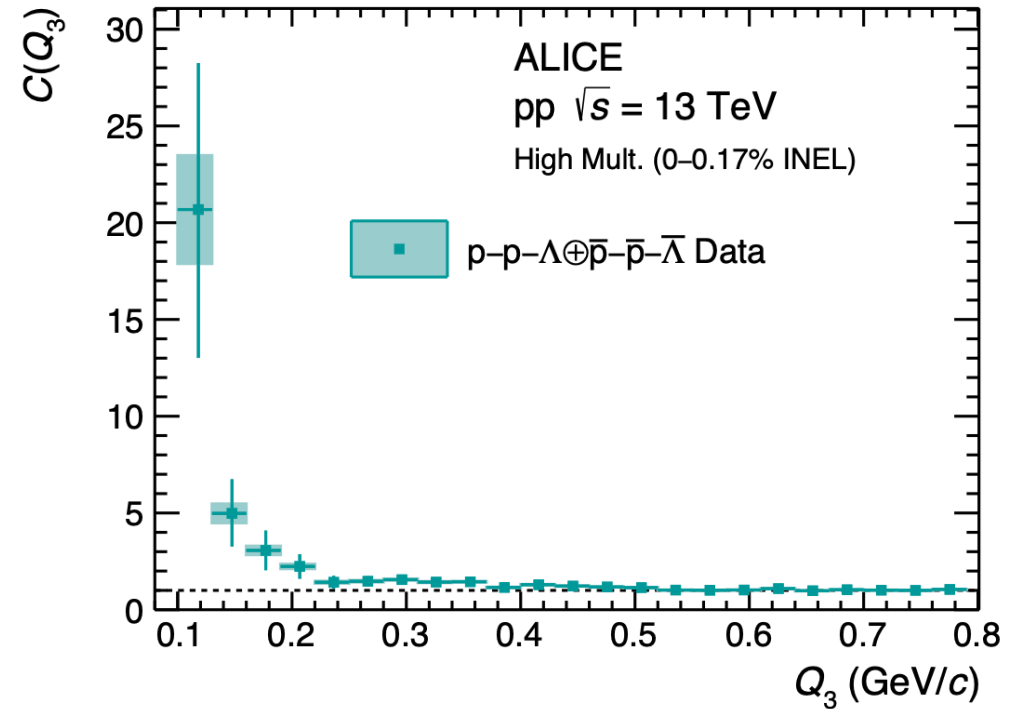
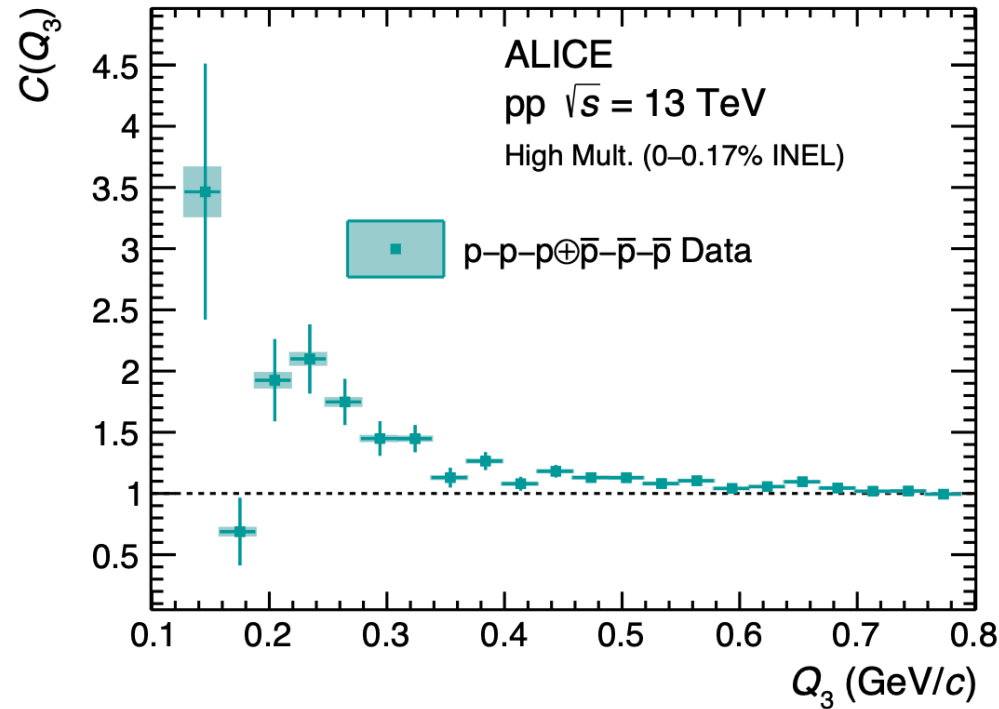
$$C(\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3) \equiv \frac{P(\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3)}{P(\mathbf{p}_1)P(\mathbf{p}_2)P(\mathbf{p}_3)} = \mathcal{N} \frac{N_{\text{same}}(Q_3)}{N_{\text{mixed}}(Q_3)}$$

$$Q_3 = \sqrt{-q_{ij}^2 - q_{jk}^2 - q_{ki}^2}$$

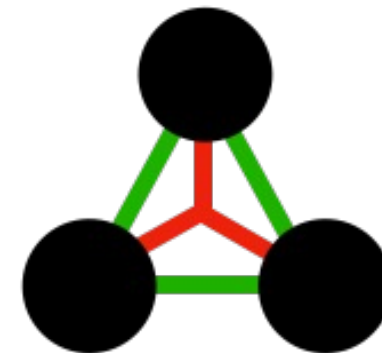
- Applied multi-pion correlations to study coherent emission  
*ALICE Coll. Phys.Rev.C 89 (2014), Phys.Rev.C 93 (2016)*

- First measurement of p-p-p and p-p- $\Lambda$  correlations in high-multiplicity pp collisions 13 TeV by ALICE  
*ALICE Coll. arXiv: 2206.03344 nucl-ex*

# Accessing the ppp and pp $\Lambda$ interactions at LHC

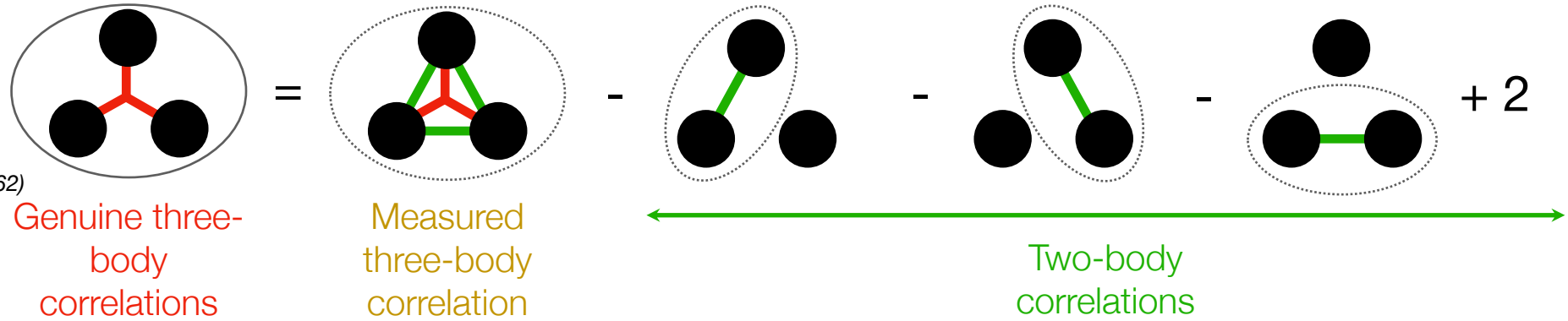


- Signal different from unity at low  $Q_3$
- How can we interpret these three-body correlations?
  - Two-body interactions
  - Three-body interactions



# TUM Cumulants in femtoscopy

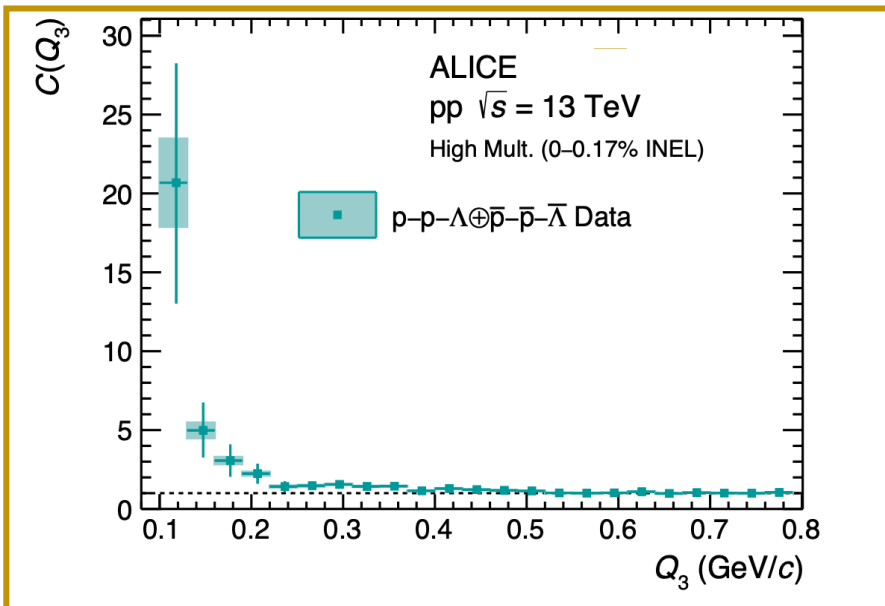
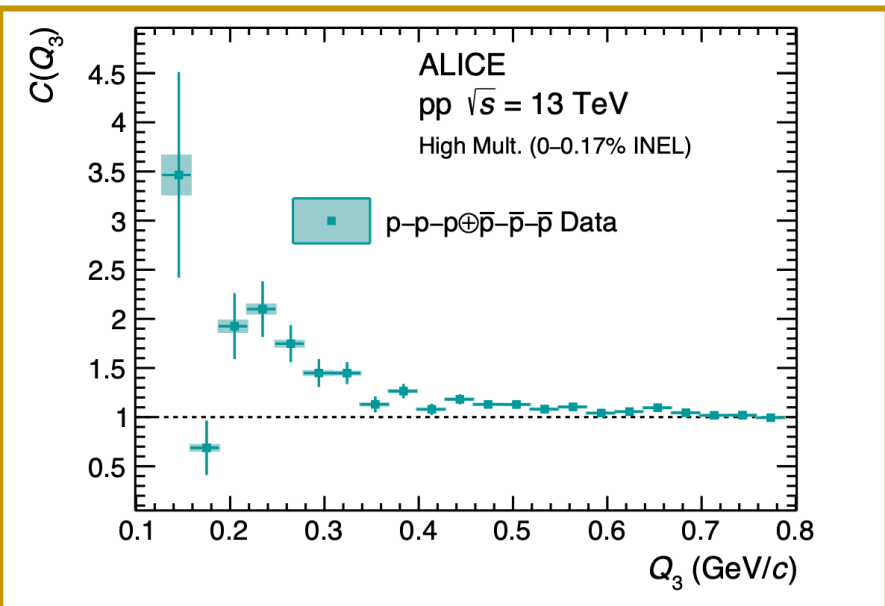
Kubo's rule<sup>(1)</sup>:



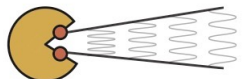
(1) R. Kubo, J. Phys. Soc. Jpn. 177 (1962)

Correlations:

$$c_3(Q_3) = C(Q_3) - C_{12}(Q_3) - C_{23}(Q_3) - C_{31}(Q_3) + 2$$

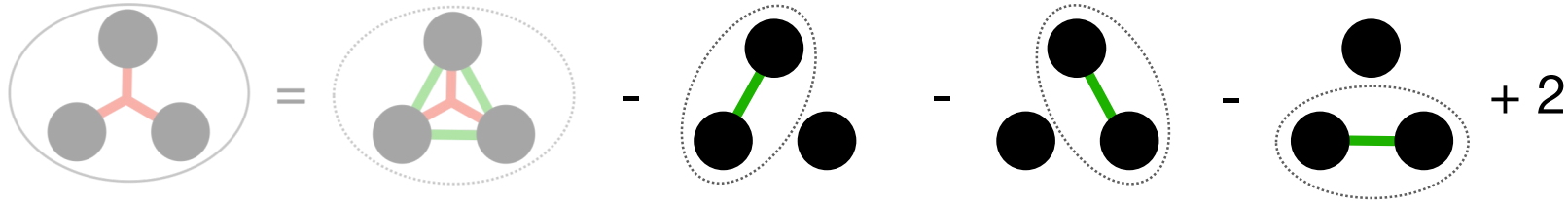


Adapted from ALICE Coll. arXiv: 2206.03344 nucl-ex

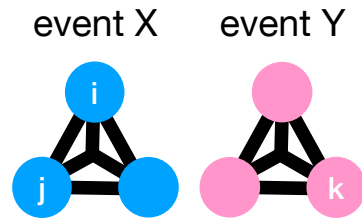




# Lower order contributions: two different approaches



- Using directly the data
- Two particles from the same event and one particle from another:



$$C_{ij}([\mathbf{p}_i, \mathbf{p}_j], \mathbf{p}_k) = \frac{N_2(\mathbf{p}_i, \mathbf{p}_j) N_1(\mathbf{p}_k)}{N_1(\mathbf{p}_i) N_1(\mathbf{p}_j) N_1(\mathbf{p}_k)}$$

ALICE Coll. arXiv: 2206.03344 nucl-ex

- Projector method
- Evaluate two-body contribution in three-body correlation function
  - Two-particle measured or theoretical correlation function  $C([\mathbf{p}_i, \mathbf{p}_j])$
- Kinematic transformation in phase space

$$C_{ij}(Q_3) = \int \underbrace{C(k_{ij}^*)}_{\text{two-body CF}} \underbrace{W_{ij}(k_{ij}^*, Q_3)}_{\text{projector}} dk_{ij}^*$$

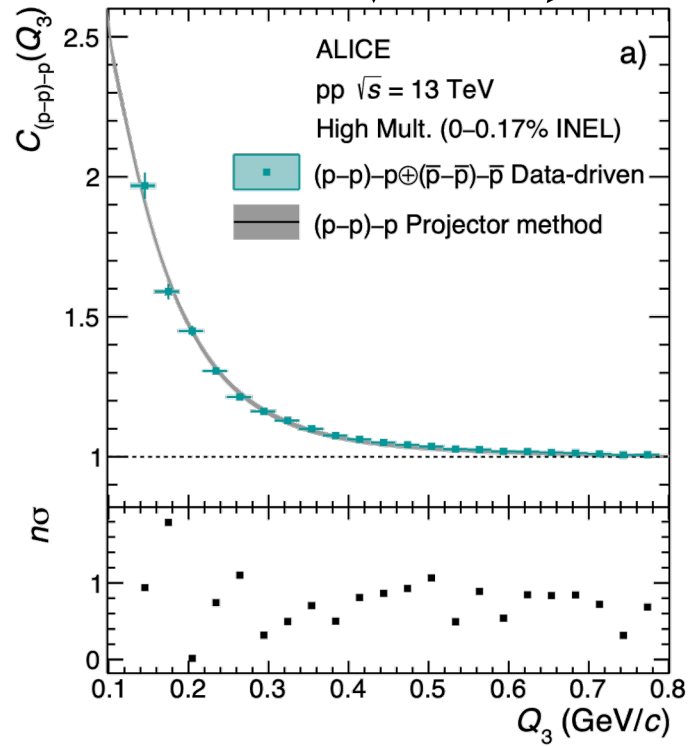
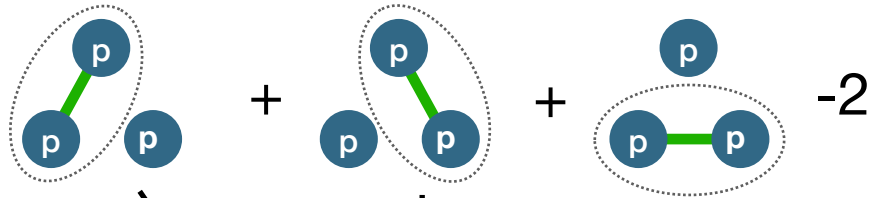
Measured/modeled 2-body correlation functions

Jacobian from 2-body to 3-body coordinates

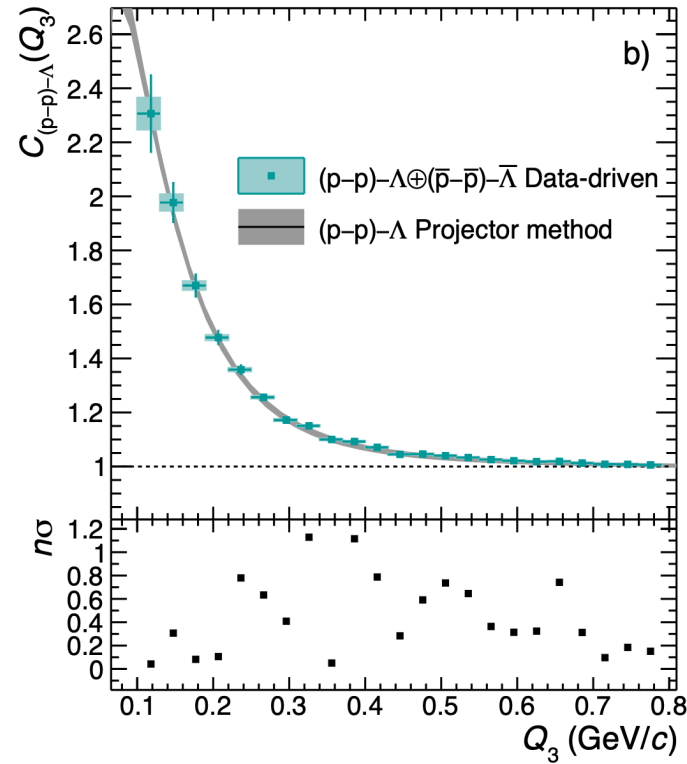
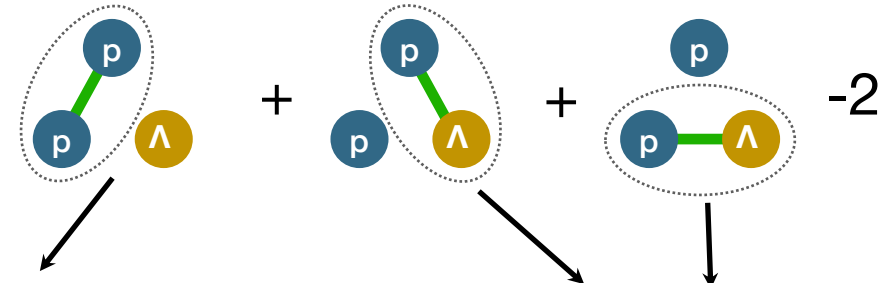
*p-p*: ALICE Coll. PLB 805 (2020) 135419  
*pA*: ALICE Coll. PLB 833 (2022), 137272

R. Del Grande, L. Serksnyte et al. EPJC 82 (2022)

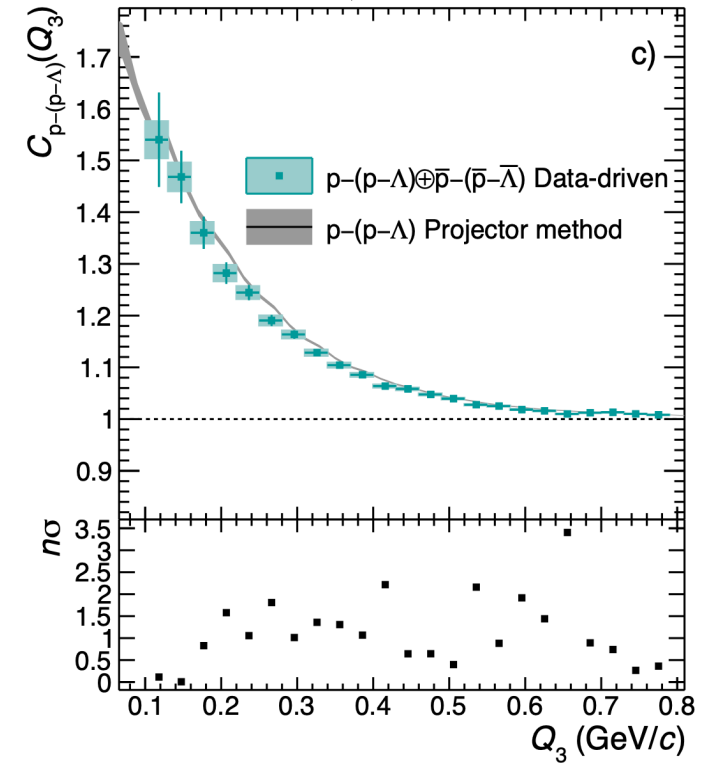
# Lower order contributions in p-p-p and p-p- $\Lambda$



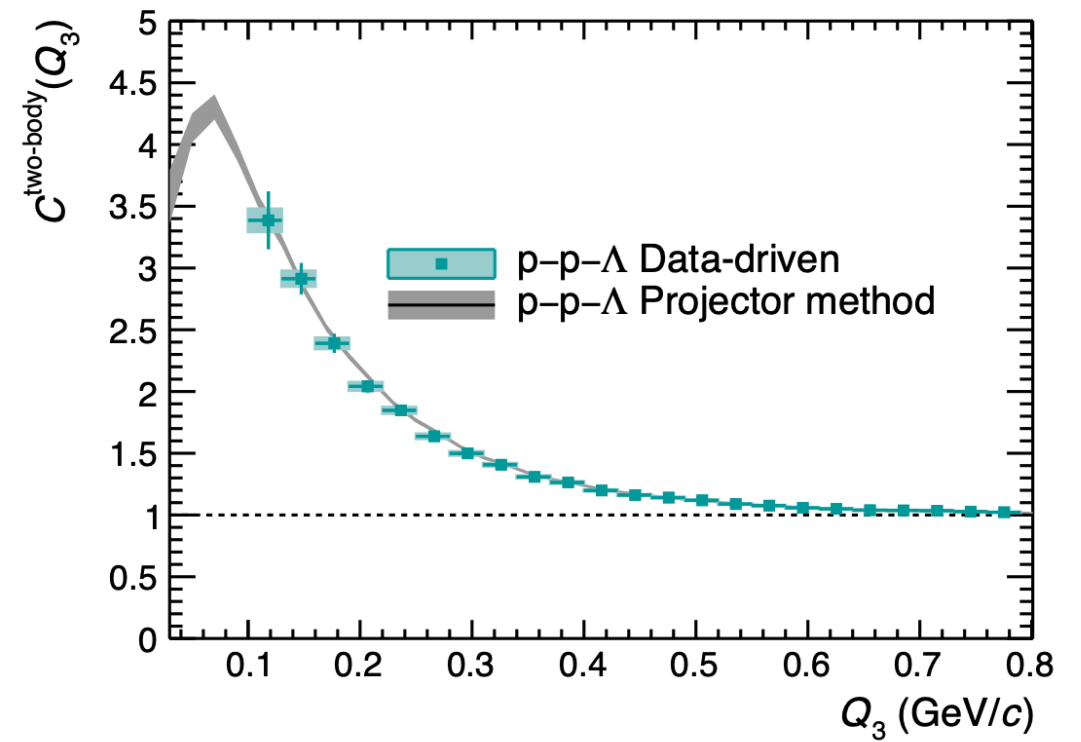
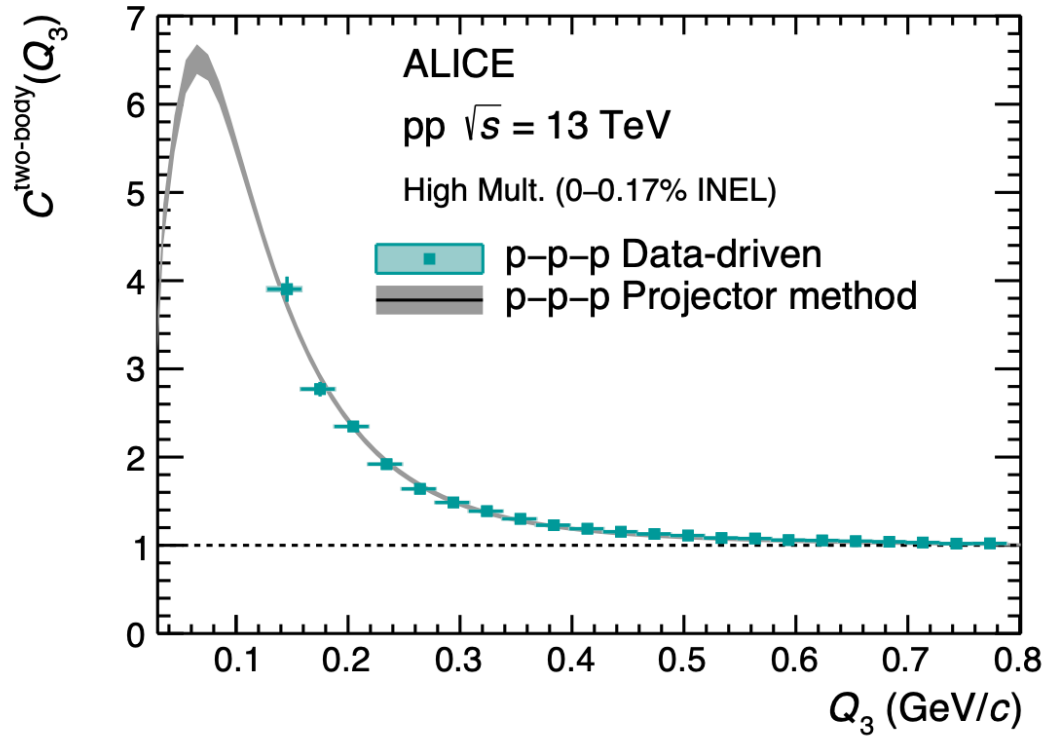
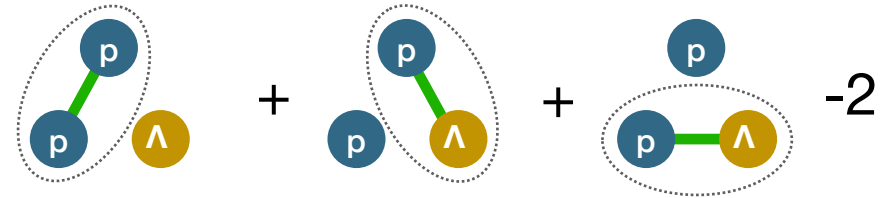
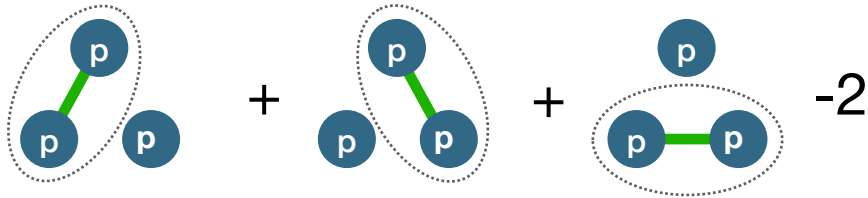
Measured p-p correlation ALICE Coll. PLB 805 (2020) 135419



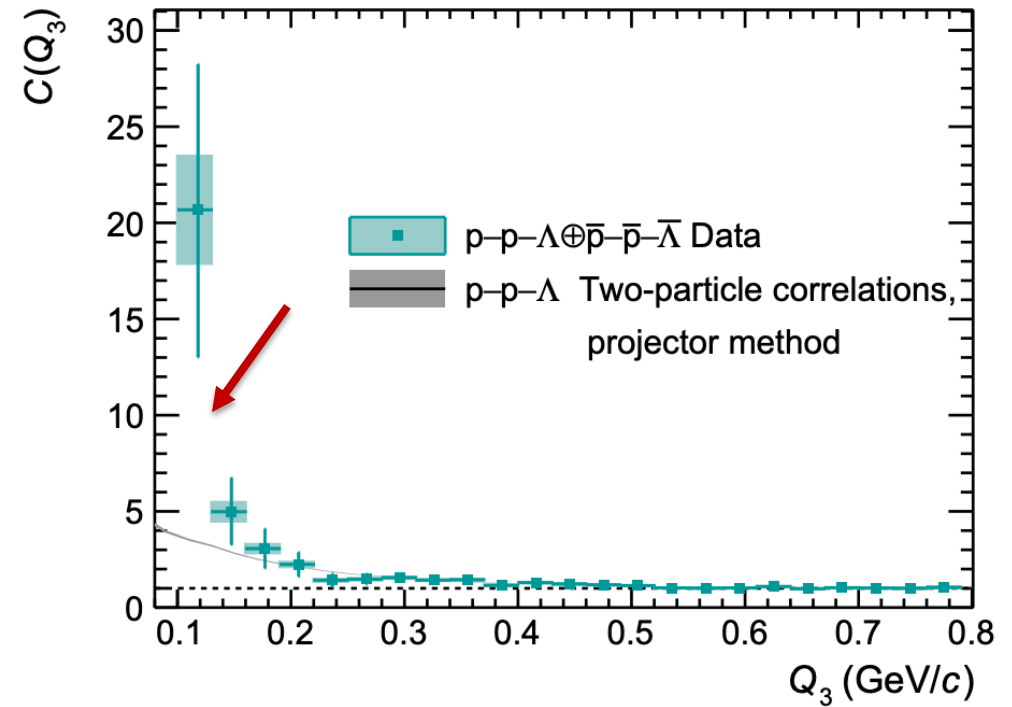
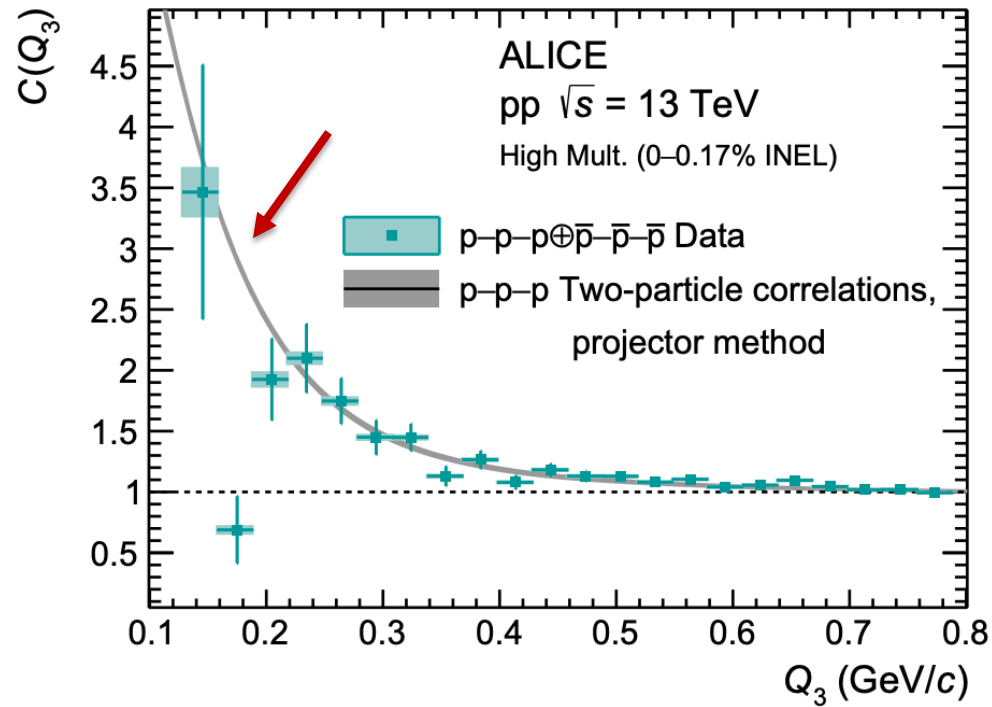
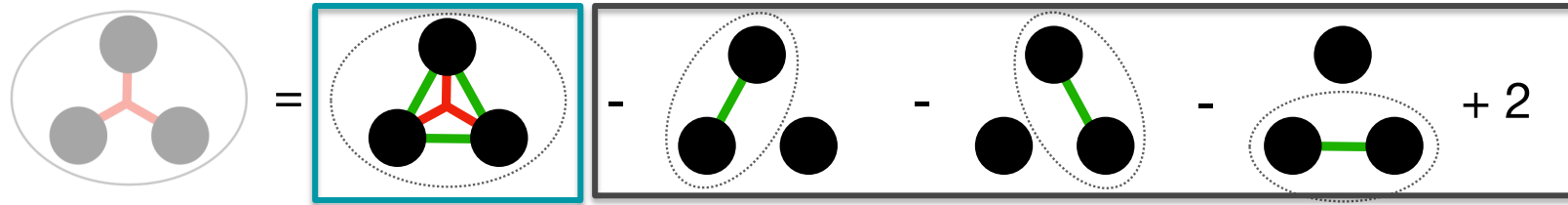
Measured p- $\Lambda$  correlation ALICE Coll. PLB 833 (2022), 137272

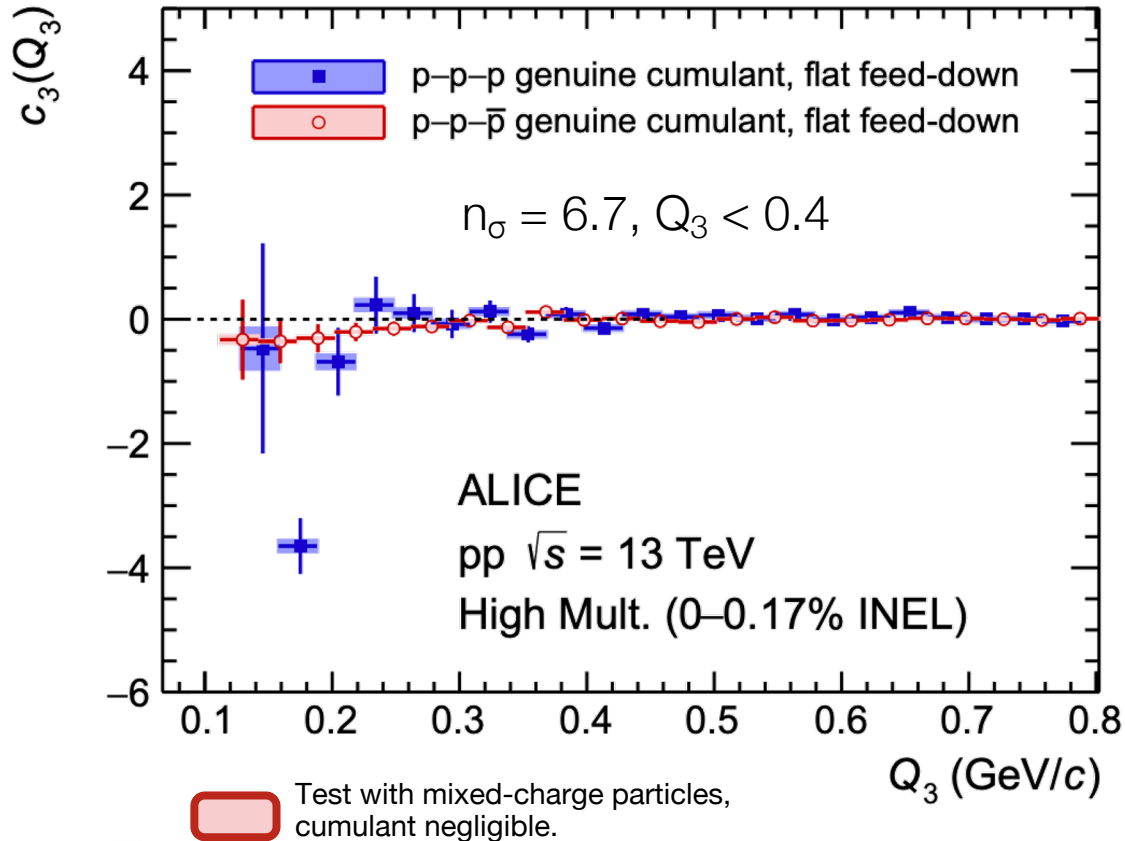
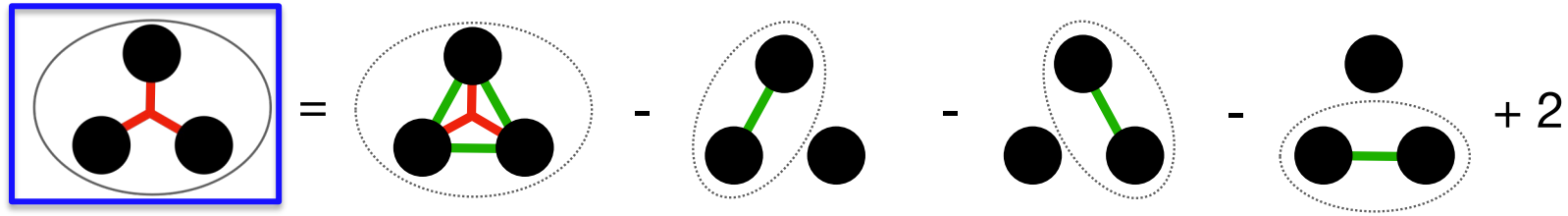


# Lower order contributions in p-p-p and p-p- $\Lambda$

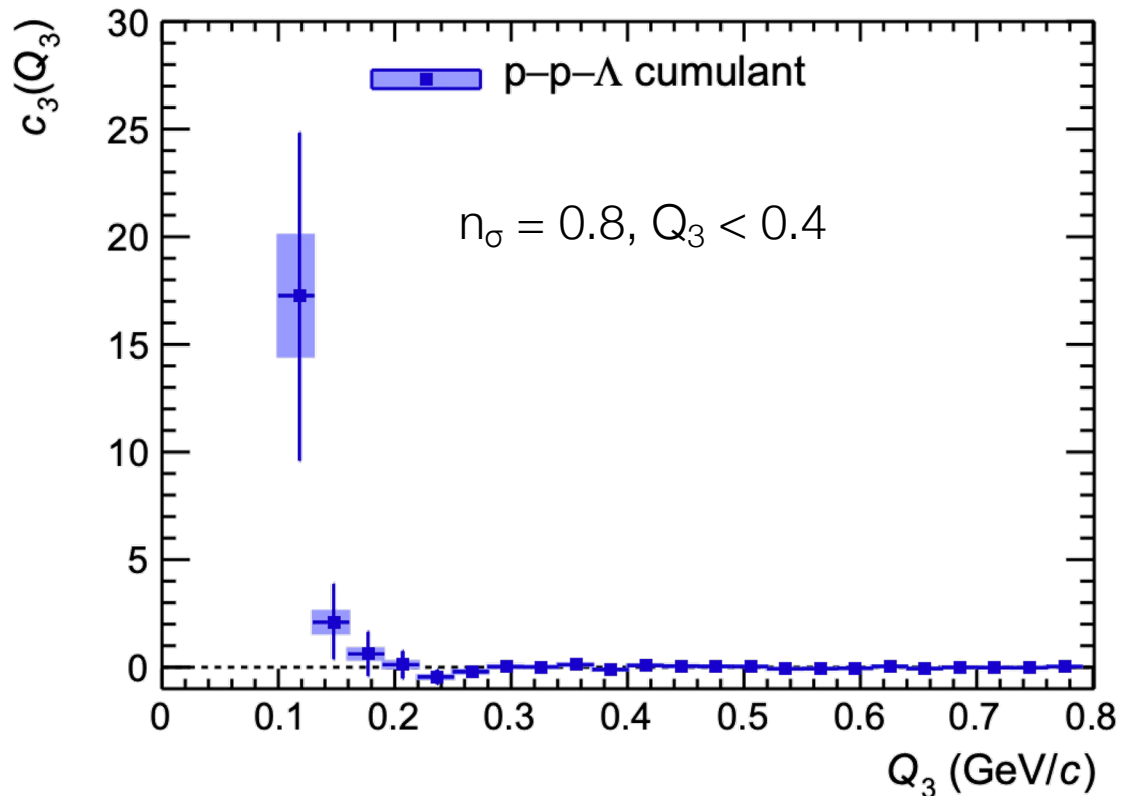
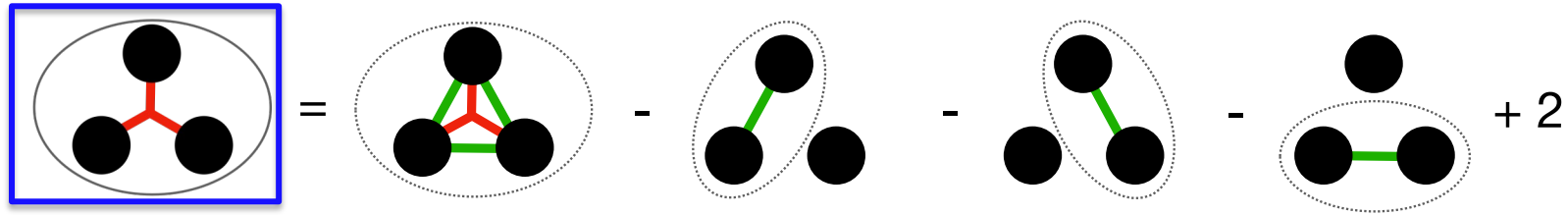


# Understanding the three-particle correlation functions



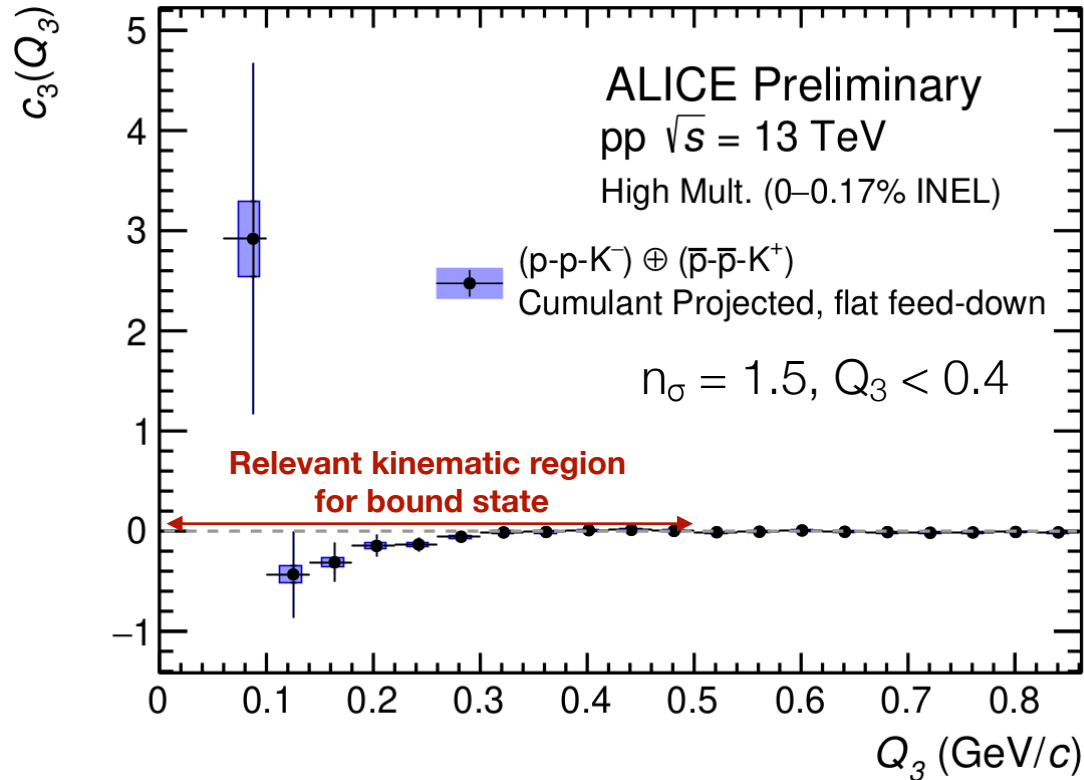
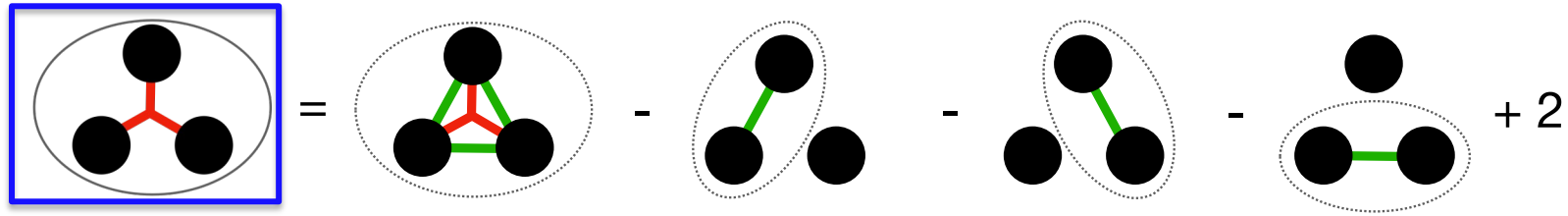


- Negative cumulant and deviation from zero  
→ Presence of a genuine 3-body effect
- Possible interpretations
  - Pauli blocking at 3-particle level
  - 3-body strong interaction
- Take-home message:
  - Significant deviation from null hypothesis
  - Ongoing collaboration with Pisa theory group (Prof. Kviesky, Prof. Maruccci and Dr. Viviani)



- Positive cumulant and deviation from zero  
→ Presence of a genuine 3-body effect
- Main contribution from three-body strong interaction  
→ Not there yet but ...
- Take-home message:  
→ No significant deviation from null hypothesis  
→ Final answer in ongoing Run 3 and future Run 4 at LHC with enhanced statistics!

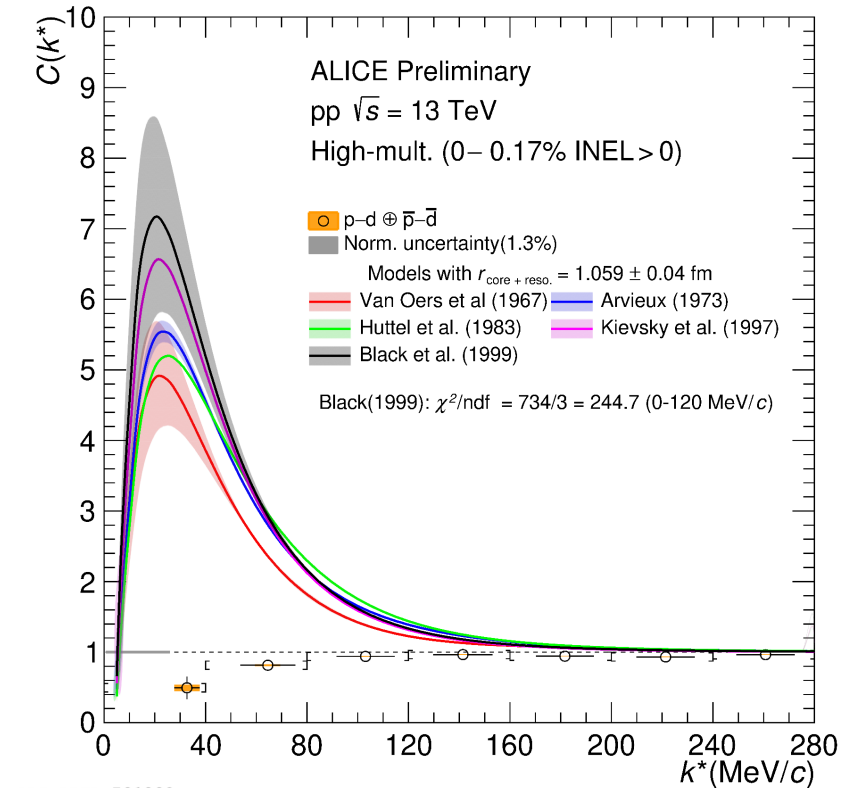
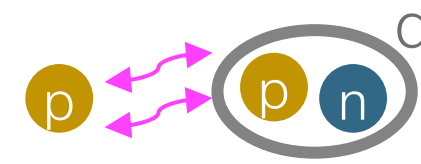
# In the meson sector: the p-p-K<sup>-</sup> cumulant



- Zero cumulant in this case  
→ No three-body effect
- Take-home message:  
→ Cumulant compatible with zero  
→ p-p-K<sup>-</sup> dynamics dominated by two-body interactions  
→ Confirming three-body strong interaction should not be relevant in the formation of exotic kaonic bound states!

# TUM Hadron-deuteron correlations

- Indirect measurements of three body forces via h-d correlations in pp collisions:
  - Presence of a composite object as the deuteron (p,n)
  - Access to spin-isospin NNN dependence
  - Large discrepancy with theoretical CF via Lednicky-Lyuboshits approach<sup>(1)</sup>
    - p and d as distinguishable point-like particles does not work!
  - Effect of three-body p-(p-n) dynamics must be included (Work in collaboration with Pisa group <sup>(2)</sup>)
  - Theoretical studies in  $\Lambda$ d correlations
    - S=1/2 state related to  ${}^3H_\Lambda$  binding energy
    - S=3/2 sensitive to spin-triplet  $\Lambda$ N
    - Dedicated ALICE measurements in Run 3 and 4



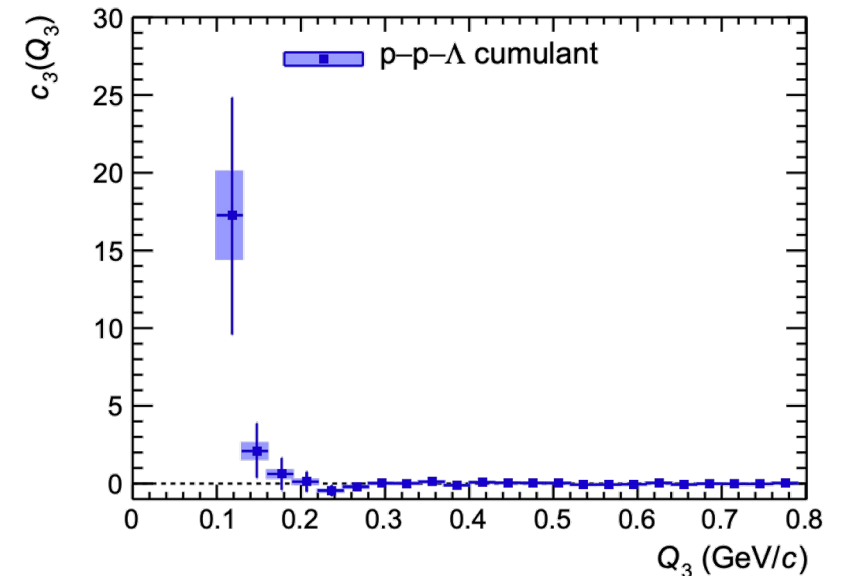
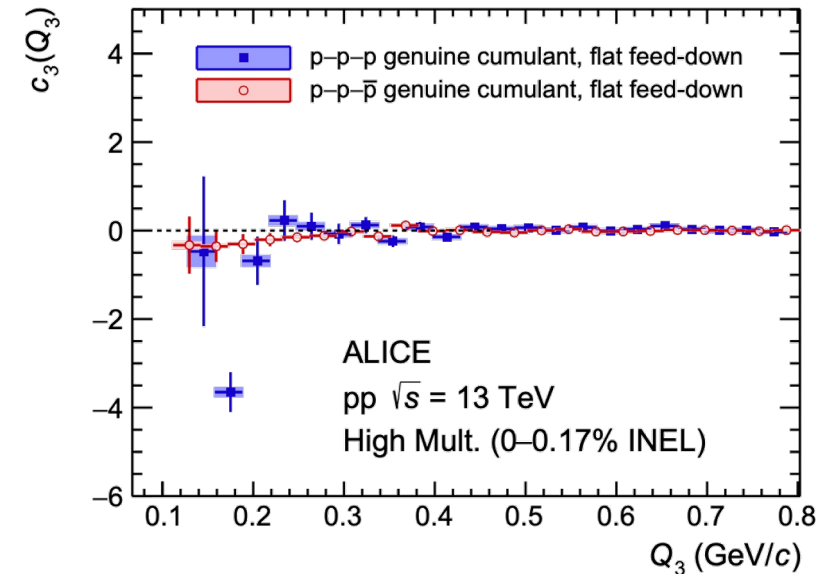
ALI-PREL-501009

- (1) R. Lednicky, *Phys. Part. Nuclei* 40, 307–352 (2009)  
 Van Oers, Brockmann et al. *Nucl. Phys. A* 561-583 (1967)  
 J. Arvieux et al. *Nucl. Phys. A* 221 253-268 (1973)  
 E. Huttel et al. *Nucl. Phys. A* 406 443-455 (1983)  
 A. Kievsky et al. *PLB* 406 292-296 (1997)  
 T.C. Black et al. *PLB* 471 103-107 (1999)
- (2) Kievsky et al, *Phys. Rev. C* 64 (2001) 024002  
 Kievsky et al, *Phys. Rev. C* 69 (2004) 014002  
 Deltuva et al, *Phys. Rev. C* 71 (2005) 064003

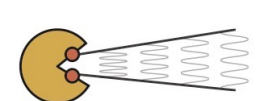


# TUM Conclusions and outlooks

1. Femtoscopy in pp collisions at the LHC applied to triplets
  - Genuine three-body cumulants isolated for the first time using the Kubo's rule
2. p-p-p and p-p- $\Lambda$  cumulants
  - p-p-p statistically significant negative cumulant, ongoing work for the theoretical interpretation and modeling
  - p-p- $\Lambda$  not significant yet, Run 3 statistics will determine the nature of the interaction
  - Important ingredient for the physics of neutron stars
3. Similar technique applied to p-p- $K^\pm$ 
  - Kaonic bound state dynamics dominated by 2-body
4. Indirect access to three-body dynamics via hadron-deuteron correlations (p-d,  $\Lambda$ -d) and extension to



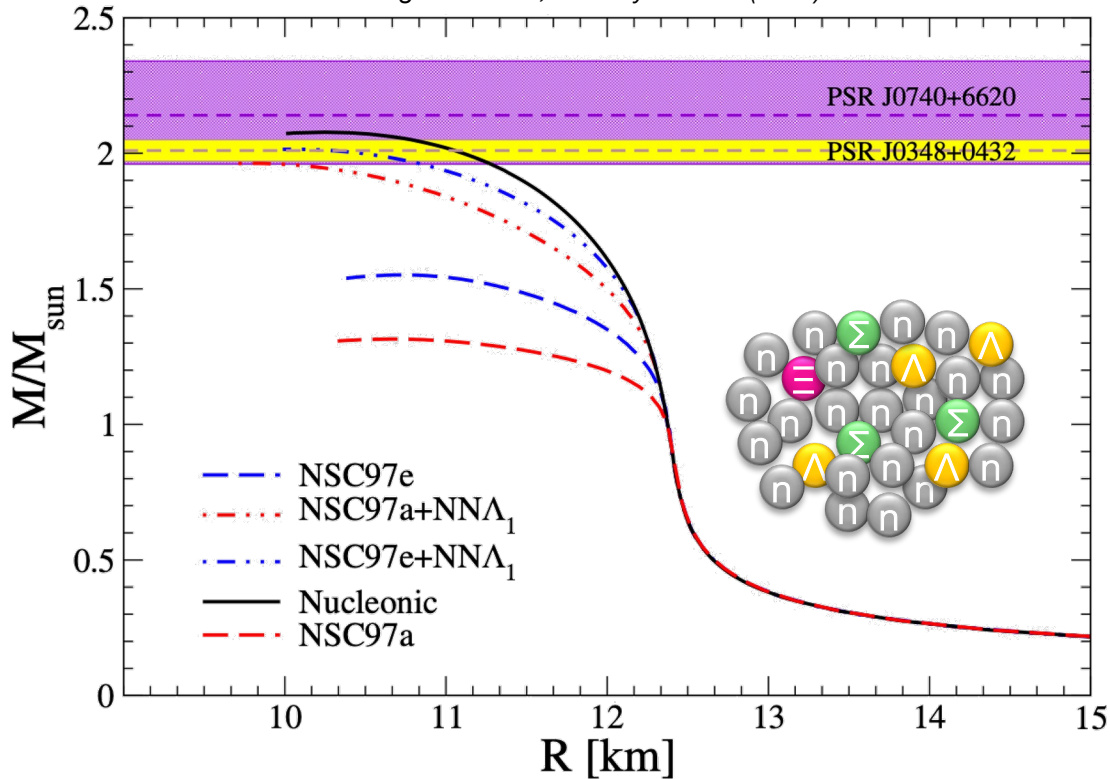
Additional slides



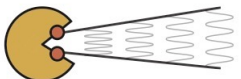
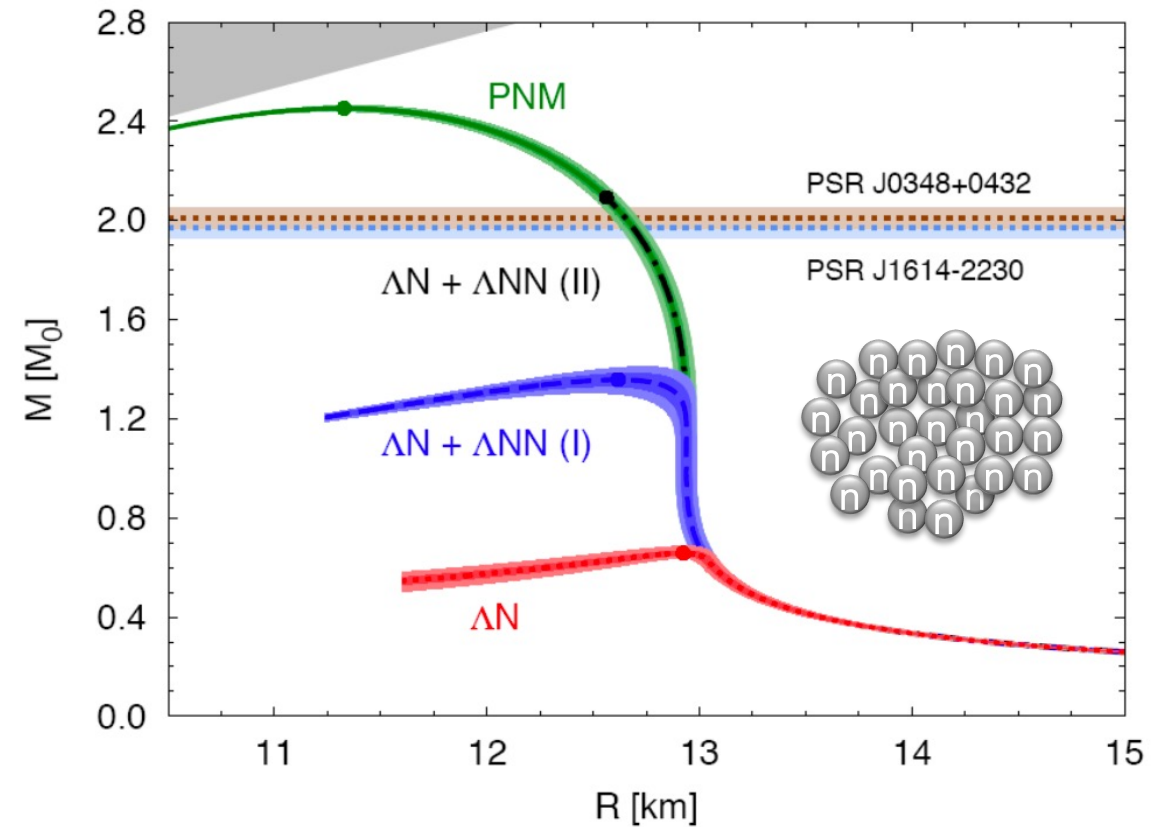
# Hyperonic three-body interactions in different EoS

- Repulsive 3-body  $\Lambda$ NN interaction can stiffen the EoS....but:
  - Effect on EoS largely model dependent  $\rightarrow$  too repulsive YNN leads to no hyperons in the NS

*D. Logoteta et al., Eur.Phys.J.A 55 (2019)*



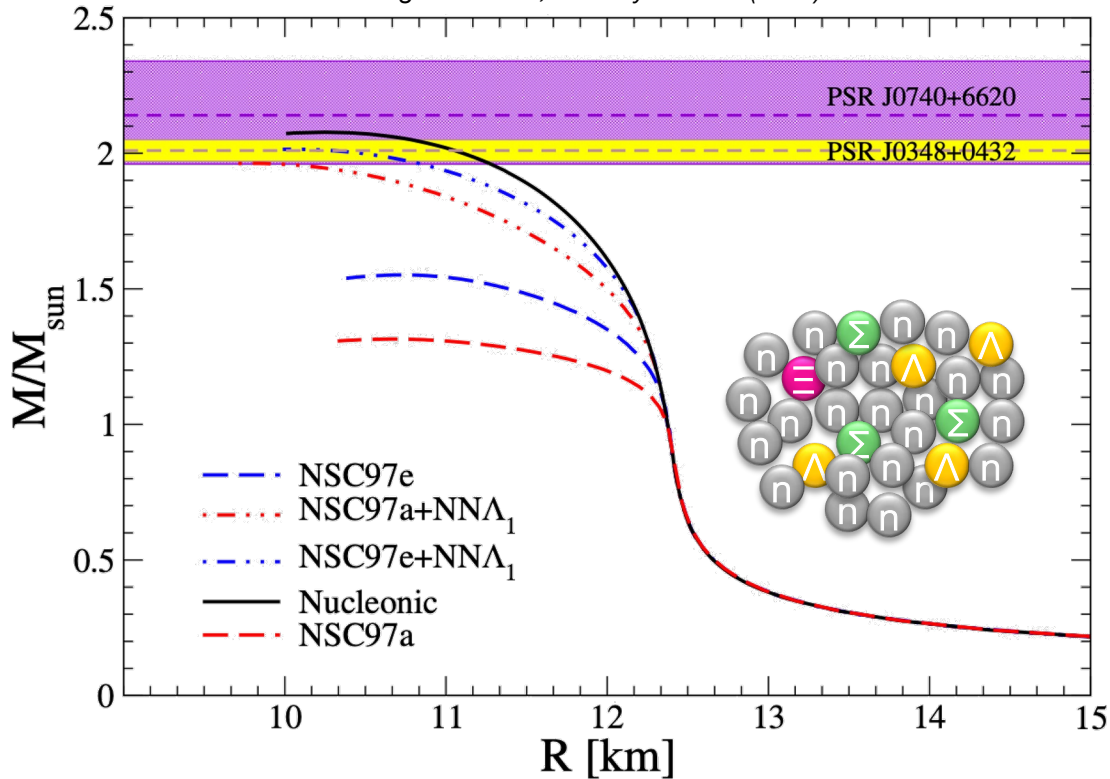
*D. Lonardoni et al., PRL 114 (2019)*



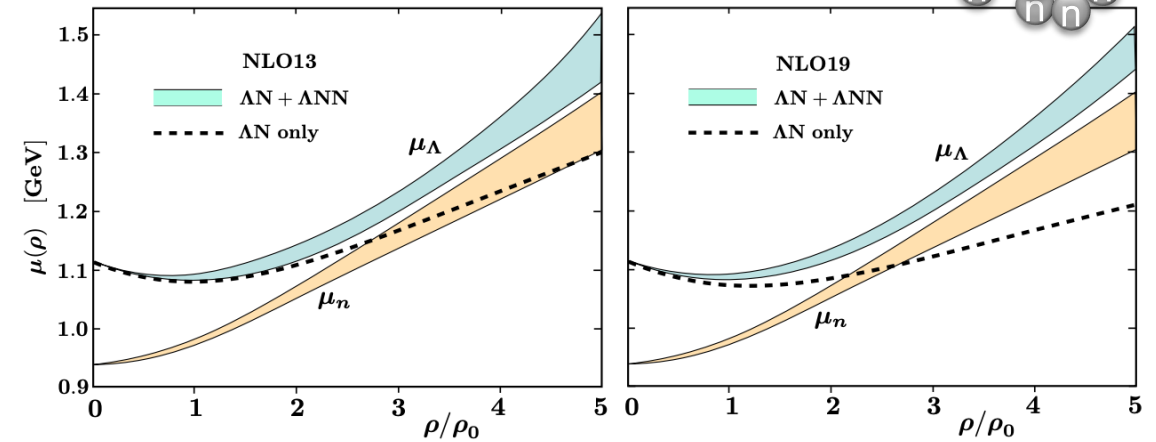
# Hyperonic three-body interactions in different EoS

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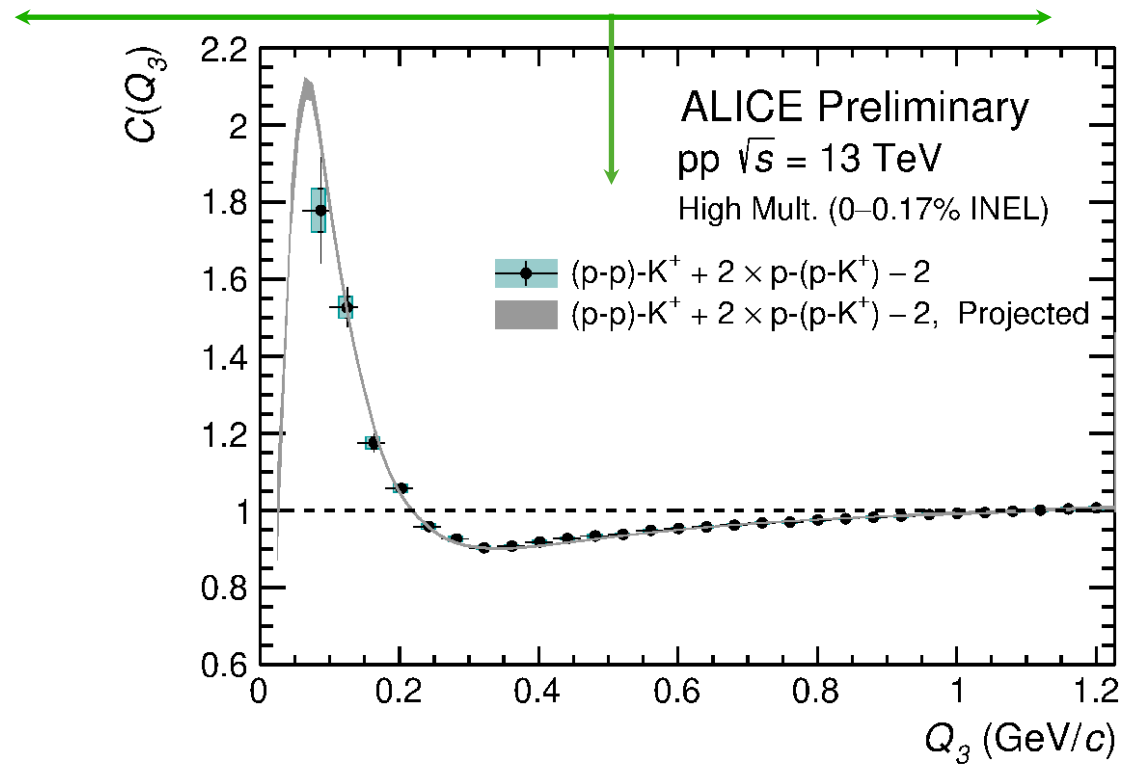
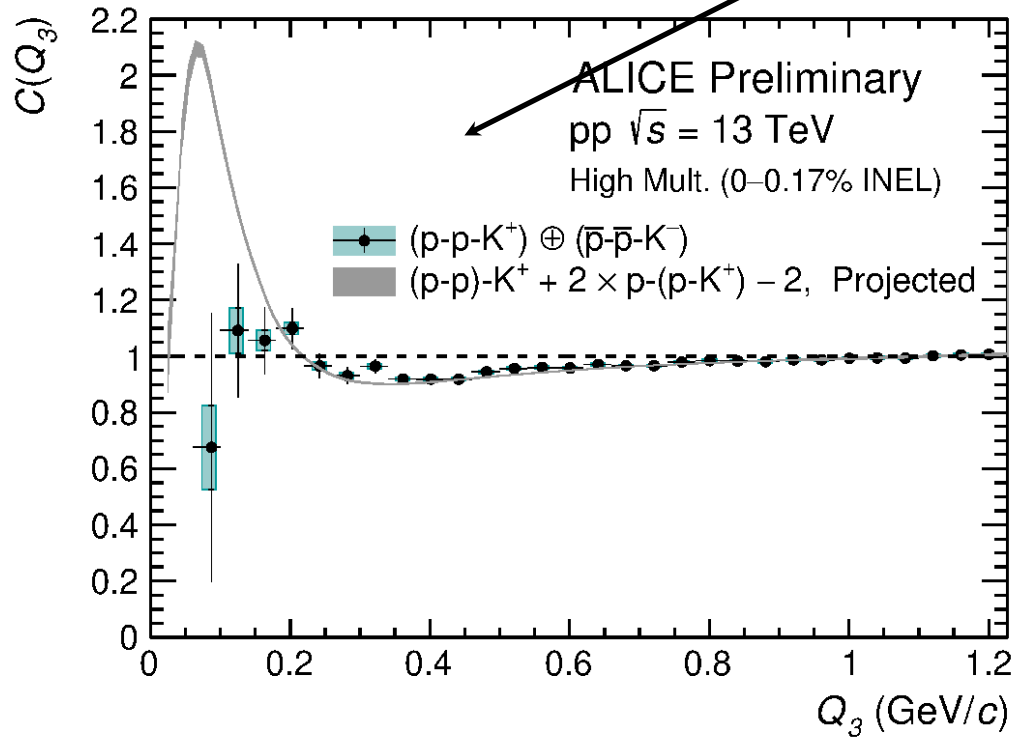
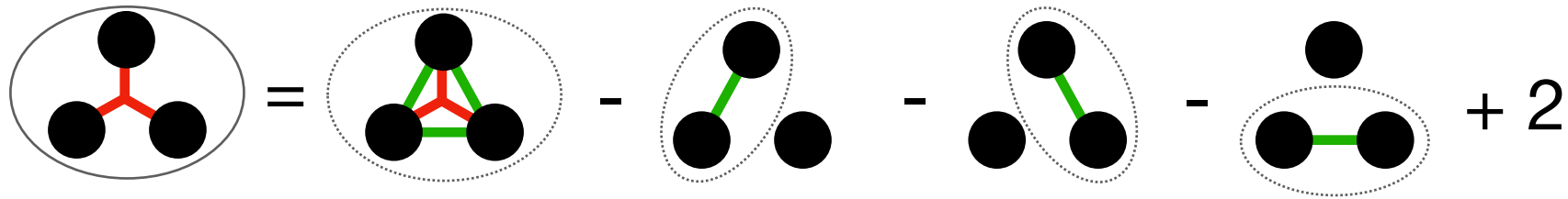
*D. Logoteta et al., Eur.Phys.J.A 55 (2019)*



*D. Gerstung et al. Eur.Phys.J.A 56 (2020)*



# TUM p-p-K<sup>+</sup> correlation function



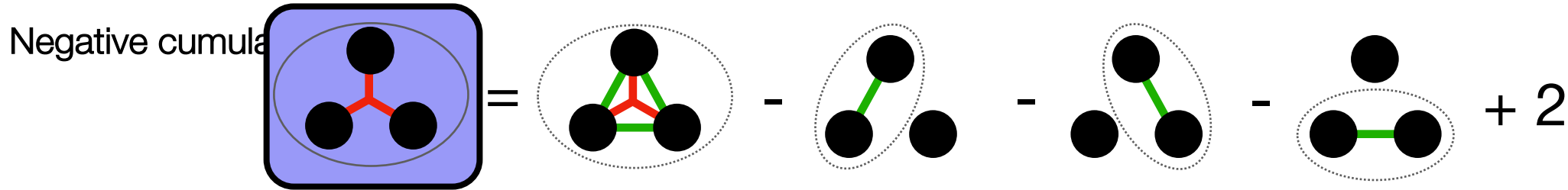
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ALI-PREL-513503

Already measured p-p [1] and newly obtained p-K<sup>+</sup> correlation

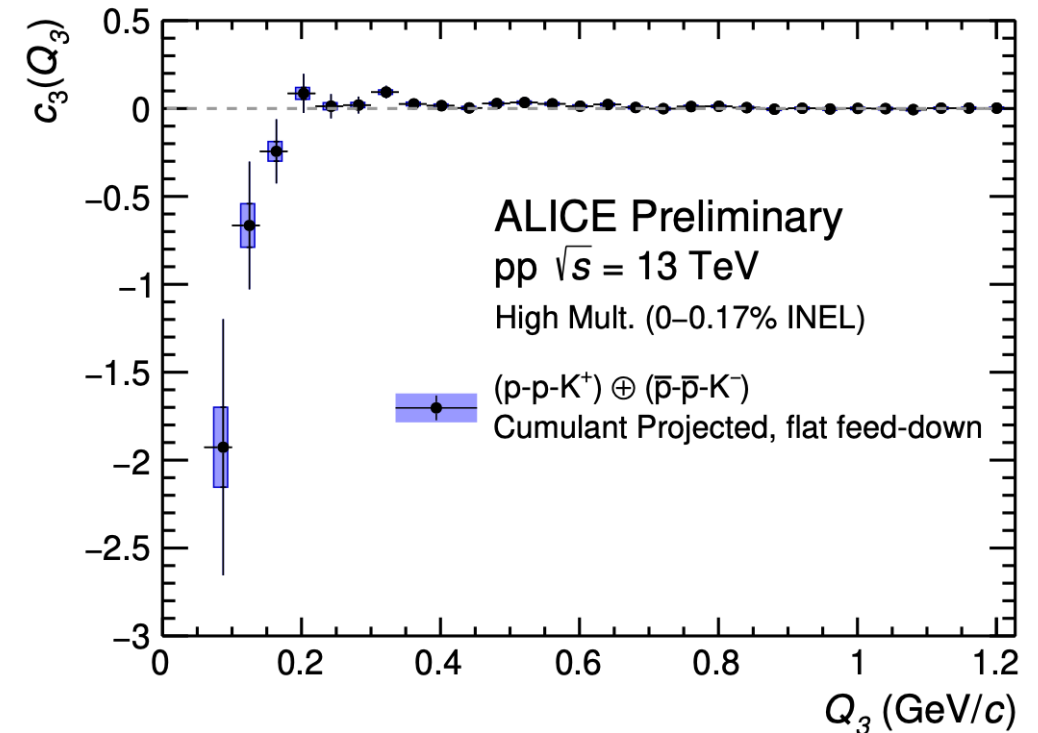
[1] PLB 805 (2020) 135419

functions used for projection.

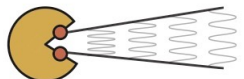


**Statistical significance:**  
 $n_\sigma = 2.3$  for  $Q_3 < 0.4 \text{ GeV}/c$

**Conclusion: the measured cumulant is compatible with zero within the uncertainties.**

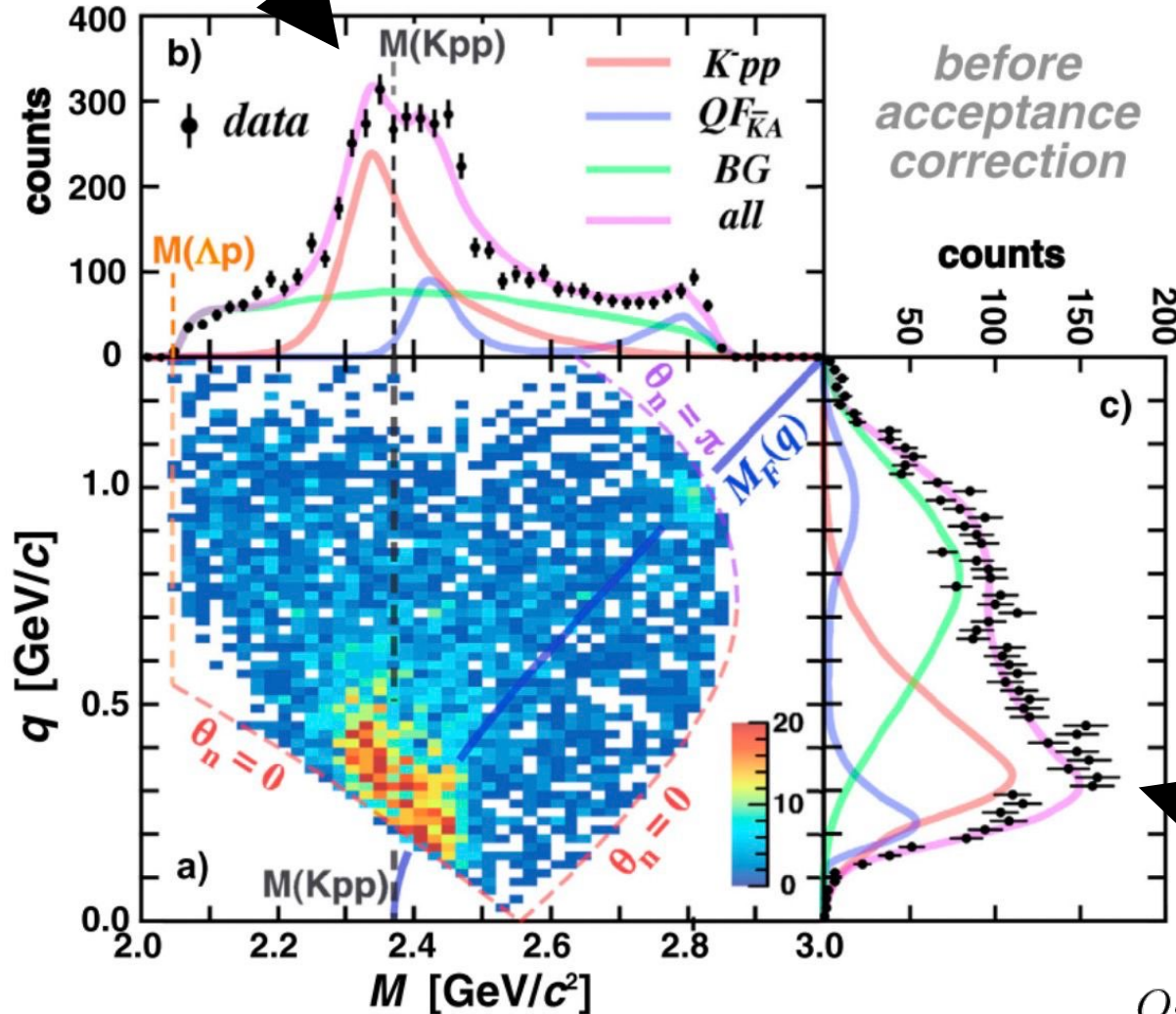


ALI-PREL-513592



# Kaonic bound state measured by E15

E15 Coll., PLB 789 (2019) 620



The E15 collaboration measured the bound state via the following decay:



The  $\Lambda p$  momentum distribution has a peak at

$$q = p_\Lambda + p_p \approx 0.35 \text{ GeV}/c$$

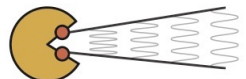
Using the momentum conservation:

$$p_{K^-} + p_p + p_p \approx 0.35 \text{ GeV}/c$$

The protons are at-rest  $\rightarrow p_K \approx 0.35 \text{ GeV}/c$

In terms of  $Q_3$  we have

$$Q_3 = 2\sqrt{k_{pK}^2 + k_{pK}^2 + k_{pp}^2} = 2\sqrt{2} k_{pK} = 4/3\sqrt{2} p_K < 0.5 \text{ GeV}/c$$



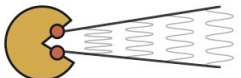
- For distinguishable pointlike particles

- Starting from the scattering parameters  $\Rightarrow$  define the s-wave two-particle relative wave function
- Considers Coulomb effects
- Assumption: p and d are pointlike particles!

$\Rightarrow$  **p–d scattering parameters** from fits to p–d scattering data

$S = 1/2$		$S = 3/4$		References
$f_0(\text{fm})$	$r_0(\text{fm})$	$f_0(\text{fm})$	$r_0(\text{fm})$	
$-1.30^{+0.20}_{-0.20}$	—	$-11.40^{+1.80}_{-1.20}$	$2.05^{+0.25}_{-0.25}$	Van Oers et al. [15]
$-2.73^{+0.10}_{-0.10}$	$2.27^{+0.12}_{-0.12}$	$-11.88^{+0.40}_{-0.10}$	$2.63^{+0.01}_{-0.02}$	Arvieux et al. [16]
-4.0	—	-11.1	—	Huttel et al. [17]
-0.024	—	-13.7	—	Kievsky et al. [18]
$0.13^{+0.04}_{-0.04}$	—	$-14.70^{+2.30}_{-2.30}$	—	Black et al. [19]

*Convention sign: In this presentation positive (negative)  $f_0$  means attractive (repulsive) interaction*





# p-d correlation with d as composite object

The three body wave function with proper treatment of 2N and 3N interaction at very short distances goes to a p-d state.

- **Three-body wavefunction for p-d:**  $\Psi_{m_2, m_1}(x, y)$  describing three-body dynamics, anchored to p-d scattering observables.

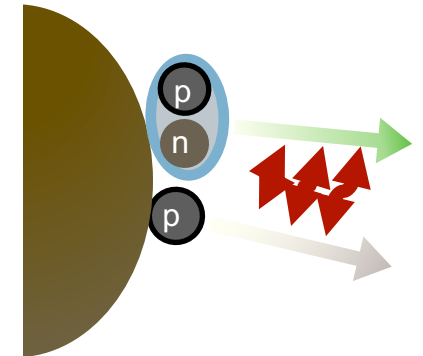
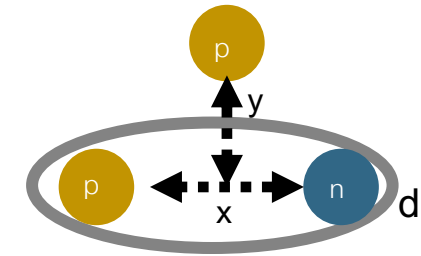
- $x$  = distance of p-n system within the deuteron
- $y$  = p-d distance
- $m_2$  and  $m_1$  deuteron and proton spin

- $\Psi_{m_2, m_1}(x, y)$  three-nucleon wave function asymptotically behaves as p-d state:

$$\Psi_{m_2, m_1}(x, y) = \underbrace{\Psi_{m_2, m_1}^{(\text{free})}}_{\text{Asymptotic form}} + \sum_{LSJ}^{J \leq \bar{J}} \underbrace{\sqrt{4\pi i^L} \sqrt{2L+1} e^{i\sigma_L} (1m_2 \frac{1}{2} m_1 |SJ_z)(LOSJ_z | JJ_z)}_{\text{Strong three-body interaction}} \tilde{\Psi}_{LSJJ_z}$$

$\tilde{\Psi}_{LSJJ_z}$  describe the configurations where the three particles are close to each other

$\Psi_{m_1, m_2}^{(\text{free})}$  an asymptotic form of p-d wave function



Kievsky et al, Phys. Rev. C 64 (2001) 024002  
 Kievsky et al, Phys. Rev. C 69 (2004) 014002  
 Deltuva et al, Phys. Rev. C 71 (2005) 064003

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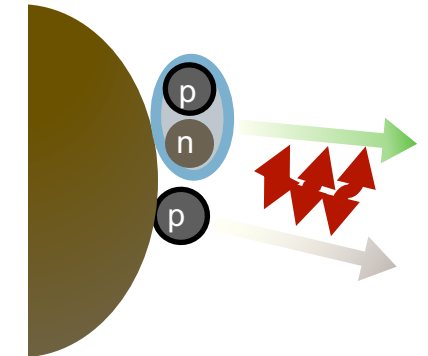
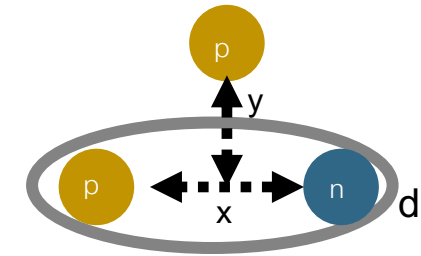
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 Kievsky et al, Phys. Rev. C 69 (2004) 014002  
 Deltuva et al, Phys. Rev. C 71 (2005) 064003

# TUM The correlation function

- **Starting with the PPN state that goes into pd state:**

- Nucleons with the Gaussian sources distributions

Single-particle Gaussian emission source

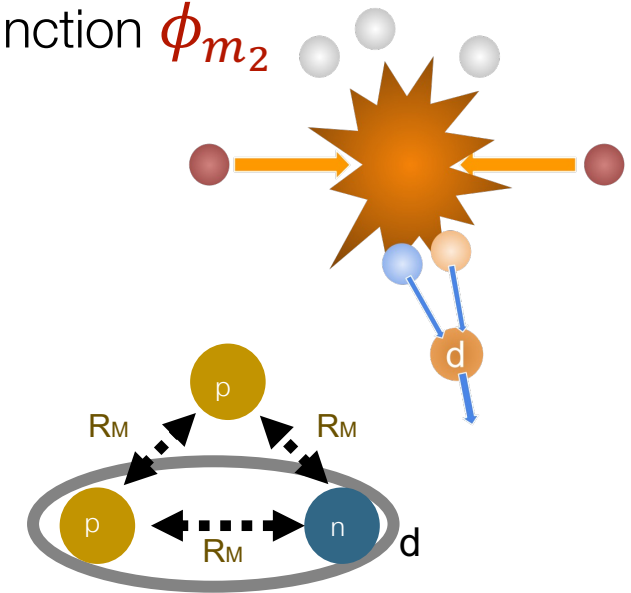
$$A_d C_{pd}(k) = \frac{1}{6} \sum_{m_2, m_1} \int d^3 r_1 d^3 r_2 d^3 r_3 S_1(r_1) S_1(r_2) S_1(r_3) |\Psi_{m_2, m_1}|^2,$$

- Where  $A_d$  is the deuteron formation probability using deuteron wavefunction  $\phi_{m_2}$

$$A_d = \frac{1}{3} \sum_{m_2} \int d^3 r_1 d^3 r_2 S_1(r_1) S_1(r_2) |\phi_{m_2}|^2,$$

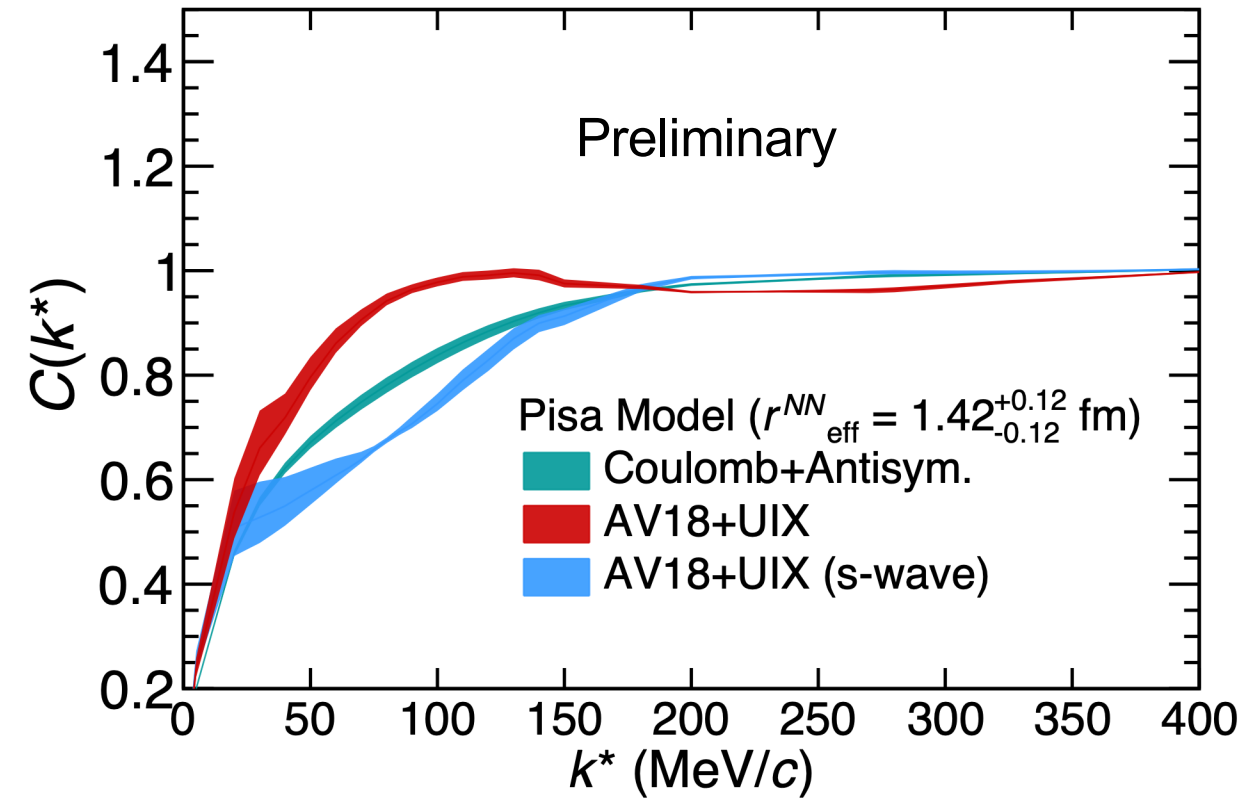
- Final definition of the correlation with p-p source size  $R_M$  :

$$A_d C_{pd}(k) = \frac{1}{6} \sum_{m_2, m_1} \int \rho^5 d\rho d\Omega \frac{e^{-\rho^2/4R_M^2}}{(4\pi R_M^2)^3} |\Psi_{m_2, m_1}|^2.$$



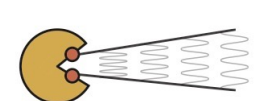
# TUM PISA model: theoretical p-d correlation

Model including NN and NNN interactions in s+d-wave

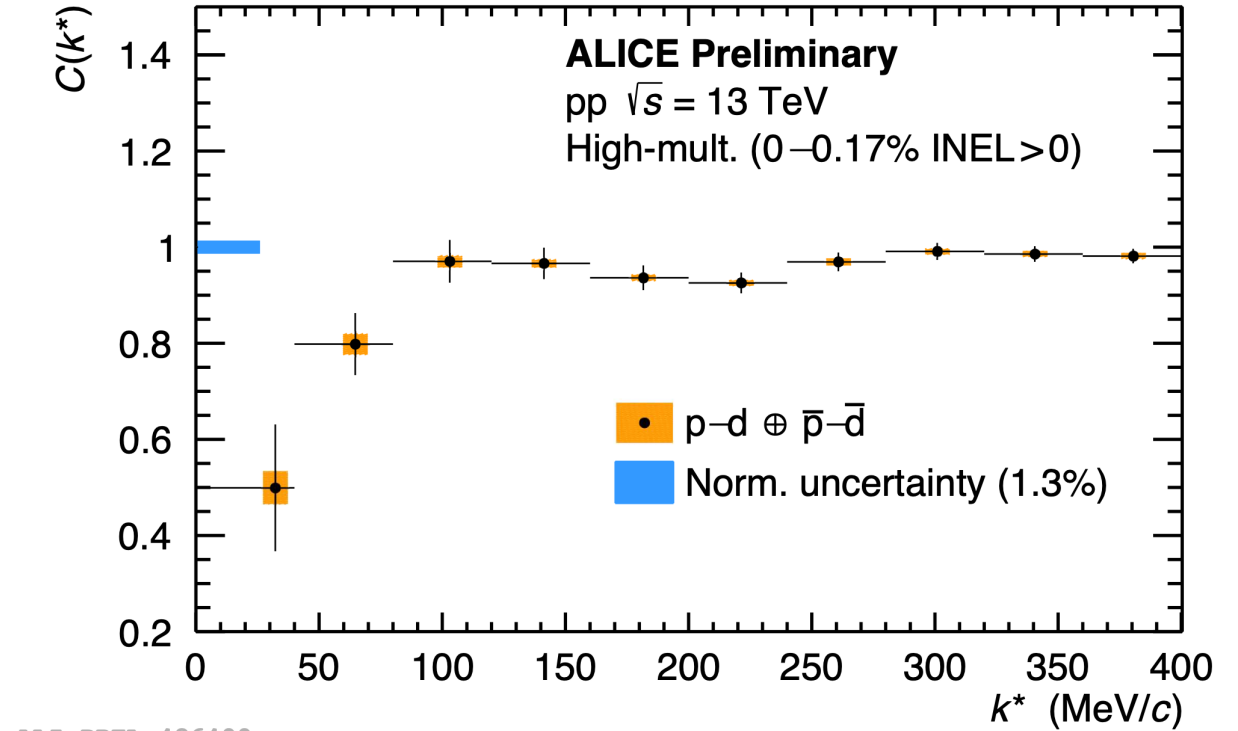
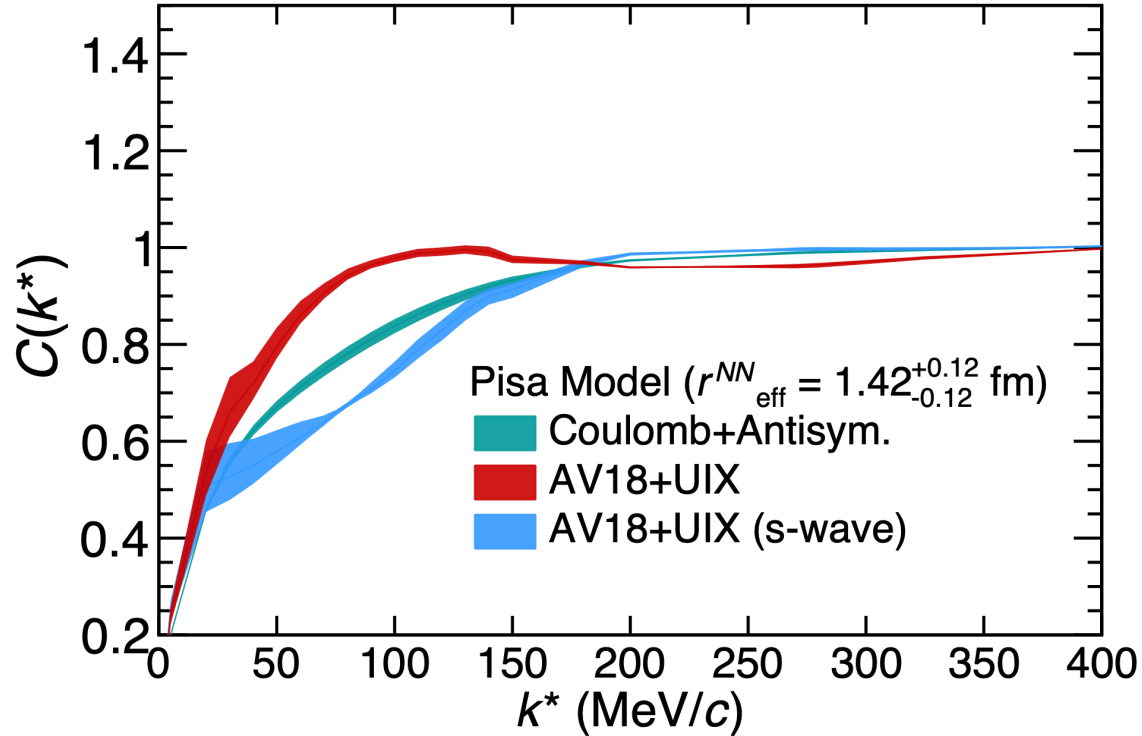


- Coulomb interaction not enough
- Coulomb +s+p+d wave:  
Source size =  $1.08 \pm 0.06 \text{ fm} \Rightarrow 1.42 \pm 0.12 \text{ fm}$  of pp source size

Calculation by Michele Viviani

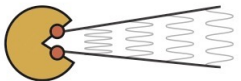


# TUM PISA model: theoretical p-d correlation



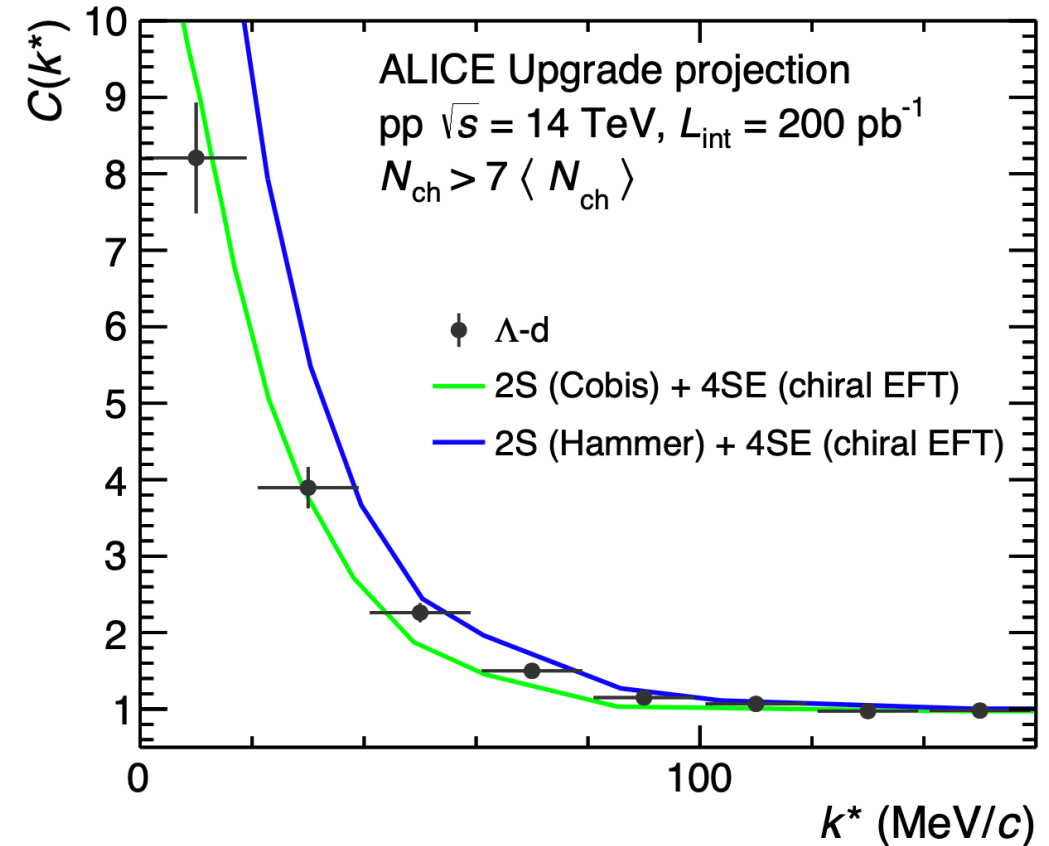
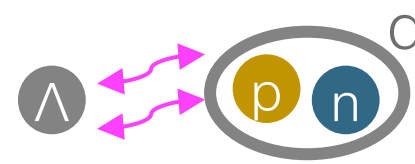
ALI-PREL-486400

- Model calculation qualitatively reproduces the data
- The p-d correlation is affected by two + three-body p-p-n interactions!

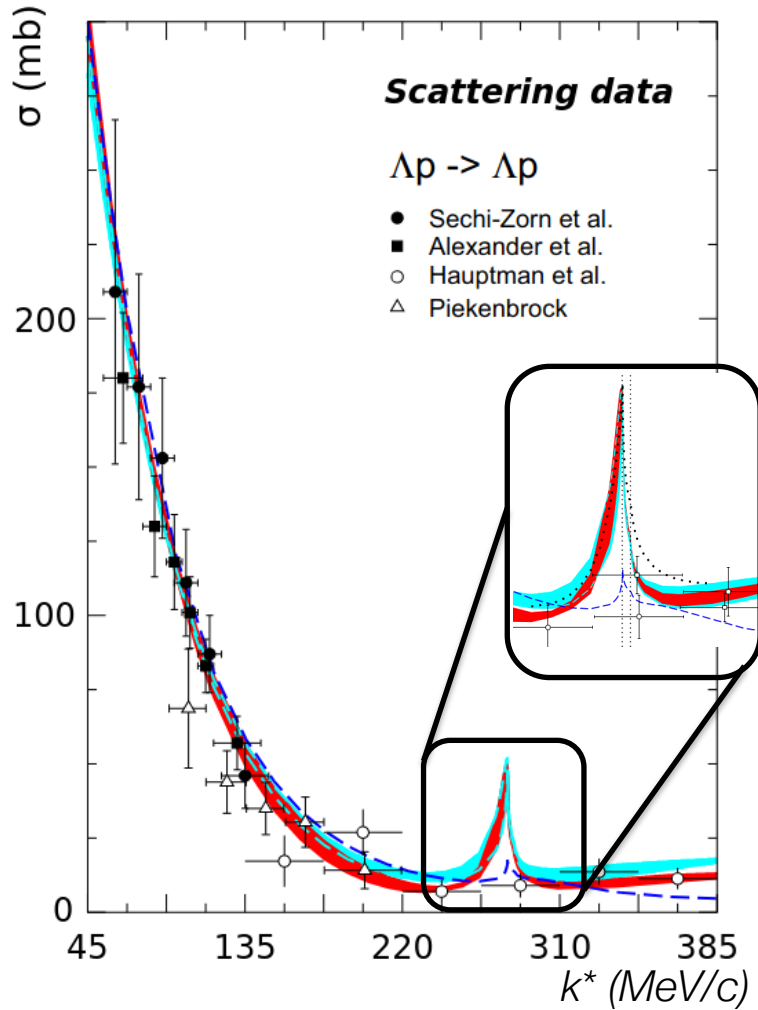


# TUM Deuteron- $\Lambda$ correlations

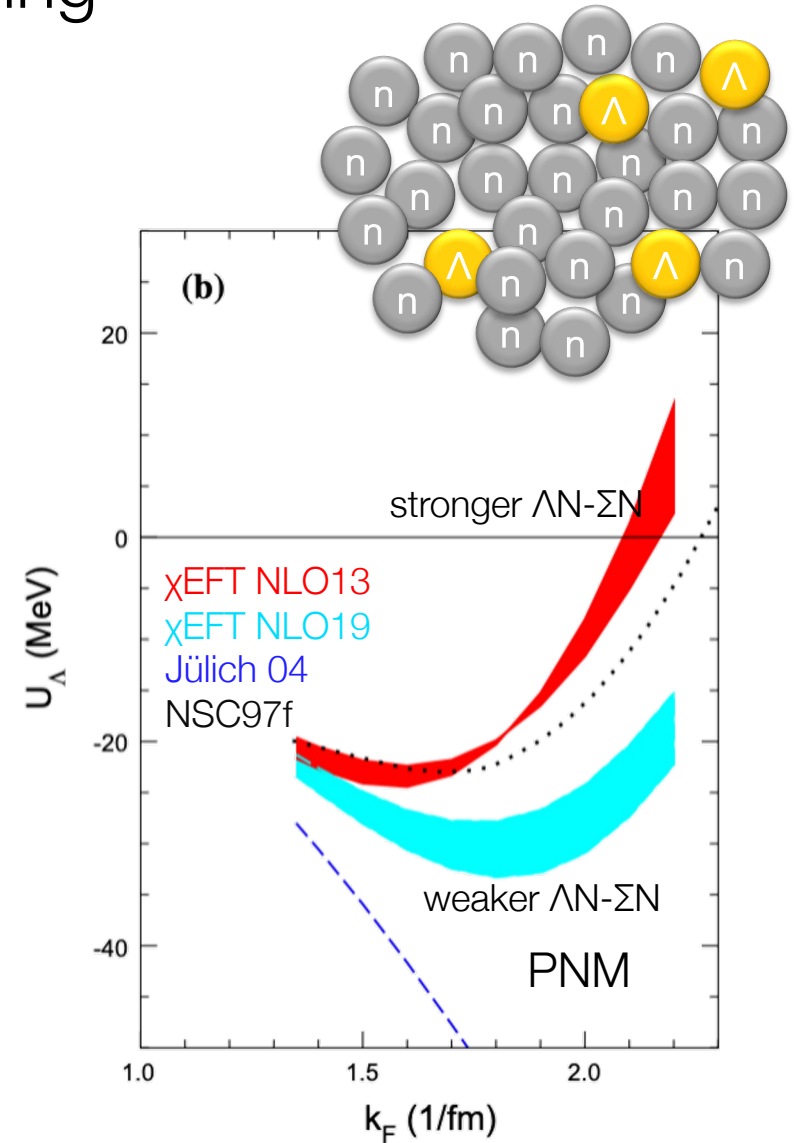
- Explorative theoretical studies by J. Haidenbauer, arXiv:2005.05012
- Scattering parameters of the doublet state can be related to hypertriton  ${}^3\text{H}_\Lambda$  binding energy
  - Synergy with direct hypertriton measurements
- Quartet state depends on the spin triplet  $\Lambda$ -N interaction
  - No experimental information on the spin dependence of the  $\Lambda$ -N interaction
  - Quartet becomes dominant for large source sizes
- Measurement in Run 2 ongoing, improved precision beyond!



ALICE Collab., Physics motivations for a high energy pp programme in ALICE after LS2  
 A. Cobis *et al.*, J. Phys. G 23 (1997) 401  
 H. Hammer, Nucl. Phys. A 705 (2002) 173  
 J. Haidenbauer, arXiv:2005.05012



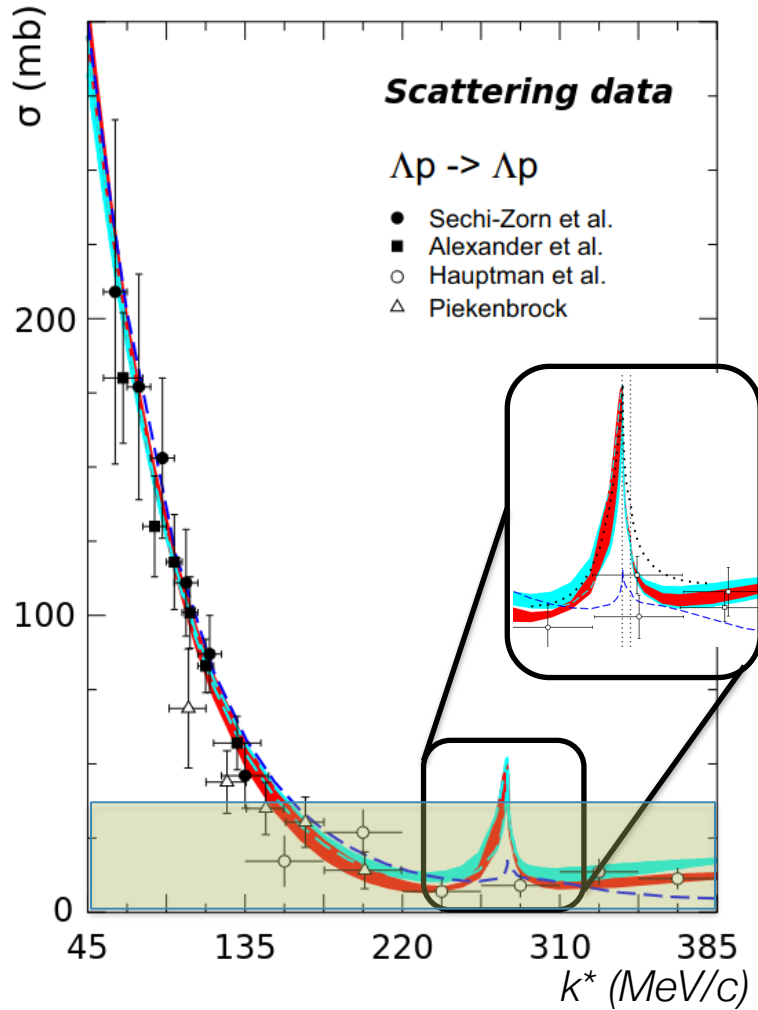
- $\Sigma$ N coupling strength relevant for EoS
- deeply affects the behaviour of  $\Lambda$  at finite density
- implications for 3-body interactions<sup>(\*)</sup>



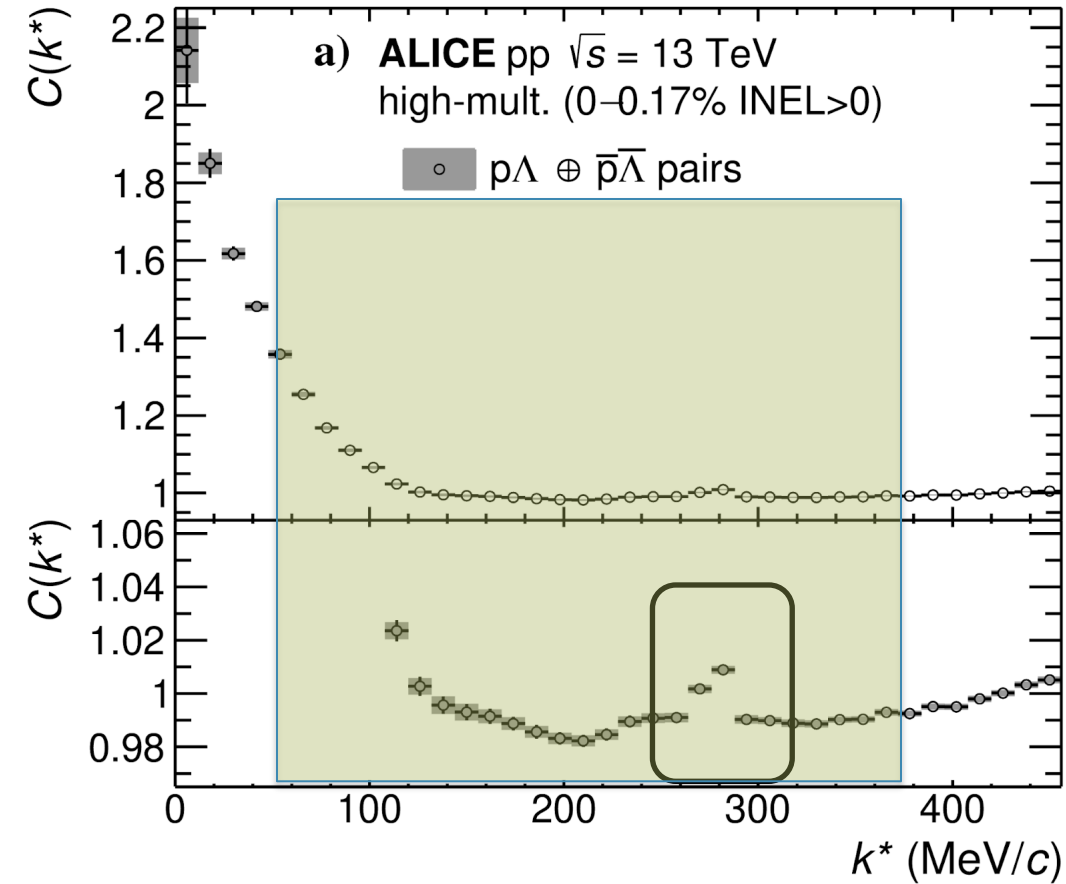
NLO13: J.Haidenbauer, N.Kaiser et al., NPA 915, 24 (2013)

NLO19: J.Haidenbauer, U. Meißner, Eur.Phys.J.A 56 (2020)

(\*)D. Gerstung et al. Eur.Phys.J.A 56 (2020) 6, 175



- Extension of kinematic range
- Measurement down to zero momentum
- Factor 20 improved precision in data (<1%)
- First experimental evidence of  $\Sigma N$  cusp in 2-body channel



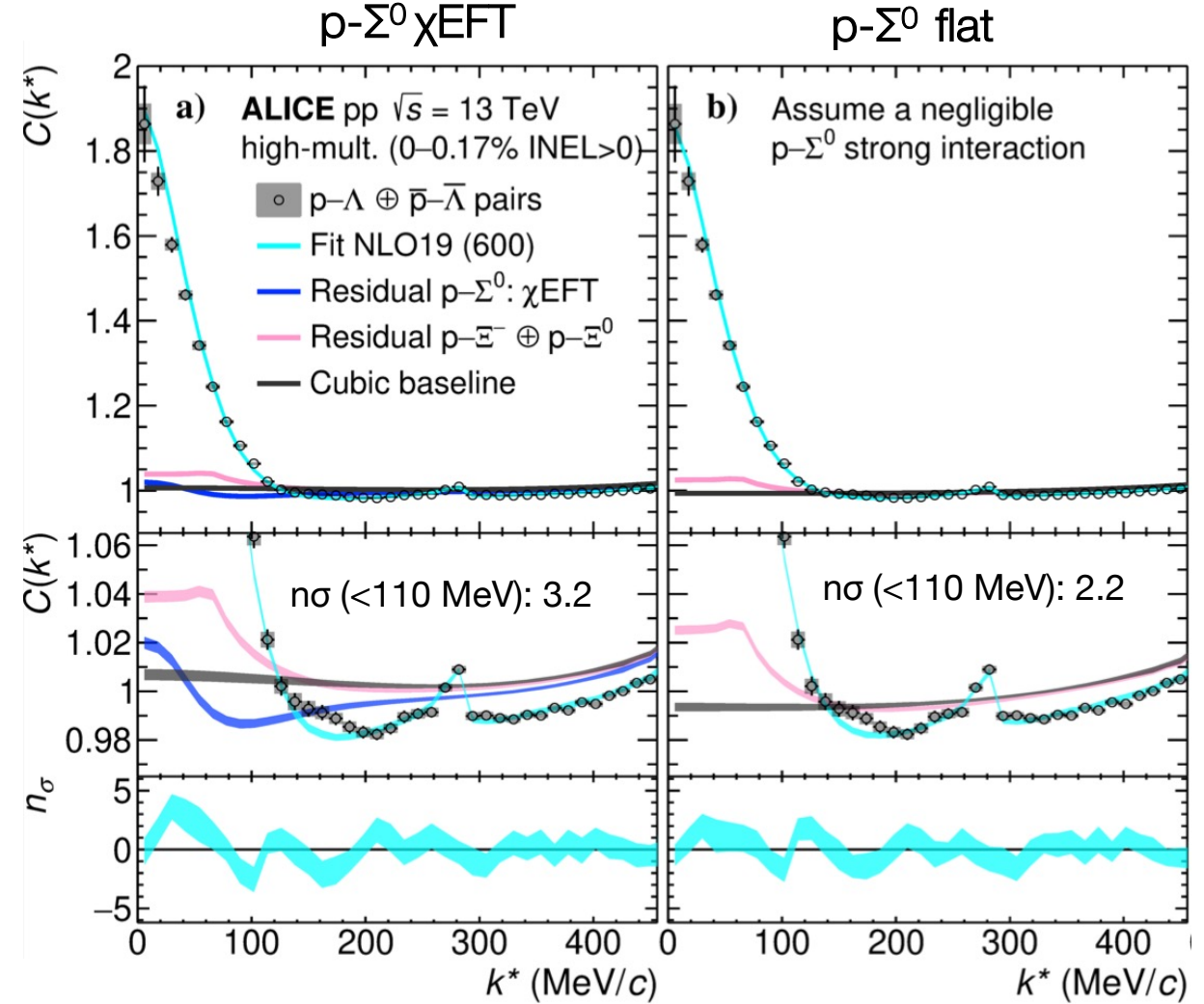
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 NLO19: J.Haidenbauer, U. Meißner, Eur.Phys.J.A 56 (2020)

(\*)D. Gerstung et al. Eur.Phys.J.A 56 (2020) 6, 175



- New scenario arising for  $\Lambda N$ - $\Sigma N$  interaction
  - NLO19 potentials favoured
    - weaker  $\Lambda N$ - $\Sigma N$  coupling
    - large attraction of  $\Lambda$  at large densities
    - $\Lambda NN$  large repulsion needed: no hyperons!<sup>(1)</sup>
  - shallow  $p$ - $\Sigma^0$  interaction favoured, in agreement with measured  $p$ - $\Sigma^0$  correlation<sup>(2)</sup>
  - Future useful measurements:
    - $p\Sigma^0$ ,  $p\Sigma^{+,-}$ <sup>(3)</sup>,  $\Lambda d$  correlations in LHC Run 3 and Run 4
    - $p\Sigma^{+,-}$  scattering data J-PARC E40

What about more strange hyperons?

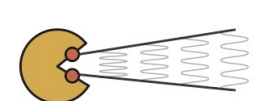


ALICE Coll. PLB 833 (2022), 137272

(1) D. Gerstung et al. Eur.Phys.J.A 56 (2020) 6, 175

(2) ALICE Coll. PLB 805 (2020) 135419

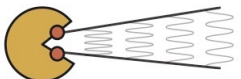
(3) J. Haidenbauer, U.G. Meißner PLB 829 (2022)



# TUM Comparison with $\chi$ EFT potentials

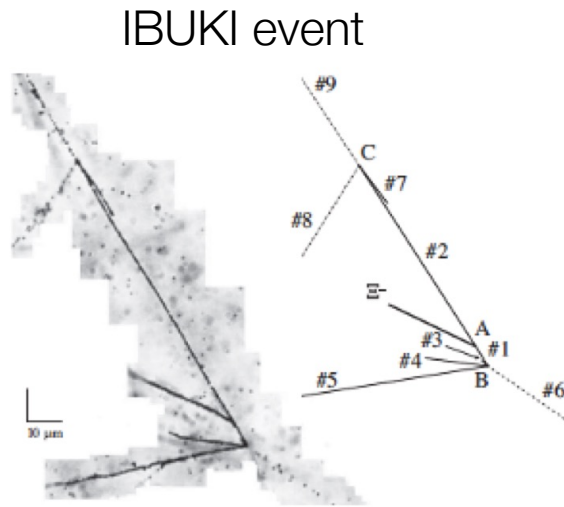
- Reduction of attraction in  $^3S_1$  improve the description of the data both at low  $k^*$  and in the  $\Sigma N$  cusp region:  $n\sigma=1.9$  (negl.  $p\Sigma$ )
- Slight sensitivity to CSB inclusion in singlet state

$p-\Sigma^0$ ( $\rightarrow$ ) $p-\Lambda$ ( $\downarrow$ )	Standard deviation ( $n_\sigma$ )	
	$\chi$ EFT	Negligible FSI
LO13-600	5.6 (7.5)	10.2 (10.3)
NLO13-500	7.3 (10.3)	4.5 (5.3)
NLO13-550	4.4 (6.5)	2.1 (2.2)
NLO13-600	5.8 (5.8)	3.3 (3.6)
NLO13-650	5.5 (5.5)	4.2 (5.0)
NLO19-500	5.6 (7.2)	3.0 (3.0)
NLO19-550	4.5 (4.3)	1.8 (2.2)
<b>NLO19-600</b>	3.7 (3.9)	1.6 (3.0)
NLO19-650	3.7 (3.7)	2.3 (4.1)



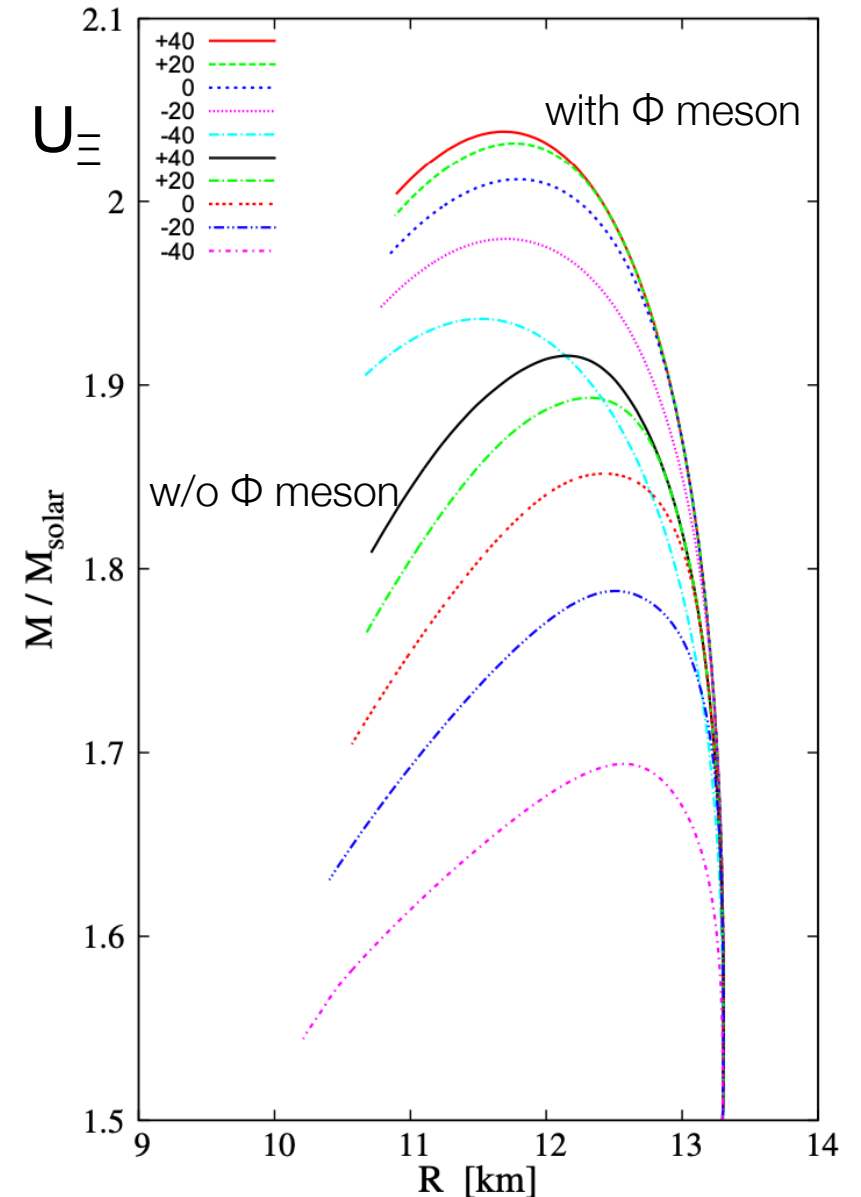
# TUM Effect of $\Xi$ hyperon in neutron stars

- $U_{\Xi}$  plays a crucial role in the EoS of NSs  
→ repulsive single-particle pot. stiffens the EoS
- continuous effort in hypernuclei measurements at J-PARC (E07, S-2S,)

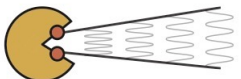


J-PARC E07 Coll. PRL 126 (2021)

- What about the  $N\Xi$  interaction in vacuum?



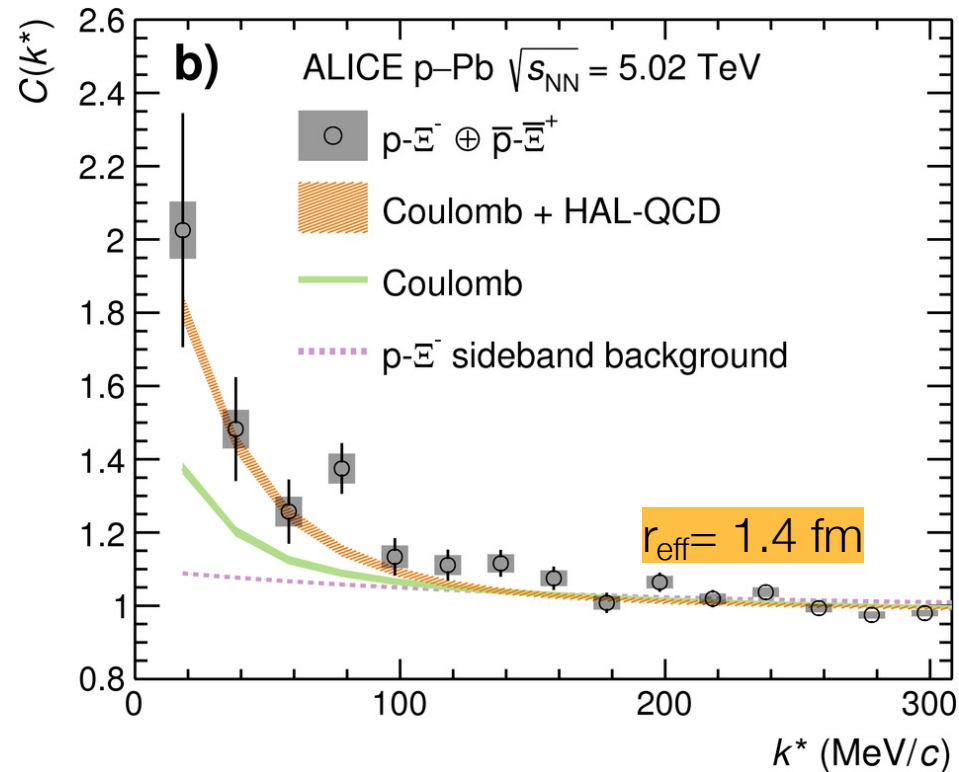
Weissenborn et al., Nucl. Phys. A 881 (2012)



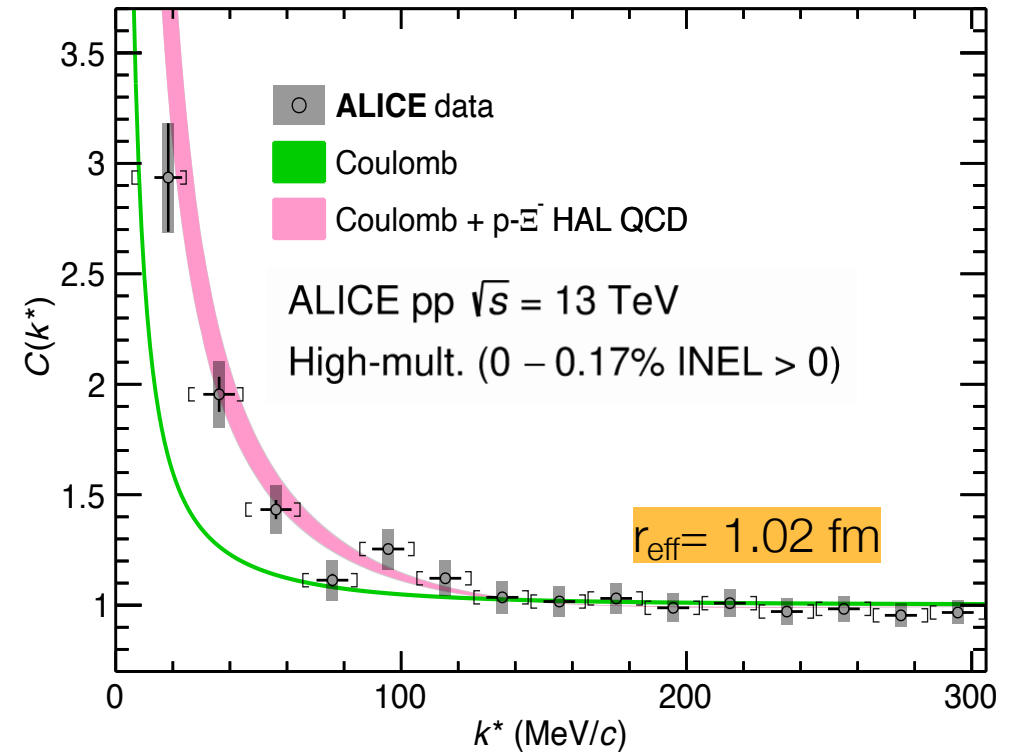
# First measurements of the $p-\Xi^-$ interaction at LHC

- Observation of the strong interaction beyond Coulomb
- Agreement with lattice calculations confirmed in pp and p-Pb colliding systems
- **At finite density HAL QCD potentials predict in PNM a slightly repulsive  $U_{\Xi} \sim +6 \text{ MeV}^{(*)} \rightarrow$  stiffening of the EoS**

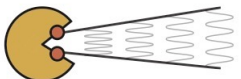
ALICE Coll, Phys. Rev. Lett 123, (2019) 112002



ALICE Coll. Nature 588, 232–238 (2020)



(\*) HAL QCD Coll., PoS INPC2016 (2016) 277



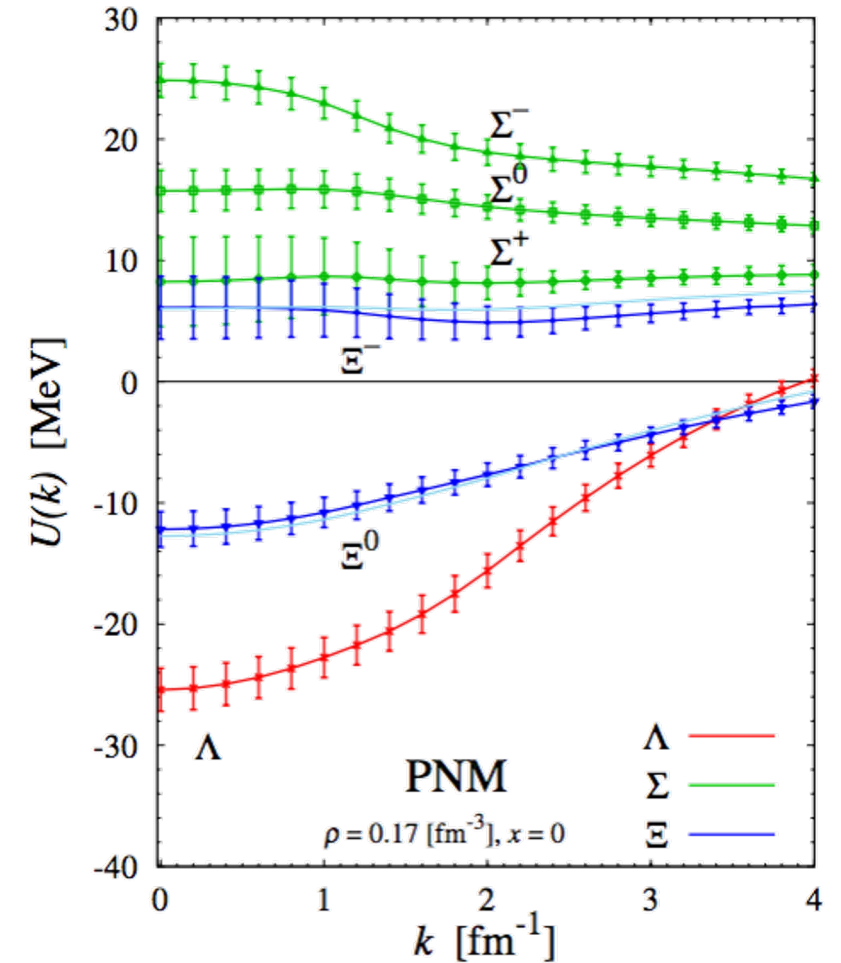
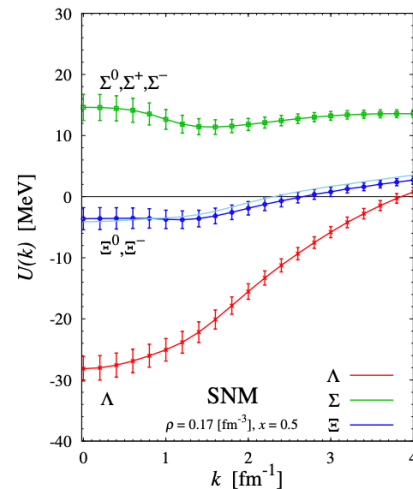
# TUM $\Xi$ single particle potential from HAL QCD calculations

HAL-QCD Collaboration, arXiv:1809.08932 (2018)

- NS environment  $\Rightarrow$  Pure Neutron Matter
- Tested HAL-QCD potential in vacuum with ALICE  $\Rightarrow$  Brueckner-Hartree-Fock many-body calculations  $\Rightarrow$   $U_{\gamma}$  single-particle potential of hyperons in nucleonic matter

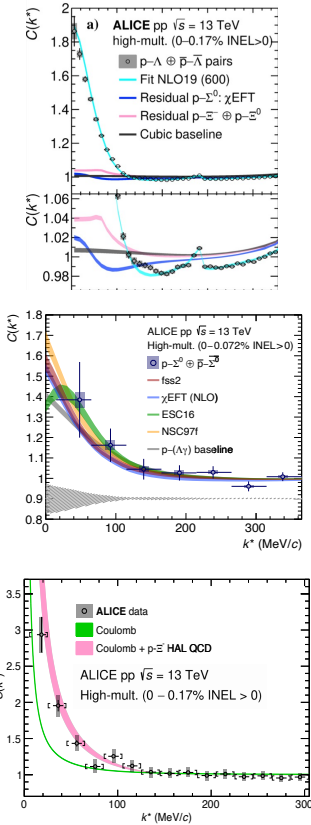
- At saturation density in PNM:  $U_{\Xi}$  slightly repulsive

- in SNM:



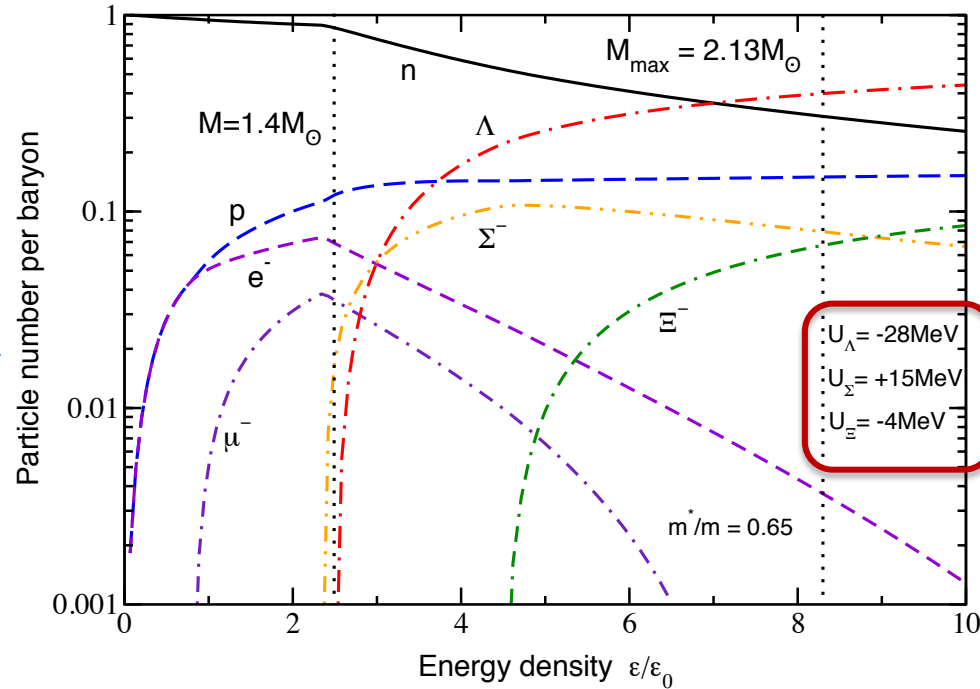
# TUM Towards a solution for the puzzle?

Correlation  
2-body interaction  
(sc. exp, hypernuclei)



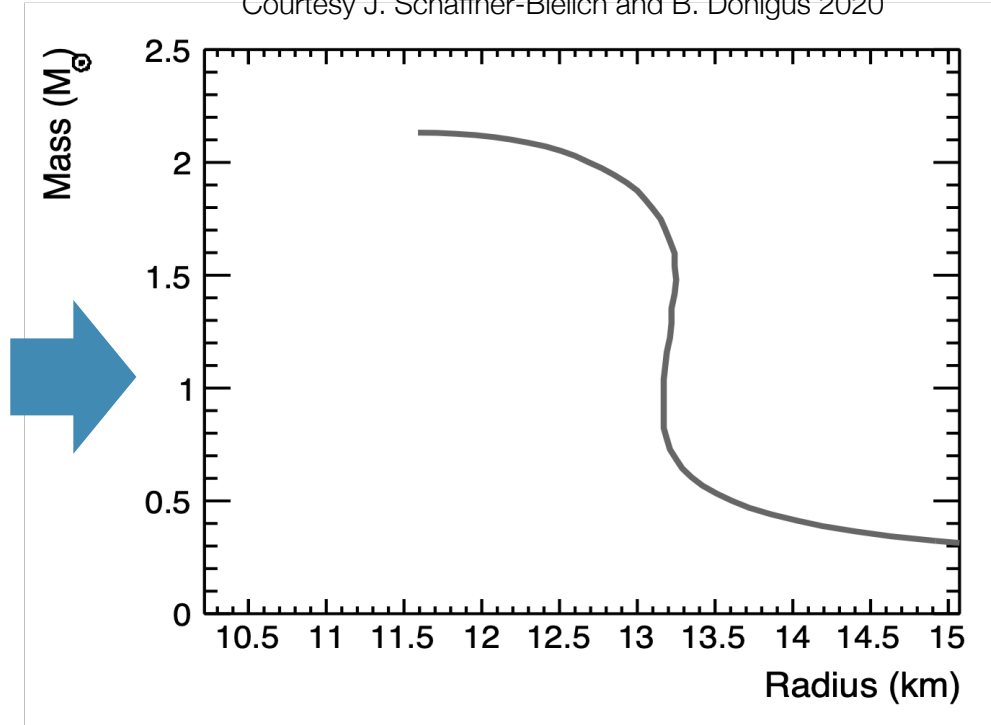
Single-particle potentials  
EoS

Courtesy J. Schaffner-Bielich 2020



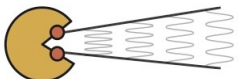
Mass vs Radius relation  
for hyperon stars

Courtesy J. Schaffner-Bielich and B. Dönig 2020



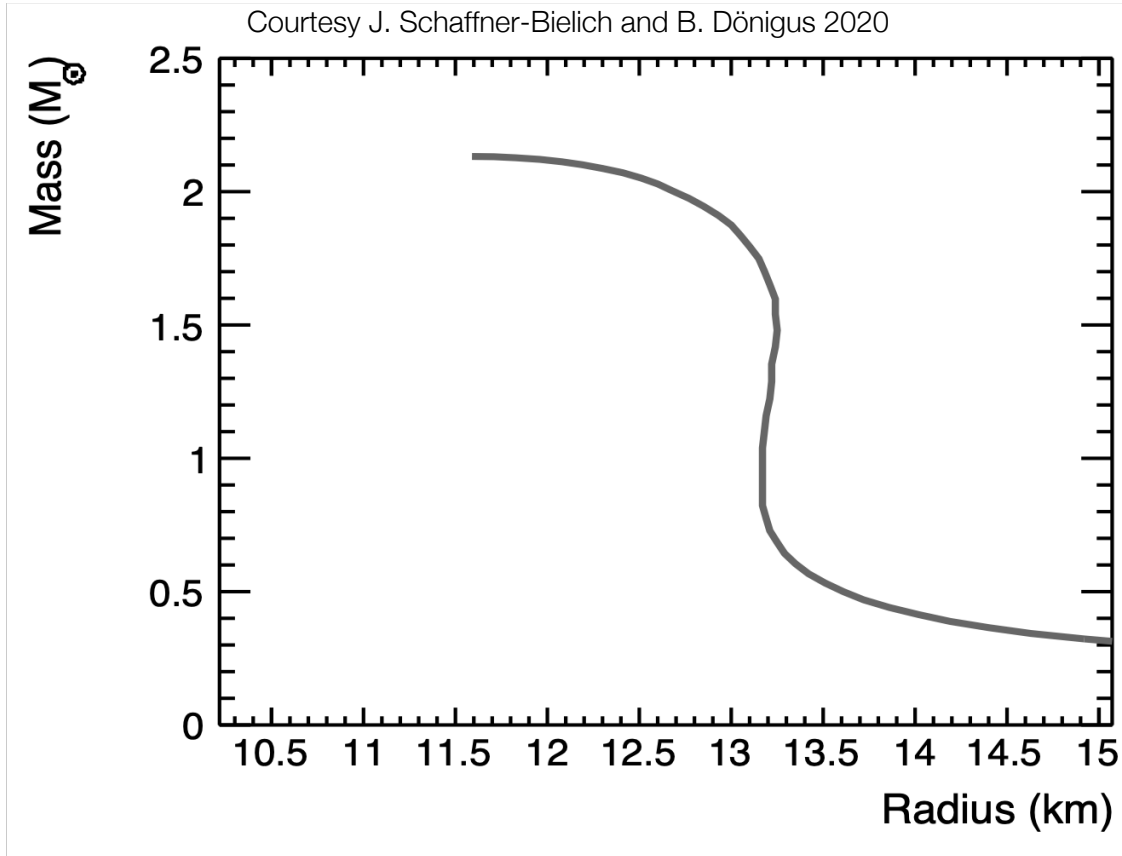
This is only an example.  
Experimental uncertainties need to be propagated and  
some interactions are missing ...

V. M. S., L. Fabbietti and O. Vazquez-Doce  
*Ann.Rev.Nucl.Part.Sci.* 71 (2021)  
S. Weissenborn et al., *J. NPA* 881 (2012)  
J. Schaffner-Bielich, I. Mishustin, *PRC* 53 (1996)  
N. Hornick et al., *PRC* 98 (2018)

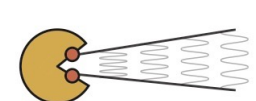
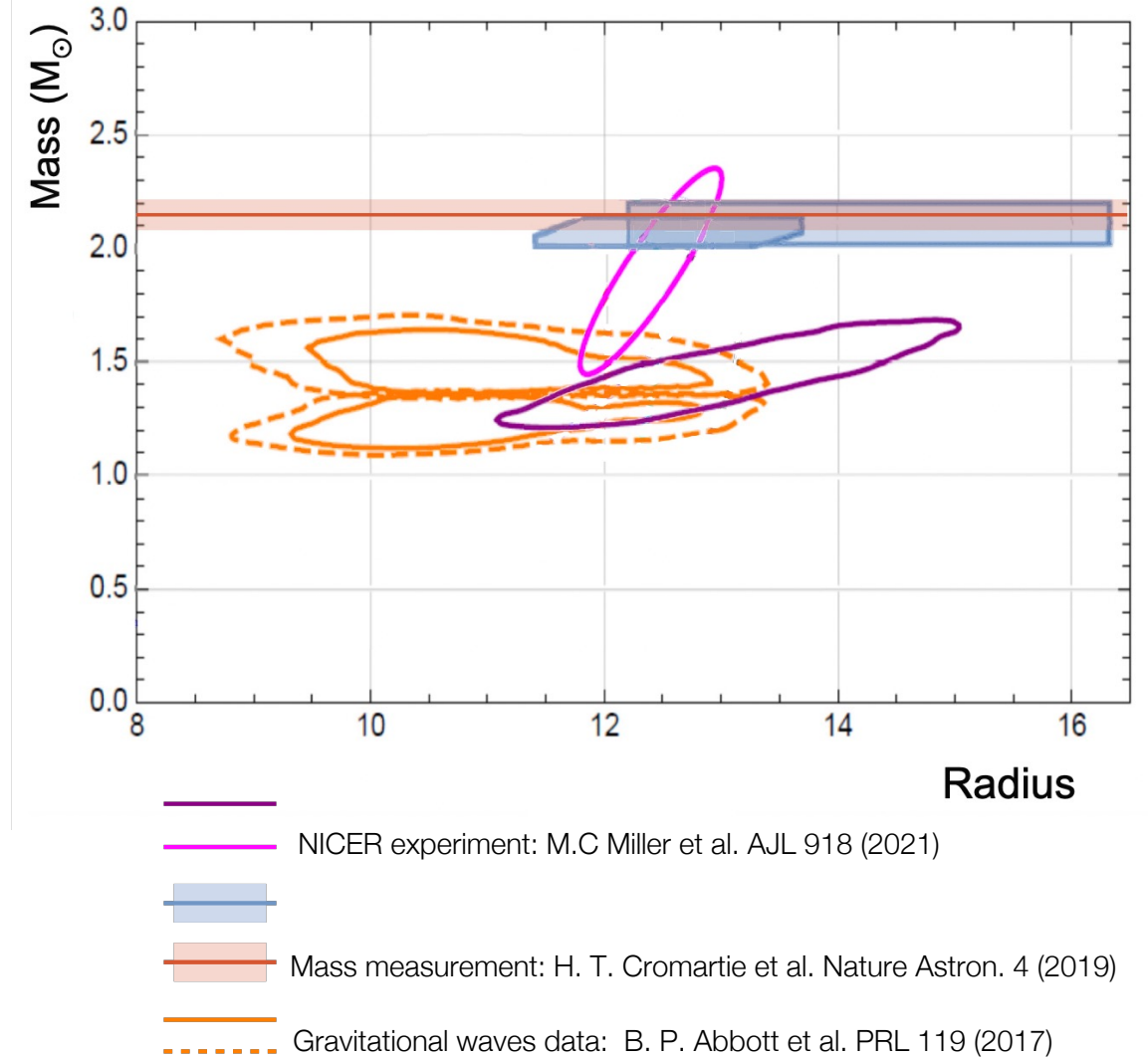


# TUM Test of the hyperon equation of state

Mass vs Radius relation  
for hyperon stars



Mass and Radius measurements  
from astrophysics



# TUM First measurement of the $p\Sigma^0$ interaction

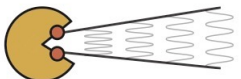
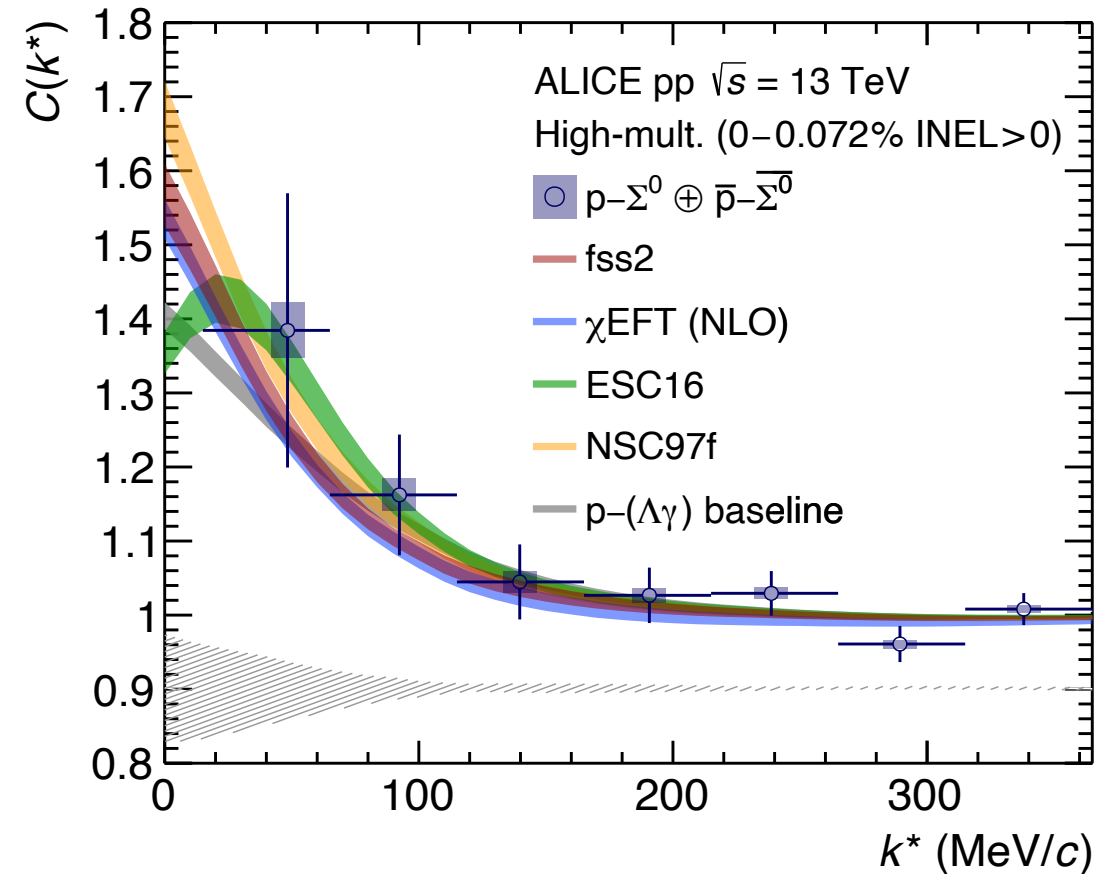
- Experimental data on hypernuclei too scarce for a final conclusion: attractive or repulsive interaction?
- Very challenging measurement via the difficult electromagnetic decay  $\Sigma^0 \rightarrow \Lambda \gamma$

- **Correlation function is above the background  $\Rightarrow$  pointing to a very shallow attractive interaction**

Model	p-( $\Lambda\gamma$ ) baseline	fss2	$\chi$ EFT	NSC97f	ESC16
$n_\sigma$ ( $k^* < 150$ MeV/c)	0.2–0.8	0.2–0.9	0.3–1.0	0.2–0.6	0.1–0.5

- Relevant for dense neutron matter is the interaction with neutrons and the interaction of  $\Sigma^{+,-}$ !
  - Disentangle the different isospin contributions
- **Larger statistics in Run3 and Run4 will definitely increase the precision and constraints on the  $\Sigma$ -N interaction**

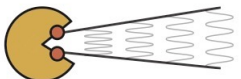
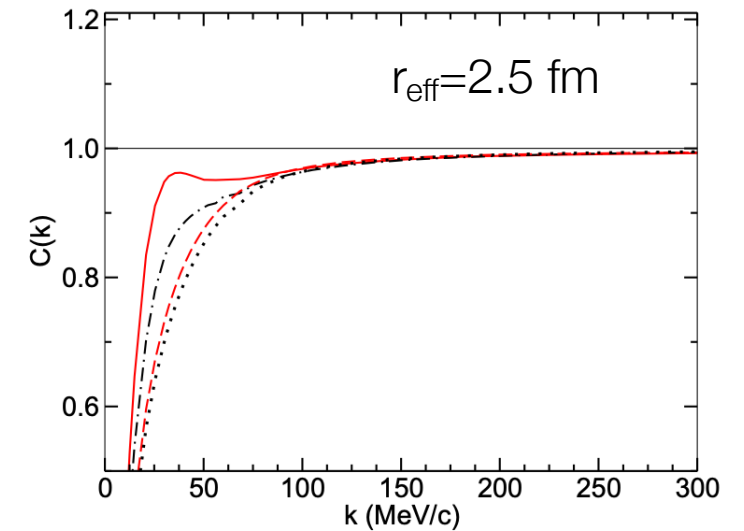
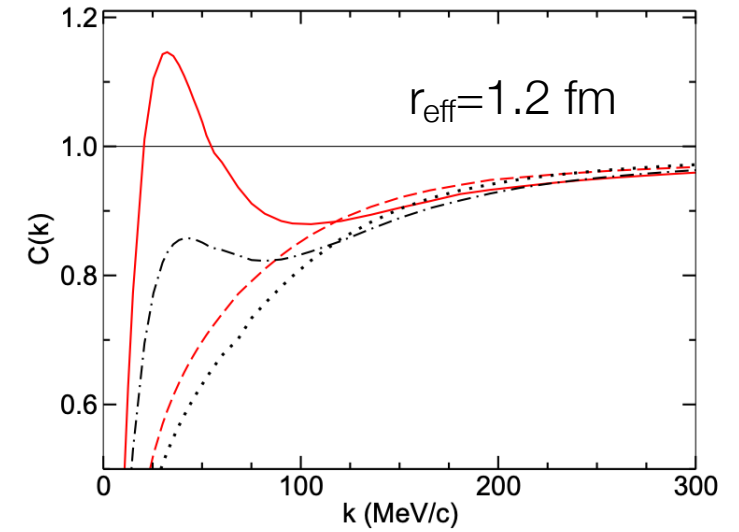
ALICE Coll. PLB 805 (2020) 135419





interaction	$a_s^c$	$r_s^c$	$a_t^c$	$r_t^c$
NLO13(600) [30]	-3.56	3.54	0.49	-5.08
NLO19(600) [31]	-3.62	3.50	0.47	-5.77
Jülich '04 [32]	-3.60	3.24	0.31	-12.2
ESC16 [33]	-4.30	3.25	0.57	-3.11
Nagels '73 [39]	$-2.42 \pm 0.30$	$3.41 \pm 0.30$	0.71	-0.78
fss2 [34]	-2.28	4.68	0.83	-1.52
NLO(sim)	-2.39	4.61	0.80	-1.25

*J.Haidenbauer, U.G. Meißner Phys.Lett.B 829 (2022)*



# TUM The projector method

- Evaluate two-body contribution in three-body correlation function
  - With no genuine three-body correlation

$$C_3(Q_3) = C_3^{12}(Q_3) + C_3^{23}(Q_3) + C_3^{31}(Q_3) - 2$$

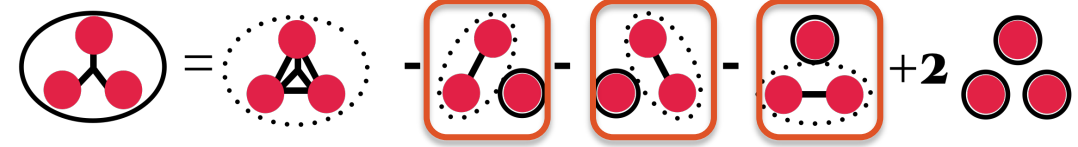
two-body CF      projector

$$C_3^{ij}(Q_3) = \int C_2(k_1^{ij}) W^{ij}(k_1^{ij}, Q_3) dk_1^{ij}$$

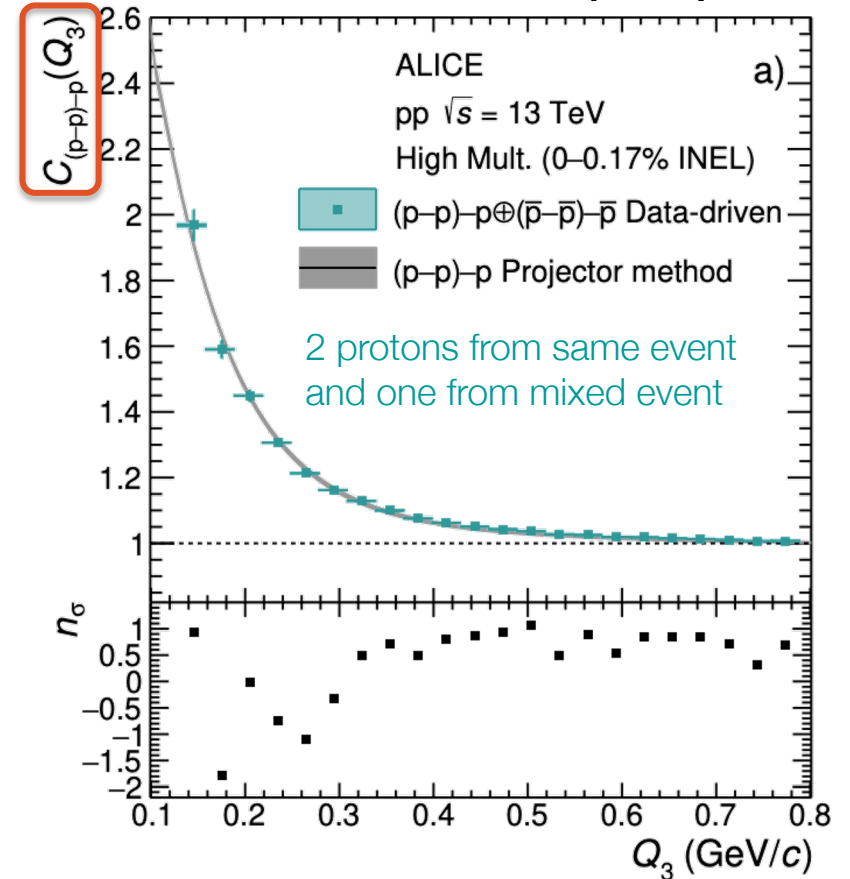
- Measured 2-body correlation functions
- same emitting source for 2-body and 3-body correlations
- Jacobian from 2-body to 3-body coordinates, integration on phase-space

- Subtract 2-body contributions with negligible uncertainties

$$Q_3 = \sqrt{q_{12}^2 + q_{22}^2 + q_{31}^2}$$



ALICE Coll. arXiv:2206.03344 [nucl-ex]



R. Del Grande, V.M.S. et al., EPJC 82 (2022)