Hyperons and Strange Mesons in Nuclei and Neutron Stars



based on L. Tolos and L. Fabbietti, Prog. Part. Nucl. Phys. 112 (2020) 103770



EMMI Workshop "Meson and Hyperon Interactions with Nuclei"

14-16 September 2022 Kitzbühel, Austria







- Hyperons and where to find them
- YN and YY interactions
- Hypernuclei
- Hyperons in matter
- Hyperons and Neutron Stars
- Present and Future

Hyperons and where to find them

A hyperon is a baryon containing one or more strange quarks

Hyperon	Quarks	(J₽)	Mass (MeV)
Δ	uds	O(I/2+)	1115
Σ^+	uus	1(1/2+)	1189
Σο	uds	1(1/2+)	1193
Σ-	dds	1(1/2+)	1197
Ξo	USS	1/2(1/2+)	1315
Ξ-	dss	1/2(1/2+)	1321
Ω-	<mark>8</mark> 55	0(3/2+)	1672

credit: Vidana

The study of hypernucleus allows for

new spectroscopy
information on strong and weak interactions between hyperons and nucleons



In Neutron Stars





YN and YY interactions

- Study strangeness in nuclear physics
- Provide input for hypernuclear physics and astrophysics



hypernucleus

Scarce YN scattering data due to the short life of hyperons and the low-density beam fluxes

 ΛN and ΣN : < 50 data points ΞN very few events

NN: > 5000 data for E_{lab}<350 MeV

Data from hypernuclei:

- more than 40 Λ-hypernuclei
 (ΛN attractive)
- few $\Lambda \Lambda$ hypernuclei
- $(\Lambda\Lambda$ weak attraction)
- few Ξ-hypernuclei(ΞN attractive)
- evidence of 1 Σ -hypernuclei ? (Σ N repulsive)

Data on femtoscopy!

Theoretical approaches to YN and YY

• Meson exchange models (Juelich/Nijmegen models)

To build YN and YY from a NN meson-exchange model imposing SU(3)_{flavor} symmetry Juelich: Holzenkamp, Holinde, Speth '89; Haidenbauer and Meißner '05 Nijmegen: Maesen, Rijken, de Swart '89; Rijken, Nagels and Yamamoto '10

 Chiral effective field theory approach (Juelich-Bonn-Munich group)
 To build YN and YY from a chiral effective Lagrangian similarly to NN
 interaction
 Juelich-Bonn-Munich: Polinder, Haidenbauer and Meißner '06; Haidenbauer, Petschauer, Kaiser, Meißner, Nogga and Weise '13

Kohno '10; Kohno '18

• Quark model potentials

To build YN and YY within constituent quark models

Fujiwara, Suzuki, Nakamoto '07 Garcilazo, Fernandez-Carames and Valcarce '07 '10

V_{low k} approach
 Garcilazo, Fernandez-Carames and Valcarce '07'10
 To calculate a "universal" effective low-momentum potential for YN and YY
 using RG techniques
 Schaefer, Wagner, Wambach, Kuo and Brown '06

• Lattice calculations (HALQCD/NPLQCD)

To solve YN and YY interactions on the lattice

HALQCD: Ishii, Aoki, Hatsuda '07; Aoki, Hatsuda and Ishii '10; Aoki et al '12 **NPLQCD:** Beane, Orginos and Savage '11; Beane et al '12

ΛN and ΣN scattering



 $T = V + V \frac{1}{E_0 - H_0 + i\eta} T$

LO: H. Polinder, J.H., U. Meißner, NPA 779 (2006) 244 NLO: J.H., N. Kaiser, et al., NPA 915 (2013) 24 Jülich '04: J.H., U.-G. Meißner, PRC 72 (2005) 044005



Hypernuclei



Double Λ hypernuclei



credit: Sanchez-Lorente

Also Ξ hypernuclei @ BNL, KEK ${}^{12}C(K^-, K^+){}^{12}_{\Xi^-}Be$ $K^- + p \rightarrow \Xi^- + K^+$

Laboratories: BNL, CERN, KEK, JLab, DAφNE, GSI, FAIR

Reactions:



Binding energy of Λ hypernuclei

Hypertriton lifetime puzzle



Binding energy of different hypernuclei as function of the mass number

Binding energy saturates at about -30 MeV for large nuclei

Single-particle model reproduces the data quite well Gal et al 2016



Expected $\tau({}^{3}_{\Lambda}H) = \tau(\Lambda)$ \Leftrightarrow observed: $\tau({}^{3}_{\Lambda}H) < \tau(\Lambda)$

Conflicting measurements by STAR and ALICE of the hypertriton lifetime triggered the revived experimental and theoretical interest

Hyperons in matter

A and Σ in dense matter

 $\mathbf{G} + \mathbf{G} + \mathbf{G} + \mathbf{G} = V + V \frac{Q_{\text{pauli}}}{E_0 - H_0} G$

 $k_F = 1.35 \text{ fm}^{-1} \ (\rho_0 = 0.166 \text{ fm}^{-3})$

	EFT LO	EFT NLO
۸ [MeV]	550 · · · 700	500 · · · 650
<i>U</i> ^(0)	-38.0 • • • -34.4	-28.2 · · · -22.4
<i>U</i> _Σ (0)	28.0 • • • 11.1	17.3 • • • 11.9

- Empirical value of Λ binding in nuclear matter ~27-30 MeV

- ΣN (I=3/2): discussion about repulsion or attraction, where ${}^{3}S_{1}$ - ${}^{3}D_{1}$ component is decisive. A repulsive ${}^{3}S_{1}$ - ${}^{3}D_{1}$ interaction is chosen in accordance to data on Σ^{-} atoms ${}_{60,6}^{-}$ and (π^{-} ,K⁺) inclusive spectra for Σ^{-} formation in heavy nuclei as well as lattice* indications

50 -10 EFT NLO 40 30 U_{Σ} (p_{Σ}=0) (MeV) 20 10 -50 -186 0.8 1.2 08 12 1.0 14 16 1.0 14 16 k_F (fm⁻¹) $k_{F} (fm^{-1})$

Haidenbauer and Meißner'15

* Nemura et al'18

Λ in dense matter: including three-body forces

Three-body forces are required to reproduce few-nucleon binding energies, scattering observables and nuclear saturation in non-relativistic many-body approaches



Λ in dense matter in χ EFT: Hyperon puzzle?



n,p,e⁻, μ^- , Λ in β -equilibrium χ EFT (NN, NNN, NNA) + meson-exchange (NY) Λ concentration is small

but still present in 2M_o NS

Only symmetric and neutron matter

 χ EFT NN, NNN,NY, NNY

 Λ in NS energetically unfavorable, but only neutrons and Λ are considered

Hyperons and Neutron Stars







- produced in core collapse
 supernova explosions, usually
 observed as pulsars
- usually refer to compact objects with M~1-2 M_{\odot} and R~10-12 Km
- extreme densities up to 5-10 ρ_0 (n₀=0.16 fm⁻³ => ρ_0 =3•10¹⁴ g/cm³)
- magnetic field : $B \sim 10^{8..16} G$
- temperature: T ~ 10 6...11 K
- observations: masses, radii, gravitational waves, cooling...



The Nucleonic Equation of State

The Equation of State (EoS) is a relation between thermodynamic variables describing the state of matter

Microscopic Ab-initio Approaches:

based on solving the many-body problem starting from two- and threebody interactions

- Variational method: APR, CBF,..
- Quantum Montecarlo : AFDMC..
- Coupled cluster expansion
- Diagrammatic: BBG (BHF), SCGF..
- Relativistic DBHF
- RG methods: SRG from *x*EFT..
- Lattice methods

Advantage: systematic addition of higher-order contributions Disadvantage: applicable up to? (SRG from $\chi EFT \sim 1-2 n_0$) Phenomenological Approaches: based on density-dependent interactions adjusted to nuclear observables and neutron star observations

- Non-relativistic EDF: Skyrme..
- Relativistic Mean-Field (RMF) and Relativistic Hartree-Fock (RHF)
- Liquid Drop Model: BPS, BBP,..
- Thomas-Fermi model: Shen
- Statistical Model: HWN,RG,HS..

Advantage: applicable to high densities beyond n₀ Disadvantage: not systematic

What about Hyperons?

First proposed in 1960 by Ambartsumyan & Saakyan

Hyperon	Mass (MeV/c ²)
Λ	1115.57 ± 0.06
Σ^+	1189.37 ± 0.06
Σ^0	1192.55 ± 0.10
Σ^{-}	1197.50 ± 0.05
Ξ^0	1314.80 ± 0.8
Ξ^-	1321.34 ± 0.14
Ω^{-}	1672.43 ± 0.14

 $p \ e^- \rightarrow n \ \nu_e$

Traditionally neutron stars were modeled by a uniform fluid of neutron rich matter in β -equilibrium $n \rightarrow p \ e^- \ \overline{\nu}_e$

but more exotic degrees of freedom are expected, such as **hyperons**, due to:

- high value of density at the center and
- the rapid increase of the nucleon chemical potential with density

Hyperons might be present at $n \sim (2-3)n_0$!!!

β-stable hyperonic matter

 μ_N is large enough to make N->Y favorable

$$n + n \rightarrow n + \Lambda$$

$$p + e^{-} \rightarrow \Lambda + v_{e^{-}}$$

$$n + n \rightarrow p + \Sigma^{-}$$

$$n + e^{-} \rightarrow \Sigma^{-} + v_{e^{-}}$$

$$\mu_i = b_i \mu_n - q_i \mu_e$$
$$\sum_i x_i q_i = 0$$





Vidana '18

The Hyperon Puzzle

The Hyperon Puzzle



Scarce experimental information:

- data from several single Λ - and few Ξ - hypernuclei, and few double Λ hypernuclei

few YN scattering data
 (~ 50 points) due to
 difficulties in preparing
 hyperon beams and no
 hyperon targets available

- YN data from femtoscopy

The presence of hyperons in neutron stars is energetically probable as density increases. However, it induces a strong softening of the EoS that leads to maximum neutron star masses < 2M_☉

Solution?

- stiffer YN and YY interactions
- hyperonic 3-body forces
- ➢ push of Y onset by ∆-isobars or meson condensates
- > quark matter below Y onset
- dark matter, modified gravity theories...

Space missions to study the interior of NS



Constraints from pulse profile modelling of rotation-powered pulsars with eXTP

and multimessenger astronomy!



Present and Future



A lot of experimental, observational and theoretical effort has been invested to study hyperons in nuclei and neutron stars

Hyperon-nucleon and hyperon-hyperon interactions are crucial for hypernuclear physics and the physics of compact objects, such as neutron stars

The presence of hyperons in neutron stars is energetically probable as density increases. However, it induces a strong softening of the EoS that leads to maximum neutron star masses $< 2M_{\odot}$ This is known as The Hyperon Puzzle

The future of hyperon physics relies on particle and nuclear experiments as well as X-ray and multimessenger astronomy









- KN interaction: $\Lambda(1405)$ resonance
- KNN bound state
- Kaons and Antikaons in matter
- Experiments and observations: from atoms to stars
- Present and Future

$\overline{K}N$ interaction: the $\Lambda(1405)$

• $\overline{K}N$ scattering in the I=0 channel is governed by the presence of the $\Lambda(1405)$ resonance, located only 27 MeV below the $\overline{K}N$ threshold

 $\pi\Sigma$

- 50's: idea originally proposed by Dalitz and Tuan
- since 90's: the study of KN scattering has been revisited by means of unitarized theories using meson-exchange models or chiral Lagrangians

 $\overline{K}N$

meson-exchange models

Mueller-Groeling, Holinde and Speth '90; Buettgen, Holinde, Mueller-Groeling, Speth and Wyborny '90; Hoffmann, Durso, Holinde, Pearce and Speth '95; Haidenbauer, Krein, Meissner and Tolos '11..

chiral Lagrangian

Kaiser, Siegl and Weise, '95; Oset and Ramos '98; Oller and Meissner '01; Lutz, and Kolomeitsev '02; Garcia-Recio et al. '03; Jido et al. '03; Borasoy, Nissler, and Weise '05; Oller, Prades, and Verbeni '05; Oller '06; Borasoy, Nissler and Weise '05; Khemchandani, Martinez-Torres, Nagahiro and Hosaka '12 Feijoo, Magas and Ramos '19; Feijoo, Gazda, Magas and Ramos '21; Ren, Epelbaum, Gegelia and Meissner '20 '21; Bruns and Cieply '22..

more channels, next-to-leading order, Born terms beyond WT (s-channel, u-channel), fits including new data, higher partial waves...

K=

11

u

11

 \mathbf{K}^{-}

d

 \mathbf{d}

u

u

(1405)

p

Double-pole structure of Λ(1405)

 $\Lambda(1405)$ results from the superposition of two poles in the complex plane,

$$T_{ij} \approx rac{g_i g_j}{z - z_R}$$

with different coupling to $\pi\Sigma$ and $\overline{K}N$ states

Pole positions for the $\Lambda(1405)$ coming from recent chiral effective models including the SIDDHARTA constraint.

Model		First Pole [MeV]	Second Pole [MeV]
NLO	Ikeda, Hyodo and Weise '12	$1424^{+7}_{-23} - i26^{+3}_{-14}$	$1381^{+18}_{-6} - i81^{+19}_{-8}$
Fit II	Guo and Oller '13	$1421^{+3}_{-2} - i 19^{+8}_{-5}$	$1388^{+9}_{-9} - i114^{+24}_{-25}$
Solutior	Nr. 2 Mai and Meissner '1	$51434^{+2}_{-2} - i 10^{+2}_{-1}$	$1330^{+4}_{-5} - i56^{+17}_{-11}$
Solutior	n Nr. 4	$1429^{+8}_{-7} - i 12^{+2}_{-3}$	$1325^{+15}_{-15} - i90^{+12}_{-18}$



the measured spectra of the $\Sigma\pi$ final states associated to the $\Lambda(1405)$ for kaon- and pion-induced reactions supports the double-pole structure of the $\Lambda(1405)$

Magas, Oset and Ramos '05

KNN bound state

if the KN interaction is so attractive, the K-nuclear clusters may form → The KNN (I=1/2) state



thoroughly addressed theoretically

Akaishi, Yamazaki, Shevchenko, Gal, Mares, Revai, Ikeda, Sato, Kamano, Dote, Hyodo, Weise, Wycech, Green, Bayar, Oset, Ramos, Yamagata-Sekihara, Barnea, Liverts, Dote, Inoue, Myo, Uchino, Hyodo, Oka..

initial claims by FINUDA, DISTO and OBELIX

that could find a conventional explanation Ramos et al '08 or not be reproduced Agakishiev et al [HADES] '15

more recent experiments did not find any Tokiyasu et al. [Spring8/LEPS] '14; Hashimoto et al [JPARC E15] '15; Vazquez-Doce et al. [AMADEUS] '16 or if found Ichikawa et al [J-PARC E27] '15; Nagae et al [J-PARC E27] '16 may have other interpretation Garcilazo et al '13

J-PARC E15 found a structure near KNN threshold Sada et al [J-PARC E15] '16 being interpreted as KNN bound state Sekihara et al '16 More recent J-PARC E15 measurements Ajimura et al '19; Yamaga et al '20 Binding energy and width of K⁻pp for different chiral and phenomenological calculations using variational, Faddeev or ccCSM+Feshbach methods. Tolos and Fabbietti '20

Work	B [MeV]	Г [MeV]	Method	Type of potential
Barnea et al.	16	41	Variational	Chiral
Dote et al.	17–23	40-70	Variational	Chiral
Dote et al.	14–50	16-38	ccCSM	Chiral
Ikeda et al.	9–16	34–46	Faddeev	Chiral
Bayar et al.	15-30	75–80	Faddeev	Chiral
Sekihara et al.	15-20	70–80	Faddeev	Chiral
Yamazaki et al.	48	61	Variational	phenomenological
Shevchenko et al.	50–70	90-110	Faddeev	Phenomenological
Ikeda et al.	60-95	45-80	Faddeev	Phenomenological
Wycech et al.	40-80	40-85	Variational	phenomenological
Dote et al.	51	32	ccCSM	Phenomenological
Revai et al.	32/ 47–54	50-65	Faddeev	Chiral/phenomenological

Binding energies **B~9-95 MeV** with decay widths **F~16-110 MeV**

Variety of values due to

- uncertainties in subthreshold extrapolation of the KN interaction

(chiral interactions give lower binding energies than phenomenological ones)

- use of variational or Faddeev calculations introduces certain approximations

(full three-body not account for in variational methods, whereas Faddeev calculations deal with separable two-body interactions), and ccCSM combines merits of variational and Faddeev but high computational cost

Antikaons in matter

Relativistic mean-field,

Quark meson coupling models...

RMF: early works based on mesonexchange picture or the chiral approach for the KN interaction on the mean-field level and fit the parameters to the KN scattering length



Phenomenological models

density dependent potentials fitted to kaonic atoms



U_{κ-}(ρ₀) ~ -100 to -200 MeV

recent K-N scattering amplitudes from χ SU(3) EFT supplemented with phenomenological terms for K-multinucleon interactions: kaonic atoms test densities $\rho < \rho_0$

Friedman and Gal '17

Unitarized theory in matter:

selfconsistent coupled-channel procedure



K spectral function in matter





Koch '94; Waas and Weise '97; Kaiser et al '97; Oset and Ramos'98; Lutz '98; Schaffner-Bielich et al '00; Ramos and Oset '00; Lutz et al '02; Tolos et al '01 '02; Jido et al '02 '03; Magas et al '05; Tolos et al '06 '08; Lutz et al '08; Cabrera et al '14...

 $\begin{array}{l} \text{Re } U_{\text{K-}}(\rho_0) \thicksim -50 \text{ to } -80 \text{ MeV} \\ \text{Im } U_{\text{K-}}(\rho_0) \gtrsim \text{Re } U_{\text{K-}}(\rho_0) \end{array}$

•s-wave $\overline{K}N$ interaction governed by $\Lambda(1405)$:

attraction due to modified $\Lambda(1405)$ in the medium using a self-consistent coupled-channel approach

•p-wave (and beyond)

contributions to KN interaction: not important for atoms but important for heavy-ion collisions due to large momentum

Experiments and observations: from HICs....

strangeness production in matter

is one of the major research domains in heavy-ion collisions from SIS/GSI to LHC and RHIC up to the future FAIR/NICA/BESII/J-PARC-HI



Iow-energy HICs:(FOPI) Ritman et al '95; Crochet et al '00; Bastid et al, '07; Zinyuk '14..
(KaoS) Menzel et al '00; Ploskon '05; Uhlig et al '05;Foerster et al '07..
(HADES) Agakishiev et al '09 '10 '11 '13 '14;
Galatyuk '17; Adamczewski-Musch '18 '19...FOPI/SIS18: K+, K⁻, $\phi(1020)$.Galatyuk '17; Adamczewski-Musch '18 '19...HADES/SIS18: K+, K*(892)⁰, K_s⁰, $\phi(1020)$, Λ, Ξ(1321),Ω..

high-energy HICs:

STAR/RHIC: K*(892)⁰, φ(1020), Ω.. ALICE/LHC: K*(892)⁰, φ(1020), Σ⁺⁻(1385), Ξ(1530)⁰... Adams et al. (STAR) '05 Aggarwal et al (STAR) '11 Kumar et al (STAR) '15 Abelev (ALICE) '15 Adam (ALICE) '16 Badala (ALICE) '17..

future:

CBM/FAIR BM@N/NICA BESII/RHIC J-PARC-HI

CBM (FAIR) Physics Book '11 NICA: http://theor0.jinr.ru/twiki-cgi/view/NICA Aggarwal et al (BES STAR White Paper) '10 JPARC: http://silver.j-parc.jp/sako/white-paper-v1.21.pdf-HI

K⁻ and K+ at high μ_{B} (FOPI/HADES @ SIS18)

KaoS: from systematics of the experimental results and detailed comparison to transport model calculations₁₅₀ Foerster et al (KaoS) '07

- K⁺ probe a soft EoS
- K⁺ and K⁻ yields are coupled $NN \rightarrow K^+YN$ by strangeness exchange: $K^-N \Leftrightarrow \pi Y$
- K⁺ and K⁻ exhibit different freeze-out conditions
- repulsion for K+ and attraction for K- seemed to be confirmed

but, for example, what is the role of $\phi \rightarrow K^+ K^-$?

Results from HADES and FOPI indicate

Zinyuk et al (FOPI)'14; Gasik et al (FOPI) '16; Piasecki et al (FOPI) '16; Adamczewski-Musch et al (HADES) '17..

- K⁺ in-medium potential is repulsive: U_{KN} (ρ₀)≈ 20...40 MeV
- K⁻ from Φ decay wash out the effects of the potential (spectra and flow!!)
- separate direct kaons (\rightarrow COSY)/elementary reactions
- more systematic, high statistic data on K⁻ production necessary



conclusions from Leifels-SQM2017

Recent results on kaon and antikaon production in HiCs using a PHSD model with in-medium strange mesons compared to KaoS, FOPI and HADES experimental data

- The nuclear effects on (anti)kaon are more prominent in the collision of large nuclei
- (Anti)kaon production is (enhanced)suppressed due to (broadening of spectral function)repulsive kaon potential
- (Anti)kaon spectrum becomes (softer)harder in nuclear matter, whereas y-distribution (shrinks)broadens
- Different behaviour of v1/v2 for antikaons and kaons due to the attractive vs repulsive character of the interaction with nucleons
- A moderate EoS (K~300 MeV) reproduces the experimental HiC data better



Song, LT, Wirth, Aichelin and Bratkovskaya '21

Experiments and observations: to stars

(MeV)

Kaon condensation in neutron stars

K⁻ feels attraction in the medium
→ Kaon condensation in neutron stars?

$$n \leftrightarrow p \ e^- \ \bar{\nu}_e \rightarrow \mu_n = \mu_p + \mu_{e^-}$$

$$n \leftrightarrow p K^- \rightarrow \mu_n = \mu_p + \mu_{K^-}$$





Antikaons are bosons. If $\mu_{K} \leq \mu_{e}$ for $\rho \geq \rho_{c}$, with ρ_{c} being a feasible density within neutron stars, antikaons will condensate

Using microscopic unitarized schemes...

The condition $\mu_{e_-} \ge m^*_{K_-}$ for a given ρ_c implies that $m_{K_-} - m^*_{K_-} (\rho_c) \approx 200, 300 \text{ MeV}.$ However, unitarized schemes based on meson-exchange models or chiral (MeV) Lagrangians predict a moderate attraction in nuclear matter

Lutz '98 Ramos and Oset '00 Tolos, Polls, Ramos '01 Tolos, Ramos and Oset '06 Tolos, Cabrera and Ramos '08 Cabrera, Tolos, Aichelin and Bratkovskaya'14

Therefore,

kaon condensation seems very unlikely within microscopic unitarized schemes



Present and Future



A lot of experimental and theoretical effort has been invested to understand the \overline{KN} interaction, that is governed by the presence of the $\Lambda(1405)$

A lot of effort has been invested in unveiling the nature of $\Lambda(1405)$, and the consequences for the formation of $\overline{K}NN$ bound state

Kaons and antikaons in matter have been also investigated in connection to strangeness in nuclear collisions and kaon condensation in neutron stars



