

Effective spectral functions via lifetime analysis

Renan Hirayama

in collaboration with

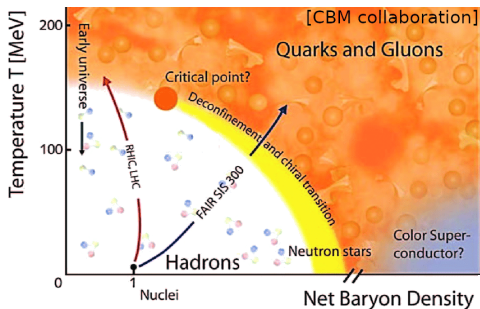
Jan Staudenmaier, Hannah Elfner

FIAS, Frankfurt am Main, Germany

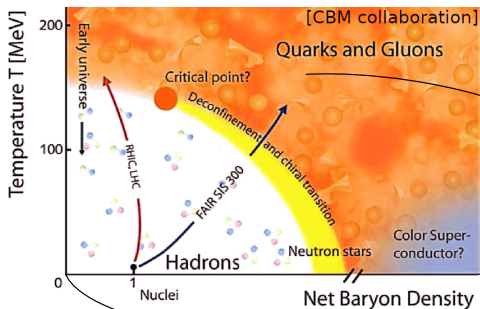
2022, May 27



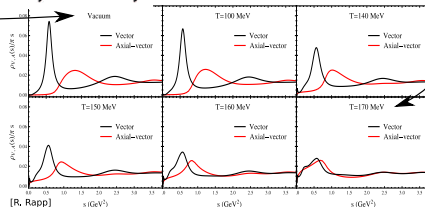
Motivation - Theory



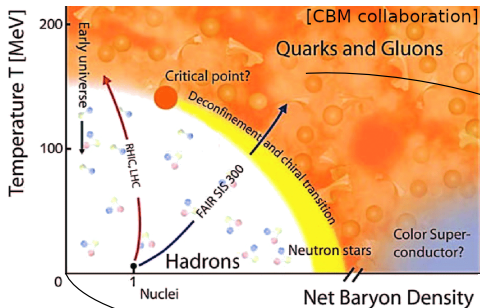
Motivation - Theory



Chiral pairs degenerate



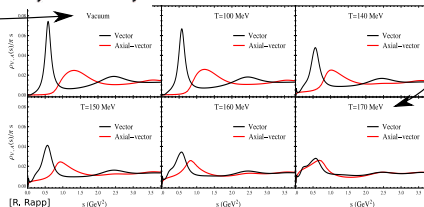
Motivation - Theory



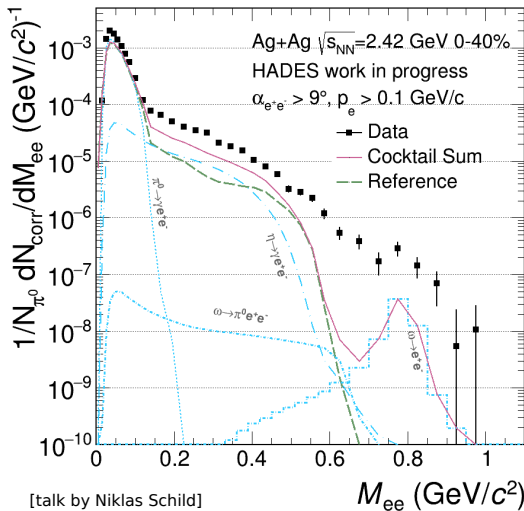
Chiral pairs degenerate

Vector : $\rho(770)$

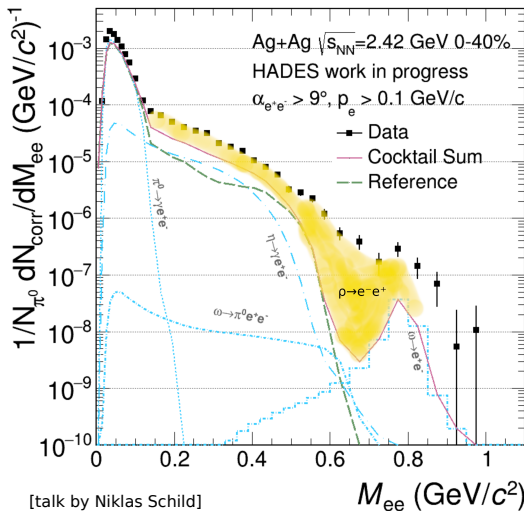
Axial : $a_1(1260)$



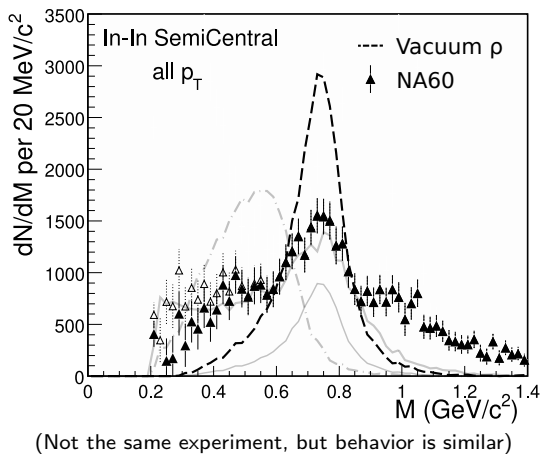
Motivation - Experiment



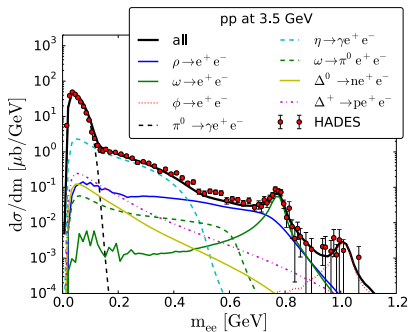
Motivation - Experiment



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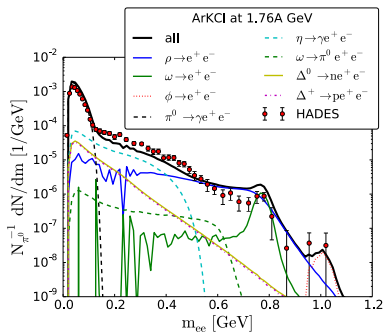
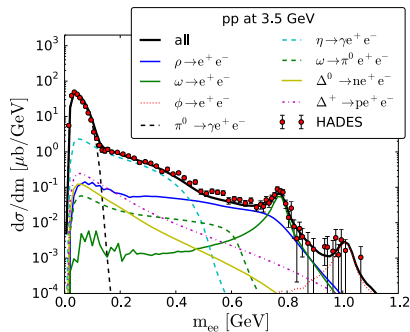


SMASH dileptons



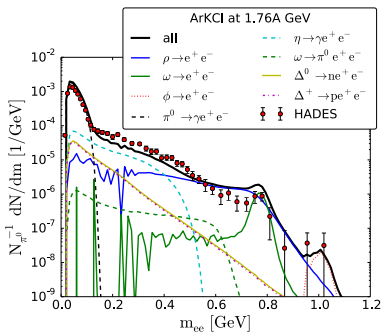
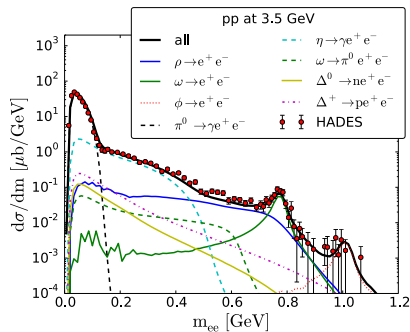
[PRC 98, 054908 (2018)]

SMASH dileptons



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SMASH dileptons



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Goal: reconstruct the ρ -meson spectral function

Addition of inelastic interactions where the ρ is absorbed

Vacuum

$$\rho^{\pm} \rightarrow \pi^0 \pi^{\mp}$$

$$\rho^0 \rightarrow \pi^+ \pi^-$$

$$\rho^0 \rightarrow l^+ l^-$$

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Medium

$$\rho N \rightarrow N^*, \Delta, \dots$$

$$\rho K \rightarrow K^*$$

$$\rho \pi \rightarrow \omega, \phi, f_1, a_1, \dots$$

Medium modifications

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Collisional broadening: *Shortening the average lifetime of resonances due to absorptions by a medium*



Simulating **M**any **A**ccelerated
Strongly-interacting **H**adrons

Transport approach: the relativistic Boltzmann equation

$$p^\mu \partial_\mu f_i(x, p) + m_i F^\alpha \partial_\alpha^p f_i(x, p) = C_i^{\text{coll}}$$



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- $f_i(x, p)$ realized as a test-particle (or more)
- F^α from MF potential

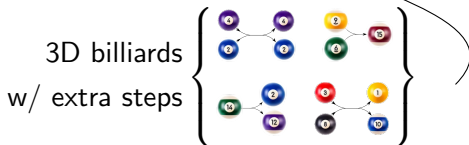


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Resonances are handled with vacuum properties:

- Decay width (=rate of decay)

$$\Gamma^{\text{vac}}(m) = \Gamma_{\rho \rightarrow \pi\pi}(m) + \Gamma_{\rho \rightarrow ll}(m)$$

No *a priori* knowledge of medium!

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- Breit-Wigner for mass sampling

$$\mathcal{A}(m) = \frac{2\mathcal{N}}{\pi} \frac{m^2 \Gamma^{\text{vac}}(m)}{(m^2 - M_0^2)^2 + m^2 \Gamma^{\text{vac}}(m)^2}$$

No *a priori* knowledge of medium!

From the interaction history:

Effective width

$$\Gamma^{\text{eff}} = \frac{1}{\langle \tau \rangle} = \left\langle \frac{\gamma}{t_f - t_i} \right\rangle$$

An effective description

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Collisional width



$$\Gamma^{\text{col}} = \Gamma^{\text{eff}} - \Gamma^{\text{vac}}$$

An effective description

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$$\Gamma^{\text{col}} = \Gamma^{\text{eff}} - \Gamma^{\text{vac}}$$

Dynamic spectral function

$$\mathcal{A}^{\text{dyn}}(m) = \frac{2\mathcal{N}'}{\pi} \frac{m^2 \Gamma^{\text{eff}}(m)}{(m^2 - M_0^2)^2 + m^2 \Gamma^{\text{eff}}(m)^2}$$

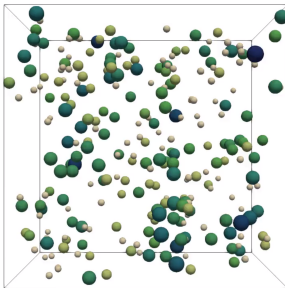
SMASH configuration: box

Box

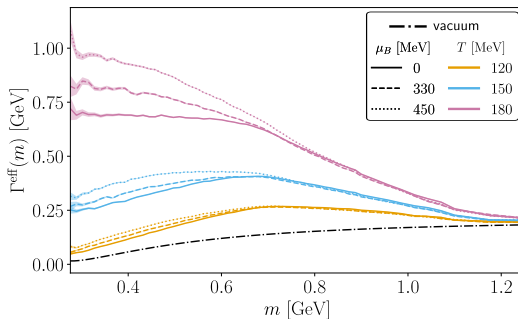
Time: 0 fm

Box Width: 10 fm

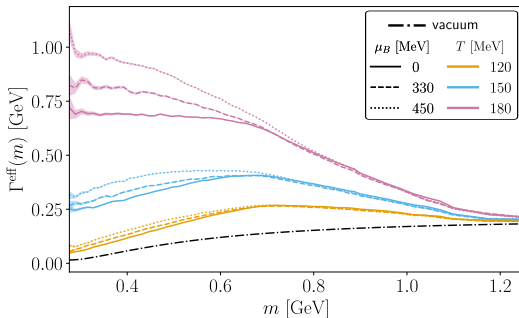
Temperature: 0.15 GeV



Thermodynamic behavior

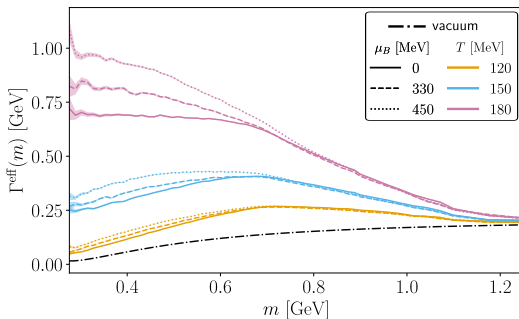


Thermodynamic behavior



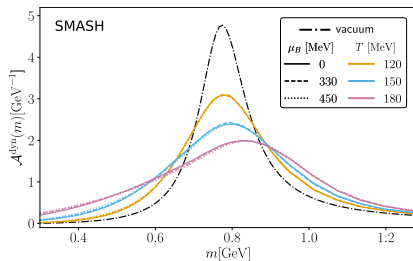
- High-mass ρ suffers little broadening (coupling $\sim 1/m^4$)
- Lower masses are mostly absorbed

Thermodynamic behavior



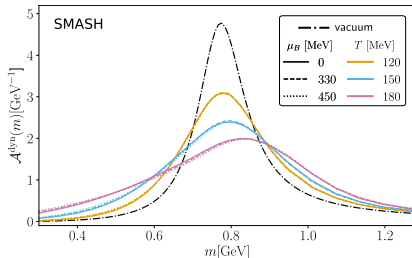
- High-mass ρ suffers little broadening (coupling $\sim 1/m^4$)
- Lower masses are mostly absorbed
- Medium temperature changes whole spectrum
- Baryo-chemical potential affects only $m \leq M_0$ GeV

SMASH

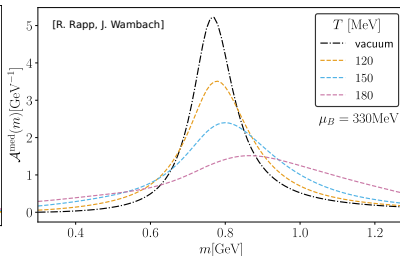


Thermodynamic behavior

SMASH

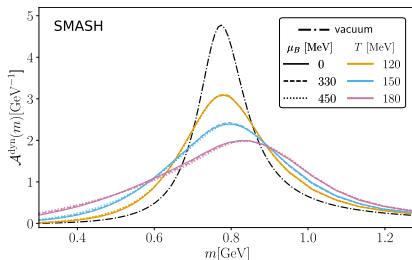


full in-medium model

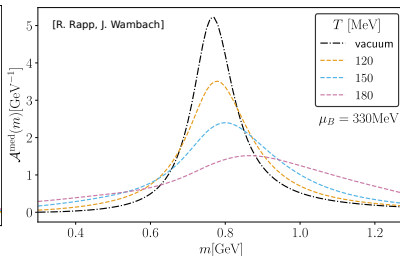


- Melting of ρ ✓
- Positive shift of peak mass ✓

SMASH



full in-medium model



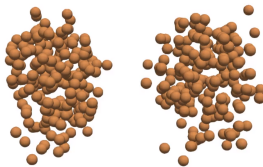
- Melting of ρ ✓
- Positive shift of peak mass ✓
- Differences: distinct processes present in SMASH ✗
(intermediate resonances, no self-energy, ...)

SMASH configuration: collider

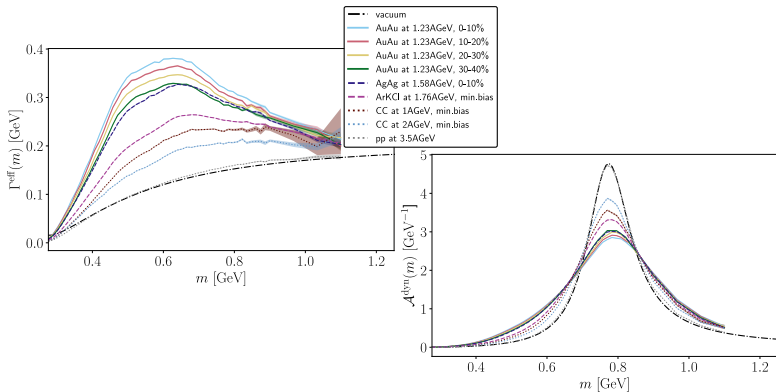
Au+Au at $E_{\text{kin}} = 1.23A \text{ GeV}$

Impact: 0.0 fm

Time: -4 fm

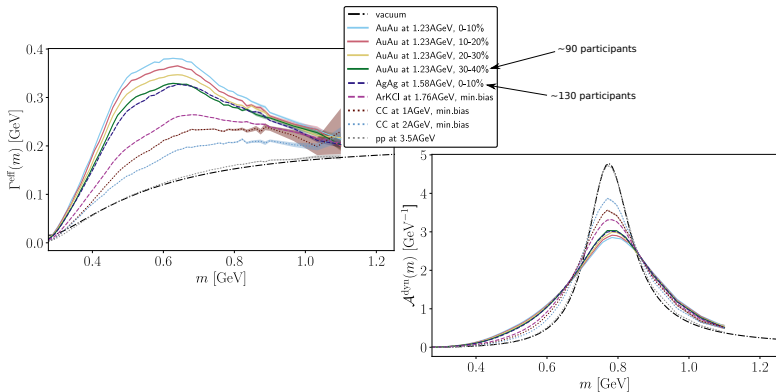


Nuclear collisions



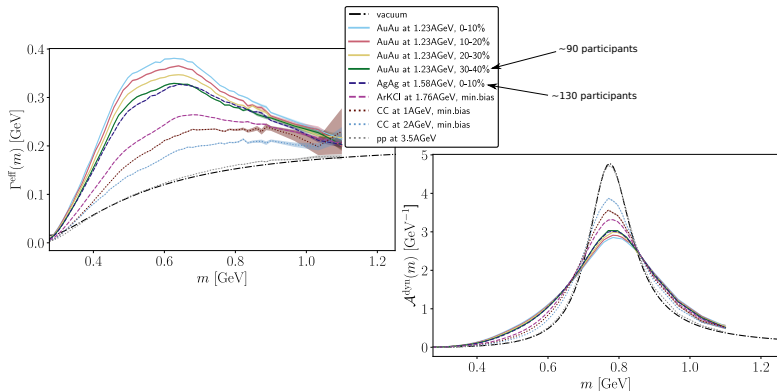
- System size dependence

Nuclear collisions



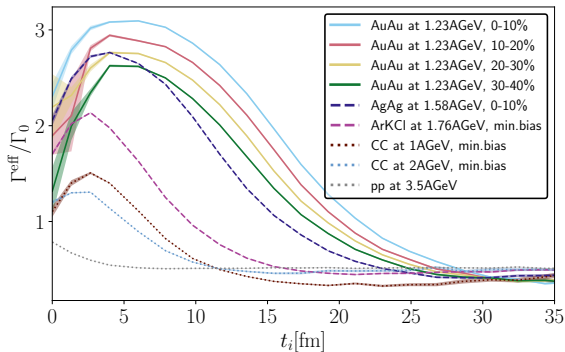
- System size dependence
- Higher energies disperse the medium

Nuclear collisions



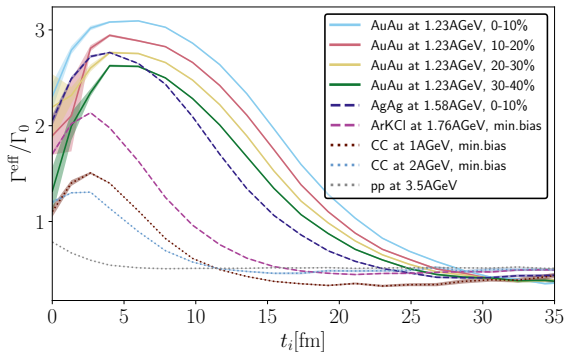
- System size dependence
- Higher energies disperse the medium
- No broadening at hadronic threshold $m = m_{2\pi}$

Nuclear collisions



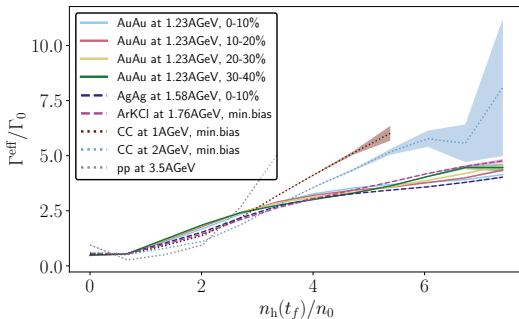
- Cronometer: how long does the medium last?

Nuclear collisions



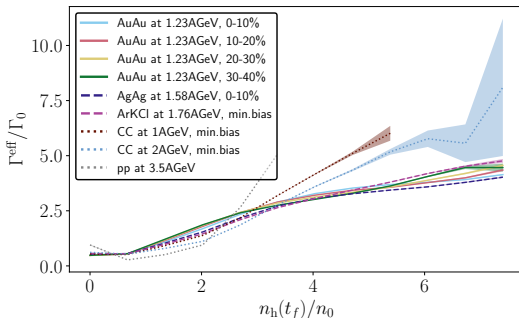
- Cronometer: how long does the medium last?
- Highlights difference between AuAu 1.23 AGeV 30-40% and AgAg 1.58 AGeV 0-10%

Nuclear collisions



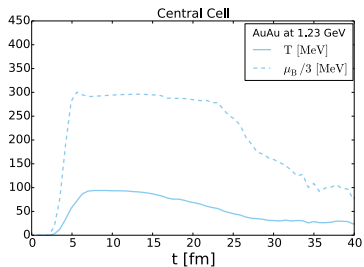
- Near universal dependence on the hadron density n_h
- Off-shell models: $\Gamma^{\text{coll}} = \gamma n_N \langle v \sigma_{VN}^{\text{tot}} \rangle$ (GiBUU, HSD)

Nuclear collisions

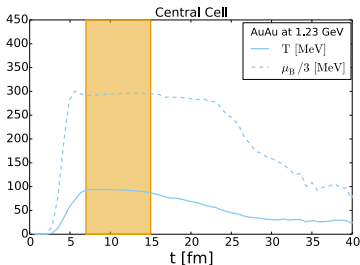


- Near universal dependence on the hadron density n_h
- Off-shell models: $\Gamma^{\text{coll}} = \gamma n_N \langle v \sigma_{VN}^{\text{tot}} \rangle$ (GiBUU, HSD)
- High densities in small systems: numerical artifacts

Non-equilibrium effects

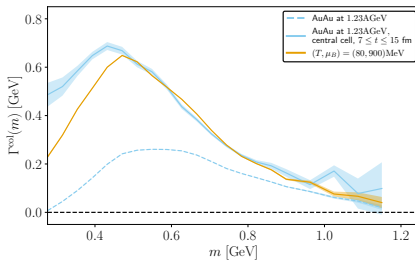
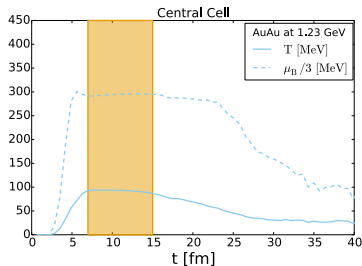


Non-equilibrium effects



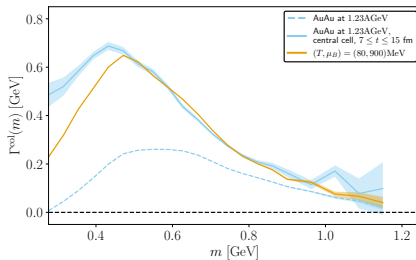
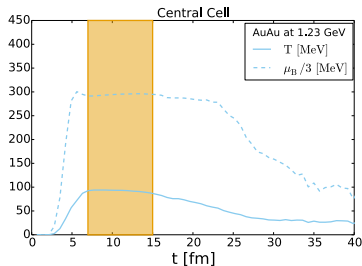
- Region of spacetime in a collision where an apparent equilibrium is attained

Non-equilibrium effects



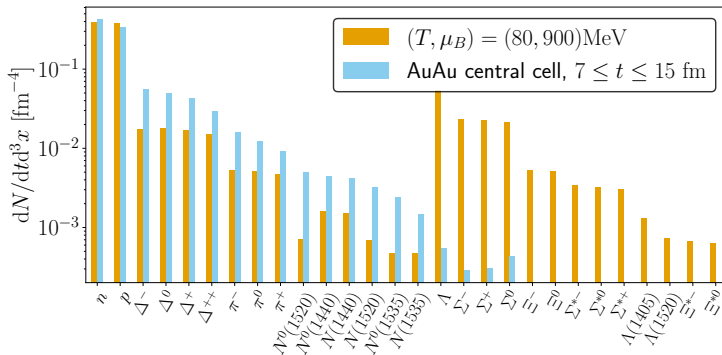
- Region of spacetime in a collision where an apparent equilibrium is attained
- Restriction to this region makes $\Gamma^{\text{col}}(m_{2\pi}) \neq 0$

Non-equilibrium effects



- Region of spacetime in a collision where an apparent equilibrium is attained
- Restriction to this region makes $\Gamma^{\text{col}}(m_{2\pi}) \neq 0$
- Similar broadening, excess for small masses

Non-equilibrium effects



- Collision systems are much more baryon-rich \Rightarrow excess in LMR
- Remaining energy in gas is filled by strange hadrons

- In SMASH: Vacuum properties $\xrightarrow{\text{medium}}$ dynamical broadening
- Equilibrium \mathcal{A}^{dyn} similar to full in-medium calculations
- Collision systems: clear setup dependence (mass number, centrality, beam energy), universality in density
- Effective width works as a cronometer for medium duration
- Same $(T, \mu_B) \not\Rightarrow$ same broadening



Have a nice lunch!

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- P. Salabura and J. Stroth, Prog. Part. Nucl. Phys. 120, 103869 (2020).
- H. van Hees and R. Rapp, Nuc. Phys. A, 806 (1-4), 339-387 (2008).
- J. Weil, et al., Phys. Rev. C94 (5), 054905 (2016).
- J. Staudenmaier, et al., Phys. Rev. C98, 054908 (2018).
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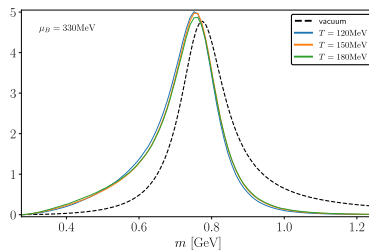
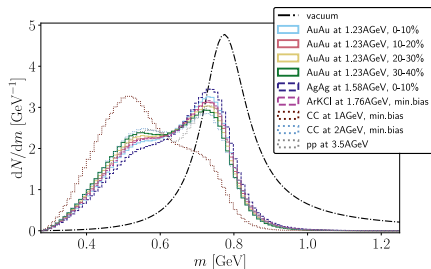
Following Manley et al.,

$$\Gamma_{\rho \rightarrow \pi\pi}^{\text{vac}}(m) = \Gamma^0 \frac{M_0}{m} \left(\frac{\frac{1}{4}m^2 - m_\pi^2}{\frac{1}{4}M_0^2 - m_\pi^2} \right)^{3/2} \left(\frac{\frac{1}{4}M_0^2 - m_\pi^2 + \Lambda^2}{\frac{1}{4}M_0^2 - m_\pi^2 + \Lambda^2} \right)$$

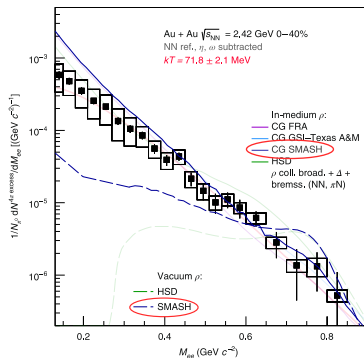
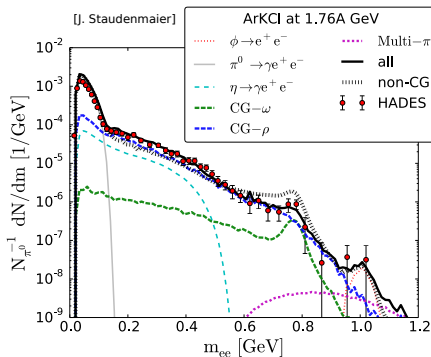
Under the *Vector Meson Dominance* model:

$$\Gamma_{\rho \rightarrow ll}^{\text{vac}}(m) = \Gamma_{\rho \rightarrow ll}^0 \left(\frac{M_0}{m} \right)^3 \left(1 + \frac{2m_l^2}{m^2} \right) \sqrt{1 - \frac{4m_l^2}{m^2}}$$

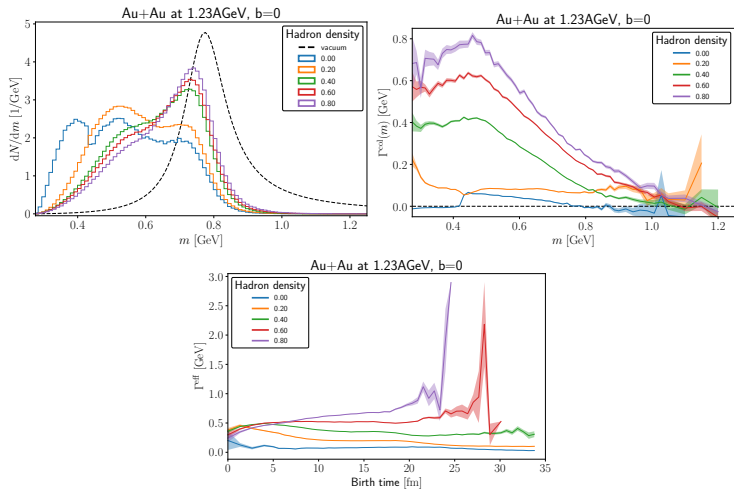
Produced ρ masses



Coarse-Grained SMASH dileptons



Density bins



Late stages of a collision

