



Latest developments of the PHQMD model

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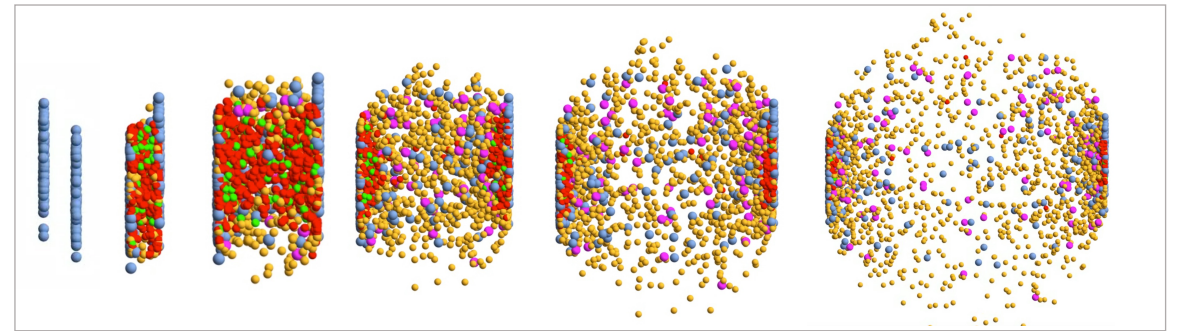
4th EMMI workshop on anti-matter, hyper-matter and exotica production at the LHC
Bologna, Italy, February 13-17, 2023

Agenda

PHQMD model & dynamic cluster formation

PHQMD results for AGS, SIS & RHIC energies

New development: Stabilisation of clusters



Cluster and hyper-cluster production in HICs with PHQMD

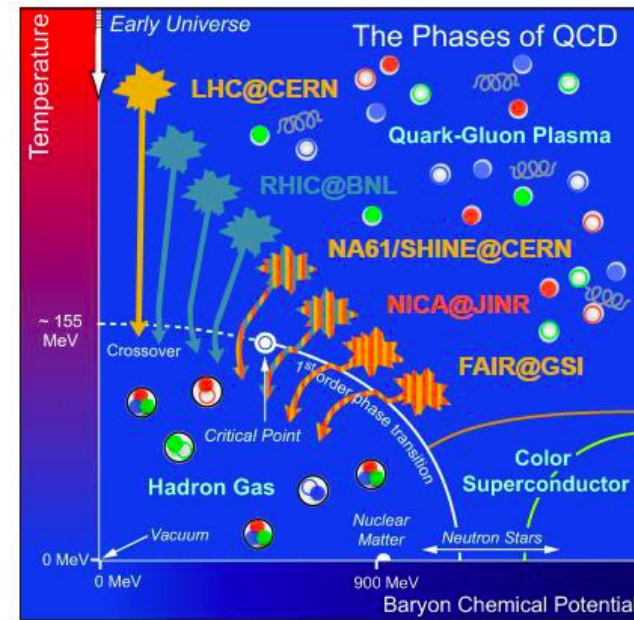
Motivation

Exploring the QCD-phase-diagram with clusters as experimental observables.

Understanding the production of clusters in relativistic heavy-ion collisions:

How can **weakly bound** clusters survive in the **hot and dense** environment of a HICs?

'Ice in fire puzzle'



Challenge

Modeling the **time evolution** of cluster formation and the origin of their production.

Modelling of cluster formation in HIC

Statistical models

- Production of nuclei depending on T and μ_B at chemical freeze-out & particle mass

Coalescence models

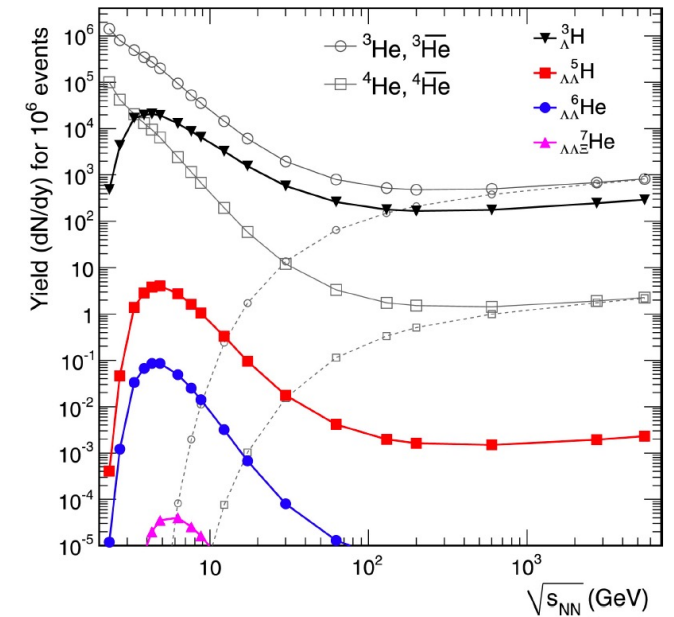
- Formation of nuclei by nucleons & hyperons that are close in coordinate and momentum spaces at freeze-out time

Hybrid models for cluster production

- sudden transition from a dynamical model to clusterisation
e.g . UrQMD + afterburner

=> no dynamical cluster formation during time evolution

=> no information on the dynamics of clusters formation & microscopic origin



A. Andronic et al., Phys. Lett. B697 (2011) 203-207.

Modeling of dynamic cluster formation in HIC

Dynamical modeling of cluster formation based on interactions.

Interaction mechanism in transport models:

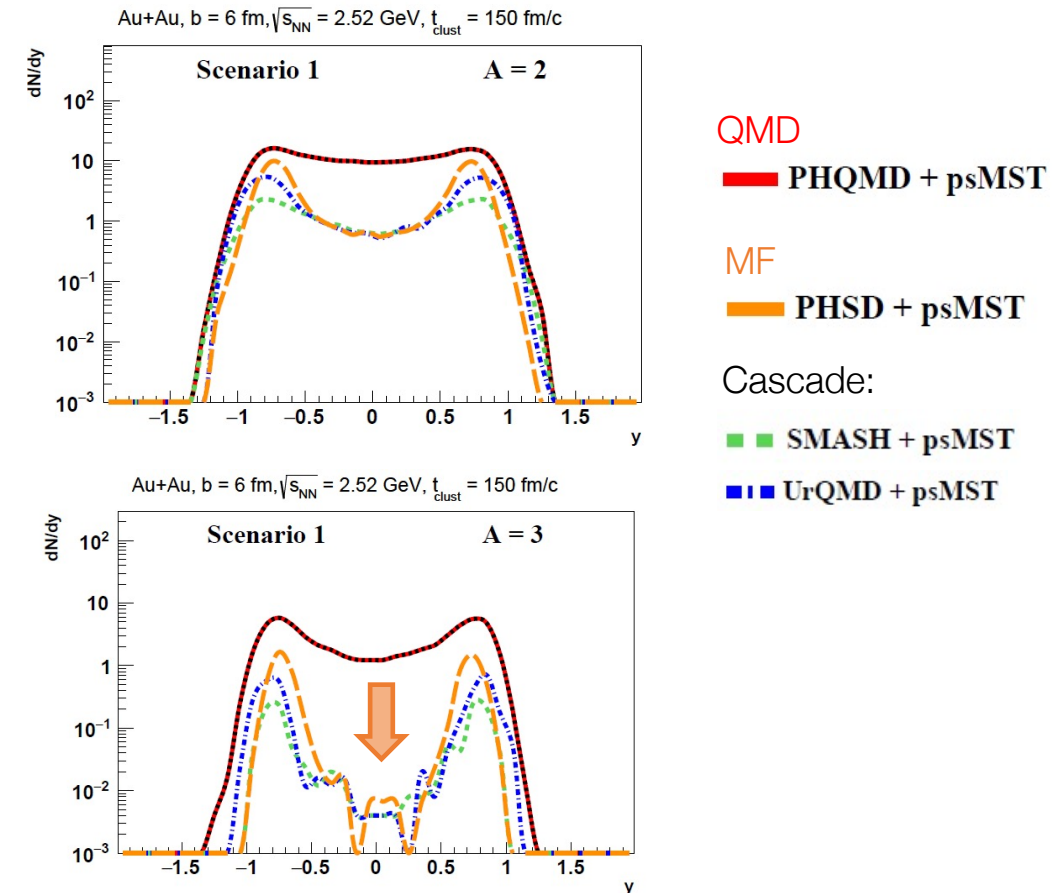
- via potential interaction - potential mechanism this talk
- by scattering - kinetic mechanism

=> Cluster formation is sensitive to nucleon dynamics.

Modeling of nucleon-nucleon potential interactions:

- QMD (quantum-molecular dynamics) – allows to keep correlations
- MF (mean-field based models) – correlations are smeared out
- Cascade – correlations by potential interactions missing

Cluster at final time

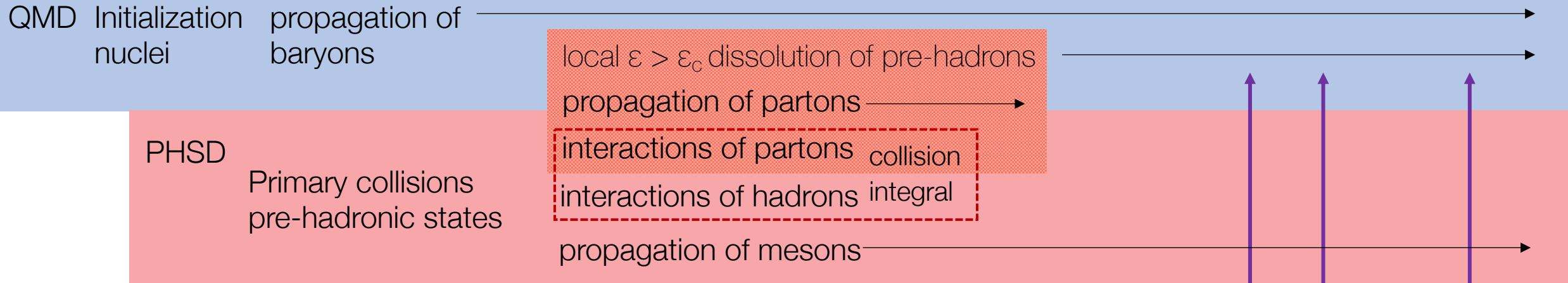
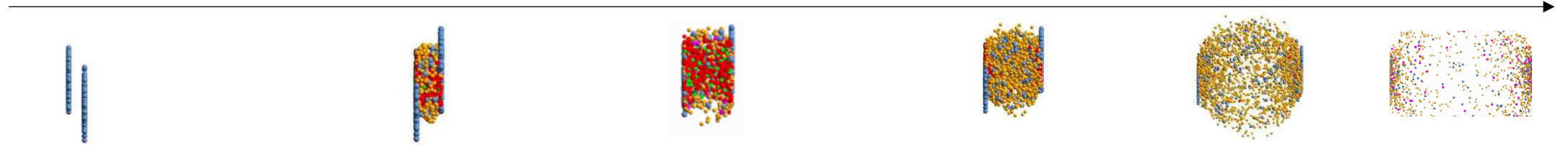


Parton-Hadron-Quantum-Molecular Dynamics

= n-body microscopic transport approach for the description of heavy-ion dynamics with dynamical cluster formation from low to ultra-relativistic energies

Relativistic considerations + Correlations between nucleons + Cluster recognition

Initial A+A collisions Formation of QGP Partonic phase Hadronization Hadronic phase



QMD propagation

= n-body transport approach

Generalized Ritz variational principle $\delta \int_{t_1}^{t_2} dt \langle \psi(t) | i \frac{d}{dt} - H | \psi(t) \rangle = 0$.

Ansatz: trial wave function for one particle “i” : Gaussian with width L centered at $\mathbf{r}_{i0}, \mathbf{p}_{i0}$ Aichelin Phys. Rept. 202 (1991)

$$\psi(\mathbf{r}_i, \mathbf{r}_{i0}, \mathbf{p}_{i0}, t) = C e^{-\frac{1}{4L} \left(\mathbf{r}_i - \mathbf{r}_{i0}(t) - \frac{\mathbf{p}_{i0}(t)}{m} t \right)^2} \cdot e^{i \mathbf{p}_{i0}(t) (\mathbf{r}_i - \mathbf{r}_{i0}(t))} \cdot e^{-i \frac{\mathbf{p}_{i0}^2(t)}{2m} t} \quad L=4.33 \text{ fm}^2$$

Assume that $\psi(t) = \prod_{i=1}^N \psi(\mathbf{r}_i, \mathbf{r}_{i0}, \mathbf{p}_{i0}, t)$ for N particles (neglecting antisymmetrization!)

Equations-of-motion (EoM) for Gaussian centers in coordinate and momentum space:

$$\dot{r}_{i0} = \frac{\partial \langle H \rangle}{\partial p_{i0}} \quad \dot{p}_{i0} = -\frac{\partial \langle H \rangle}{\partial r_{i0}}$$

QMD potentials and EoS

expectation value of Hamiltonian

$$\langle H \rangle = \sum_i \langle H_i \rangle = \sum_i (\langle T_i \rangle + \sum_{j \neq i} \langle V_{i,j} \rangle) = \sum_i (\sqrt{p_{i0}^2 + m^2} - m) + \sum_i \langle V_{Skyrme}(\mathbf{r}_{i0}, t) + V_{coul}(\mathbf{r}_{i0}, t) \rangle$$

Skyrme interaction ('static'): Effective density dependent nucleon-nucleon interactions

$$\langle V_{Skyrme}(\mathbf{r}_{i0}, t) \rangle = \alpha \left(\frac{\rho_{int}(\mathbf{r}_{i0}, t)}{\rho_0} \right) + \beta \left(\frac{\rho_{int}(\mathbf{r}_{i0}, t)}{\rho_0} \right)^\gamma$$

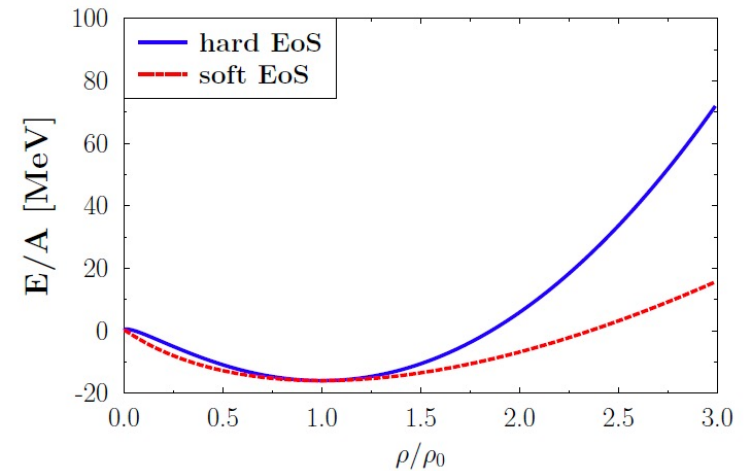
Interaction density (with relativistic extension):

$$\rho_{int}(\mathbf{r}_{i0}, t) \rightarrow C \sum_j \left(\frac{4}{\pi L} \right)^{3/2} e^{-\frac{4}{L}(\mathbf{r}_{i0}^T(t) - \mathbf{r}_{j0}^T(t))^2} e^{-\frac{4\gamma_c^2 m}{L}(\mathbf{r}_{i0}^L(t) - \mathbf{r}_{j0}^L(t))^2}$$

Parameter of the nuclear equation of state in PHQMD

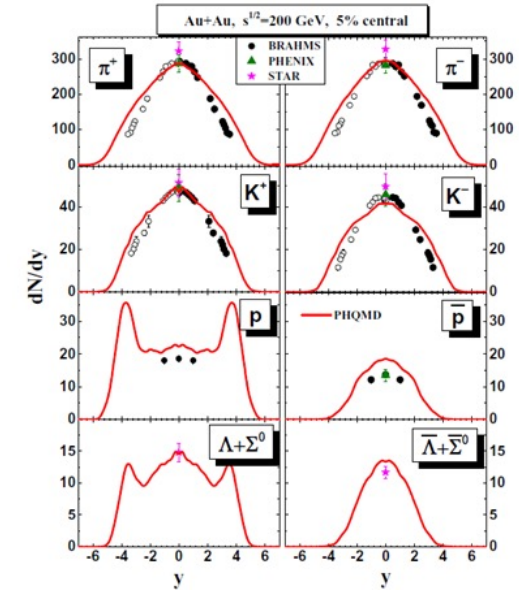
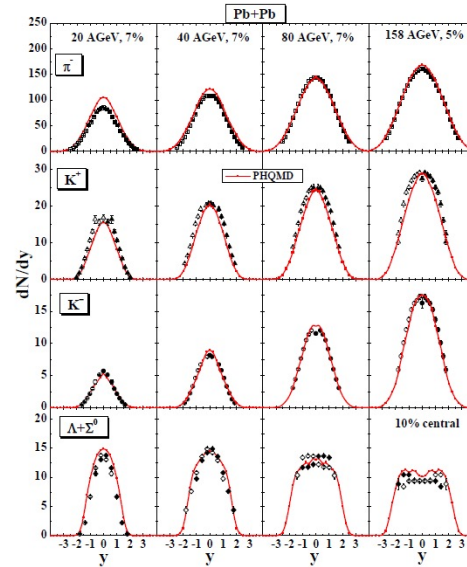
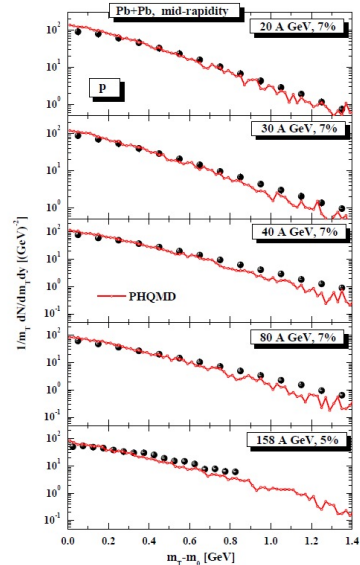
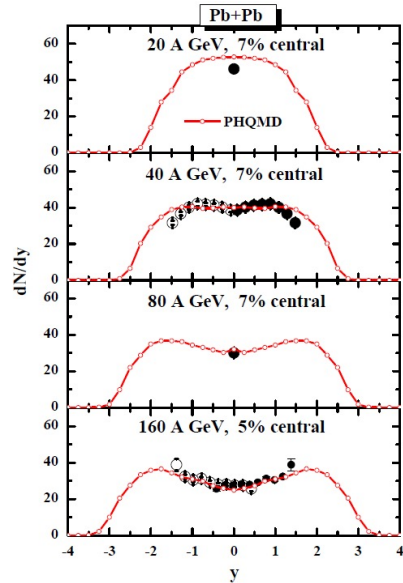
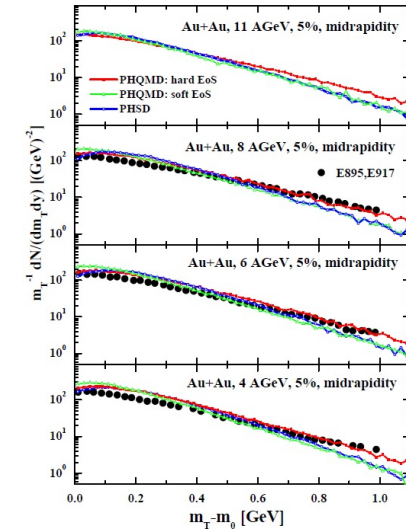
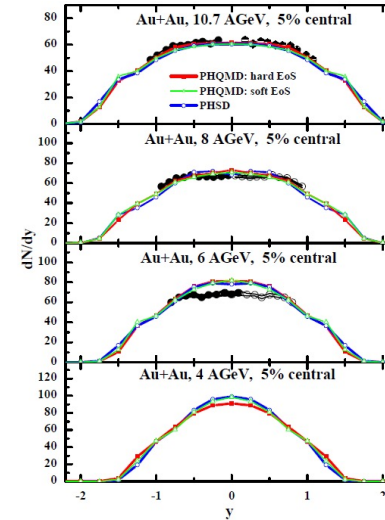
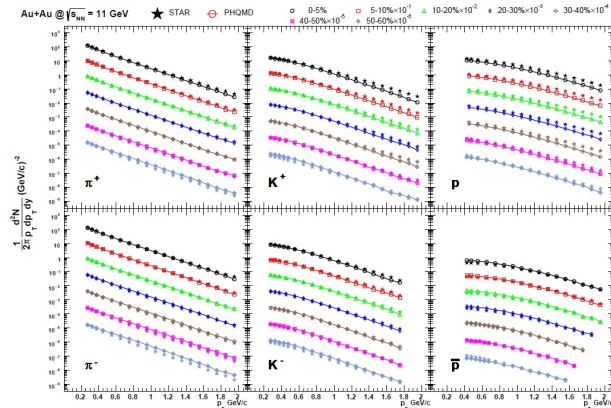
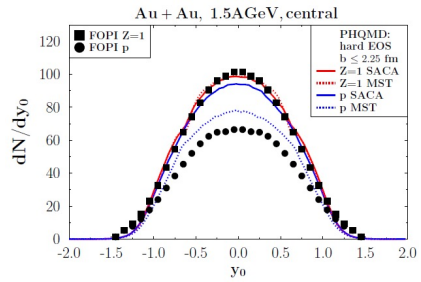
	α [MeV]	β [MeV]	γ	K [MeV]	compression modulus K
S	-390	320	1.14	200	$K = -V \frac{dP}{dV} = 9\rho^2 \frac{\partial^2(E/A(\rho))}{(\partial\rho)^2} \Big _{\rho=\rho_0}$
H	-130	59	2.09	380	

EoS for infinite matter at rest



*Work in progress: implementation of momentum-dependent potential (M. Winn)

Highlights: PHQMD ,bulk' dynamics from SIS to RHIC



PHQMD: J. Aichelin et al., PRC 101 (2020) 044905

=> PHQMD provides a good description of hadronic 'bulk' observables from SIS to RHIC energies.

Minimum Spanning Tree (MST)

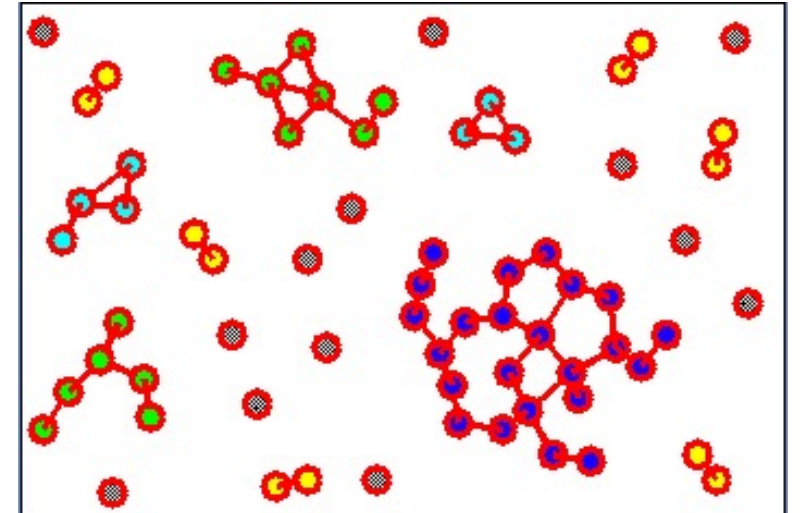
Cluster criterion: distance of nuclei

Algorithm: search for accumulations of particles in coordinate space

1. Two particles i & j are bound if:

$$|r_i - r_j| < 4.0 \text{ fm}$$

2. Particle is bound to cluster if bound with at least one particle of cluster



Remark: additional momentum cuts lead to a small changes: particles with large relative momentum are mostly not at the same position (V. Kireyeu, Phys.Rev.C 103 (2021) 5)

Cluster stability over time

Semiclassical model QMD:

Clusters not described as ‘quantum objects’

with minimal average kinetic energy for nucleons

=> cluster-nucleon can have enough kinetic energy to escape cluster

=> “bound” cluster at time t can spontaneously dissolve at $t + \Delta t$

Skyrme potential is not relativistic:

calculation of cluster binding energy in cluster-system after Lorentz-boost

=> baryons have different times in cluster-frame

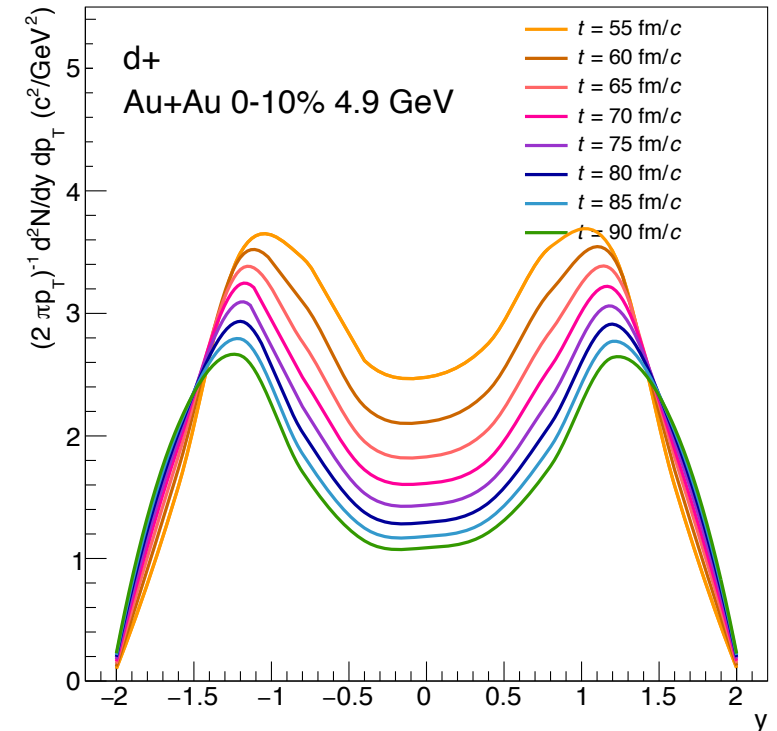
=> sign of binding energy can change

=> cluster multiplicities decrease over time

=> time selection needed for every nuclei specie for correct multiplicities

=> form of y -, p_T distribution & ratio of particles don't change over time

New: Stabilisation procedure aMST



Relativistic fragment formation time for spectra

Consideration of time dilatation for the cluster identification:

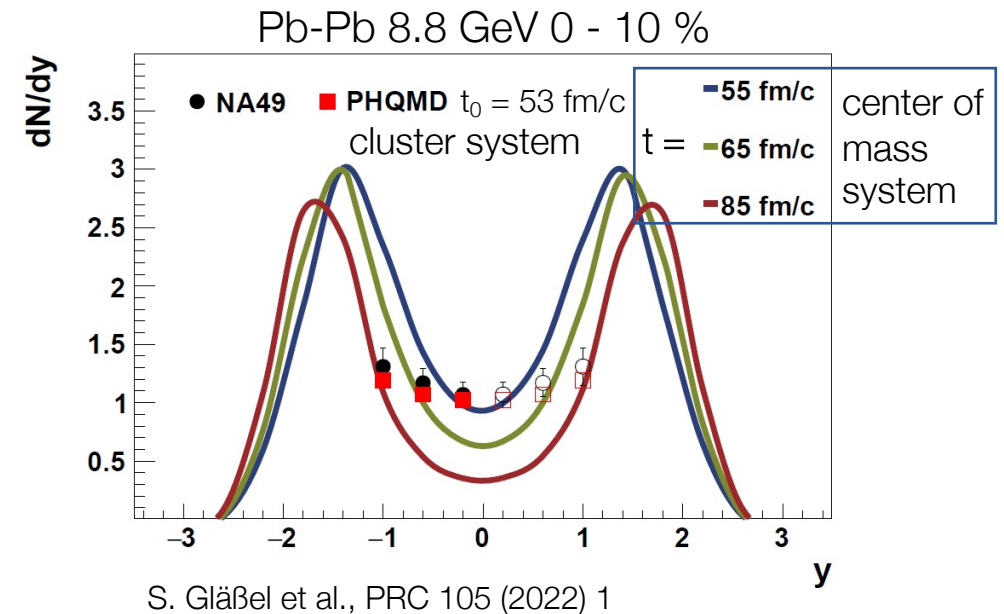
time t_0 in cluster rest-system is delayed vs. time t in the center of mass system depending on y

=> create spectra at the same time in the cluster rest-system t_0

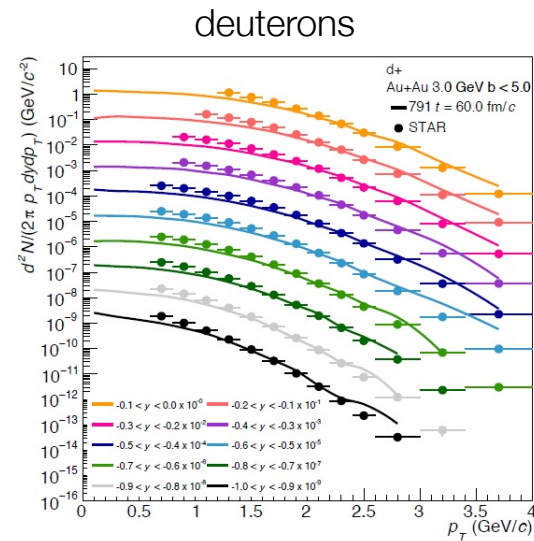
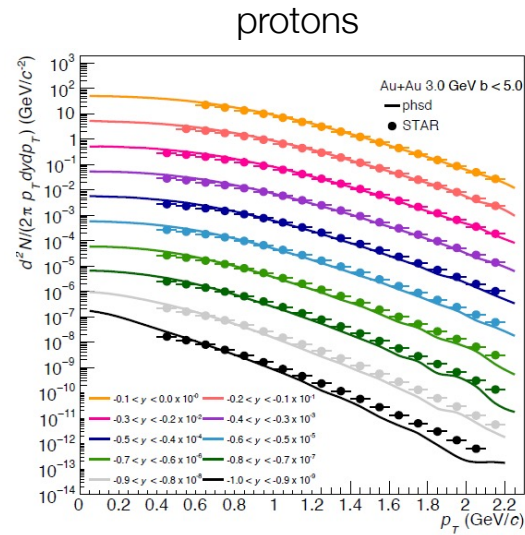
=> transformation: $t = t_0 \cosh(y)$

center of mass system cluster system

↑ ↑

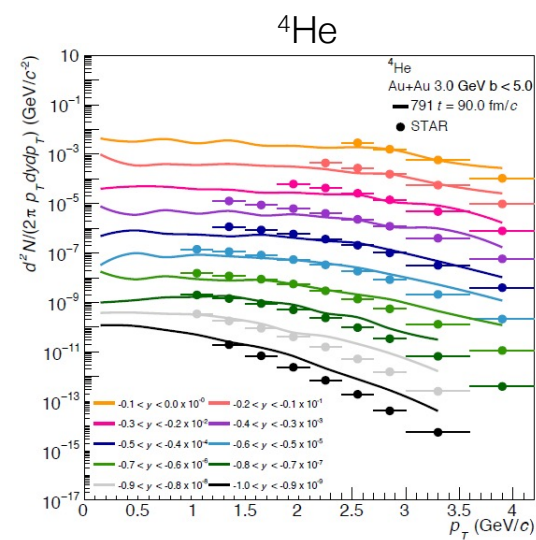
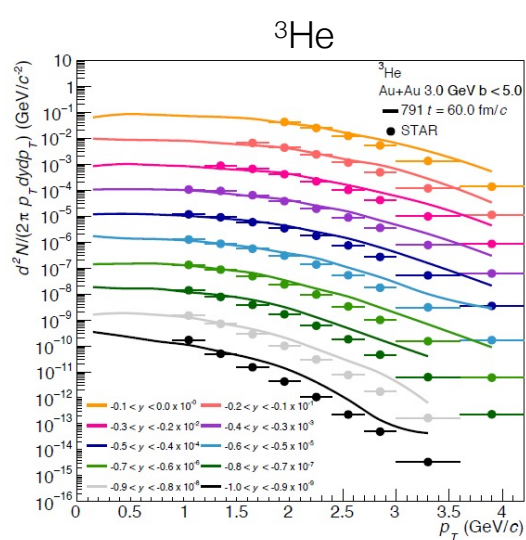
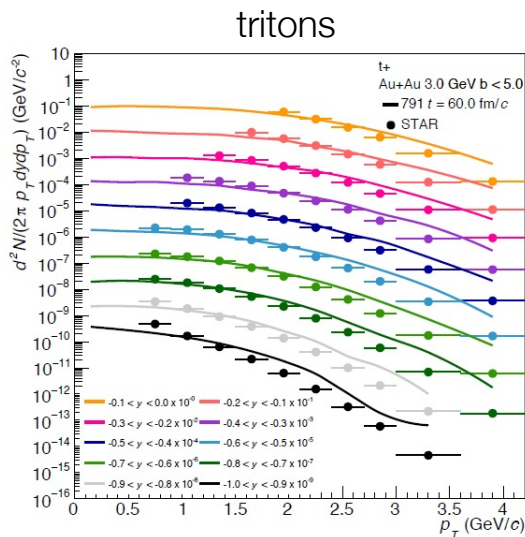


Light cluster production at $\sqrt{s}_{NN} = 3$ GeV

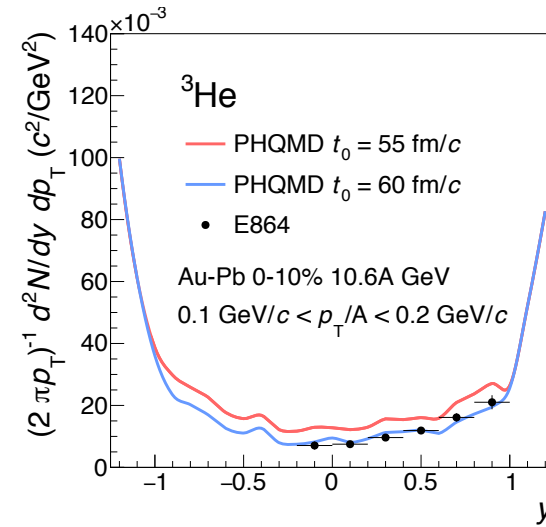
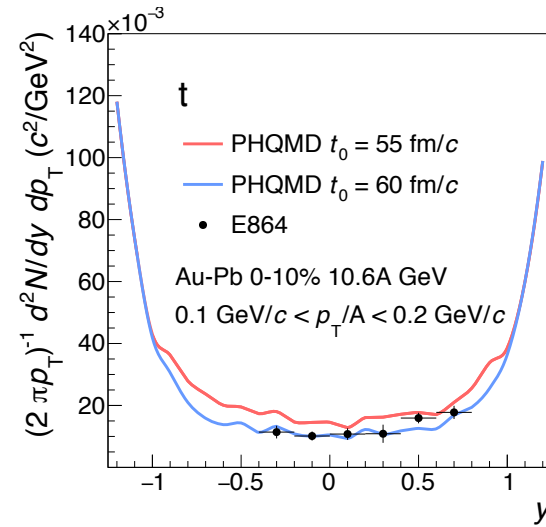
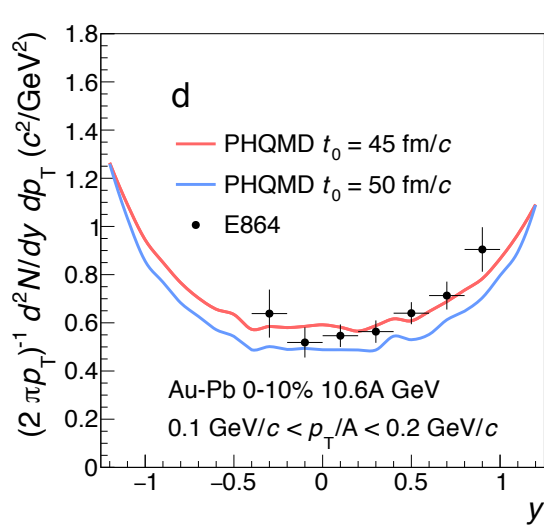


=> PHQMD reproduces STAR* p_T spectra for protons, deuterons, tritons ^3He and ^4He .

* (preliminary) STAR data – talk by Hui Liu at QM'2022

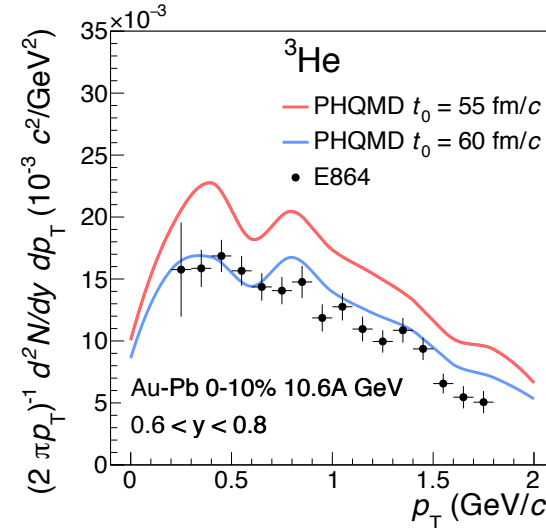
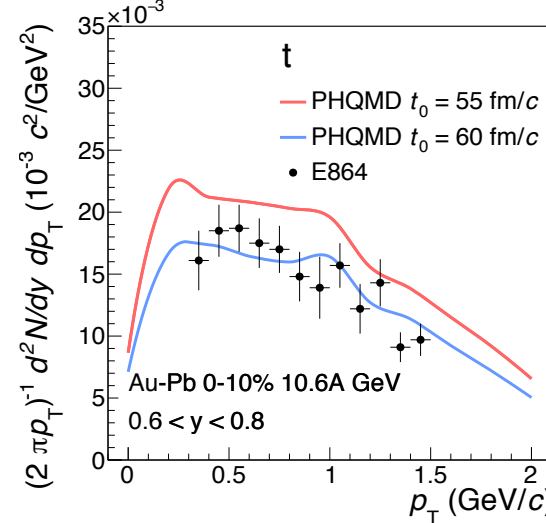
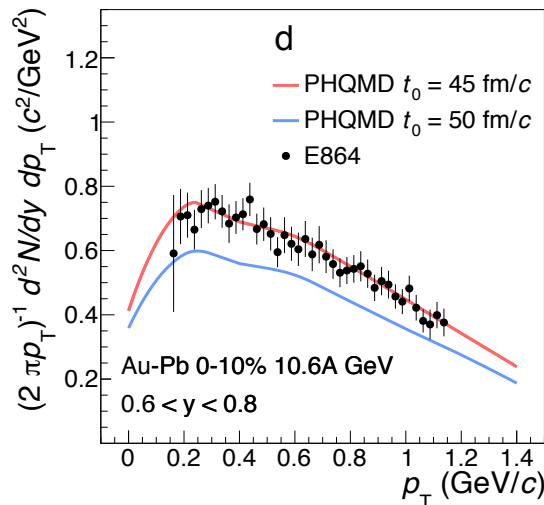


Light cluster production at AGS energies



$$t = t_0 \cosh(y)$$

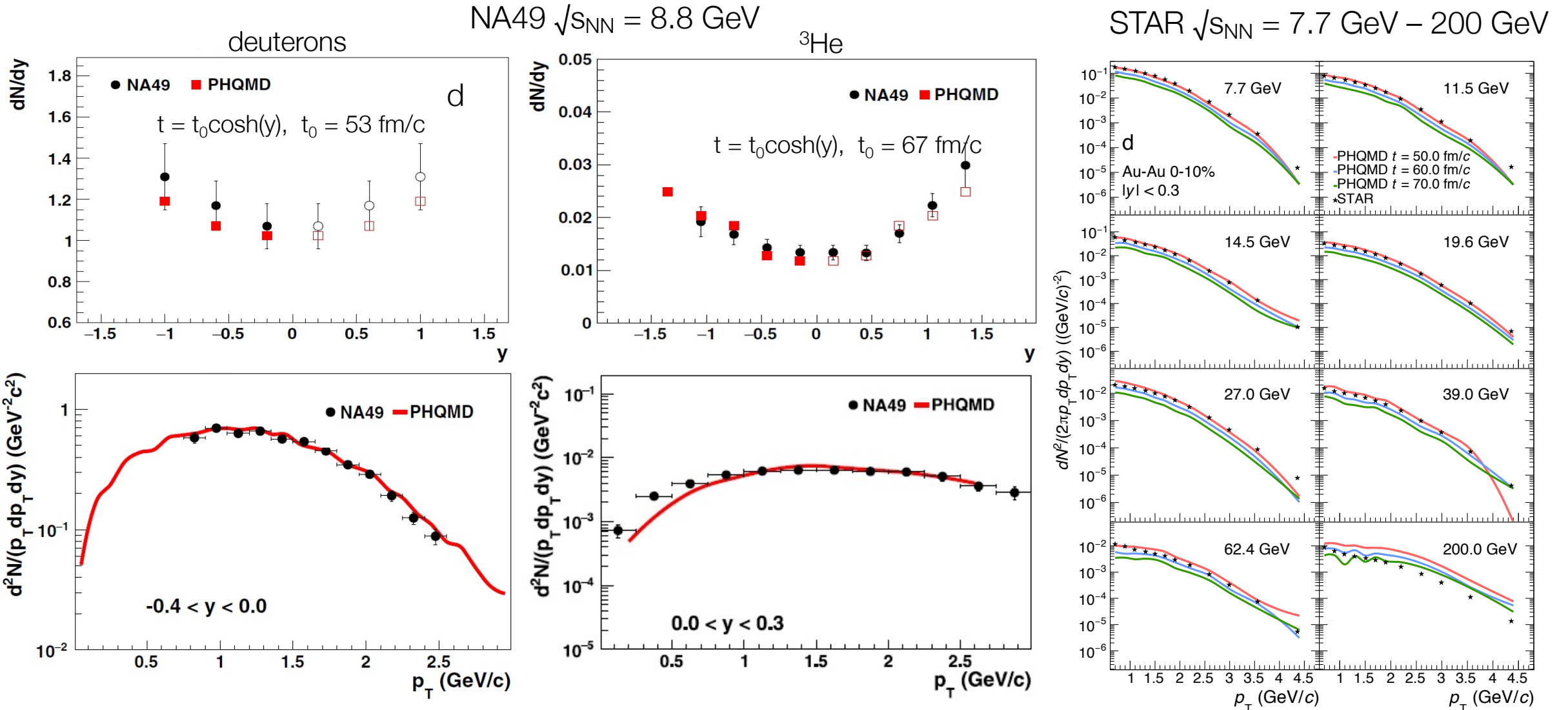
=> y - & p_T -spectra for deuterons, tritons and ^3He from E864* at $\sqrt{s_{\text{NN}}} = 4.9$ GeV are nicely described by PHQMD.



S. G. et al., PRC 105 (2022) 1

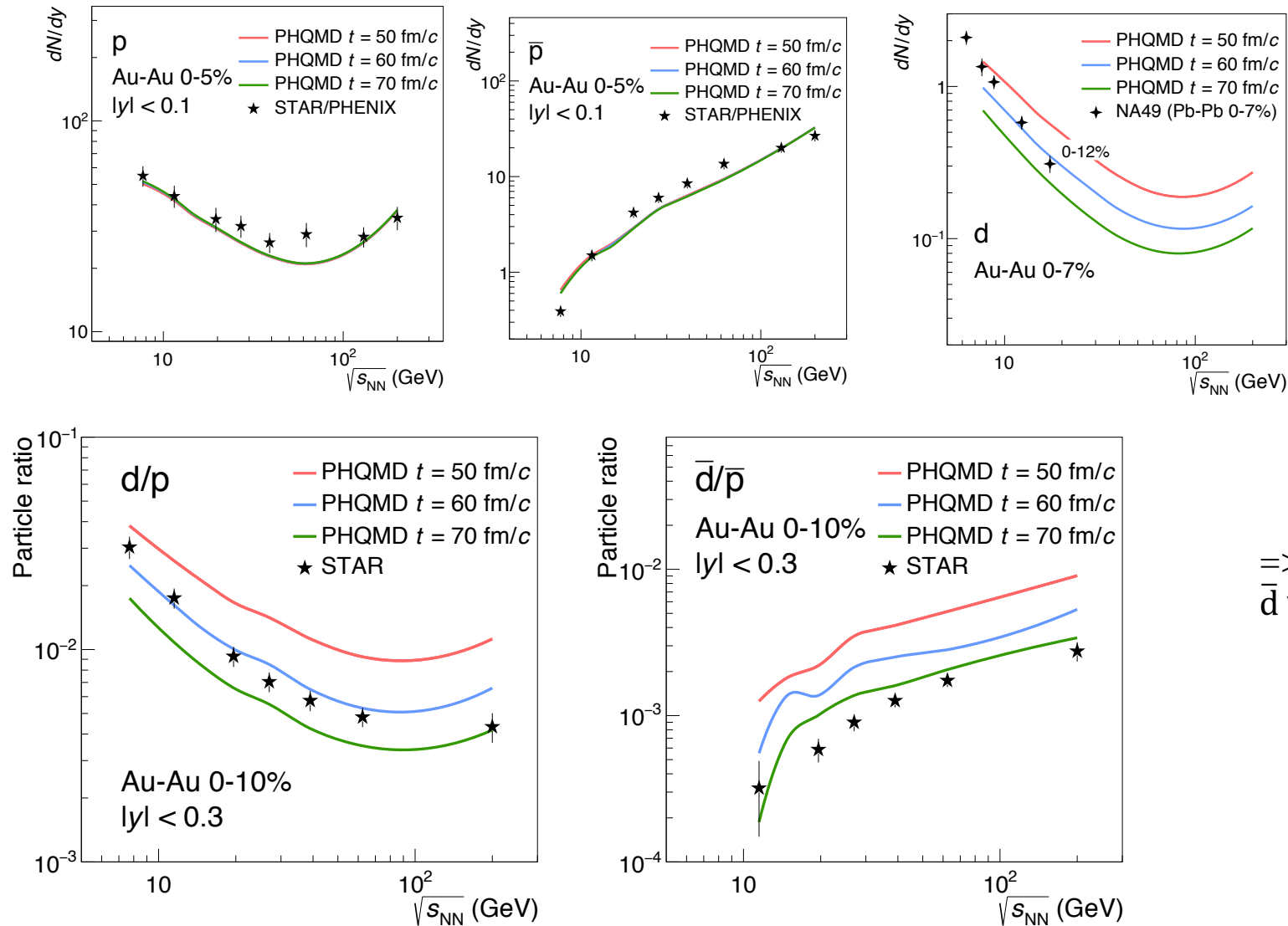
*T. A. Armstrong et al: „Measurements of light nuclei production in pbeam=11.5A GeV/c Au+Pb heavy-ion collisions“, PHYSICAL REVIEW C, VOLUME 61, 064908 (2002)

Cluster production at SPS & RHIC energies



=> The PHQMD results for d and ^3He agree with NA49 and STAR data.

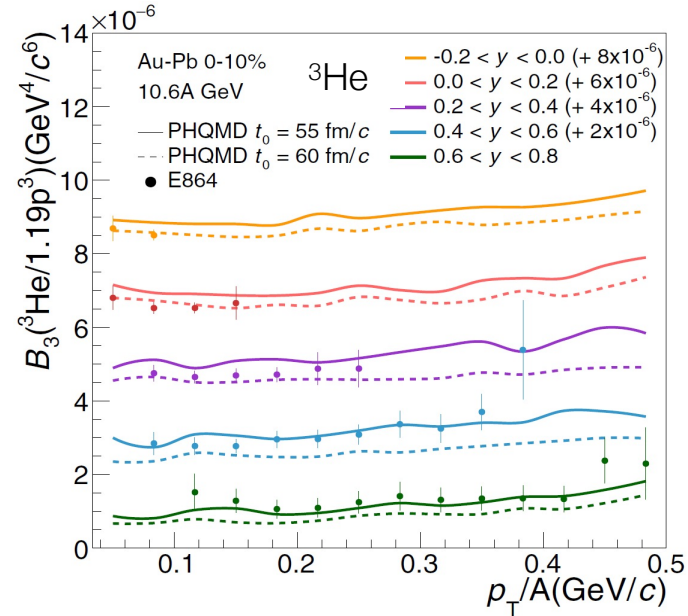
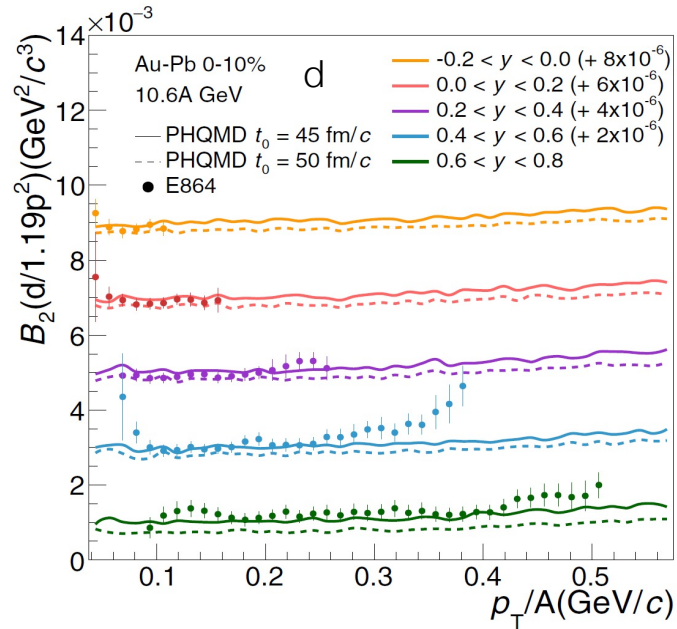
Excitation function of multiplicity of p , \bar{p} , d & p/d -ratio



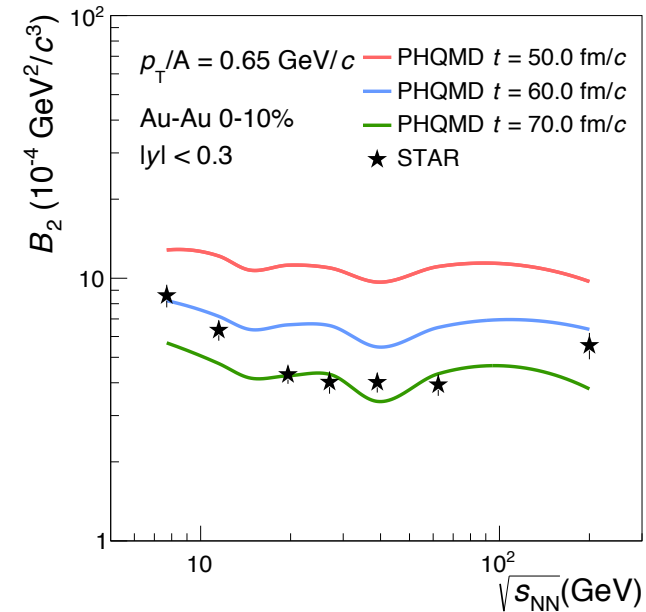
=> The p , \bar{p} yields at $y \sim 0$ are stable, the d , \bar{d} yields are best described at $t = 60-70$ fm/c.

Coalescence parameter for d and ^3He

E864 $\sqrt{s_{\text{NN}}} = 4.9 \text{ GeV}$



STAR $\sqrt{s_{\text{NN}}} = 7.7 \text{ GeV} - 200 \text{ GeV}$



Coalescence parameter:

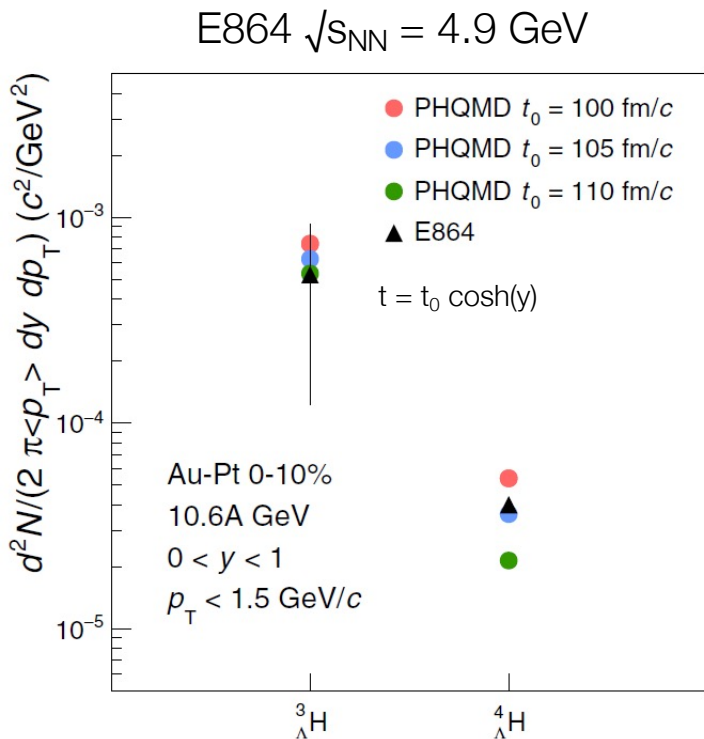
$$B_2 = \frac{E_d \frac{d^3 N_d}{d^3 P_d}}{\left(E_p \frac{d^3 N_p}{d^3 P_p} \Big|_{p_p = P_d/2} \right)^2}$$

=> Good description of B_A for energies between 5 – 200 GeV.

=> No remarkable structure for B_A as a function of p_T , only slight increase.

=> Probability that baryons with p_T/A form a cluster with size A more or less independent of p_T , only increases slightly.

Hypernuclei production at $\sqrt{s_{NN}} = 3.0$ and 4.9 GeV



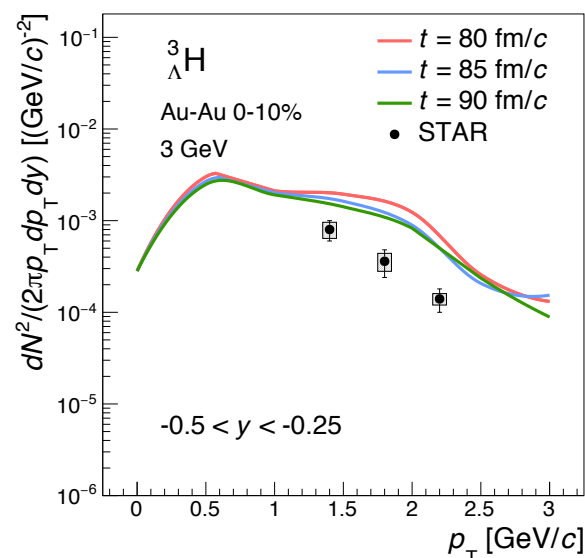
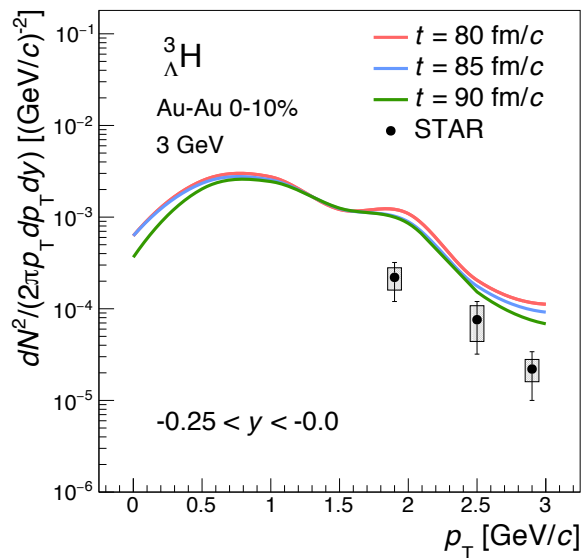
Assumption for nucleon-hyperon potential: $V_{NL} = 2/3 V_{NN}$

=> trend of the experimental STAR*

${}^3_{\Lambda}H$ & ${}^4_{\Lambda}H$ p_T -spectra at $\sqrt{s_{NN}}=3$ is produced well

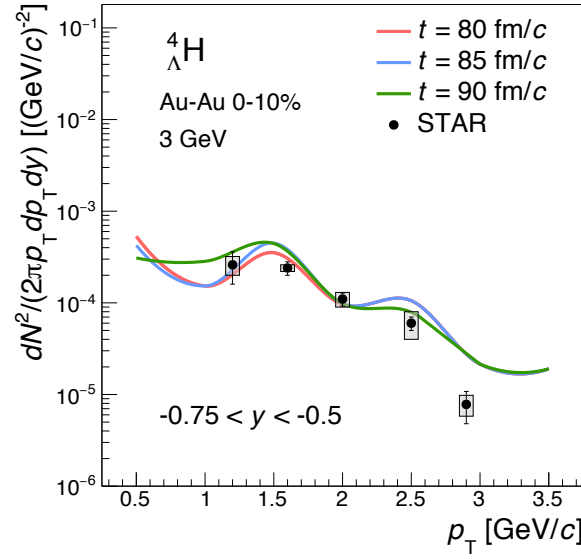
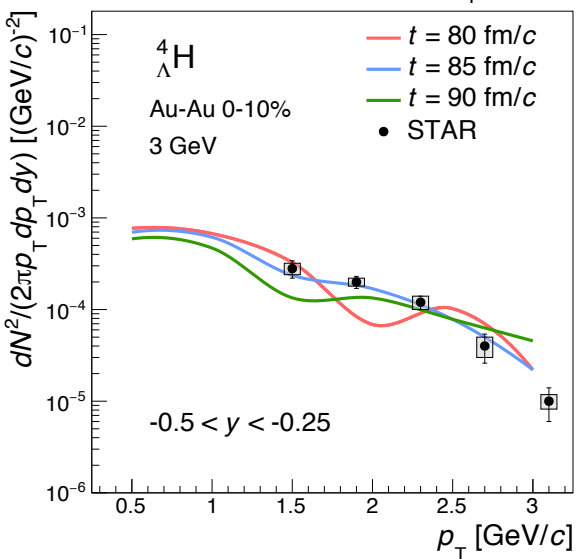
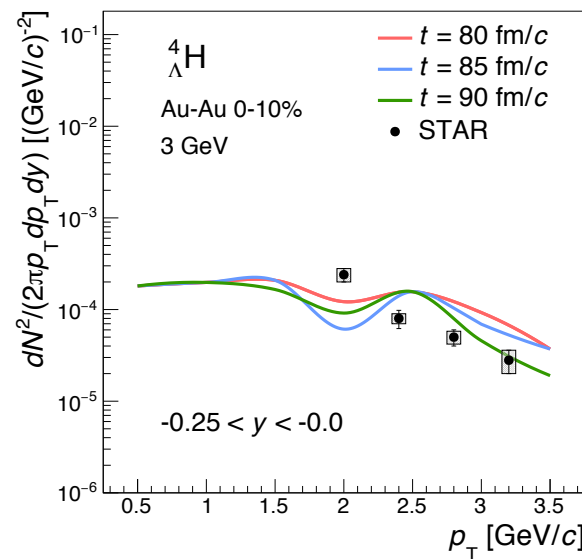
=> yields are slightly overpredicted

STAR $\sqrt{s_{NN}} = 3.0$ GeV



=> Reasonable description of hypernuclei production at $\sqrt{s_{NN}} = 3.0$ GeV

*Yue-Hang Leung: First results of H3L & H4L (dN/dy , c_T , v_1) from 3 GeV Au+Au collisions with the STAR detector (CPOD2021)



Stabilisation procedure Advanced MST

Clusters are not stable over time

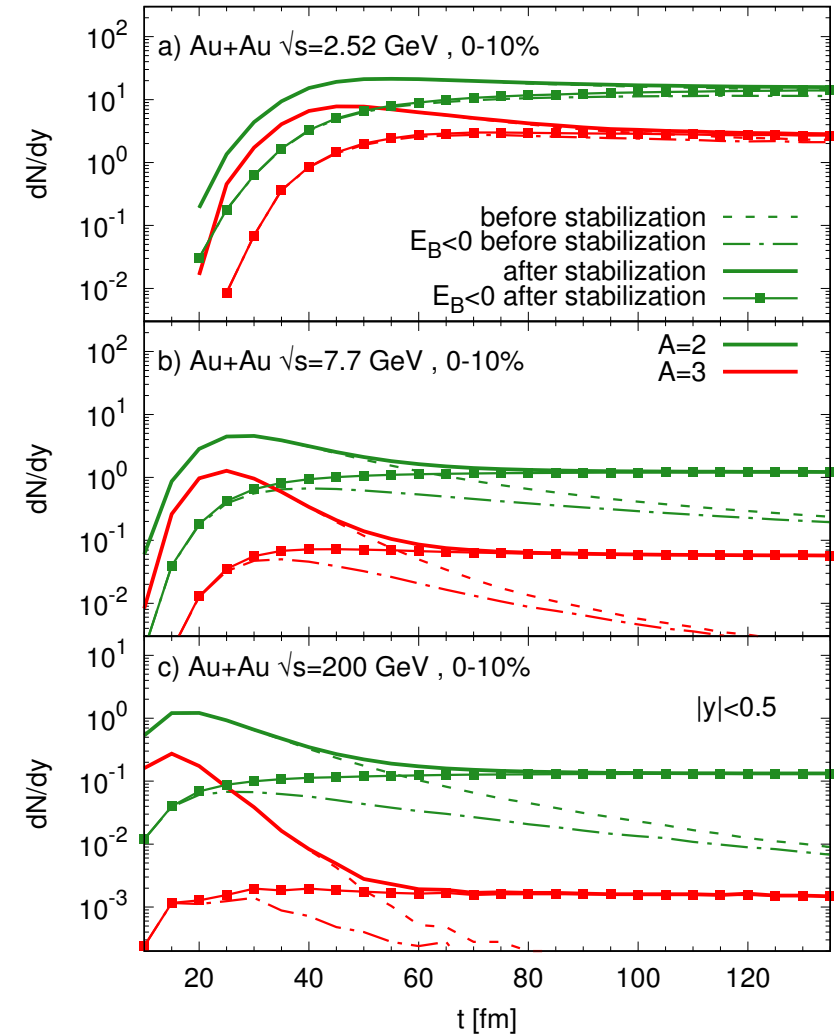
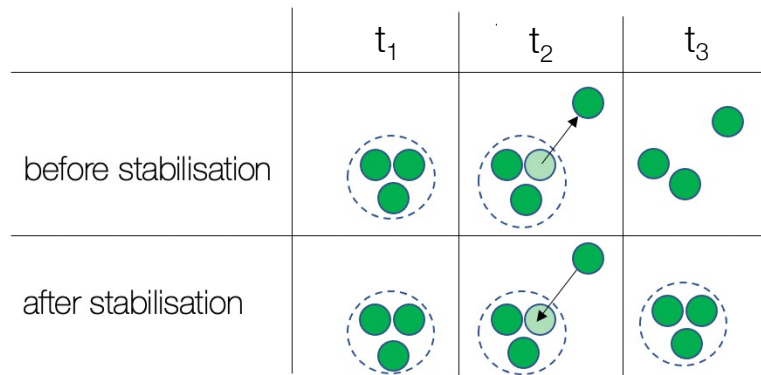
- 1) Clusters not described as 'quantum objects' : cluster can dissolve
- 2) Skyrme potential is not relativistic: binding energy can change

Advanced MST

- Consider asymptotic state: clusters and free nucleons
- Track the freezeout-time for each baryon = time of the last collision

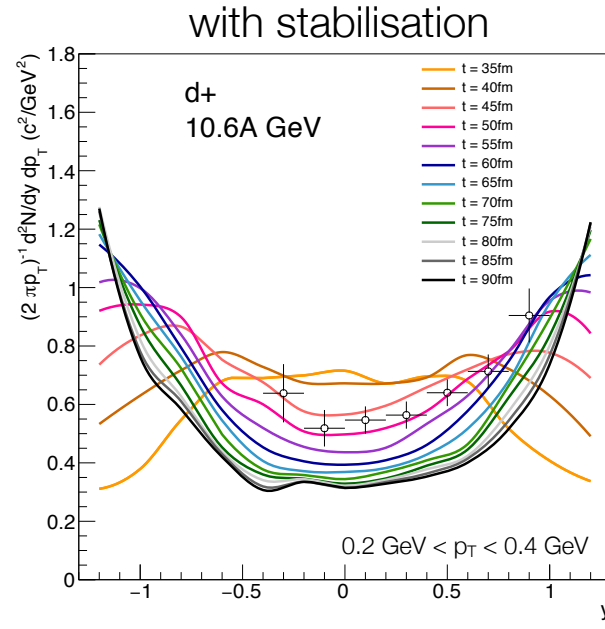
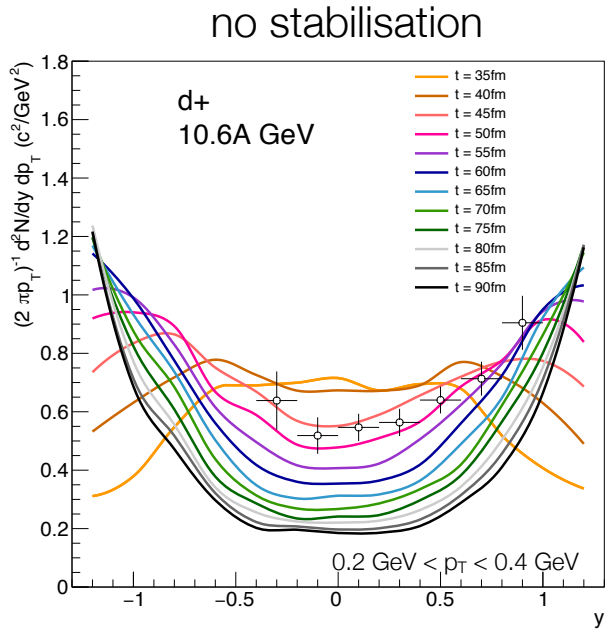
- a) If time of cluster disintegration > baryon freeze-out time (and if $E_B < 0$) :
=> restore cluster
- b) If sign of binding energy changes after baryon freeze-out time:
=> restore E_B

+ extra condition: $E_B < 0$ negative binding energy for identified clusters



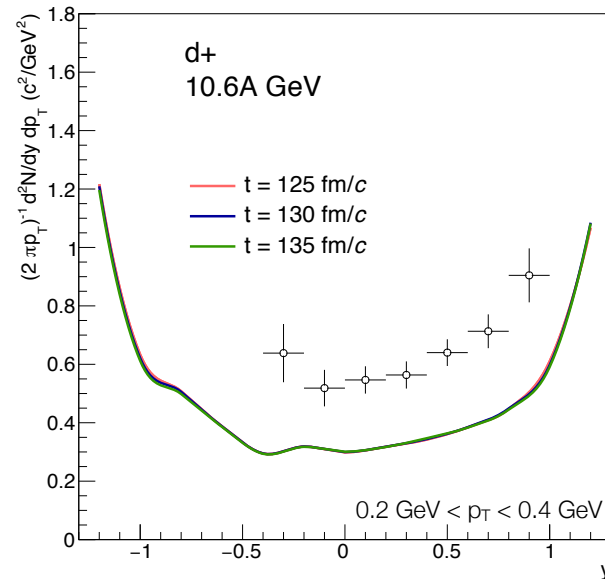
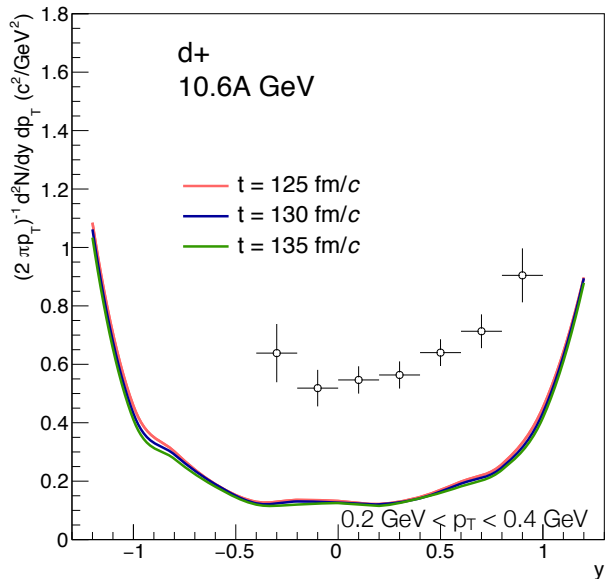
dN/dy time evolution for deuteron at $\sqrt{s_{NN}} = 4.9$ GeV

35 fm < t < 90 fm



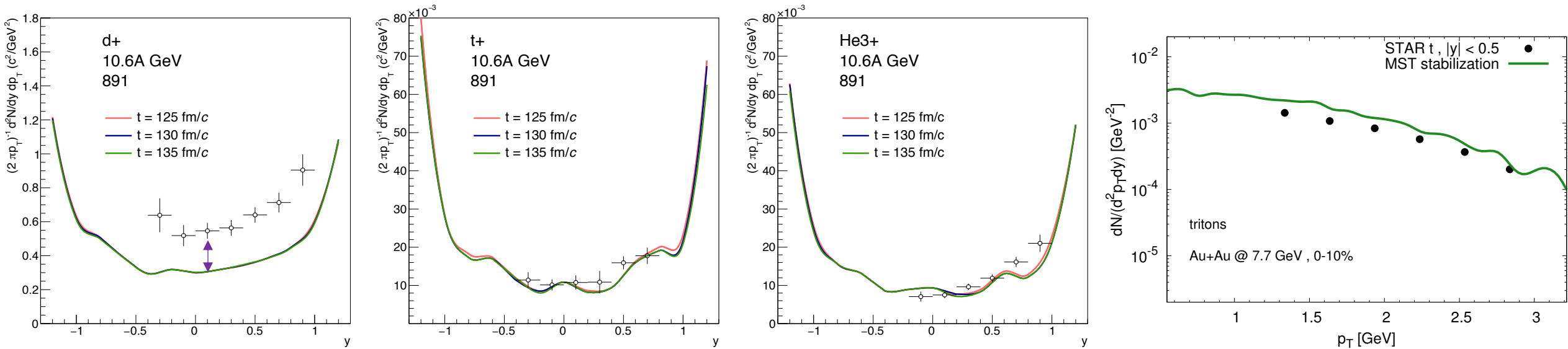
=> With stabilisation,
multiplicity starts to drop
slower at around 60 fm/c.

125 fm < t < 135 fm



=> At the final time, deuteron
multiplicity is ~2 times higher
with stabilisation.

Stable light nuclei at $\sqrt{s_{NN}} = 4.9$ and 7.7 GeV with aMST



PHQMD with stabilisation procedure fits the experimental data at $\sqrt{s_{NN}} = 4.9$ and 7.7 GeV very well for triton and ^3He .

Deuterons are underestimated => contribution of deuterons formed by inelastic scattering.

See talk on Tues, 4.30 pm by Elena Bratkovskaya.

Summary & Conclusion

PHQMD

- is a **microscopic n-body transport approach** to describe HIC and **dynamical cluster formation**
- is a combined model: **PHSD + QMD + MST | SACA**
- transports baryons with QMD to keep **potential interactions** among baryons and allow cluster formation
- identifies with MST the clusters based on the distance between baryons

=> reproduces cluster and hypercluster data on dN/dy , dN/dp_T and ratios from AGS to top RHIC energies

- The newly developed stabilisation routine in **advanced MST** overcomes time dependency of cluster multiplicities.

Outlook

- study of hyper-nuclei with more realistic potentials
- extension to LHC energies

Talk on Tuesday, 4.30 pm by Elena Bratkovskaya:

How can weakly bound clusters survive in the hot and dense environment of a heavy ion collision?

Deuteron production based on kinetic interactions / scattering

THANK YOU FOR YOUR ATTENTION.