QCD Symmetries and Etaprime-Nucleus Interactions

Steven Bass

The eta' and eta mesons are connected to non-perturbative glue

About 300-400 MeV of their masses comes from gluonic (extra to Goldstone boson) dynamics beyond DChSB.

How is this glue manifest in eta-prime nucleus interactions with the eta-prime effective mass shift observed ~-40 MeV at nuclear matter density?

Medium modifications and proton spin physics: the GDH sum-rule in medium

EMMI in Kitzbühel, September 14, 2022

The η - and η' -nucleus interactions and the search for η , η' - mesic states

Steven D. Bass, Volker Metag* and Pawel Moskal

Experimental results

• CBELSA/TAPS measurement of n' potential in nuclei

$$V(\rho = \rho_0) = -(40 \pm 6(\text{stat}) \pm 15(\text{syst}))\text{MeV}$$

$$W(\rho = \rho_0) = -(13 \pm 3(\text{stat}) \pm 3(\text{syst}))\text{MeV}$$

- n' properties modified in medium
- Search for n' bound states in nuclei, WASA@GSI, BGOegg, ...



Protons and mesons from QCD

In QCD all hadron physics is emergent from more fundamental quarks and gluons

- Protons including their mass, spin ...
- Pions as messengers (exchange particles) of nuclear forces
- Pions are special because of chiral symmetry
- Confinement and DChSB in the infrared



 Gluonic degrees of freedom manifest in the n and n' mesons and their interactions (plus perhaps glueballs)

Chiral symmetry

• QCD Lagrangian with massless quarks exhibits chiral symmetry

$$\mathcal{L}_{QCD} = \bar{\psi}_L \Big(i\hat{\partial} + g\hat{A} \Big) \psi_L + \bar{\psi}_R \Big(i\hat{\partial} + g\hat{A} \Big) \psi_R - m_q \Big(\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L \Big) - \frac{1}{4} G_{\mu\nu} G^{\mu\nu}$$

$$\left(\begin{array}{c} u_L \\ d_L \end{array}\right) \ \mapsto \ e^{i\frac{1}{2}\vec{\alpha}.\vec{\tau}\gamma_5} \left(\begin{array}{c} u_L \\ d_L \end{array}\right) \quad , \quad \left(\begin{array}{c} u_R \\ d_R \end{array}\right) \ \mapsto \ e^{i\frac{1}{2}\vec{\beta}.\vec{\tau}\gamma_5} \left(\begin{array}{c} u_R \\ d_R \end{array}\right)$$

• Noether currents

$$J^{(3)}_{\mu 5} = \left[ar{u} \gamma_{\mu} \gamma_{5} u - ar{d} \gamma_{\mu} \gamma_{5} d
ight] \qquad \qquad \partial^{\mu} J^{(3)}_{\mu 5} = 2m_{u} ar{u} i \gamma_{5} u - 2m_{d} ar{d} i \gamma_{5} d ar{d} i \gamma_$$

• No parity doublets in hadron spectrum \rightarrow Spontaneous Chiral symmetry breaking: non zero condensate $\langle vac | \bar{q}q | vac \rangle < 0$ spontaneously breaks the symmetry

 \rightarrow Nonet of near massless Goldstone bosons with J^P = O^-

• Identify with pion, kaon, eta with meson mass squared proportional to m_q

$$m_{\eta_8}^2 = rac{4}{3}m_{
m K}^2 - rac{1}{3}m_{\pi}^2$$

... where is the singlet boson?

Eta and Etaprime masses

Mass matrix

٠

٠

$$M_{\eta-\eta'}^2 = \begin{pmatrix} \frac{4}{3}m_{\rm K}^2 - \frac{1}{3}m_{\pi}^2 & -\frac{2}{3}\sqrt{2}(m_{\rm K}^2 - m_{\pi}^2) \\ \\ -\frac{2}{3}\sqrt{2}(m_{\rm K}^2 - m_{\pi}^2) & [\frac{2}{3}m_{\rm K}^2 + \frac{1}{3}m_{\pi}^2 + \tilde{m}_{\eta_0}^2] \end{pmatrix}$$

$$egin{array}{rcl} ert \eta
angle &=& \cos heta \; ert \eta_8
angle - \sin heta \; ert \eta_0
angle \ ert \eta'
angle &=& \sin heta \; ert \eta_8
angle + \cos heta \; ert \eta_0
angle \end{array}$$

$$m_{\eta',\eta}^2 = (m_{\rm K}^2 + \tilde{m}_{\eta_0}^2/2) \pm \frac{1}{2} \sqrt{(2m_{\rm K}^2 - 2m_{\pi}^2 - \frac{1}{3}\tilde{m}_{\eta_0}^2)^2 + \frac{2}{3}}$$

• Eigenvalues

Diagonalize

- $$\begin{split} m_{\eta',\eta}^2 &= (m_{\rm K}^2 + \tilde{m}_{\eta_0}^2/2) \pm \frac{1}{2} \sqrt{(2m_{\rm K}^2 2m_{\pi}^2 \frac{1}{3}\tilde{m}_{\eta_0}^2)^2 + \frac{1}{9}\tilde{m}_{\eta_0}^4} \\ m_{\eta}^2 + m_{\eta'}^2 &= 2m_{K}^2 + \tilde{m}_{\eta_0}^2. \end{split}$$
- With no glue: chiral symmetry "predicts" eigenstates with masses 300 MeV "too small"

» "eta"
$$\left(\frac{1}{\sqrt{2}}|\bar{u}u+\bar{d}d\rangle\right)$$
 degenerate with the pion
» "etaprime" $|\bar{s}s\rangle$ with mass $\sqrt{2m_K^2 - m_\pi^2}$

Glue and the eta-prime

Etaprime has a strong affinity to glue

- SU(3) breaking means the gluon anomaly is important to both the eta and eta-prime
 - Extra gluonic mass term is associated with the QCD axial anomaly

$$J_{\mu5} = \left[ar{u}\gamma_{\mu}\gamma_{5}u + ar{d}\gamma_{\mu}\gamma_{5}d + ar{s}\gamma_{\mu}\gamma_{5}s
ight]$$

$$\partial^{\mu} J_{\mu 5} = \sum_{k=1}^{f} 2i \left[m_k \bar{q}_k \gamma_5 q_k \right] + N_f \left[\frac{\alpha_s}{4\pi} G_{\mu\nu} \tilde{G}^{\mu\nu} \right]$$

plus gluon topology (note the difference with "perturbative glue")

Chirality and anomalous glue

- Perturbative QCD conserves chirality for massless quarks
- Confinement and vacuum tunneling processes (instantons, ...) connect left and right handed quarks





 $\Gamma_\mu D^{\mu\nu}\Gamma_\nu \propto 1/k^4$

Axial anomaly primer

- Axial anomaly
 - Cannot regularize ultraviolet behaviour of the VVA (vector vector axial-vector) amplitude preserving gauge invariance and chiral symmetry at the same time
 - Must choose current conservation at the vector or at the axial-vector vertices but cannot keep both !
 - Gauge invariance wins !





Symmetries and anomalies

- Symmetries and UV regularization
- Need to define "infinite" momentum consistent with how nature works



Symmetries and anomalies

- Symmetries and UV regularization
- Need to define "infinite" momentum consistent with how nature works



• Famous examples: $\pi^0 \rightarrow 2\gamma$, n' mass in QCD, cosmological constant

Glue in etaprime physics

- Three important places it can contribute
 - » Gluonic potential associated with QCD vacuum gives the etaprime a big mass
 - > The etaprime has a large singlet component
 → coupling to gluonic intermediate states (OZI violation)
 - > Gluonic Fock components in the etaprime wavefunction
 → radiativel generated through perturbative QCD
 [Acta Phys Pol B Proc Suppl 2 (2009) 11]
 - » Mixing with higher mass pseudo-scalar glueball(s)

U(1) extended chiral Lagrangian

- Low energy effective Lagrangian
 - constructed to reproduce the axial anomaly in the anomalous divergence equation and the gluonic mass term for the singlet boson

$$\mathcal{L} = \frac{F_{\pi}^2}{4} \operatorname{Tr} \left(\partial^{\mu} U \partial_{\mu} U^{\dagger} \right) + \frac{F_{\pi}^2}{4} \operatorname{Tr} M \left(U + U^{\dagger} \right) + \frac{1}{2} i Q \operatorname{Tr} \left[\log U - \log U^{\dagger} \right] + \frac{3}{\tilde{m}_{\eta_0}^2 F_0^2} Q^2.$$

$$U = \exp\{i(\phi/F_{\pi} + \sqrt{2/3}\eta_0/F_0)\}\$$

• Q is the topological charge density and the gluonic potential yields the gluonic contribution to the etaprime mass term

$$\frac{1}{2}iQ \operatorname{Tr}[\log U - \log U^{\dagger}] + \frac{3}{\tilde{m}_{\eta_0}^2 F_0^2} Q^2 \mapsto -\frac{1}{2} \tilde{m}_{\eta_0}^2 \eta_0^2.$$

• Couple to sigma mean field and repeat ...

$$\mathcal{L}_{\sigma Q} = Q^2 g_{\sigma}^Q \sigma$$

$$\tilde{m}_{\eta_0}^2 \mapsto \tilde{m}_{\eta_0}^{*2} = \tilde{m}_{\eta_0}^2 \frac{1+2x}{(1+x)^2} < \tilde{m}_{\eta_0}^2$$

where

$$x = \frac{1}{3} g_{\sigma}^Q \sigma \tilde{m}_{\eta_0}^2 F_0^2.$$

Eta(prime) bound states in nuclei

[SDB + AW Thomas, Phys Lett B634 (2006) 368, Acta Phys Pol B 45 (2014) 627]

- New experiments + big effort ...
- Binding energies and effective masses in nuclei are sensitive to
 - Coupling to scalar sigma field in the nuclei in mean field approx.
 - Nucleon-nucleon and nucleon-hole excitations in the medium
- TH: Solve for the meson self-energy in the medium

$$k^2-m^2={\rm Re}~\Pi(E,\vec{k},\rho)$$

$$\Pi(E,\vec{k},\rho)\Big|_{\{\vec{k}=0\}} = -4\pi a\rho \left(1 + \frac{A}{A-1}\frac{\mu}{M}\right)$$

- Where a is the "eta(prime)-nucleon scattering length"

QCD Inspired Models

- Quark Meson Coupling Model:
 - Can vary the mixing angle !
 - Use large eta and eta' masses to treat the eta and eta' as MIT Bags embedded in the medium with coupling between the light-quarks and the sigma mean field in the nucleus

Solve for in-medium mass and binding energy

- \rightarrow Extract an "effective" scattering length for the model
- \rightarrow Increases with increasing singlet component in the eta !

	$m [{ m MeV}]$	$m^*~[{\rm MeV}]$	$\operatorname{Re} a$ [fm]
η_8	547.75	500.0	0.43
$\eta (-10^{\circ})$ $\eta (-20^{\circ})$	547.75	449.3	$\begin{array}{c} 0.04 \\ 0.85 \end{array}$
η_0	958	878.6	0.99
$\eta' (-10^{\circ}) = \eta' (-20^{\circ})$	958 958	899.2 921.3	$\begin{array}{c} 0.74 \\ 0.47 \end{array}$
· 、 /	I		

 $m_{\eta}^*/m_{\eta} \approx 1 - 0.17 \ \rho/\rho_0$ $m_{\eta'}^*/m_{\eta'} \approx 1 - 0.05 \ \rho/\rho_0.$

- Eta-prime mass shift induced by glue; otherwise strange state without large coupling to sigma mean-field in the nucleus
- Excellent agreement with EP results from CBELSA/TAPS

Medium modifications II

- Nucleon spin structure is modified in medium
- GDH physics with polarized photoproduction
 - Scattering on polarized proton in a polarized nucleus

$$\int_{M^2}^{\infty} \frac{\mathrm{d}s_{\gamma p}}{s_{\gamma p} - M^2} (\sigma_{\mathrm{P}} - \sigma_{\mathrm{A}}) = 2\pi^2 \alpha_{\mathrm{QED}} \kappa^2 / M^2$$

- Biggest contribution comes from the Δ excitation.
 - Effective hadron masses reduced in medium
 - Expected enhancement in nucleon (anomalous) magnetic moments
- How do resonance, especially the Δ excitation, contributions behave with increasing nuclear density?
- In principle, large effects (relative to the deep inelastic case)

GDH sum-rule from proton

- GDH integral with biggest contribution from Δ M1-magnetic transition
- Theory predicts + 205 μ b, nearly all in isovector channel with κ = 1.79
- MAMI+ELSA gave + 226 \pm 5 \pm 12 µb for photon energies up to 2.9 GeV
- High-energy part extratarcted from low-Q2 data from mainly Jlab and COMPASS [Bass, Skurzok, Moskal in PRC 2018]

$$\int_{s_0}^{\infty} \frac{ds}{s - M^2} (\sigma_P - \sigma_A) = -15 \pm 2 \ \mu \mathrm{b}.$$

(with interesting effective Regge intercept value) gives

$$\int_{M^2}^{\infty} \frac{ds}{s - M^2} (\sigma_P - \sigma_A) = 211 \pm 13 \ \mu \text{b}$$

Estimates – idealized case at ρ_0

- Proton mass ~ reduced about -30 MeV at ρ_0 (Oset et al)
- Take ~ 20% quenching in $g_A^{(3)}$ together with (QMC relation)

$$\mu_N^*/\mu_N \sim g_A^{(3)}/g_A^{*(3)}$$

- Gives expected enhancement in GDH integral by factor ~ 2 at ρ_0
 - order of magnitude bigger than 20% effect in polarized DIS

$$\int_{0}^{1} \mathrm{d}x g_{1}^{(p-n)}\left(x, Q^{2}\right) = \frac{g_{\mathrm{A}}^{(3)}}{6} C_{\mathrm{NS}}\left(Q^{2}\right)$$

Deuteron

$$\int_{\nu_{\rm th}}^{\infty} \frac{d\nu}{\nu} (\sigma_P - \sigma_A) = 4\pi^2 \alpha_{\rm QED} S \kappa^2 / M^2$$

- K_p = 1.79
- $K_n = -1.91$
- K_d = 0.86
- (Photo-disintregration contribution to GDH)
- Small binding, -2.2 MeV
- D-state probability (5 ± 1)%

³He, K = -8.37

⁷Li, K = +4.61



Conclusions + Outlook

- Eta-primes in medium
 - mass shift ~-40MeV at p_0 observed by CBELSA/TAPS in good agreement with QMC prediction
 - Catalysed by anomalous glue, otherwise "pure" strange state
- Experiments: can we bound states with EP reach to this precision?
- What about eta bound states?
- GDH physics with polarized photoproduction
 - Effective hadron masses reduced in medium
 - Expected enhancement in nucleon (anomalous) magnetic moments
 - In principle factor 2 bigger enhancement cf. Polarised DIS case.
 Targets: d, ³He, ⁷Li... Experiments please!