Pionic atoms and chiral symmetry

Meson Science Laboratory, RIKEN
Kenta Itahashi

Nishi, KI et al., arXiv: 2204.05568

Pionic atom unveils hidden structure of QCD vacuum

Takahiro Nishi[1], Kenta Itahashi[1,*] DeukSoon Ahn[1,2], Georg P.A. Berg[3], Masanori Dozono[1], Daijiro Etoh[4], Hiroyuki Fujioka[5], Naoki Fukuda[1], Nobuhisa Fukunishi[1], Hans Geissel[6], Emma Haettner[6], Tadashi Hashimoto[1], Ryugo S. Hayano[7], Satoru Hirenzaki[8], Hiroshi Horii[7], Natsumi Ikeno[9], Naoto Inabe[1], Masahiko Iwasaki[1], Daisuke Kameda[1], Keichi Kisamori[10], Yu Kiyokawa[10], Toshiyuki Kubo[1], Kensuke Kusaka[1], Masafumi Matsushita[10], Shin'ichiro Michimasa[10], Go Mishima[7], Hiroyuki Miya[1], Daichi Murai[1], Hideko Nagahiro[8], Megumi Niikura[7], Naoko Nose-Togawa[11], Shinsuke Ota[10], Naruhiko Sakamoto[1], Kimiko Sekiguchi[4], Yuta Shiokawa[4], Hiroshi Suzuki[1], Ken Suzuki[12], Motonobu Takaki[10], Hiroyuki Takeda[1], Yoshiki K. Tanaka[1], Tomohiro Uesaka[1], Yasumori Wada[4], Atomu Watanabe[4], Yuni N. Watanabe[7], Helmut Weick[6], Hiroki Yamakami[5], Yoshiyuki Yanagisawa[1], and Koichi Yoshida[1]

Pionic atoms and chiral symmetry

- **Dominant symmetry of the vacuum in low-energy QCD.**
- **Spontaneous breakdown due to the non-perturbative nature of the strong interaction.**
- **Non-trivial structure of the QCD vacuum.**

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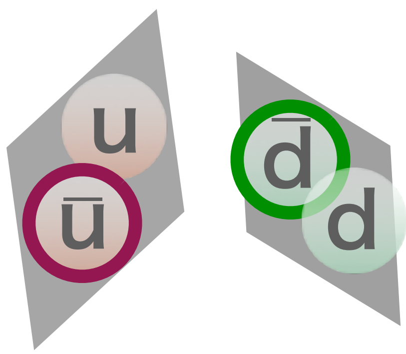
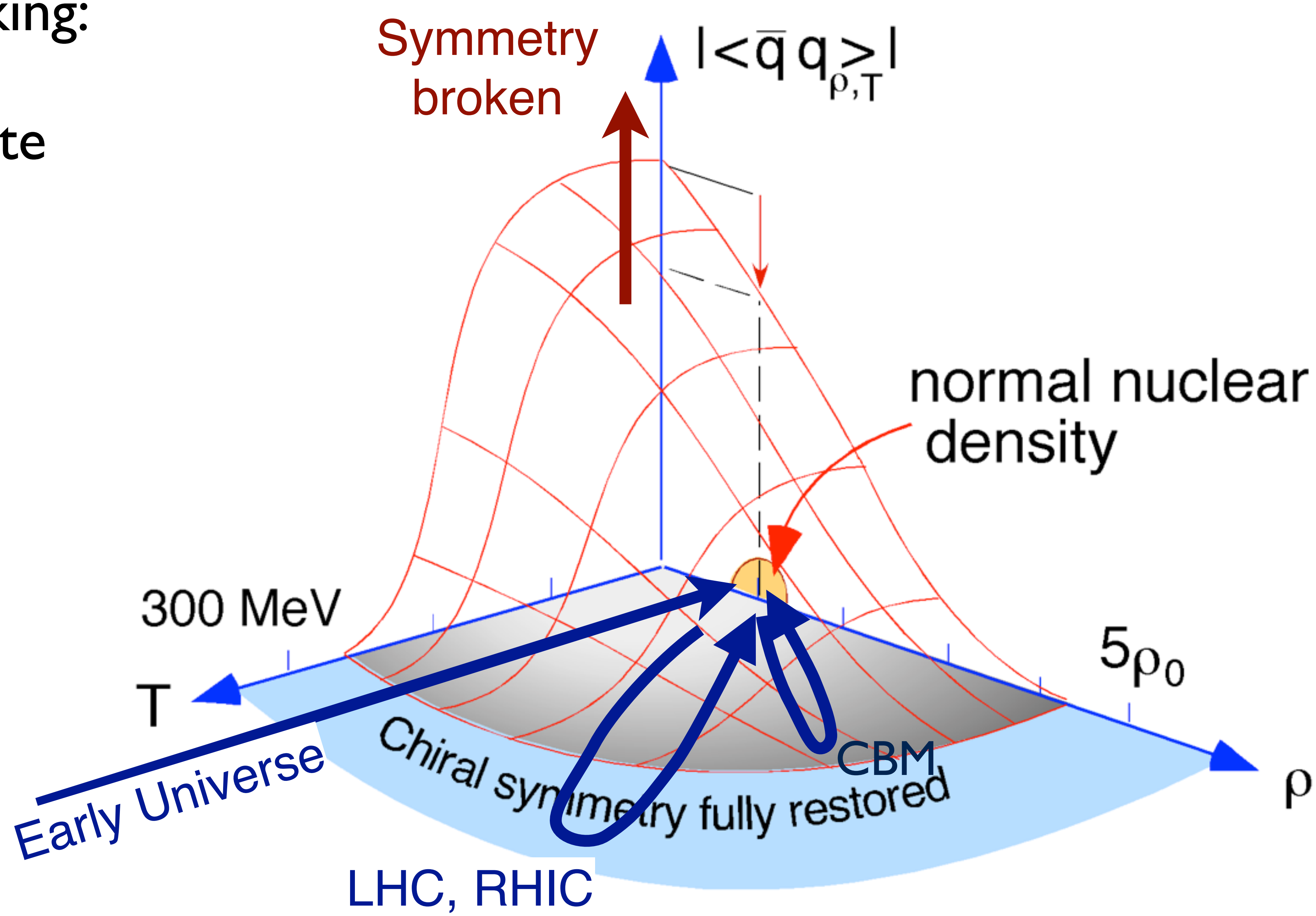
Chiral condensate, order parameter of chiral symmetry

One of order parameters of
χ-symmetry breaking:

Chiral condensate

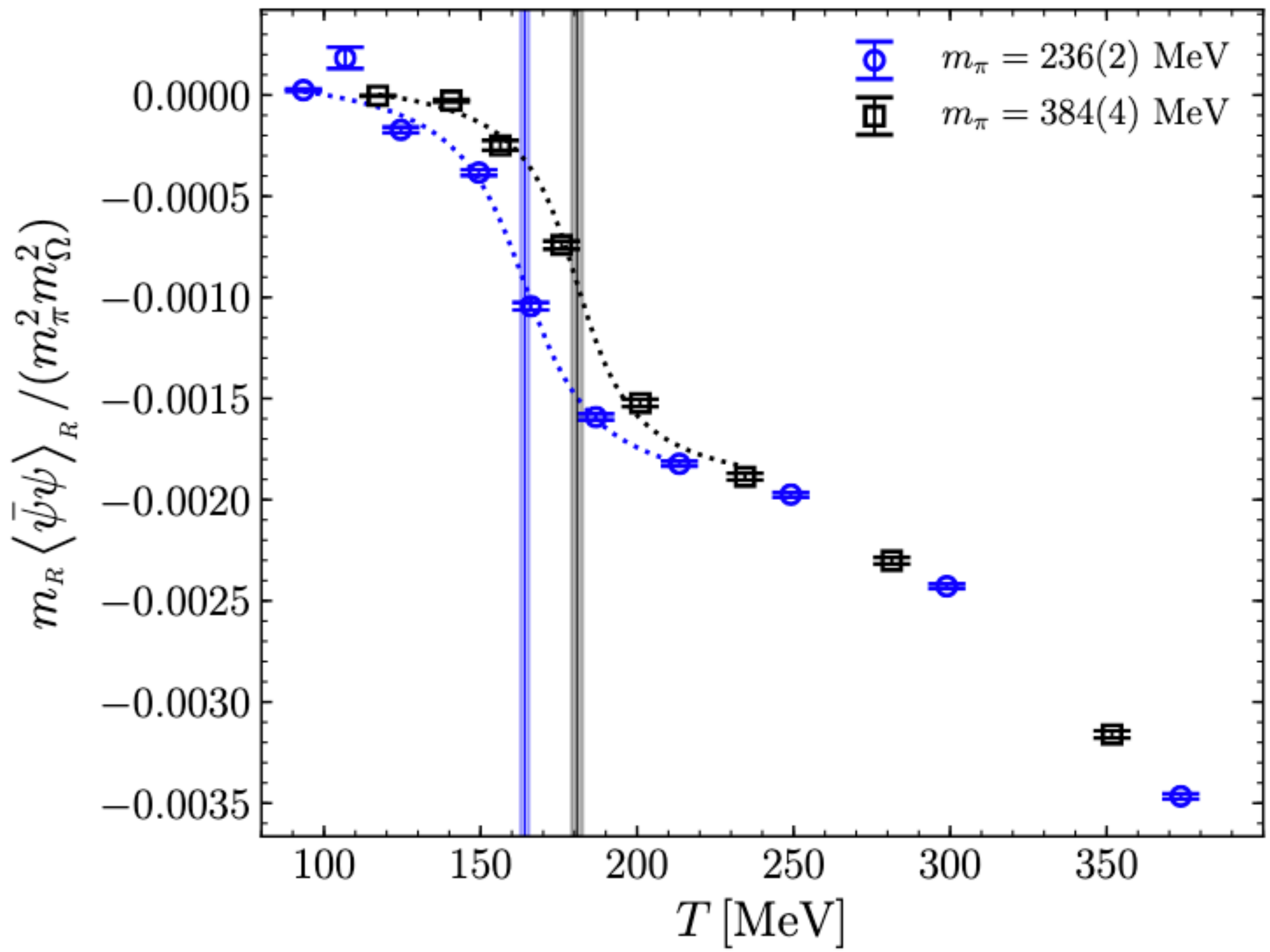




Analysis of material properties
of QCD vacuum

W.Weise,
NPA553(93)59.

Lattice QCD calculated T dependence of chiral condensate



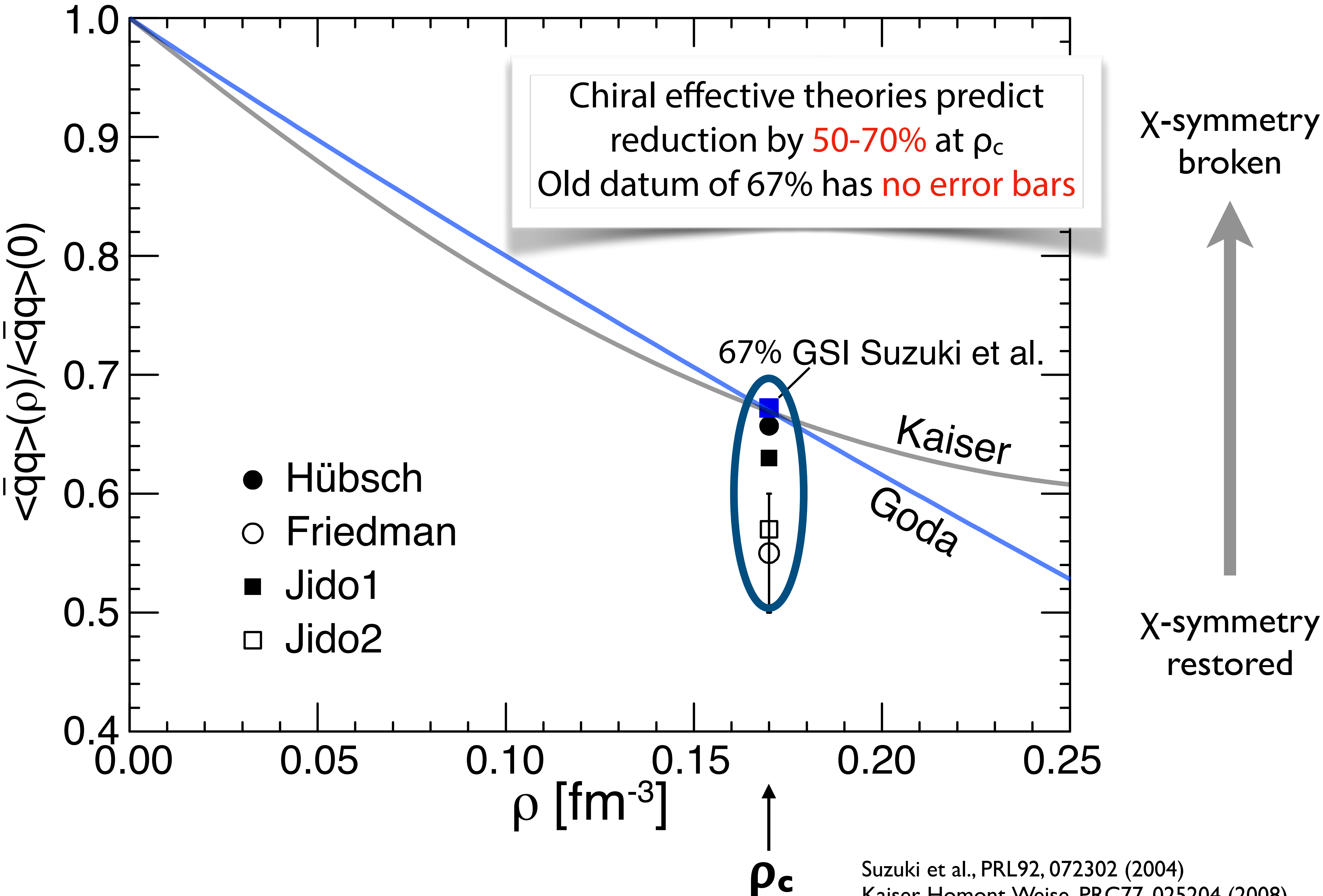
χ-symmetry broken

χ-symmetry restored

Remark: sign problem makes it difficult for lattice to approach non-zero ρ region

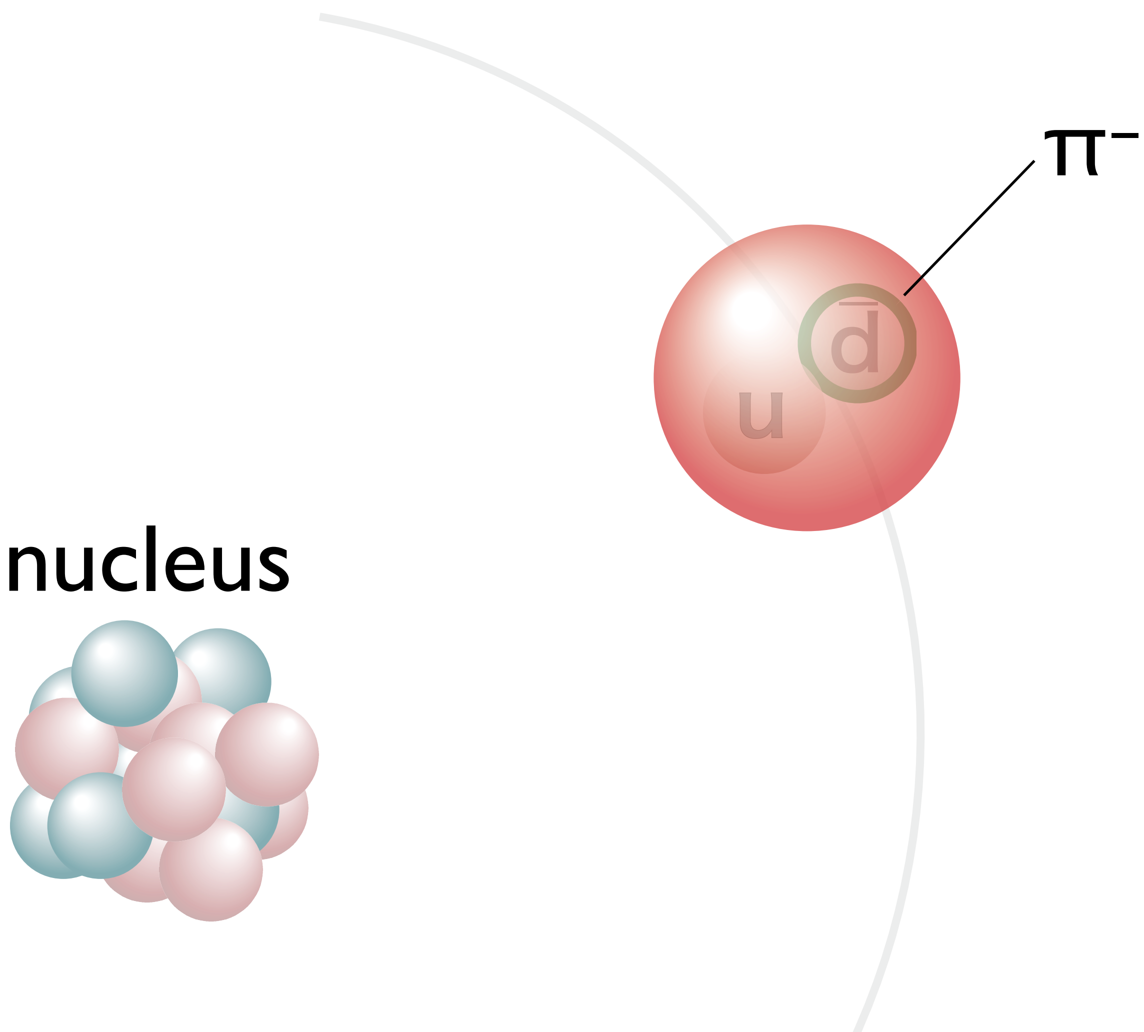
Jon-Ivar Skullerud
PRD105(2022)034504

ρ dependence of chiral condensate



Suzuki et al., PRL92, 072302 (2004)
Kaiser, Homont, Weise, PRC77, 025204 (2008)
Goda and Jido, PRC88, 065204 (2013)
Huebsch, Jido, PRC104, 015202 (2021)
Friedman, Gal, PLB792, 340 (2019)
Jido, Hatsuda, Kunihiro, PLB670, 109 (2008)

Pionic atoms

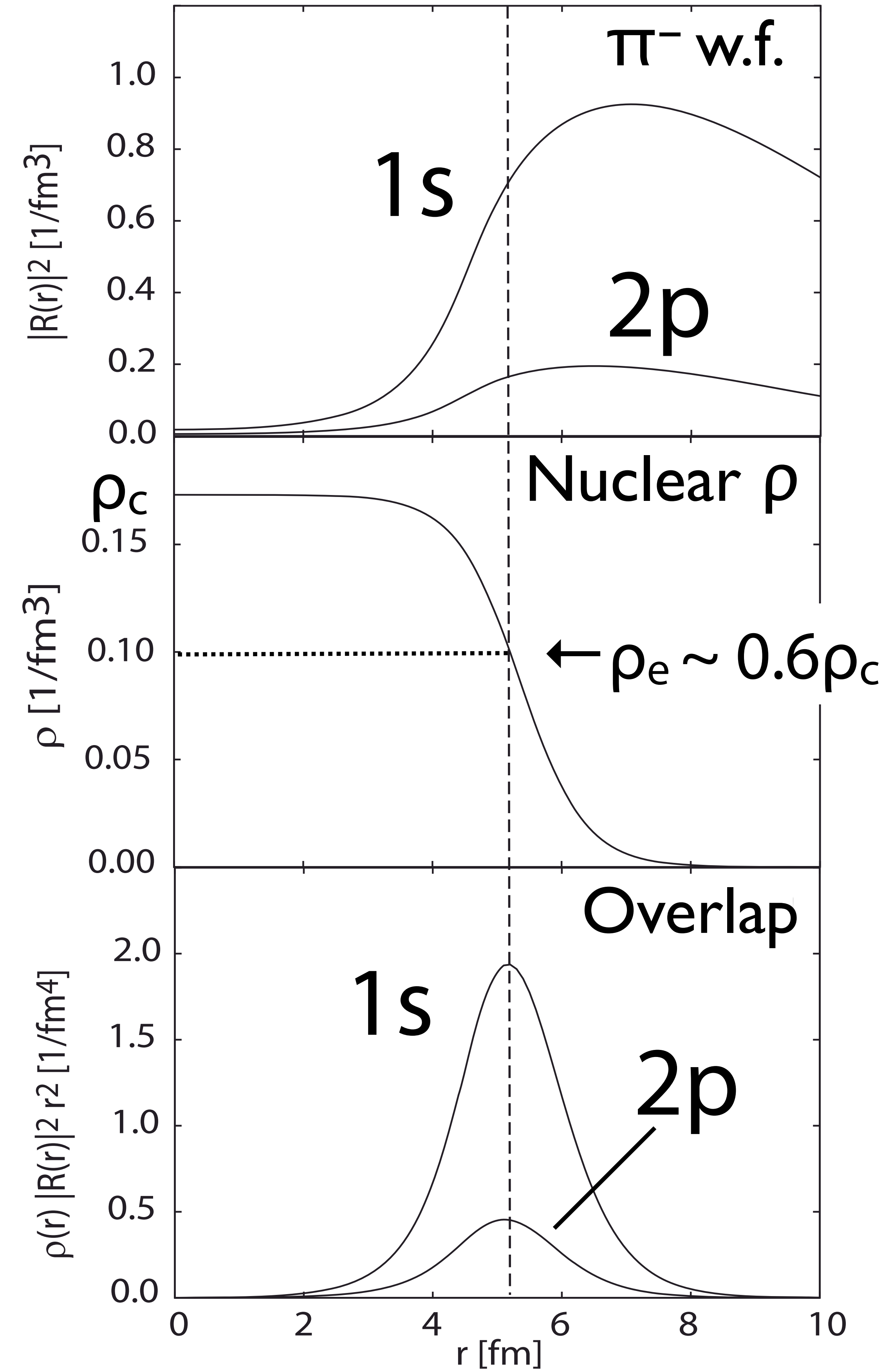


Ericson-Ericson potential

$$U_{\rm opt}(r) = U_s(r) + U_p(r),$$

$$U_s(r) = b_0\,\rho + \mathbf{b_1}\,(\rho_n - \rho_p) + B_0\,\rho^2$$

$$U_p(r) = \frac{2\pi}{\mu}\vec{\nabla}\cdot[c(r) + \varepsilon_2^{-1}C_0\rho^2(r)]L(r)\vec{\nabla}$$



Ikeno et al., PTP126 (2011) 483

Pion-nucleus interaction

Overlap between
pion w.f. and nucleus
→ π works as a probe
at $\rho_e \sim 0.6\rho_c$

π-nucleus interaction is changed
for wavefunction renormalization
of medium effect

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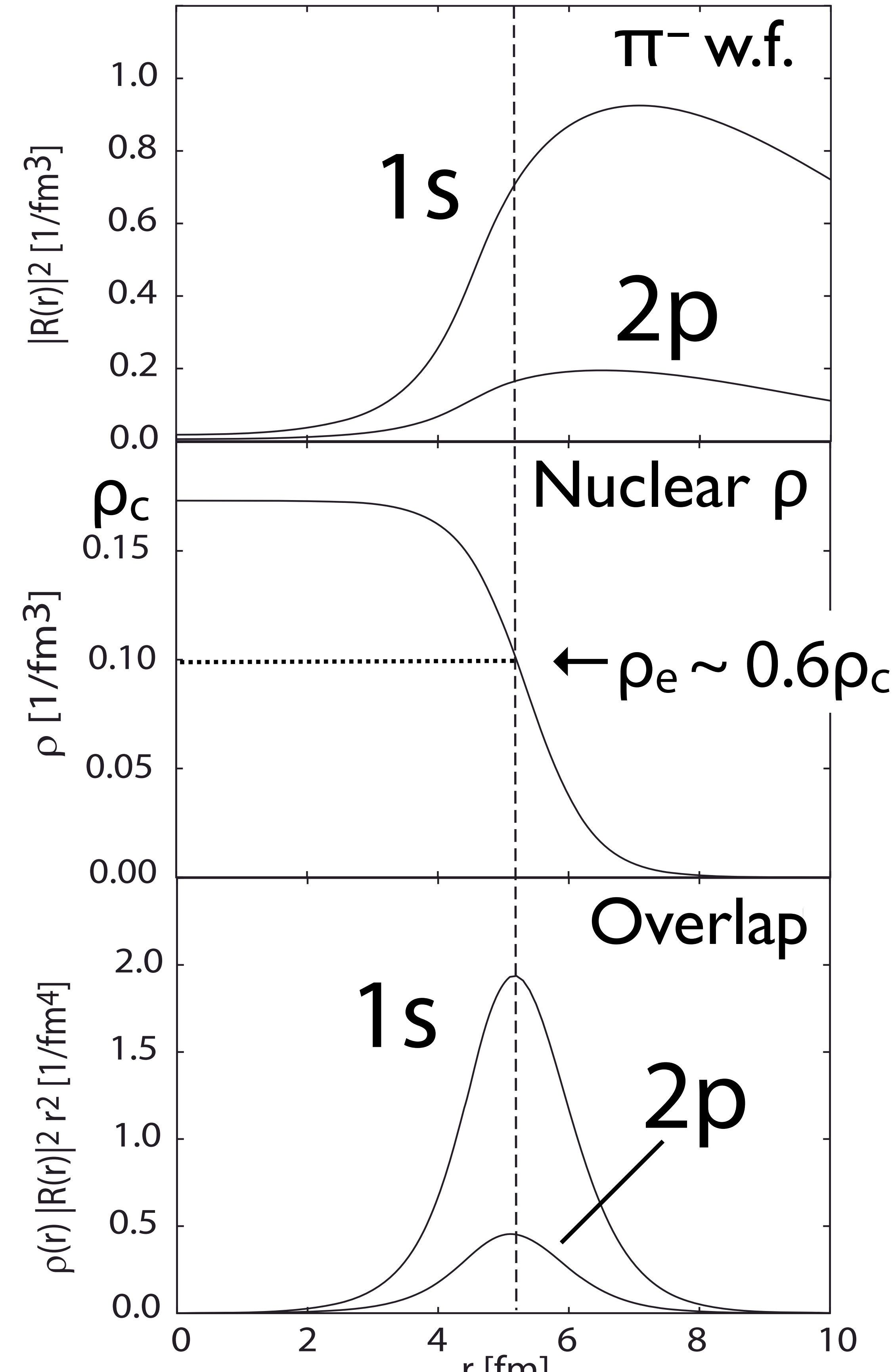

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In-medium Glashow-Weinberg relation

$$\frac{\langle\bar{q}q\rangle^*}{\langle\bar{q}q\rangle} \simeq \left(\frac{b_1}{b_1^{\rm v}}\right)^{1/2}\left(1-\gamma\,\frac{\rho}{\rho_0}\right)$$

$\gamma = 0.184 \pm 0.003$

Jido, Hatsuda, Kunihiro, PLB670, 109 (2008)

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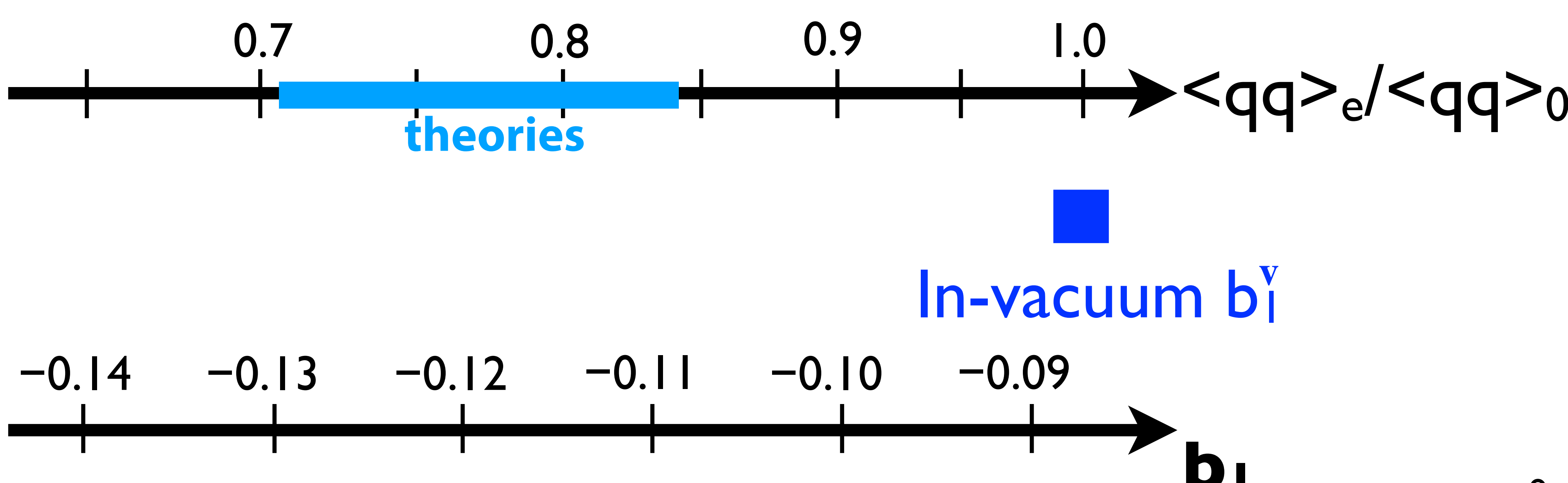
$\gamma = 0.184 \pm 0.003$

Jido, Hatsuda, Kunihiro, PLB670, 109 (2008)

Pionic hydrogen and deuterium

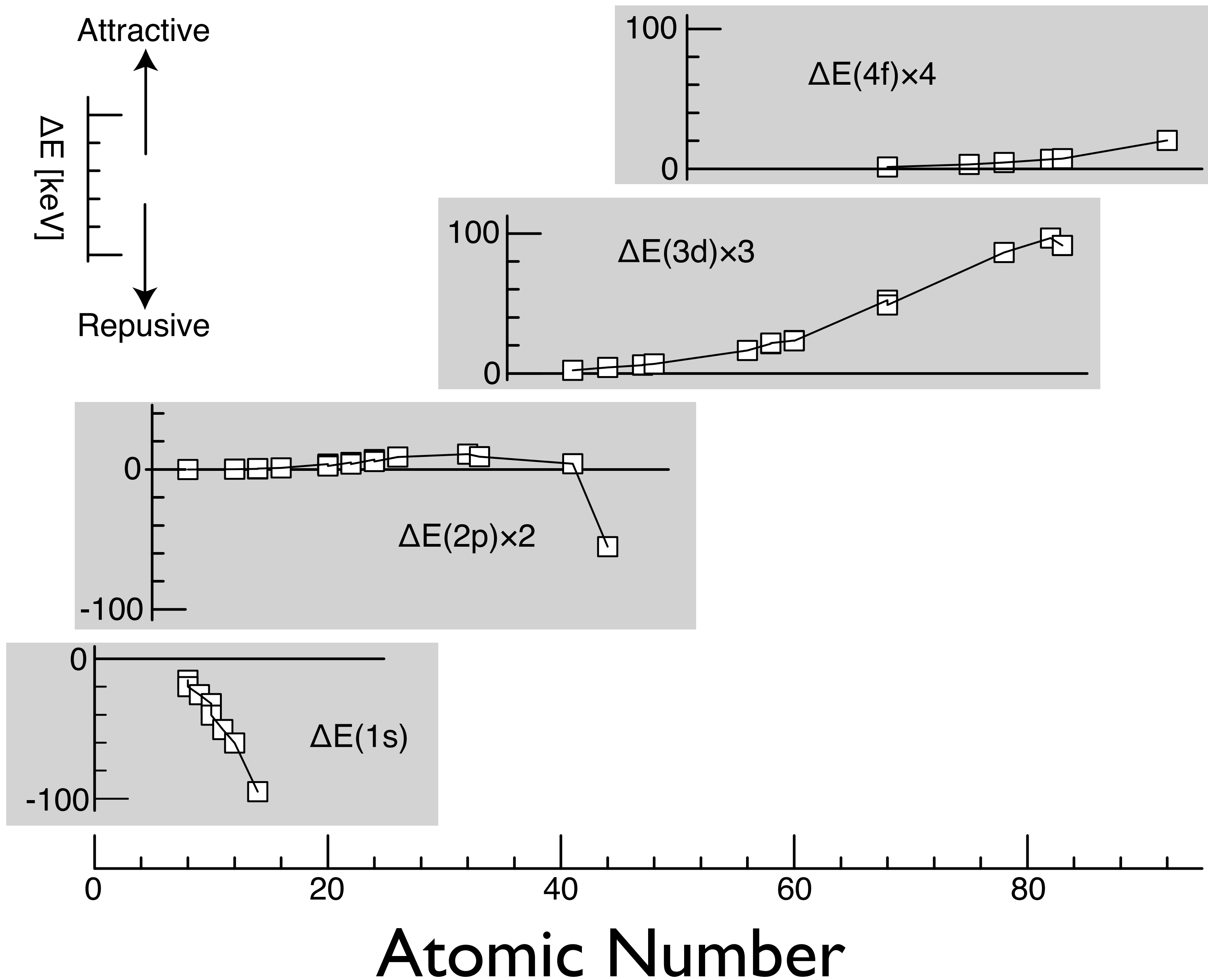
$b_1^{\mathrm{v}} = 0.0882 \pm 0.0014 \pm 0.0006$

Hirtl et al., EPJA57, 70 (2021)



Isovector
interaction

Level shifts in pionic X-ray measurements



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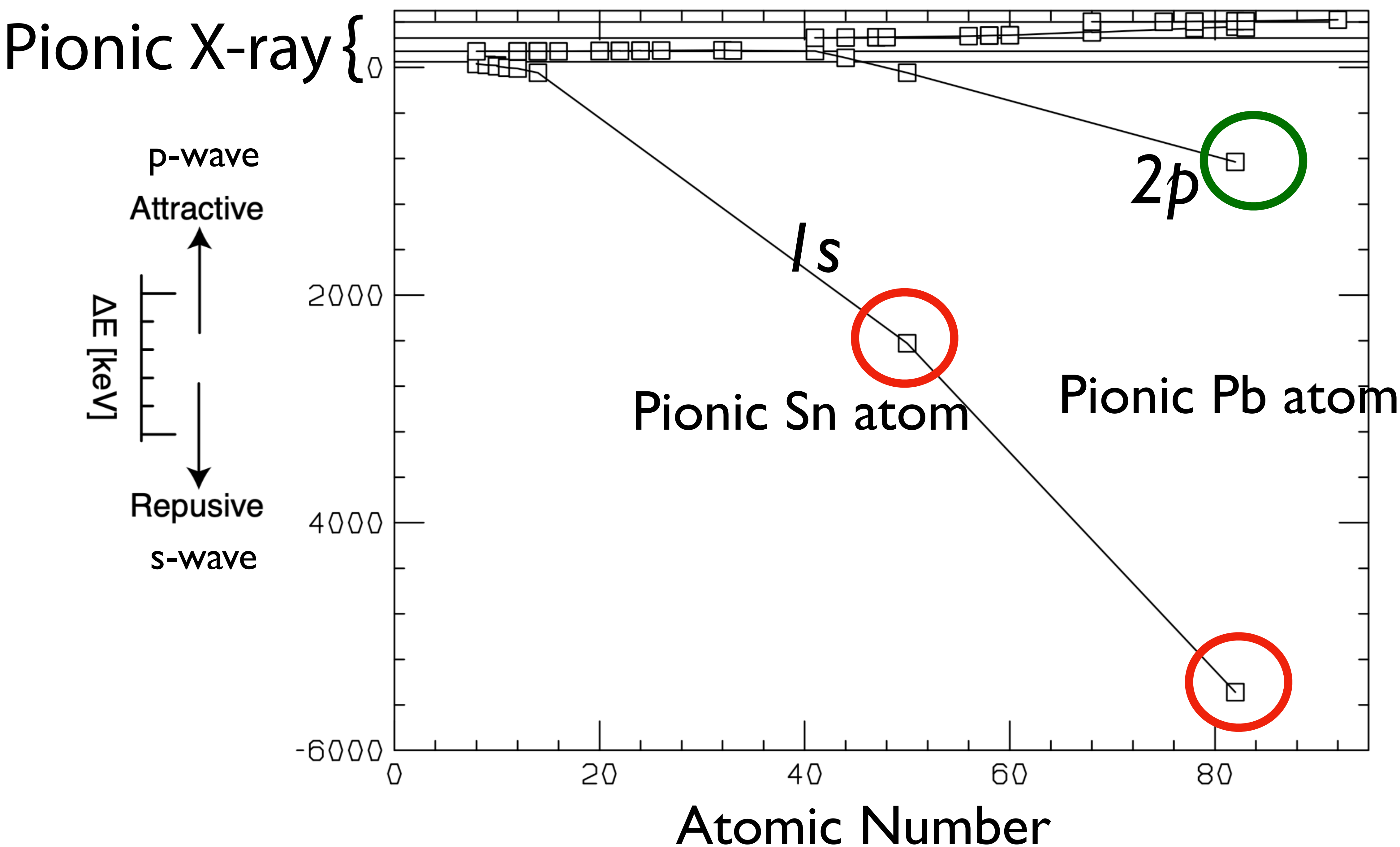
$$U_{\rm opt}(r) = U_s(r) + U_p(r),$$

$$U_s(r) = b_0\,\rho + \mathbf{b_1}\,(\rho_n - \rho_p) + B_0\,\rho^2$$ → s-wave = repulsive = negative shift

$$U_p(r) = \frac{2\pi}{\mu}\vec{\nabla}\cdot[c(r) + \varepsilon_2^{-1}C_0\rho^2(r)]L(r)\vec{\nabla}$$ → p-wave = attractive = positive shift

Deeply bound pionic atoms
Level shifts

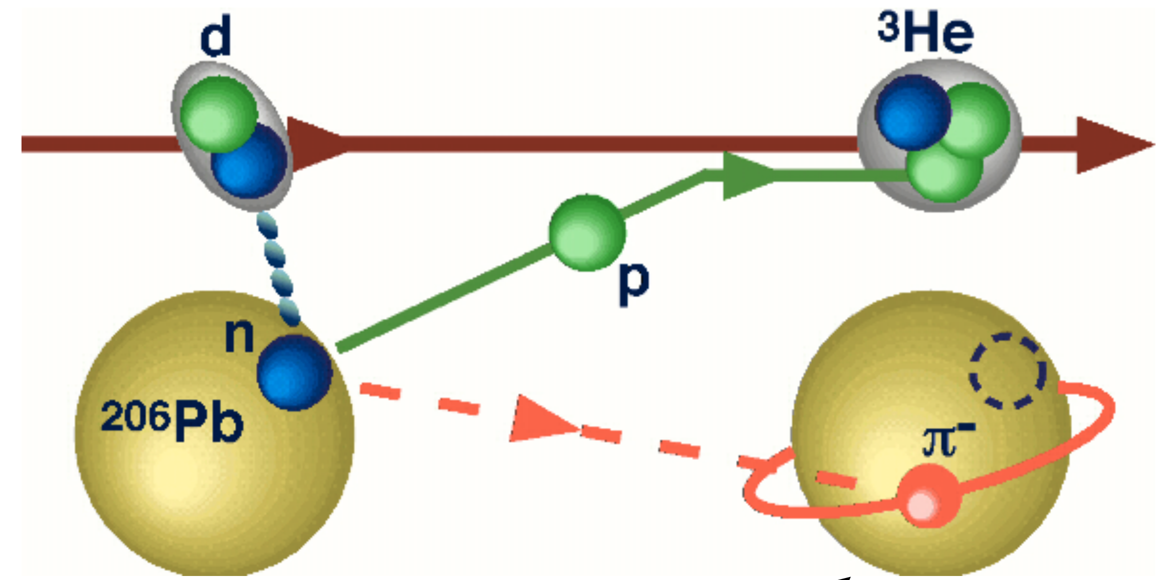


Deeply bound atoms have "super" repulsive shifts and provide *s*-wave information

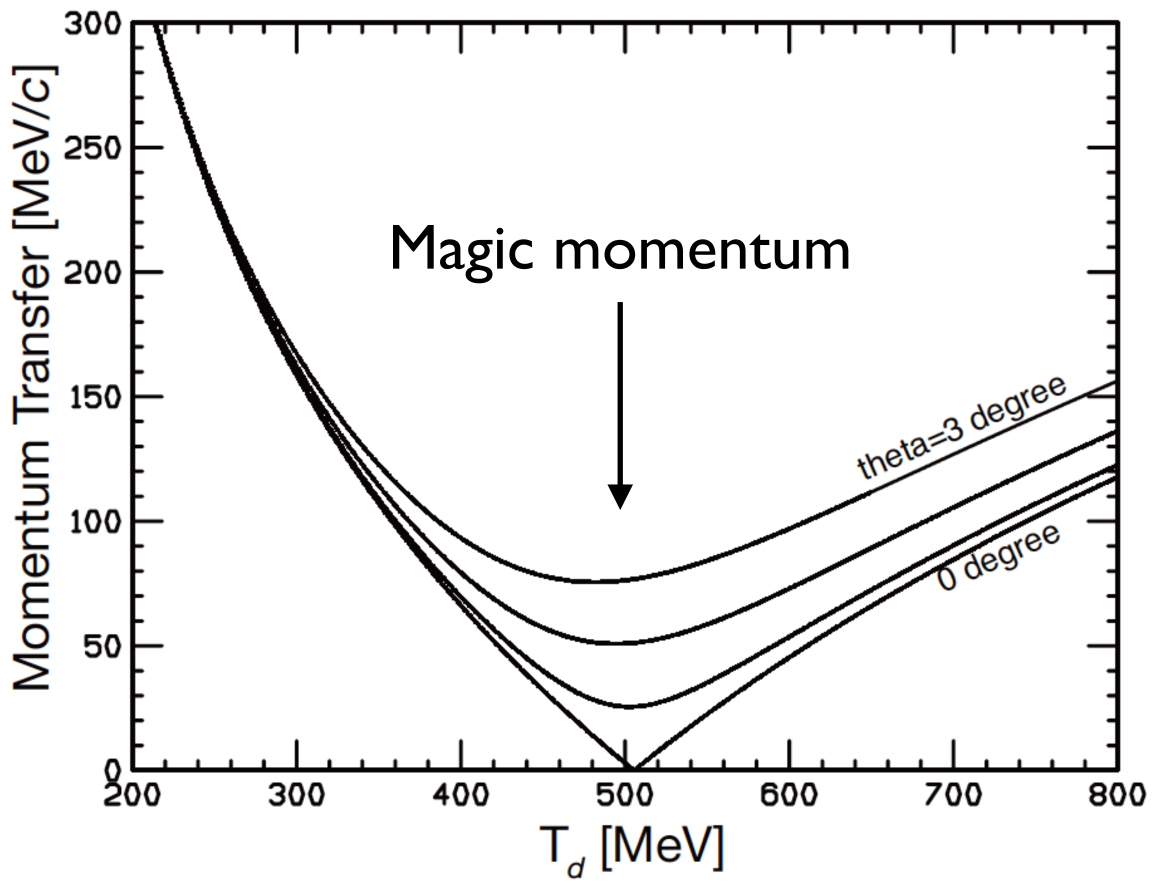
Spectroscopy of pionic atoms in (d,^{3}He) reactions

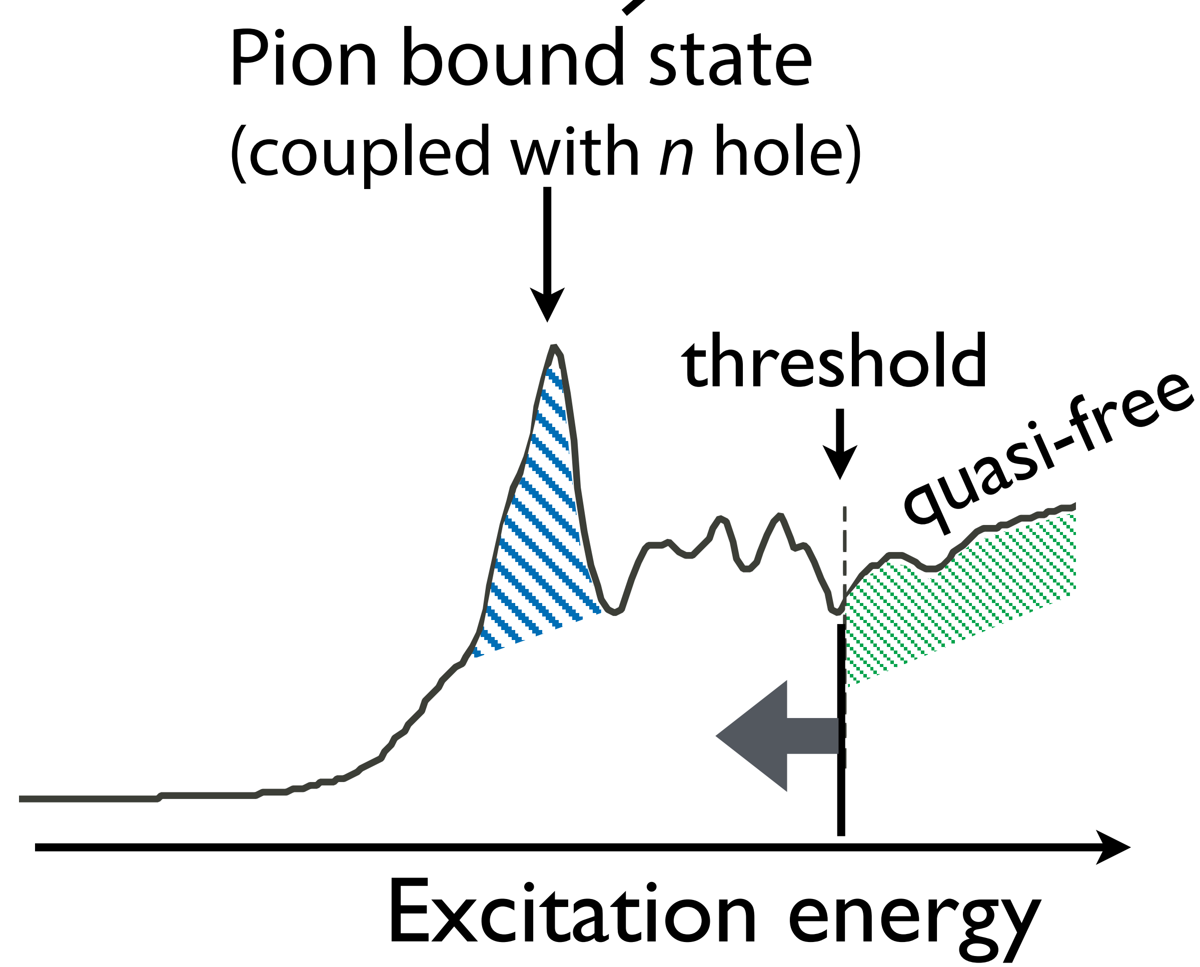
Missing mass spectroscopy to measure excitation spectrum of pionic atoms

Direct production of pionic atoms



Momentum transfer





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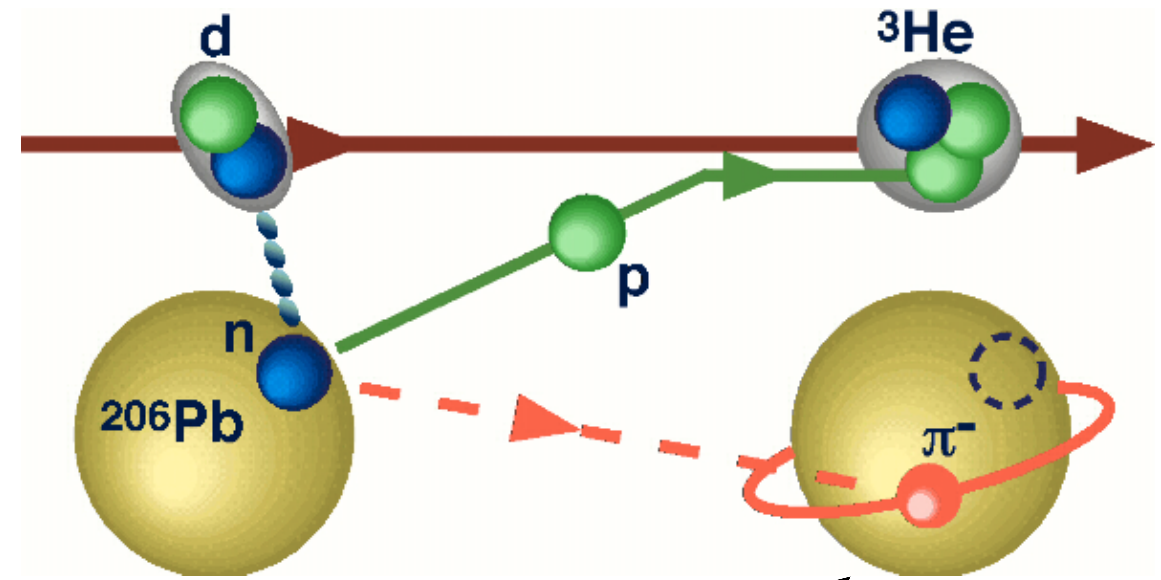
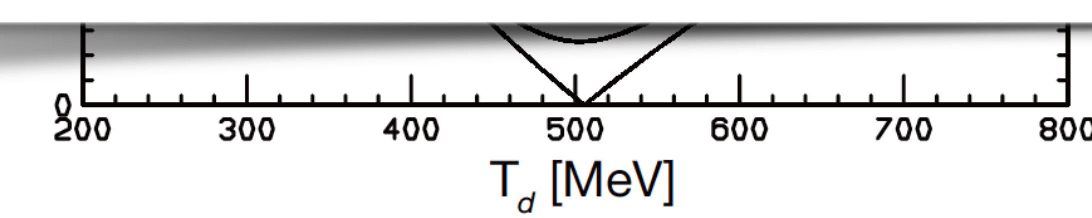


Momentum transfer

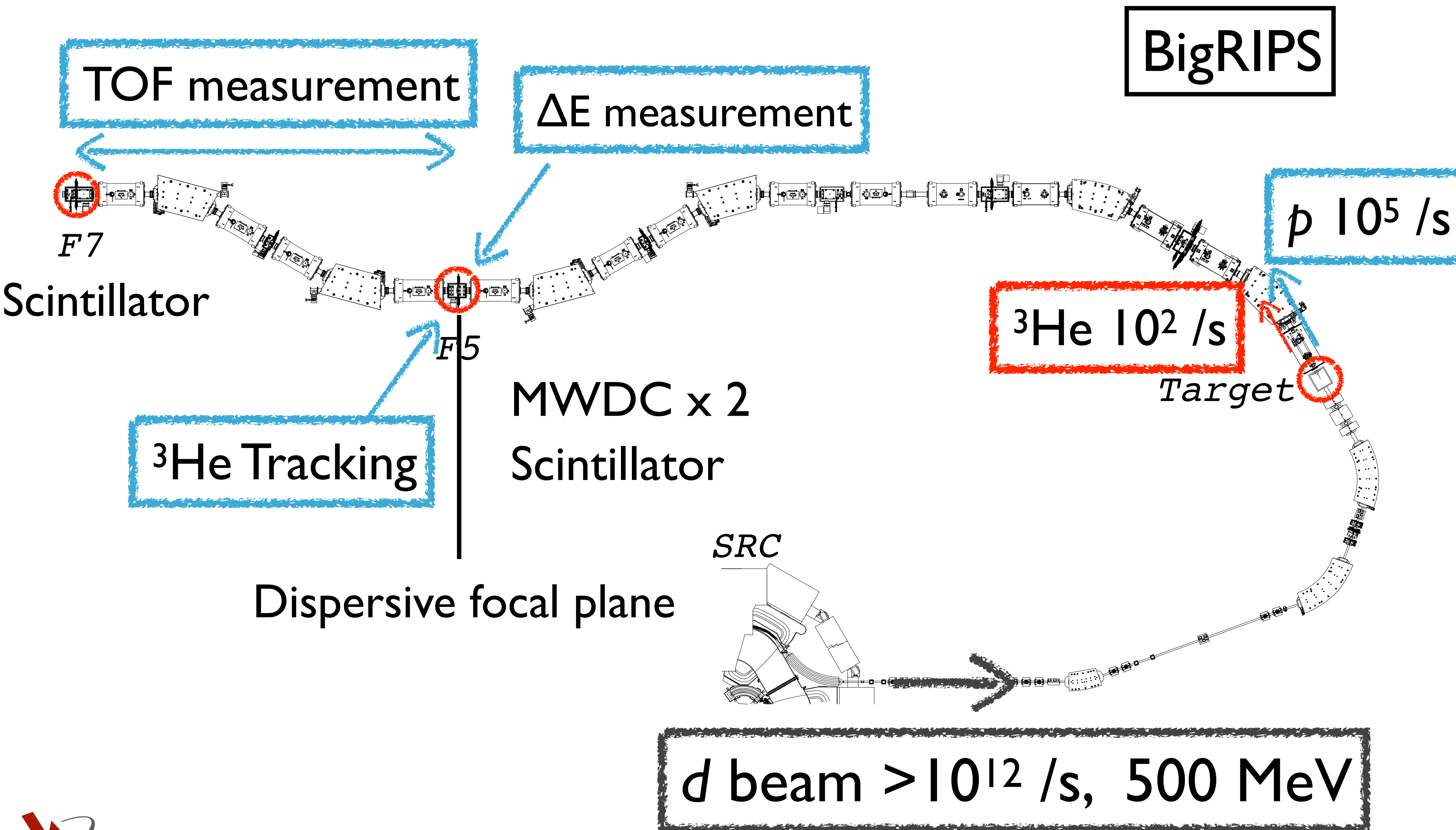


Keys for high quality information

- Quantum states of the pionic atoms are determined i.e. both pionic levels and nuclear excitation levels.
- Observable is Lorentz invariant.

Excitation energy

(d,^{3}He) Reaction Spectroscopy in RIBF



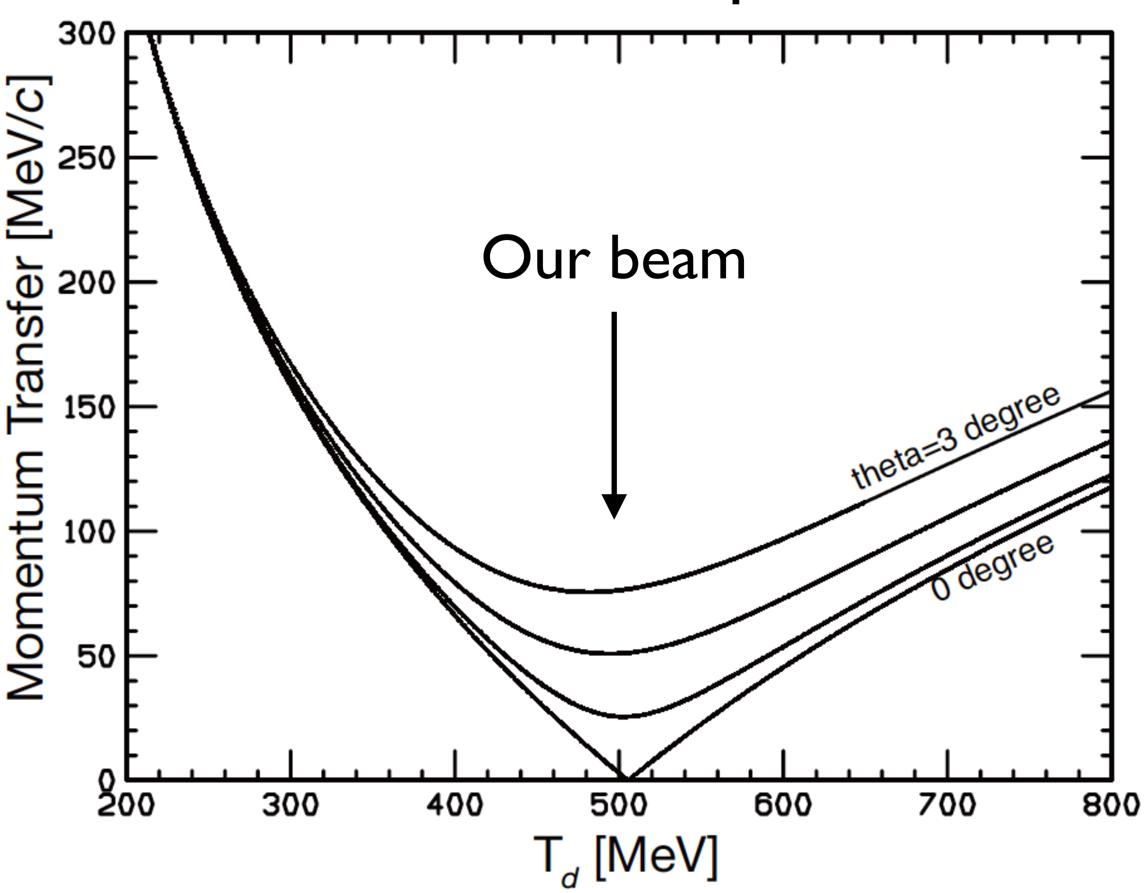
Pionic ^{121}Sn atom

Pilot run

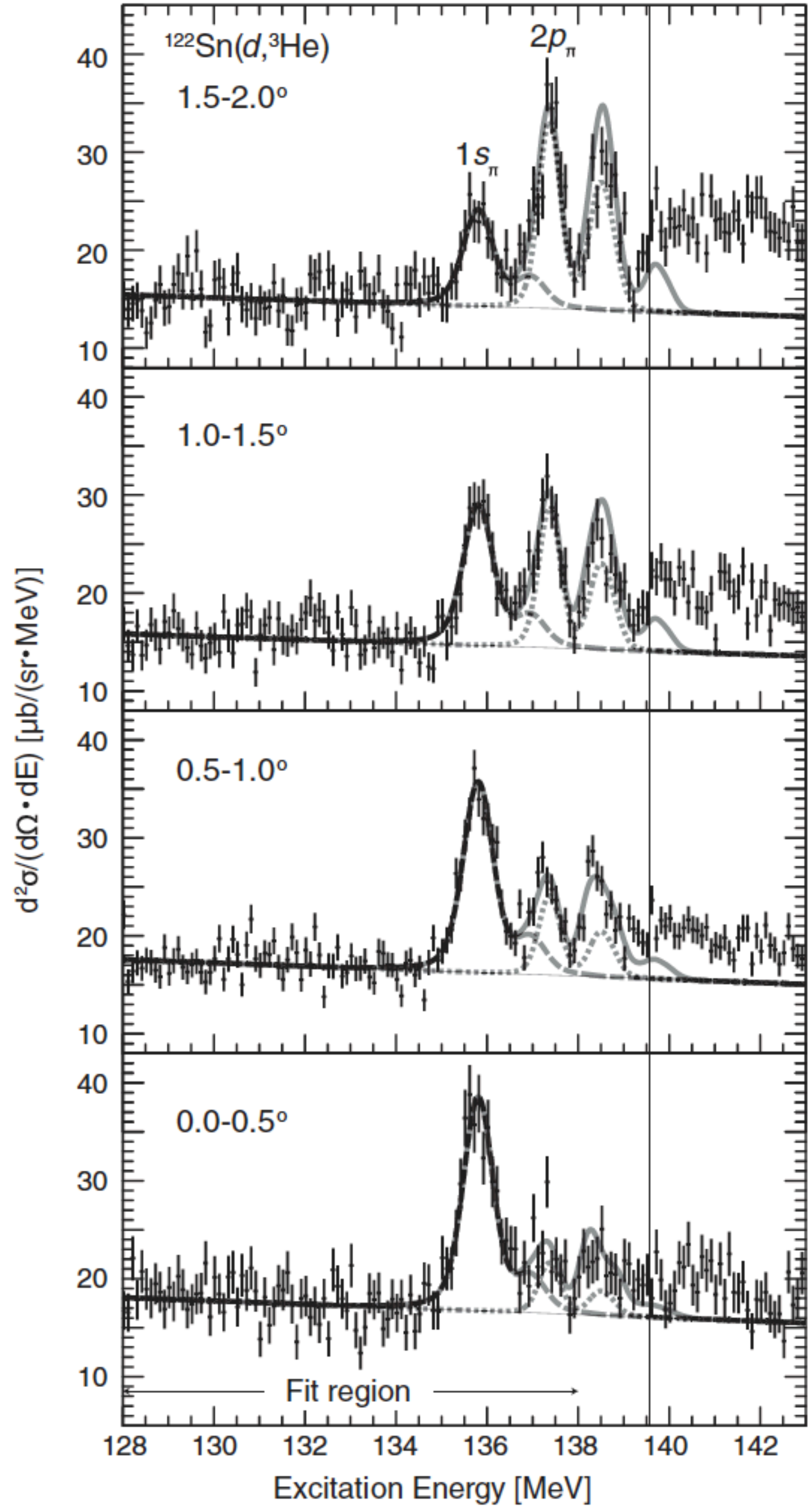
15 hours DAQ in 2010

First observation of θ dependence of π atom cross section

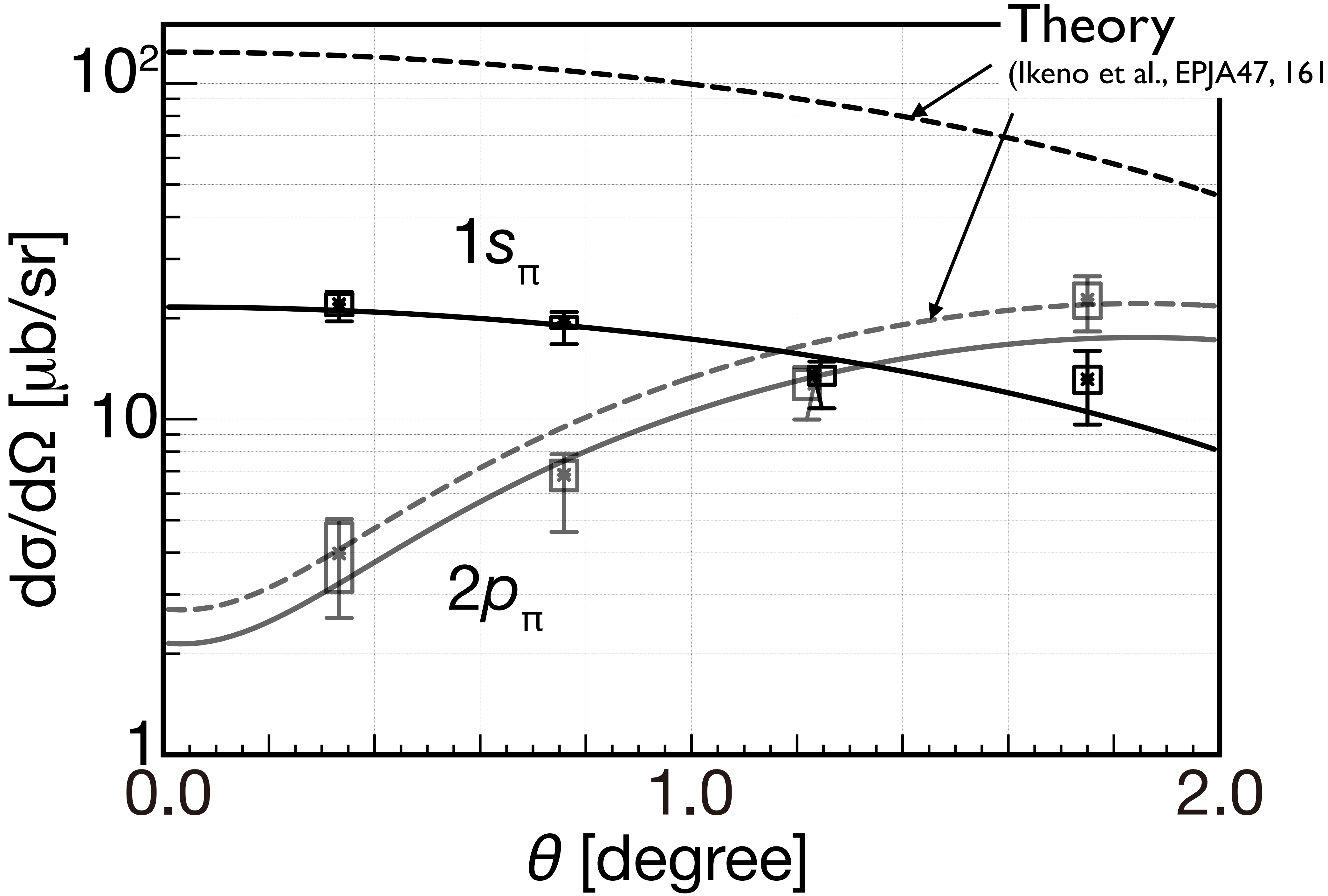
$\Delta L \sim R \times q$



T. Nishi KI et al., PRL120, 152505 (2018)



1s and 2p pionic atom cross sections in (d,^{3}He)



θ dependence is well reproduced.
Theory calculates 5x larger cross section for 1s

T. Nishi KI et al., PRL120, 152505 (2018)

Pionic ^{121}Sn atom

Pilot run

15 hours DAQ in 2010

First simultaneous 1s and 2p observation

$$B_{1s} = 3.828 \pm 0.013(\text{stat})^{+0.036}_{-0.033}(\text{syst})\ \text{MeV}$$
$$\Gamma_{1s} = 0.252 \pm 0.054(\text{stat})^{+0.053}_{-0.070}(\text{syst})\ \text{MeV}$$
$$B_{2p} = 2.238 \pm 0.015(\text{stat})^{+0.046}_{-0.043}(\text{syst})\ \text{MeV}$$

Resolution 394 keV (FWHM)

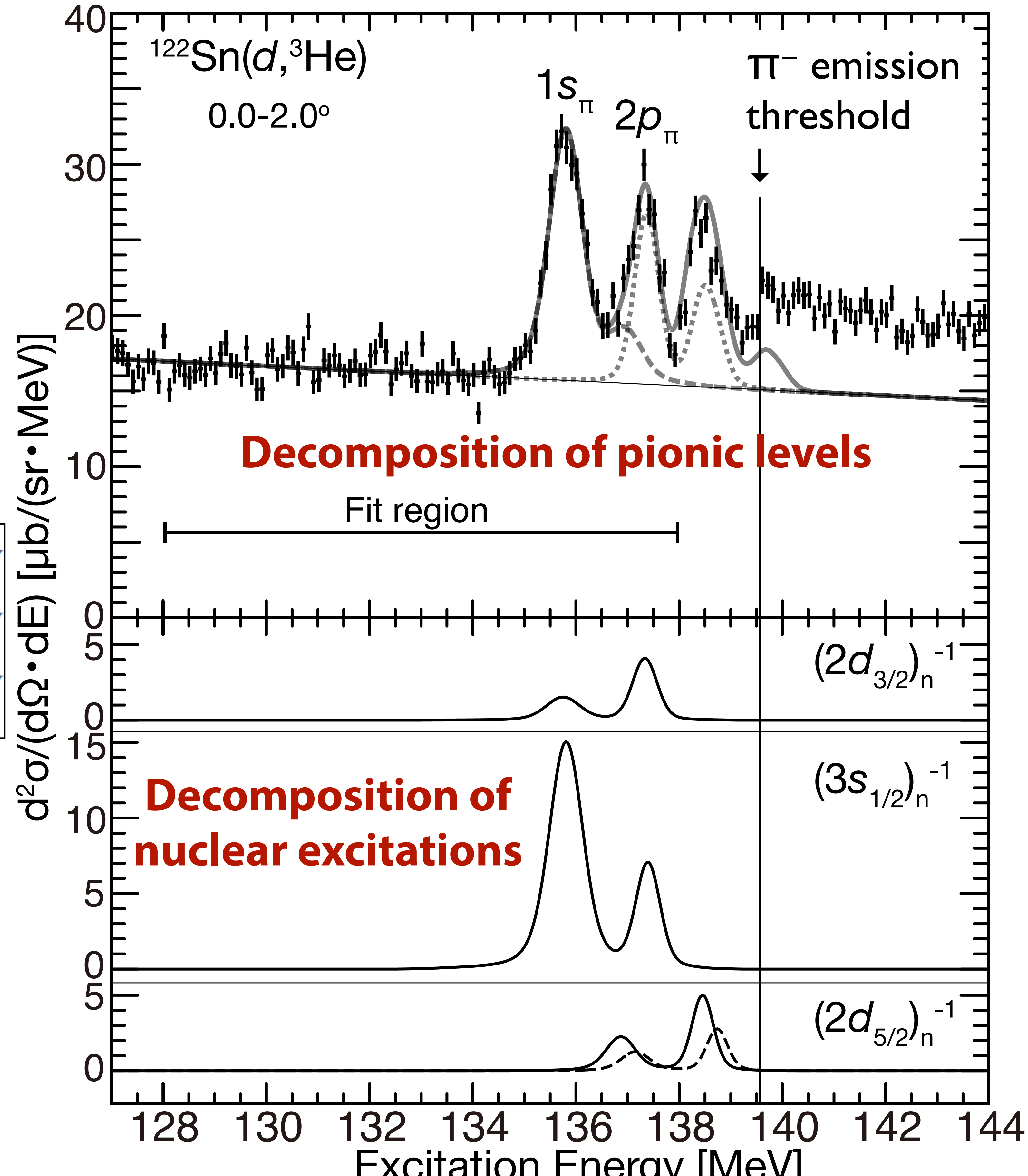
Theories

B_{1s}= 3.787–3.850 MeV

Γ_{1s}= 0.306–0.324 MeV

B_{2p}= 2.257–2.276 MeV

T. Nishi KI et al., PRL120, 152505 (2018)



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T. Nishi KI et al., PRL120, 152505 (2018)

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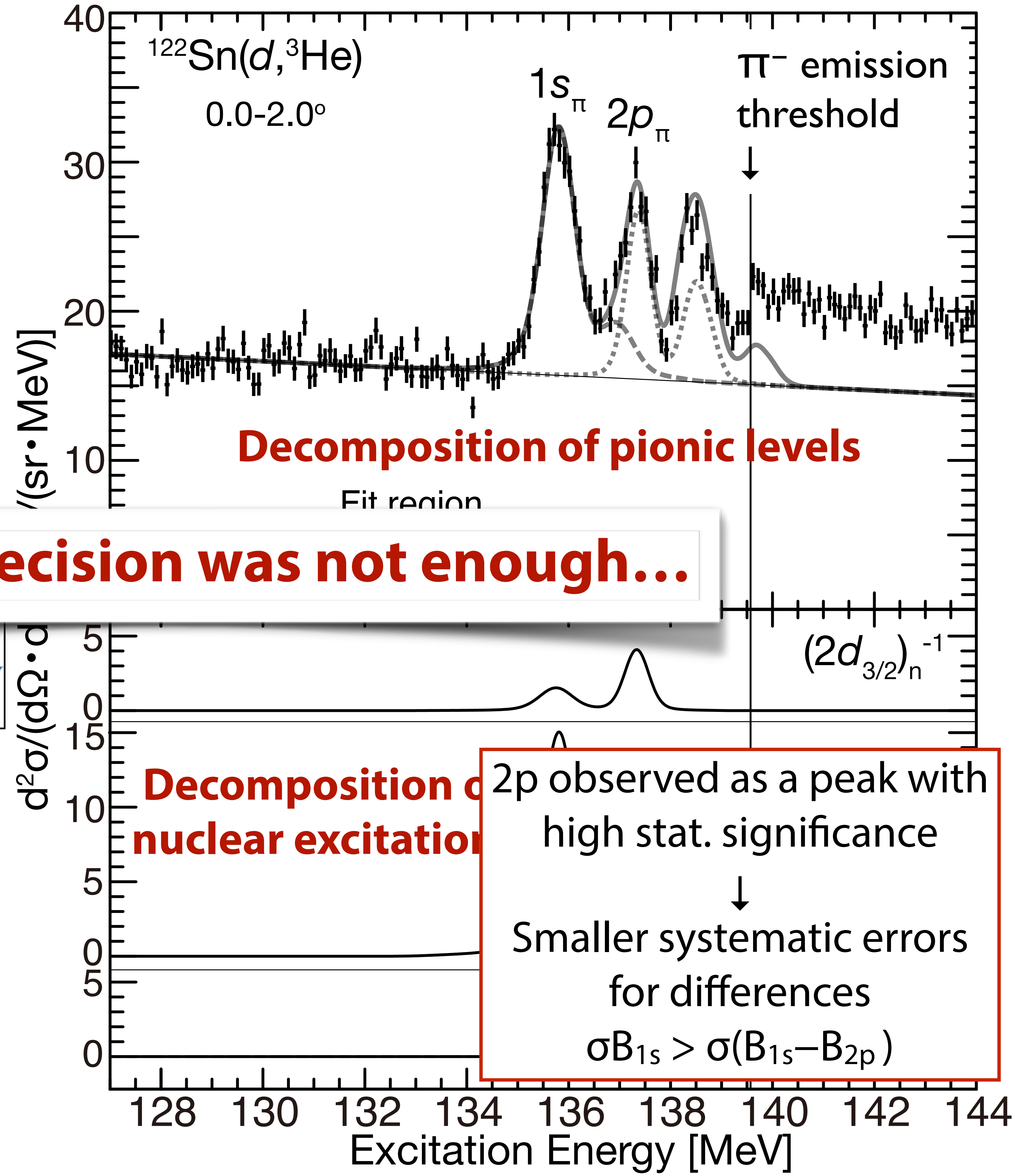
$B_{1s} = 3.828 \pm 0.0$ …

$\Gamma_{1s} = 0.252 \pm 0.0$ …

$B_{2p} = 2.238 \pm 0.015(\text{stat})^{+0.046}_{-0.043}(\text{syst})$ MeV

However, precision was not enough…



2p observed as a peak with high stat. significance

↓

Smaller systematic errors for differences

$\sigma B_{1s} > \sigma(B_{1s} - B_{2p})$

Resolution 394 keV (FWHM)

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B_{1s}= 3.787–3.850 MeV

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High Precision Spectrum of $^{122}Sn(d,{}^3He)$ in 2014 run

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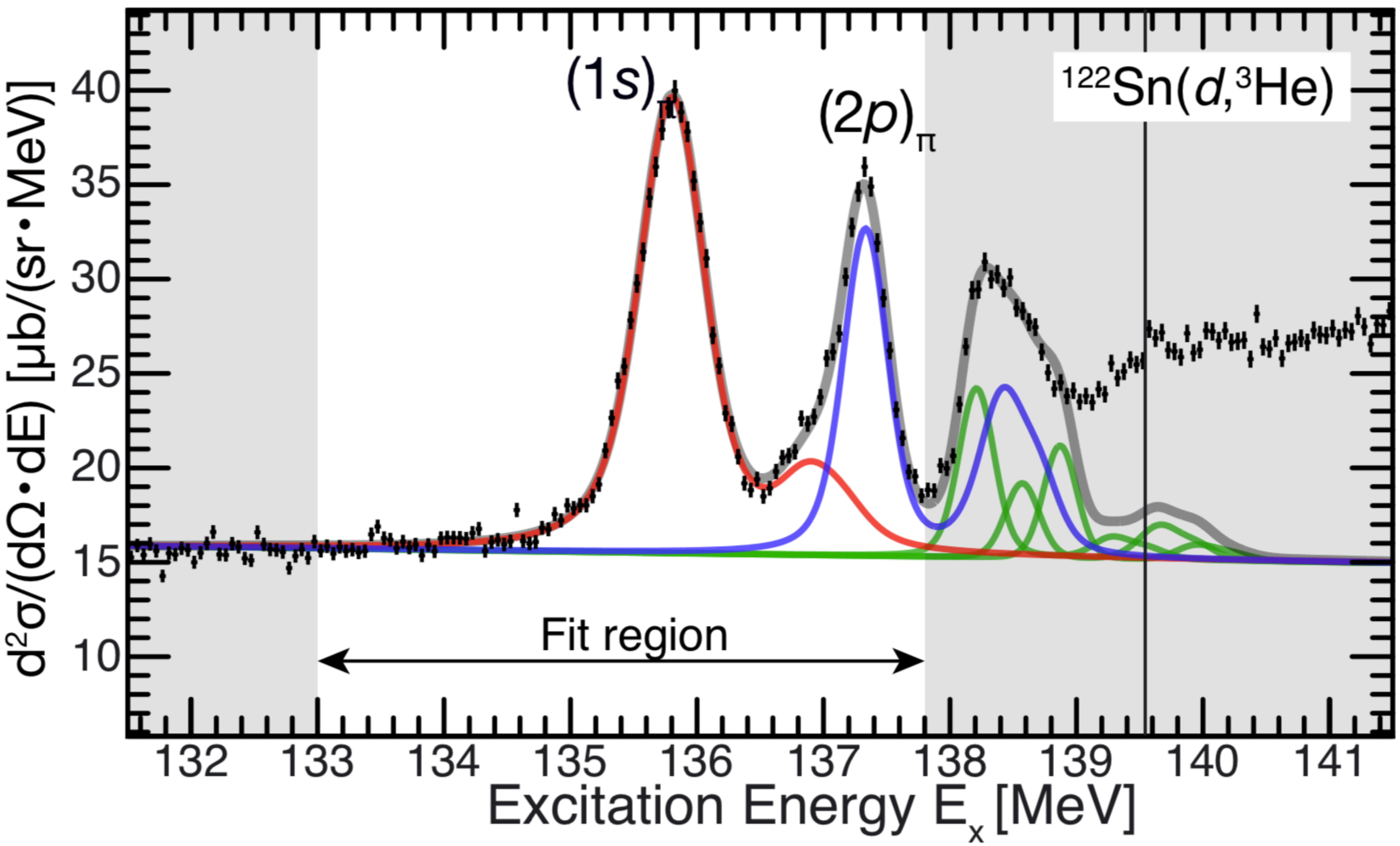


2p observed as a peak with high stat. significance

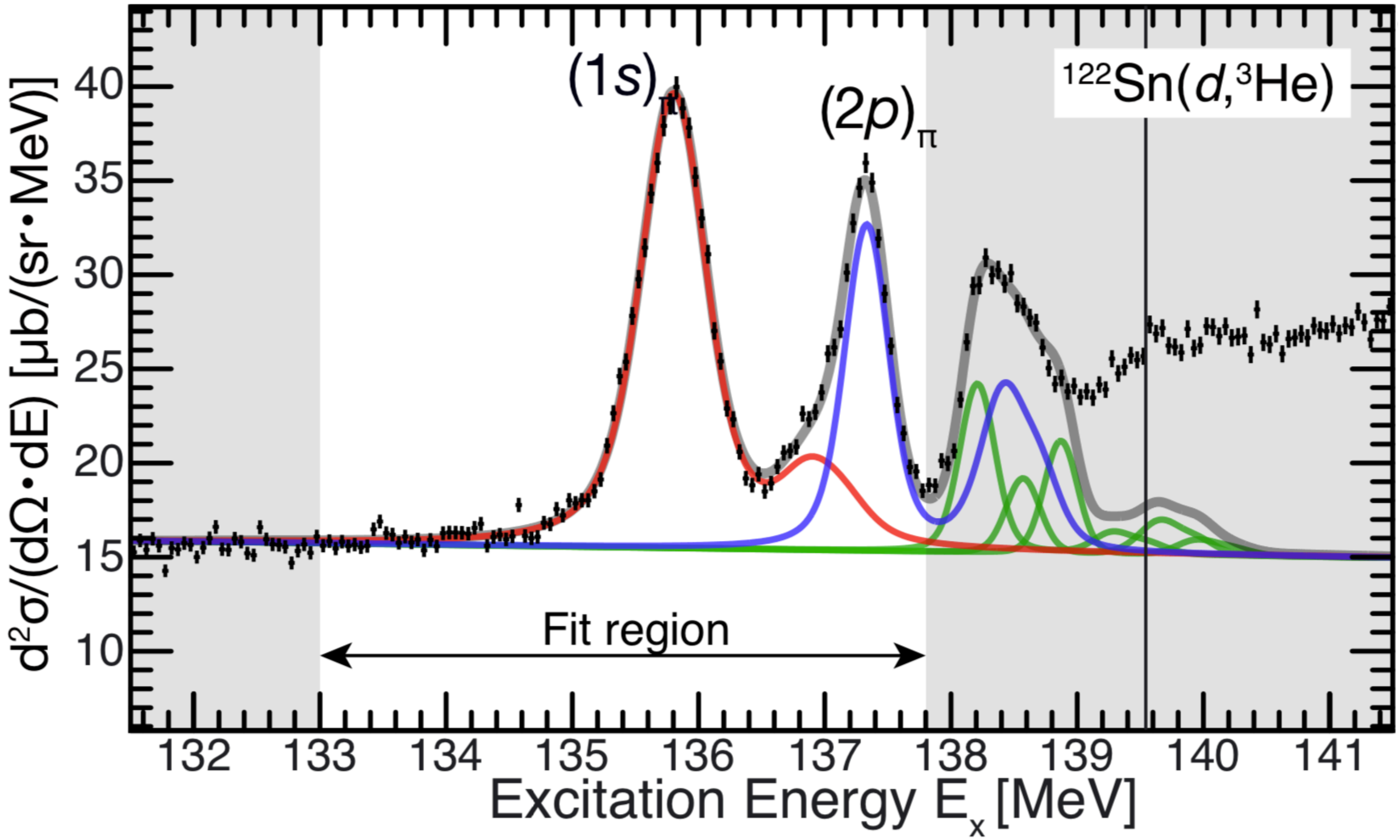
↓

Smaller systematic errors for differences

$\sigma B_{1s} > \sigma(B_{1s} - B_{2p})$

Under review
arXiv: 2204.05568

High Precision Spectrum of ^{122}Sn(d,^{3}He) in 2014 run



	[keV]	Statistical	Systematic
$B_\pi(1s)$	3831	±3	$+78-76$
$B_\pi(2p)$	2276	±3	$+84-83$
$B_\pi(1s)-B_\pi(2p)$	1555	±4	±12
$\Gamma_\pi(1s)$	316	±12	$+36-39$
$\Gamma_\pi(2p)$	164	±17	$+41-32$
$\Gamma_\pi(1s)-\Gamma_\pi(2p)$	152	±20	$+28-36$

Best resolution 287 keV (FWHM)

2p observed as a peak with high stat. significance
↓
Smaller systematic errors for differences
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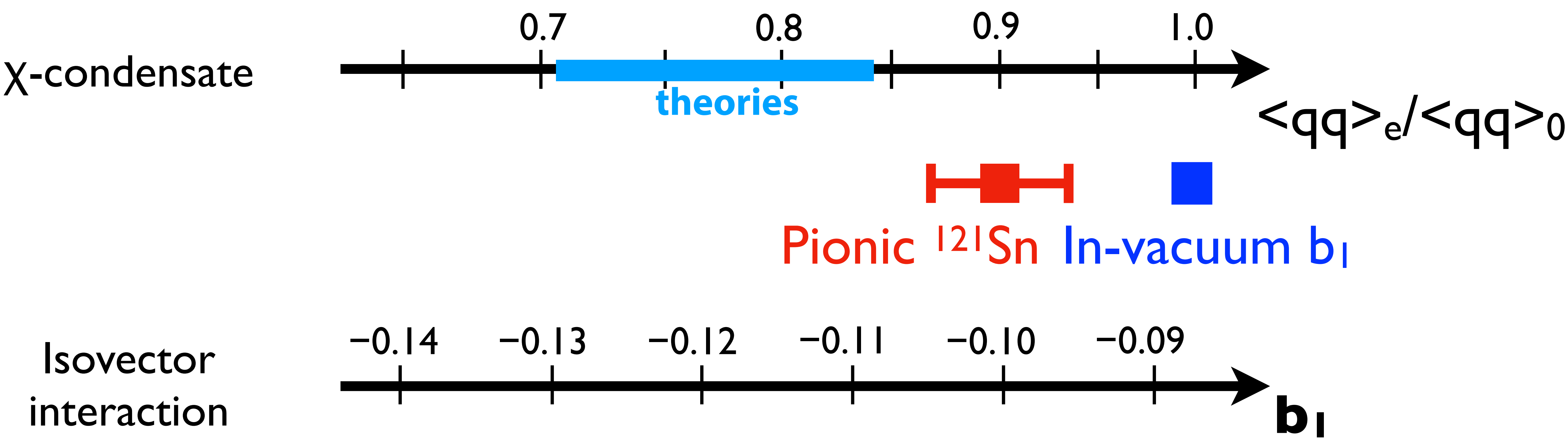
Under review
arXiv: 2204.05568

Deduced b_1 and chiral condensate at ρ_c



Before corrections: b_1 = -0.1005

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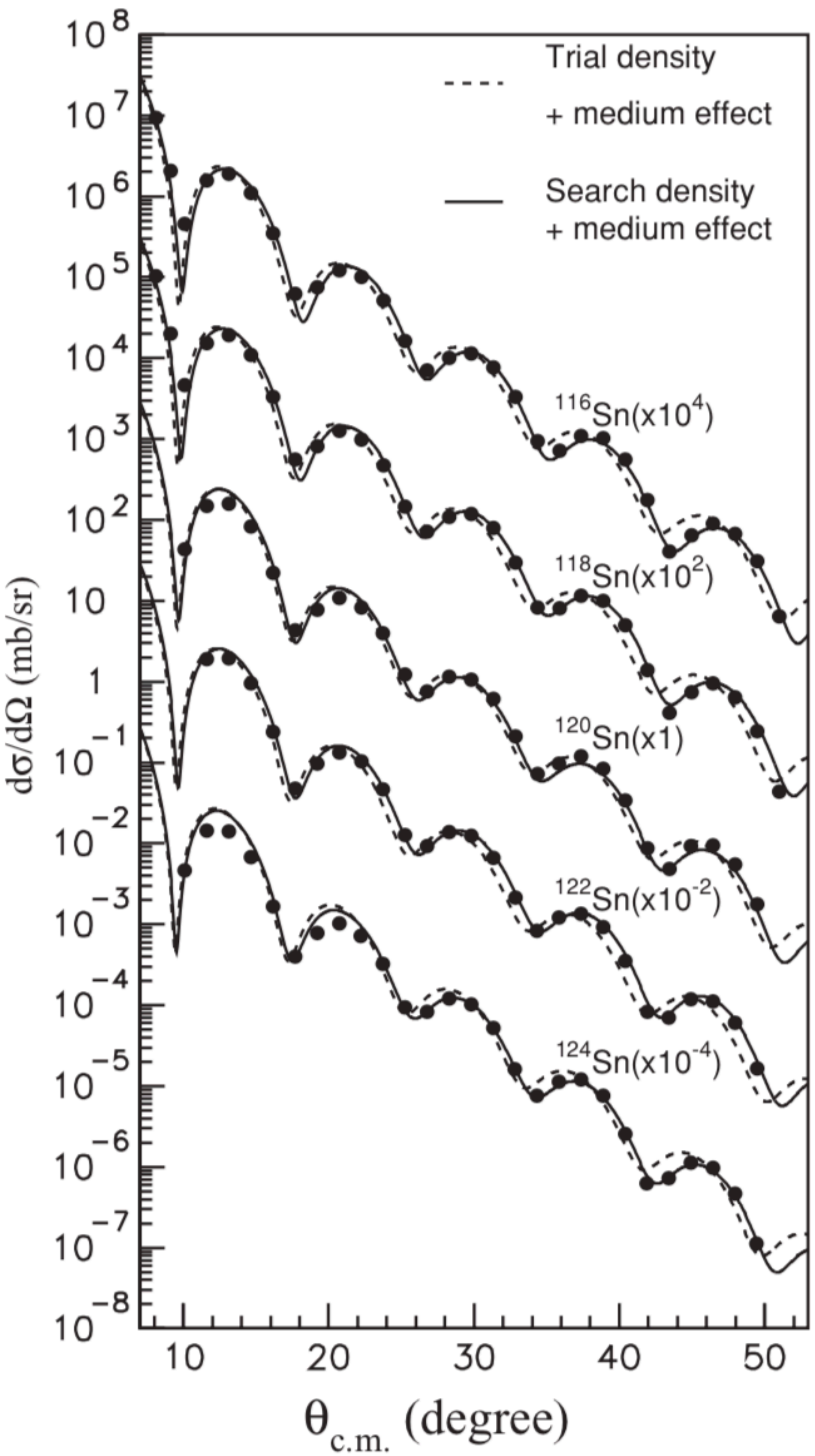
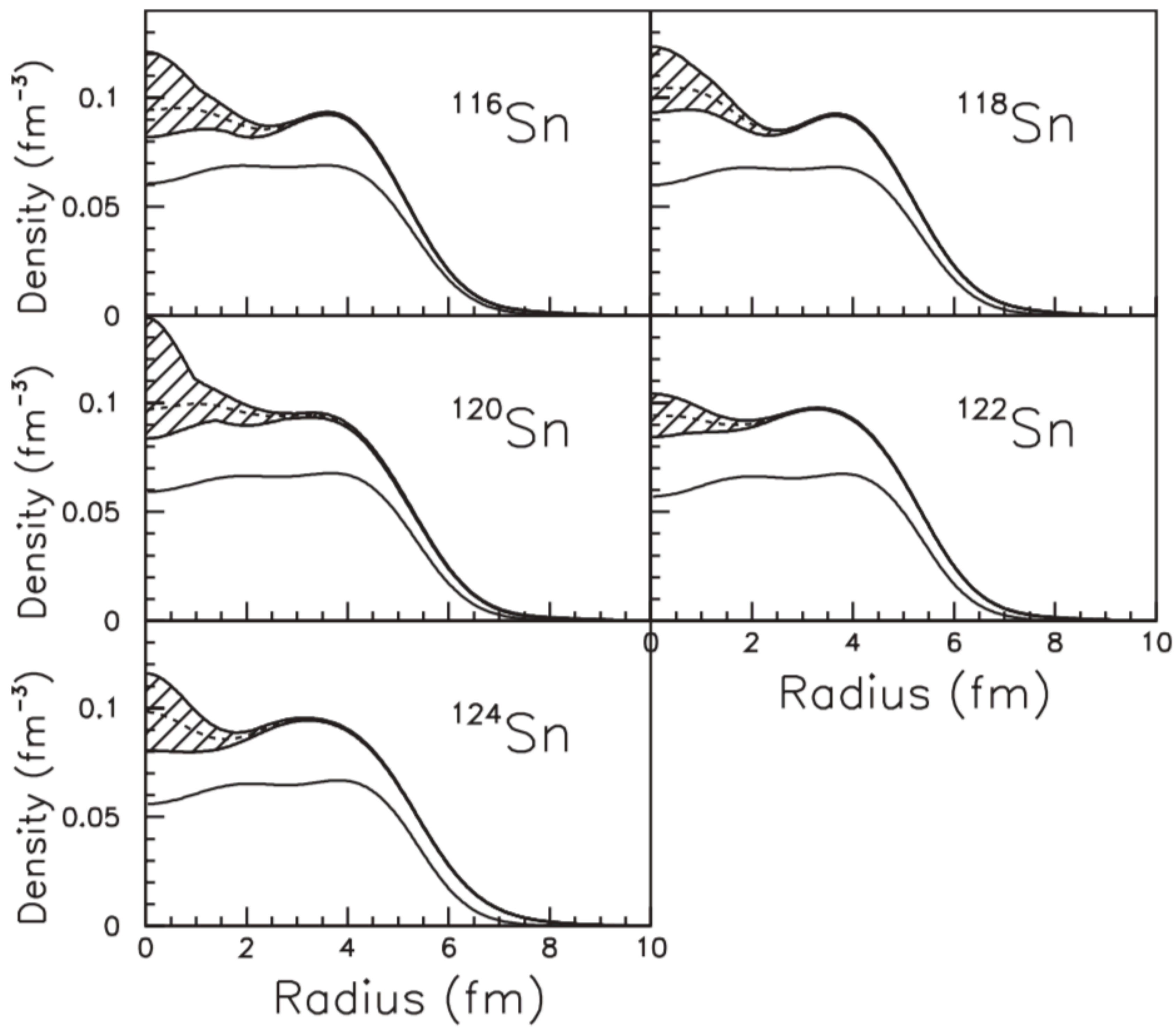
ξ	$\rho_n(r)$	Abs.	C.S.	Res.	$b_1\,[m_\pi^{-1}]$	$\mathrm{Im}B_0\,[m_\pi^{-4}]$
0	2pF	ρ^2	Neff	–	−0.1005	0.0469
1	2pF	ρ^2	Neff	–	−0.0997	0.0473
1	Osaka	ρ^2	Neff	–	−0.1148	0.0473
1	Osaka	$pp+2np$	Neff	–	−0.1162	0.0474
1	Osaka	$pp+2np$	Green	–	−0.1194	0.0474
1	Osaka	$pp+2np$	Green	✓	−0.1210	0.0474

ξ : short-range correction, LLE
ρ : neutron density distribution
Abs. : representation of absorption term
C.S. : cross section calculation method
Res.: residual interaction

Measured nuclear density distribution of Sn isotopes

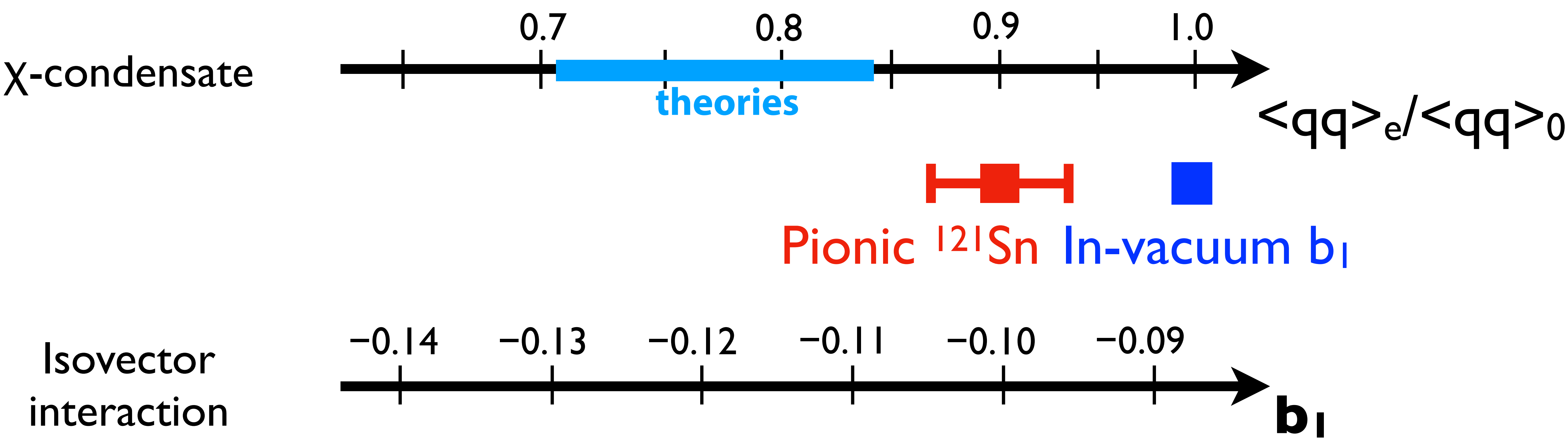
Sn(p,p') reaction at RCNP, Osaka





Terashima et al., PHYSICAL REVIEW C **77**, 024317 (2008)

Deduced b_1 and chiral condensate at ρ_e



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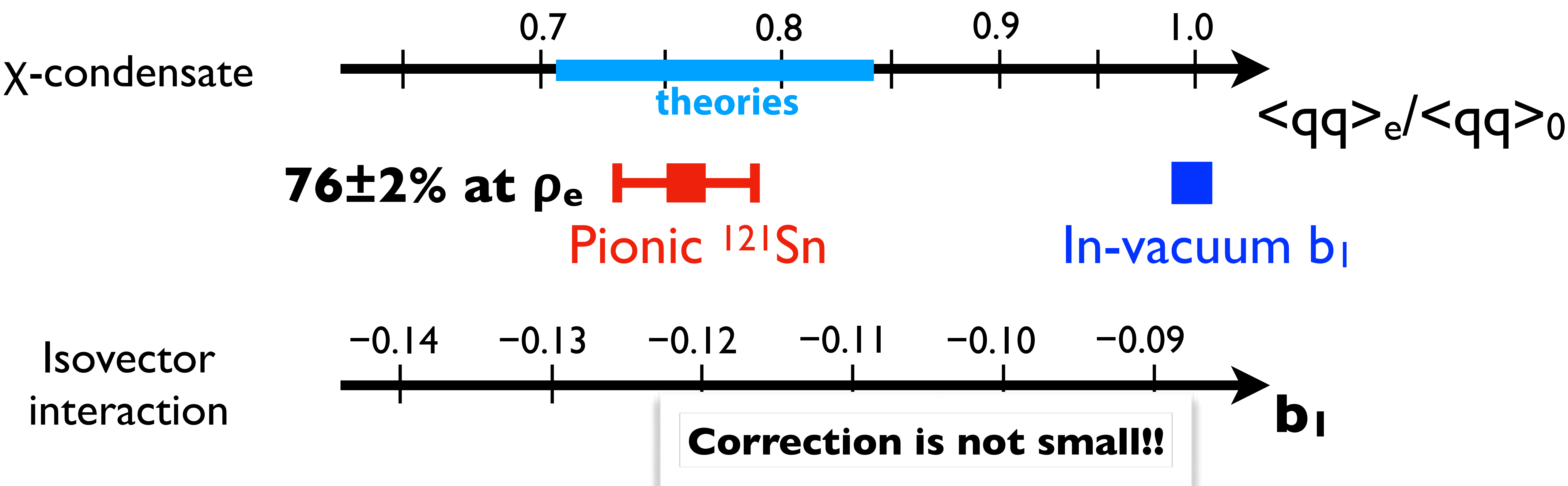

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After corrections: b_1 = -0.1210±0.0063

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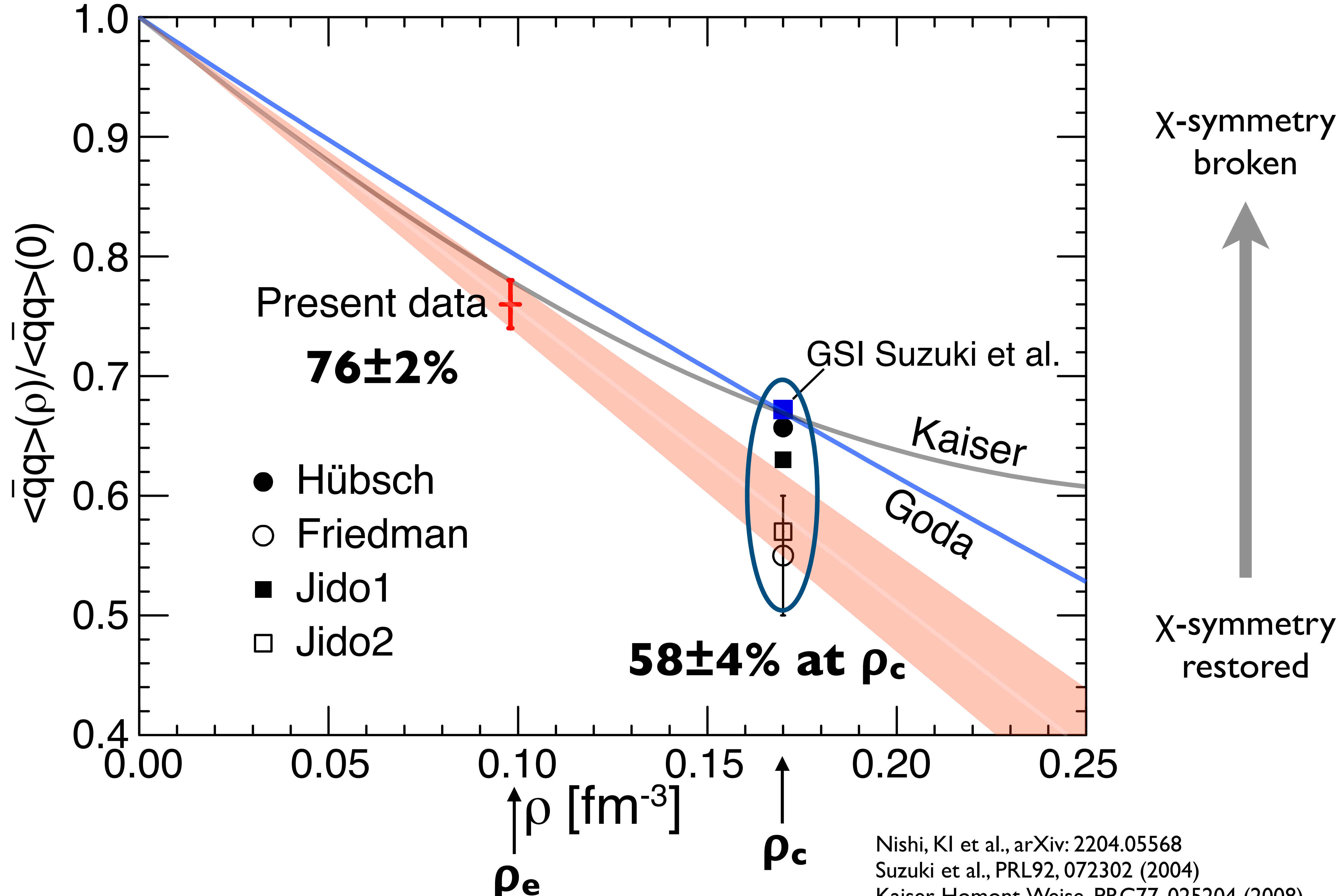
Res.: residual interaction

ρ dependence of chiral condensate



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Summary

- Chiral condensate at the normal nuclear density is evaluated to be reduced by 58±4%. We evaluated chiral condensate with errors at the well-defined density for the first time by pionic atom spectroscopy.
- The binding energies and widths of the 1*s* and 2*p* states in Sn121 were determined with unprecedented precision. Difference between the 1*s* and 2*p* values reduces the systematic errors drastically.
- Recent theoretical progress was adopted for the evaluation, which directly relates the chiral condensate and the pion-nucleus interaction.
- We calculated various corrections for the first time and applied them. The application made a large jump of the deduced chiral condensate. After the corrections, the chiral condensate ratio was deduced to be 58±4% with much higher reliability.
- We plan measurement of ρ dependence of chiral condensate in systematic study.