



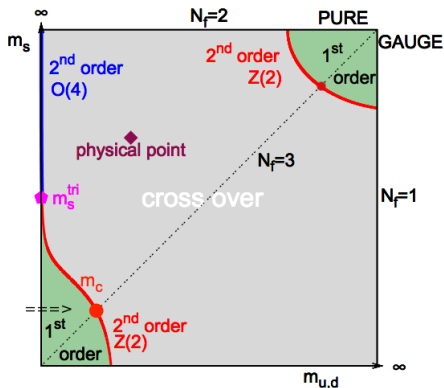
BERGISCHE
UNIVERSITÄT
WUPPERTAL

THE UPPER RIGHT CORNER OF THE COLUMBIA PLOT WITH STAGGERED FERMIONS

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and

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D. Sexty

Fairness 2022



Forcrand et. al. 1702.00330

Columbia Plot for $\mu = 0$

- $m_q \rightarrow \infty$: Pure gauge theory (quenched case)
- Latent heat in cont. lim. Shirogane PRD 16, Borsanyi PRD 22 $\Rightarrow$ 1st order
- Phys. quark masses and $\mu = 0$: Crossover
- Which value takes the (heavy) critical mass m_c ?

What are the reasons for the structures?

QCD Lagrangian

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu D_\mu - m)\psi - \frac{1}{4}F_{\mu\nu}^a F_a^{\mu\nu}$$

- Exact global and local SU(3) gauge theory

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Center Symmetry $\mathcal{Z}_3$ for $m \rightarrow \infty$

- Invariance under $U_4(\vec{x}, t_0) \rightarrow zU_4(\vec{x}, t_0)$, $z \in \{1, e^{\pm i2\pi/3}\}$, $U \in \text{SU}(3)$
- SSB at high temperatures (deconfinement)
- Order parameter Polyakov loop: $|\langle P \rangle| \sim e^{-F/T}$

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Chiral symmetry for $m \rightarrow 0$

- $\mathcal{L}_F = \bar{\psi}_L i \not{D} \psi_L + \bar{\psi}_R i \not{D} \psi_R$
- SSB at moderate temperatures
- Order parameter chiral condensate $\langle \bar{\psi} \psi \rangle = \langle \bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L \rangle$

Chiral observables

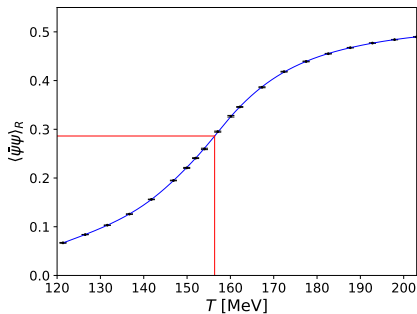
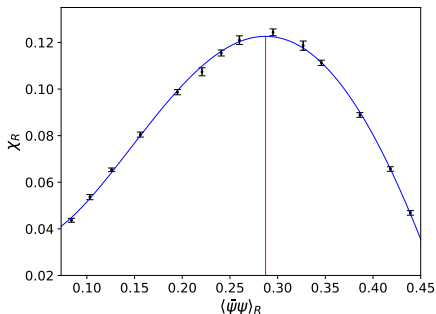
$$\langle \bar{\psi}\psi \rangle = \frac{T}{V} \frac{\partial \log Z}{\partial m}$$

$$\langle \bar{\psi}\psi \rangle_R = - [\langle \bar{\psi}\psi \rangle_T - \langle \bar{\psi}\psi \rangle_{T=0}] \frac{m}{f_\pi^4}$$

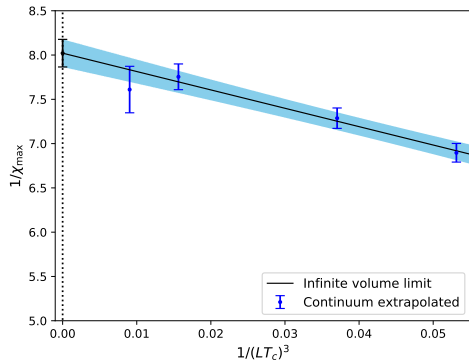
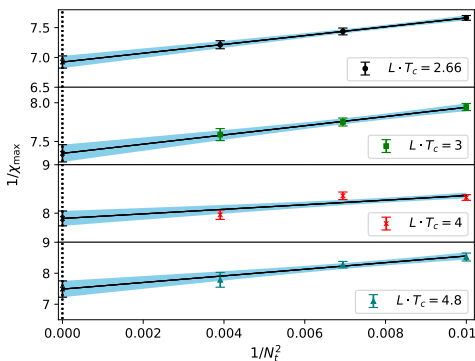
$$\chi = \frac{T}{V} \frac{\partial^2 \log Z}{\partial m^2}$$

$$\chi_R = [\chi_T - \chi_{T=0}] \frac{m^2}{f_\pi^4}$$

How to determine precisely the inflection point of $\langle \bar{\psi}\psi \rangle$ or $\chi_{\max}$?



What is the type of transition?



Analysis

- No real phase transition
- Weak crossover¹ for $\mu = 0$ (early universe)

¹First time: Fodor et. al. Nature 2006

Dependence of the strength and width on μ

Definition of the transition width (Gaussian)

$$(\Delta T)^2 = -\chi(T_c) \left[\frac{d^2\chi}{dT^2}(T_c) \right]^{-1}$$

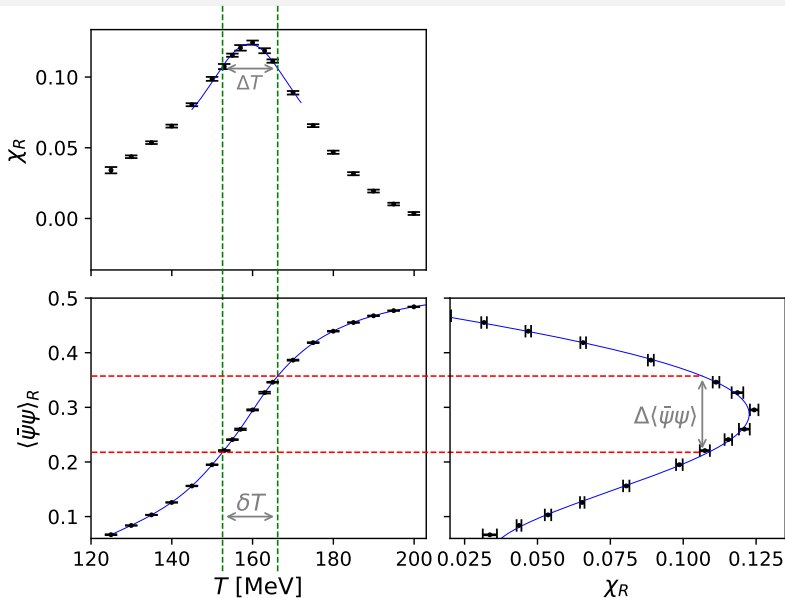
- Second derivative without appropriate fitfunction?

$$(\Delta \langle \bar{\psi}\psi \rangle)^2 = -\chi(T_c) \left(\frac{d^2\chi}{d\langle \bar{\psi}\psi \rangle^2} \bigg|_{\langle \bar{\psi}\psi \rangle_c} \right)^{-1}$$

$$\delta T = \langle \bar{\psi}\psi \rangle^{-1} \left(\langle \bar{\psi}\psi \rangle_c + \frac{\Delta \langle \bar{\psi}\psi \rangle}{2} \right) - \langle \bar{\psi}\psi \rangle^{-1} \left(\langle \bar{\psi}\psi \rangle_c - \frac{\Delta \langle \bar{\psi}\psi \rangle}{2} \right)$$

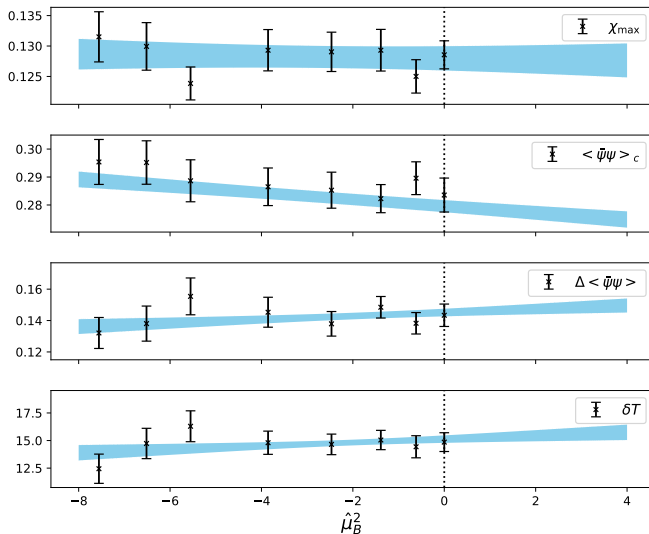
Let's visualize the approach

What is ΔT and δT ?

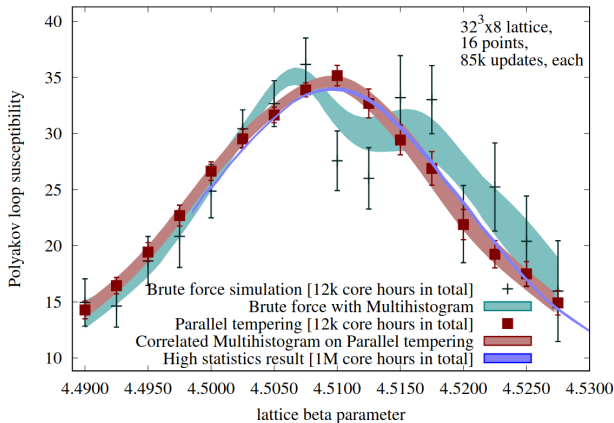


Continuum extrapolated observables as functions of $\hat{\mu}_B^2$

Borsanyi, R.K. et. al. PRL (2020)

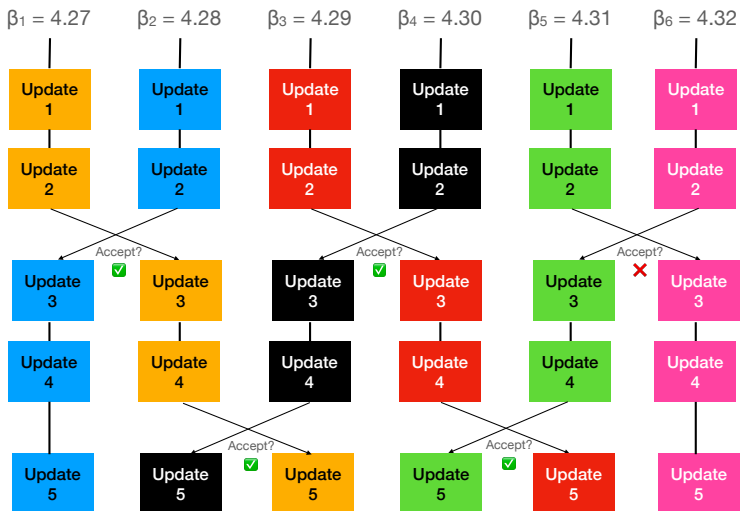


Quenched case: (Super-)critical slowing down



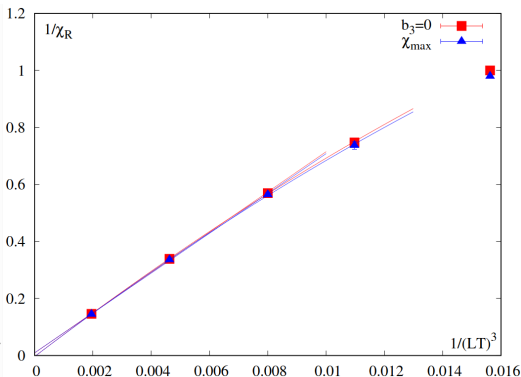
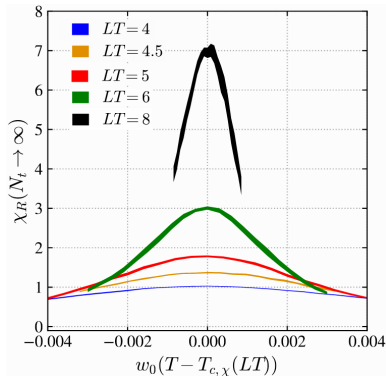
Tempering algorithms

- Multiple simulations at different β [Marinari 9205018] [Jo6 9810032]
- Swap configurations between subensembles $\Rightarrow$ reduced auto. corr.



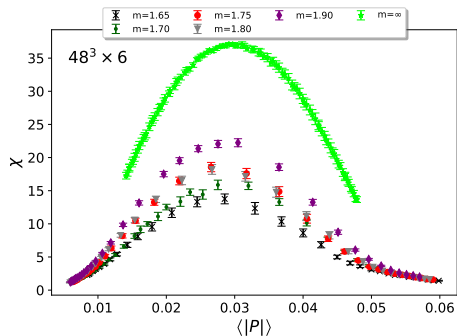
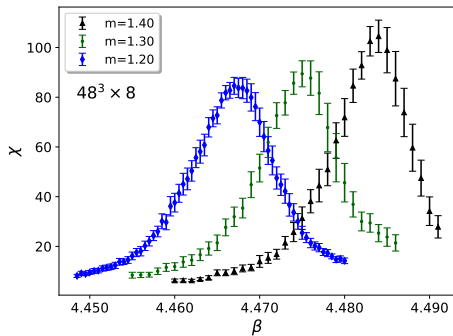
Polyakov loop and its susceptibility

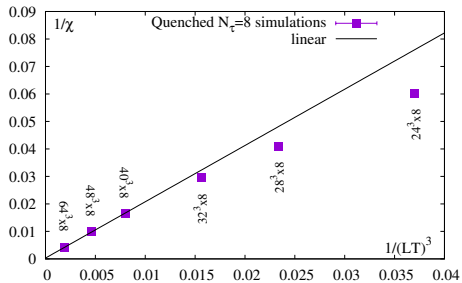
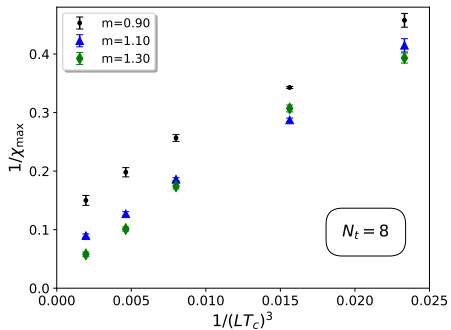
$$P = \frac{1}{N_s^3} \sum_{\vec{x}} P_{\vec{x}} = \frac{1}{N_s^3} \sum_{\vec{x}} \text{tr} \left[\prod_{\tau} U_4(\vec{x}, \tau) \right] \quad \chi = N_s^3 \left(\langle |P|^2 \rangle - \langle |P| \rangle^2 \right)$$

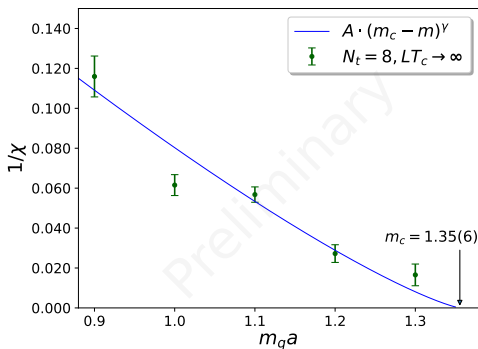
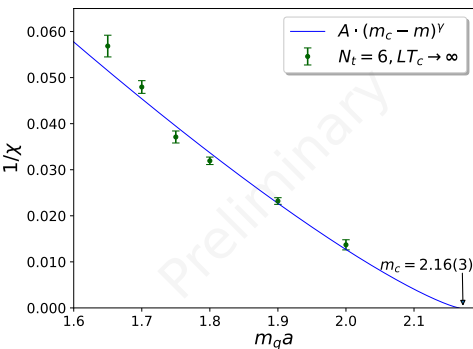


Borsanyi, R.K. et. al. PRD (2022)

Heavy but finite masses



Volume scaling for $N_t = 8$ 

Determination of the critical mass m_c 

Critical region

- $\chi_{\max}^{-1}$ follows a power law near the transition
- m_q represents symmetry breaking field

$$\chi_{\max}^{-1}(LT_c \rightarrow \infty) = A \cdot (m_c - m)^\gamma$$

Summary

Results of dynamical quark simulations

- β tempering algorithms to reduce auto-correlation
- $\chi(\langle |P| \rangle)$ allows precise peak determination
- Estimate of the critical mass: $N_t = 6$ $m_c = 2.16$; $N_t = 8$ $m_c = 1.35$

Outlook

- Systematic error analysis
- Continuum limit of the critical mass

Thank you for your attention!

Tempering algorithm

Background

- Tempering simulation is collection of sub-simulations
- Overall phase space as direct product of sub-ensemble phase spaces
- Equilibrium distribution of configs. in individual sub-ensembles overlap
- Independent Markov-processes in each sub-ensemble
- Swap configs. between sub-ensemble pairs

⇒ Auto-corr. is reduced

Tempering algorithm

Transition between sub-ensembles

- a is config. in sub-ensemble i , b config. in sub-ensemble j

$(a, b) \rightarrow (b, a) \Leftrightarrow$ swap accepted

$(a, b) \rightarrow (a, b) \Leftrightarrow$ swap rejected

Detailed balance condition

$$P_s(i, j)e^{-\mathcal{H}_i(a)}e^{-\mathcal{H}_j(b)} = P_s(j, i)e^{-\mathcal{H}_j(a)}e^{-\mathcal{H}_i(b)}$$

$$P_s(i, j) = \min(1, e^{\Delta\mathcal{H}})$$

$$\Delta\mathcal{H} = \{\mathcal{H}_j(a) + \mathcal{H}_i(b)\} - \{\mathcal{H}_i(a) + \mathcal{H}_j(b)\}$$

Cubic spline Interpolation

$$s_i(x) = a_i + b_i x + c_i x^2 + d_i x^3 \implies 4(N - 1) \text{ conditions}$$

Smoothness conditions

$$s_i(x_{i+1}) = s_{i+1}(x_{i+1})$$

$$s'_i(x_{i+1}) = s'_{i+1}(x_{i+1})$$

$$s''_i(x_{i+1}) = s''_{i+1}(x_{i+1})$$

$$\implies 3(N - 2) \text{ constraints}$$

Interpolation condition

$$S(x_j) = y_j$$

$$\implies N \text{ conditions}$$

Endpoints

$$s''_0(x_0) = 0 = s''_{N-2}(x_{N-1}) \implies 2 \text{ conditions}$$