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Double-beta decay:

and new results

from EXO-200

SPP 2012, Groningen, Jun 2012

## Double-beta decay:

a second-order process only detectable if first order beta decay is energetically forbidden



#### Candidate nuclei with Q>2 MeV

Candidate	Q	Abund.
	(MeV)	(%)

<sup>48</sup> Ca→ <sup>48</sup> Ti	4.271	0.187
<sup>76</sup> Ge→ <sup>76</sup> Se	2.040	7.8
<sup>82</sup> Se→ <sup>82</sup> Kr	2.995	9.2
$^{96}$ Zr $\rightarrow$ <sup>96</sup> Mo	3.350	2.8
$^{100}Mo \rightarrow ^{100}Ru$	3.034	9.6
<sup>110</sup> Pd→ <sup>110</sup> Cd	2.013	11.8
$^{116}Cd \rightarrow ^{116}Sn$	2.802	7.5
<sup>124</sup> Sn→ <sup>124</sup> Te	2.228	5.64
<sup>130</sup> Te→ <sup>130</sup> Xe	2.533	34.5
<sup>136</sup> Xe→ <sup>136</sup> Ba	2.458	8.9
$^{150}Nd \rightarrow ^{150}Sm$	3.367	5.6

### There are two varieties of $\beta\beta$ decay

2v mode: a conventional 2<sup>nd</sup> order process in nuclear physics


## "Dirac" neutrinos

(some "redundant" information but the "good feeling" of things we know...)

## "Majorana" neutrinos

(more efficient description, no lepton number conservation, new paradigm...)





Which way Nature chose to proceed is an experimental question

→ But the alternative is only meaningful/testable for massive particles... which we now know neutrinos are!

DoubleBeta decay

### Our knowledge of the v mass pattern



The connection of v masses with cosmological measurements is particularly interesting because it ties together very different fields. We need both, the connection between the two is the interesting part!

DoubleBeta decay

### In the last 10 years there has been a transition

1) From a few kg detectors to 100s or 1000s kg detectors → Think big: qualitative transition from cottage industry to large experiments

2) From "random shooting" to the knowledge that at least the inverted hierarchy will be tested

Discovering Ovββ decay: → Discovery of the neutrino mass scale → Discovery of Majorana particles → Discovery of Majorana masses → Discovery of lepton number violation

### If $0v\beta\beta$ is due to light v Majorana masses

$$\left\langle m_{\nu}\right\rangle^{2} = \left(T_{1/2}^{0\nu\beta\beta} G^{0\nu\beta\beta}(E_{0},Z) \left|M_{GT}^{0\nu\beta\beta} - \frac{g_{\nu}^{2}}{g_{A}^{2}}M_{F}^{0\nu\beta\beta}\right|^{2}\right)^{-1}$$

$$M_{F}^{\,0
uetaeta}$$
 and  $M_{GT}^{\,0
uetaeta}$ 

 $G^{0
uetaeta}$ 

 $T_{1/2}^{0\nu\beta\beta}$ 

can be calculated within particular nuclear models

a known phasespace factor

is the quantity to be measured

$$\langle m_{v} \rangle = \sum_{i=1}^{3} \left| U_{e,i} \right|^{2} m_{i} \mathcal{E}_{i}$$

effective Majorana v mass ( $\varepsilon_i = \pm 1$  if CP is conserved)



Note, however, that to discover Majorana neutrinos and lepton number violation the value of the nuclear matrix element is inessential!

 $\rightarrow$  0v $\beta\beta$  decay always implies new physics

This is comforting for the ones of us spending their time building experiments!

### Simplified List of Limits for BBOv decay

Candidate	Detector		Present	<m> (eV)</m>
nucleus	type	(kg yr)	T <sub>1/2</sub> <sup>0νββ</sup> (yr)	
48 <b>Ca</b>			>5.8*10 <sup>22</sup> (90%CL)	
<sup>76</sup> Ge	Ge diode	47.7	>1.9*10 <sup>25</sup> (90%CL)	<0.35
<sup>82</sup> Se			>2.1*10 <sup>23</sup> (90%CL)	
<sup>96</sup> Zr			>9.2*10 <sup>21</sup> (90%CL)	
<sup>100</sup> Mo	Foil.Geiger	tubes	>5.8*10 <sup>23</sup> (90%CL)	
<sup>116</sup> Cd			>1.7*10 <sup>23</sup> (90%CL)	
<sup>128</sup> Te			>1.1*10 <sup>23</sup> (90%CL)	
<sup>130</sup> Te	TeO <sub>2</sub> cryo	~12	>3*10 <sup>24</sup> (90%CL)	<0.19-0.68
<sup>136</sup> Xe	Xe scint	~4.5	>1.2*10 <sup>24</sup> (90%CL)	<del>→1.1-2.9</del>
	Xe TPC	32.3	>1.6*10 <sup>25</sup> (90%CL)	<0.14-0.38
<sup>150</sup> Nd			>1.8*10 <sup>22</sup> (90%CL)	
<sup>160</sup> <b>Gd</b>			>1.3*10 <sup>21</sup> (90%CL)	

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## $\beta\beta0\nu$ discovery claim



Fit model:

6 gaussians + linear bknd.

Fitted excess @  $Q_{\beta\beta}$  28.75 ± 6.86.

Claimed significance: 4.2  $\sigma$ 

$$T_{1/2} = 2.23^{+0.44}_{-0.31} \cdot 10^{24} yr$$
$$\langle m_{\nu} \rangle = 0.32 \pm 0.03 \ eV$$

[H.V.Klapdor-Kleingrothaus and I.Krivosheina, Mod.Phys.Lett. A21 (2006) 1547]

However, this is a very controversial matter

See e.g. Strumia+Vissani Nucl Phys B726 (2005) 294

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Need very	y large f	iducial
mass	(tons) of	isotopically
separa	ted mate	erial
(except	t for <sup>130</sup>	Ге)

[using natural material typically means that 90% of the source produced background but not signal]

This is expensive and provides encouragement to use the material in the best possible way:

For no bkgnd 
$$\langle m_{\nu} \rangle \propto 1 / \sqrt{T_{1/2}^{0\nu\beta\beta}} \propto 1 / \sqrt{Nt}$$

Candidate	Q (MeV)	Abund. (%)
<sup>48</sup> Ca→ <sup>48</sup> Ti	4.271	0.187
<sup>76</sup> Ge→ <sup>76</sup> Se	2.040	7.8
<sup>82</sup> Se→ <sup>82</sup> Kr	2.995	9.2
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<sup>150</sup> Nd→ <sup>150</sup> Sm	3.367	5.6

For statistical bkgnd subtraction

$$\langle m_{\nu} \rangle \propto 1 / \sqrt{T_{1/2}^{0\nu\beta\beta}} \propto 1 / (Nt)^{1/4}$$

### How to "organize" an experiment: the source



- High Q value reduces backgrounds and increases the phase space & decay rate,
- Large abundance makes the experiment cheaper
- A number of isotopes have similar matrix element performance

### How to "organize" an experiment: the technique

- Final state ID: 1) "Geochemical": search for an abnormal abundance
  - of (A,Z+2) in a material containing (A,Z)
  - 2) "Radiochemical": store in a mine some material (A,Z)
    - and after some time try to find (A,Z+2) in it
    - + Very specific signature
    - + Large live times (particularly for 1)
    - + Large masses
    - Possible only for a few isotopes (in the case of 1)
    - No distinction between Ov, 2v or other modes
- "Real time": ionization or scintillation is detected in the decay
  - a) "Homogeneous": source=detector
  - b) "Heterogeneous": source # detector
    - + Energy/some tracking available (can distinguish modes)
    - + In principle universal (b)
    - Many  $\gamma$  backgrounds can fake signature
    - Exposure is limited by human patience

# Shielding a detector from gammas is difficult because the absorption cross section is small.



Example: y interaction length in Ge is 4.6 cm, comparable to the size of a germanium detector.

Shielding *BB* decay detectors is much harder than shielding Dark Matter ones We are entering the "golden era" of *BB* decay experiments as detector sizes exceed int lengths

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**DoubleBeta decay** 



### <u>The two can be separated in a detector with</u> <u>sufficiently good energy resolution</u>

Topology and particle ID are also important to recognize backgrounds

### Some experiments in preparation

(~approved or under construction, in addition a number of R&D efforts)

Isotope	Experiment	Main principle	Fid mass	Status	Lab
	Majorana <sup>†</sup>	Eres,2site tag, Cu shield	30-60kg	Construction	SUSEL
<sup>76</sup> Ge	Gerda <sup>†</sup>	Eres,2site tag, LAr shield	34.3 kg	Data taking	G Sasso
	MaGe/GeMa	See above	~1ton	Planning	?
<sup>150</sup> Nd	SNO+	Size/shielding	44 kg	Construction	SNOlab
<sup>82</sup> Se	SuperNEMO <sup>‡</sup>	Tracking	100 kg	Planning	Canfranc Frejus
<sup>130</sup> Te*	CUORE	E Res.	204 kg	Construction	G Sasso
<sup>136</sup> Xe	KamLAND-Zen	Size/shielding	400 kg	Data taking	Kamioka
1362	EXO	Tracking/Eres	150 kg	Data taking	WIPP
<sup>130</sup> Xe		Ba tag, Track/Eres	1-10ton	Planning	SNOlab?

\* No isotopic enrichment in baseline design

<sup>+</sup> Plan to merge efforts for ton-scale experiment

SPP 2012, Groningen Jun 2 <sup>†</sup> Non-homogeneous detector

It is very important to understand that a healthy neutrinoless double-beta decay program requires more than one isotope. This is because:

- There could be unknown gamma transitions and a line observed at the "end point" in one isotope does not necessarily imply that Ovßß decay was discovered
- Nuclear matrix elements are not very well known and any given isotope could come with unknown liabilities
- Different isotopes correspond to vastly different experimental techniques
- 2 neutrino background is different for various isotopes
- The elucidation of the mechanism producing the decay requires the analysis of more than one isotope

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## Xe is ideal for a large experiment

- No need to grow crystals
- Can be re-purified during the experiment
- No long lived Xe isotopes to activate
- Can be easily transferred from one detector to another if new technologies become available
- Noble gas: easy(er) to purify
- <sup>136</sup>Xe enrichment easier and safer:
  - noble gas (no chemistry involved)
  - centrifuge feed rate in gram/s, all mass useful
  - centrifuge efficiency  $\sim \Delta m$ . For Xe 4.7 amu
- Only known case where final state identification appears to be not impossible
  - $\rightarrow$  elominate all non- $\beta\beta$  backgrounds

## • <sup>129</sup>Xe is a hyperpolarizable nucleus, under study for NMR

tomography... a joint enrichment program ?

# The EXO-200 TPC Two almost identical halves reading ionization and 178 nm scintillation, each with:

- 38 U triplet wire channels (charge)
- 38 V triplet wire channels, at 60° (induction)
- 234 large Avalanche PhotoDiodes (in gangs of 7)
- Triplet pitch 9 mm
- Wire planes 6 mm apart and 6 mm from APDs
  - Signals digitized at 1 MS/s, ±1024s around trigger
- Drift field 376 V/cm
  - Field shaping rings: copper
  - Supports: acrylic
  - Light reflectors/diffusers: Teflon
  - APD support plane: copper; Au (Al) coated for contact (light reflection)
  - Central cathode, U+V wires: photoetched phosphor bronze
  - Flex cables for bias/readout: copper on kapton, no glue
  - Vast material screening program
  - → Goal: 40 cnts/2y in 0vββ ±2σ ROI, 140 kg LXe 20

-40 cm-

Copper vessel 1.37 mm thick
175 kg LXe, 80.6% enr. in <sup>136</sup>Xe
Copper conduits (6) for:
APD bias and readout cables
U+V wires bias and readout
LXe supply and return
Epoxy feedthroughs at cold and warm doors
Dedicated HV bias line

EXO-200 detector: Characterization of APDs: Materials screening: JINST 7 (2012) P05010 NIM A608 68-75 (2009) NIM A591, 490-509 (2008)

## The EXO-200 Detector



## Data taking phases and Xenon Purity

	Run I	Run 2 (this analysis)
Period	May 21, 11 – Jul 9, 11	Sep 22, 11 - Apr 15,12
Live Time	752.7 hr	2,896.6 hr
Exposure	3.2 kg-yr	32.5 kg-yr
Publ.	PRL 107 (2011) 212501	arXiv:1205:5608 (May 2012)

Xenon gas is forced through heated Zr getter by a custom ultraclean pump.

Electron lifetime  $T_e$ :  $\rightarrow$  measure ionization signal attenuation as a function of drift time for the full-absorption peak of  $\gamma$  ray sources

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At Te = 3 ms:

- drift time <110 μs

- loss of charge: 3.6%

at full drift length
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Ultraclean pump:

Rev Sci Instr. 82 (10) 105114

Xenon purity with mass spec:

NIM A675 (2012) 40

Gas purity monitors:

NIM A659 (2011) 215
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23







### T<sub>1/2</sub> = (2.11 ± 0.04 stat ± 0.21 sys) · 10<sup>21</sup> yr

[Ackerman et al Phys Rev Lett 107 (2001) 212501]

In significant disagreement with previous limits:  $T_{1/2} > 1.0 \cdot 10^{22}$  yr (90% C.L.) (R. Bernabei *et al.* Phys. Lett. B 546 (2002) 23)

T<sub>1/2</sub> > 8.5 · 10<sup>21</sup> yr (90% C.L.) (Yu. M. Gavriljuk *et al.*, Phys. Atom. Nucl. 69 (2006) 2129)

Later confirmed by KamLAND-ZEN T<sub>1/2</sub>=(2.38 ± 0.02stat ± 0.14sys) · 10<sup>21</sup> yr [A.Gando et al. Phys Rev C 85 (2012) 045504]

## **Combining Ionization and Scintillation**



# **Energy Calibration**



## Source Data/MC Agreement



- Single site fraction agrees to within 8.5%
- Source activities measured to within 9.4%

## **Rn Content in Xenon**



Long-term study shows a constant source of <sup>222</sup>Rn dissolving in <sup>enr</sup>LXe: 360 ± 65 µBq (Fid. vol.)

**DoubleBeta decay** 

## Low Background 2D SS Spectrum



#### Events removed by diagonal cut:

- $\alpha$  (larger ionization density  $\rightarrow$  more recombination  $\rightarrow$  more scintillation light)
- events near detector edge  $\rightarrow$  not all charge is collected

### Low Background Spectrum



Low background run livetime: 120.7 days

Active mass: 98.5 kg LXe (79.4kg <sup>136</sup>LXe)

Exposure: 32.5 kg.yr

Vetos dead time: 8.6%

 ββ2ν

 ββ0ν (90% CL Limit)

 <sup>40</sup>K LXe Vessel

 <sup>54</sup>Mn LXe Vessel

 <sup>60</sup>Co LXe Vessel

 <sup>65</sup>Zn LXe Vessel

 <sup>232</sup>Th LXe Vessel

 <sup>238</sup>U LXe Vessel

 <sup>135</sup>Xe Active LXe

 <sup>222</sup>Rn Active LXe

 <sup>222</sup>Rn Inactive LXe

 <sup>214</sup>Bi Cathode Surface

 <sup>222</sup>Rn Air Gap



#### Low background spectrum zoomed around the Ovßß region of interest (ROI)



 ββ2ν

 ββ0ν (90% CL Limit)

 40K LXe Vessel

 54Mn LXe Vessel

 60Co LXe Vessel

 65Zn LXe Vessel

 232Th LXe Vessel

 238U LXe Vessel

 135Xe Active LXe

 222Rn Active LXe

 222Rn Inactive LXe

 214Bi Cathode Surface

 222Rn Air Gap

 Data

 Total

No Ov signal observed in the ROI

Use likelihood fit to establish limit

## Background counts in $\pm 1,2 \sigma$ ROI



	Expected eve		ents from fit	
	ť	1 σ	±2	2σ
<sup>222</sup> Rn in cryostat air-gap	1.9	±0.2	2.9	±0.3
<sup>238</sup> U in LXe Vessel	0.9	±0.2	1.3	±0.3
<sup>232</sup> Th in LXe Vessel	0.9	±0.1	2.9	±0.3
<sup>214</sup> Bi on Cathode	0.2	±0.01	0.3	±0.02
All Others	~0.2		~0.2	
Total	4.1	±0.3	7.5	±0.5
Observed		1	;	5
Background index b (kg <sup>-1</sup> yr <sup>-</sup> <sup>1</sup> keV <sup>-1</sup> )	1.5.10	<sup>-3</sup> ± 0.1	1.4.10	<sup>-3</sup> ± 0.1

## Limits on $T_{1/2}^{Ov\beta\beta}$ and $\langle m_{\beta\beta} \rangle$



# Summary

- Several new experiments started taking data in the last year
- As usual, new experiments are much more powerful than the previous generation
- Expect rapid progress for the next few years
- Stay tuned for more results!



# The EXO collaboration





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## Systematics and sensitivity



Error breakout: expected 90% CL limit given absolute knowledge (0 error) of a given parameter or set of parameters

Term	%
Fiducial Volume	12.34
β scale	9.32
SS / (SS + MS)	0.93
<sup>232</sup> Th LXe Vessel	0.11
<sup>238</sup> U LXe Vessel	0.04
<sup>222</sup> Rn Air Gap	0.04
Calibration offsets	0.04

Distribution of Ovßß T1/2 90% CL



From estimated background, expect to quote a 90% CL upper limit on T<sub>1/2</sub> :

2	1.6	$x \ 10^{25} \ yr$	6.5%	of the time
2	7	x 10 <sup>24</sup> yr	50%	of the time