

# Exploring the nucleon structure: Fragmentation Functions from e+e-annihilation experiment

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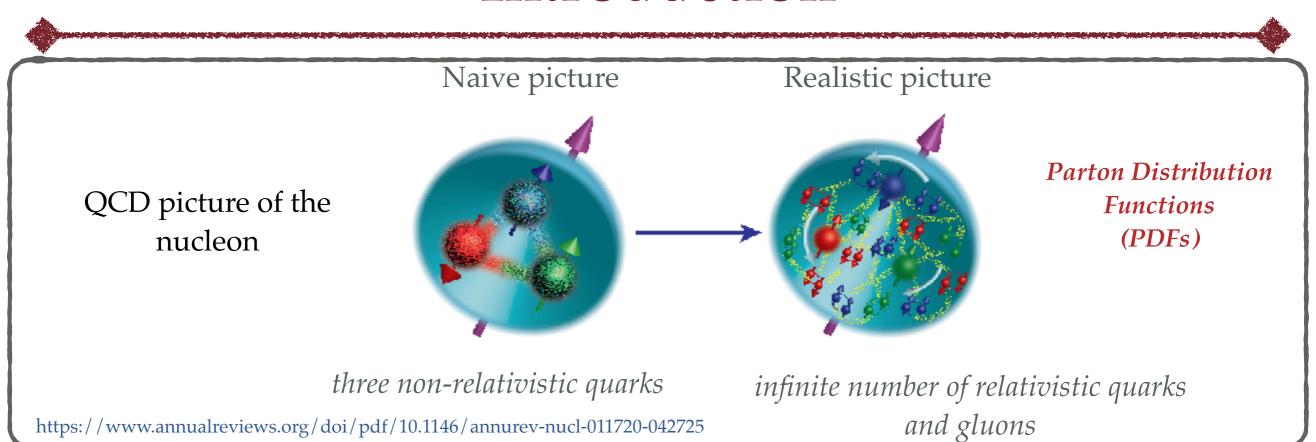




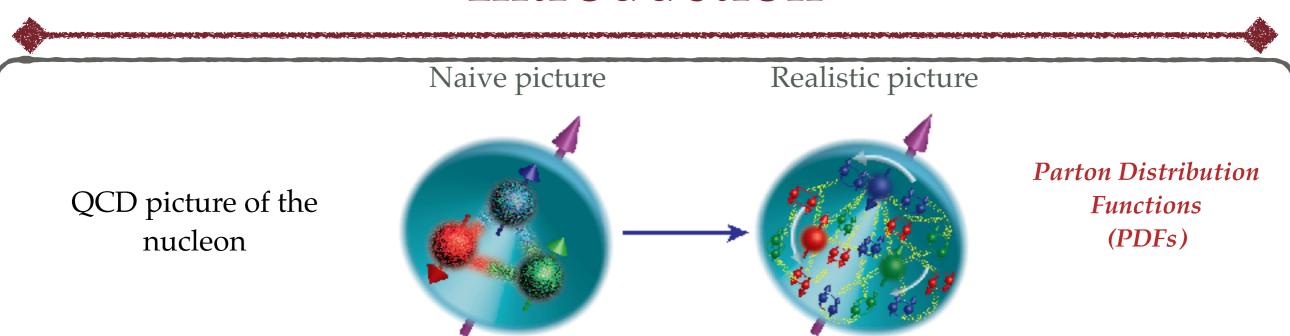
HOTEL & SPA MEDITERRANEAN VILLAGE, PARALIA (PIERIA, GREECE)

THURSDAY, MAY 26, 2022

## Introduction



#### Introduction



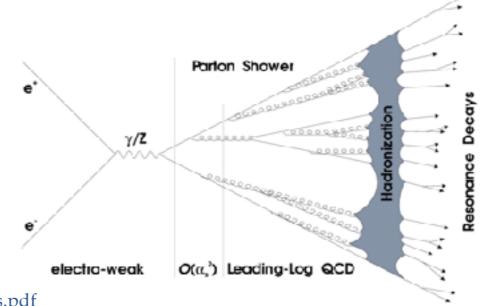
three non-relativistic quarks

https://www.annual reviews.org/doi/pdf/10.1146/annurev-nucl-011720-042725

infinite number of relativistic quarks and gluons

#### Hadron formation

- How many particles and how many jets created?
- What fraction of the initial parton momenta do they carry?

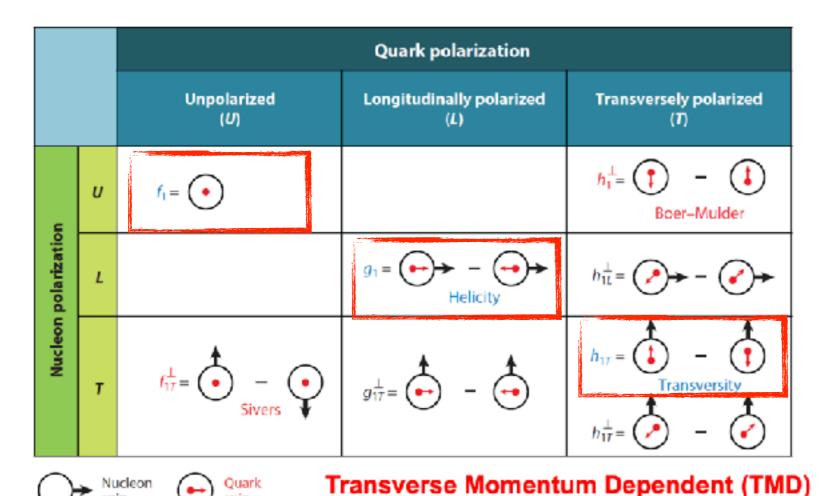


Fragmentation
Functions
(FFs)

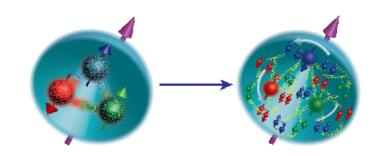
https://pdg.lbl.gov/2019/reviews/rpp2019-rev-frag-functions.pdf

#### Universal and non-perturbative objects

## Parton Distribution Functions



TMD independent

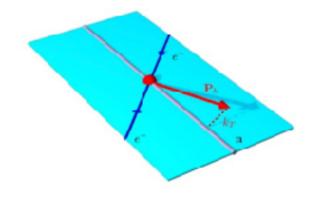


f<sub>1</sub>, g<sub>1</sub>, and h<sub>1</sub>: three leading-twist PDFs

- probability interpretation;
- complete description of the nucleon structure at leading-twist level;
- provide 2-dim (collinear) picture of the nucleon structure
- TMD: 3-D picture

#### **h1: Transversity PDF**

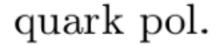
- Describe the distribution of quark's transverse spin in a transversely polarized nucleon
- Chiral-odd



# Fragmentation Functions

**FFs:** formation of colourless hadrons starting from a coloured partonic initial state;

• probability that a parton i fragments into an hadron h carrying away a fraction z of the parton's momentum



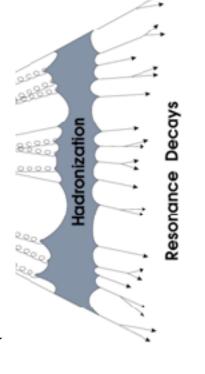
hadron pol.		U	L	T
	U	$D_1$		$H_1^\perp$
	L		$G_{1L}$	$H_{1L}^{\perp}$
	$\mathbf{T}$	$D_{1T}^{\perp}$	$G_{1T}$	$H_1, H_{1T}^{\perp}$

#### H<sub>1</sub><sup>⊥</sup>: Collins FF

· Chiral-odd function

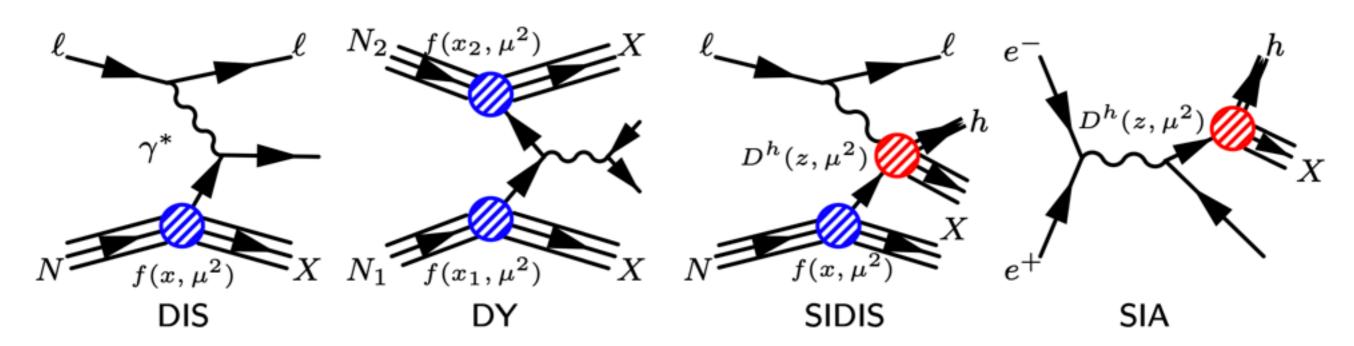
**Transverse Momentum Dependent (TMD)** FFs ⇒ to study the spin-dependent observables

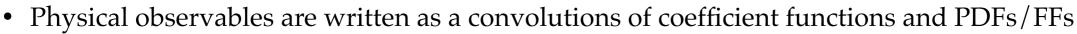
- tools to investigate the 3Dstructure of nucleons
- when only spinless hadrons  $(\pi, K)$  are considered, we have only  $D_1$  and  $H_1^{\perp}$
- TMDs evolution



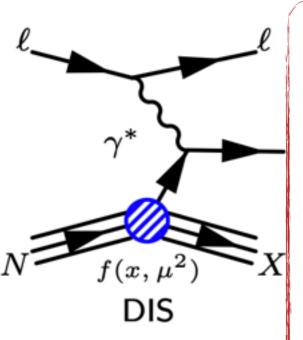
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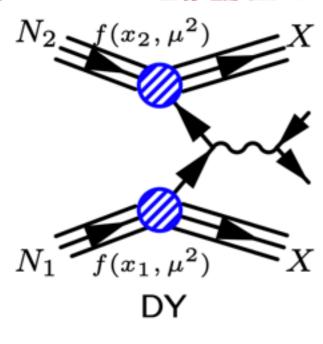
- Physical observables are written as a convolutions of coefficient functions and PDFs/FFs
- *QCD factorization theorem*: separate the cross sections into a perturbatively calculable scattering contributions and non-perturbative one encoded in PDFs/FFs



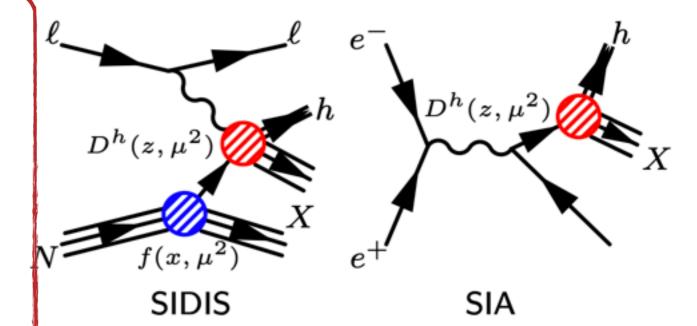


• *QCD factorization theorem*: separate the cross sections into a perturbatively calculable scattering contributions and non-perturbative one encoded in PDFs/FFs



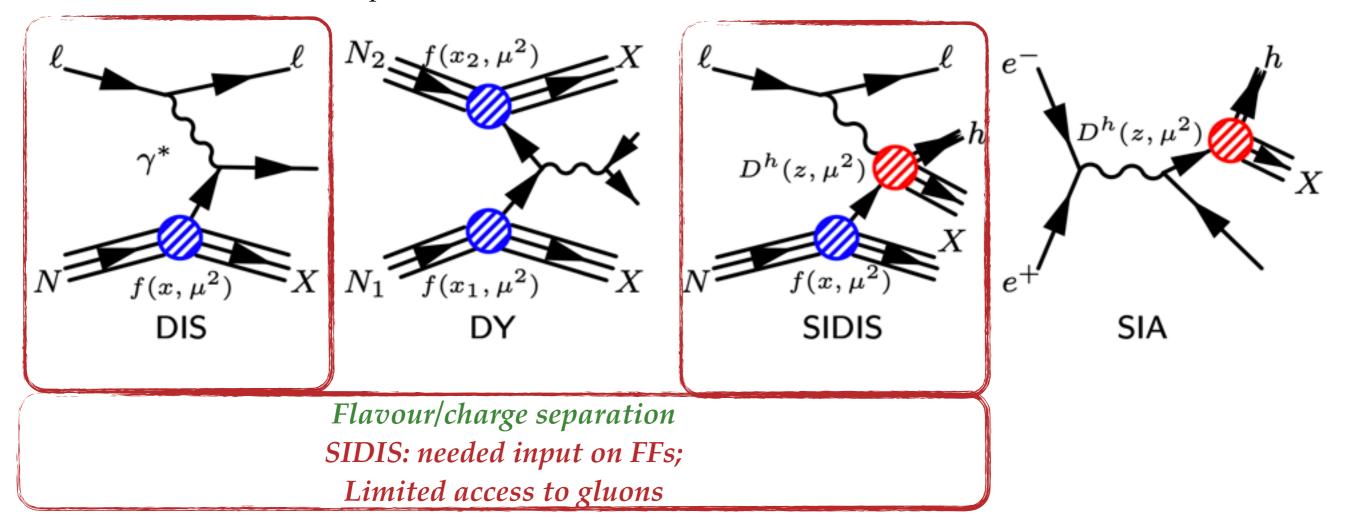


Direct access to PDFs;
Sensitive to gluon FFs;
large z behaviour;
Large theoretical
uncertainties;
no sensitive to heavy
quarks



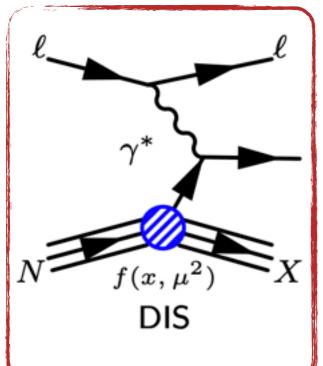


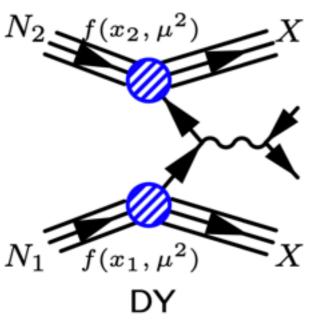
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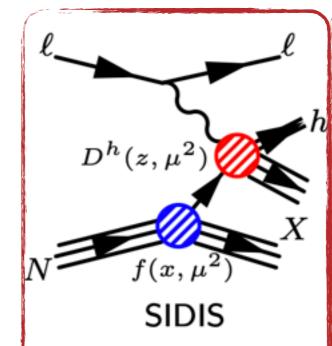


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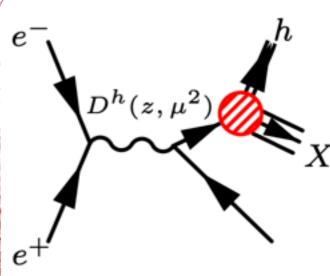
- Physical observables are written as a convolutions of coefficient functions and PDFs/FFs
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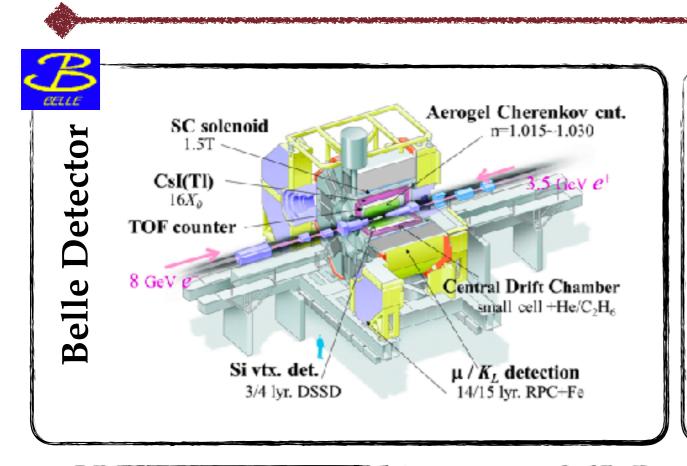
Flavour/charge separation SIDIS: needed input on FFs; Limited access to gluons

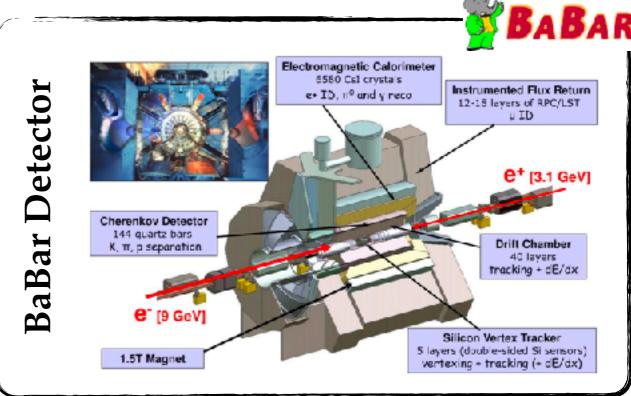


#### SIA

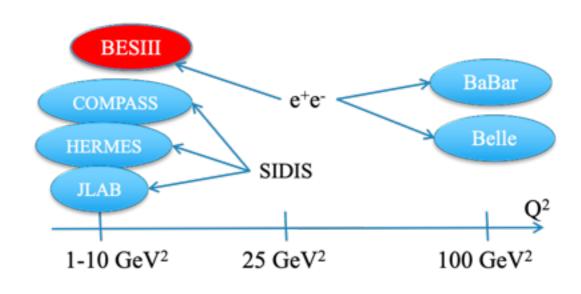
Clean process, high
statistic;
Sensitivity to heavy
quarks;
Limited in flavour
separation;
Limited access to gluon
FFs

# Belle, BaBar, BESIII detectors

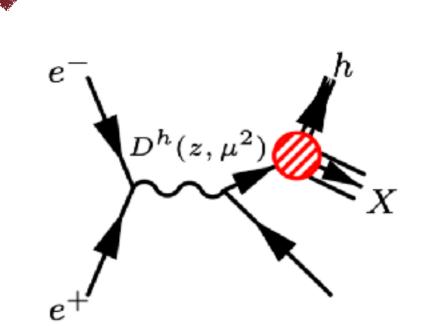


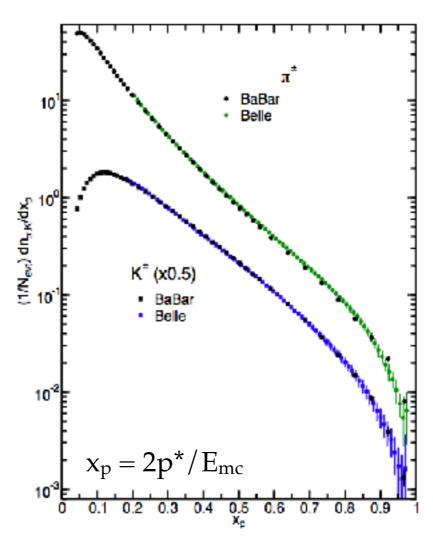






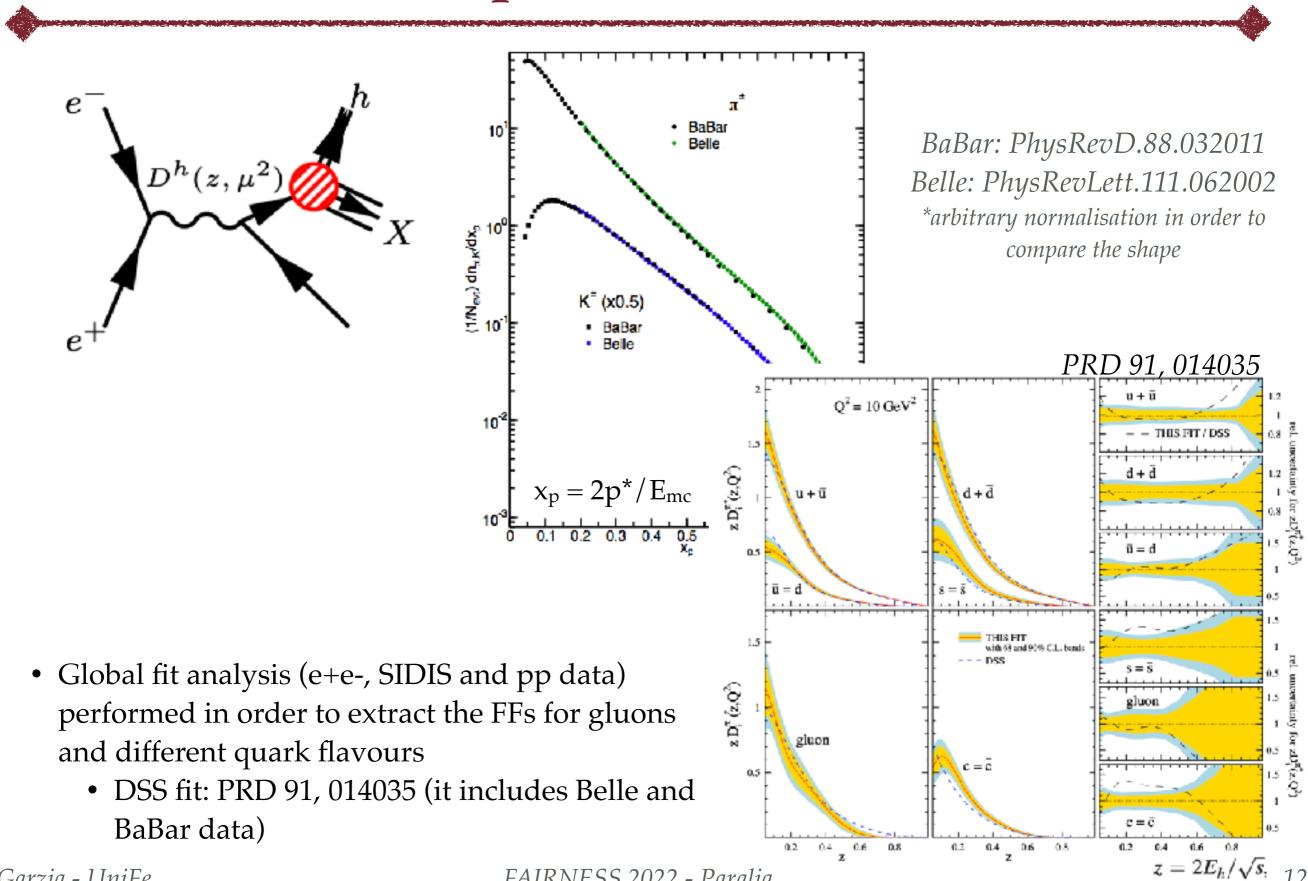
# Unpolarised FFs





BaBar: <u>PhysRevD.88.032011</u>
Belle: PhysRevLett.111.062002
\*arbitrary normalisation in order to compare the shape

# Unpolarised FFs



# Inclusive light hadrons productions

- Data set used:  $0.9 \text{ fb}^{-1} @ \Upsilon(4S)$  and  $3.6 \text{ fb}^{-1}$  at 10.54 GeV
- measured both *prompt* and *conventional* hadron cross section
  - *prompt*: primary hadrons or products of a decay chain where all particles have a lifetime shorter than  $10^{-11}$  s
  - *conventional*: includes weak decay products of KS and strange baryons
- Scaled momentum distribution:  $x_p = 2p^*/E_{mc}$  (0.2<p<5.27 GeV/c)

#### **JETSET model**:

- represent the color field between the parton by a "string", and according to an iterative algorithm breaks the string into several pieces, each corresponding to a primary hadron
- large number of free parameters (models many hadron species)

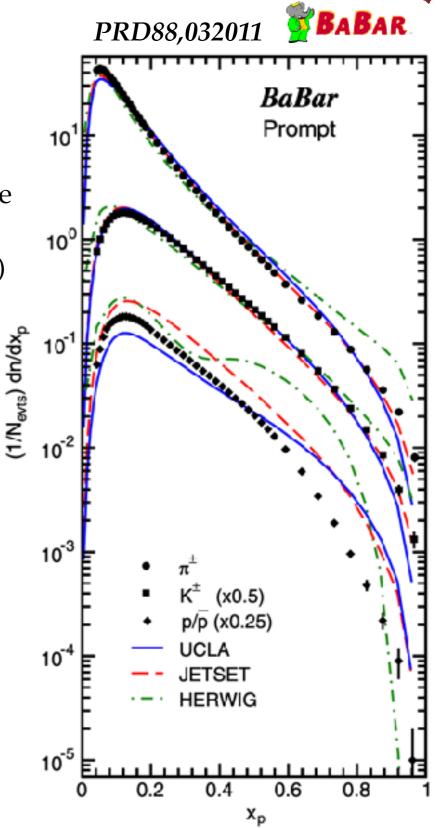
#### **HERWIG model**:

- splits the gluons produced into  $q\overline{\overline{q}}$  pairs, combines these quark and antiquark locally to form colorless "clusters", and decay these "clusters" into primary hadrons
- few free parameters

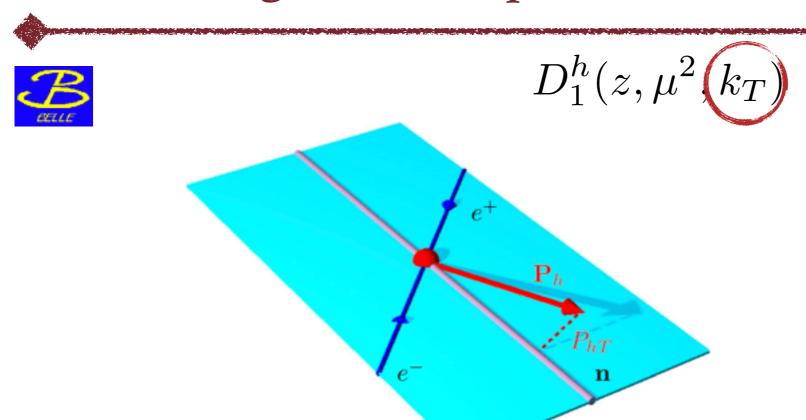
#### **UCLA model**:

- generates whole events according to weights derived from phase space and Clebsch-Gordan coefficients
- few free parameters

Good qualitative description of the bulk spectra (no the details)



#### TMD single-hadron production cross sections @ Belle

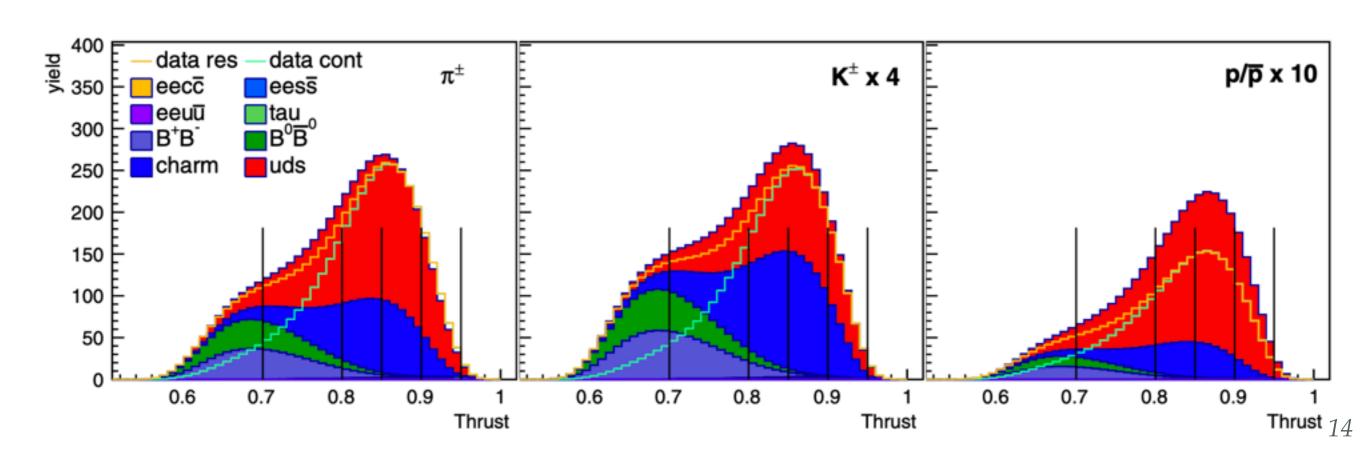


PRD **99**, 112006

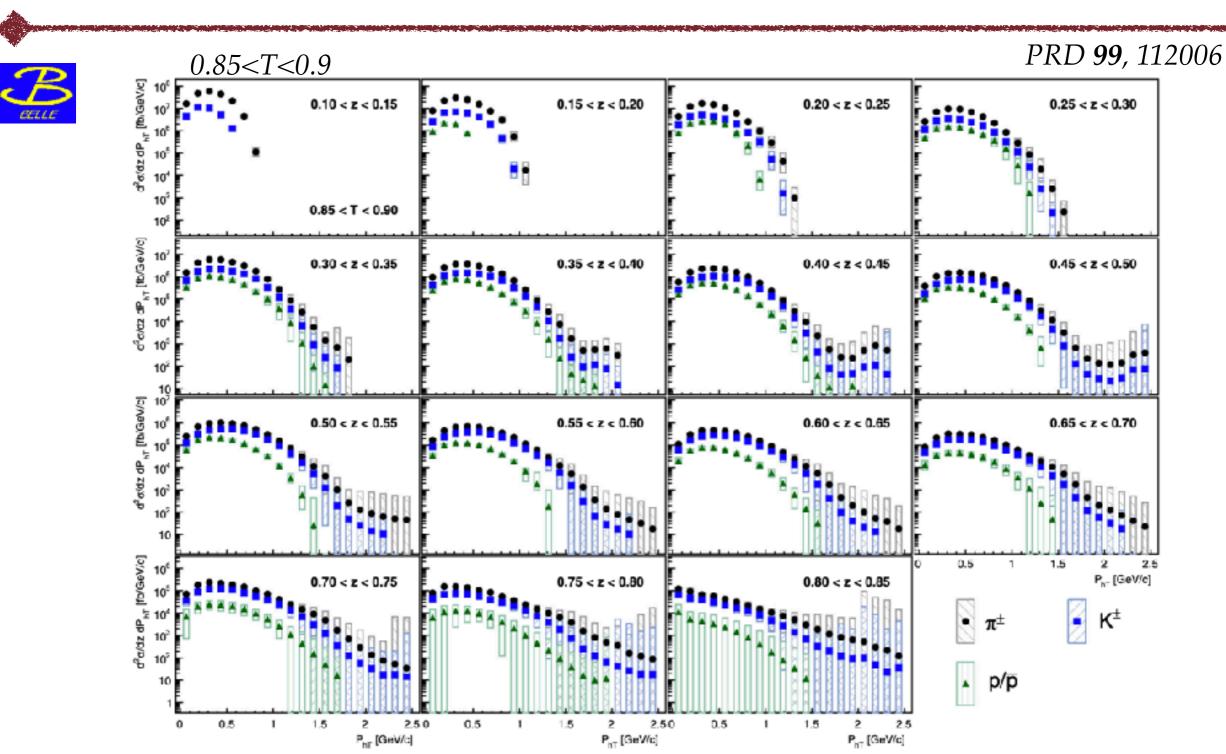
Definition of thrust event shape variable:

$$T \stackrel{\text{max}}{=} \frac{\sum_{h} |\mathbf{P}_{h}^{\text{CMS}} \cdot \hat{\mathbf{n}}|}{\sum_{h} |\mathbf{P}_{h}^{\text{CMS}}|}.$$

$$z = 2E_h/\sqrt{s}$$
,  $\mathbf{p}_T = -z\mathbf{k}_T$ 

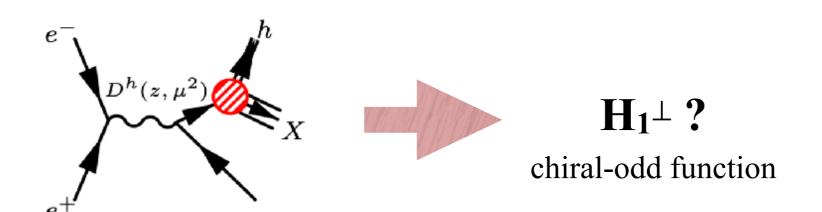


#### TMD single-hadron production cross sections @ Belle

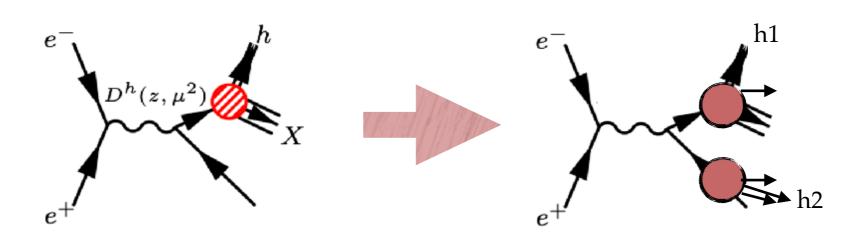


First direct measurements for pions, kaons and protons: better understanding of transverse spin phenomena and important role in the future electron-ion collider to pin down the transverse-momentum structure of nucleons and its transverse spin

## Polarised FF: the Collins effect



#### Polarised FF: the Collins effect



q↑
$$\rightarrow$$
hX:  $D_1^{q\uparrow}(z, \mathbf{P}_\perp; s_q) = D_1^q(z, P_\perp) + \frac{P_\perp}{zM_h} H_1^{\perp q}(z, P_\perp) \mathbf{s}_q \cdot (\mathbf{k}_q \times \mathbf{P}_\perp)$ 
Unpolarized FF Collins FF [NPB 396, 161 (1993)]:

related to the probability that a transversely polarized quark (q<sup>↑</sup>) fragments into a spinless hadron

- Evolution of TMD objects
- Global analysis (PRD 78,032011 (2007); PRD 87,094019 (2013), PRD 91,014034 (2015)):
  - combines Semi Inclusive Deep Inelastic Scattering (SIDIS) and e+e-data
  - extraction of  $H^{\perp_1}$  and transversity parton distributions  $h_1$  for the "u" and "d" quarks

## Collins effect in e+e- annihilation

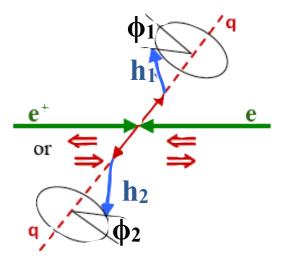
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In e<sup>+</sup>e<sup>-</sup> $\rightarrow$ q $\overline{q}$ , spins unknown, but  $s_q \parallel s_{\overline{q}}$ 

- exploit this correlation by using hadrons in opposite jets
- define **favored**  $(u \rightarrow \pi^+, d \rightarrow \pi^-)$  and **disfavored**  $(d \rightarrow \pi^+, u \rightarrow \pi^-, s(\overline{s}) \rightarrow \pi^{\pm})$  FFs

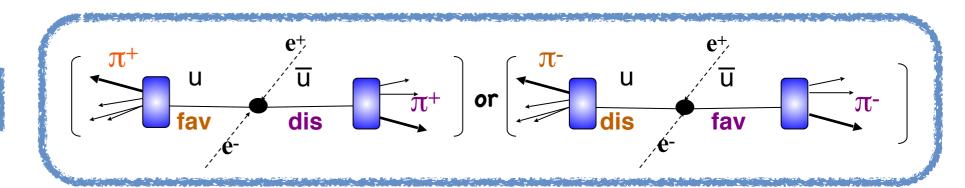
$$e^+e^- \rightarrow q\bar{q} \rightarrow h_1h_2X \quad (q=u,d,s) \Rightarrow \sigma \propto \cos(\phi_1 + \phi_2)H_1^{\perp (h_1)} \times H_1^{\perp (h_2)}$$



**Azimuthal modulation wrt the quark spin direction: Collins effect (or Collins asymmetry)** 

Example: Like  $\pi\pi$  pairs (L)

Collins asymmetry for  $\pi\pi$ 



## Collins effect in e+e- annihilation

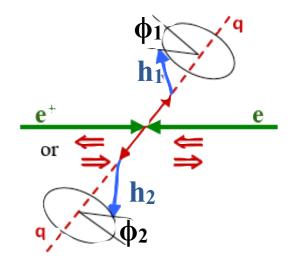
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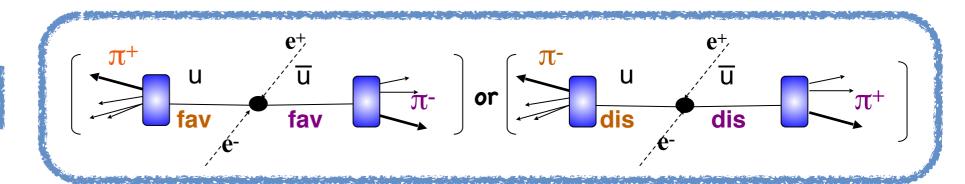
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**Azimuthal modulation wrt the quark spin direction: Collins effect (or Collins asymmetry)** 

Example: Unlike  $\pi\pi$  pairs (U)

Collins asymmetry for  $\pi\pi$ 



#### Collins effect in e<sup>+</sup>e<sup>-</sup> annihilation

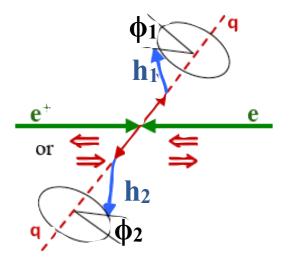




In e<sup>+</sup>e<sup>-</sup> $\rightarrow$ q $\overline{q}$ , spins unknown, but  $s_q \parallel s_{\overline{q}}$ 

- · exploit this correlation by using hadrons in opposite jets
- define **favored** ( $u \rightarrow \pi^+$ ,  $d \rightarrow \pi^-$ ) and **disfavored** ( $d \rightarrow \pi^+$ ,  $u \rightarrow \pi^-$ ,  $s(\overline{s}) \rightarrow \pi^{\pm}$ ) FFs

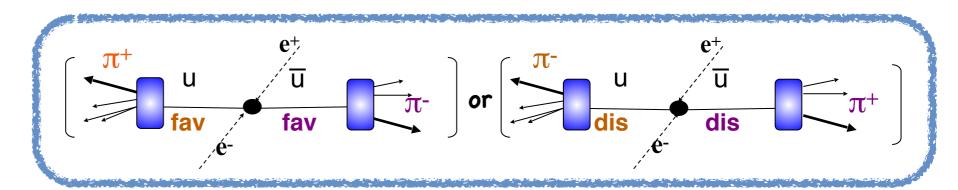
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**Azimuthal modulation wrt the quark spin direction: Collins effect (or Collins asymmetry)** 

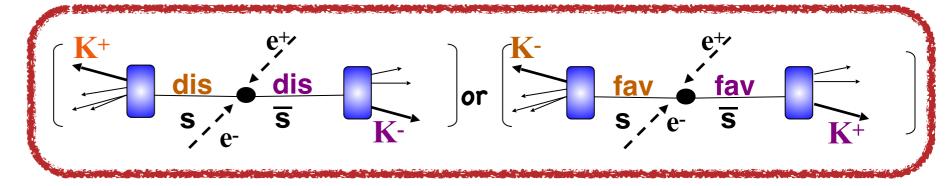
Example: Unlike  $\pi\pi$  pairs (U)

Collins asymmetry for  $\pi\pi$ 



Example: Unlike KK pairs (U)

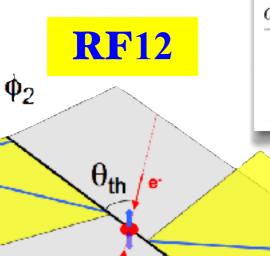
Collins asymmetry for KK:



#### Reference Frames



See D. Boer, NPB 806, 23 (2009)



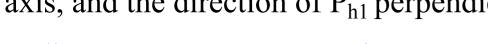
$$\begin{split} \frac{d\sigma(e^+e^- \to h_1 h_2 X)}{d\Omega dz_1 dz_2 d\phi_1 d\phi_2} &= \sum_{q,\bar{q}} \frac{3\alpha^2}{Q^2} \frac{e_q^2}{4} z_1^2 z_2^2 \left[ (1 + \cos^2\theta) D_1^{q,(0)}(z_1) \bar{D}_1^{q,(0)}(z_2) + \right. \\ &+ \left. \sin^2(\theta) \cos(\phi_1 + \phi_2) H_1^{\perp,(1),q}(z_1) \bar{H}_1^{\perp,(1),q}(z_2) \right] \end{split}$$

All quantities in e+e- center of mass

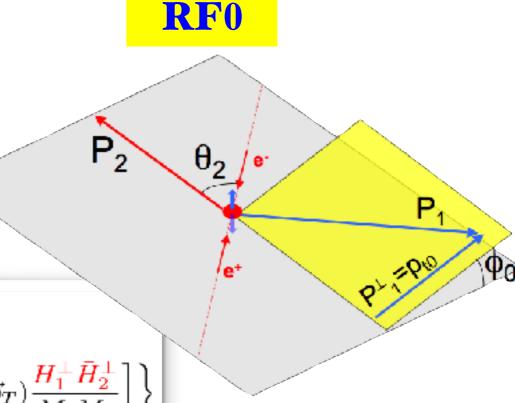
 $\theta$ : angle between the e<sup>+</sup>e<sup>-</sup> axis and the thrust axis;  $\phi_{1,2}$ : azimuthal angles between  $P_{h1(h2)}$  and the scattering plane

 $\theta_2$ : angle between the e<sup>+</sup>e<sup>-</sup> axis and  $P_{h2}$ ;  $\phi_0$ : angle between the plane spanned by  $P_{h2}$  and the e<sup>+</sup>e<sup>-</sup> axis, and the direction of  $P_{h1}$  perpendicular to  $P_{h2}$ .

n



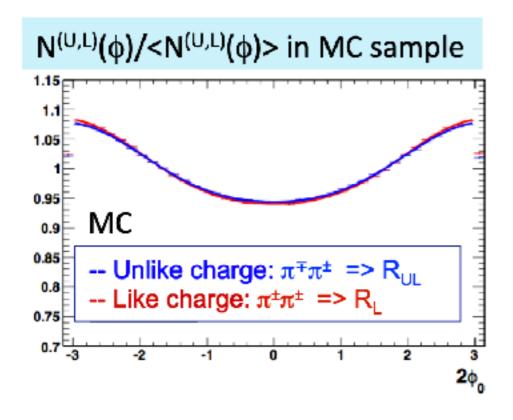
All quantities in e+e<sup>-</sup> center of mass 
$$\begin{vmatrix} d\sigma(e^{+}e^{-} \rightarrow h_{1}h_{2}X) \\ d\Omega dz_{1}dz_{2}d^{2}\vec{q}_{T} \end{vmatrix} = \frac{3\alpha^{2}}{Q^{2}}z_{1}^{2}z_{2}^{2}\left\{A(y)\mathcal{F}[D_{1}\bar{D_{2}}] + B(y)\cos(2\phi_{0})\mathcal{F}\left[(2\hat{h}\cdot\vec{k}_{T}\hat{h}\cdot\vec{p}_{T} - \vec{k}_{T}\cdot\vec{p}_{T})\frac{H_{1}^{\perp}\bar{H}_{2}^{\perp}}{M_{1}M_{2}}\right]\right\}$$



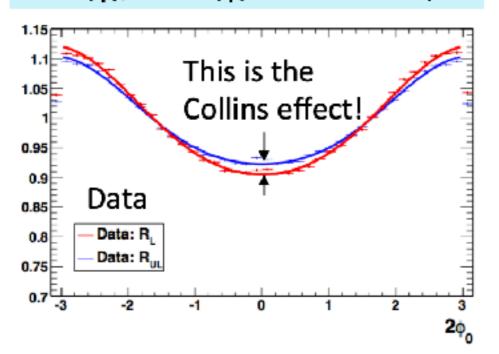
P12

# Asymmetry Measurements from Data





#### $N^{(U,L)}(\phi)/\langle N^{(U,L)}(\phi) \rangle$ in data sample



• Collins asymmetries extracted from fit to the normalised azimuthal distribution

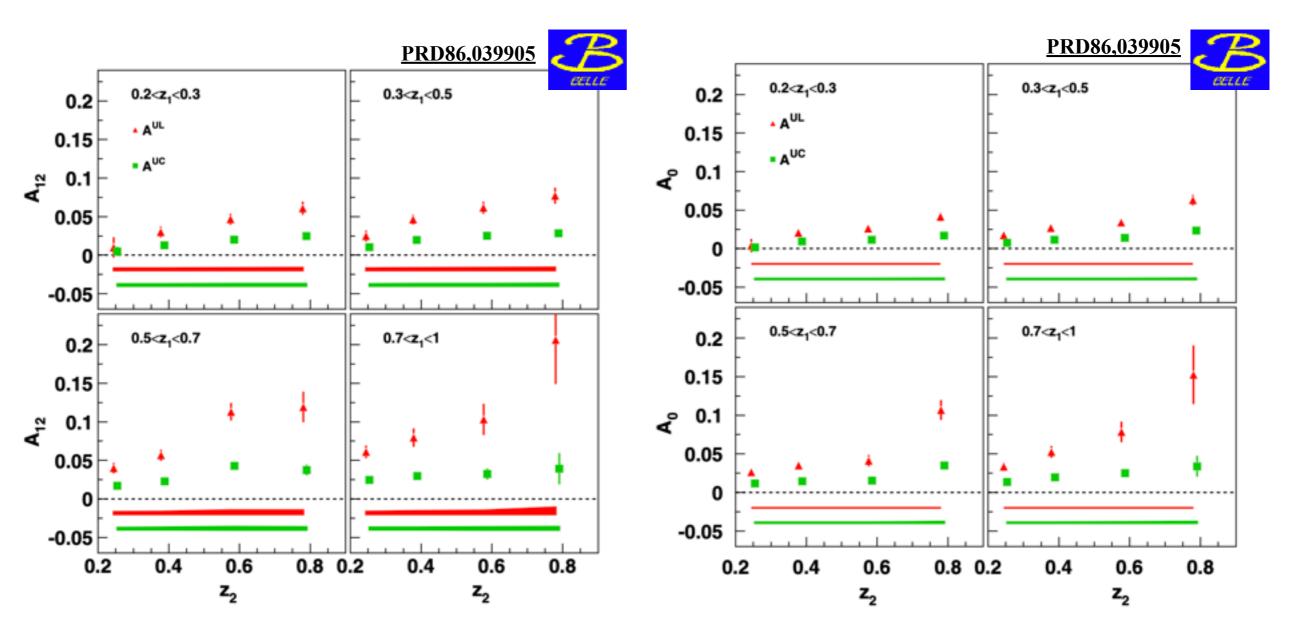
$$R_{\alpha} = \frac{N(\phi_{\alpha})}{\langle N_{\alpha} \rangle} = a + b \cdot \cos(\phi_{\alpha})$$

- MC generator does not include polarised FFs as the Collins FF
  - modulation observed in MC sample produced by detector acceptance
  - correction of these effects would bring too large systematic uncertainties
- Ratio of U and L distributions
  - Collins effect is not sensitive to the electric charge
  - U and L sigli different in data due to the different contribution of favoured and disfavoured FFs

## Collins FFs Results

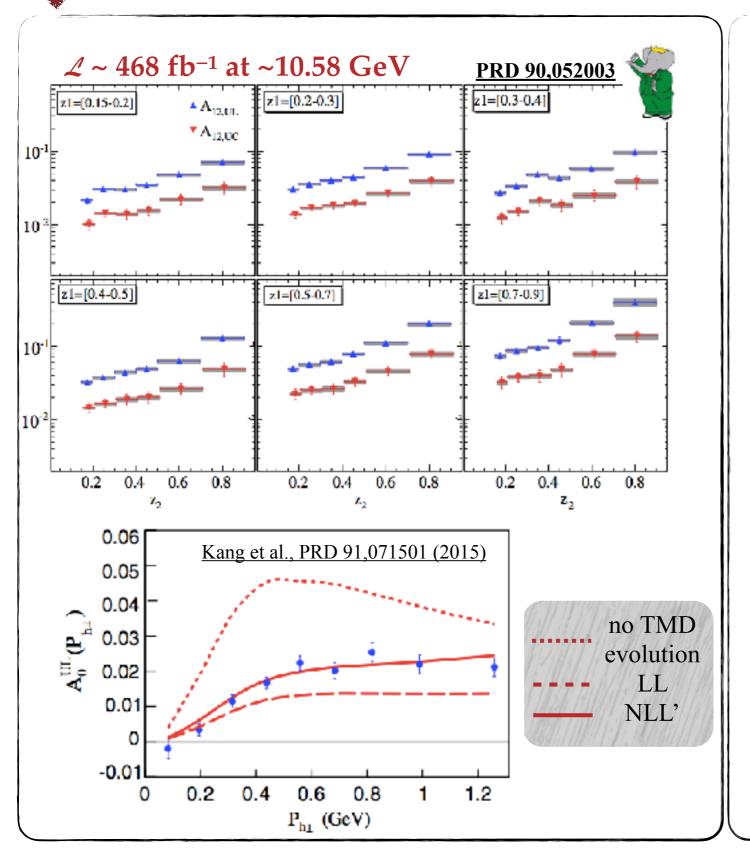


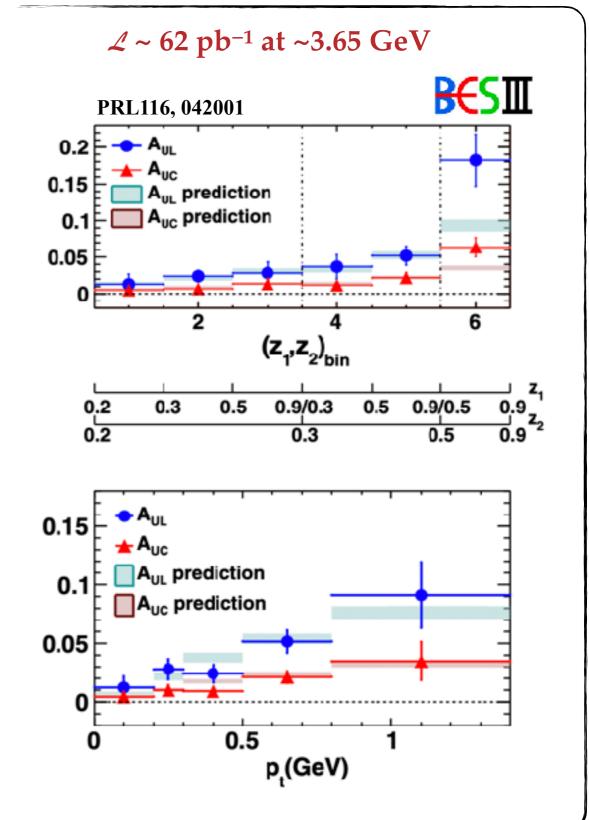
#### First measurement of Collins effects in e+e- annihilation (PRL 96,232002) !!!



- $\mathcal{L} \sim 547 \text{ fb}^{-1} \text{ at } \sim 10.58 \text{ GeV}$
- Non-zero asymmetries increasing with z

## Collins FFs Results

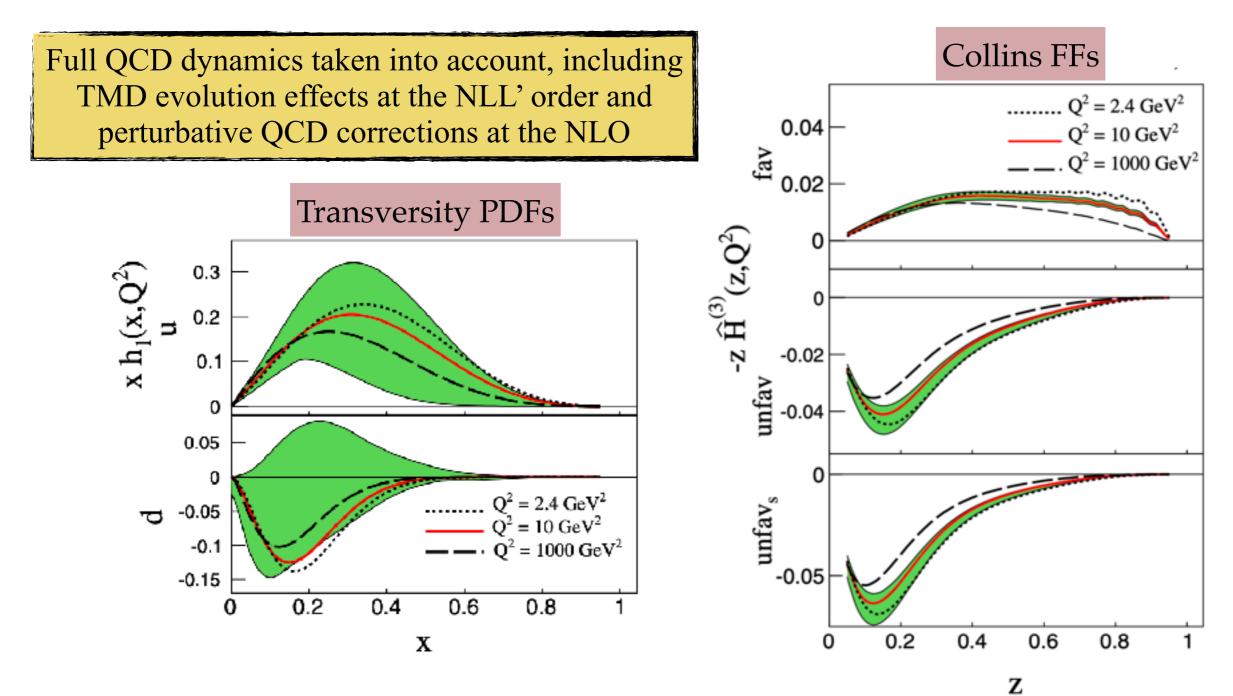




# Extraction of h<sub>1</sub> and Collins FF

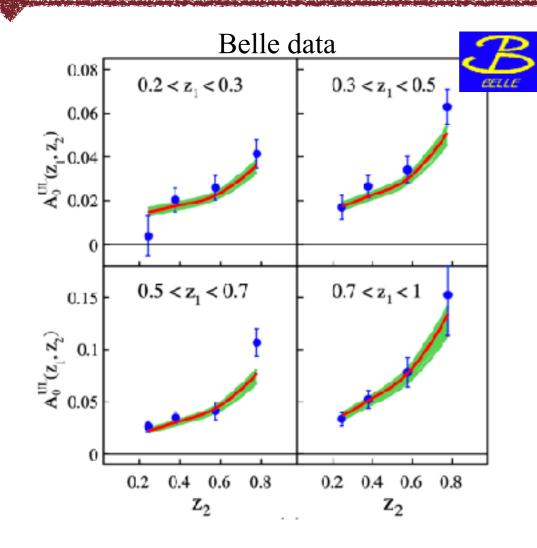
PRD**93**,014009

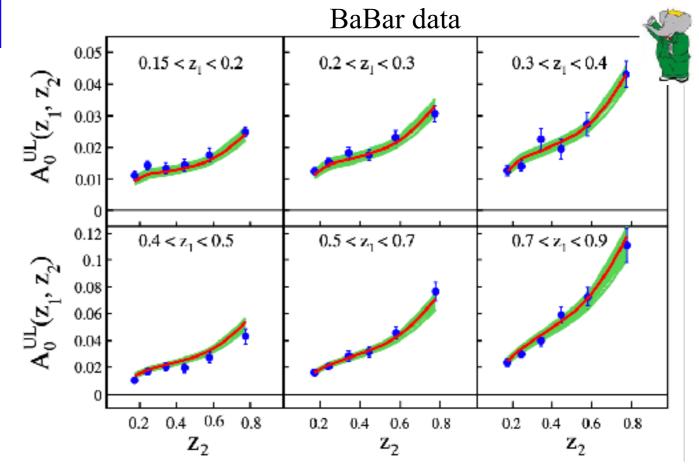
- First extraction: Anselmino et al. PRD75,054032
- Global analysis with SIDIS (HERMES (PRL103,152002), COMPASS (PLB744,250; PLB673,127)) + BaBar RF0 data (PRD90,052003) and Belle RF0 (PRD78,032011) data



# Extraction of Collins FF: data description

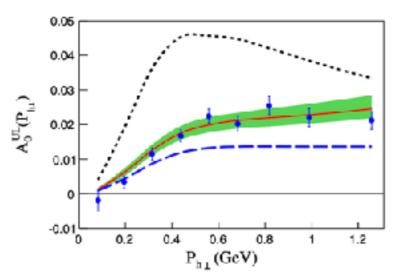
PRD**93**,014009

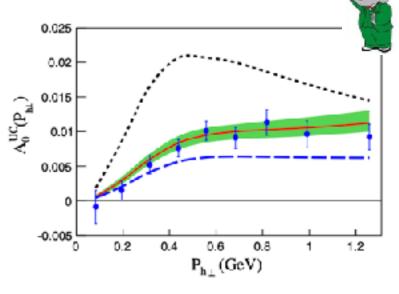




NLL' calculationsLL calculationno TMD evolution

e<sup>+</sup>e<sup>-</sup> data are very sensitive to the inclusion of higher order corrections



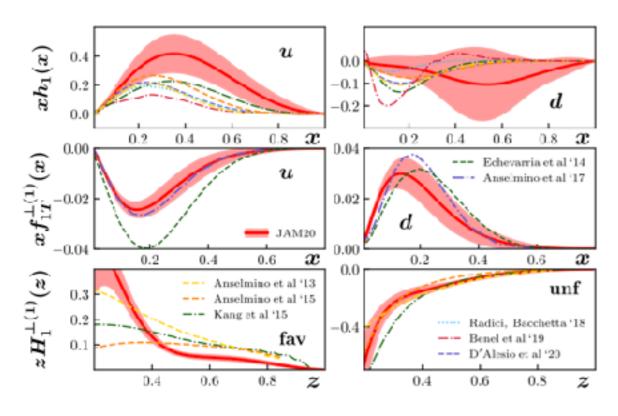


# First Global Analysis

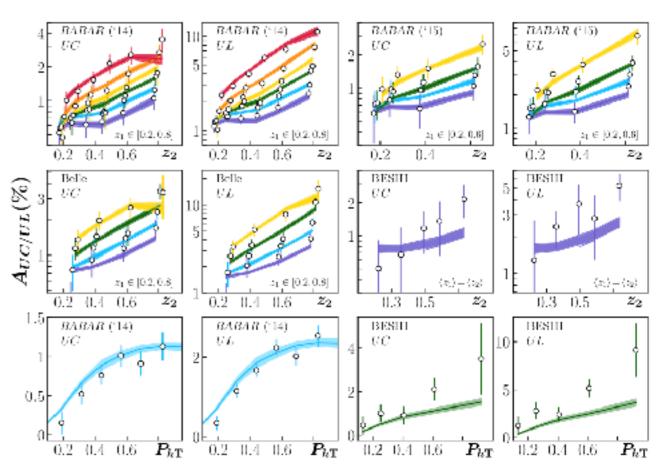
PRD102,054002

First simultaneous QCD global analysis with SIDIS, e+e- annihilation, DY and proton-proton collisions

- Test of universality
- Indication that transverse-spin asymmetries in high-energy collisions have a common origin
- · Extracted quark tensor charges are in excellent agreement with lattice QCD



<u>Anselmino'13, Anselmino'15, Anselmino'17, Benel'19, D'Alessio'20, Echevarria'14, Kang'15, Radici-Bacchetta'18</u>

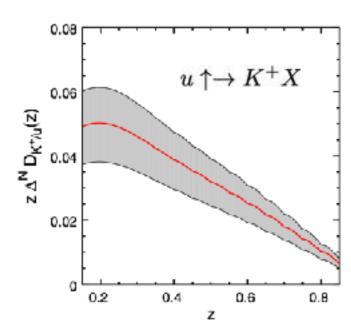


Comparison between theory and data (<u>BaBar</u>, <u>Belle</u>, <u>BESIII</u>)

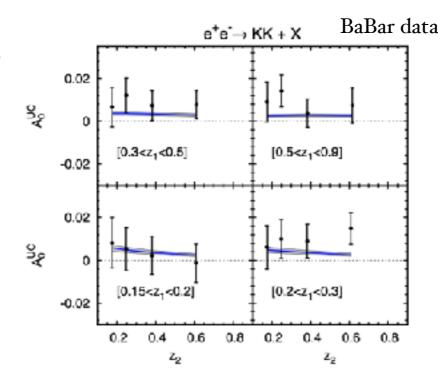
## Extraction of Collins FF for Kaons

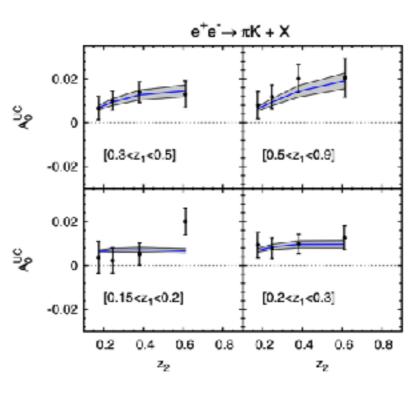
PRD **93**,034025

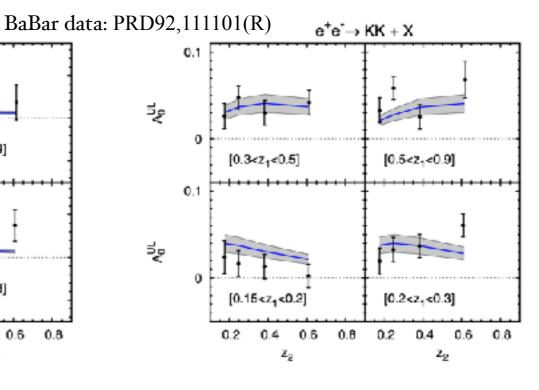
- Pions FF from previous analysis
- Similar parametrisation for kaon favoured and disfavoured FFs as used for pions

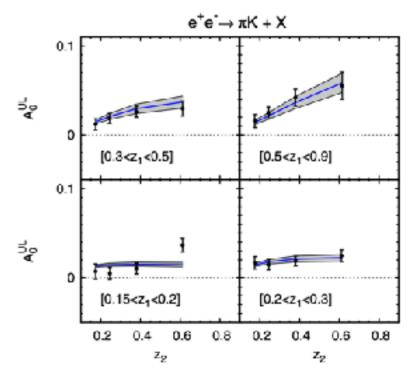


- Positive favored Collins function and a disfavoured contribution compatible with zero
- Large uncertainty
- No definitive conclusions can be extracted for strange Collins FF nor on the disfavoured ones



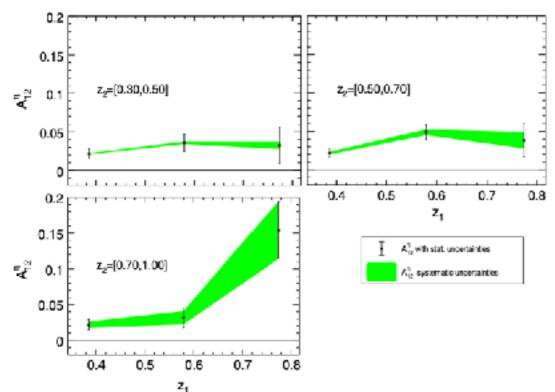


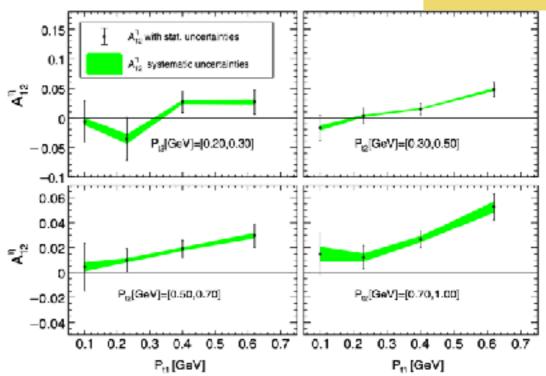




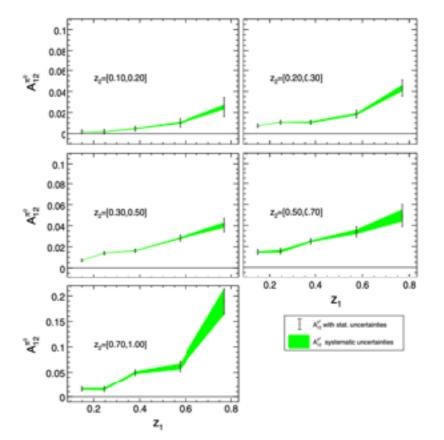
#### Azimuthal asymmetries of back-to-back $\pi^{\pm}\eta/\pi^{\pm}\pi^{0}$ from Belle







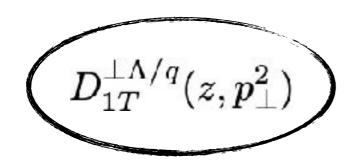
- RF12 (thrust reference frame) only
- $\pi^0$ :
  - Significant asymmetries that rise with z
  - Continuous rise with Pt1 consistent with a linear behaviour
- η:
  - Larger uncertainties
  - The rise with z is much less pronounced
  - Rise of the asymmetry with increasing pt
  - Charm contribution  $\sim 20-30\%$  larger than neutral pions sample
  - for those bins with similar charm contributions  $\eta$  and  $\pi^0$  asymmetries are fully consistent



# $\Lambda/\overline{\Lambda}$ hyperon polarization at Belle

PRL **122**,042001

•  $e^+e^- \rightarrow \Lambda(\overline{\Lambda})X$  and  $e^+e^- \rightarrow \Lambda(\overline{\Lambda})h^\pm X$  using 800.4 fb<sup>-1</sup> collected by Belle at or near  $\sqrt{s}=10.58$  GeV



==> production of transversely polarized  $\Lambda$  hyperons from unpolarised quark

==> chiral-even function (the sign of this function can be easily determined)

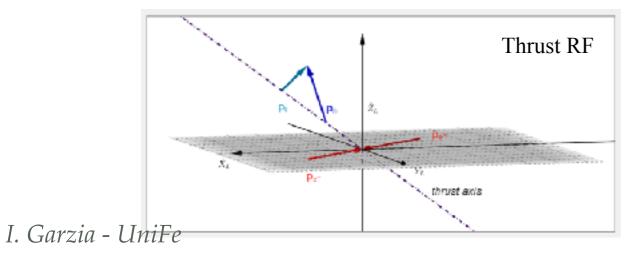
Distribution of protons from  $\Lambda$  decays with a transverse polarisation P:

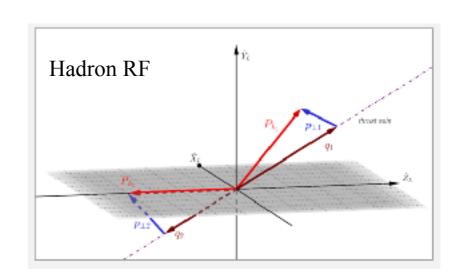
$$\frac{1}{N}\frac{dN}{d\cos\theta} = 1 + \alpha P\cos\theta,$$

$$\begin{split} P_{\Lambda}(z,p_{\perp}) &= \frac{\sum_{q} e_{q}^{2} \frac{4z p_{\perp} \sqrt{s}}{z^{2} s + 4 p_{\perp}^{2}} \Delta^{N} D_{\Lambda^{\uparrow}/q}(z,p_{\perp})}{\sum_{q} e_{q}^{2} 2\pi D_{\Lambda/q}(z,p_{\perp})} \\ &\times \frac{\int d(\cos\theta) \sin(2\theta)}{\int d(\cos\theta) (1 + \cos^{2}\theta)} \end{split} \quad \begin{array}{c} \text{PRD 100,014029} \end{split}$$

 $\theta$ : angle between the axis  $\stackrel{\wedge}{n} \propto \stackrel{\wedge}{T} \times \stackrel{\wedge}{p}_{\Lambda}$  and the proton momentum in the  $\Lambda$  rest frame;  $\alpha$  is the world average value of the parity-violating decay asymmetry for the  $\Lambda$  (from PDG)

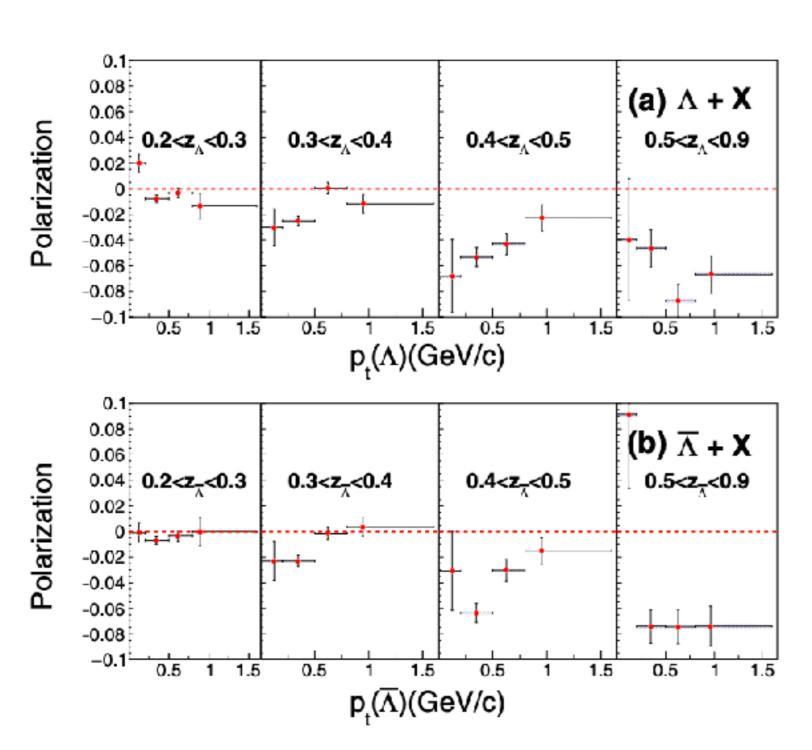
==> Significant transverse polarization is observed





# $\Lambda/\overline{\Lambda}$ hyperon polarization at Belle: results

PRL **122**,042001

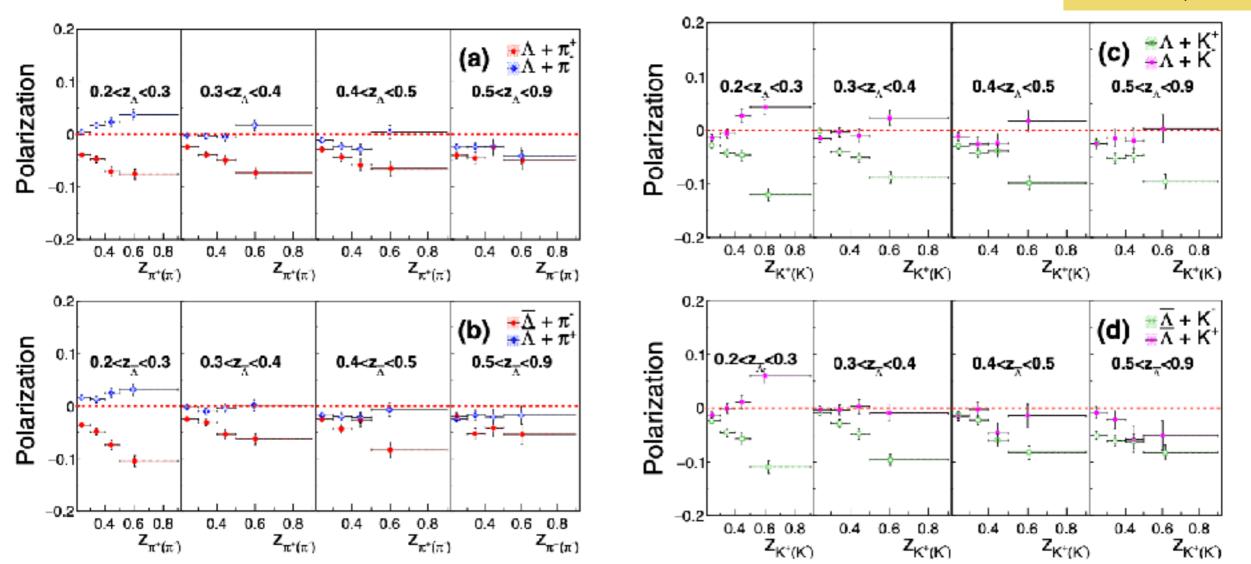


#### Complex pt dependence

- $z_{\Lambda}$ >0.5: increasing asymmetry with pt; opposite behaviour for intermediate  $z_{\Lambda}$ 
  - different quark flavour contributions in different kinematic regions (s quark contribution is dominant in the highest z<sub>Λ</sub> bin; u quark and charm contribution in the intermediate regions)
  - large charm contribution in intermediate z-bins

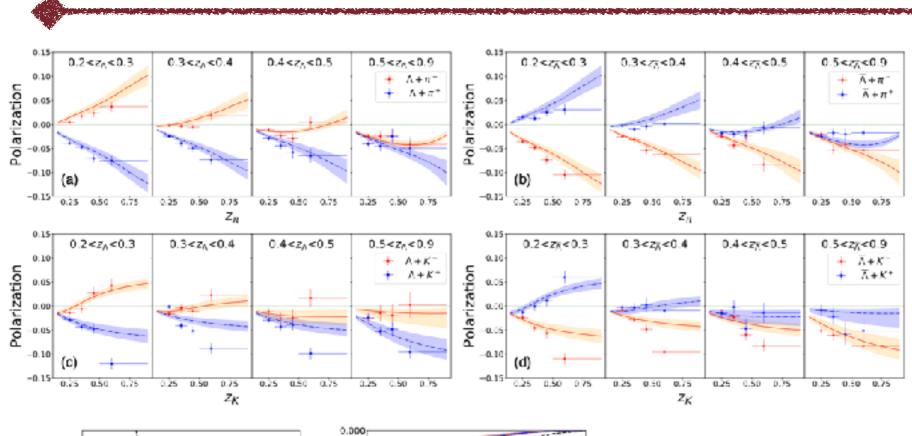
# $\Lambda/\overline{\Lambda}$ hyperon polarization at Belle: results

PRL **122**,042001



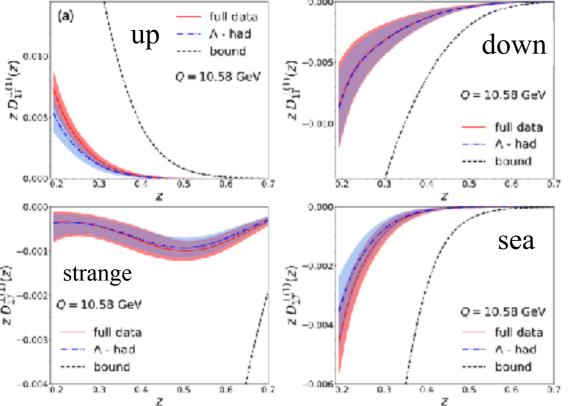
- low  $z_{\Lambda}$ : opposite site and increasing magnitude
- Sensitivity to the flavour dependence by selecting a light hadron in the opposite hemisphere:
  - $\Lambda$  polarization from s quark is negative: negative asymmetry observed in  $\Lambda K^+X$  at high  $z_{\Lambda}$ , where s to  $\Lambda$  fragmentation dominates
  - In  $\Lambda \pi^- X$  and  $\Lambda K^- X$   $u \rightarrow \Lambda$  fragmentation dominates at low  $z_{\Lambda}$ : u fragmentation to  $\Lambda$  is positive

## First extraction of $\Lambda$ polarised FF from Belle data



First extraction from Belle data within a TMD approach: D'Alessio, Murgia, Zaccheddu

PRD 102,114015



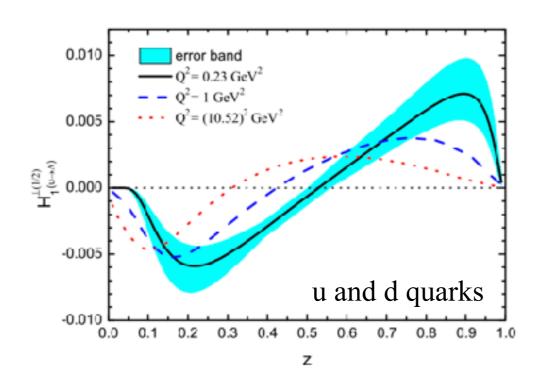
- Clear separation in flavour achieved
- The needed of sea-quark polarised FF is well supported
- Extracted the first pt dependence

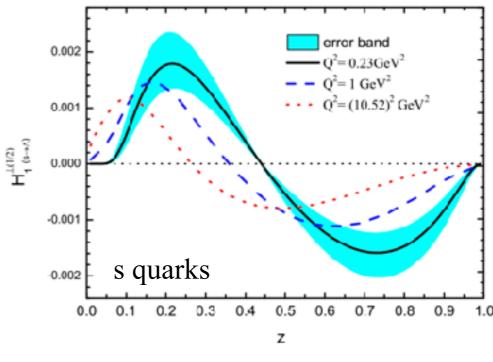
Higher statistic and complementary studies in other processes needed for a deeper understanding of this TMD-FF and solving the long-standing puzzle of the observed spontaneous transverse hyperon polarisation

# A hyperon Collins FFs prediction

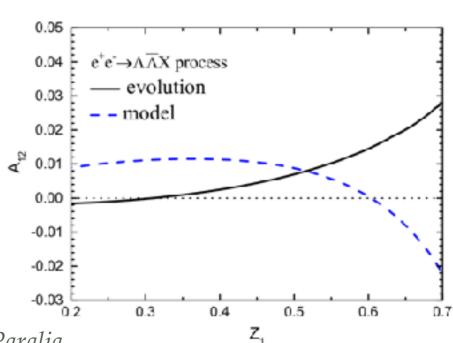
PRD97,114015

- Theoretical expectation from Collins FFs for  $\Lambda$  hyperon:  $e^+e^- \rightarrow \Lambda \overline{\Lambda} X$ 
  - Process accessible in e<sup>+</sup>e<sup>-</sup> facilities
  - Collins FF: complementary information to  $D_{1T^{\perp}}$





- Azimuthal asymmetry prediction in RF12 frame for  $0.2 < z_1 < 0.7$  (integrated over  $z_2$ )
- Increasing asymmetry with z, of the order of several percent

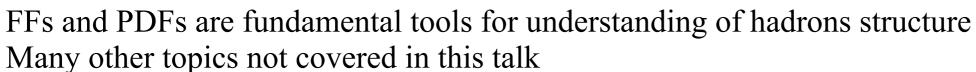


## Summary and Conclusions

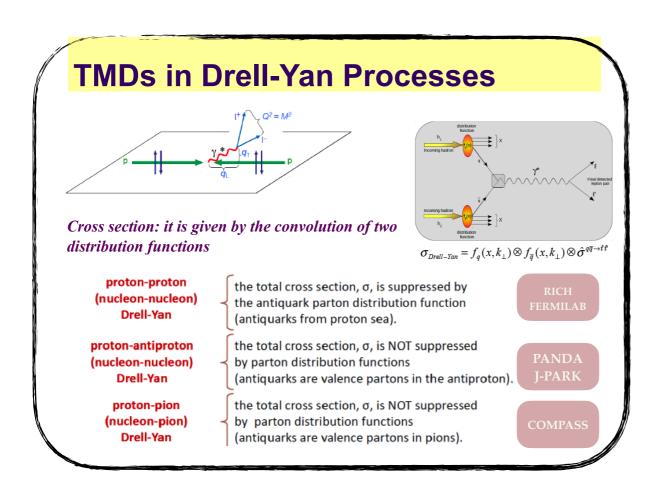
FFs and PDFs are fundamental tools for understanding of hadrons structure Many other topics not covered in this talk

- Di-hadron FF studies: PRL 107,072004
- Form Factors
  - Timelike and Spacelike region
- Generalized Parton Distribution functions (GPDs) and Transition Distribution Amplitudes (TDA)

# Summary and Conclusions

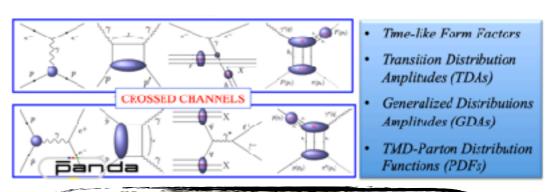


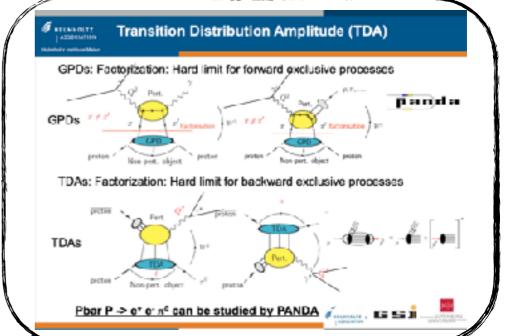
- Di-hadron FF studies: PRL 107,072004
- Form Factors
  - Timelike and Spacelike region
- Generalized Parton Distribution functions (GPDs) and Transition Distribution Amplitudes (TDA)



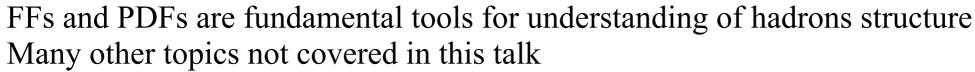


pan)da



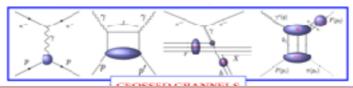


#### Summary and Conclusions



- Di-hadron FF studies: PRL 107,072004
- Form Factors
  - Timelike and Spacelike region
- Generalized Parton Distribution functions (GPDs) and Transition Distribution Amplitudes (TDA)

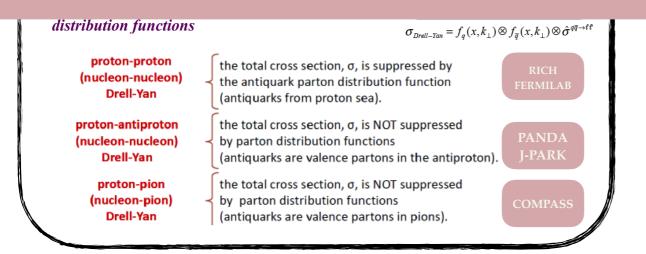


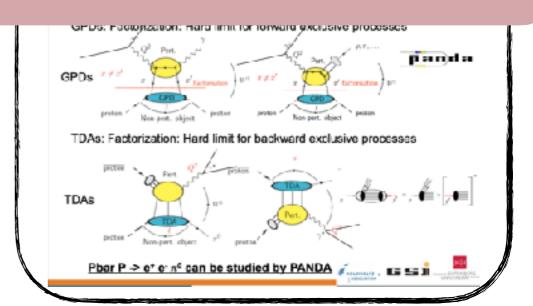


- · Time-like Form Factors
- Transition Distribution Amplitudes (TDAs)

#### TMDe in Drall-Van Processes

## Thanks for your attention









# Backup

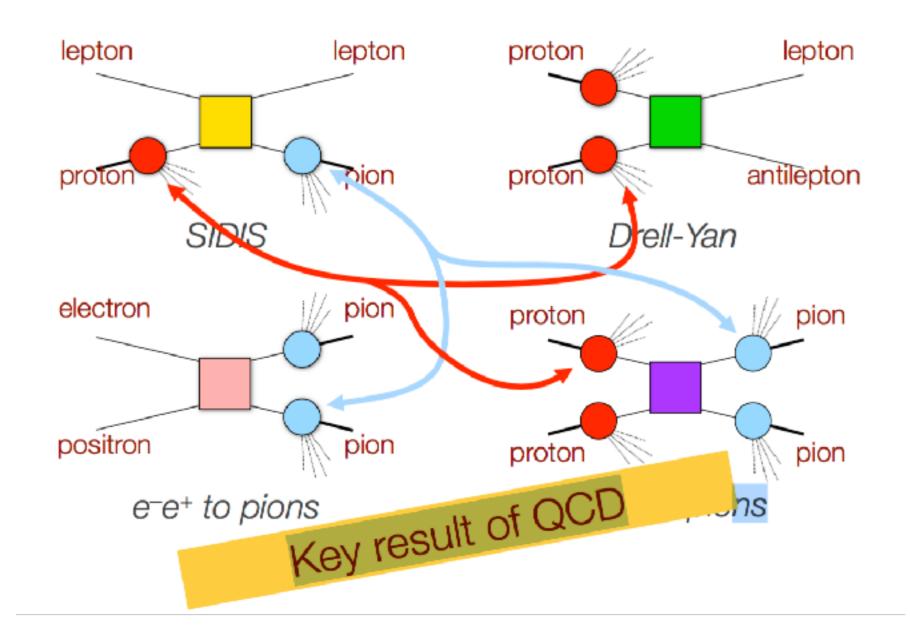




#### **Universality**

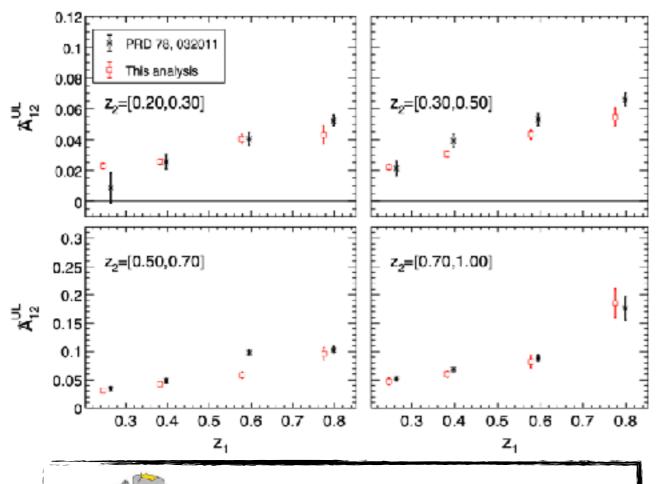
FF: independent of how the parton was produced

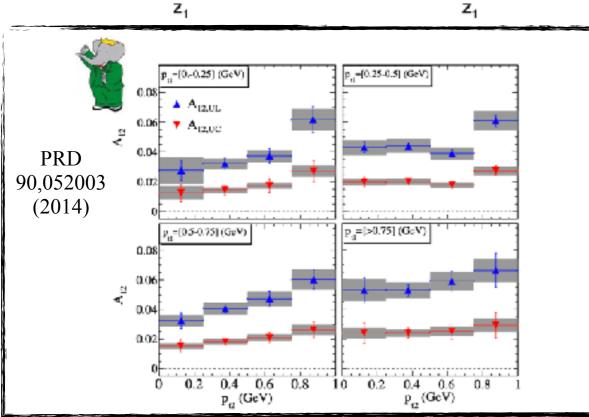
PDF: process independent nature



#### Azimuthal asymmetries of back-to-back $\pi^{\pm}$ - $(\pi^0, \eta, \pi^{\pm})$ from Belle





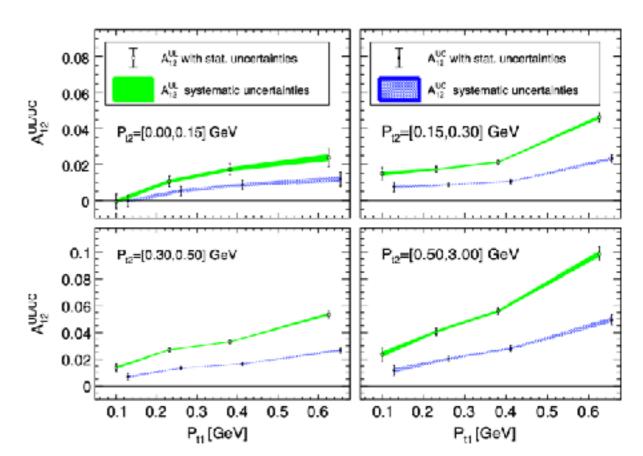


New Belle analysis: RF12 only

- No charm correction: only the estimated charm fractions are provided
- bin-by-bin generated-thrust axis correction (not the qqbar axis)
- Pt smering correction
- MC asymmetry subtraction

Only for results comparison:

- Vanishing charm asymmetry
- No thrust axis correction
- Kinematic factors taken into account



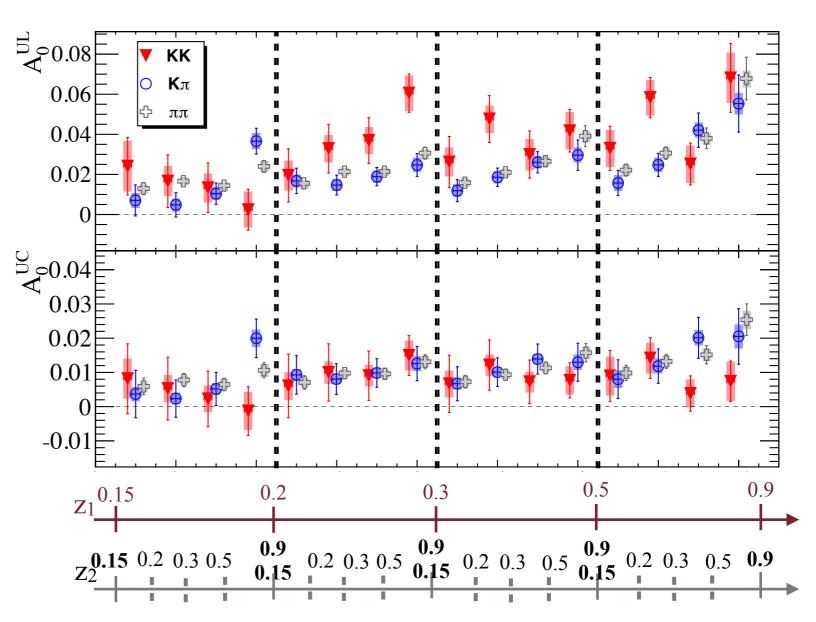
#### Results: RF0



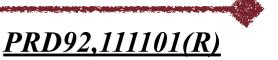
Simultaneous measurement of KK,  $K\pi$  and  $\pi\pi$  Collins asymmetries

- all corrections are applied
- Rising of the asymmetry as a function of z:
  - more pronounced for U/L
- AUL KK asymmetry slightly higher than pion asymmetry for high z
- \* KK asymmetry consistent with zero at lower z

Note that A<sup>UL</sup> and A<sup>UC</sup> asymmetries are obtained using the same data sample, and are strongly correlated

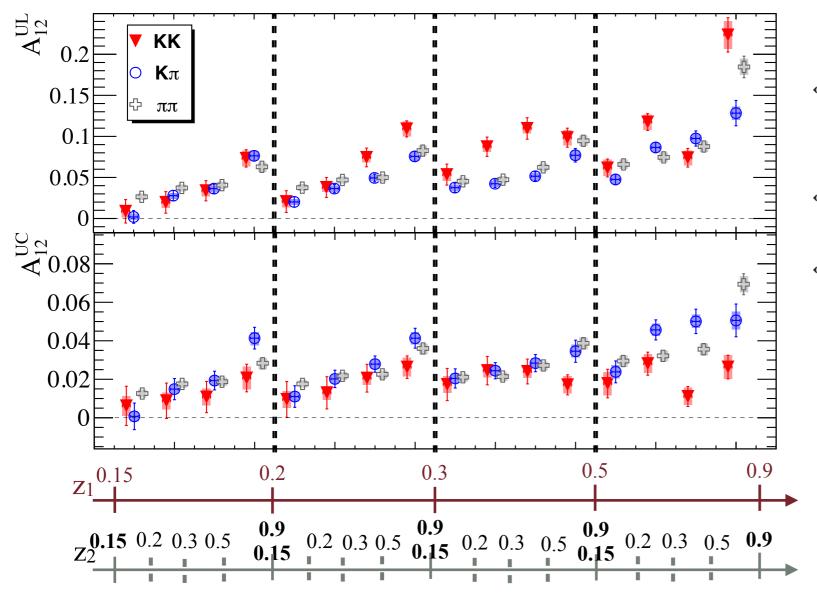


#### Results: RF12



Simultaneous measurement of KK,  $K\pi$  and  $\pi\pi$  Collins asymmetries

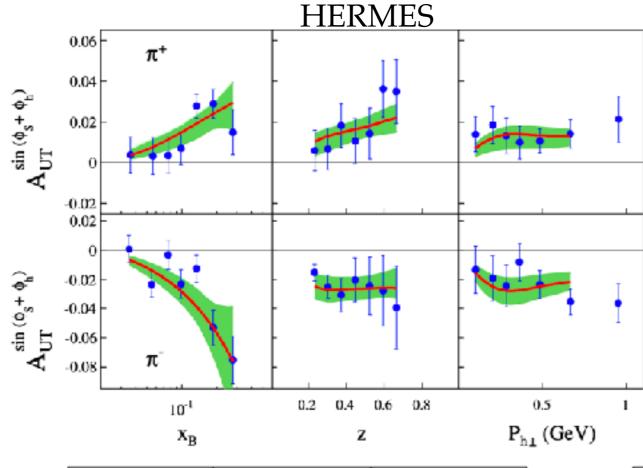
• all corrections are applied



- Rising of the asymmetry as a function of z:
  - more pronounced for U/L
- A<sup>UL</sup> KK asymmetry slightly higher than pion asymmetry for high z
- \* KK asymmetry consistent with zero at lower z

Note that A<sup>UL</sup> and A<sup>UC</sup> asymmetries are obtained using the same data sample, and are strongly correlated

PRD93,014009



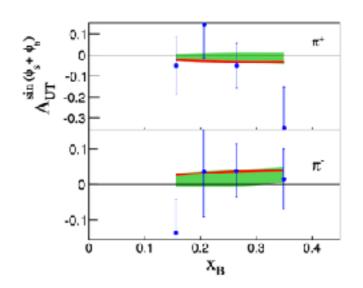
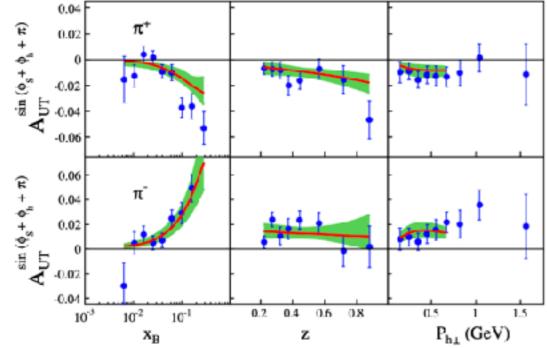
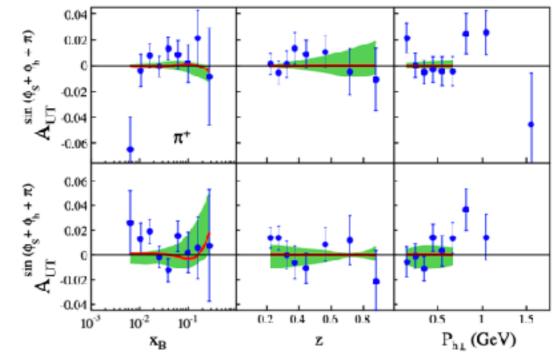


FIG. 14. Description of Collins asymmetries measured by the JLab's HALL A [9] as a function of  $x_B$  on the effective neutron target. The shaded region corresponds to our estimate of a 90% C.L. error band.



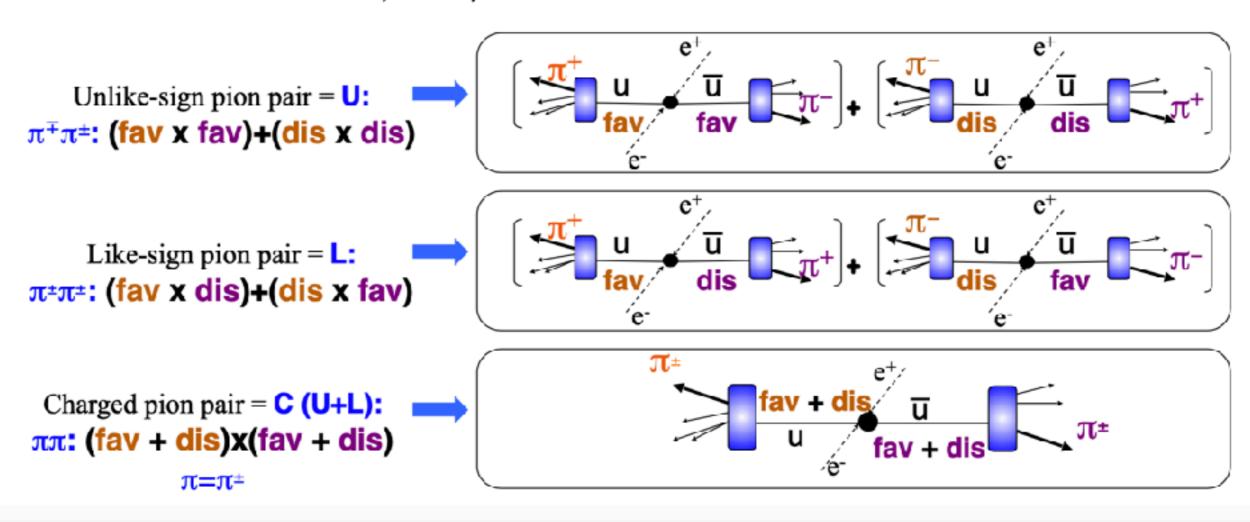


**COMPASS** 

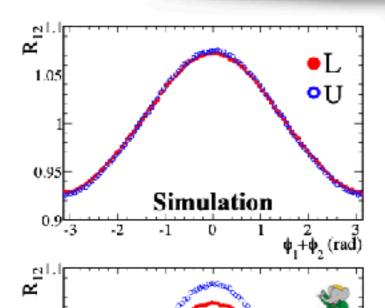
## Favored and Disfavored processes

Different combinations of charged pions ⇒ sensitivity to favored or disfavored FFs

- favored process: fragmentation of a quark of flavor q into a hadron with a valence quark of the same flavor: i.e.:  $u \rightarrow \pi^+$ ,  $d \rightarrow \pi^-$
- disfavored for  $d\rightarrow \pi^+$ ,  $u\rightarrow \pi^-$ , and  $s\rightarrow \pi^\pm$



# Raw asymmetries and Double Ratios



1.05

- Collins asymmetry:
- consider all the U (unlike) and L (like) sign pion pairs
- make histograms of  $\phi_{\alpha} = \phi_1 + \phi_2$  or  $2\phi_0$  ( $\alpha = 12,0$ )
- The MC generator (JETSET) does not include the Collins effect
- → flat distribution is expected
- strong modulation due to acceptance of the detector
- but similar distribution for U and L pairs
- Data shows difference between U and L distributions, that can be ascribed to the Collins effect



Acceptance effects can be reduced by performing the ratio of U/L sign pion pairs (or U/C):

φ<sub>1</sub>+φ<sub>2</sub> (rad)

- MC: consistent with a flat distribution

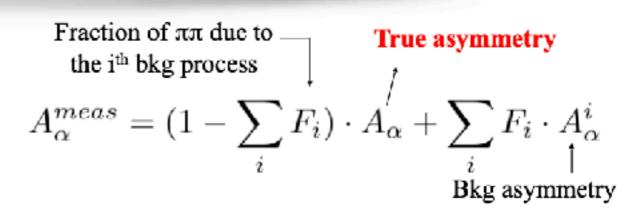
Data

- Data: cosine modulation clearly visible

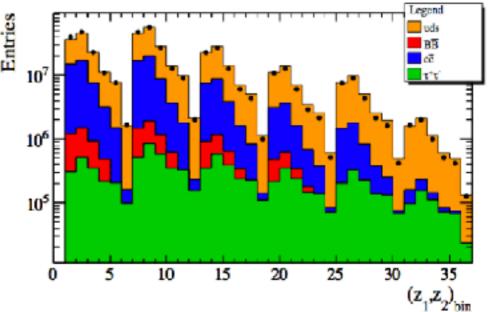
$$\frac{R_{\alpha}^{U}}{R_{\alpha}^{L(C)}} = \frac{N^{U}(\phi_{\alpha})/< N^{U}(\phi_{\alpha})>}{N^{L(C)}(\phi_{\alpha})/< N^{L(C)}(\phi_{\alpha})>} \rightarrow B_{\alpha}^{UL(UC)} + \underbrace{A_{\alpha}^{UL,(UC)} \cdot \cos(\phi_{\alpha})}^{-3} \cdot \cos(\phi_{\alpha})$$

# Backgrounds: contributions and corrections

- In each bin, the data sample includes pairs from
  - signal uds events
  - $B\overline{B}$  events (small, mostly at low z)
  - $\overline{CC}$  events (important at low/medium z)
  - $\tau^+\tau^-$  events (important at high z)



- We must calculate these quantities:
  - F<sub>i</sub> using MC sample; we assign MC-data difference in each bin as systematic error
  - $A^{B\overline{B}}$  must be zero; we set  $A^{B\overline{B}} = 0$
  - $A^{\tau}$  small in simulation; checked in data; we set  $A^{\tau} = 0$

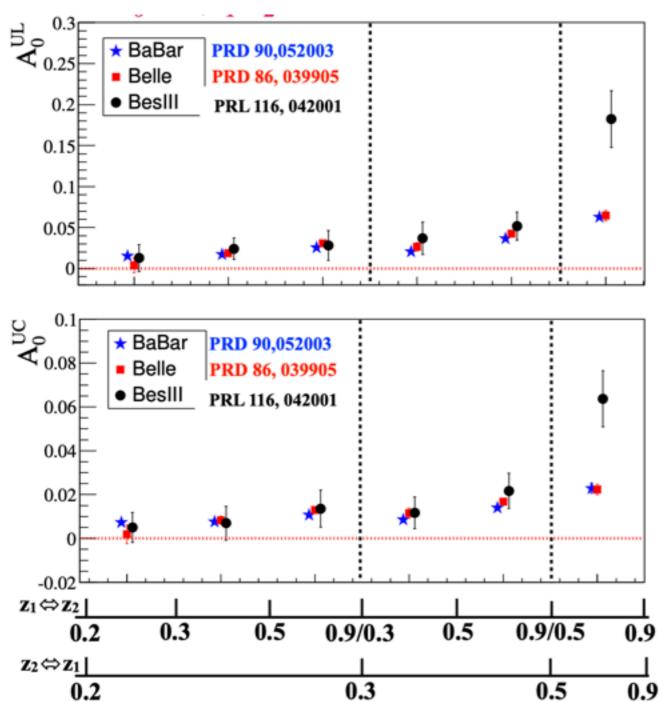


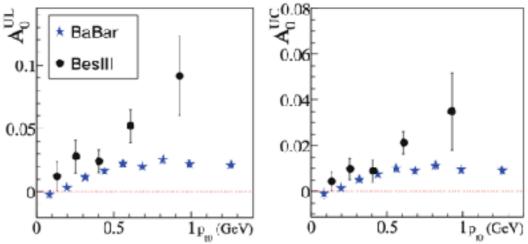
- Charm background contribution is about 30% on average
  - Both fragmentation processes and weak decays can introduce azimuthal asymmetries
  - We used a D\*=-enhanced control sample to estimate its effect

$$\begin{cases} A_{\alpha}^{meas} &= (1 - F_c - F_B - F_\tau) \cdot A_{\alpha} + F_c \cdot A_{\alpha}^{ch} \\ A_{\alpha}^{D^*} &= f_c \cdot A_{\alpha}^{ch} + (1 - f_c - f_B) \cdot A_{\alpha} \end{cases}$$

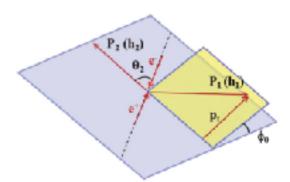
### Collins FFs Results Comparison



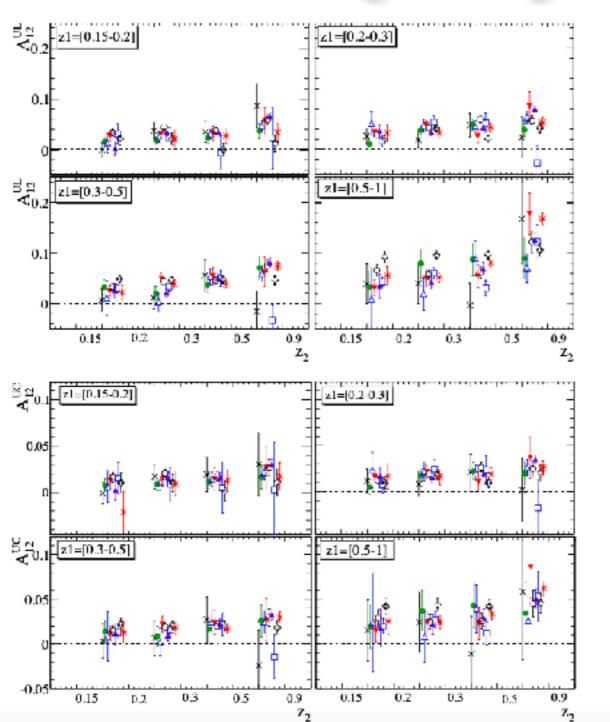




• Larger asymmetry from BESIII data, in agreement with the prediction (PRD93,014009)



## 4-D: asymmetry vs. $(z_{1,z_2})x(p_{t1,p_{t2}})$



We study the asymmetries in the RF12 frame in a four-dimensional space:

$$(z_1,z_2,p_{t1},p_{p2})$$

- We use 4 z<sub>i</sub> and 3 p<sub>t</sub> intervals
- Test to probe the factorization of the Collins fragmentation functions
- Powerful tools to access p<sub>t</sub> z correlation

$$(p_{t1}, p_{t2}) = [0.,0.25][0.,0.25]$$

$$(\mathbf{p}_{t1}, \mathbf{p}_{t2}) = [0., 0.25][0.25, 0.5]$$

$$(p_{t1},p_{t2})=[0.25,0.5][0.,0.25]$$

$$(\mathbf{p}_{i1},\mathbf{p}_{i2}) = [0.25,0.5][0.25,0.5]$$

$$(p_{i1}p_{i2})=[>0.5][0.,0.25]$$

$$(\mathbf{p}_{11},\mathbf{p}_{12})=[>0.5][0.25,0.5]$$

$$(p_{t1}, p_{t2}) = [0., 0.25][>0.5]$$

$$(p_{t1}, p_{t2}) = [0.25, 0.5] > 0.5]$$

$$(p_{ij},p_{ij})=[>0.5][>0.5]$$

#### Azimuthal asymmetries of back-to-back $\pi^{\pm}$ - $(\pi^0, \eta, \pi^{\pm})$ from Belle

PRD100,092008

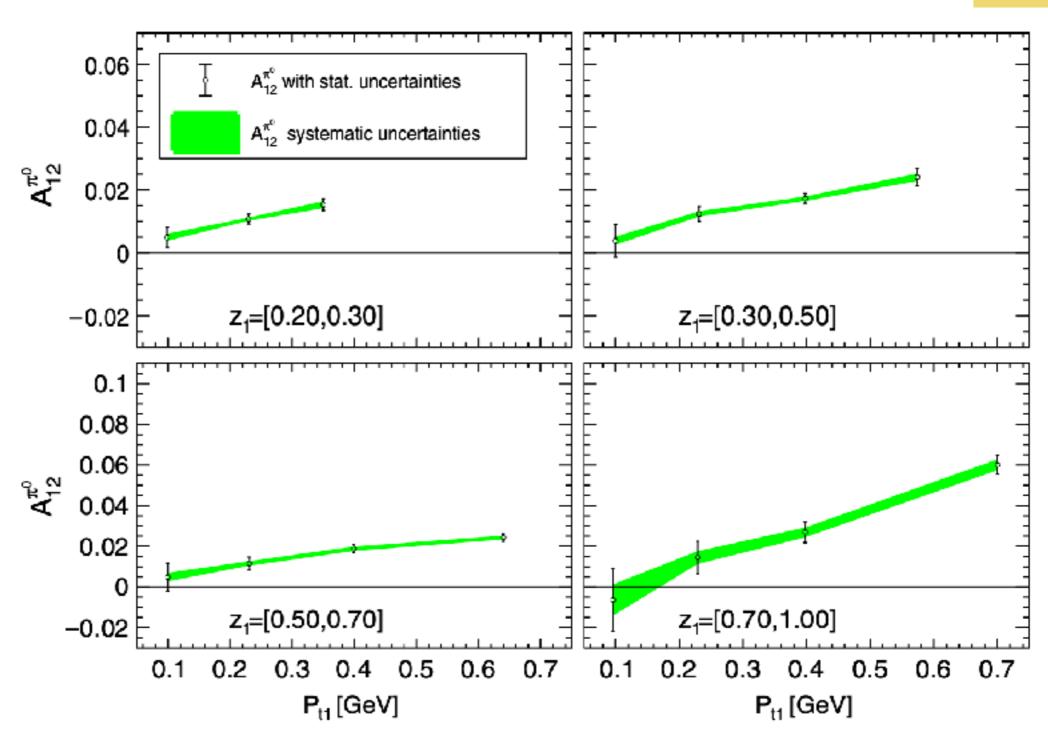
$$\mathcal{R}_{12}^{UL} \approx 1 + \cos(\phi_{12}) \frac{\sin^{2}(\theta)}{1 + \cos^{2}(\theta)} \times \left\{ \frac{5(H_{1}^{\perp, \text{fav}} \otimes H_{1}^{\perp, \text{fav}} + H_{1}^{\perp, \text{dis}} \otimes H_{1}^{\perp, \text{dis}}) + 2H_{1, s \to \pi}^{\perp, \text{dis}} \otimes H_{1, s \to \pi}^{\perp, \text{dis}}}{5(D_{1}^{\text{fav}} \otimes D_{1}^{\text{fav}} + D_{1}^{\text{dis}} \otimes D_{1}^{\text{dis}}) + 2D_{1, s \to \pi}^{\text{dis}} \otimes D_{1, s \to \pi}^{\text{dis}}} - \frac{10H_{1}^{\perp, \text{fav}} \otimes H_{1}^{\perp, \text{dis}} + 2H_{1, s \to \pi}^{\perp, \text{dis}} \otimes H_{1, s \to \pi}^{\perp, \text{dis}}}{10D_{1}^{\text{fav}} \otimes D_{1}^{\text{dis}} + 2D_{1, s \to \pi}^{\text{dis}} \otimes D_{1, s \to \pi}^{\text{dis}}} \right\},$$

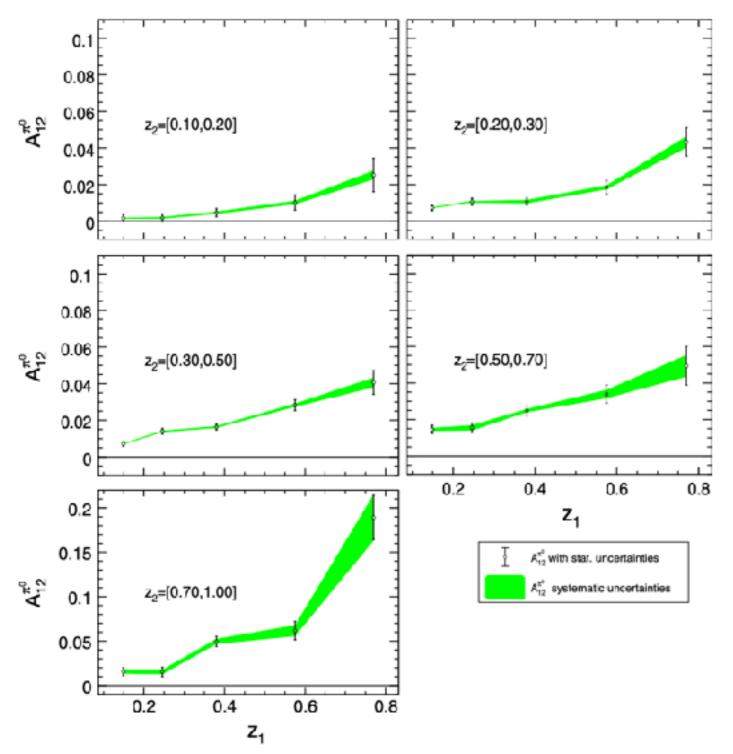
$$\mathcal{R}_{12}^{UC} \approx 1 + \cos(\phi_{12}) \frac{\sin^2(\theta)}{1 + \cos^2(\theta)} \times \left\{ \frac{5(H_1^{\perp, \text{fav}} \otimes H_1^{\perp, \text{fav}} + H_1^{\perp, \text{dis}} \otimes H_1^{\perp, \text{dis}}) + 2H_{1, s \to \pi}^{\perp, \text{dis}} \otimes H_{1, s \to \pi}^{\perp, \text{dis}}}{5(D_1^{\text{fav}} \otimes D_1^{\text{fav}} + D_1^{\text{dis}} \otimes D_1^{\text{dis}}) + 2D_{1, s \to \pi}^{\text{dis}} \otimes D_{1, s \to \pi}^{\text{dis}}} \right. \\ \left. - \frac{5(H_1^{\perp, \text{fav}} + H_1^{\perp, \text{dis}}) \otimes (H_1^{\perp, \text{fav}} + H_1^{\perp, \text{dis}}) + 4H_{1, s \to \pi}^{\perp, \text{dis}} \otimes H_{1, s \to \pi}^{\perp, \text{dis}}}{5(D_1^{\text{fav}} + D_1^{\text{dis}}) \otimes (D_1^{\text{fav}} + D_1^{\text{dis}}) + 4D_{1, s \to \pi}^{\text{dis}} \otimes D_{1, s \to \pi}^{\text{dis}}} \right\},$$

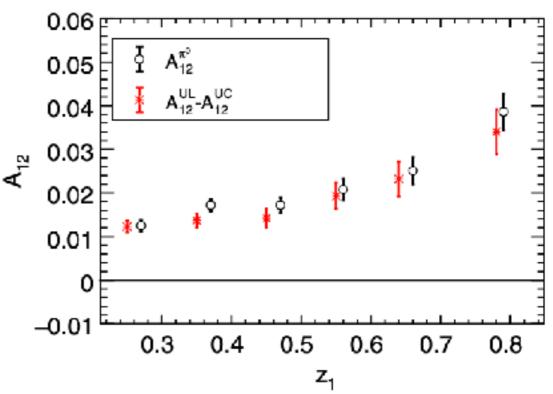
$$\begin{split} \mathcal{R}_{12}^{\pi^0} &= \frac{R_{12}^{0\pm}}{R_{12}^L} \approx 1 + \cos(\phi_{12}) \frac{\sin^2(\theta)}{1 + \cos^2(\theta)} \times \left\{ \frac{5(H_1^{\perp, \text{fav}} + H_1^{\perp, \text{dis}}) \otimes (H_1^{\perp, \text{fav}} + H_1^{\perp, \text{dis}}) + 4H_{1, s \to \pi}^{\perp, \text{dis}} \otimes H_{1, s \to \pi}^{\perp, \text{dis}}}{5(D_1^{\text{fav}} + D_1^{\text{dis}}) \otimes (D_1^{\text{fav}} + D_1^{\text{dis}}) + 4D_{1, s \to \pi}^{\text{dis}} \otimes D_{1, s \to \pi}^{\text{dis}}} \right\} \\ &- \frac{10H_1^{\perp, \text{fav}} \otimes H_1^{\perp, \text{dis}} + 2H_{1, s \to \pi}^{\perp, \text{dis}} H_{1, s \to \pi}^{\perp, \text{dis}}}{10D_1^{\text{fav}} \otimes D_1^{\text{dis}} + 2D_{1, s \to \pi}^{\text{dis}} \otimes D_{1, s \to \pi}^{\text{dis}}} \right\}. \end{split}$$

$$\mathcal{R}_{12}^{\eta} = \frac{R_{12}^{\eta \pm}}{R_{12}^{L}} \approx 1 + \cos(\phi_{12}) \frac{\sin^{2}(\theta)}{1 + \cos^{2}(\theta)} \times \left\{ \frac{5(H_{1}^{\perp, \text{fav}_{\eta}} + H_{1}^{\perp, \text{dis}_{\eta}}) \otimes (H_{1}^{\perp, \text{dis}} + H_{1}^{\perp, \text{fav}}) + 4H_{1, s \to \eta}^{\perp, \text{dis}} \otimes H_{1, s \to \pi}^{\perp, \text{dis}}}{5(D_{1}^{\perp, \text{fav}_{\eta}} + D_{1}^{\perp, \text{dis}_{\eta}}) \otimes (D_{1}^{\perp, \text{dis}} + D_{1}^{\perp, \text{fav}}) + 4D_{1, s \to \eta} \otimes D_{1, s \to \pi}^{\text{dis}}} \right\} \\ - \frac{10H_{1}^{\perp, \text{fav}} \otimes H_{1}^{\perp, \text{dis}} + 2H_{1, s \to \pi}^{\perp, \text{dis}} \otimes H_{1, s \to \pi}^{\perp, \text{dis}}}{10D_{1}^{\perp, \text{fav}} \otimes D_{1}^{\perp, \text{dis}} + 2D_{1, s \to \pi}^{\text{dis}} \otimes D_{1, s \to \pi}^{\text{dis}}} \right\}.$$

PRD100,092008



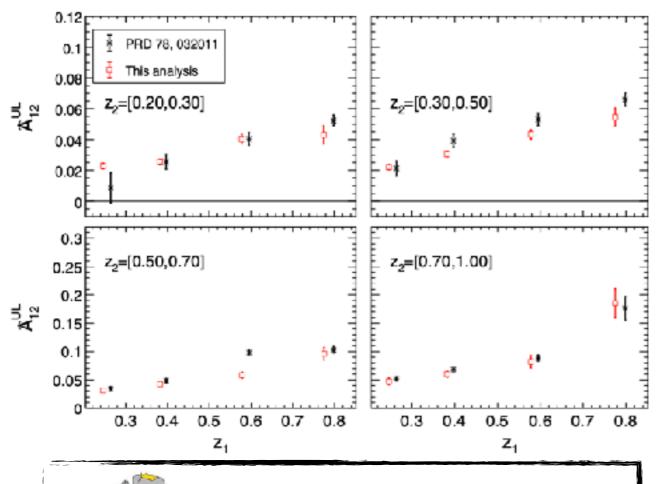


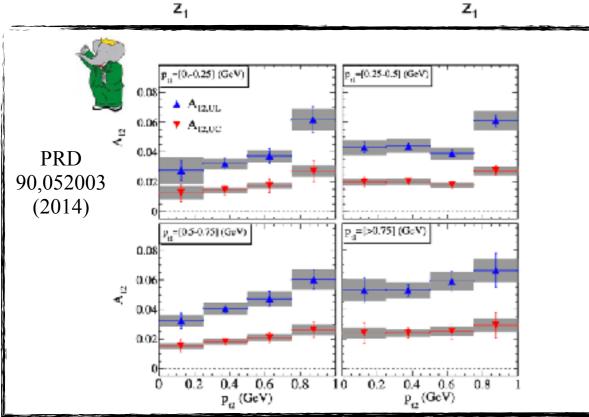


• Good agreement between  $A^{UL}$ - $A^{UC}$  and  $A^{\pi 0}$  as expected from isospin relation (taking into account statistical and systematic uncertainties)

#### Azimuthal asymmetries of back-to-back $\pi^{\pm}$ - $(\pi^0, \eta, \pi^{\pm})$ from Belle





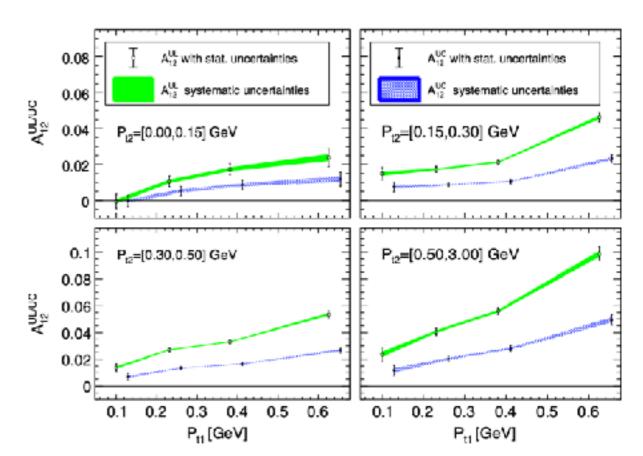


New Belle analysis: RF12 only

- No charm correction: only the estimated charm fractions are provided
- bin-by-bin generated-thrust axis correction (not the qqbar axis)
- Pt smering correction
- MC asymmetry subtraction

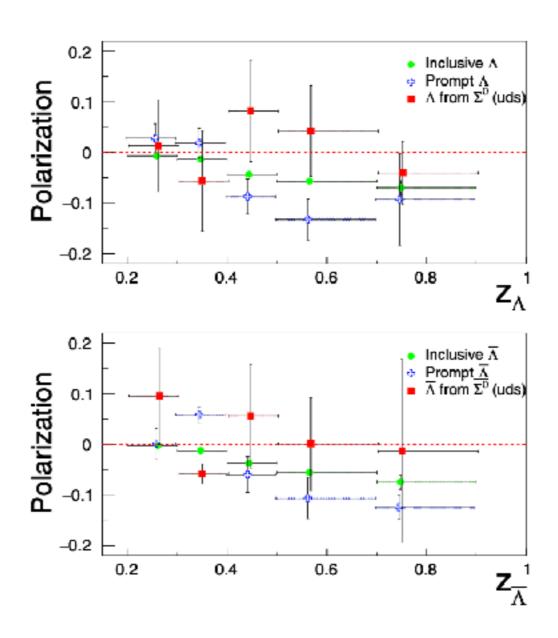
Only for results comparison:

- Vanishing charm asymmetry
- No thrust axis correction
- Kinematic factors taken into account



## $\Lambda/\overline{\Lambda}$ hyperon polarization at Belle

PRL122,042001



• Charm corrected unfolded transverse polarisation also obtained ==> large statistical uncertainties

## $\overline{\Lambda}/\Lambda$ hyperon polarization



PRD **100**,014029

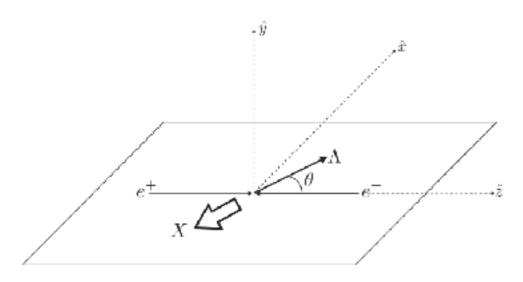
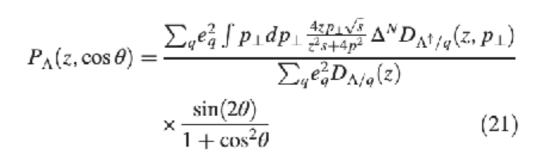
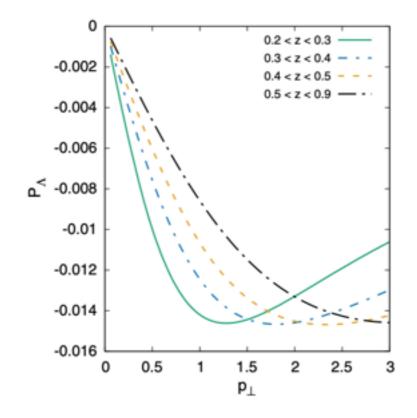
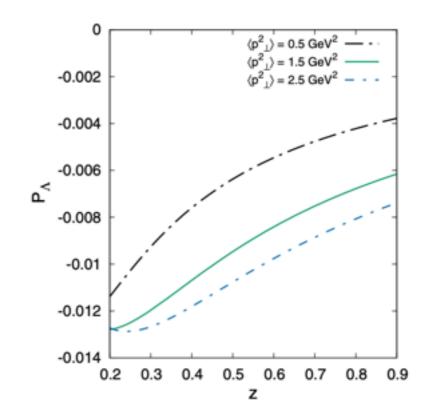


FIG. 1. Definition of the c.m. reference frame for the process  $e^+e^- \to \Lambda X$ .



$$P_{\Lambda}(z, p_{\perp}) = \frac{\sum_{q} e_{q}^{2} \frac{4z p_{\perp} \sqrt{s}}{z^{2} s + 4p_{\perp}^{2}} \Delta^{N} D_{\Lambda^{\uparrow}/q}(z, p_{\perp})}{\sum_{q} e_{q}^{2} 2\pi D_{\Lambda/q}(z, p_{\perp})} \times \frac{\int d(\cos \theta) \sin(2\theta)}{\int d(\cos \theta) (1 + \cos^{2} \theta)}$$
(22)





#### Helicity Formalism: TMD Fragmentation Functions for quarks

Quark Polarisation States determine the Hadron ones

$$\rho_{\lambda_h,\lambda_h'}^{h,S_h} \hat{D}_{h/q,s_q}(z,\mathbf{k}_{\perp h}) = \sum_{\lambda_q,\lambda_q'} \rho_{\lambda_q,\lambda_q'}^{q,s_q} \hat{D}_{\lambda_q,\lambda_q'}^{\lambda_h,\lambda_h'}(z,\mathbf{k}_{\perp h})$$

$$P_j^h \hat{D}_{h/q,s_q} = \hat{D}_{S_j/q,s_q}^h - \hat{D}_{-S_j/q,s_q}^h = \Delta \hat{D}_{S_j/s_q}^h.$$

Helicity density matrix

$$\rho^{i,s_i}_{\lambda_i,\lambda_i'} = \frac{1}{2} \begin{pmatrix} 1 + P_z^i & P_x^i - iP_y^i \\ P_x^i + iP_y^i & 1 - P_z^i \end{pmatrix}$$

#### Amsterdam notation:

$$\begin{array}{ll} \text{Collins} & \Delta^N D_{h/q^\uparrow} \longleftrightarrow H_1^\bot \\ \\ \text{Polarizing} & \Delta D_{S_Y/q}^h \longleftrightarrow D_{1T}^\bot \end{array} \quad \begin{cases} \Delta^- D_{S_Y/s_T}^{h_1} \\ \Delta D_{S_X/s_T}^{h_1} \end{cases} \quad \leftrightarrow \begin{cases} H_1 \\ H_{1T}^\bot \end{cases}$$

Marco Zaccheddu – Università degli Studi di Cagliari

#### 8 independent Fragmentation Functions

		Hadron		
	Pol. States	U	L	T
Q u a r k	U	$(\hat{D}_{h/q})$		$(\Delta \hat{D}^h_{S_Y/q})$
	L		$\Delta \hat{D}^{h/q}_{S_Z/s_L}$	$\Delta \hat{D}_{S_X/s_L}^{h/q}$
	Т	$(\Delta^N D_{h/q^\uparrow})$	$\Delta \hat{D}_{S_Z/s_T}^{h/q}$	$\Delta \hat{D}_{S_X/s_T}^{h/q} / \Delta^- \hat{D}_{S_Y/s_T}^{h/q}$

#### **SIDIS** cross section

$$\frac{d\sigma^{h}}{dx\,dy\,d\phi_{S}\,dz\,d\phi\,d\mathbf{P}_{h\perp}^{2}} = \frac{\alpha^{2}}{xyQ^{2}}\frac{y^{2}}{2\left(1-\epsilon\right)}\left(1+\frac{\gamma^{2}}{2x}\right)$$

$$\left\{ F_{\text{UU,T}} + \epsilon F_{\text{UU,L}} + \sqrt{2\epsilon (1+\epsilon)} \cos(\phi) F_{\text{UU}}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{\text{UU}}^{\cos(2\phi)} \right]$$

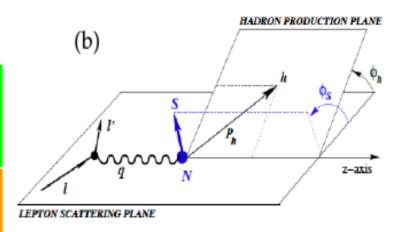
+ 
$$\lambda_l \left[ \sqrt{2\epsilon (1 - \epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right]$$

+ 
$$S_L$$
  $\left[\sqrt{2\epsilon (1+\epsilon)} \sin(\phi) F_{\mathrm{UL}}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{\mathrm{UL}}^{\sin(2\phi)}\right]$ 

+ 
$$S_L \lambda_l \left[ \sqrt{1 - \epsilon^2} F_{LL} + \sqrt{2\epsilon (1 - \epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right]$$

+ 
$$S_T$$
  $\left[\sin(\phi - \phi_S) \left(F_{\text{UT},\text{T}}^{\sin(\phi - \phi_S)} + \epsilon F_{\text{UT},\text{L}}^{\sin(\phi - \phi_S)}\right) + \epsilon \sin(\phi + \phi_S) F_{\text{UT}}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{\text{UT}}^{\sin(3\phi - \phi_S)} + \sqrt{2\epsilon (1 + \epsilon)} \sin(\phi_S) F_{\text{UT}}^{\sin(\phi_S)} + \sqrt{2\epsilon (1 + \epsilon)} \sin(2\phi - \phi_S) F_{\text{UT}}^{\sin(2\phi - \phi_S)}\right]$ 

$$+ S_{T} \lambda_{l} \left[ \sqrt{1 - \epsilon^{2}} \cos (\phi - \phi_{S}) F_{LT}^{\cos (\phi - \phi_{S})} + \sqrt{2\epsilon (1 - \epsilon)} \cos (\phi_{S}) F_{LT}^{\cos (\phi_{S})} + \sqrt{2\epsilon (1 - \epsilon)} \cos (2\phi - \phi_{S}) F_{LT}^{\cos (2\phi - \phi_{S})} \right]$$



Unpolarized proton and lepton

Longitudinally polarized proton, unpolarized lepton

Longitudinally polarized proton, polarized lepton

Transversely polarized proton, unpolarized lepton

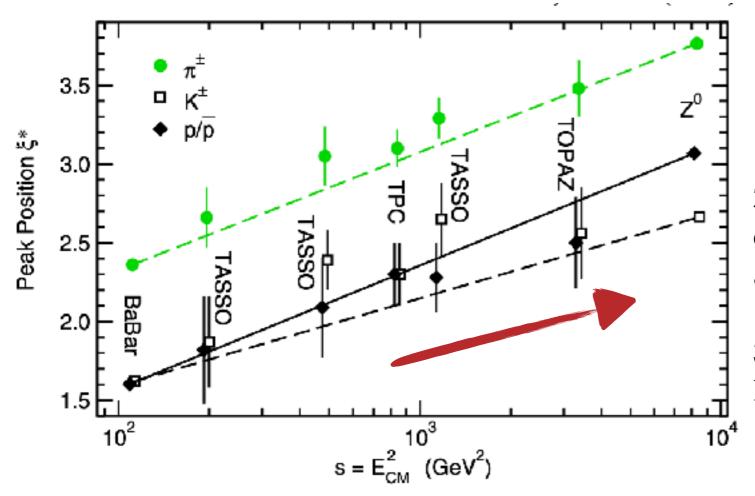
Transversely polarized proton, polarized lepton

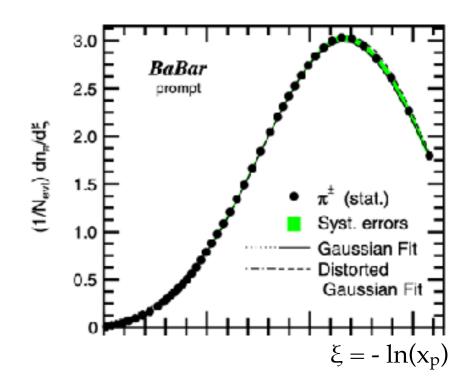
JHEP093(07)

Single hadron FF							
Unpolarized ingredients	Polarized ingredients	Flavor sensitivity					
Single hadron cross sections: $e^+e^- \rightarrow hX$ $D^h_{1,q}(z,Q^2)$ PRL111 (2013) 062002	Azimuthal asymmetries: $e^+e^- \rightarrow (h)(h)X$ , $\cos(\phi_1 + \phi_2)$ $H_{1,q}^{\perp(1)h}(z,Q^2)$ PRL 96 (2006) 232002	Unpol SIDIS, pp: $\frac{d\sigma}{dz}$ $e^{+}e^{-} \rightarrow (h)(h)X$ PRD92 (2015) 092007 PRD101(2020) 092004 and scale dependence					
Transverse momentum dependent FFs: $e^+e^- \rightarrow (h)X$ $D_{1,q}^h(z,k_T,Q^2)$ PRD 99 (2019) 112006	Transverse momentum dependent asymmetries $e^+e^- \rightarrow (h)(h)X$ , $\cos(\phi_1 + \phi_2), Q_t$ $H_{1,q}^{\perp h}(z, k_T, Q^2)$ PRD100 (2019) 92008	Polarizing $\Lambda$ fragmentation $ \frac{PRL \ 122 \ (2019), \ 042001}{D_{1,q}^{\perp h}(z,k_T,Q^2)} $					
Dihadron FF (IFF)							
Unpolarized ingredients	Polarized ingredients	Flavor sensitivity					
Dihadron cross sections $e^+e^- \rightarrow (hh)X$ $D_{1,q}^{h_1h_2}(z,m,Q^2)$ PRD96 (2017) 032005	Azimuthal asymmetries: $e^+e^- \to (hh)(hh)X$ , $\cos(\phi_1 + \phi_2)$ , $H_{1,q}^{h_1,h_2,\triangleleft}(z,Q^2,M_h)$ PRL107 (2011) 072004	Unpol SIDIS, pp: $\frac{d^2\sigma}{dzdm}$					

## Test of MLLA+LHPD QCD predictions

Modified Leading Logarithmic Approximation (MLLA) with Local Parton-Hadron Duality (LHPD) ansatz (Z.Phys.C27,65): 1) the multiplicity distributions vs  $\xi = -\ln(x_p)$  should be Gaussian near the peak



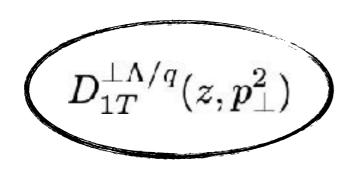


- 2) the peak position should decrease exponentially with increasing hadron mass at a given  $E_{cm}$ 
  - observed that  $\xi_K \approx \xi_p$
- 3)  $\xi^*$  should increase logarithmically with  $E_{mc}$  for a given hadron type
  - similar slope for pions and protons, different for kaons
  - possibly due to flavour composition changing with  $E_{cm}$

# $\Lambda/\overline{\Lambda}$ hyperon polarization at Belle

PRL **122**,042001

•  $e^+e^- \rightarrow \Lambda(\overline{\Lambda})X$  and  $e^+e^- \rightarrow \Lambda(\overline{\Lambda})h^\pm X$  using 800.4 fb<sup>-1</sup> collected by Belle at or near  $\sqrt{s}=10.58$  GeV

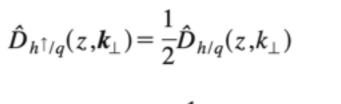


==> production of transversely polarized  $\Lambda$  hyperons from unpolarised quark

==> chiral-even function (the sign of this function can be easily determined)

Distribution of protons from  $\Lambda$  decays with a transverse polarisation P:

$$\frac{1}{N}\frac{dN}{d\cos\theta} = 1 + \alpha P\cos\theta,$$



along ↑ direction
PRD 100,014029

hadron polarisation vector

$$+\frac{1}{2}\Delta^{N}D_{h^{\uparrow/q}}(z,k_{\perp})\frac{\hat{\boldsymbol{P}}_{h}\cdot(\boldsymbol{p}_{q}\times\boldsymbol{k}_{\perp})}{|\boldsymbol{p}_{q}\times\boldsymbol{k}_{\perp}|}$$

 $\theta$ : angle between the axis  $\stackrel{\wedge}{n} \propto \stackrel{\wedge}{T} \times \stackrel{\wedge}{p}_{\Lambda}$  and the proton momentum in the  $\Lambda$  rest frame;  $\alpha$  is the world average value of the parity-violating decay asymmetry for the  $\Lambda$  (from PDG)

==> Significant transverse polarization is observed

