

# **Approaching a SU(3) Energy Density Functional**

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# Agenda

- **SU(3) symmetry and octet BB-interactions**
- **Exploiting SU(3) relations for a covariant Octet-EDF**
- **Baryon mean-fields in asymmetric nuclear matter**
- **New effect:  $\Lambda$ - $\Sigma$  mixing induced by the nuclear isovector mean-field**
- **Summary and outlook**

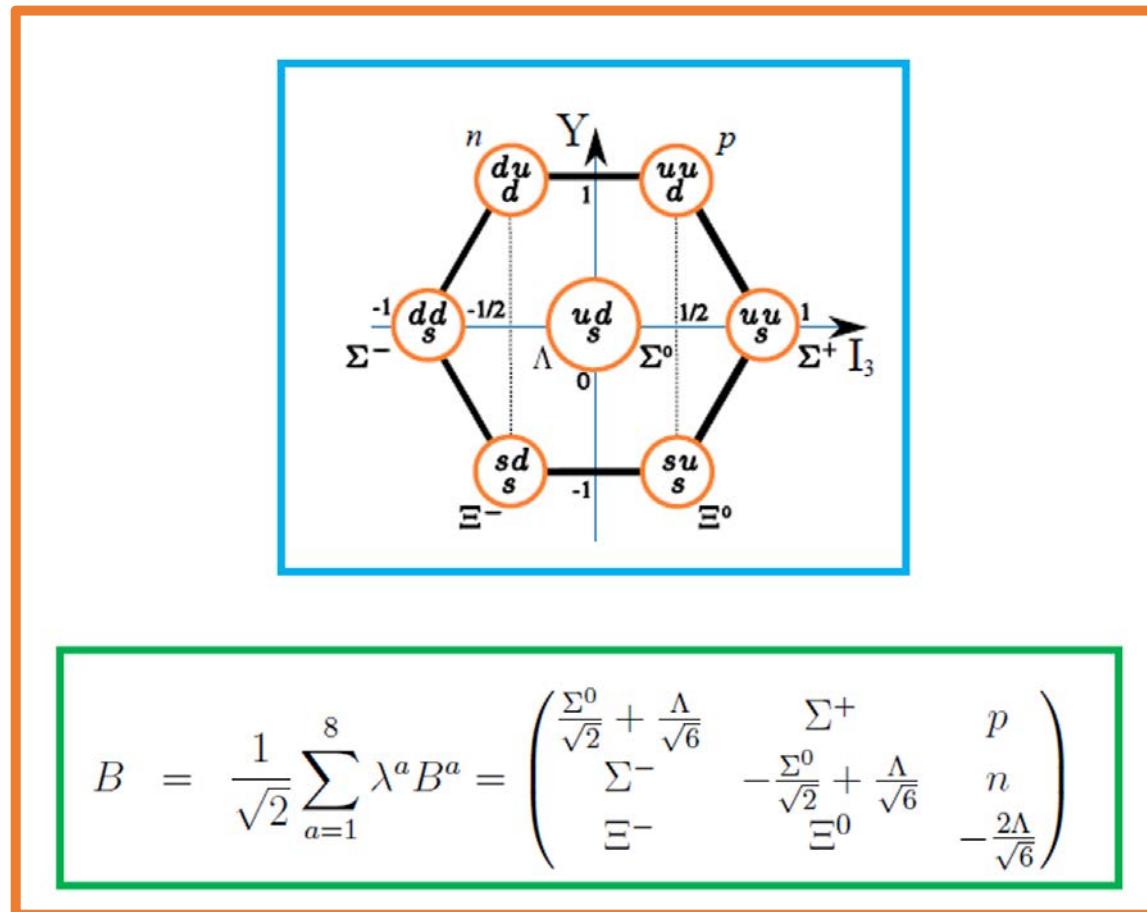
H. Lenske, M. Dhar, EPJ Web Conf. 271 (2022) 05003 (Hyp 2022, Prague), 2208.04916 [nucl-th]

Theoretical Background:

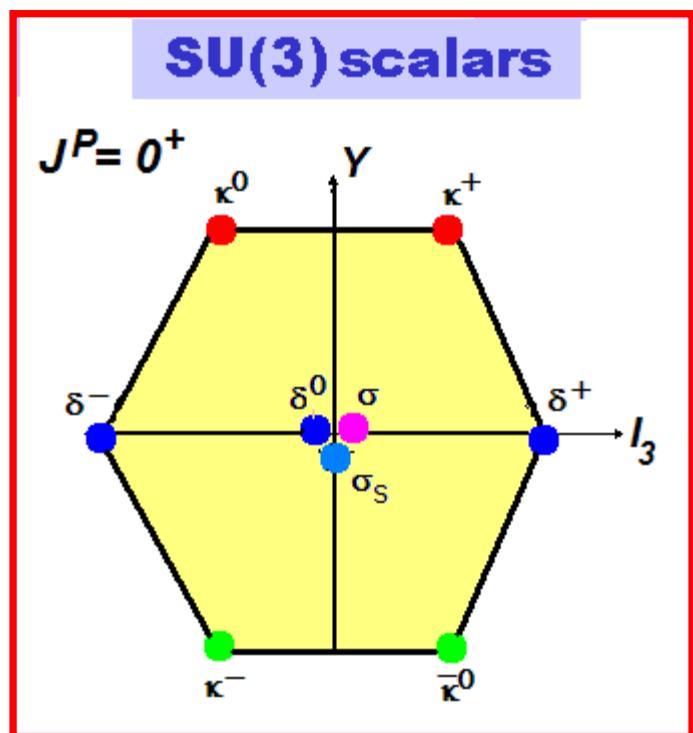
H. Lenske, M. Dhar, *Lect. Notes Phys.* 948 (2018) 161

H. Lenske, M. Dhar, Th. Gaitanos, Xu Cao, *Prog. Part. Nucl. Phys.* 98 (2018) 119

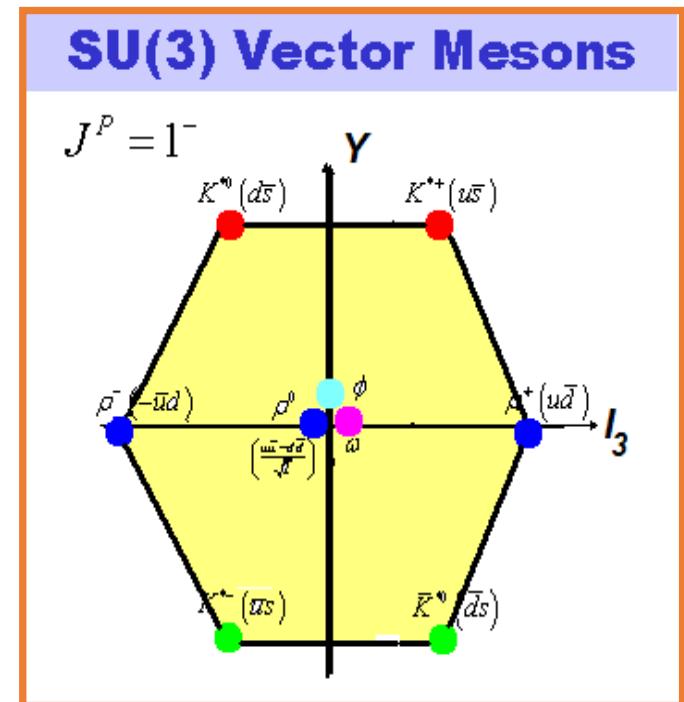
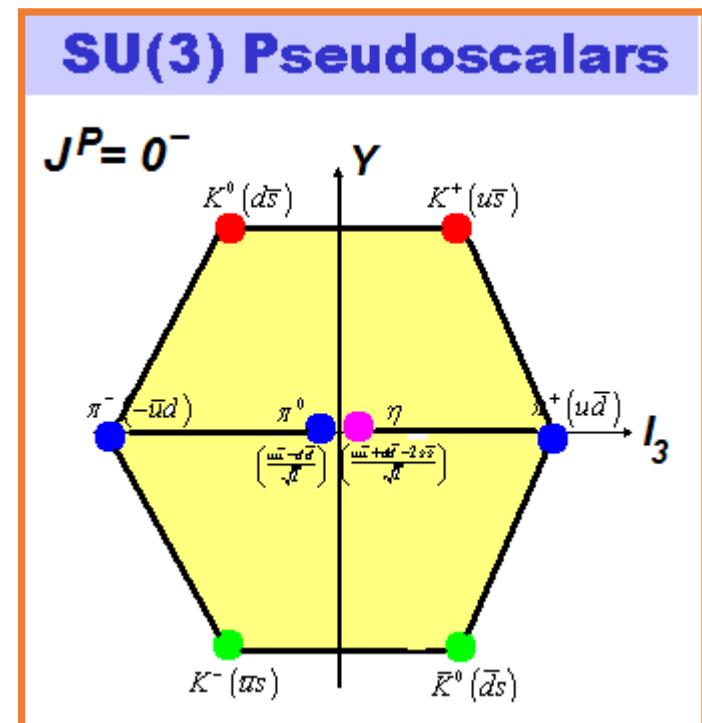
# SU(3) Scheme of Octet Baryon Interactions



# The Octet/Nonet Mesons



Nuclear Mean-Fields



## Lagrangian Density of BB-Octet Interactions

$$\mathcal{L}_{int}^{\mathcal{P}} = -\sqrt{2} \left\{ g_D [\bar{B}B\mathcal{P}_8]_D + g_F [\bar{B}B\mathcal{P}_8]_F \right\} - g_S \frac{1}{\sqrt{3}} [\bar{B}B\mathcal{P}_1]_S$$

### P=Pseudoscalar, Vector, and Scalar Meson Exchange

anti-symmetric  $[\bar{B}, B] = \bar{B}B - B\bar{B}$  and symmetric  $\{\bar{B}, B\} = \bar{B}B + B\bar{B}$  configurations

$$[\bar{B}B\mathcal{P}]_D = \text{Tr}(\{\bar{B}, B\} \mathcal{P}_8) , \quad [\bar{B}B\mathcal{P}]_F = \text{Tr}([\bar{B}, B] \mathcal{P}_8) , \quad [\bar{B}B\mathcal{P}]_S = \text{Tr}(\bar{B}B)\text{Tr}(\mathcal{P}_1)$$

### Octet Baryon Physics:

3 sets of 3 Fundamental Coupling Constants  $\{g_D, g_F, g_S\}$  fix the in total 48 BB'M Vertices

# BBM Vertices under Singlet-Octet Mixing: (mean-field producing ) Vector Couplings

Vertex	Coupling constant
$NN\omega$	$g_N^\omega = g_S \cos(\theta) + \sqrt{\frac{3}{2}} g_F \sin(\theta) - \frac{1}{\sqrt{6}} g_D \sin(\theta)$
$NN\phi$	$g_N^\phi = g_S \sin(\theta) - \sqrt{\frac{3}{2}} g_F \cos(\theta) + \frac{1}{\sqrt{6}} g_D \cos(\theta)$
$NN\rho$	$g_N^\rho = \sqrt{2}(g_F + g_D)$
$\Lambda\Lambda\omega$	$g_\Lambda^\omega = g_S \cos(\theta) - \sqrt{\frac{2}{3}} g_D \sin(\theta)$
$\Lambda\Lambda\phi$	$g_\Lambda^\phi = g_S \sin(\theta) + \sqrt{\frac{2}{3}} g_D \cos(\theta)$
$\Sigma\Sigma\omega$	$g_\Sigma^\omega = g_S \cos(\theta) + \sqrt{\frac{2}{3}} g_D \sin(\theta)$
$\Sigma\Sigma\phi$	$g_\Sigma^\phi = g_S \sin(\theta) - \sqrt{\frac{2}{3}} g_D \cos(\theta)$
$\Sigma\Sigma\rho$	$g_\Sigma^\rho = \sqrt{2}g_F$
$\Lambda\Sigma\rho$	$g_{\Lambda\Sigma}^\rho = \sqrt{\frac{2}{3}} g_D$
$\Xi\Xi\omega$	$g_\Xi^\omega = g_S \cos(\theta) - \sqrt{\frac{3}{2}} g_F \sin(\theta) - \frac{1}{\sqrt{6}} g_D \sin(\theta)$
$\Xi\Xi\phi$	$g_\Xi^\phi = g_S \sin(\theta) + \sqrt{\frac{3}{2}} g_F \cos(\theta) + \frac{1}{\sqrt{6}} g_D \cos(\theta)$
$\Xi\Xi\rho$	$g_\Xi^\rho = \sqrt{2}(g_F - g_D)$

## Guiding Principles for a SU(3) DFT Mean-Field Dynamics

- Interactions inherit SU(3) symmetry from NN (BB) scattering data
- Incorporate SU(3)-breaking by use of physical masses and empirical coupling constants
- Free space, in-medium Bethe-Salpeter, and vertex equations conserve the fundamental symmetries

### The SU(3) DFT/EDF Program:

- Three in-medium couplings are needed to fix the scalar and vector sets of  $\{g_D, g_F, g_S\}$
- For known  $g_{NN\omega}(\rho)$ ,  $g_{NN\rho}(\rho)$  and imposing  $g_{NN\phi}(\rho)=0 \rightarrow$  vector  $g_D(\rho), g_F(\rho), g_S(\rho)$
- For known  $g_{NN\sigma}(\rho)$ ,  $g_{NN\delta}(\rho)$  and imposing  $g_{NN\sigma'}(\rho)=0 \rightarrow$  scalar  $g_D(\rho), g_F(\rho), g_S(\rho)$
- Fix the mixing angle – ideal mixing  $\tan(\theta)=1/\sqrt{2}$  ( $\theta \sim 35^\circ$ )
- → SU(3) BBM vector and scalar mean-field couplings

# Covariant SU(3) Density Functional

$$\mathcal{L}_{int}^{DF} = -\sqrt{2} \sum_{\mathcal{M} \in \{\mathcal{P}, \mathcal{S}, \mathcal{V}\}} \left\{ g_D^{*(\mathcal{M})}(\hat{\rho}) [\bar{\mathcal{B}}\mathcal{B}\mathcal{P}_8]_D + g_F^{*(\mathcal{M})}(\hat{\rho}) [\bar{\mathcal{B}}\mathcal{B}\mathcal{P}_8]_F - g_S^{*(\mathcal{M})}(\hat{\rho}) \frac{1}{\sqrt{6}} [\bar{\mathcal{B}}\mathcal{B}\mathcal{P}_1]_S \right\}$$

$\hat{\rho} = F(\bar{\Psi}, \Psi)$

**Meson and Baryon Field Equations:**

$$(\partial_\mu \partial^\mu + m_{\mathcal{M}}^2) \Phi_{\mathcal{M}}^s = \sum_{BB'} g_{BB'\mathcal{M}}^*(\hat{\rho}) \rho^{BB's}, \quad (\partial_\mu \partial^\mu + m_{\mathcal{M}}^2) V_{\mathcal{M}}^\lambda = \sum_{BB'} g_{BB'\mathcal{M}}^*(\hat{\rho}) \rho^{BB'\lambda}$$

$$(\gamma_\mu (p^\mu - \Sigma_B^\mu(\hat{\rho})) - M_B + \Sigma_B^{(s)}(\hat{\rho})) \Psi_B = 0.$$

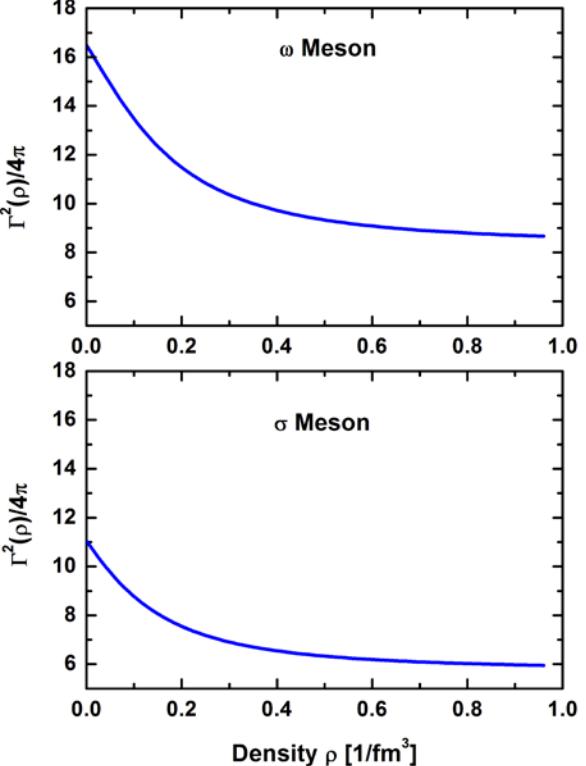
**Evaluated in mean-field approximation**

$$\Phi = \langle 0 | \Phi | 0 \rangle + \delta\Phi \quad ; \quad V^\mu = \langle 0 | V^\mu | 0 \rangle + \delta V^\mu \quad ; \quad \hat{\rho} = \langle 0 | \hat{\rho} | 0 \rangle + \delta\rho \quad ; \quad g_{BB'M}(\hat{\rho}) = g_{BB'M}(\rho_B) + \delta\rho \frac{\partial g_{BB'M}(\rho_B)}{\partial \rho_B}$$

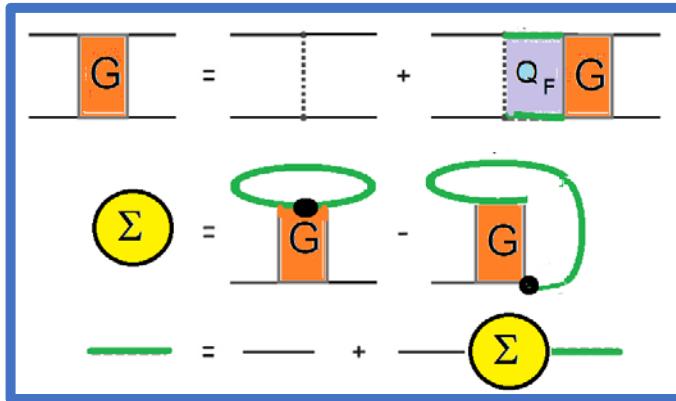
$$\phi = \langle 0 | \Phi | 0 \rangle \quad ; \quad V^0 = \langle 0 | V^\mu | 0 \rangle \delta^{\mu 0} \quad ; \quad \rho_B \equiv \langle 0 | \hat{\rho} | 0 \rangle \quad ; \quad \delta\rho = \hat{\rho} - \rho_B$$

# The Practitioner's Approach

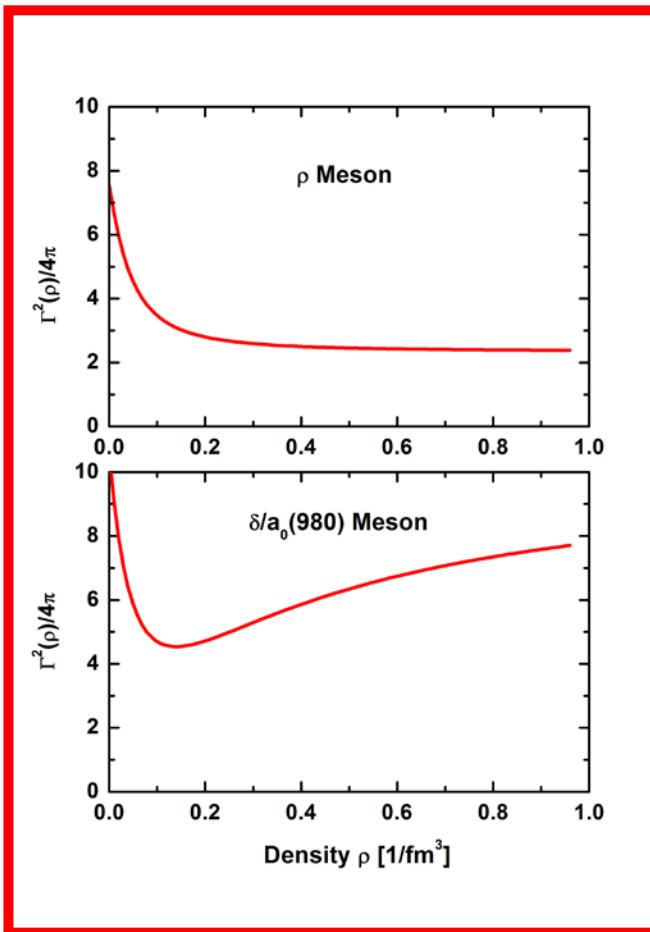
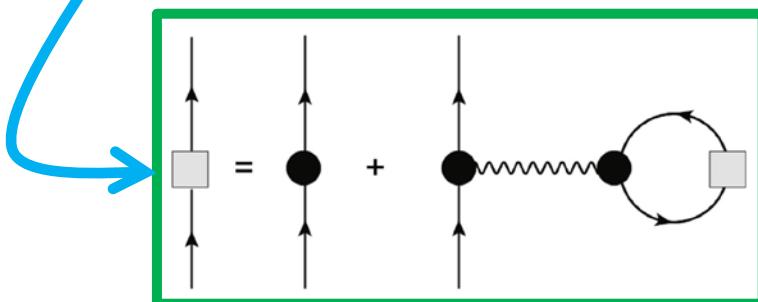
# Covariant Nuclear EDF from DBHF Interactions



Isoscalar Vertices



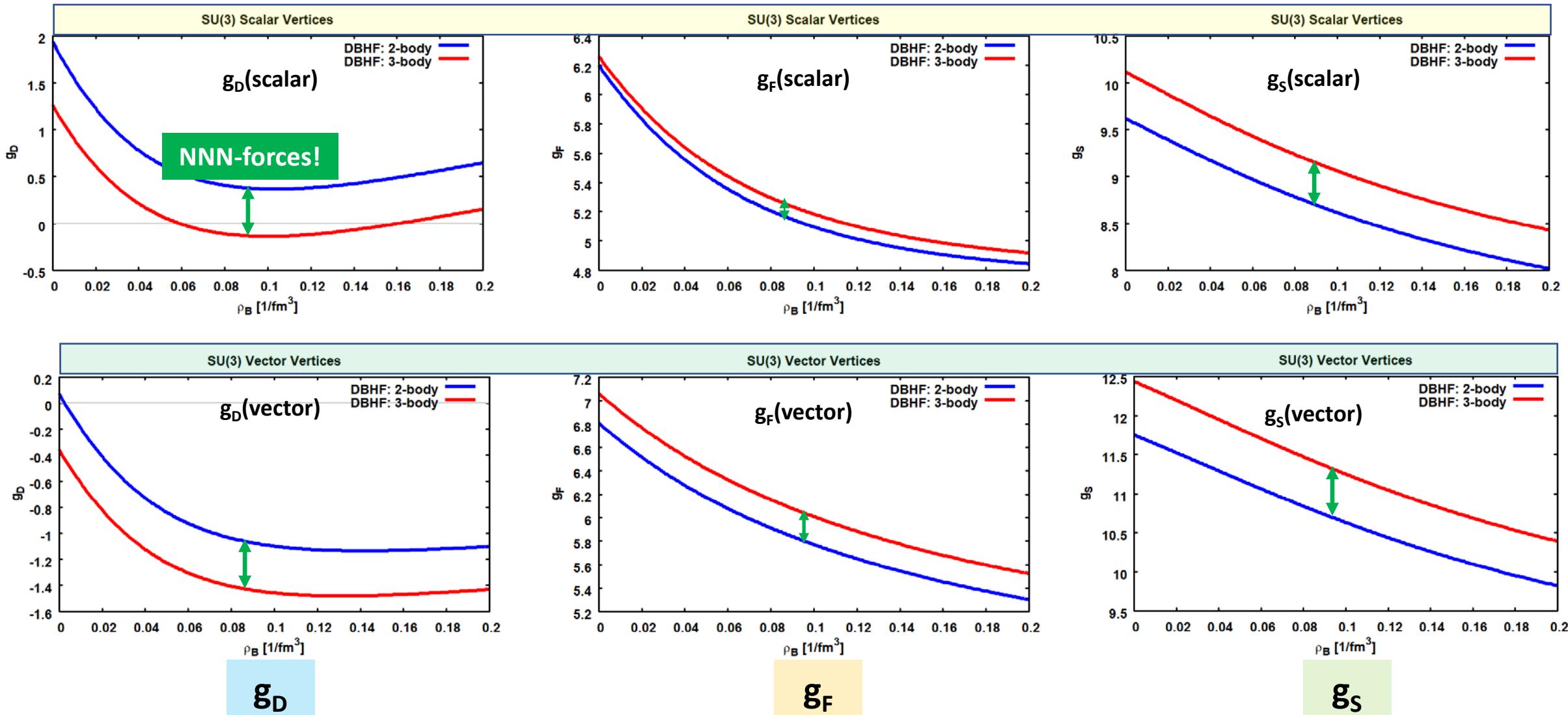
$$\Gamma_{BB'a}(q_s, k_F) \simeq \frac{1}{1 - \int dq' V_a G^* Q_F} |_{BB'} g_{BB'a}.$$



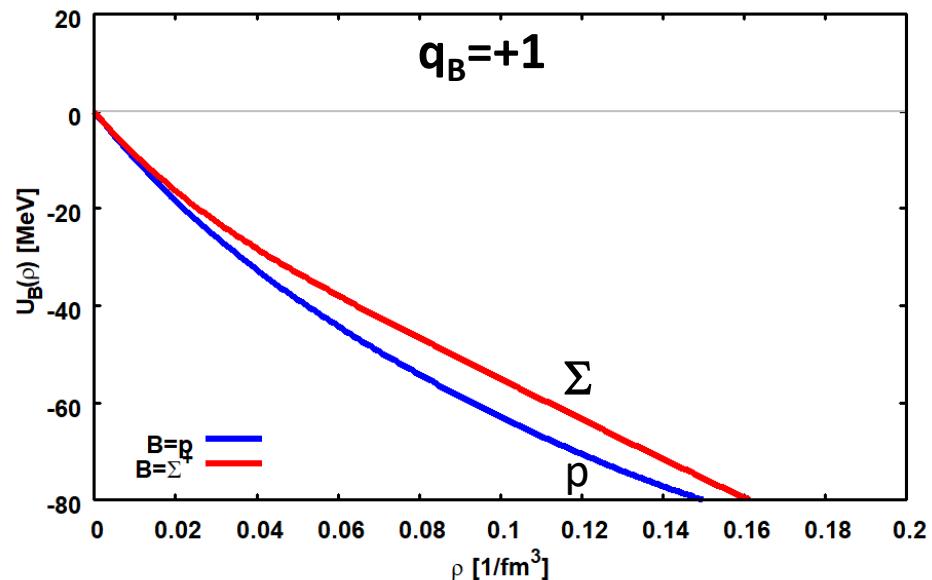
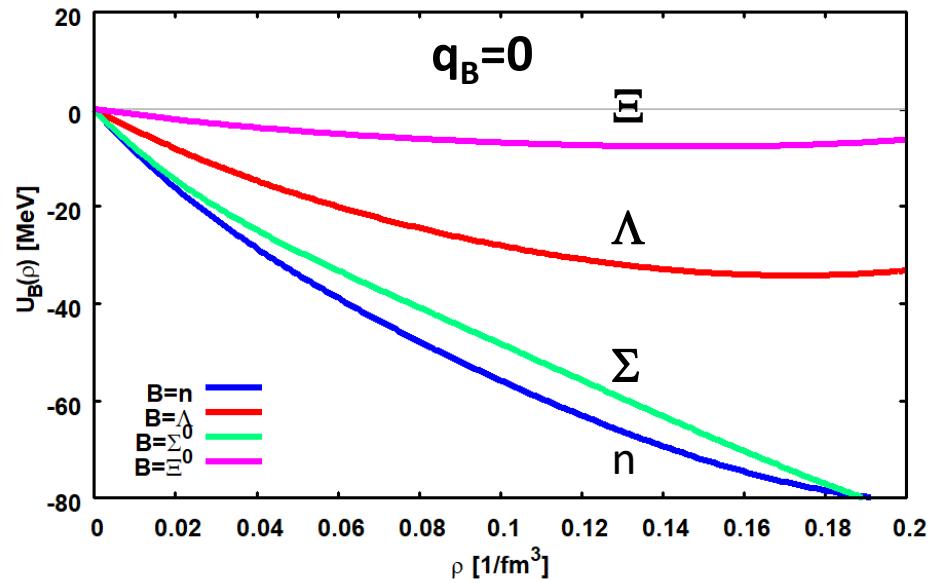
Isovector Vertices

# GI-DBHF SU(3) Vertices in Infinite Nuclear Matter

Ideal Mixing:  $\theta=35.26^\circ$



# SU(3) Baryon Mean-Fields

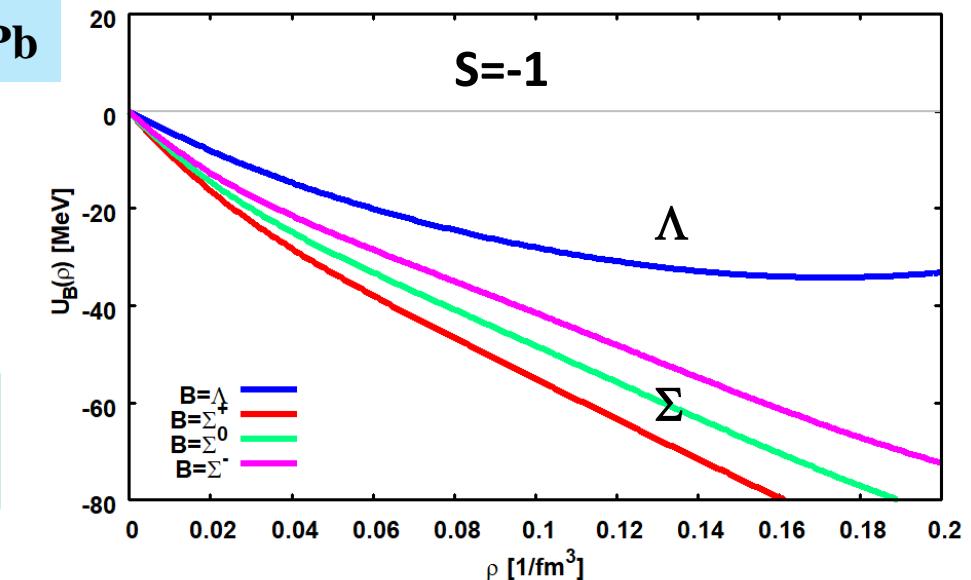
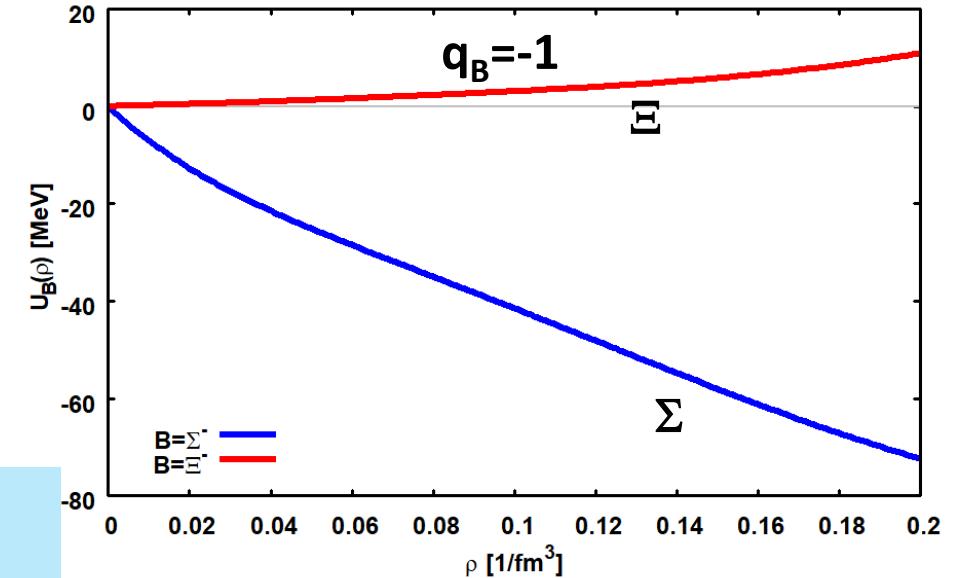


$Z/A=0.4$

as e.g.

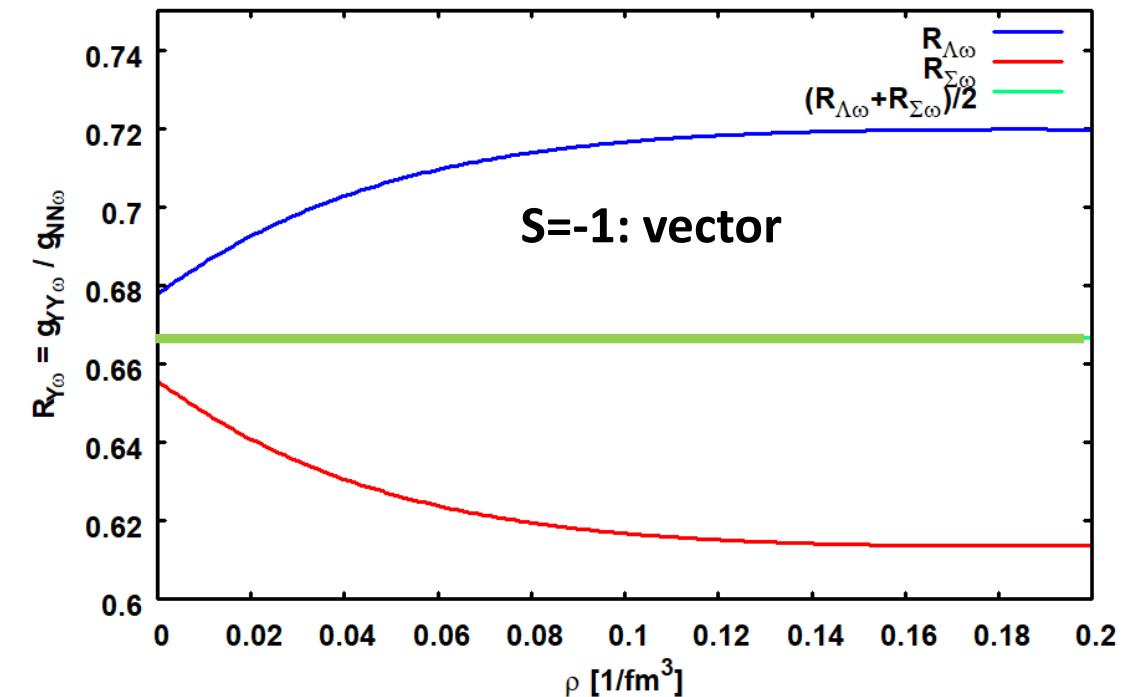
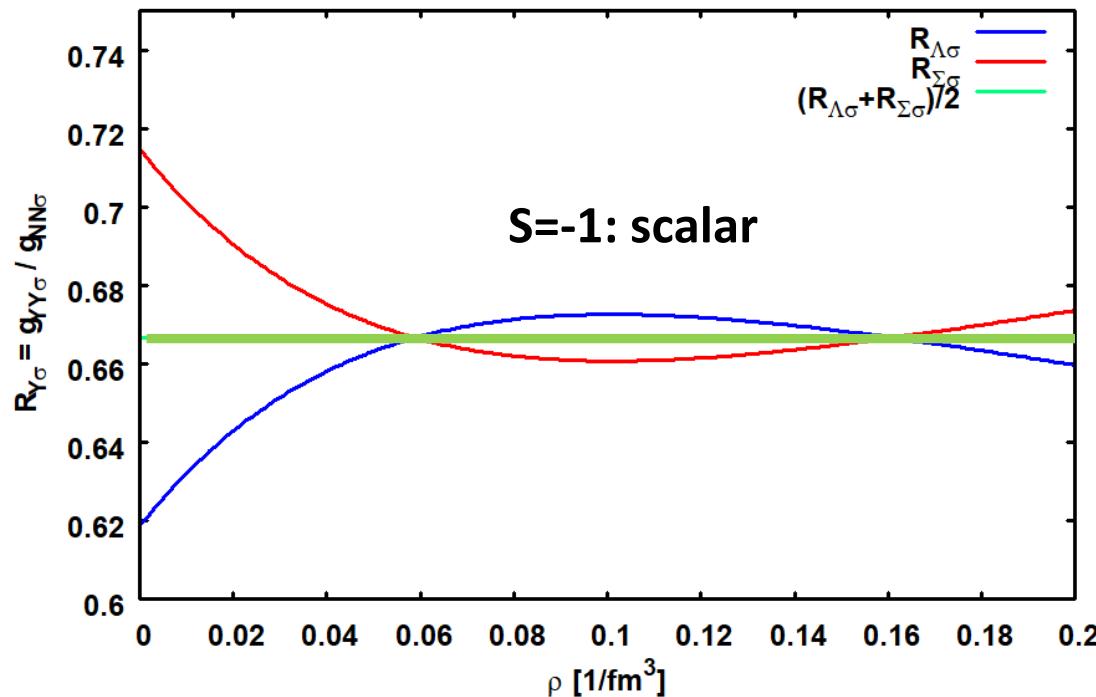
$^{10}\text{Be}, ^{48}\text{Ca}, ^{124}\text{Sn}, ^{208}\text{Pb}$

Similarity:  
 $\{\Sigma^+, \Sigma^0\} \sim \{p, n\}$



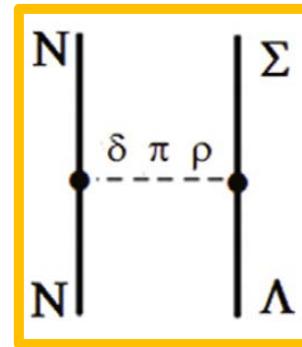
# Hyperon Couplings and „Quark Scaling“

$$R_{YM} = \frac{g_{YYm}}{g_{NNm}} \approx \frac{N_Y(u,d\text{-quarks})}{N_Y(\text{quarks})} \underset{s=-1}{=} \frac{2}{3}$$



In-Medium Quark Scaling Hypothesis recovered for  
Y=0,-1 Hypercharge-Multiplets

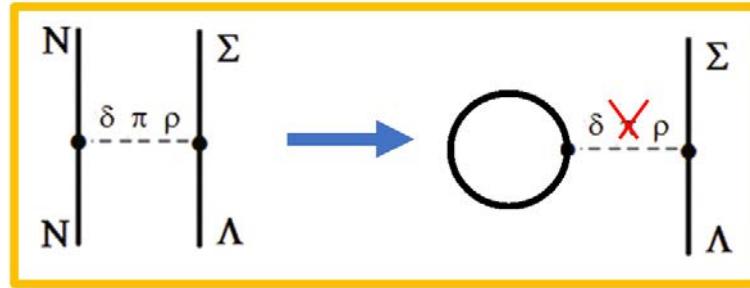
# $\Lambda$ - $\Sigma$ Mixing by Isovector Interactions



$$\mathcal{L} = \sum_{b=\Lambda,\Sigma} \bar{\psi}_b \left( \gamma_\mu \left( p^\mu + V_b^\mu \right) - M_b - \phi_b \right) \psi_b + \bar{\psi}_\Lambda U_{\Lambda\Sigma} \psi_\Sigma + h.c.$$

$$U_{\Lambda\Sigma} = \left( g_{\Lambda\Sigma\delta} \boldsymbol{\phi}_\delta + g_{\Lambda\Sigma\rho} \mathbf{V}_\rho^\mu \right) \cdot \boldsymbol{\tau}_{\Lambda\Sigma}$$

# $\Lambda$ - $\Sigma$ Mixing in Asymmetric Nuclear Matter Induced by the Static Isovector Mean-Field

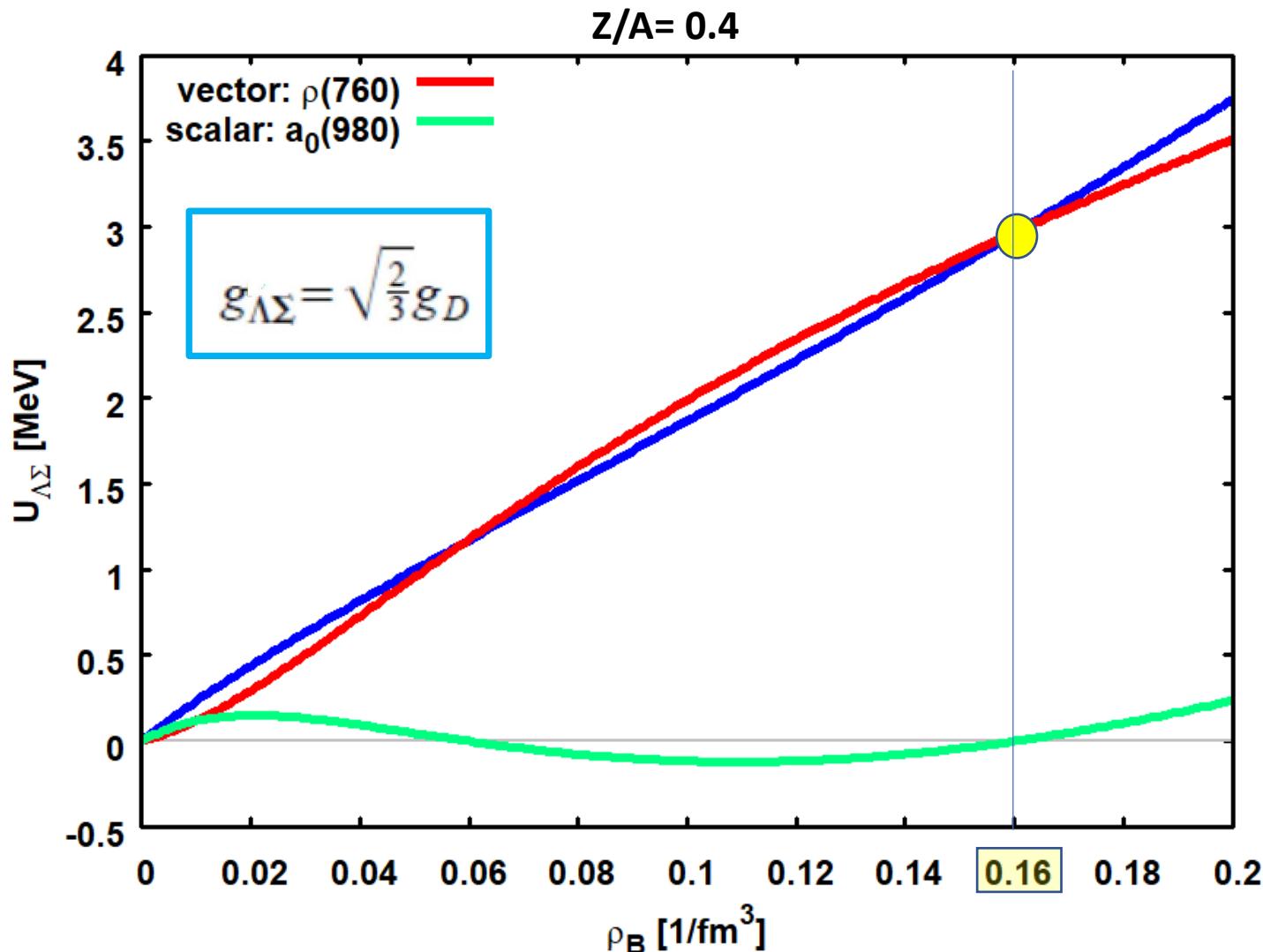


$$U_{\Lambda\Sigma}(\rho_B) = U_\delta^{(NN)}(\rho_B) \left( \frac{g_{\Lambda\Sigma\delta}}{g_{NN\delta}} \right) + U_\rho^{(NN)}(\rho_B) \left( \frac{g_{\Lambda\Sigma\rho}}{g_{NN\rho}} \right)$$

Mean-Field Induced Mixing

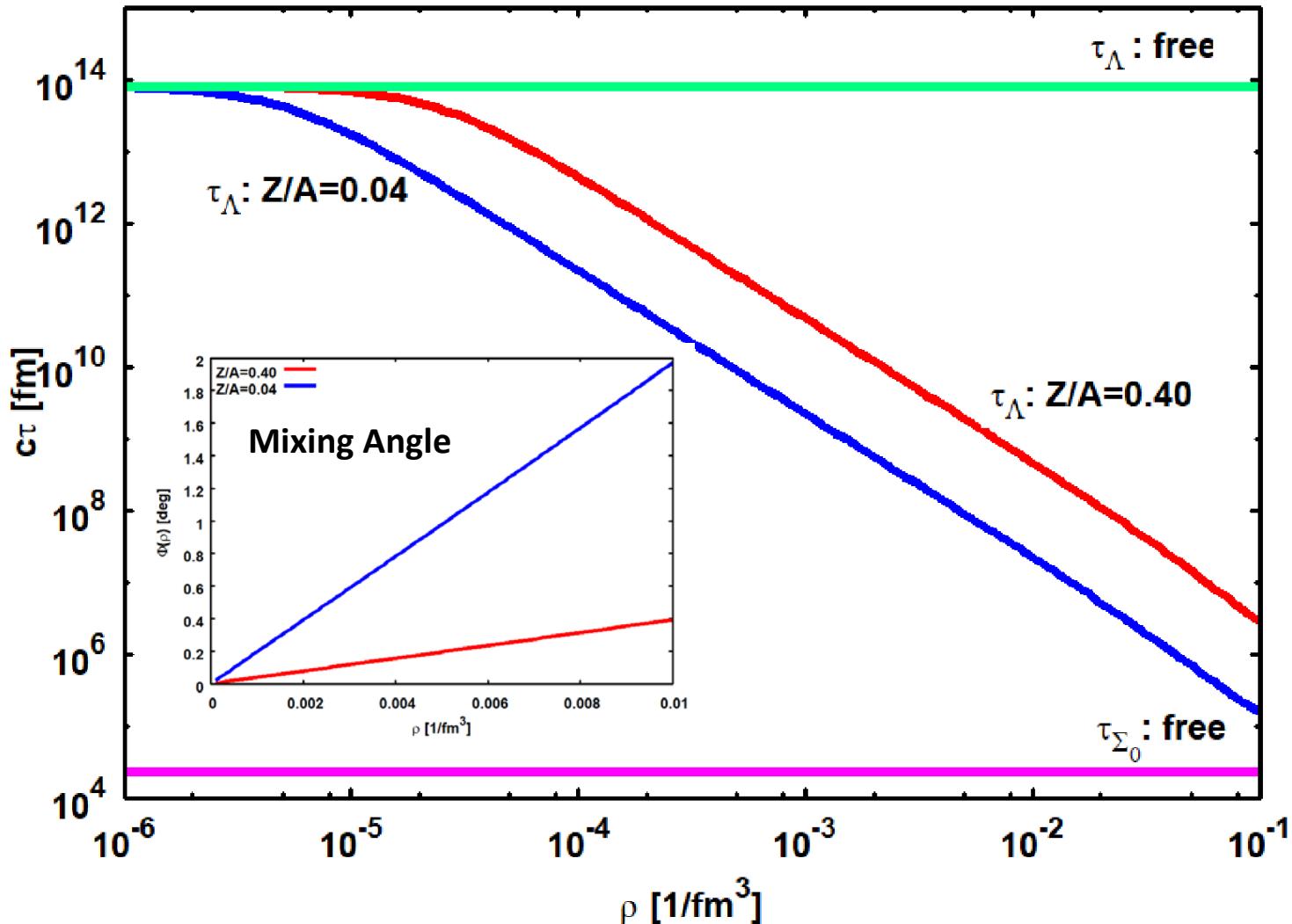
$$\begin{pmatrix} H_{\Lambda A} - E & U_{\Lambda\Sigma} \\ U_{\Lambda\Sigma}^\dagger & H_{\Lambda A} + m_{\Sigma\Lambda} - E \end{pmatrix} \begin{pmatrix} [\phi_\Lambda \otimes |A\rangle]_{I_A N_A} \\ [\phi_\Sigma \otimes |A\rangle]_{I_A N_A} \end{pmatrix} = \mathbf{0}$$

# $\Lambda$ - $\Sigma$ Mean-Field Mixing SU(3) Potential in Asymmetric Nuclear Matter

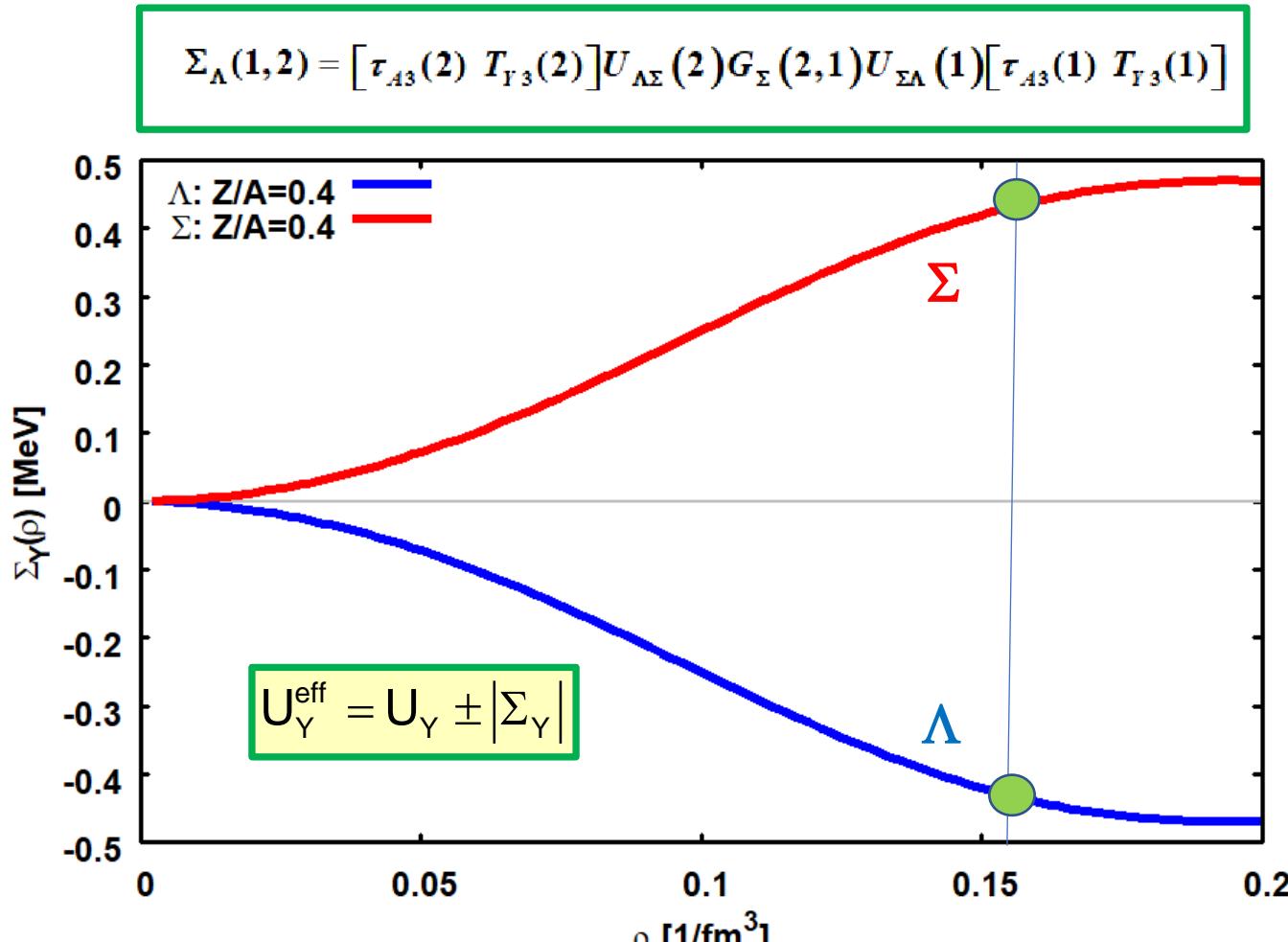


# In-Medium $\Lambda$ - $\Sigma$ Mixing and the $\Lambda$ Lifetime

$$\left| \Lambda^*(\rho) \right\rangle = \cos(\Phi) \left| \Lambda \right\rangle - \sin(\Phi) \left| \Sigma \right\rangle \quad \left| \Sigma^*(\rho) \right\rangle = \cos(\Phi) \left| \Sigma \right\rangle + \sin(\Phi) \left| \Lambda \right\rangle$$



# $\Lambda$ and $\Sigma$ Self-Energies Induced by the Isovector Mean-Field



- Depth of effective  $\Lambda$  potential increases
- Depth of effective  $\Sigma$  potential decreases

# Summary and Outlook

- Octet BB interactions in the SU(3) scheme
- Derivation of a SU(3)-based EDF from NN interactions
- „Quark Scaling“ in  $Y=0,-1$  ( $S=-1,-2$ ) hypercharge multiplets
- $\Lambda-\Sigma$  mixing induced by the isovector mean-field
- to come: Hypernuclear spectra, magnetic moments, neutron star EoS...

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