

Key Uncertainties in Astrophysical r-Process Nucleosynthesis

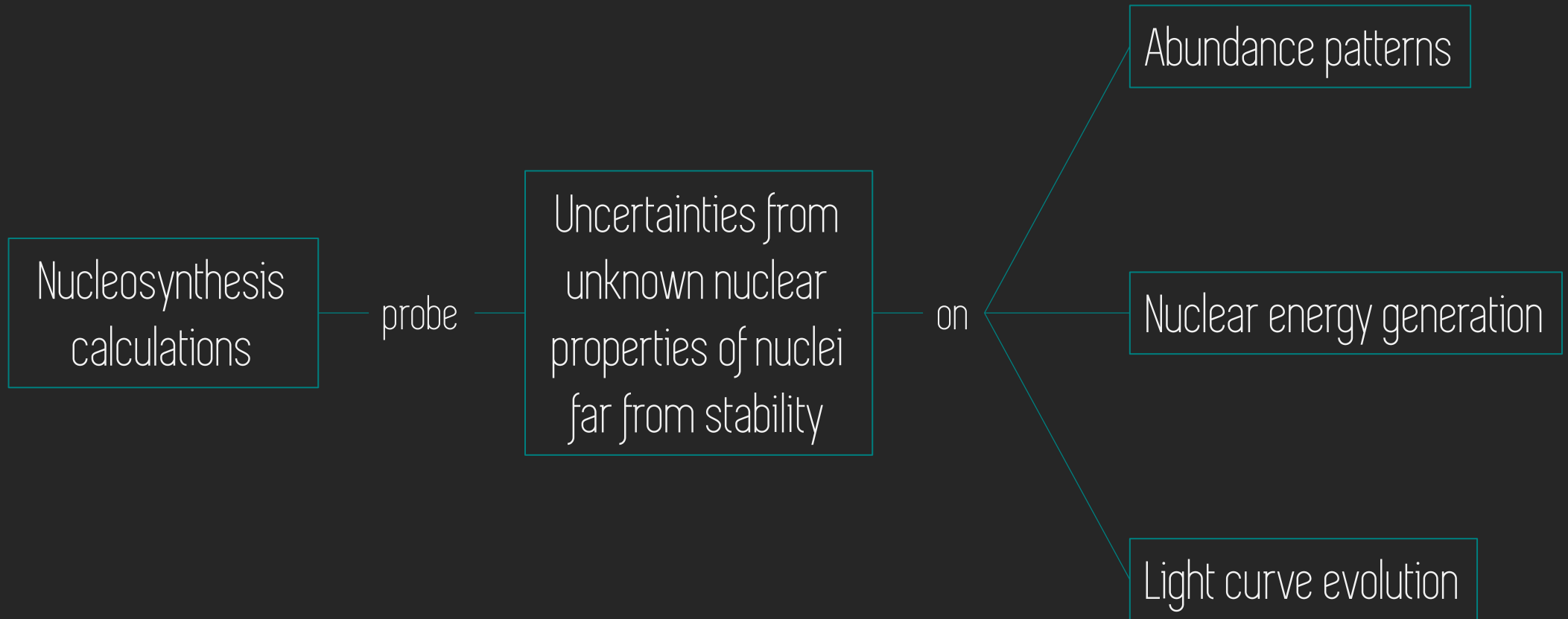
Kelsey Lund

19 October 2022

In collaboration with G. McLaughlin, Y. Zhu, J. Barnes and many others

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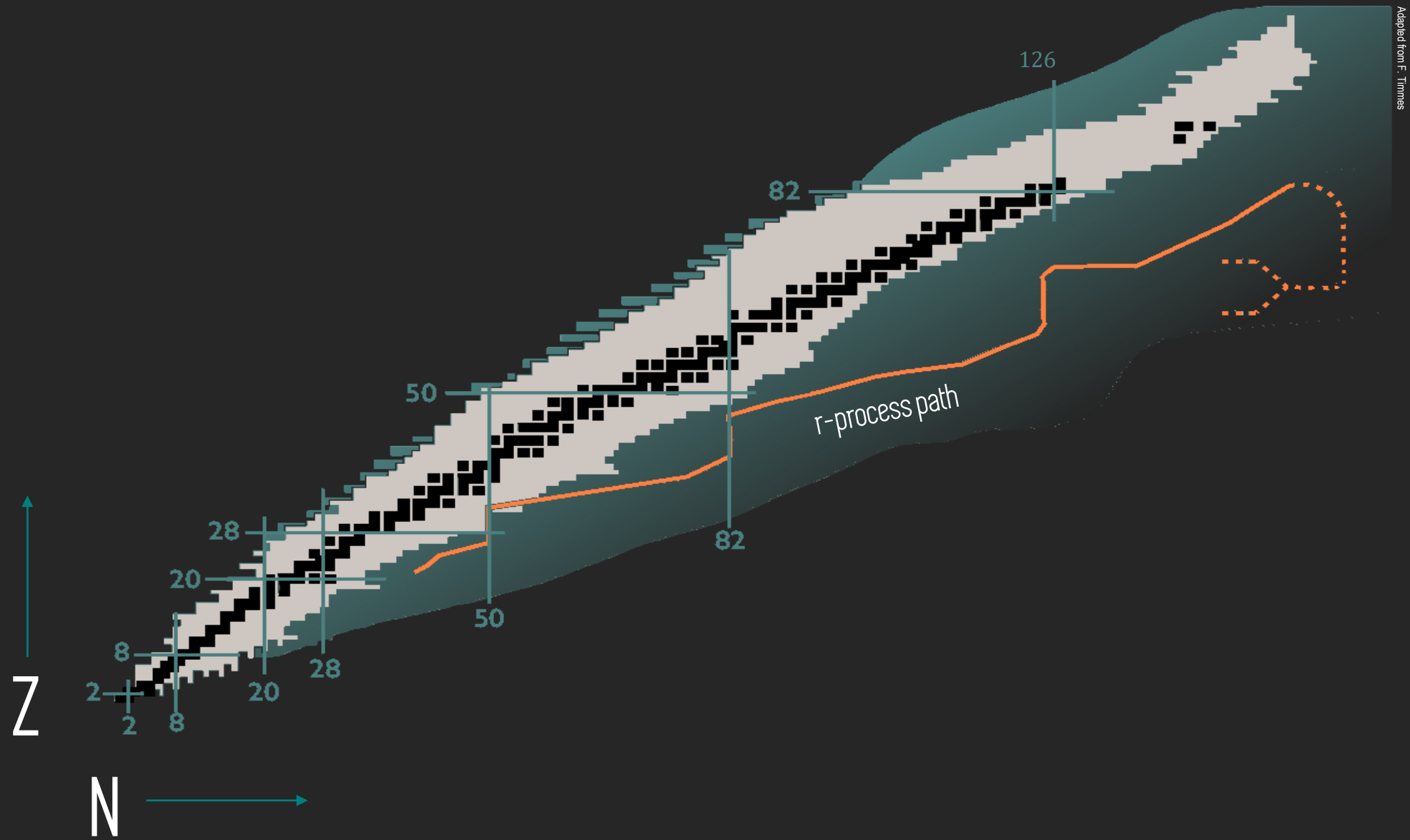


Miller+ 2019

arXiv: 2208.06373, 2010.03668, 2010.11182

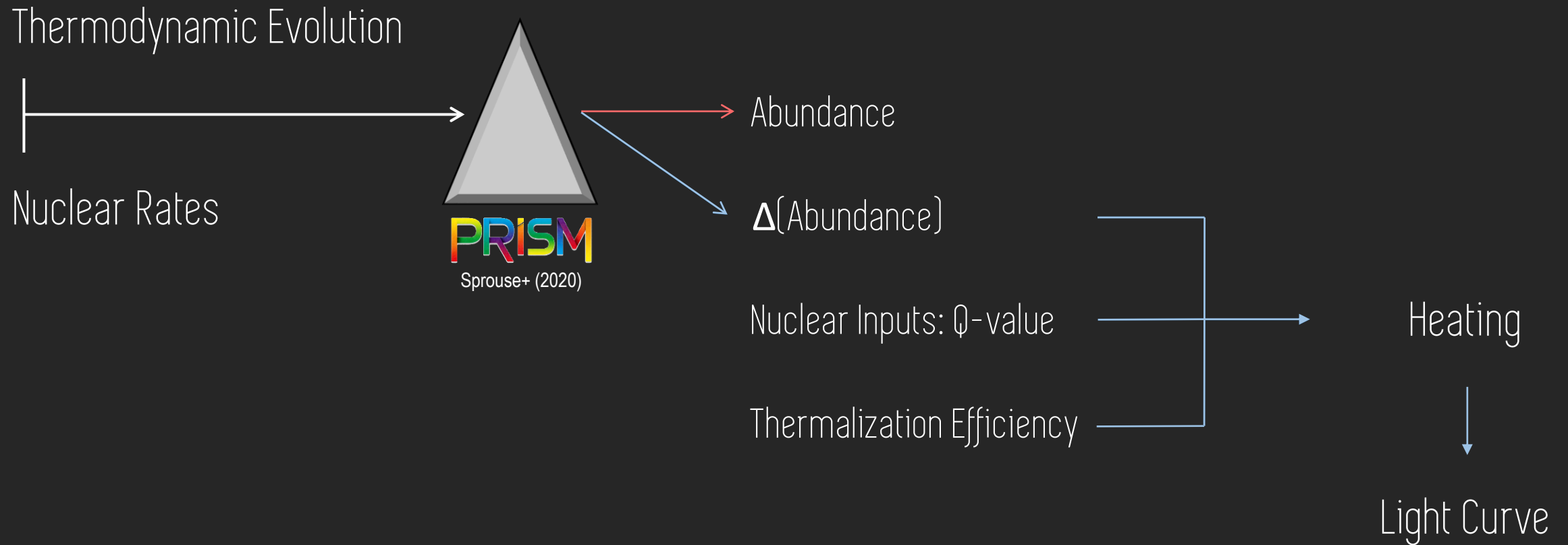
Sources of Nuclear Uncertainty

See also, e.g.
 Mumpower+ 2016
 Vassh+ 2019
 Giuliani+ 2020
 Nikas+ 2020
 Wang+ 2020
 Zhu+ 2021
 Barnes+ 2021
 Kullmann+ 2022
 and many others

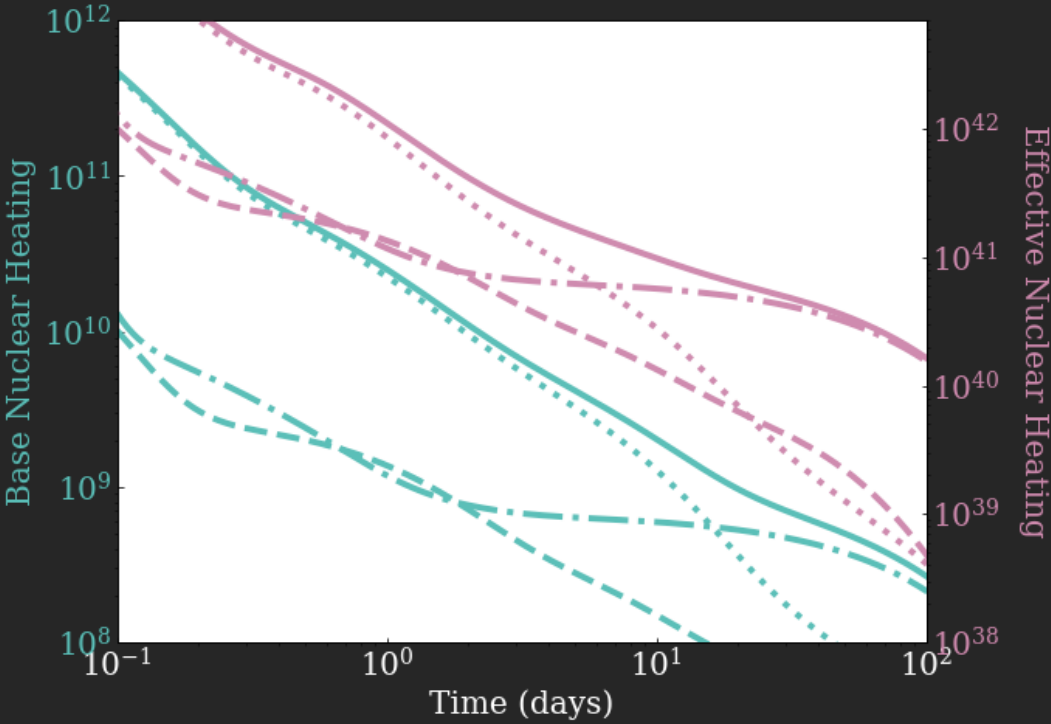


Adapted from F. Timmes

PRISM: A Sparse Matrix Solver



Nuclear Heating



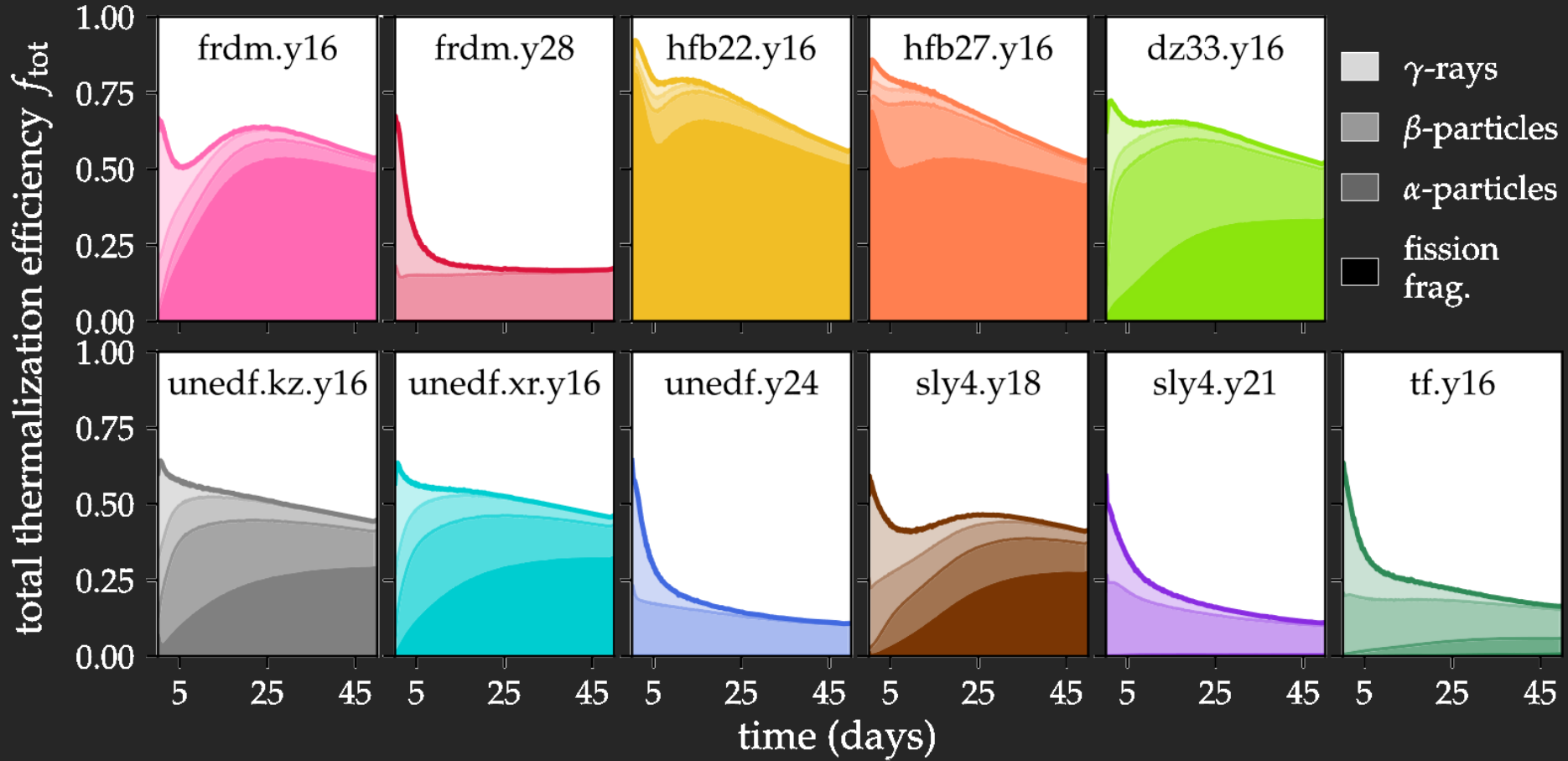
$$\dot{Q}(t) = \sum_i f_i(M_{ej}, v_{ej}, t) \dot{q}_i(t) M_{ej}$$



Thermalization efficiency: how effectively decay products can heat ejecta (function of time, ejecta mass, and characteristic velocity)

Thermalization based on Kasen & Barnes (2019)

See Barnes+ 2021



Light Curve Shell Model

Shell model for ejecta: the mass of each shell, M_v , depends on the velocity, v , of that shell (100 shells evenly distributed between $0.1c$ and $0.4c$)

Time evolution of the energy of a shell:

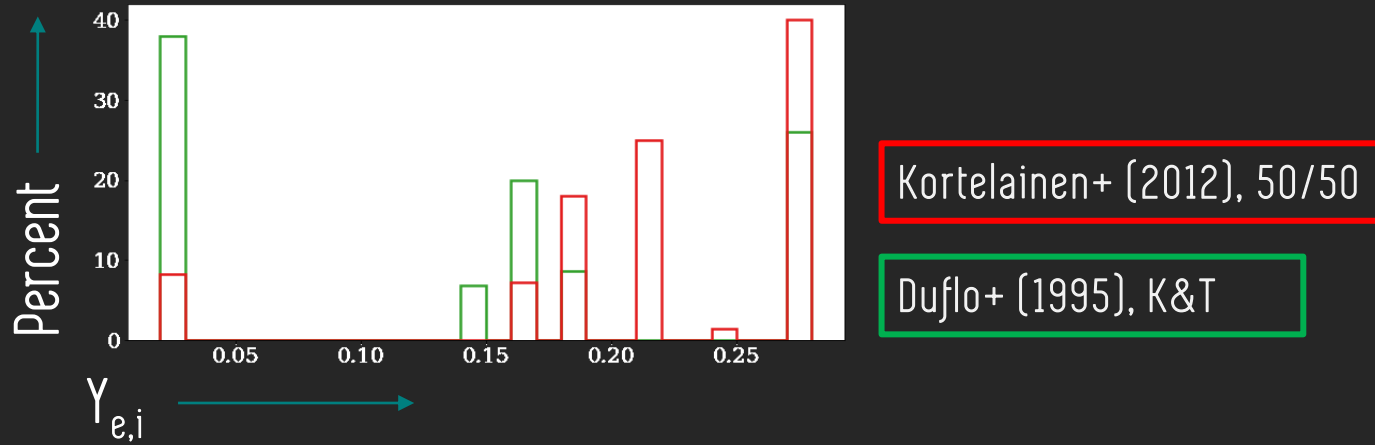
$$\frac{dE_v}{dt} = \underbrace{\frac{M_v}{M_{ej}} \dot{Q}(t, v)}_{\text{Effective heating}} - \underbrace{\frac{E_v}{t}}_{\text{Adiabatic expansion}} - \underbrace{\frac{E_v}{t_{d,v} + t_{lc,v}}}_{\text{Luminosity}}$$

$t_{d,v}$: Diffusion timescale (depends on opacity)
 $t_{lc,v}$: Light-crossing timescale

Combined Trajectories



Finding a Linear Combination



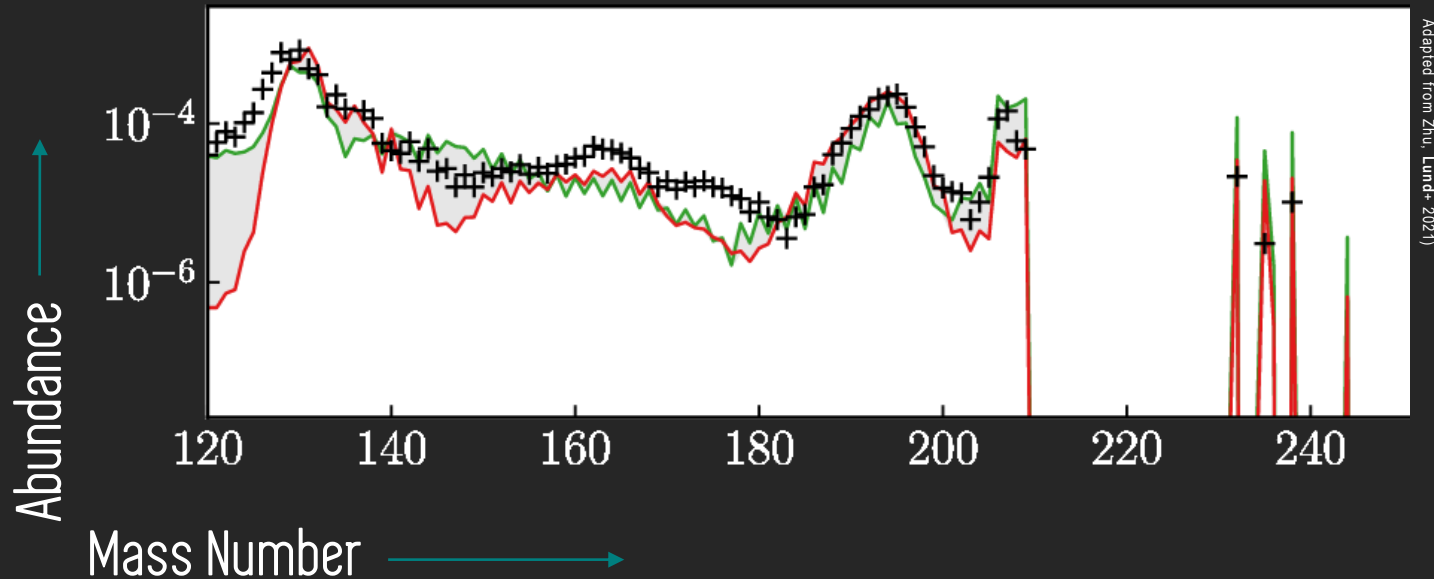
Kortelainen+ (2012), 50/50

Duflo+ (1995), K&T

Linear combination to fit to solar abundance pattern.

- Select from single- Y_e trajectories [0.02-0.28]

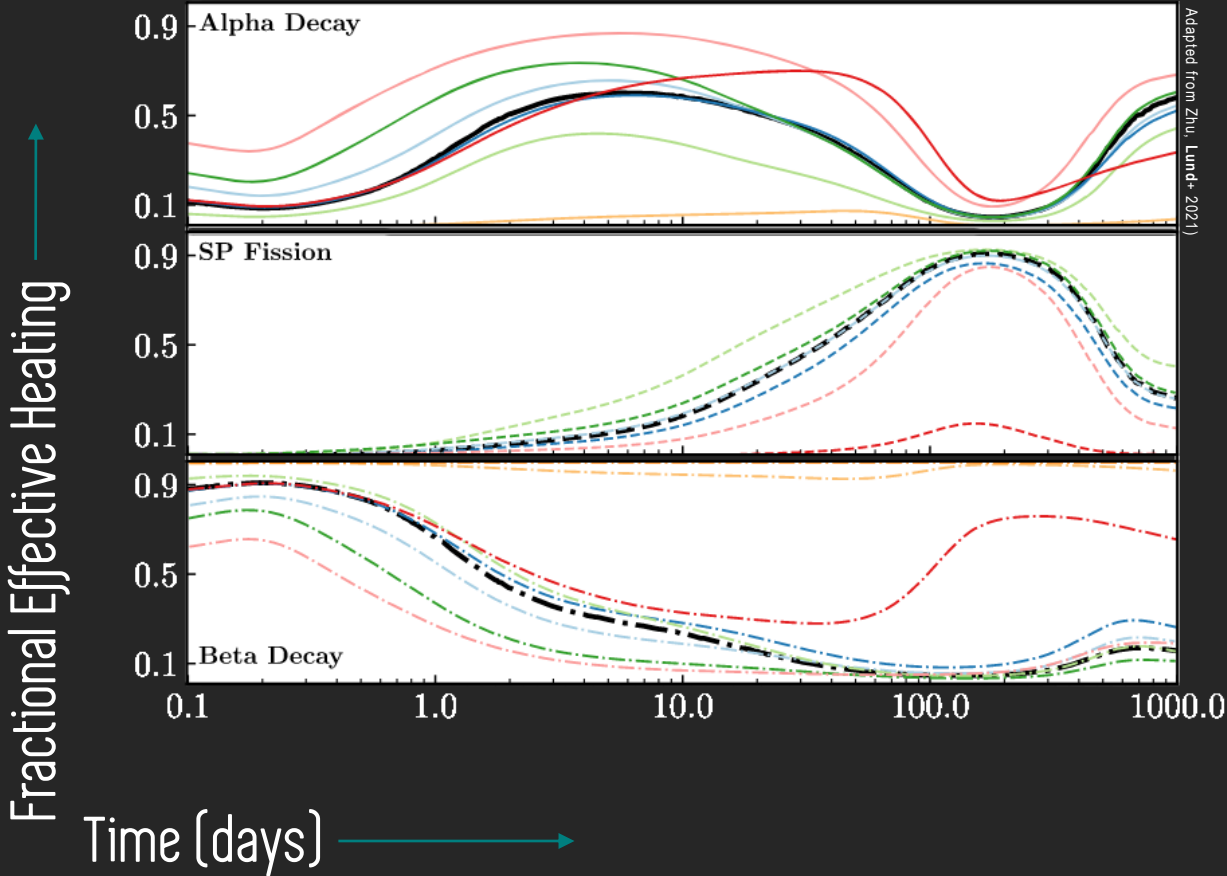
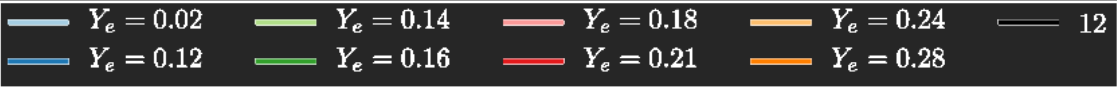
- Necessary to combine low *and* high Y_e material



Adapted from Zhu, Lund+ 2021)

Fractional Nuclear Heating

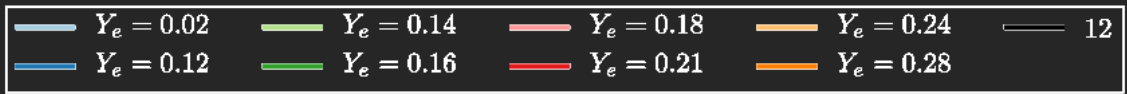
DZ33 + K&T



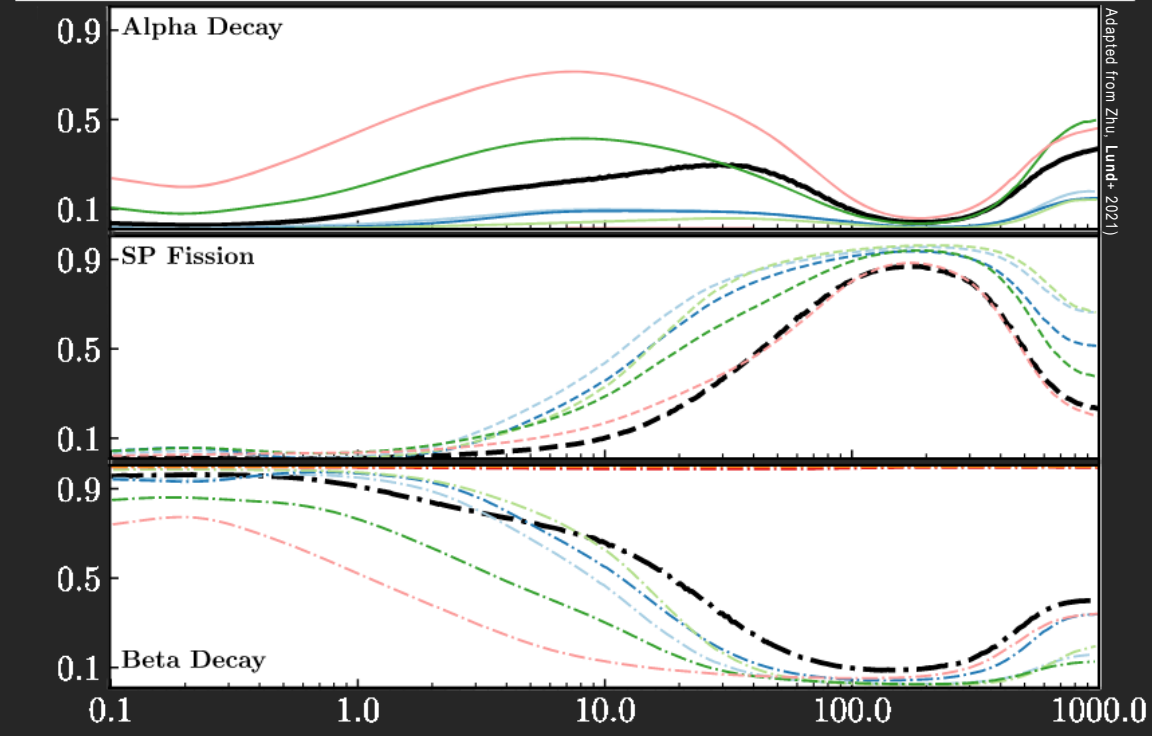
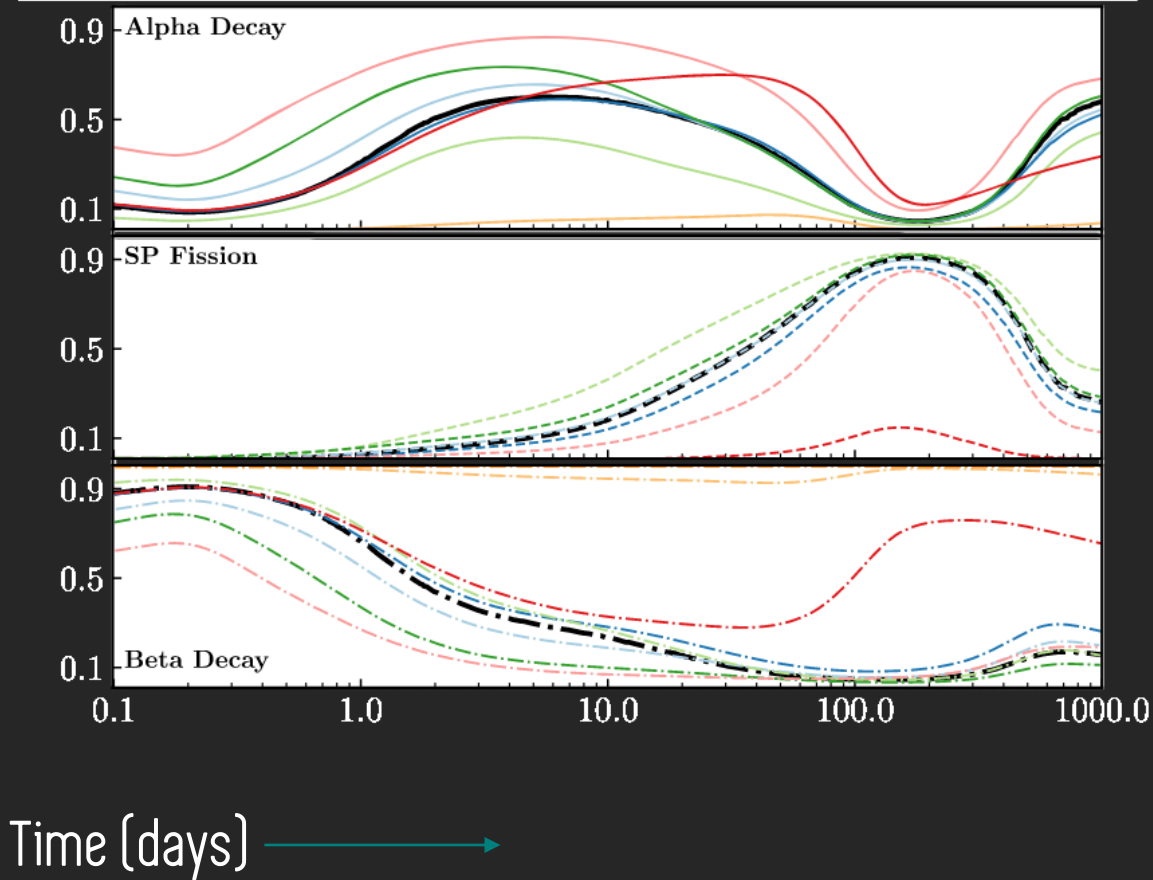
Fractional Nuclear Heating

DZ33 + K&T

UNEDF1 + 50/50

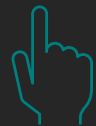


Fractional Effective Heating

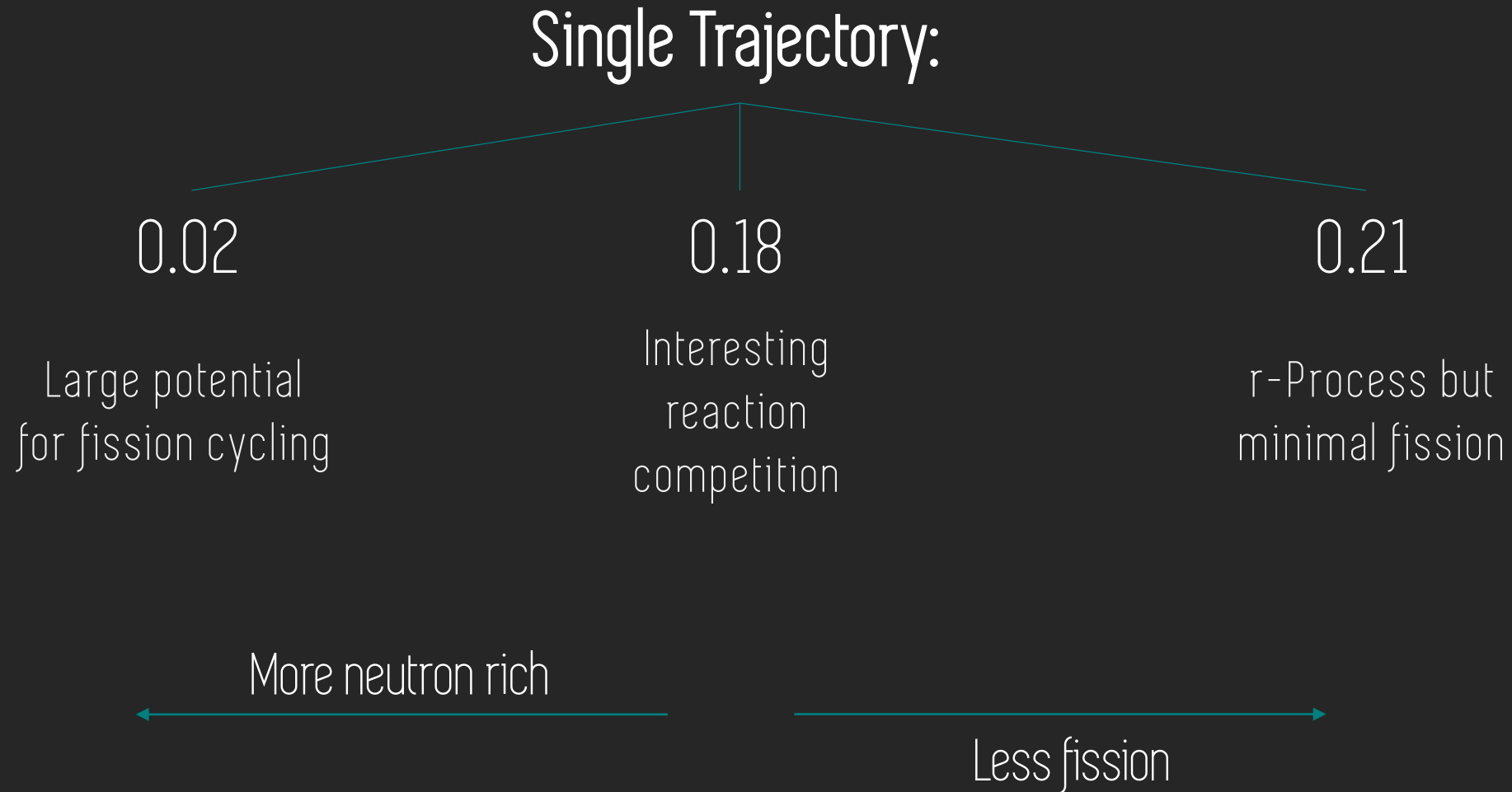


Adapted from Zhu, Lund+ 2021

Single-Ye Trajectories



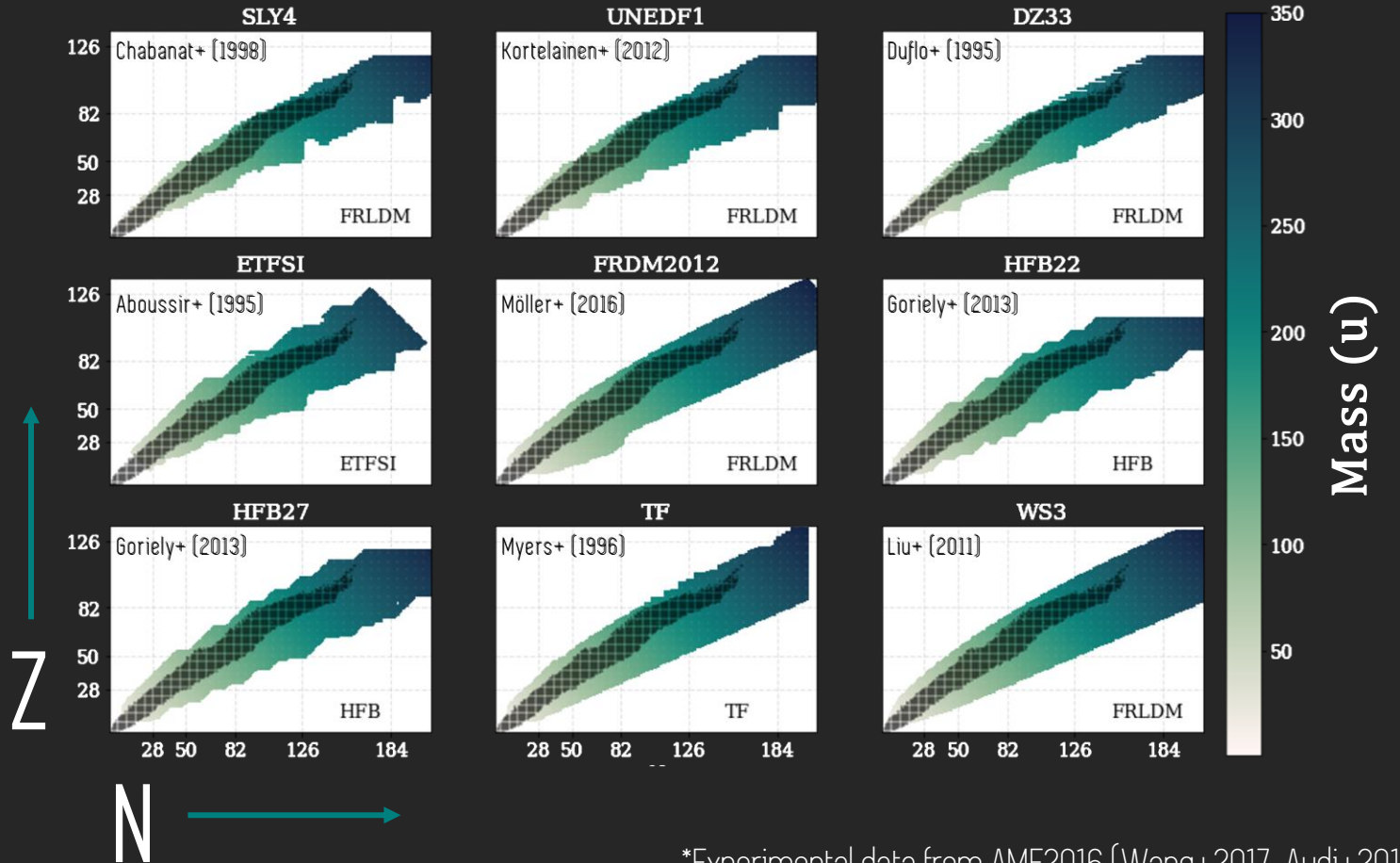
Astrophysical Uncertainty: Initial Y_e



Nuclear Uncertainty: Mass Model

Most basic nuclear property:
mass

Each mass model associated
with fission barrier height
model



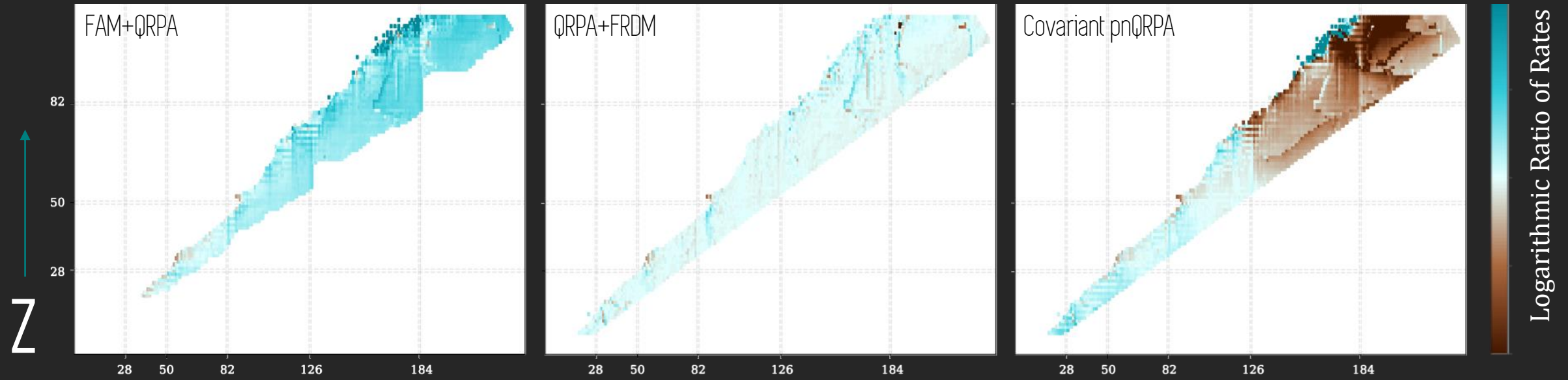
*Experimental data from AME2016 (Wang+2017, Audi+2017)

Effect of nuclear mass: Mendoza-Temis+2015, Mumpower+2015,2016, Barnes+2016, Wu+2019, Nikas+2020, Zhu+2021, Barnes+2021, Kullmann+2022

Nuclear Uncertainty: Beta Decay Rates

← ~ slower rate*

Logarithmic ratio compared to Möller+ 2003



N →
Ney (NES)
Ney+ 2020

Möller (MLR)
Möller+ 2019

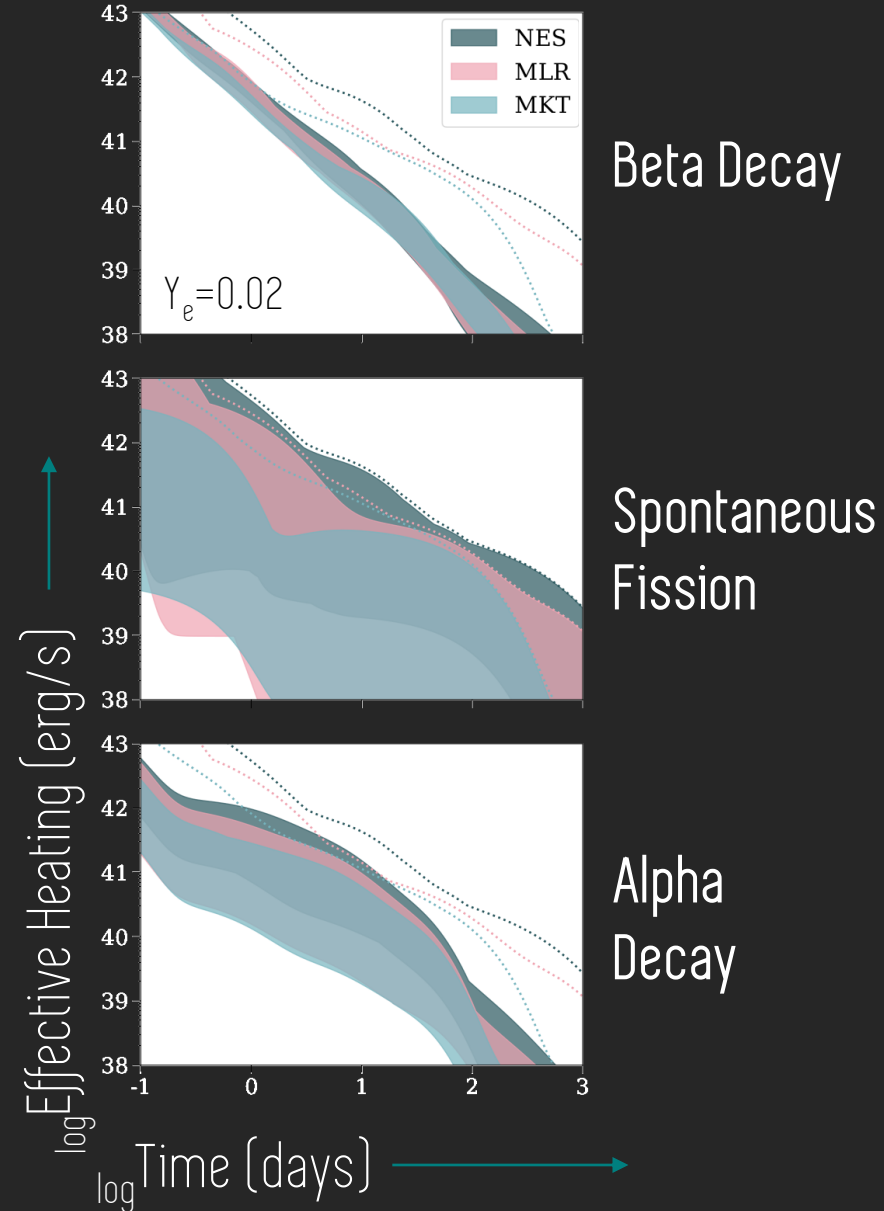
Marketin (MKT)
Marketin+ 2016

Effect of beta decay rates: Marketin+ 2016, Nikas+ 2020, Kullmann+ 2022

Nuclear Heating

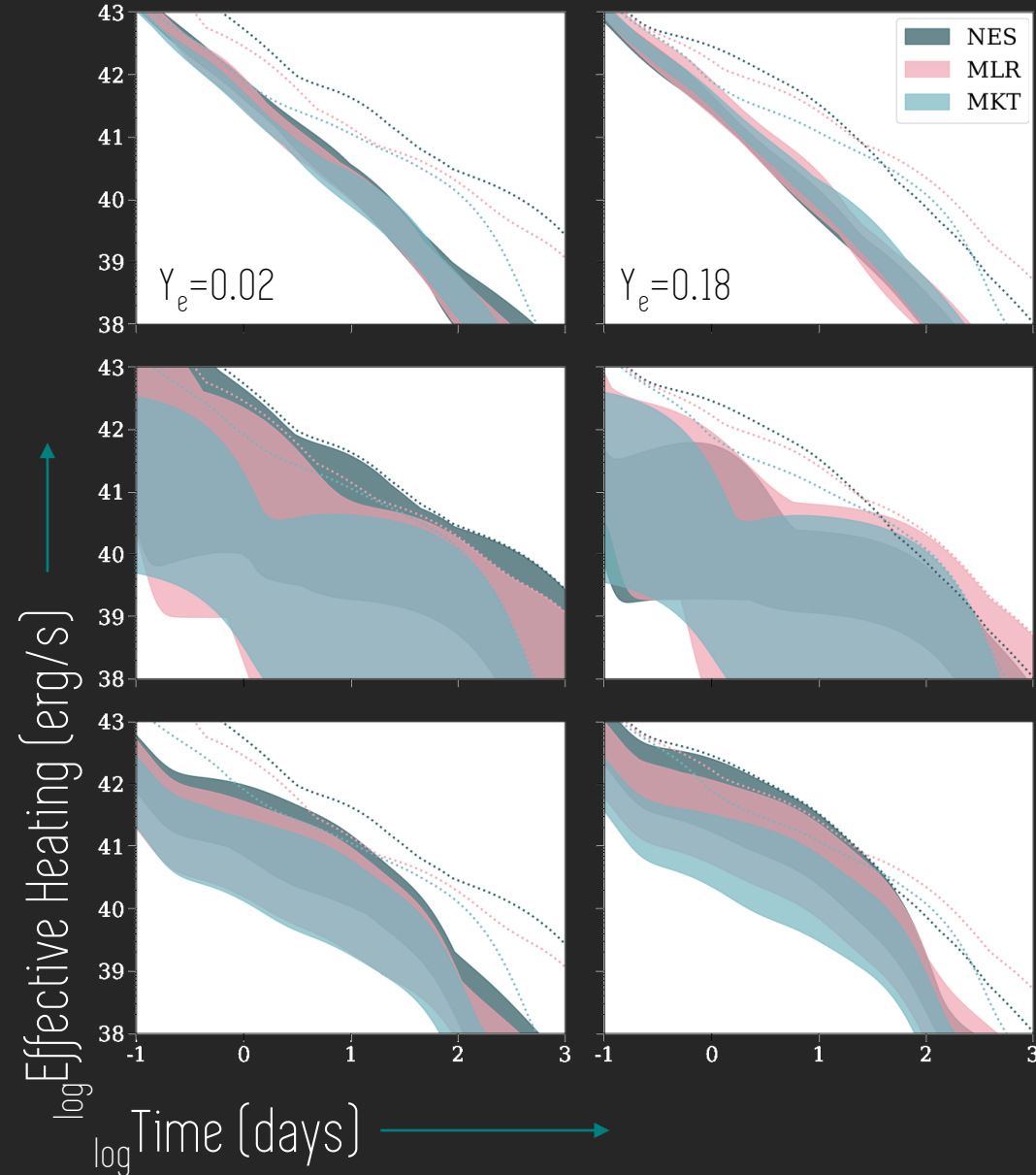


Nuclear Heating by Reaction Type



- Upper limit of heating uncertainty can be set by fission
- Beta models differ in behavior of dominating fission heating

Nuclear Heating by Reaction Type



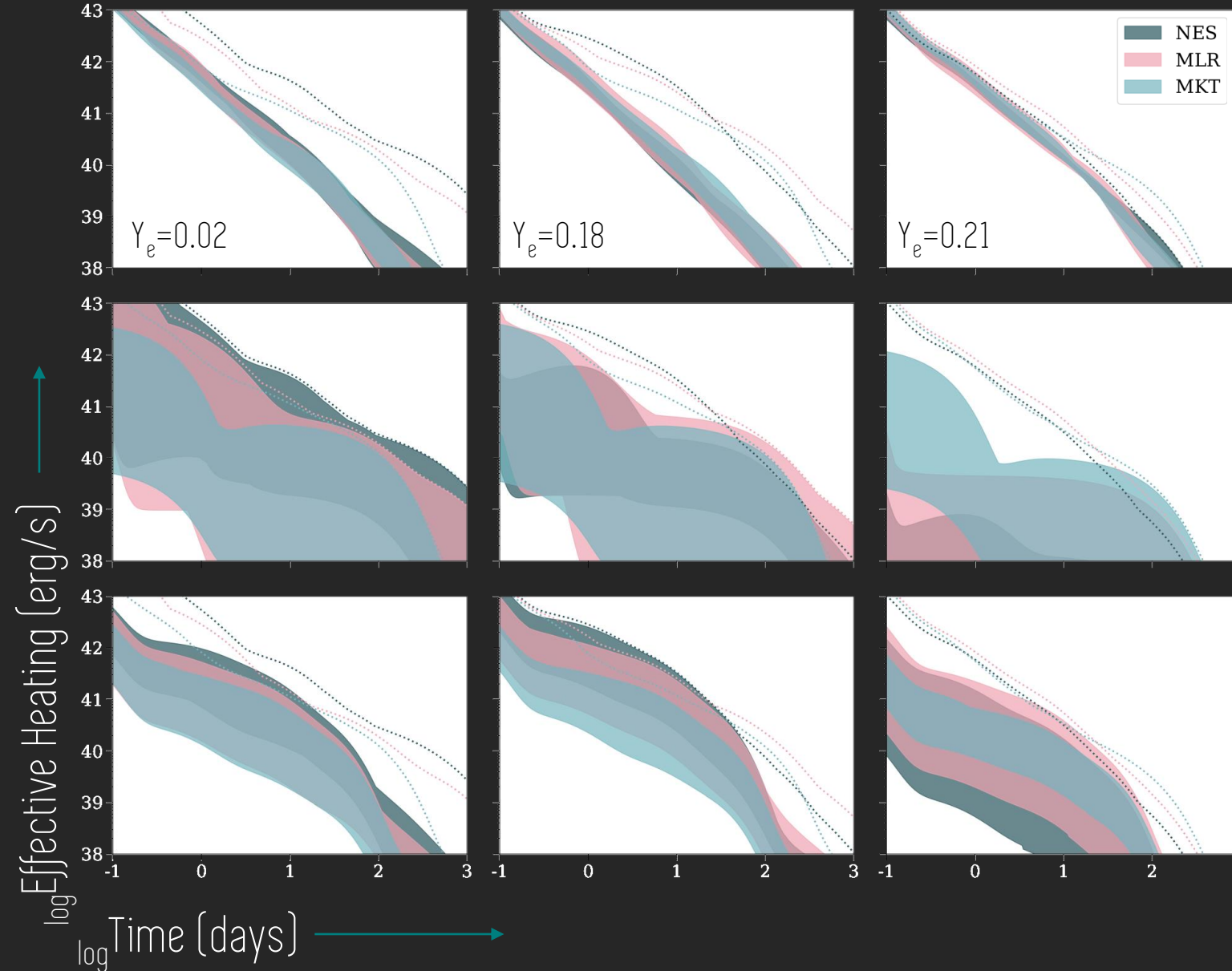
Beta Decay

Spontaneous Fission

Alpha Decay

- Alpha heating becomes more important 10-100 days
- Beta models differ in predicting when alpha tends to dominate heating+ late-time tail shape of fission

Nuclear Heating by Reaction Type



Beta Decay

- Much more overlap, total heating tends to be set by beta, then alpha decay

Sp. Fission

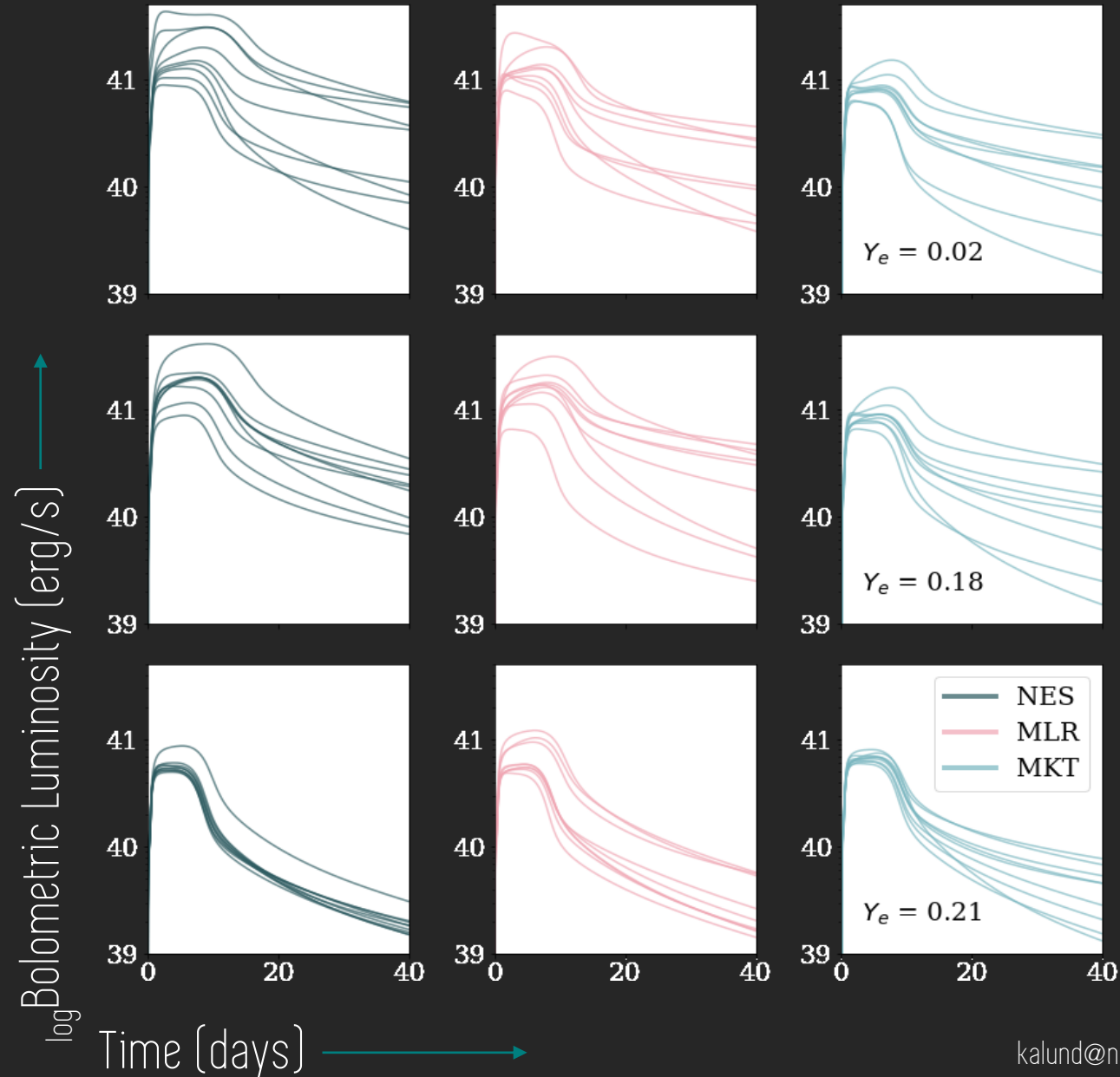
- Overall effect on beta decay heating is small

Alpha Decay

Light Curve



Light Curves



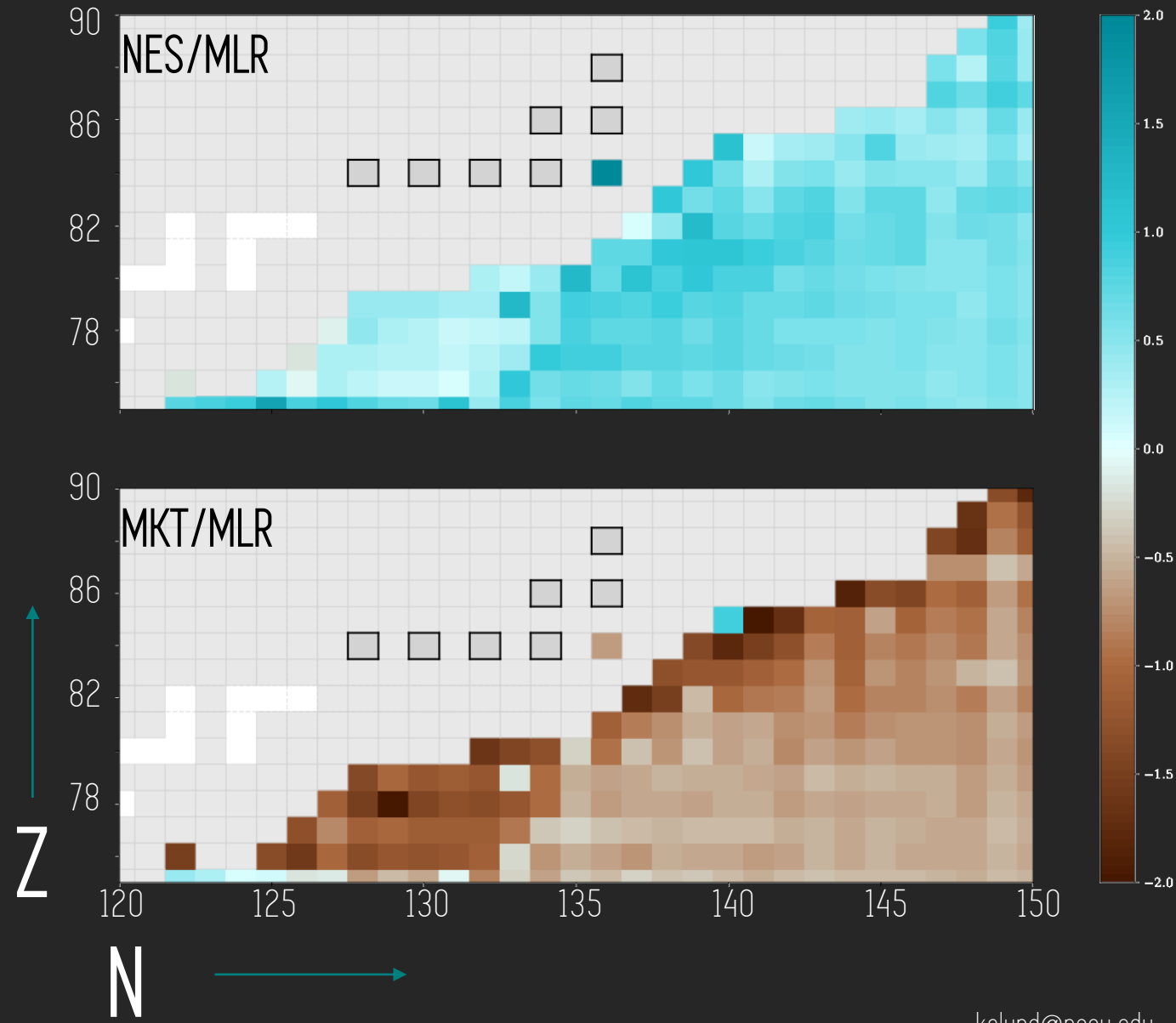
Diversity in heating sources changes
shape of light curves

Generally slower rates (NES) yield
more heating and brighter light curves

Nuclear Heating (revisited)



Alpha Decay Heating



Differences in beta decay rates affect heating from alpha heaters with measured decay times, especially:

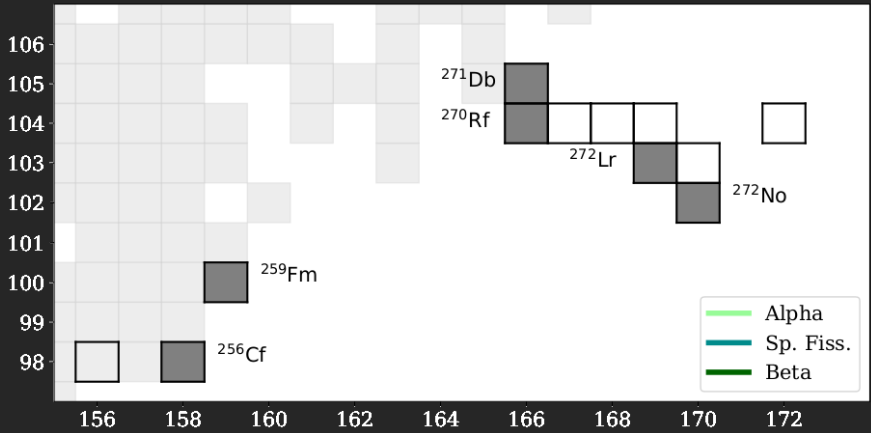
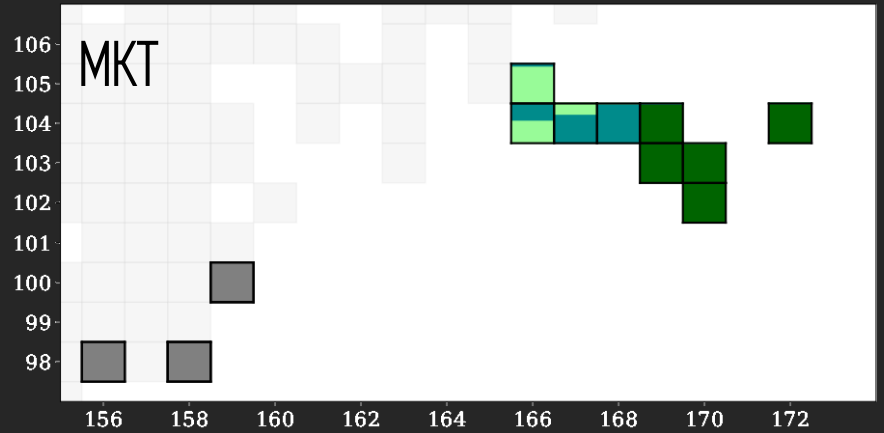
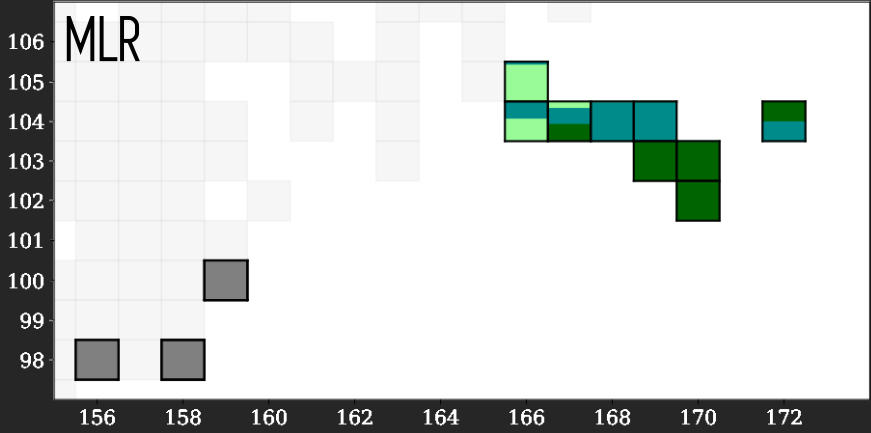
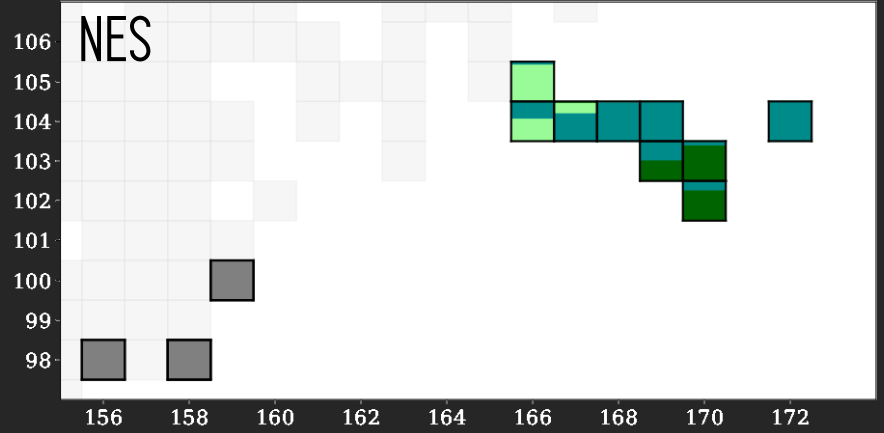


Spontaneous Fission (et al.) Heating

Theoretical branching ratios affect spontaneous fission heating

Z ↑

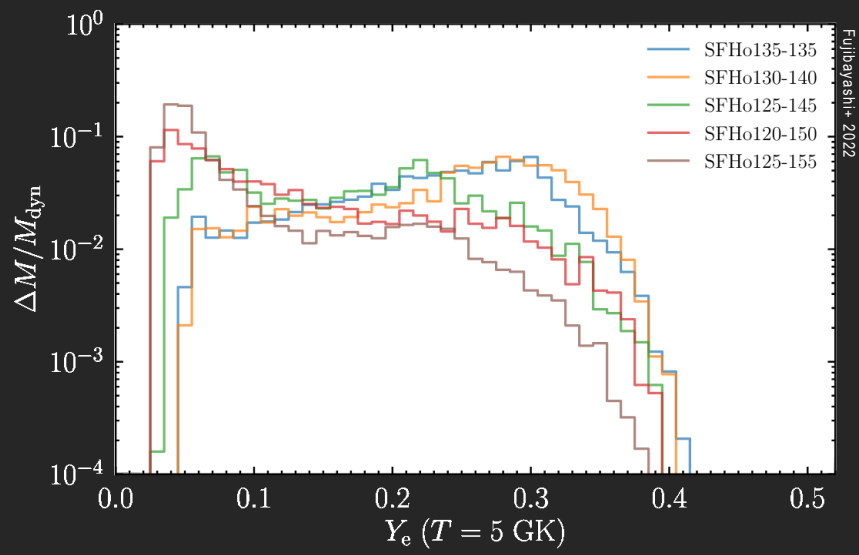
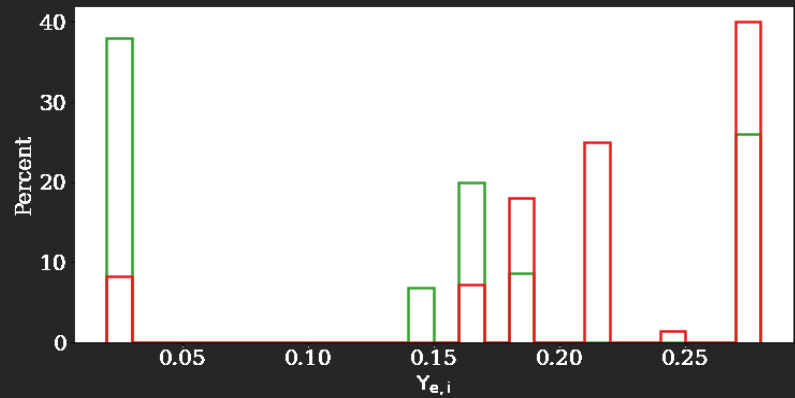
N →



Combined Trajectories (revisited)

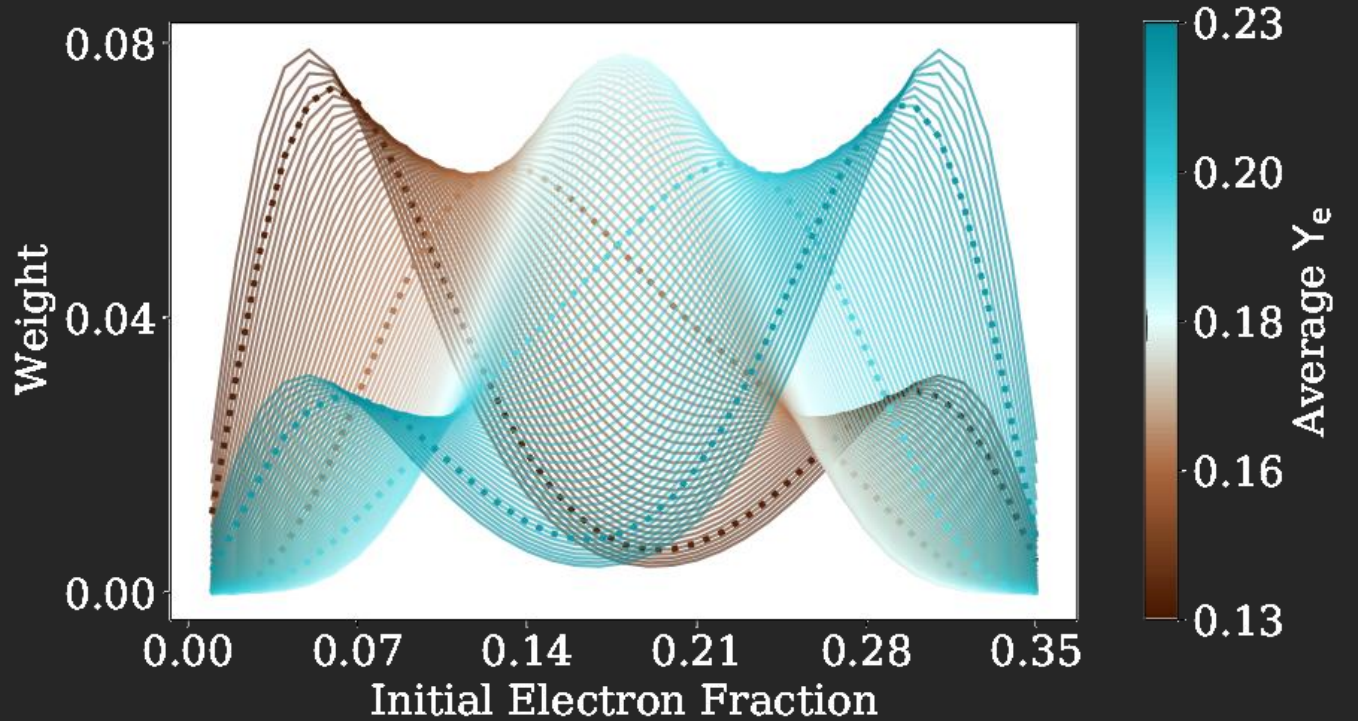
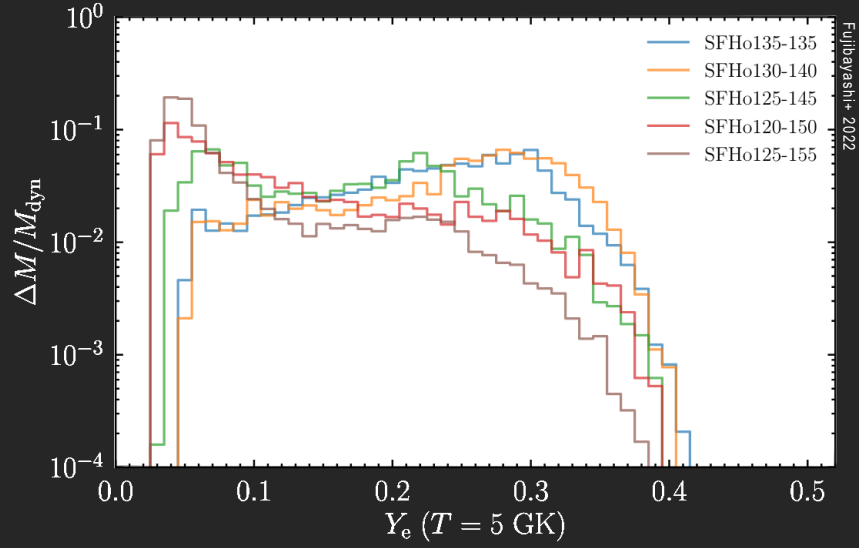
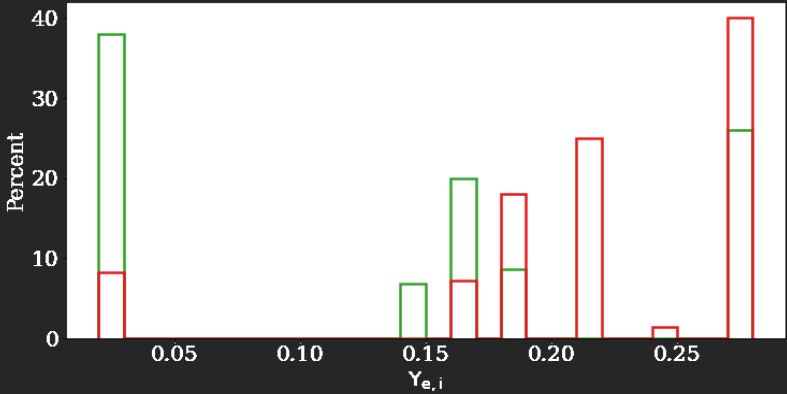


Multiple Y_e components



Fujibayashi+ 2022

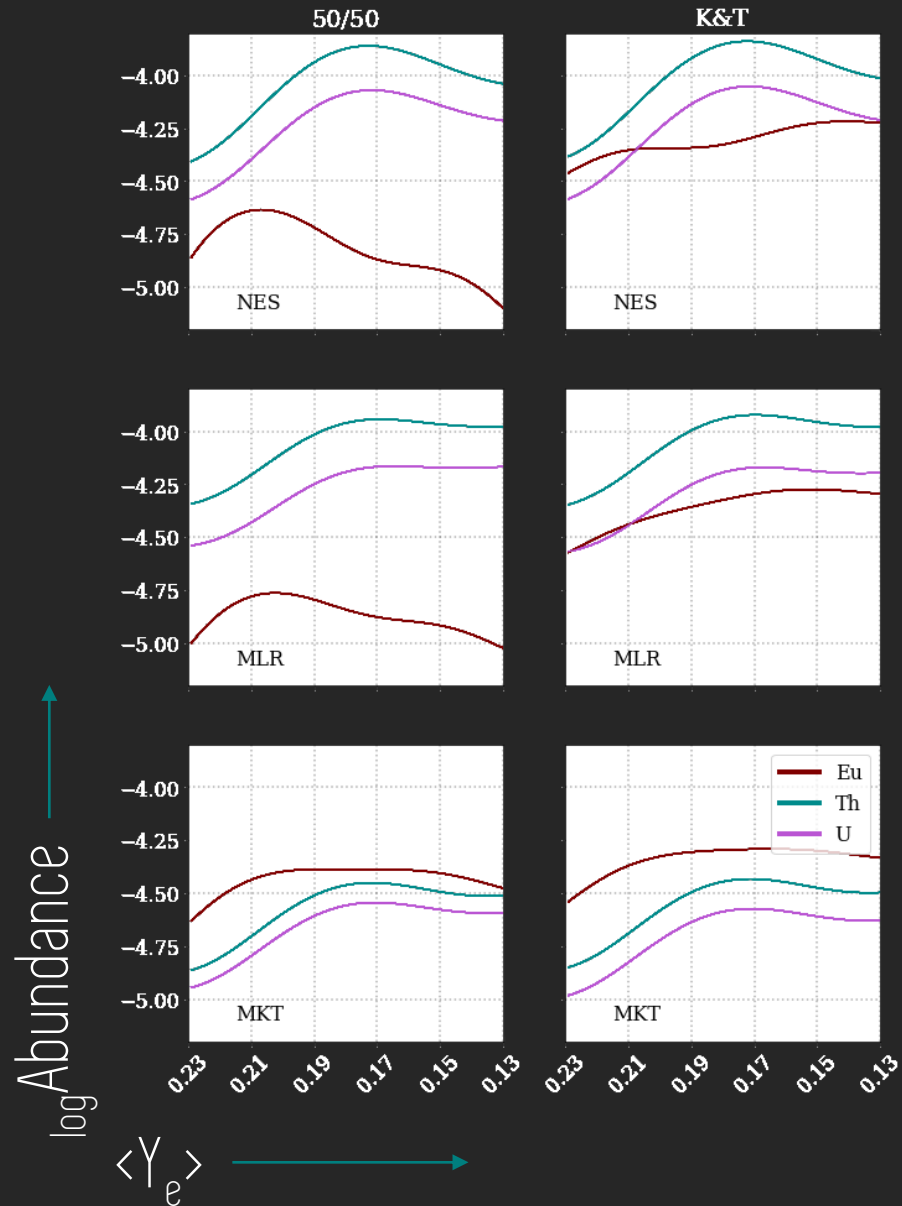
Multiple Y_e components



Increasing actinide enhancement ↓

Star Name	$\log_{\epsilon}(\text{Eu})$	$\log_{\epsilon}(\text{Th})$	$\log_{\epsilon}(\text{U})$	Reference
HE1523-0901	-0.62	-1.2	-2.06	Frebel+2007
CS29497-004	-0.66	-1.16	-2.20	Hill+2017
J2038-0023	-0.75	-1.24	-2.14	Placco+2017
CS31082-001	-0.72	-0.98	-1.92	Siquiera Mello+2013
J0954+5246	-1.19	-1.31	-2.13	Holmbeck+2018

Abundances for Cosmochronometry



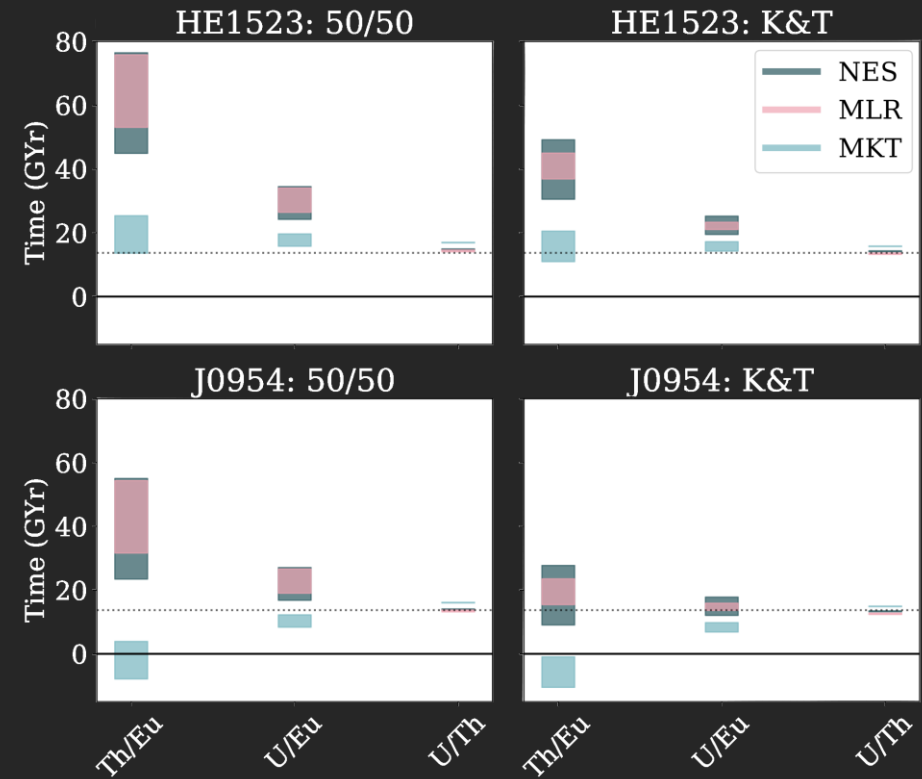
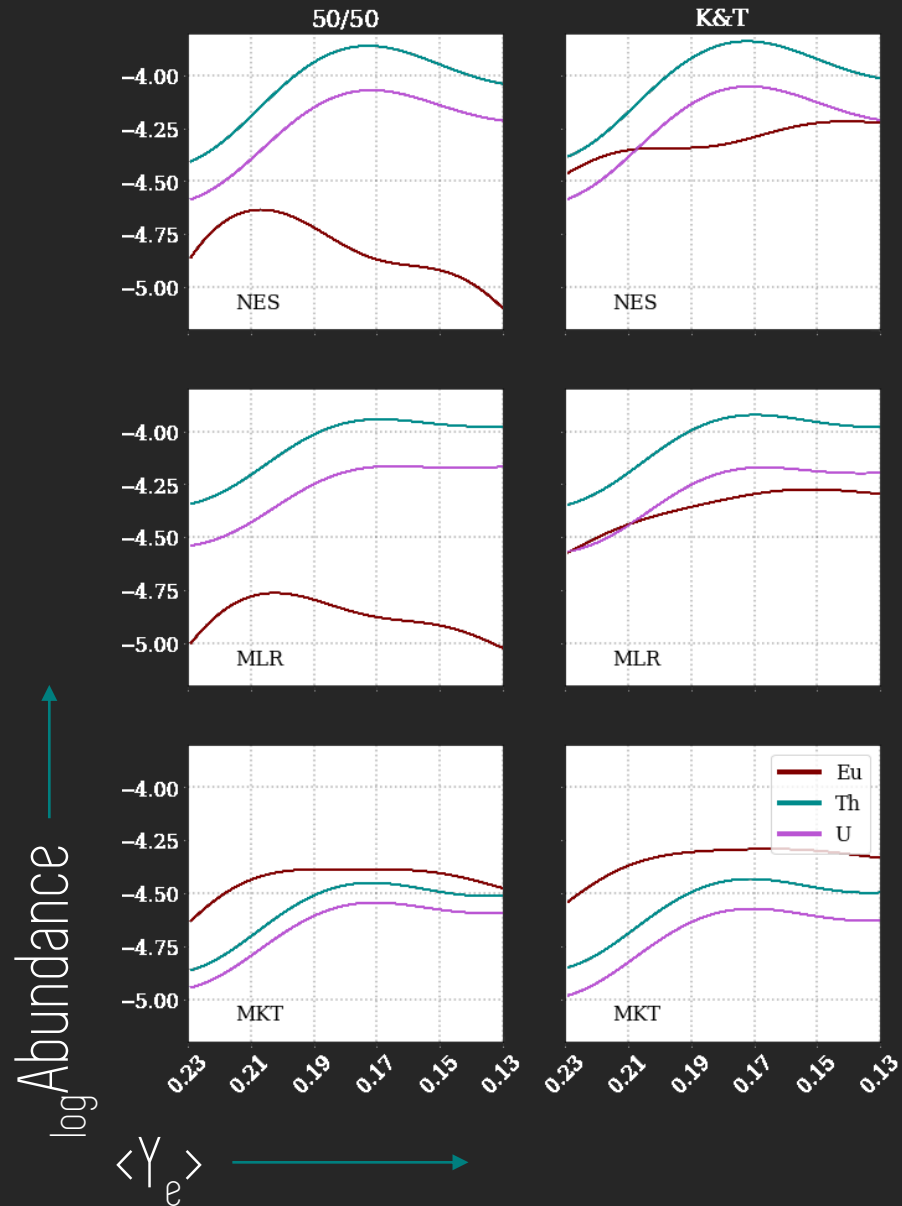
$$t = 46.67 \text{ Gyr} \left[-\log_{\epsilon} \left(\frac{\text{Th}}{\text{Eu}} \right)_{\text{obs}} + \log_{\epsilon} \left(\frac{\text{Th}}{\text{Eu}} \right)_0 \right]$$

$$t = 14.84 \text{ Gyr} \left[-\log_{\epsilon} \left(\frac{\text{U}}{\text{Eu}} \right)_{\text{obs}} + \log_{\epsilon} \left(\frac{\text{U}}{\text{Eu}} \right)_0 \right]$$

$$t = 21.80 \text{ Gyr} \left[-\log_{\epsilon} \left(\frac{\text{U}}{\text{Th}} \right)_{\text{obs}} + \log_{\epsilon} \left(\frac{\text{U}}{\text{Th}} \right)_0 \right]$$

^{232}Th & ^{238}U : produced exclusively via r-process ($t_{1/2} = 14 \text{ Gyr}, 4.486 \text{ Gyr}$ respectively)

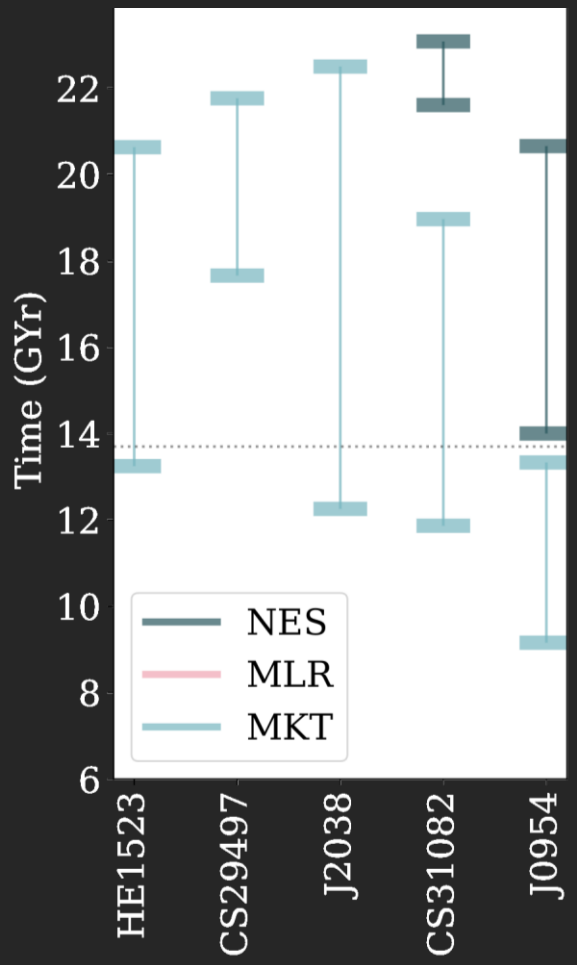
Abundances for Cosmochronometry



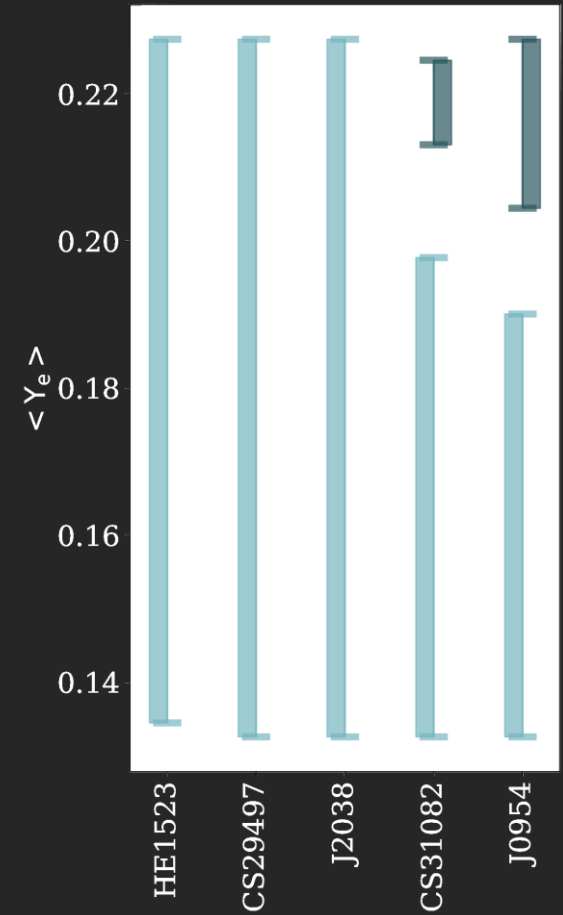
- Europium production highly sensitive to average Y_e and fission yield
- Underabundance of actinides can lead to negative age predictions

Error Bars from Observations

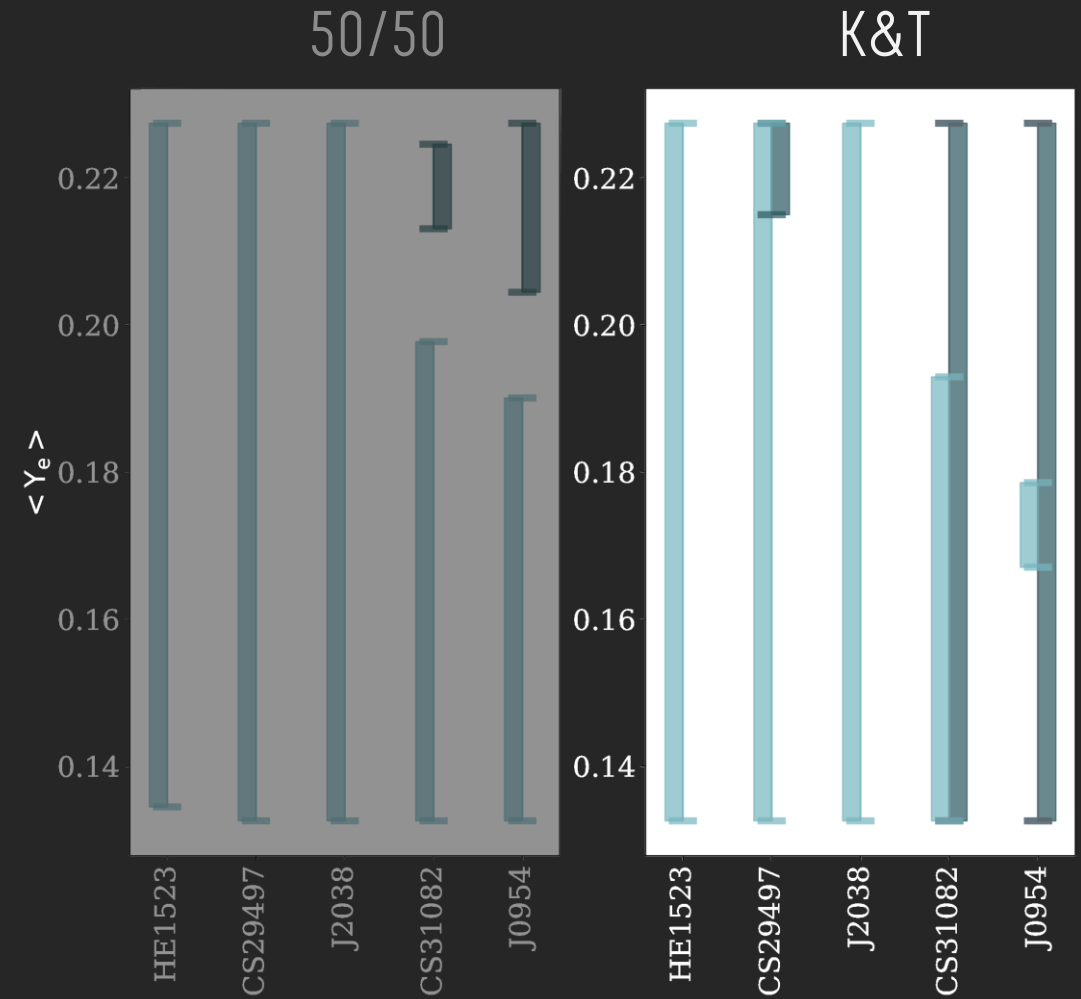
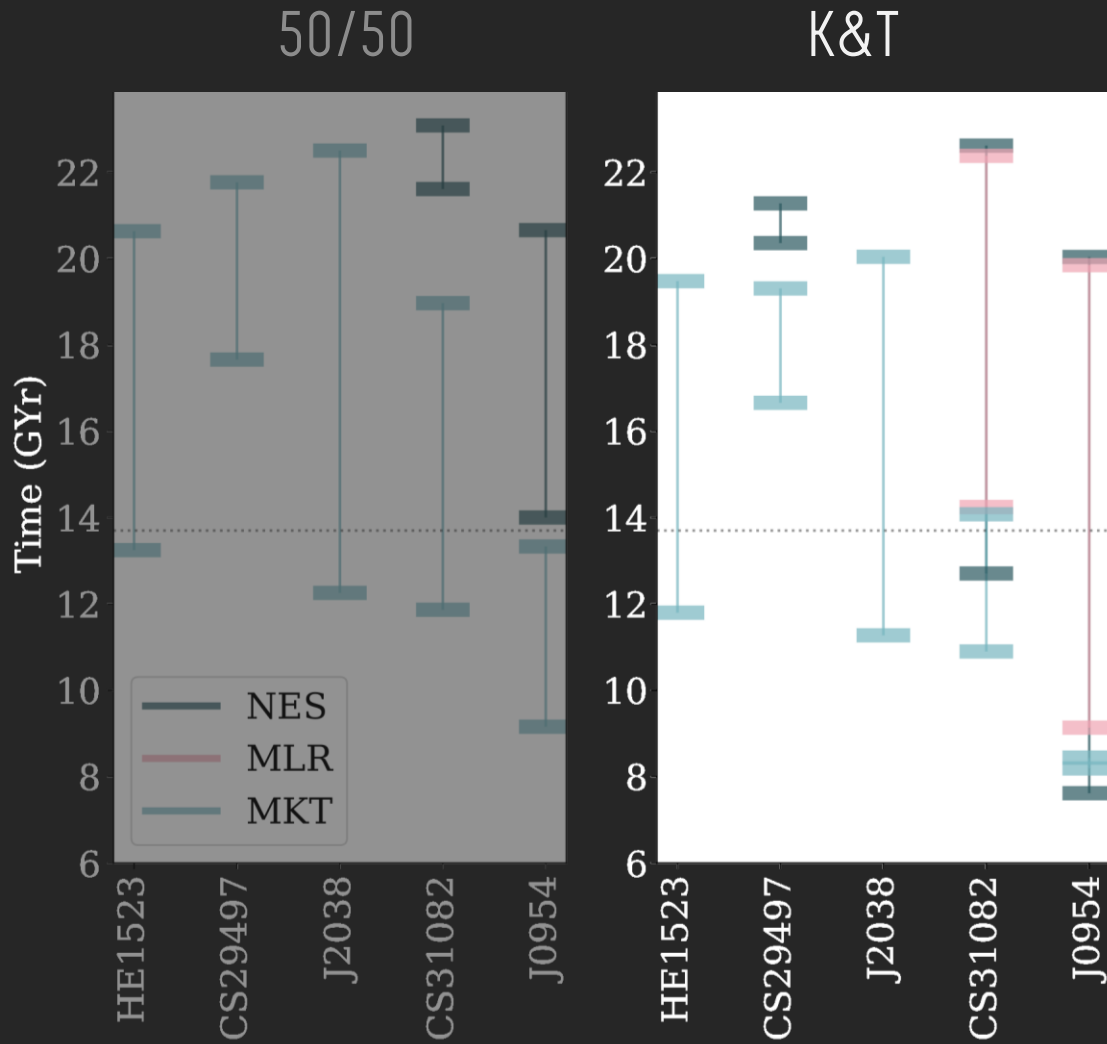
50/50



50/50



Error Bars from Observations



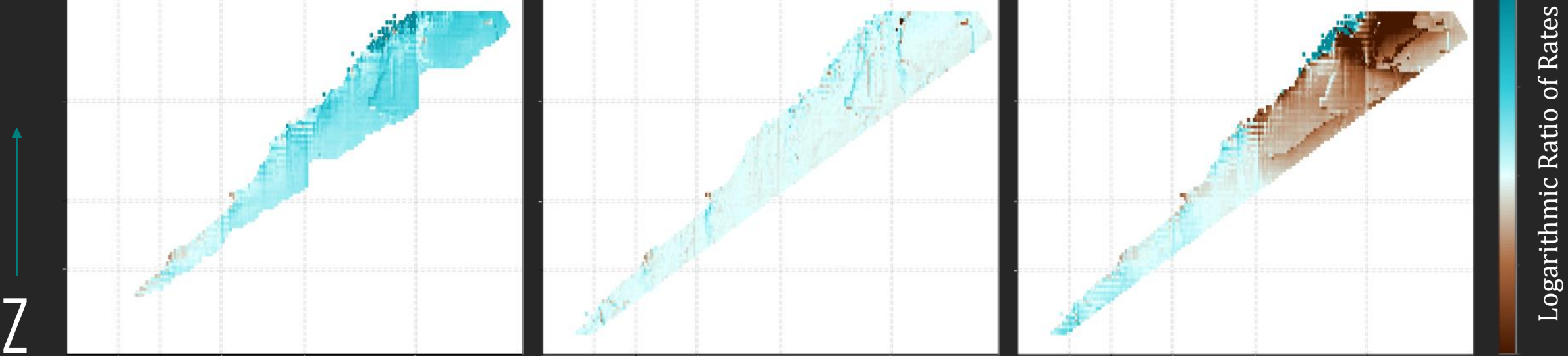
There is a wealth of physics in the unknown properties of nuclei far from stability that can impact key kilonova related quantities

- Identified key measured and unmeasured nuclei important for nuclear heating on light curve-relevant time scales.
- Explored a variety of theoretical nuclear models as a source of uncertainty for nuclear energy generation.
- Probed lanthanide/actinide abundances for cosmic dating of r-process enhanced metal-poor stars.

Nucleosynthesis Appendix

Appendix: Beta Decay Rates

Logarithmic ratio compared to Möller+ 2003



N →
Ney (NES)
Ney+ 2020

Möller (MLR)
Möller+ 2019

Marketin (MKT)
Marketin+ 2016

FAM+QRPA

QRPA+FRDM

Covariant DFT

Appendix: Thermalization Efficiency

Kasen & Barnes 2019 ApJ [arXiv: 1807.03319]

Assume thermalization of massive particles: $f = (1 + \tau)^{-n}$

Electrons ($n=1$): $t_{th,\beta} = 12.9 M_{0.01}^{2/3} v_{0.2}^{-2}$ days

γ -rays: $t_{\gamma} = 0.3 M_{0.01}^{1/2} v_{0.2}^{-1}$

and

$$f_{\gamma}(t) = 1 - \exp\left[-\frac{t_{\gamma}^2}{t^2}\right]$$

α -decay: $t_{th,\alpha} = 2t_{th,\beta}$

and

$$f_{\alpha}(t) = \left(1 + \frac{t}{t_{th,\alpha}}\right)^{-1.5}$$

Fission: $t_{th,f} = 4t_{th,\beta}$

and

$$f_f(t) = \left(1 + \frac{t}{t_{th,f}}\right)^{-1}$$

Appendix: Opacity (i)

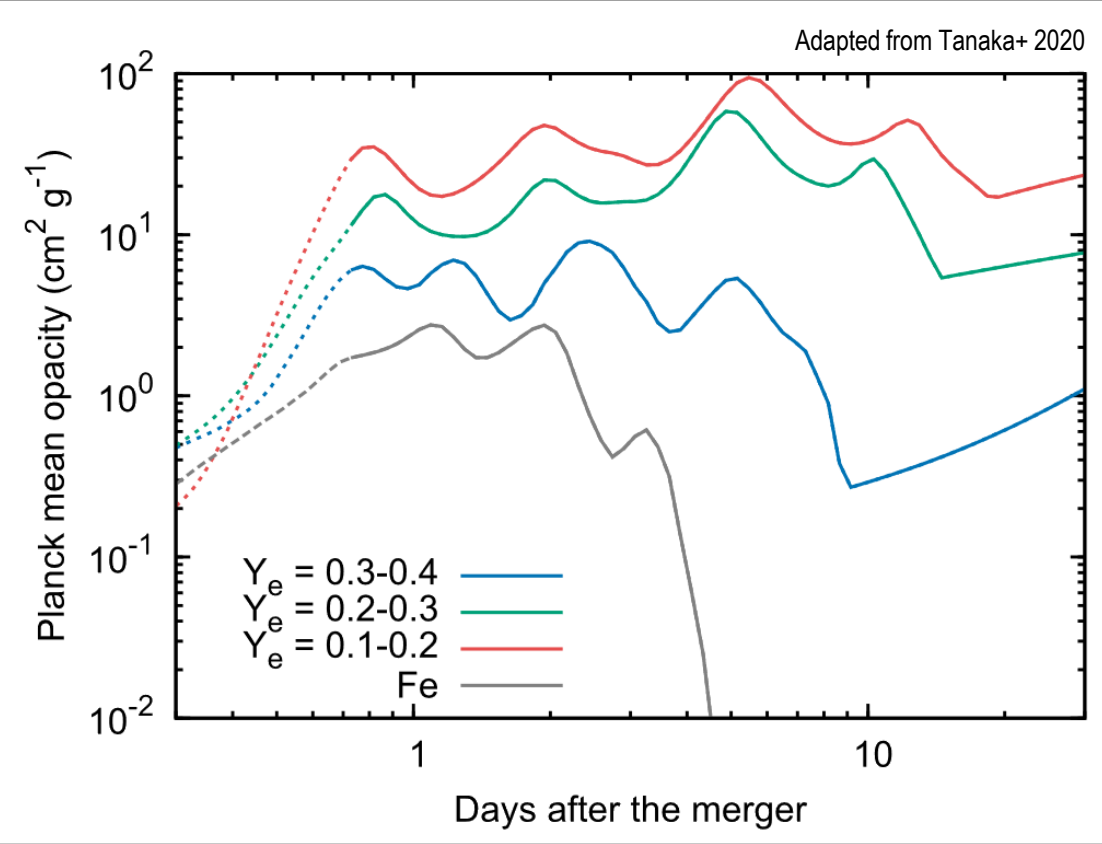
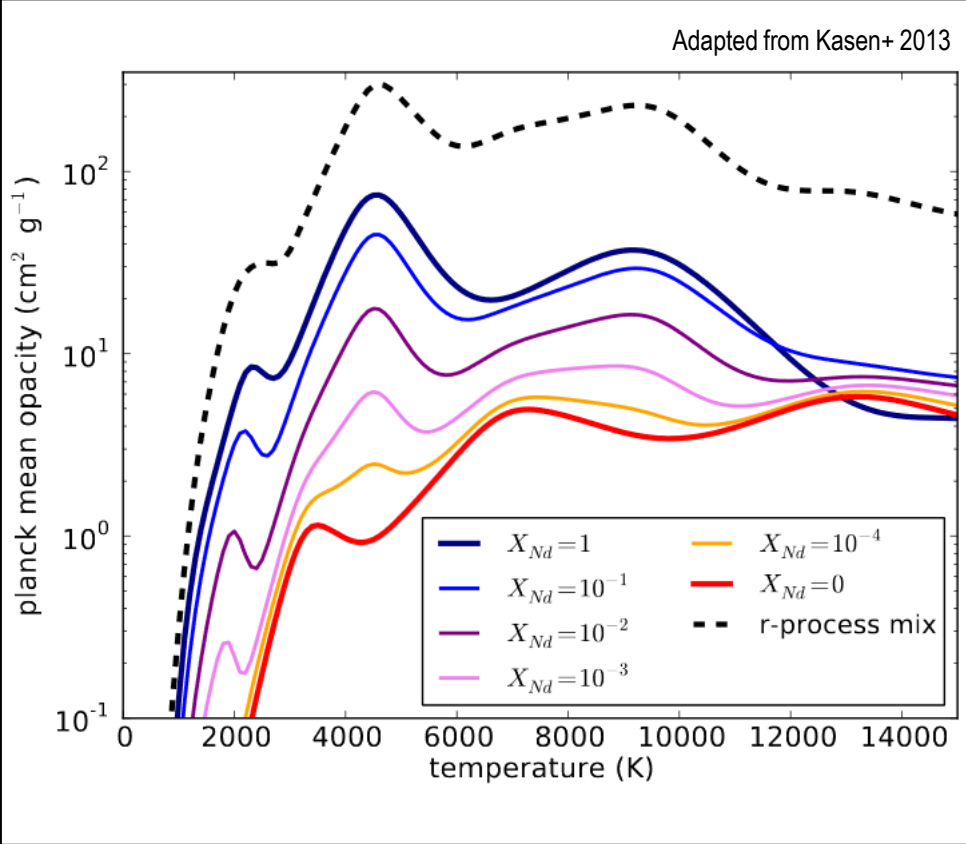
Diffusion timescale is dependent on opacity; in this model, the opacity is temperature dependent.

$$\kappa = \begin{cases} \kappa_{max} \left(\frac{T}{4000\text{K}} \right)^{5.5} & T < 4000\text{K} \\ \kappa_{max} & \text{otherwise} \end{cases} \quad \text{with } \kappa_{max} = 100 \text{ cm}^2\text{g}^{-1}$$

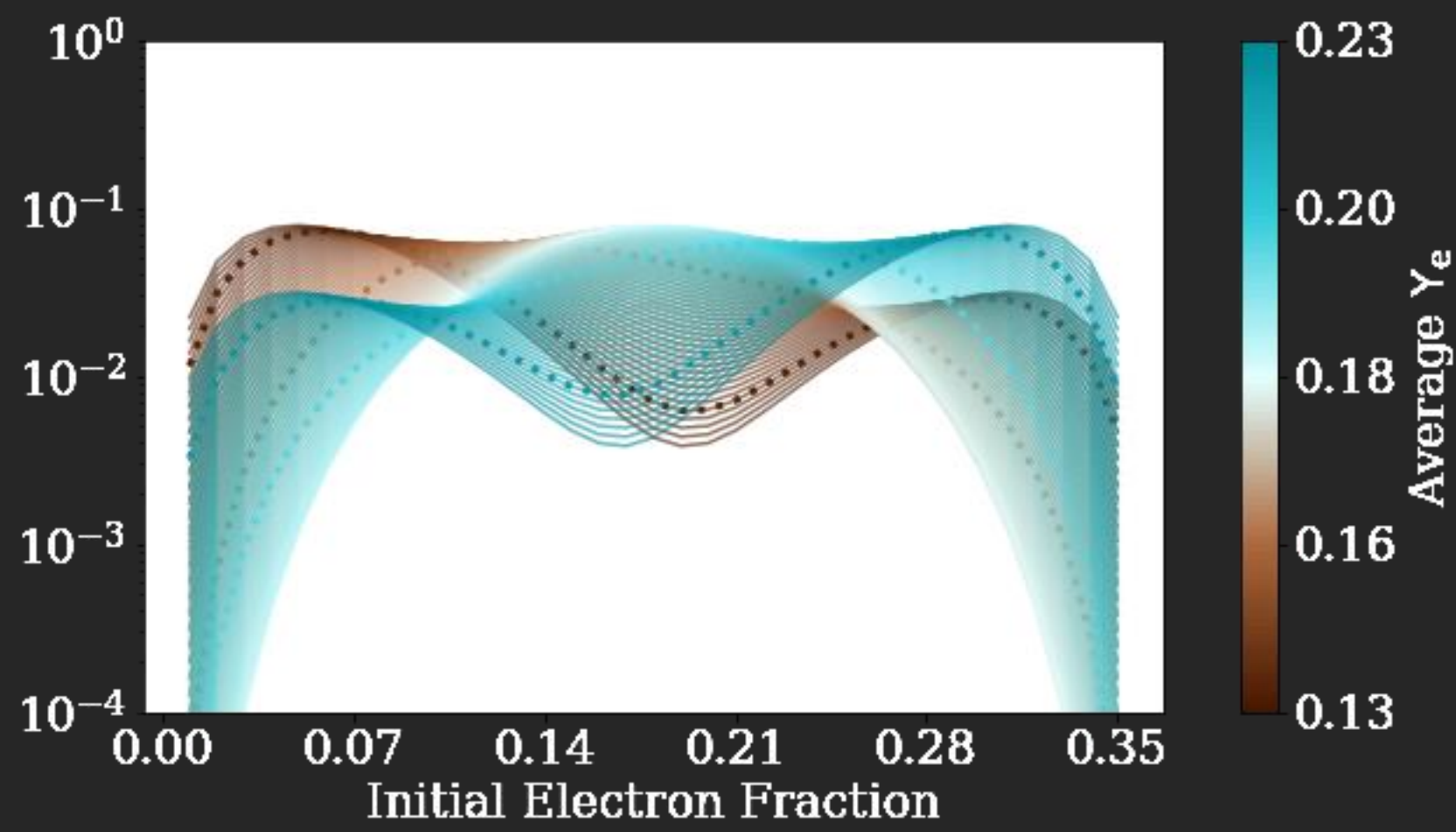
Technically, κ_{max} should be calculated uniquely for each composition. For now, consider mainly that a larger opacity should

- extend the light curve in time, and
- decrease the height of the light curve peak

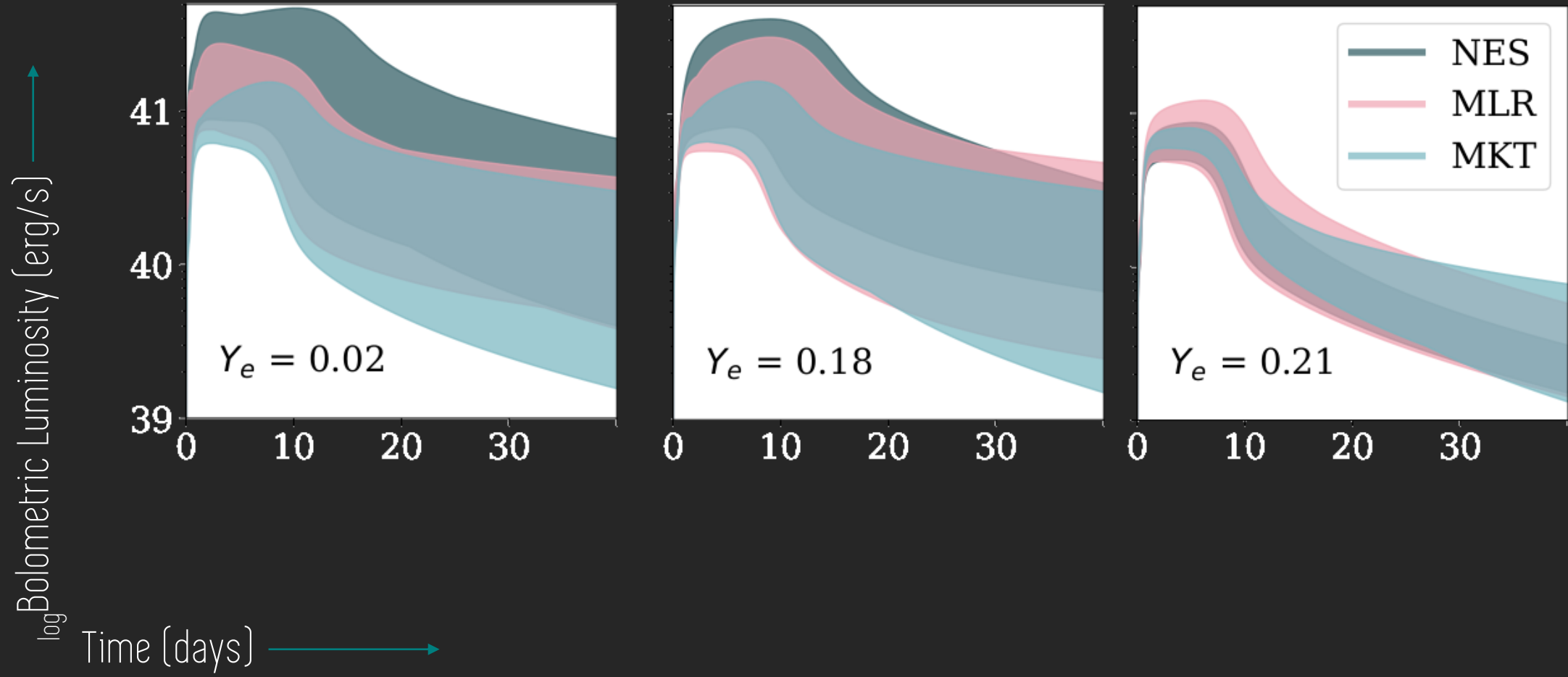
Appendix: Opacity (ii)



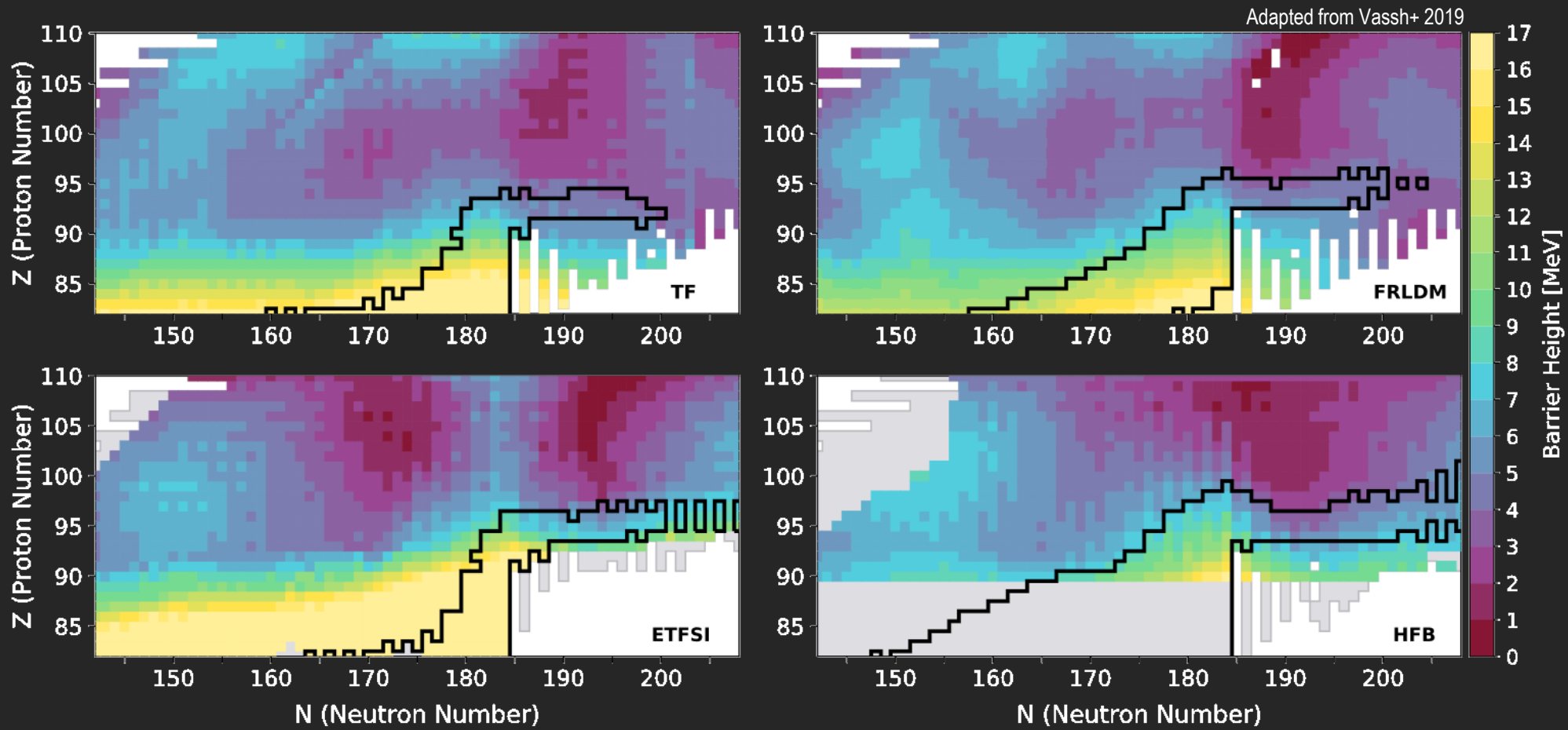
Appendix: Linear Combinations



Backup: Light Curves



Appendix: Fission Barriers



Appendix: Reaction Network

The evolution of the abundance of species i can be described by:

$$\frac{dY_i}{dt} = \underbrace{\sum_j N_j^i \lambda_j Y_j}_{\text{One-body}} + \underbrace{\sum_{j,k} N_{j,k}^i \rho N_A \langle j, k \rangle Y_j Y_k}_{\text{Two-body}} + \underbrace{\sum_{j,k,l} N_{j,k,l}^i \rho^2 N_A^2 \langle j, k, l \rangle Y_j Y_k Y_l}_{\text{Three-body}}$$

Decay Rate Thermal Cross-Section