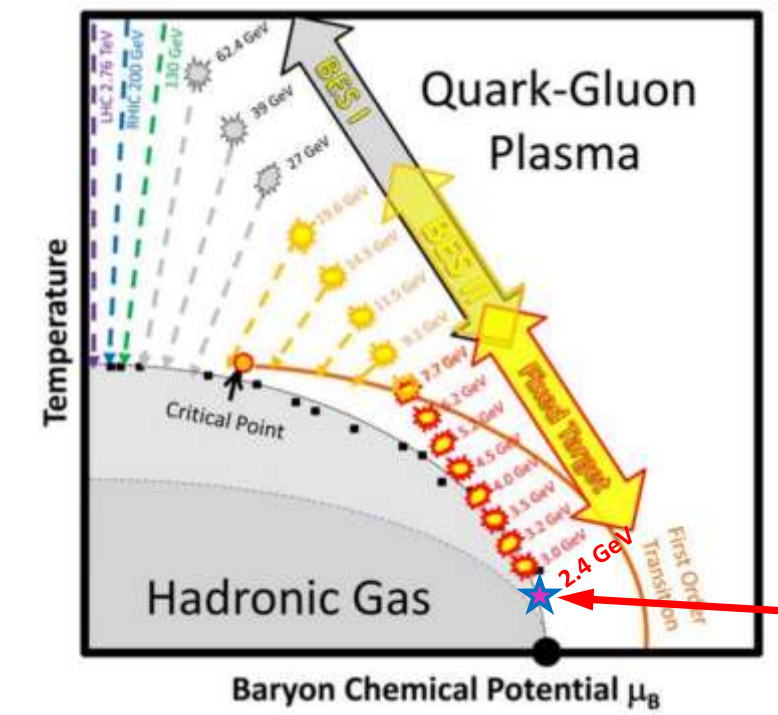


Critical fluctuations studied with HADES

30/11/2021 FANI-2021 | R. Holzmann (GSI) for the HADES collaboration

Exploring the QCD phase diagram through HI collisions



- proton detection in HADES
- centrality estimators
- cumulants & correlators
- **volume fluctuation corrections**
- conclusions & outlook

HADES Au+Au at $\sqrt{s_{NN}} = 2.41$ GeV
→ Adamczewski-Musch et al. PRC 102 (2020)

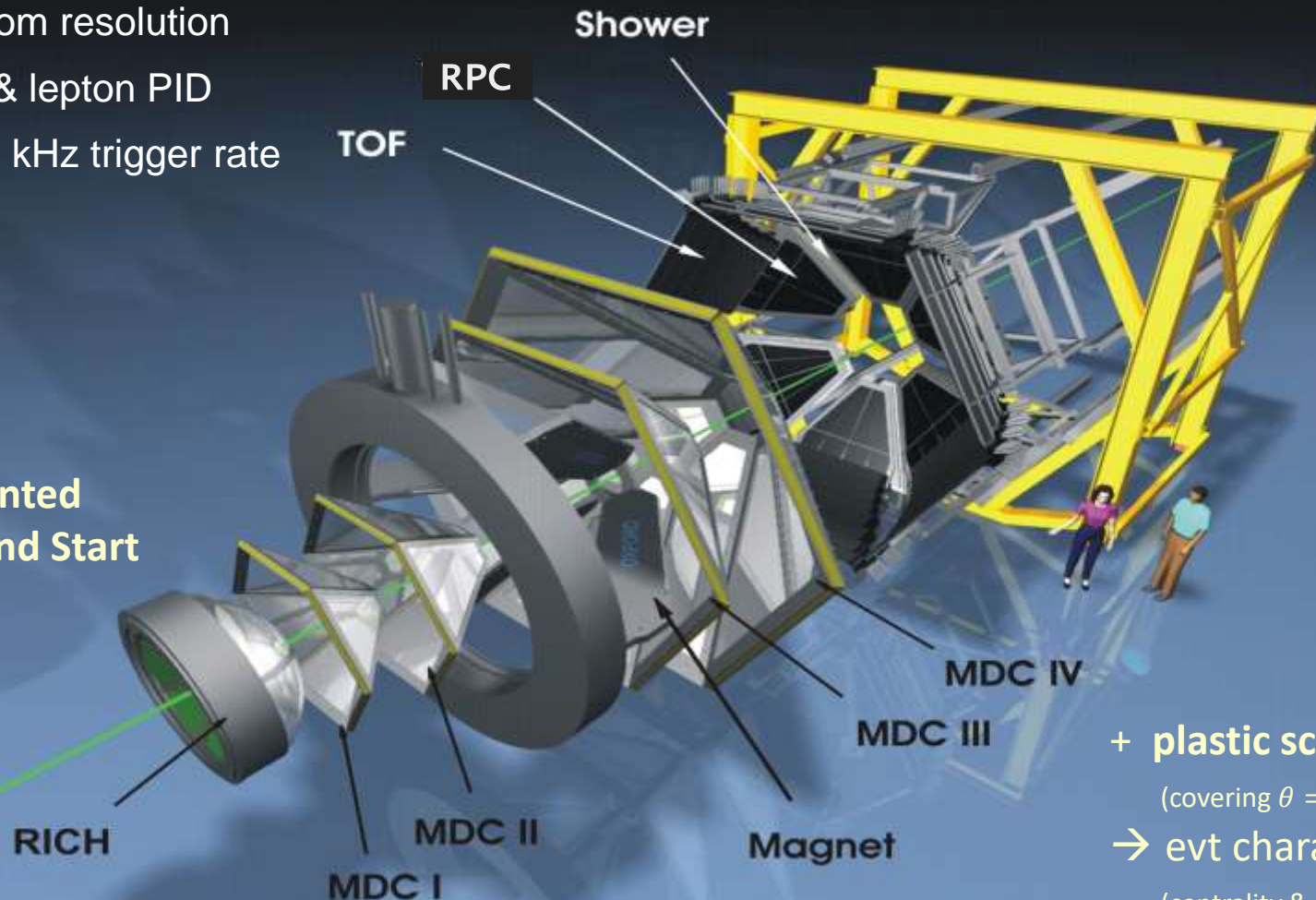
Ag+Ag at $\sqrt{s_{NN}} = 2.55$ GeV
→ analysis in progress ...

The HADES detector at GSI

High Acceptance DiElectron Spectrometer
(setup used in Au+Au run)

- large acceptance
- 2-3% mom resolution
- hadron & lepton PID
- up to 20 kHz trigger rate

+ segmented diamond Start



+ plastic scint. FW
(covering $\theta = 0.5^\circ - 7.5^\circ$)
→ evt characterization
(centrality & event plane)

Event cleaning in HADES

Segmented gold target:

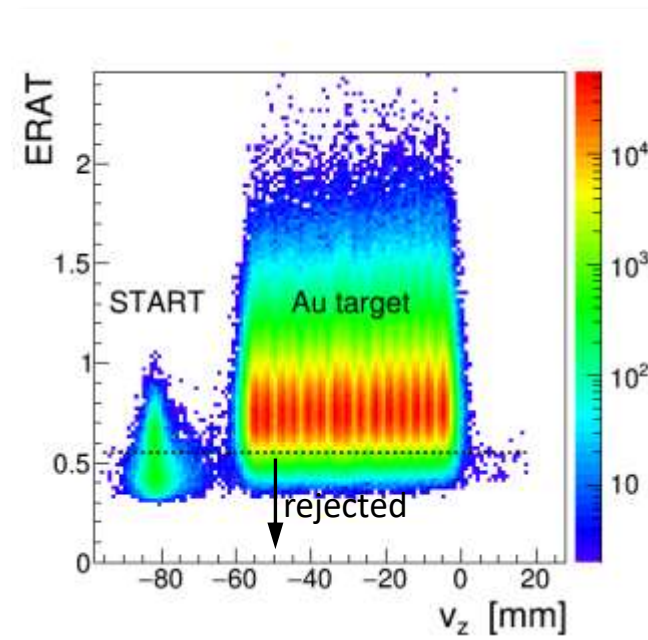
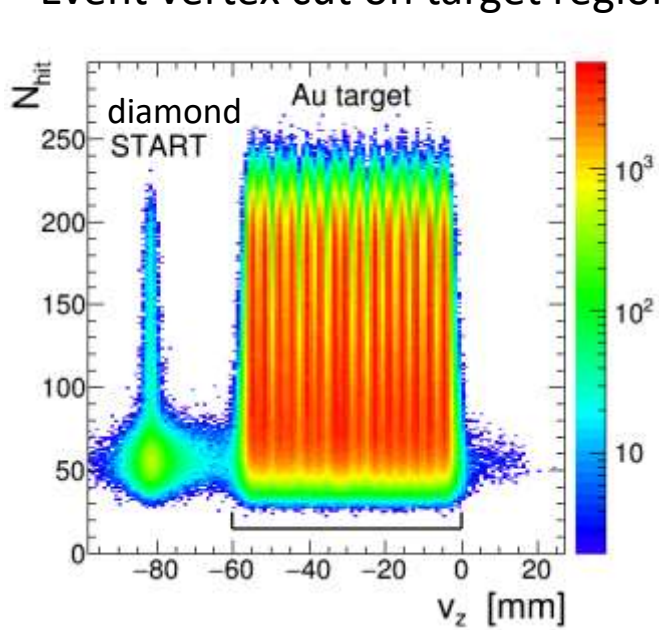
- ^{197}Au material
- 15 discs of $\varnothing = 2.2$ mm mounted on kapton strips
- $\Delta z = 3.6$ mm
- 2.0% interaction prob.



Kindler et al.,
NIM A 655 (2011) 95

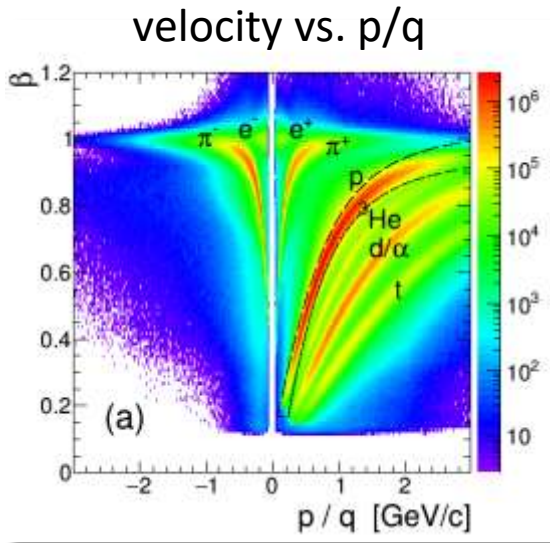
Remove Au+C bkgd on the kapton with a cut on $ERAT = \sum E_t / \sum E_l$

Event vertex cut on target region



→
beam direction

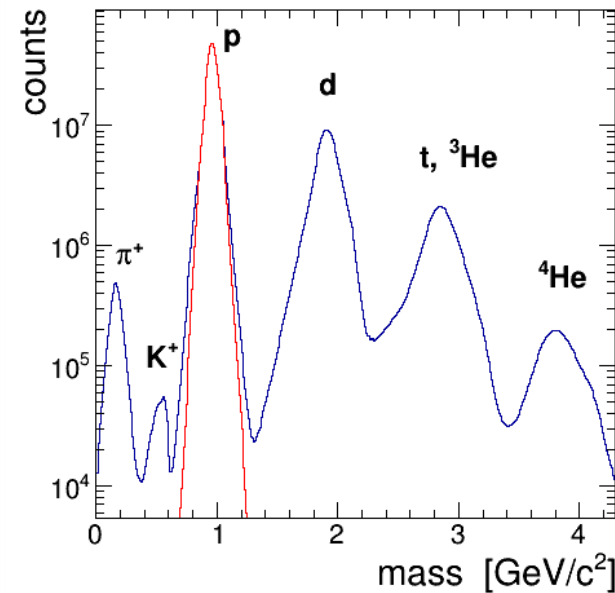
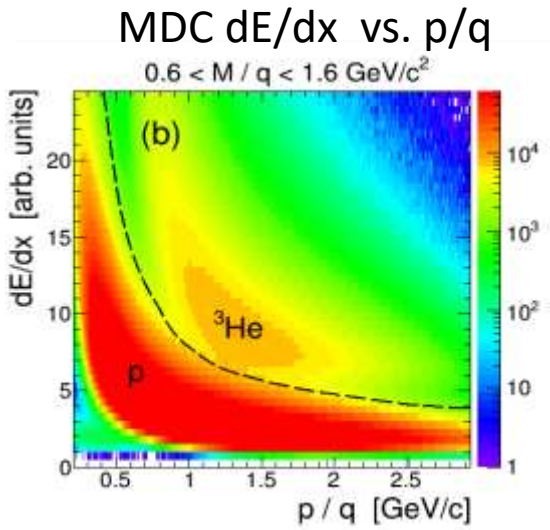
Particle ID in HADES



Hadron ID based on

- ToF
- momentum
- dE/dx

Mass spectrum and **accepted protons**



Proton fluctuation signal purity

within same evt	■ Proton pid impurities	$\leq 10^{-3}$
	■ Weak decays, e.g. $\Lambda \rightarrow p + \pi^{-}$, $\Sigma \rightarrow p + \pi^{-}$	$\leq 6.5 \cdot 10^{-4}$
	■ Knock-out (spallation) protons <ul style="list-style-type: none">■ from secondary reactions in target / target holder■ 50% pp, 45% np, <5% πp (Geant3 + GCalor)	$\leq 3 \cdot 10^{-3}$
different evt classes	■ Au + C reactions on target holder (8 μ m kapton) foils <ul style="list-style-type: none">■ suppressed by trigger & centrality selection■ asymmetric rapidity distribution $\gamma > \gamma_0$	$\leq 10^{-3}$
	■ Event pile-up (central evt + min. bias evt)	$\leq 3 \cdot 10^{-5}$

relative contribution

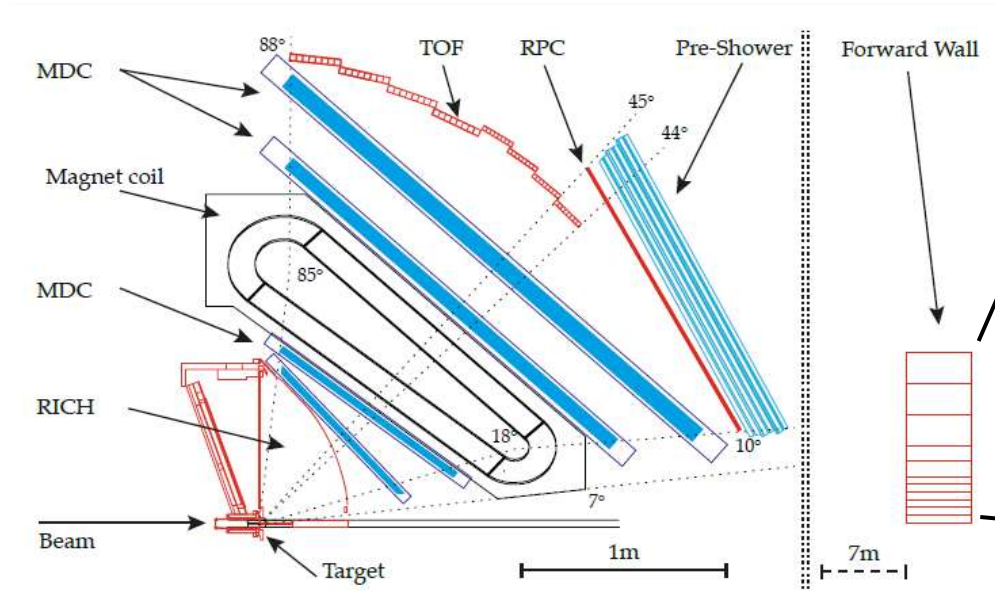
Centrality selection with the Forward Wall

In 1.23 GeV/u Au+Au collisions:

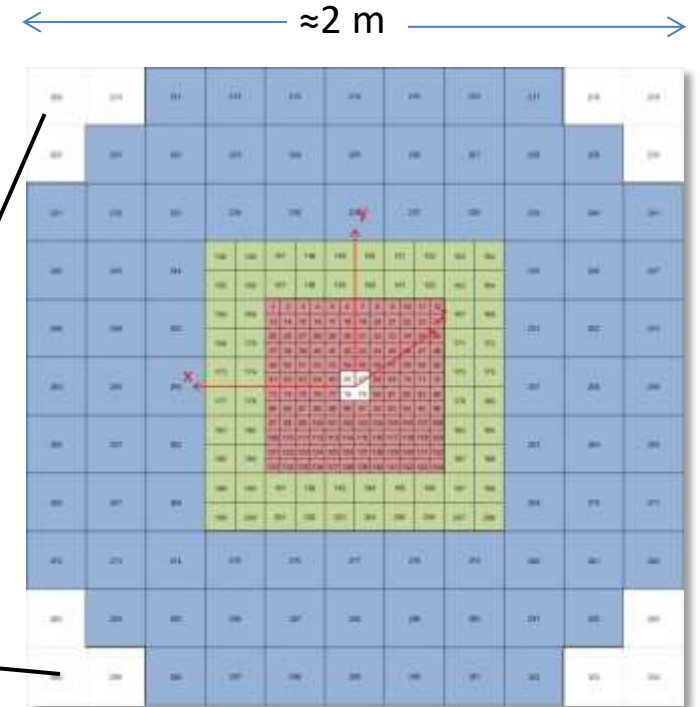
- protons & clusters dominate
 - centrality selection based on
 - hit mult in TOF & RPC
 - or track mult
 - or **FW sum of charges**
- ➔ reduce auto-correlations!

FW made of plastic scintillator tiles covering polar angles $\theta = 0.5^\circ - 7.5^\circ$ i.e. a pseudorapidity of $\eta = 2.7 - 5.4$ (HADES itself covers $y \approx 0 - 1.8$)

➔ Used for event-plane reconstruction



cross section of 1/6 HADES sector

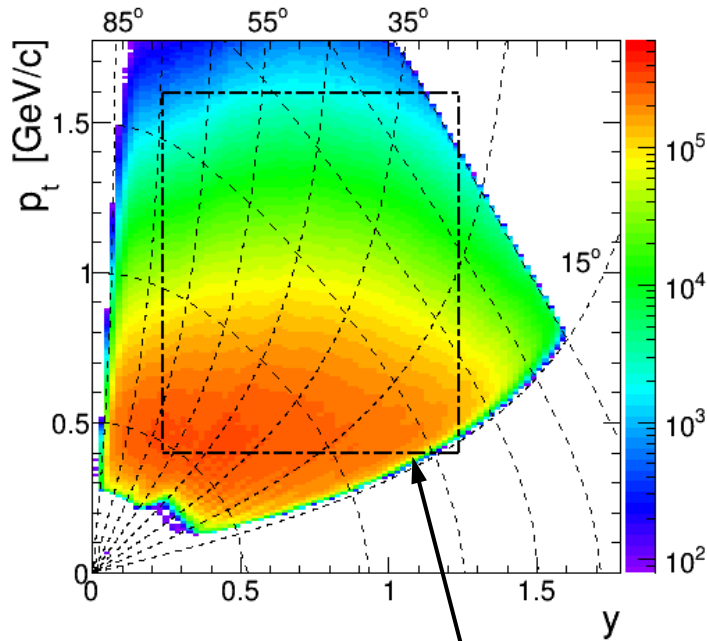


4x4, 8x8, 16x16 cm² tiles

Proton distributions in Au+Au at $\sqrt{s} = 2.41 \text{ GeV}$

HADES $y - p_t$ coverage for protons

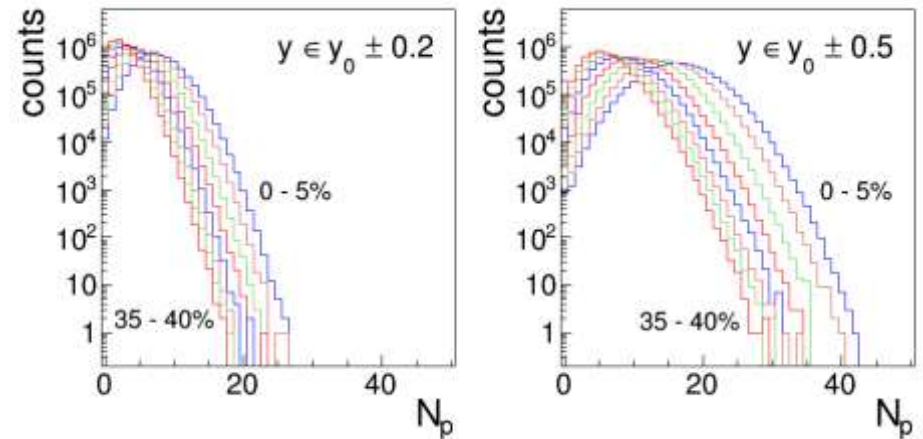
$$y_0 = 0.74$$



Useful acceptance
for fluctuation analysis

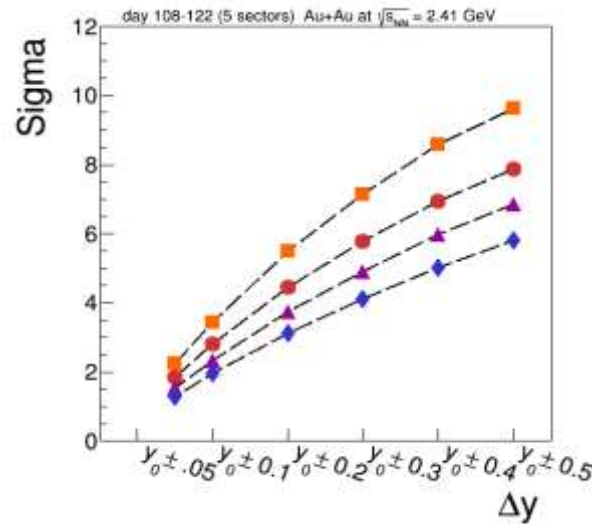
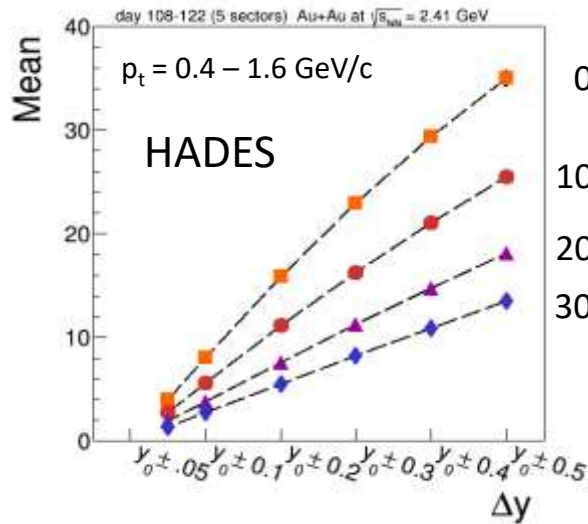
$$\left\{ \begin{array}{l} y = y_0 \pm 0.5 \\ p_t = 0.4 - 1.6 \text{ GeV}/c \end{array} \right.$$

Proton multiplicity distributions



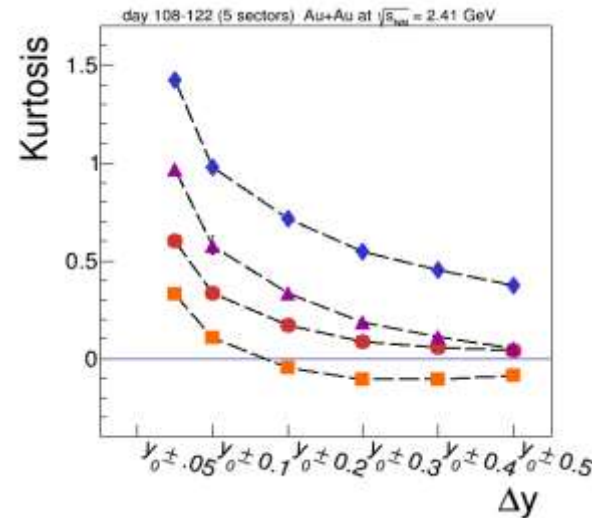
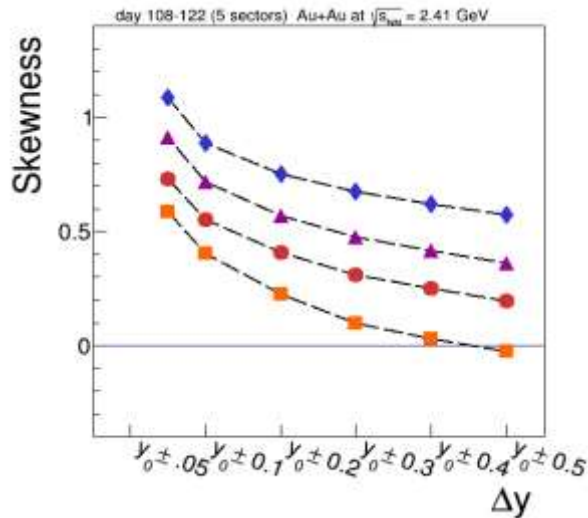
Analysis based on $1.6 \cdot 10^8$ Au+Au events
divided into 5%-centrality bins in the range
of the 0 - 40% most central events

E-by-E efficiency-corrected moments in Au+Au



From the moments we can calculate

- factorial moments
- cumulants
- reduced cumulants
- factorial cumulants
- etc.

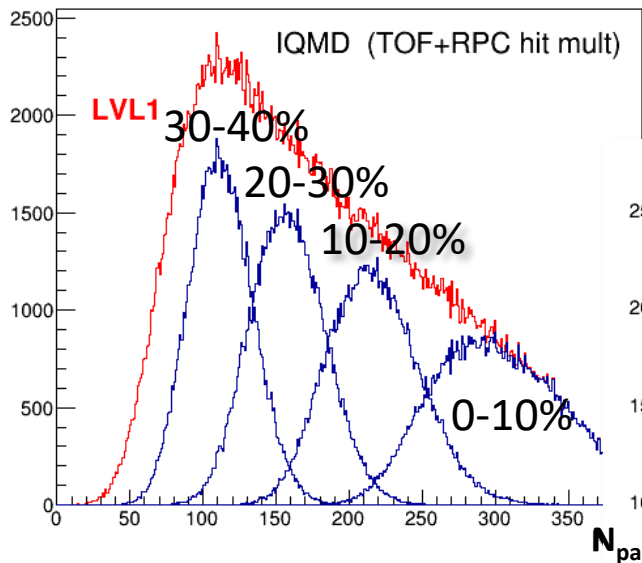


but ...

... are dominated by e-by-e variations of the source volume!

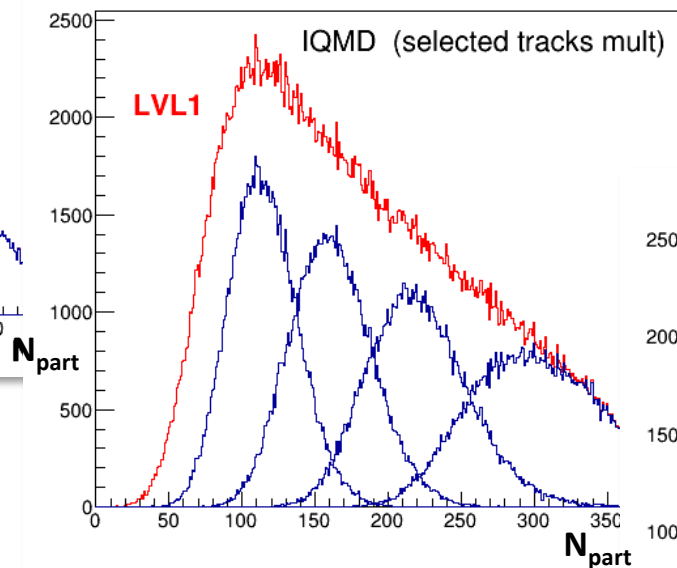
Centrality estimators in HADES

N_{hit} centrality



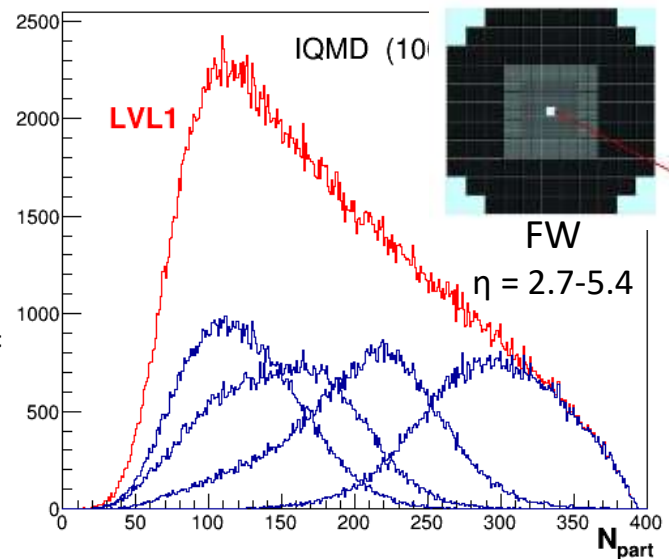
see also EPJA 54 (2018) 85

N_{trk} centrality



Based on a simulation with
IQMD + MST clusterizer
(evts provided by Y. Leifels)

Σ_{FW} centrality



→ Centrality selections lead to large volume fluctuations $\delta V \equiv \delta N_{part}$, characterized by volume cumulants v_n

→ We use FW for fluct. analysis

Volume fluctuation corrections (VFC)

V. Skokov, B. Friman & K. Redlich PRC 88 (2013), A. Rustamov et al. NPA 960 (2017),
Sugiura, Nonaka & Esumi PRC 100 (2019), Esumi & Nonaka NIM A987 (2021)

Averaging the volume-dependent proton distribution moments

$$\langle N_{prot}^n \rangle = \int P(V) \underbrace{\sum N_{prot}^n P(N_{prot}|V)}_{\langle N_{prot}^n \rangle_V} dV$$

one obtains volume contributions to the observed reduced cumulants c_n :

observed true

$$c_1 = \kappa_1 \tag{1}$$

$$c_2 = \kappa_2 + \kappa_1^2 v_2 \tag{2}$$

$$c_3 = \kappa_3 + \kappa_1^3 v_3 + 3\kappa_1 \kappa_2 v_2 \tag{3}$$

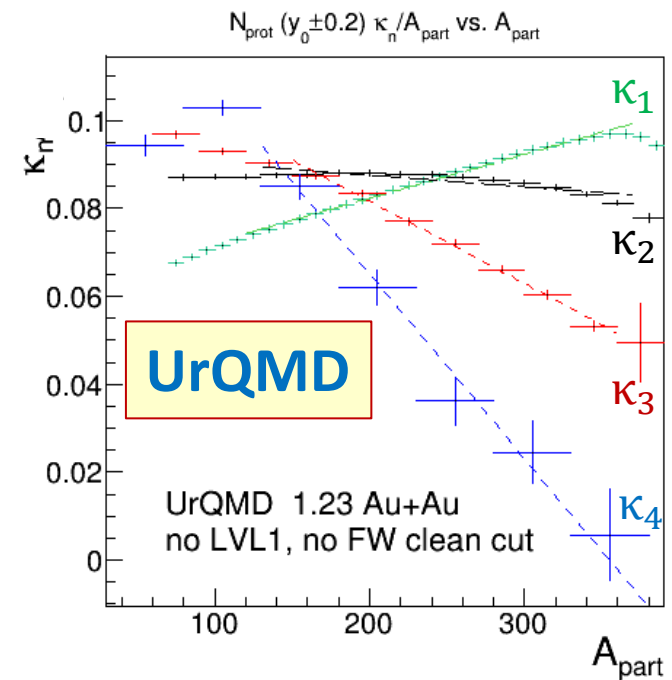
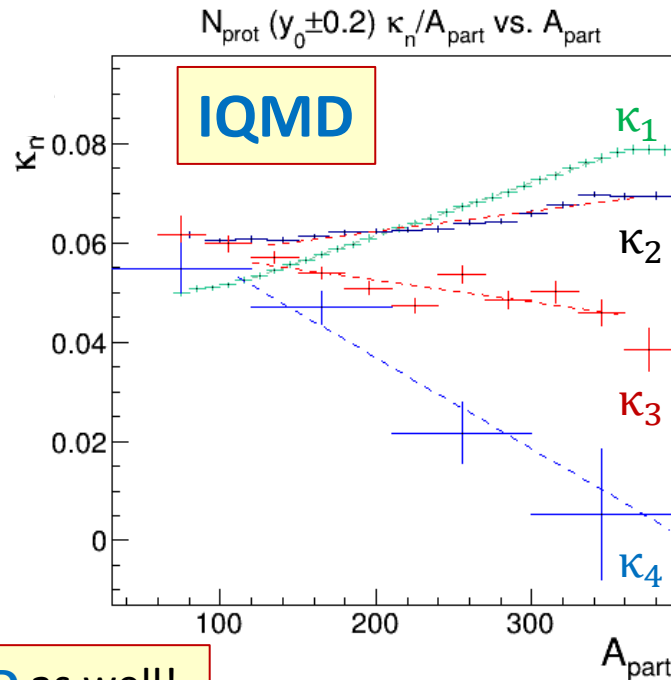
$$c_4 = \kappa_4 + \kappa_1^4 v_4 + 6\kappa_1^2 \kappa_2 v_3 + (4\kappa_1 \kappa_3 + 3\kappa_2^2) v_2 \tag{4}$$

where κ_n are true reduced cumulants, c_n are observed cumulants and v_n are volume cumulants, and assuming all κ_n are constant!

Reduced proton cumulants in transport models

Definition: $\kappa_n = K_n/V$ with $V = A_{\text{part}}$

See also Sugiura, Nonaka & Esumi
PRC 100 (2019)



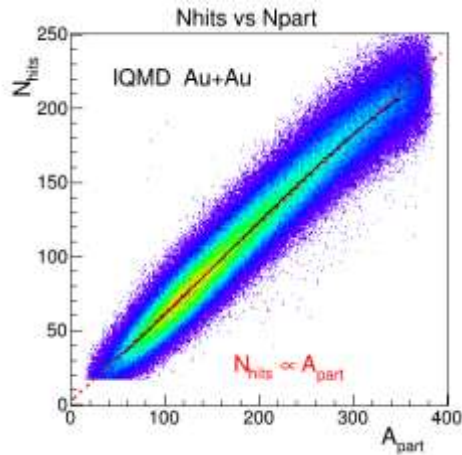
→ HSD as well!

→ Skokov et al. assumption ($\kappa_n = \text{cst}$) is not fulfilled!

→ Extend formalism to NLO: $\kappa_n \rightarrow \kappa_n + \kappa'_n \cdot (V - \bar{V})$
and N2LO: $\kappa_n \rightarrow \kappa_n + \kappa'_n \cdot (V - \bar{V}) + \kappa''_n (V - \bar{V})^2$

details given in Adamczewski-Musch et al. PRC 102 (2020)

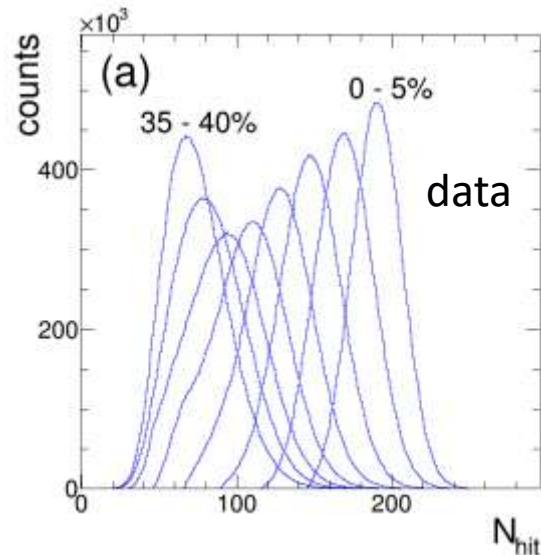
Ad-hoc approach: N_{hit} as a proxy for N_{part}



IQMD simulation shows that N_{hit} is proportional to N_{part}

→ use N_{hit} as proxy for vol. flucs.
i.e. rescale & adjust the v_n

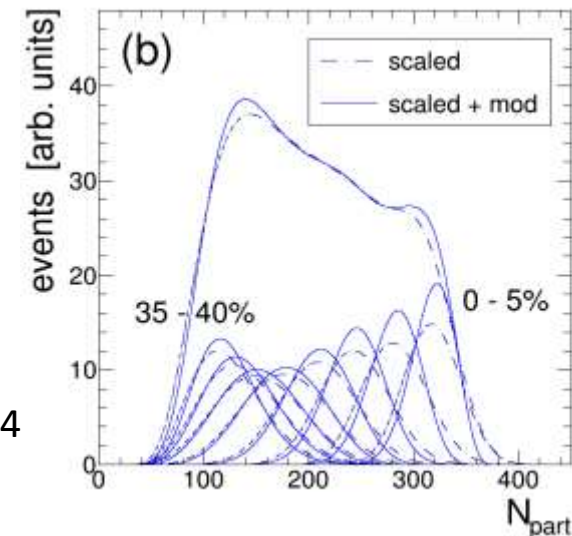
Observed N_{hit} distributions (selected on FW)



fit IQMD cumulants
to adjust scaled v_n
→
morph N_{hit} into N_{part}

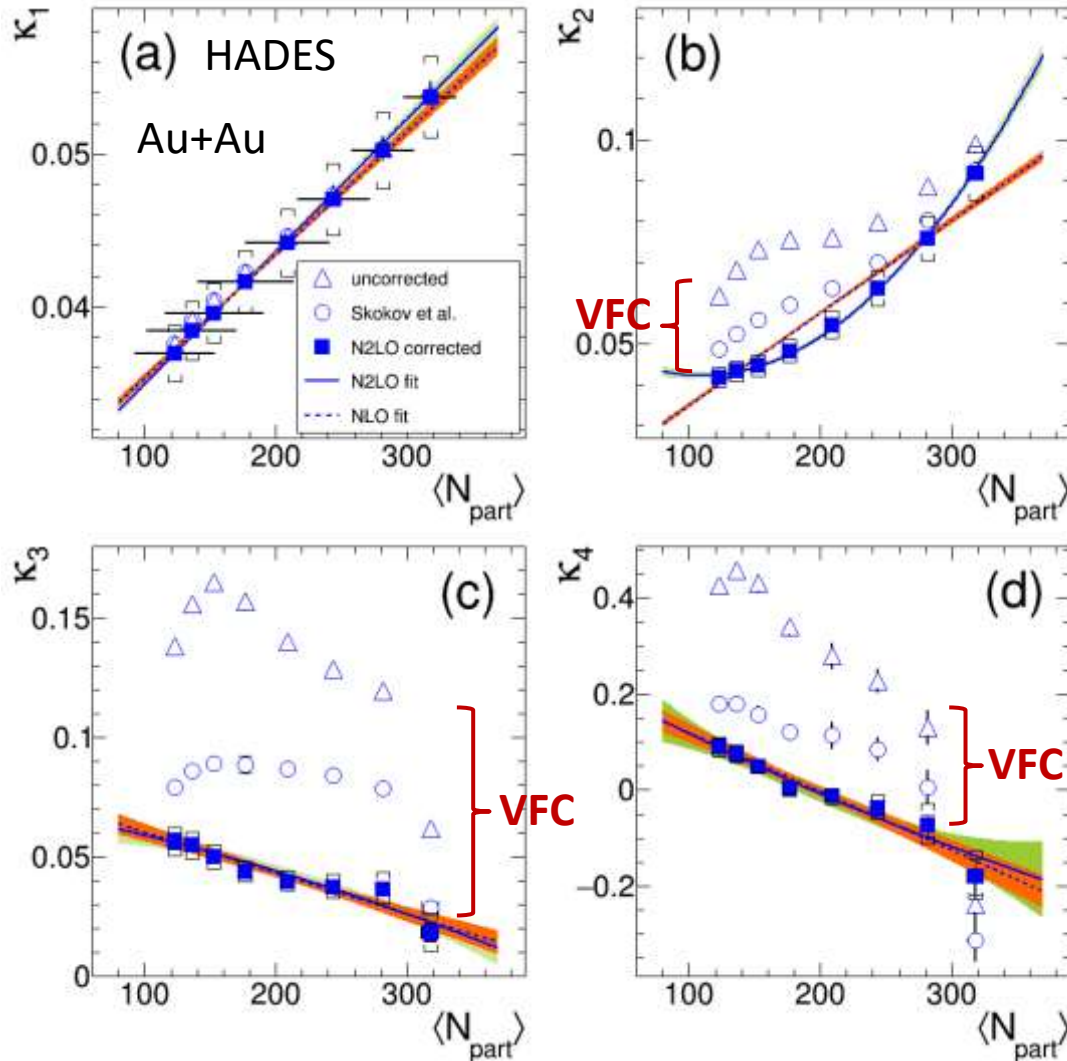
PRC 102 (2020) 024914

Reconstructed N_{part} distributions



Reduced proton cumulants in Au+Au data: $y = y_0 \pm 0.2$

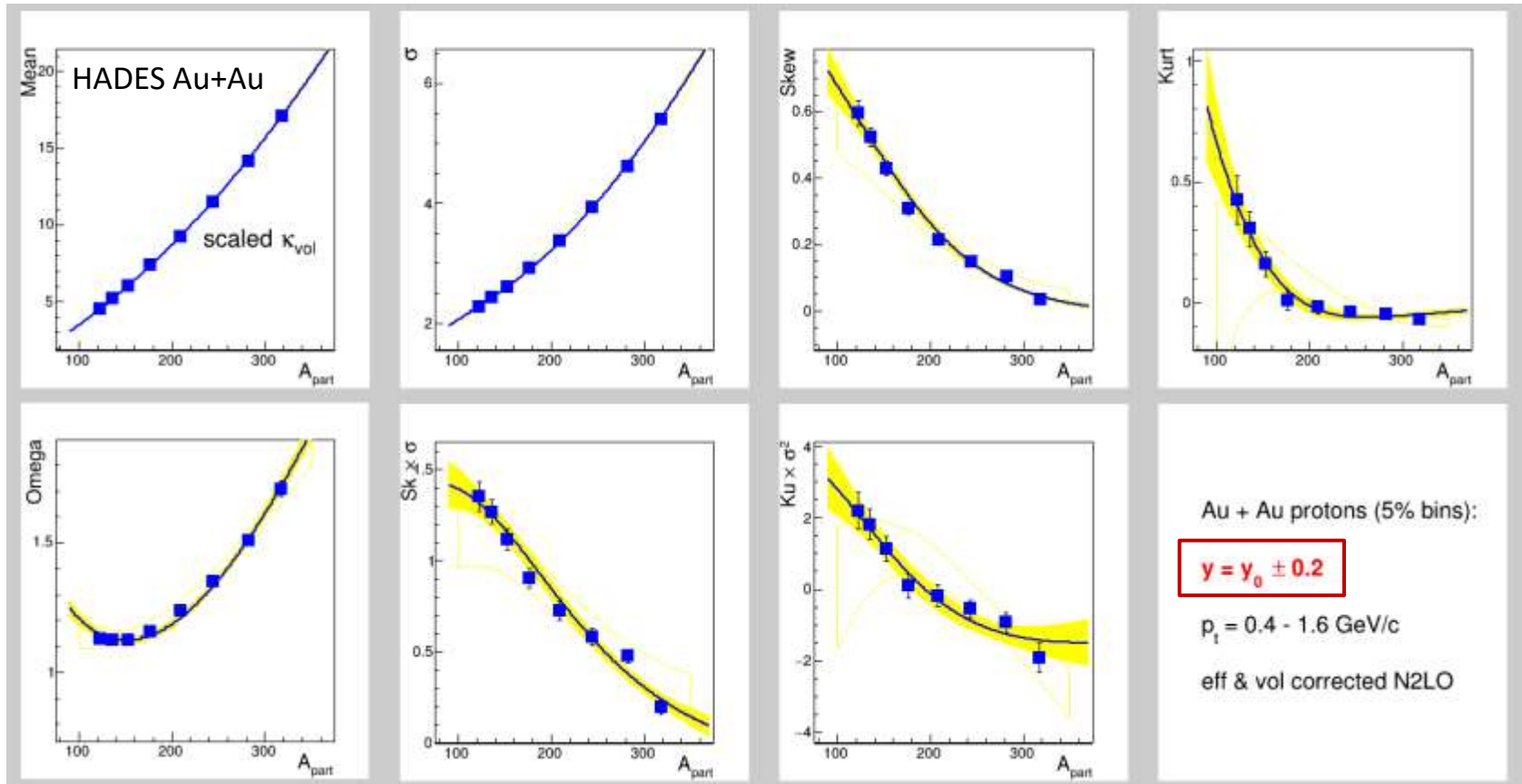
effect of volume fluctuation corrections



- Vol. fluct. Corr. (**VFC**) are important!
- NLO terms are needed!
- Eventually, also N2LO terms

Volume correction of proton nb. moments:

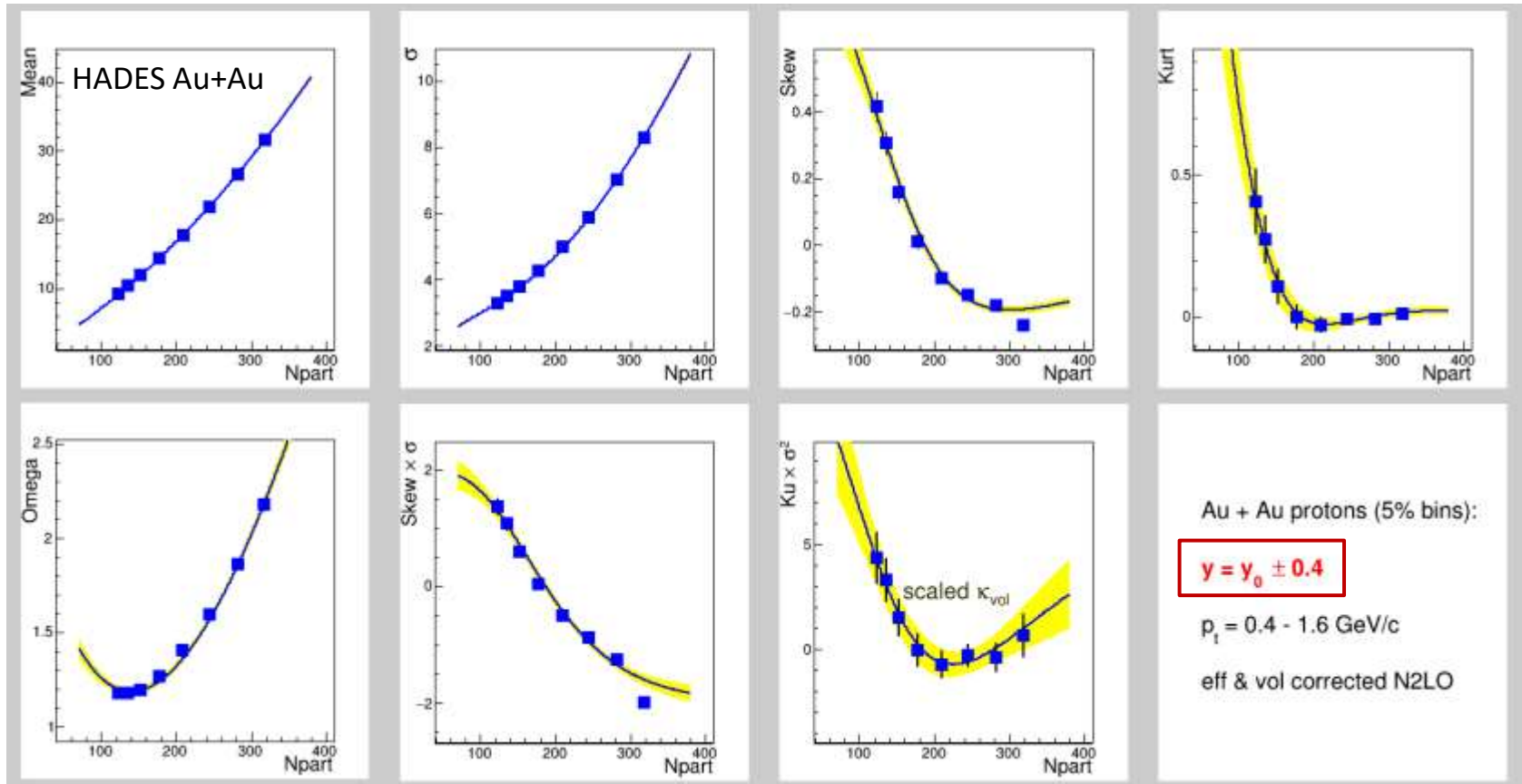
IQMD scaled & adjusted $N_{hit} \kappa_{vol}$



HADES, PRC 102 (2020)

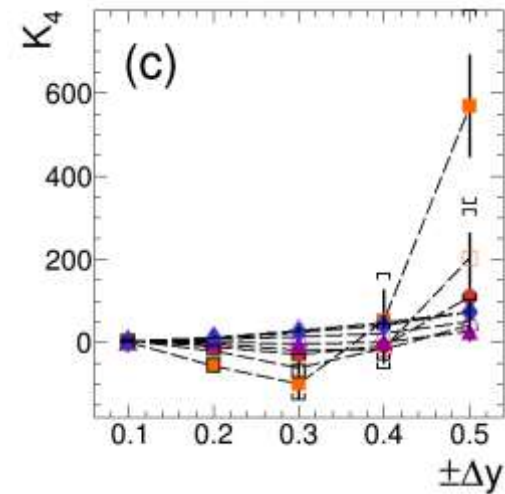
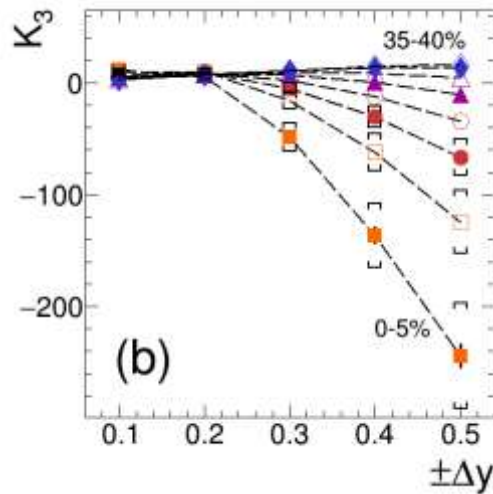
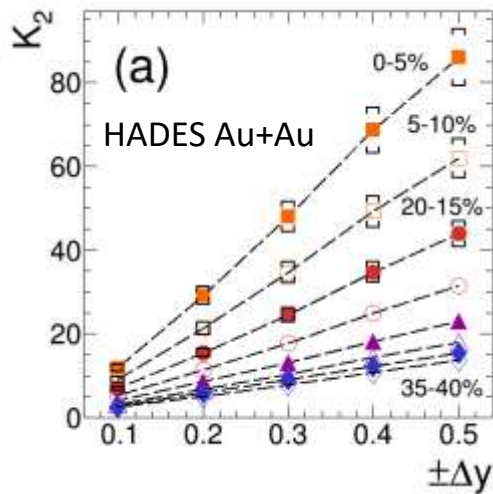
Volume correction of proton nb. moments:

IQMD scaled & adjusted $N_{hit} \kappa_{vol}$

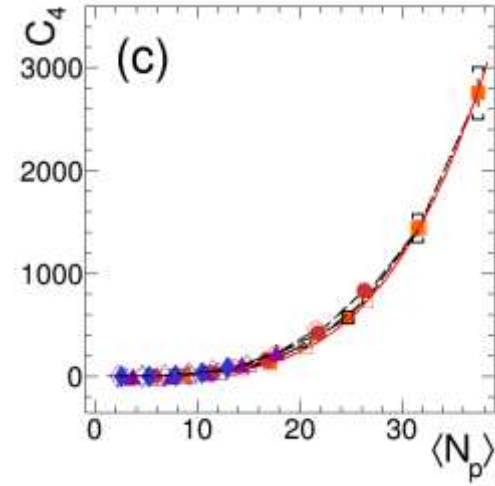
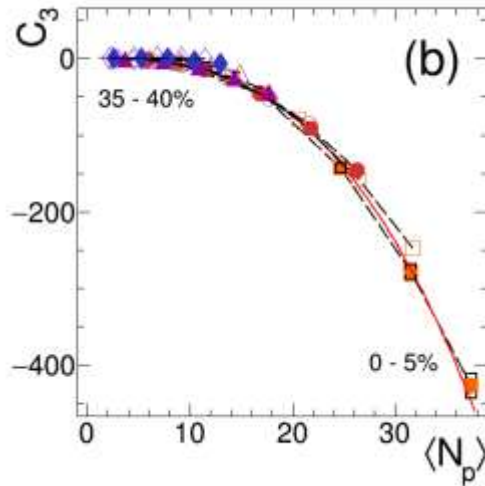
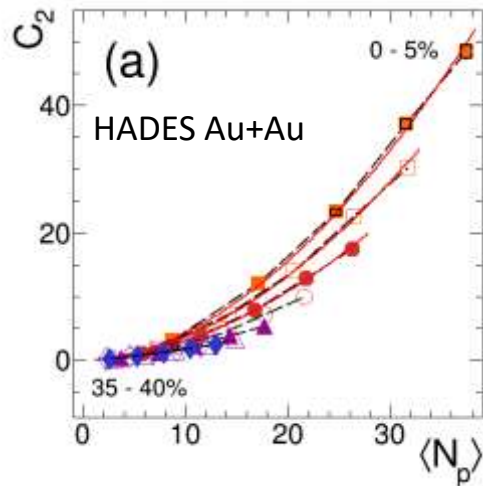


HADES, PRC 102 (2020)

Au+Au data: N2LO corrected cumulants & correlators



eff. & vol.
corrected
proton-nb
cumulants



correlators
a.k.a factorial
cumulants

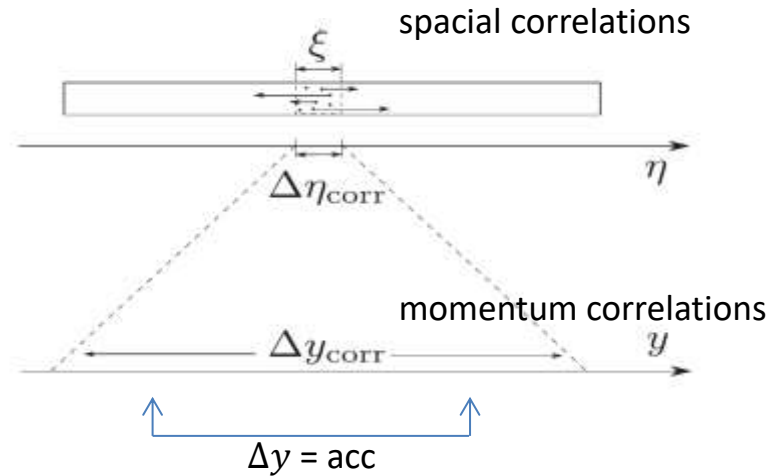
$\langle N \rangle$ scaling of volume-corrected correlators C_k

B. Ling & M. Stephanov PRC 93 (2016)

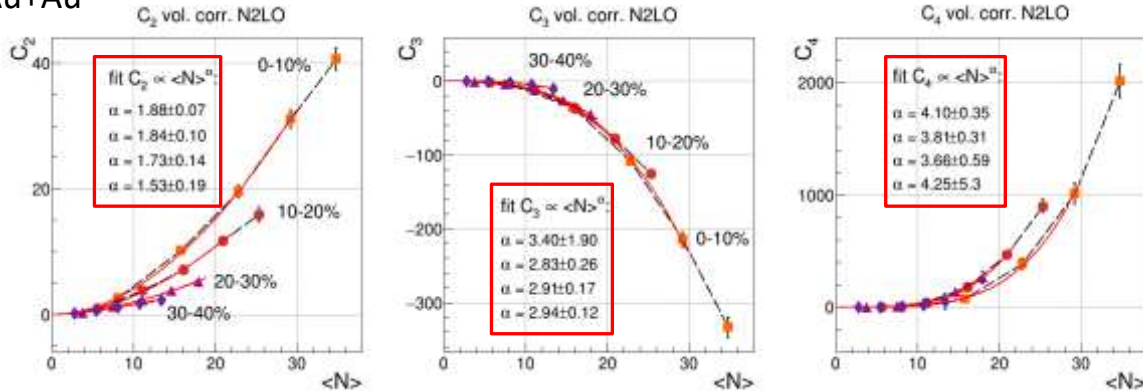
consider two extreme scenarios:

(1) $\Delta y_{\text{corr}} \ll \Delta y \rightarrow C_n \propto \Delta y$

(2) $\Delta y_{\text{corr}} \gg \Delta y \rightarrow C_n \propto (\Delta y)^n$



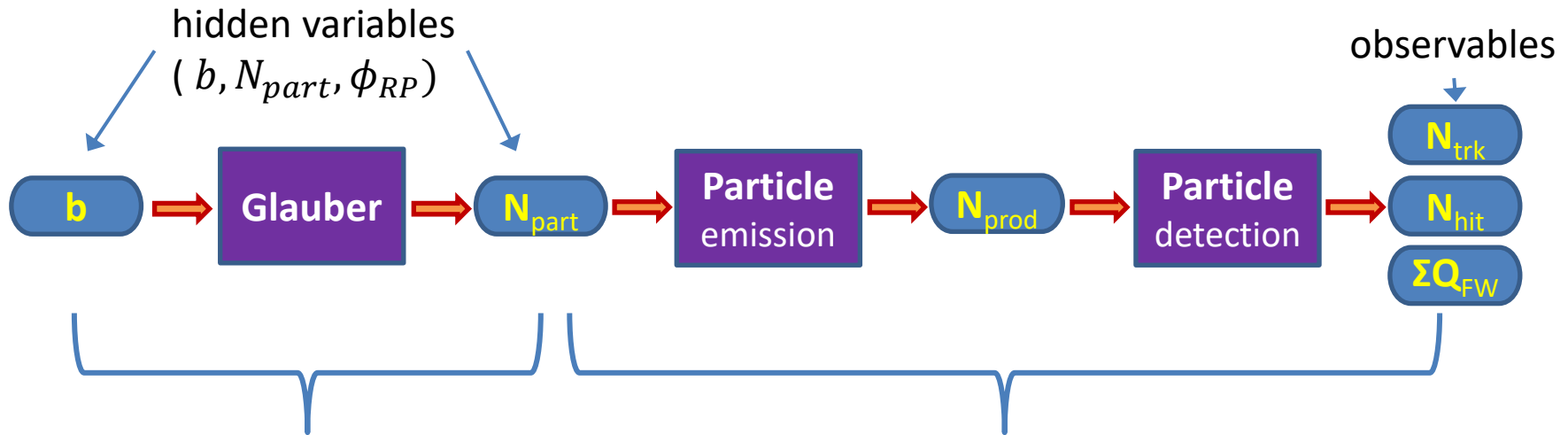
HADES Au+Au



HADES Au+Au
10% centrality bins

→ observed scaling favors long-range correlations ($\Delta y_{\text{corr}} > 1$)

Alternative $N_{hit} \rightarrow N_{part}$ transformations based on a Glauber Monte Carlo



A minimal model:



e.g. $\text{Poisson}(\lambda N_{part}) \times (1 - \alpha N_{part}^2)$

→ compute the $\kappa(N_{part})$ from the $\kappa(N_{hit})$ and expand N_{part}

→ do a Bayesian reconstruction of the N_{part} distribution

True Poissonian process: $N_{hit} = \text{Poisson}(\lambda N_{part})$

Applying **total cumulance** to $X = N_{hit} = \text{Poisson}(X|_{Z=N_{part}})$ we obtain an analytic relation between the cumulants of N_{hit} and N_{part} :

$$\kappa_n[N_{hit}] = \sum_{i=1}^n \lambda^i S_2(n, i) \kappa_i[N_{part}] \quad (1)$$

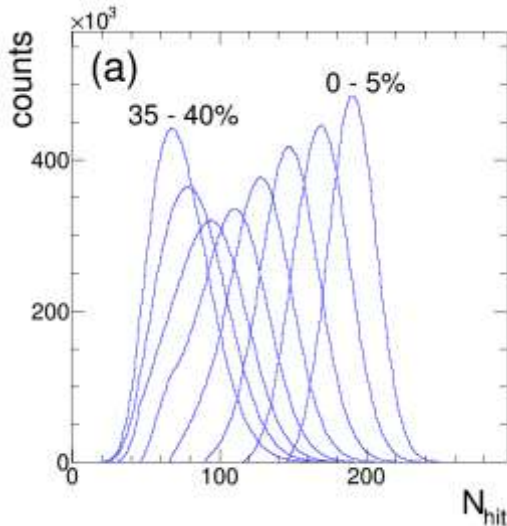
With the inverse:

$$\kappa_n[N_{part}] = \sum_{i=1}^n \lambda^{-i} S_1(n, i) \kappa_i[N_{hit}] \quad (2)$$

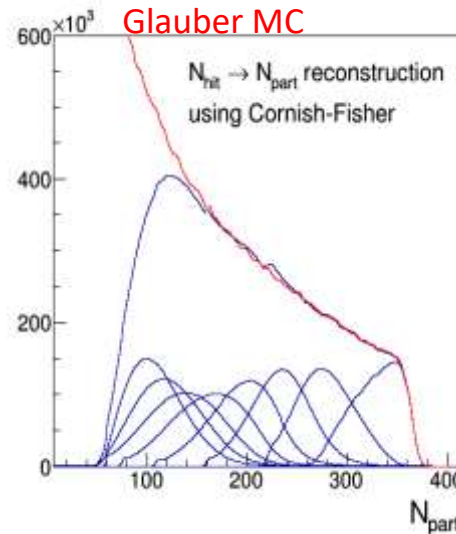
See also
Broniowski & Olszewski, PRC95 (2017)

where S_1 and S_2 are Stirling numbers

Observed Au+Au N_{hit} distributions



Poisson



$N_{hit} \rightarrow N_{part}$ transformation
using total cumulance
+
Cornish-Fisher expansion
of the N_{part} distributions

Bayesian reconstruction of centrality

PHYSICAL REVIEW C **97**, 014905 (2018)

Relating centrality to impact parameter in nucleus-nucleus collisions

Sruthy Jyothi Das,^{1,2} Giuliano Giacalone,¹ Pierre-Arnaury Monard,¹ and Jean-Yves Ollitrault¹

PHYSICAL REVIEW C **98**, 024902 (2018)

Editors' Suggestion

Reconstructing the impact parameter of proton-nucleus and nucleus-nucleus collisions

Rudolph Rogly,^{1,2} Giuliano Giacalone,¹ and Jean-Yves Ollitrault¹

PHYSICAL REVIEW C **104**, 034609 (2021)

Model independent reconstruction of impact parameter distributions for intermediate energy heavy ion collisions

J. D. Frankland,^{1,7} D. Gruyer,² E. Bonnet,³ B. Borderie,⁴ R. Bougault,² A. Chhibi,¹ J. E. Ducret,³ D. Durand,² Q. Fable,² M. Henri,¹ J. Lemarié,¹ N. Le Neindre,² I. Lombardo,⁵ O. Lopez,² L. Manduci,^{2,6} M. Piarog,^{2,7} J. Quicray,⁷ G. Verde,^{3,8} E. Vient,⁷ and M. Vigilante⁹
(INDRA Collaboration)

Ollitrault et al.

- validated with simulations and applied to LHC data

INDRA collab.

- validated with low-energy GANIL data $E_{\text{beam}} < 100 \text{ MeV/u}$

Apply Bayes' theorem \rightarrow

$$P(B|A) = P(A|B) P(B)/P(A)$$

Setting $A = N_{\text{hit}}$, $B = N_{\text{part}}$ with

$P(B|A) \leftrightarrow$ prob of N_{part} for given N_{hit}

$P(A|B) \leftrightarrow$ prob of N_{hit} for given N_{part}

$P(A) \leftrightarrow$ min bias N_{hit} distribution

$P(B) \leftrightarrow$ min bias N_{part} distribution

\leftarrow to be reconstructed

\leftarrow Glauber fit to N_{hit} data

\leftarrow data

\leftarrow Glauber

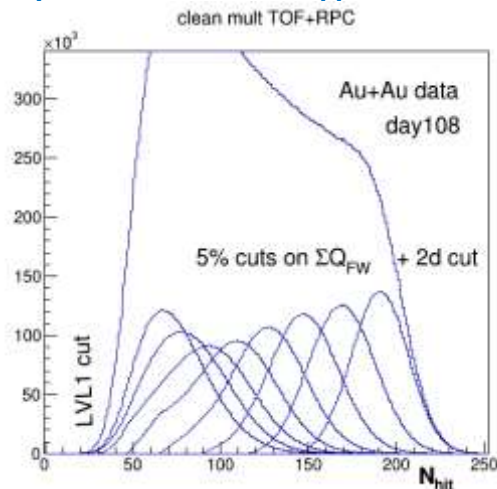
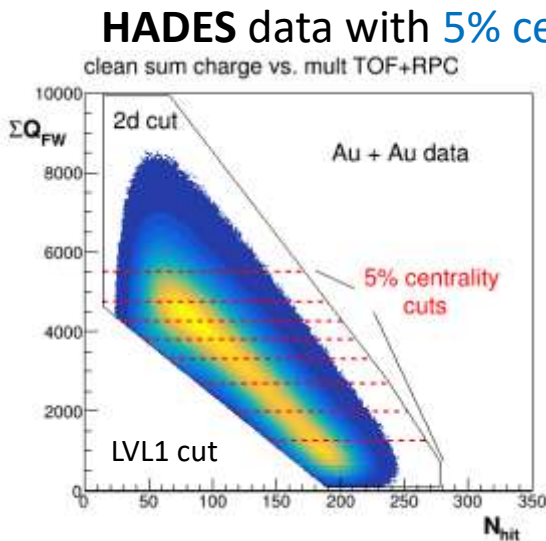
} minimal model (Glauber MC)

Glauber-based Monte Carlo of Au+Au data

for details, see HADES centrality paper EPJA 54 (2018)

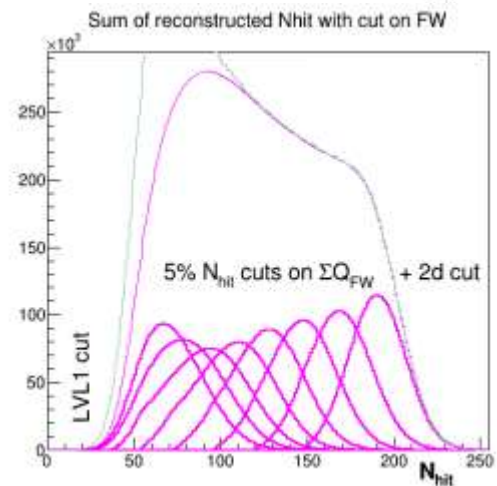
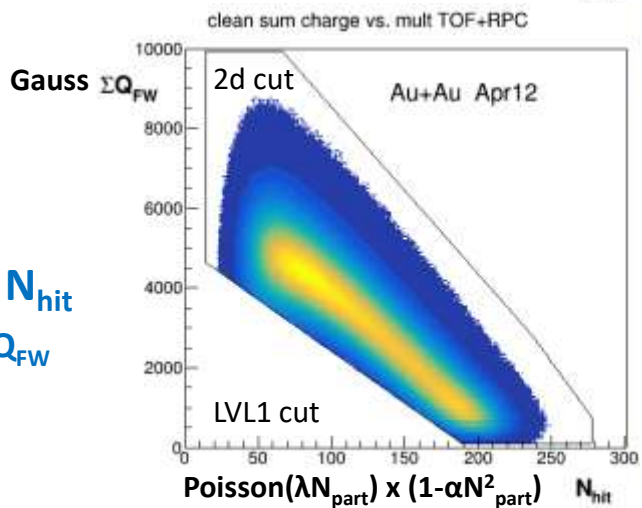
Au+Au data

- LVL1 trigger
- cleaning cuts



Monte Carlo

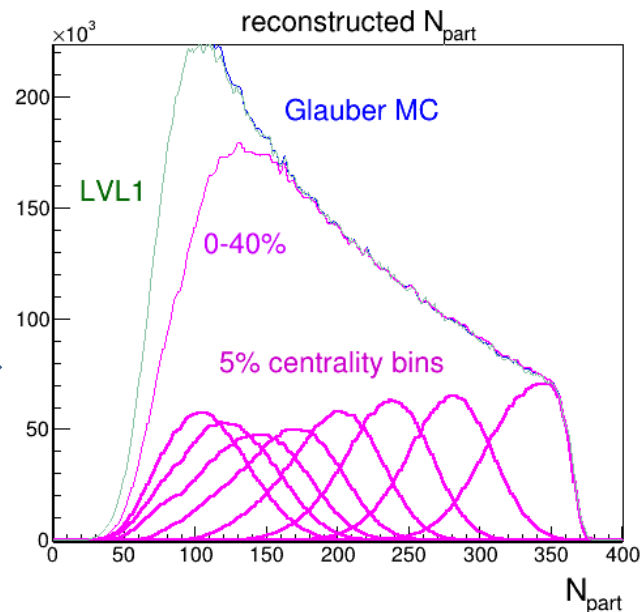
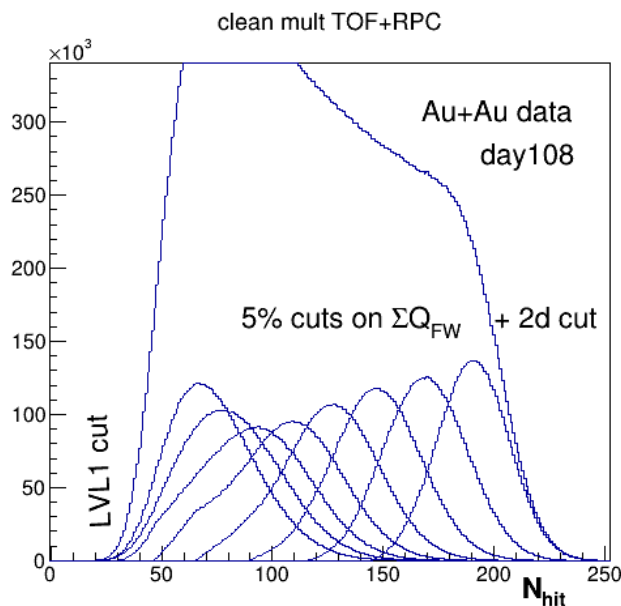
- Glauber
- Poissonian N_{hit}
- Gaussian ΣQ_{FW}



Bayesian reconstruction of centrality

For the $N_{hit} \rightarrow N_{part}$ reconstruction we follow Frankland et al. PRC 104 (2021):

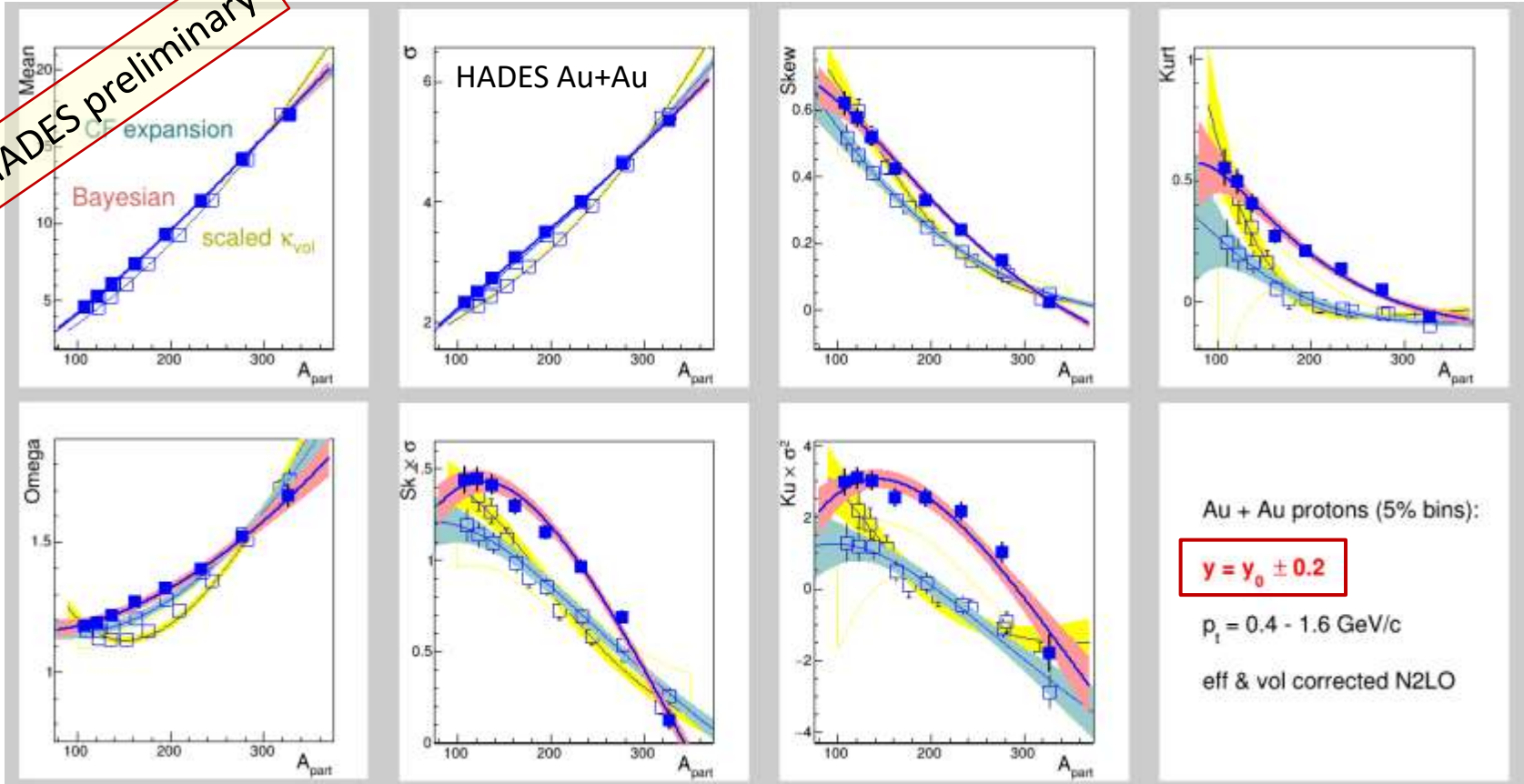
$$P(N_{part}|FW_c) = \frac{\sum_{N_{hit}} P(N_{hit}|N_{part}) P(N_{part}) w(FW_c \& LVL1 \& 2D)}{\sum_{N_{hit}} P(N_{hit}|FW_c)}$$



Volume correction of proton nb. moments:

Bayesian $N_{hit} \rightarrow N_{part}$

HADES preliminary

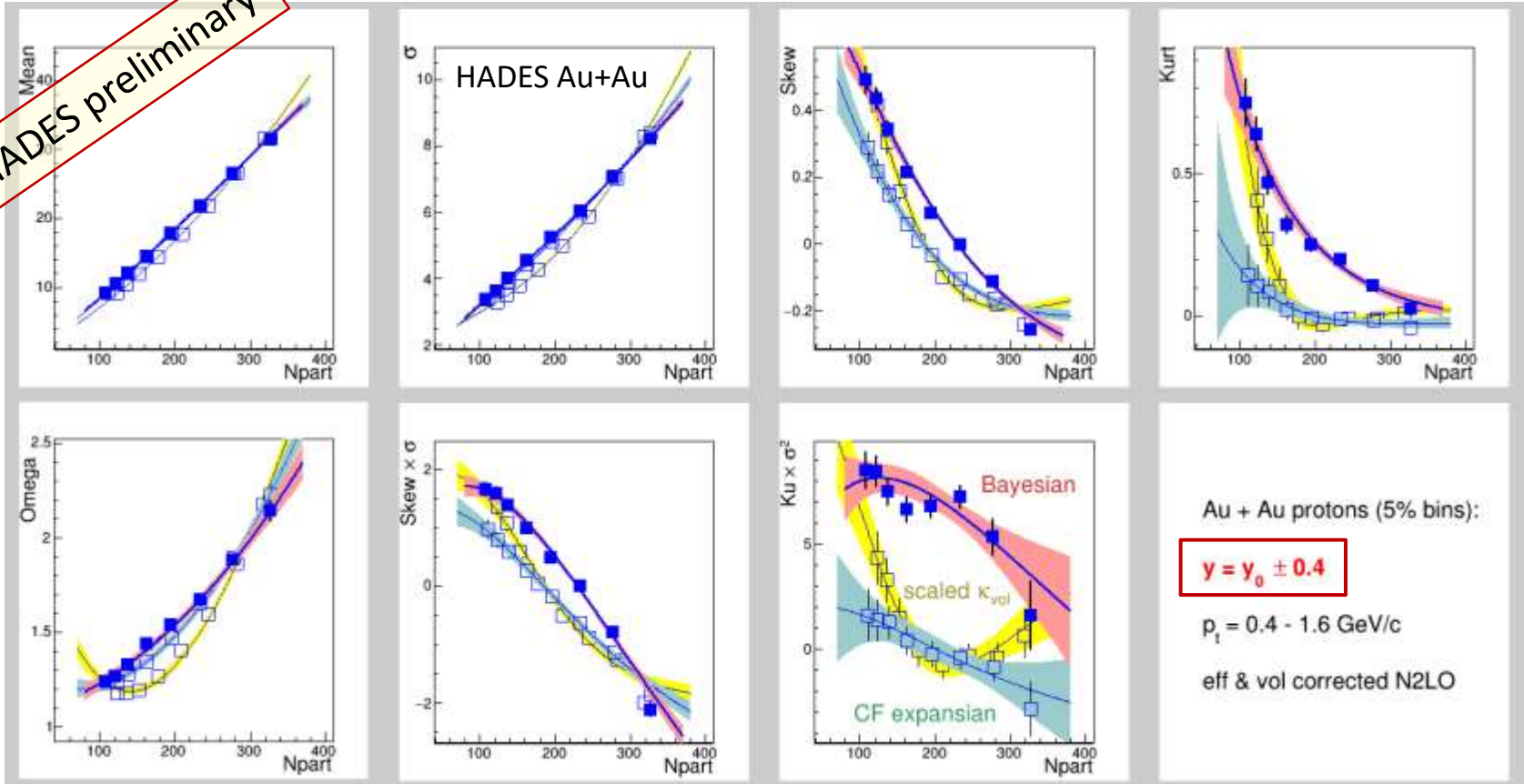


HADES re-analysis November 2021

Volume correction of proton nb. moments:

Bayesian $N_{\text{hit}} \rightarrow N_{\text{part}}$

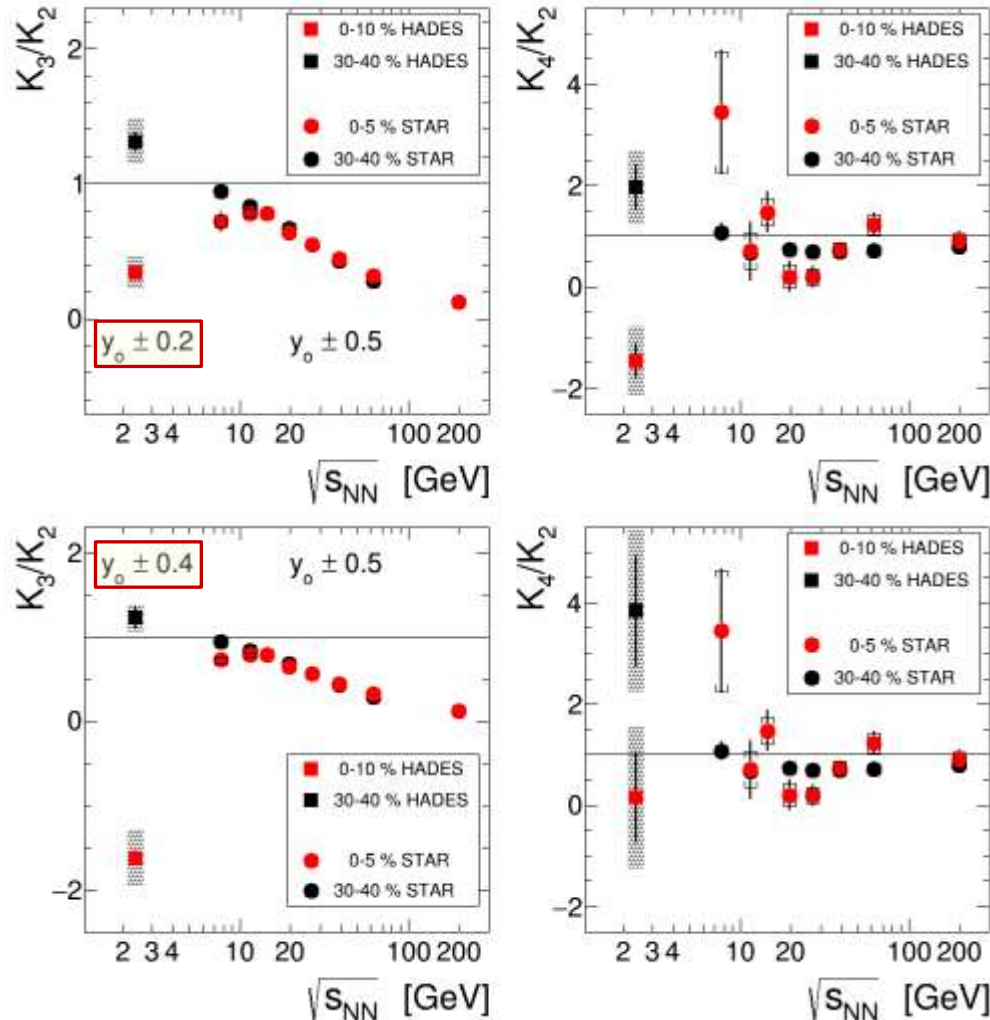
HADES preliminary



HADES re-analysis November 2021

HADES vs STAR: status of 2020

IQMD rescaled N_{hit} (used in HADES PRC 102 paper)

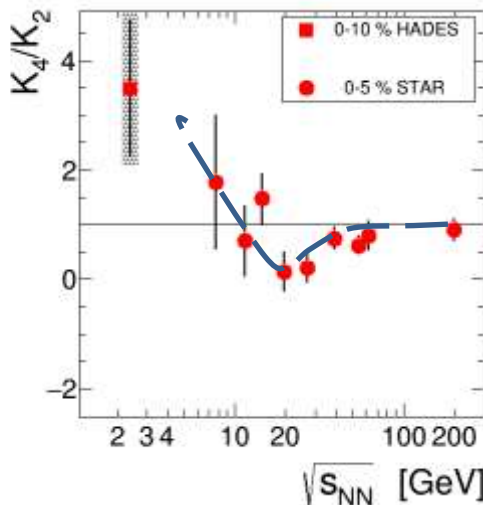
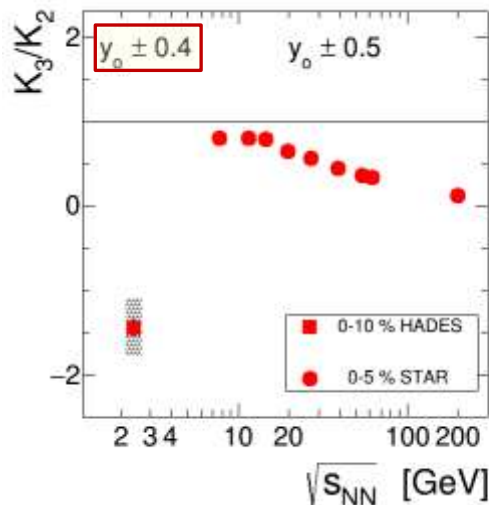
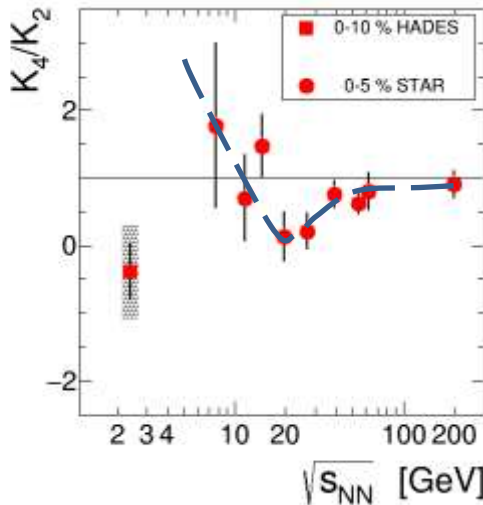
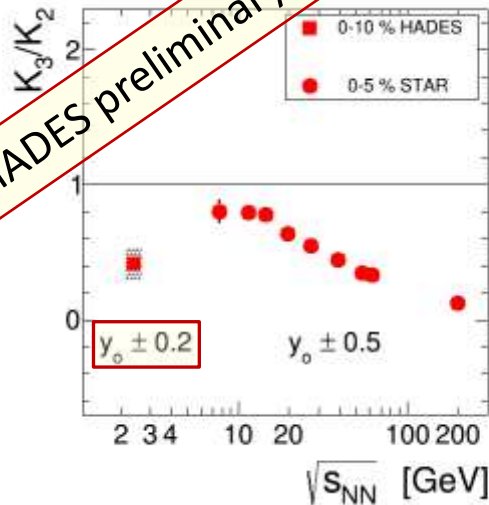


Note that STAR updated their data in 2021!
→ PRL & PRC 104

Comparison with STAR 2021

Bayesian $N_{\text{hit}} \rightarrow N_{\text{part}}$

HADES preliminary



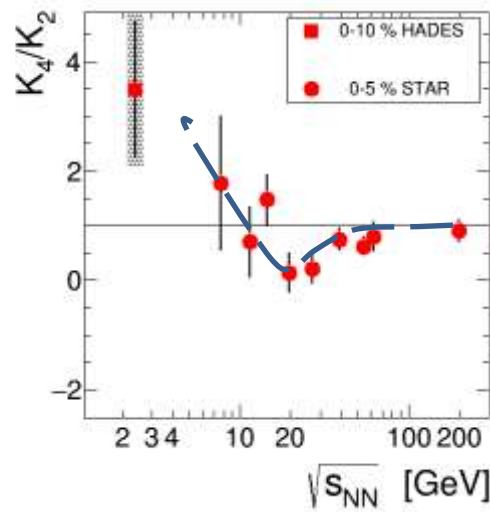
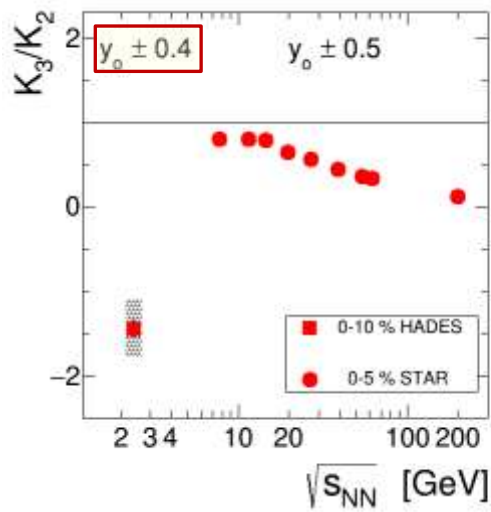
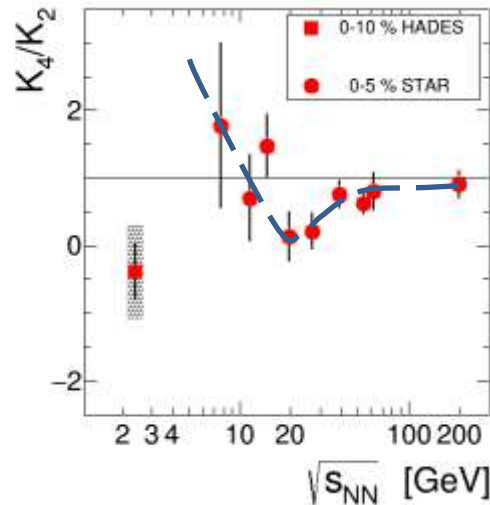
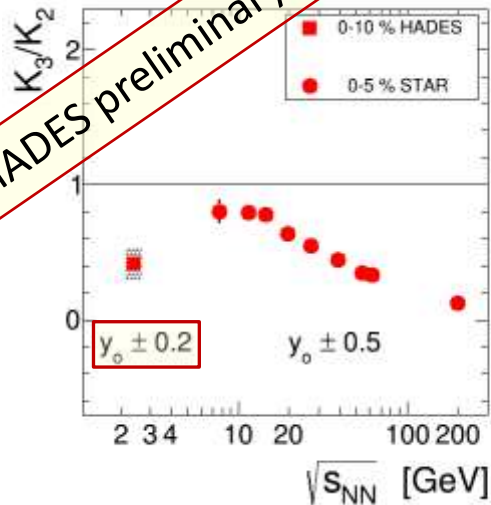
STAR sees a „non-monotonic“ trend of K_4/K_2 with N_{part} which they interpret as a sign of possible critical behavior

PRL (2021), PRC 104 (2021)

Comparison with STAR 2021

Bayesian $N_{\text{hit}} \rightarrow N_{\text{part}}$

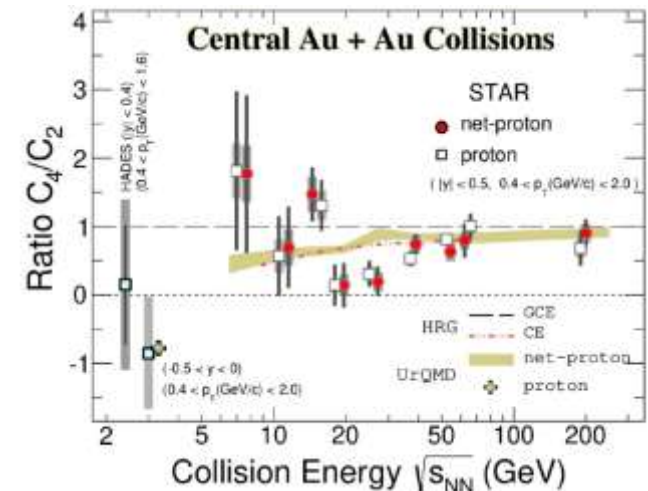
HADES preliminary



STAR sees a „non-monotonic“ trend of K_4/K_2 with N_{part} which they interpret as a sign of possible critical behavior

PRL (2021), PRC 104 (2021)

STAR 3 GeV result:



Hadron Resonance Gas + van der Waals forces

PHYSICAL REVIEW C 98, 024910 (2018)

Critical point of nuclear matter and beam-energy dependence of net-proton number fluctuations

Volodymyr Vovchenko,^{1,2} Lijia Jiang,² Mark I. Gorenstein,^{2,3} and Horst Stoecker^{1,2,4}

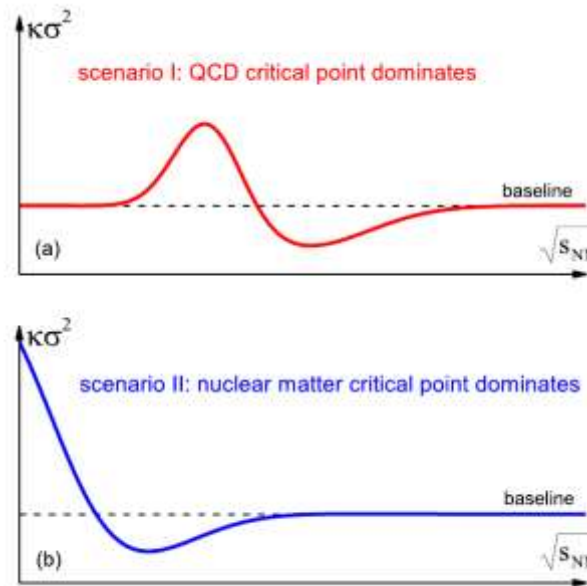
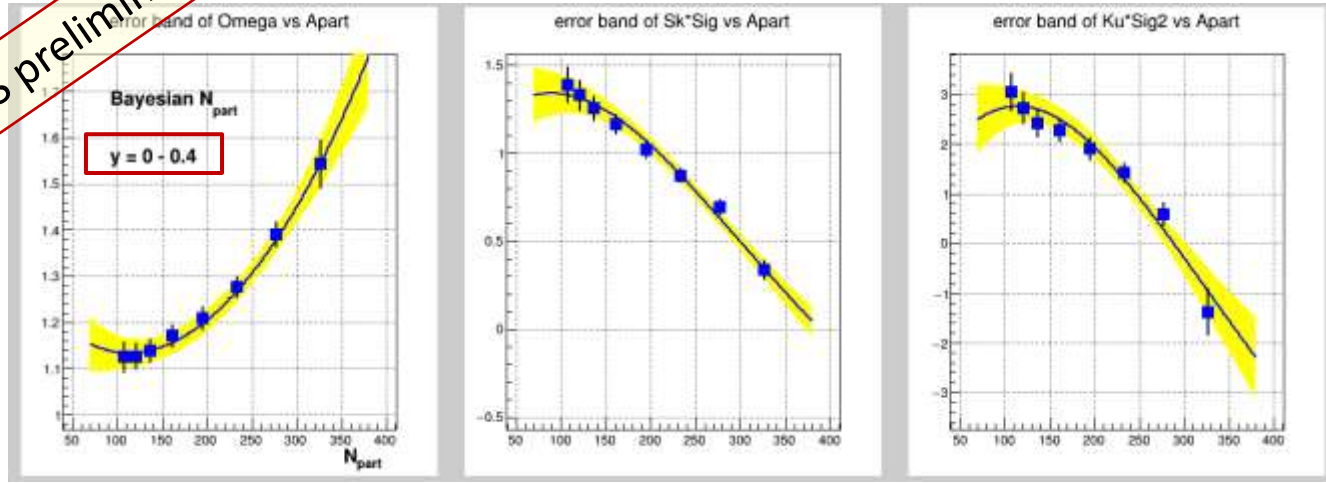


Fig. 3

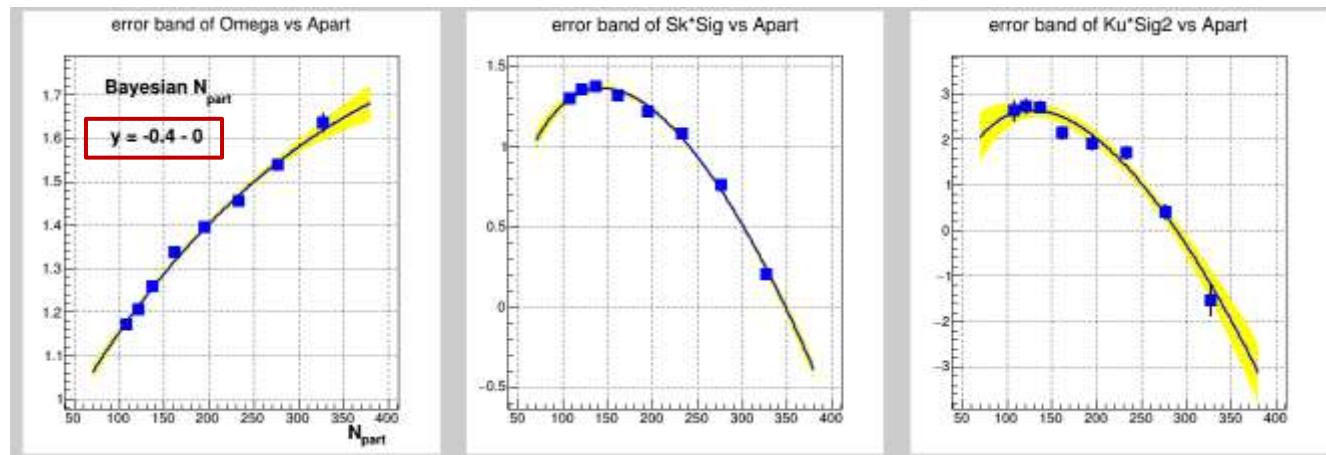
Proton cumulants: forward vs backward rapidities

Bayesian $N_{hit} \rightarrow N_{part}$

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HADES Au+Au
forward rapidities



HADES Au+Au
backward rapidities

HADES re-analysis November 2021

→ Systematics!

Summary and Outlook:

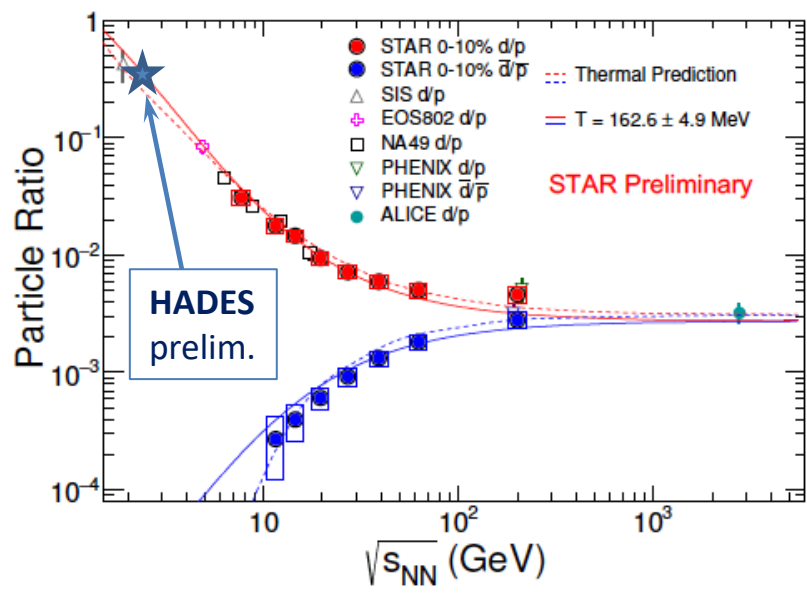
- Analyzed proton number fluctuations with HADES in 1.23 AGeV Au+Au collisions
- Applied full (N2LO) volume corrections to observed proton cumulants using different $N_{\text{hit}} \rightarrow N_{\text{part}}$ transforms
- Found indications for strong long-range correlations (>1 in rapidity)
 - absolutely need to control volume corrections!
 - remnant effects of liquid-gas phase transition?
 - bound vs free protons?

 - very high-statistics Ag+Ag data are available for analysis
 - beam-energy scan <1 AGeV (SIS18) & >2 AGeV (SIS100)

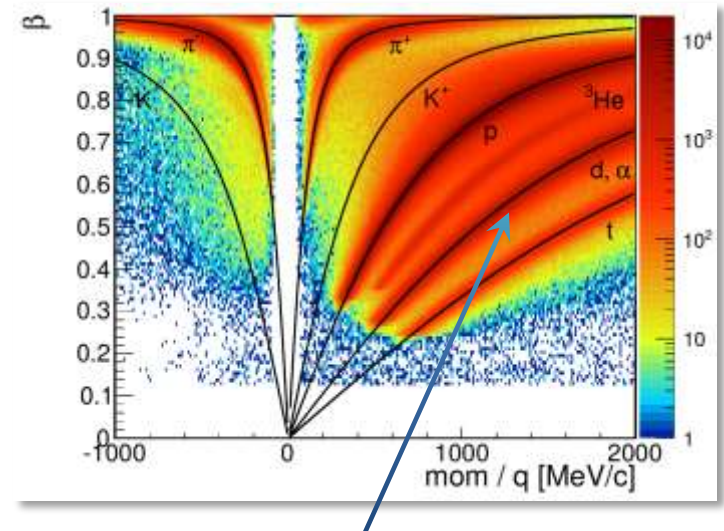
Extra slides

How do bound protons contribute?

Systematics of d/p from STAR collaboration (QM2017)



HADES 1.23 GeV/u Au+Au data



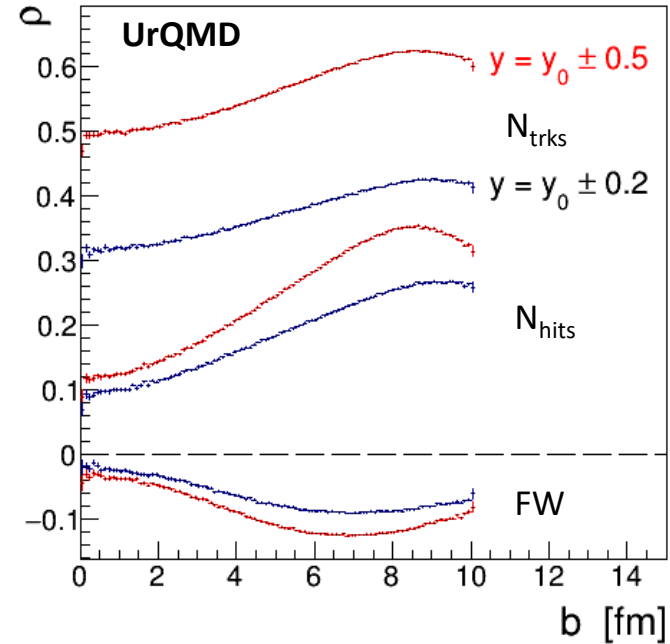
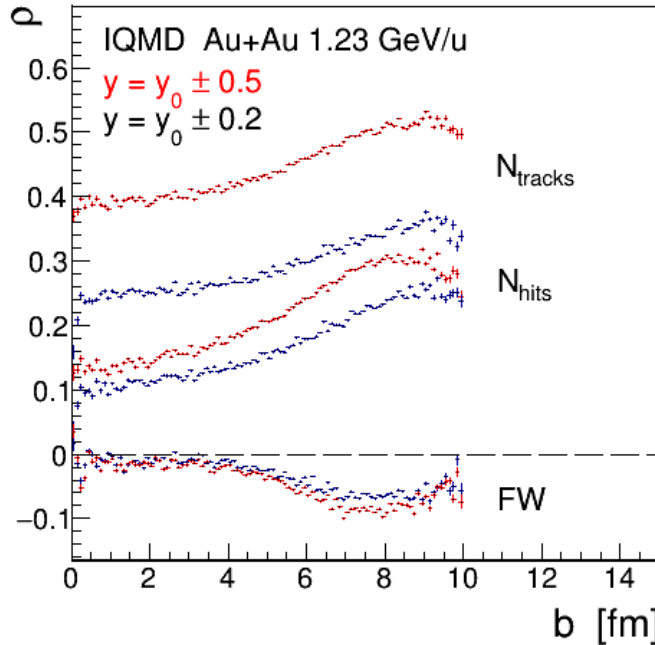
$d/p \approx 0.37$ in 0-10% most central

Large fraction of protons are bound in nuclei: d, t, He, etc.

How do they contribute to proton-number fluctuations?

- ➔ Investigation of light nucleus production in Au+Au is ongoing (M. Szala & M. Lorenz)

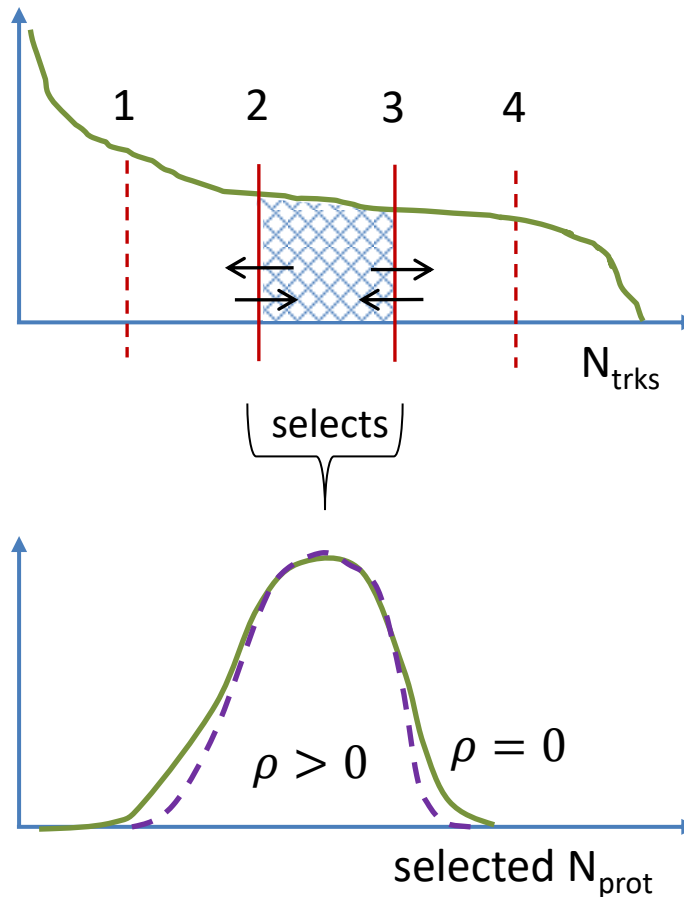
Correlations of N_{prot} & centrality selection



Pearson's linear correlation coefficient: $\rho_{xy} = \frac{\text{cov}(x,y)}{\sqrt{\text{var}(x) \text{var}(y)}}$

→ Non-zero correlations have a damping effect volume fluctuations!

Correlations & volume fluctuations



A centrality selection on nb. of tracks is problematic

If N_{trks} and N_{prot} are correlated, a cut on N_{trks} will reduce the width of the selected proton distribution

Volume fluctuation cumulants v_n are modified by correlations!

Extended (NLO) Skokov volume corrections

$$c_1 = \kappa_1 + \kappa'_1 v_2 \quad (1)$$

$$c_2 = \kappa_2 + \kappa_1^2 v_2 + \kappa'_2 v_2 + 2\kappa_1 \kappa'_1 V_2 + 2\kappa_1 \kappa'_1 v_3 \\ + 2\kappa_1'^2 v_2 V_2 + \kappa_1'^2 V_1 V_2 + 2\kappa_1'^2 V_3 + \kappa_1'^2 v_4 \quad (2)$$

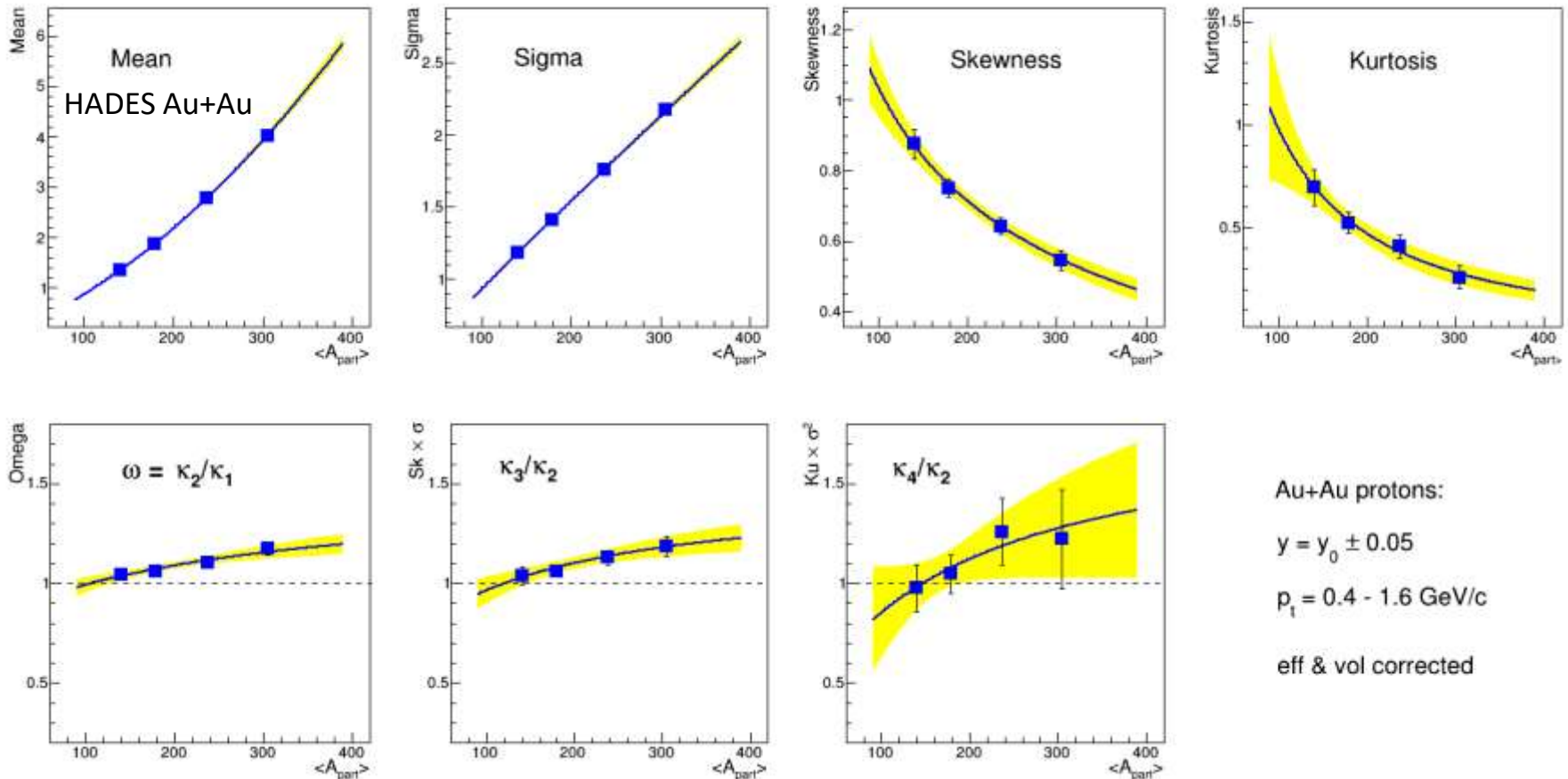
$$c_3 = \kappa_3 + \kappa_1^3 v_3 + 3\kappa_1 \kappa_2 v_2 + 3(\kappa_1 \kappa'_2 + \kappa'_1 \kappa_2) v_3 + 6\kappa_1' (\kappa_1^2 + \kappa_2') v_2 V_2 \\ + 3\kappa_1' (\kappa_1^2 + 2\kappa_2') V_3 + 3\kappa_1' (\kappa_1^2 + \kappa_2') v_4 + 12\kappa_1 \kappa_1'^2 V_2^2 + 3\kappa_1 \kappa_1'^2 V_1 V_3 \\ + 24\kappa_1 \kappa_1'^2 v_2 V_3 + 6\kappa_1 \kappa_1'^2 V_4 + 3\kappa_1 \kappa_1'^2 v_5 + 3(\kappa_1 \kappa'_2 + \kappa'_1 \kappa_2) V_2 \\ + 8\kappa_1'^3 v_2 V_2^2 + 6\kappa_1'^3 V_1 V_2^2 + 10\kappa_1'^3 v_3 V_3 + \kappa_1'^3 V_1^2 V_3 + 24V_2 V_3 \kappa_1'^3 \\ + 3\kappa_1'^3 V_1 V_4 + 12\kappa_1'^3 v_2 V_4 + 3\kappa_1'^3 V_5 + \kappa_1'^3 v_6 + 3\kappa_1' \kappa_2' V_1 V_2 + \kappa_3' v_2 \quad (3)$$

$$c_4 = \kappa_4 + \kappa_1^4 v_4 + 6\kappa_1^2 \kappa_2 v_3 + (4\kappa_1 \kappa_3 + 3\kappa_2^2) v_2 + 24(\kappa_1^3 \kappa_1' + 4\kappa_1 \kappa_1' \kappa_2' + 2\kappa_1'^2 \kappa_2) v_2 V_3 \\ + 4(\kappa_1^3 \kappa_1' + 6\kappa_1 \kappa_1' \kappa_2' + 3\kappa_1'^2 \kappa_2) V_4 + 2(2\kappa_1^3 \kappa_1' + 6\kappa_1 \kappa_1' \kappa_2' + 3\kappa_1'^2 \kappa_2) v_5 \\ + 48(\kappa_1^2 \kappa_1'^2 + \kappa_1'^2 \kappa_2') v_2 V_2^2 + 12(4\kappa_1^2 \kappa_1'^2 + 5\kappa_1'^2 \kappa_2') v_3 V_3 \\ + 72(\kappa_1^2 \kappa_1'^2 + 2\kappa_1'^2 \kappa_2') V_2 V_3 + 6(\kappa_1^2 \kappa_1'^2 + 3\kappa_1'^2 \kappa_2') V_1 V_4 + 72(\kappa_1^2 \kappa_1'^2 + \kappa_1'^2 \kappa_2') v_2 V_4 \\ + 6(2\kappa_1^2 \kappa_1'^2 + 3\kappa_1'^2 \kappa_2') V_5 + 6(\kappa_1^2 \kappa_1'^2 + \kappa_1'^2 \kappa_2') v_6 \\ + 2(6\kappa_1^2 \kappa_2' + 12\kappa_1 \kappa_1' \kappa_2 + 4\kappa_1' \kappa_3' + 3\kappa_2'^2) v_2 V_2 \\ + 2(3\kappa_1^2 \kappa_2' + 6\kappa_1 \kappa_1' \kappa_2 + 4\kappa_1' \kappa_3' + 3\kappa_2'^2) V_3 + 2(3\kappa_1^2 \kappa_2 + 2\kappa_1 \kappa_3' + 2\kappa_1' \kappa_3 + 3\kappa_2 \kappa_2') v_3 \\ + (6\kappa_1^2 \kappa_2' + 12\kappa_1 \kappa_1' \kappa_2 + 4\kappa_1' \kappa_3' + 3\kappa_2'^2) v_4 + 96\kappa_1 \kappa_1'^3 V_2^3 + 96\kappa_1 \kappa_1'^3 V_3^2 \\ + 288\kappa_1 \kappa_1'^3 v_3 V_2^2 + 72\kappa_1 \kappa_1'^3 V_1 V_2 V_3 + 4\kappa_1 \kappa_1'^3 V_1^2 V_4 + 144\kappa_1 \kappa_1'^3 V_2 V_4 \\ + 128\kappa_1 \kappa_1'^3 v_3 V_4 + 12\kappa_1 \kappa_1'^3 V_1 V_5 + 72\kappa_1 \kappa_1'^3 v_2 V_5 + 12\kappa_1 \kappa_1'^3 V_6 + 4\kappa_1 \kappa_1'^3 v_7 \\ + 24(2\kappa_1 \kappa_1' \kappa_2' + \kappa_1'^2 \kappa_2) V_2^2 + 6(2\kappa_1 \kappa_1' \kappa_2' + \kappa_1'^2 \kappa_2) V_1 V_3 \\ + 2(2\kappa_1 \kappa_3' + 2\kappa_1' \kappa_3 + 3\kappa_2 \kappa_2') V_2 + 48\kappa_1^4 v_2 V_2^3 + 48\kappa_1^4 V_1 V_2^3 + 48\kappa_1^4 V_1 V_3^2 \\ + 240\kappa_1^4 v_2 V_3^2 + 32\kappa_1^4 v_4 V_4 + 288\kappa_1^4 V_2^2 V_3 + 24\kappa_1^4 V_1^2 V_2 V_3 + \kappa_1^4 V_1^3 V_4 \\ + 144\kappa_1^4 v_4 V_2^2 + 72\kappa_1^4 V_1 V_2 V_4 + 128\kappa_1^4 V_3 V_4 + 4\kappa_1^4 V_1^2 V_5 + 72\kappa_1^4 V_2 V_5 \\ + 56\kappa_1^4 v_3 V_5 + 6\kappa_1^4 V_1 V_6 + 24V_2 V_6 \kappa_1^4 v_2 V_6 + 4\kappa_1^4 V_7 + \kappa_1^4 v_8 + 36\kappa_1^2 \kappa_2' V_1 V_2^2 \\ + 6\kappa_1'^2 \kappa_2' V_1^2 V_3 + 4\kappa_1' \kappa_3' V_1 V_2 + 3\kappa_2'^2 V_1 V_2 + \kappa_4' v_2 \quad (4)$$

Where the v_n are reduced volume, i.e. N_{part} , cumulants, and the κ_n and κ_n' are obtained by fitting eq. (1) - (4) to the measured proton cumulants c_n

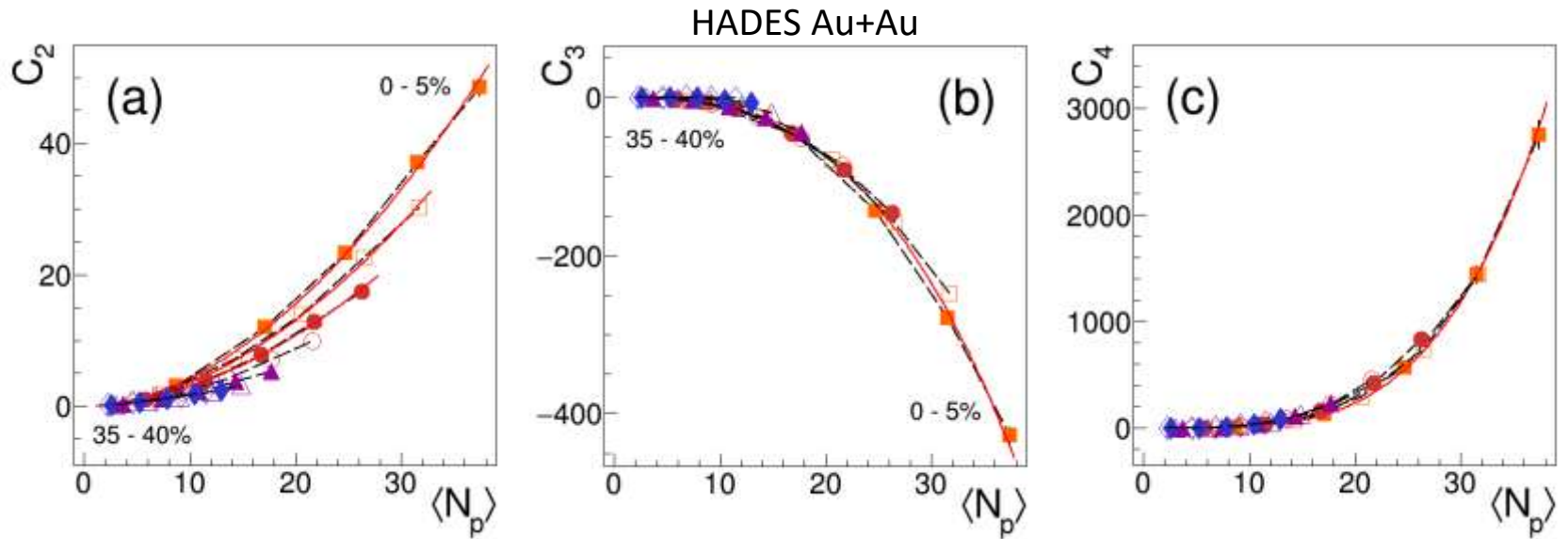
Note that v_n terms appear now up to 8th order!

NLO corr. proton moments in Au+Au data: $y = y_0 \pm 0.05$



→ approaching Poisson limit
in narrow phase-space bin!

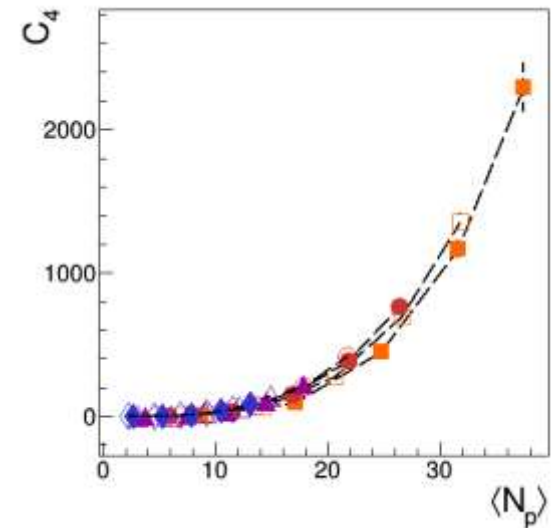
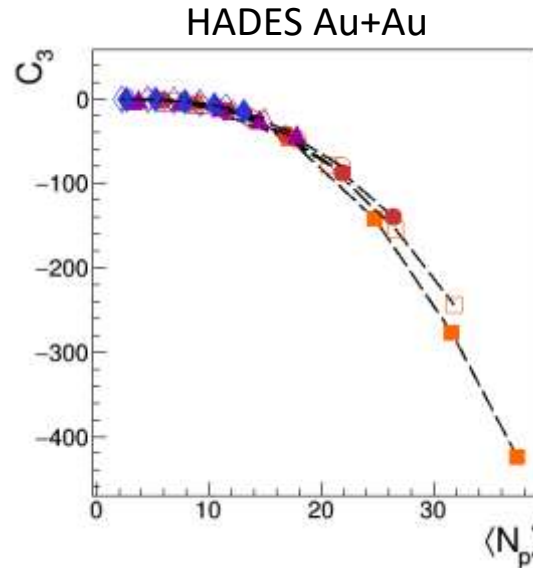
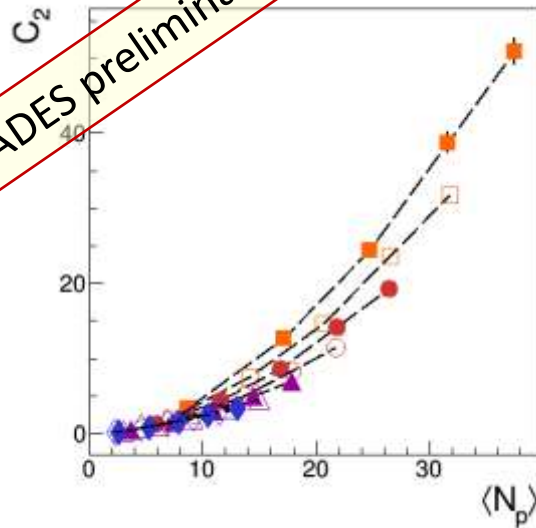
Re-analysis of the proton fluctuations: proton correlators (IQMD based N_{part})



As published in PRC 102 (2020)

Re-analysis of the proton fluctuations: proton correlators (Poisson+CF of N_{part})

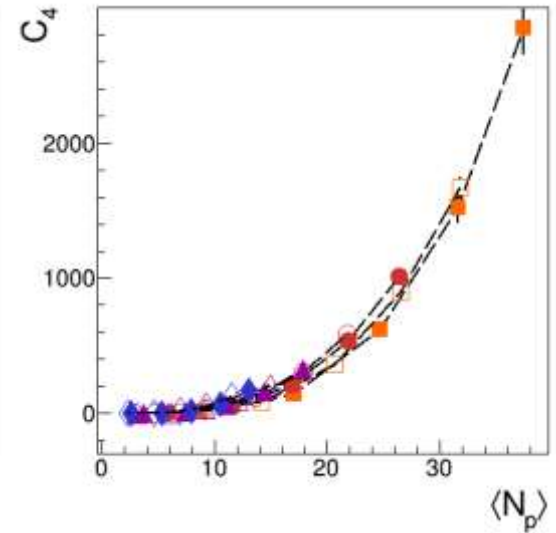
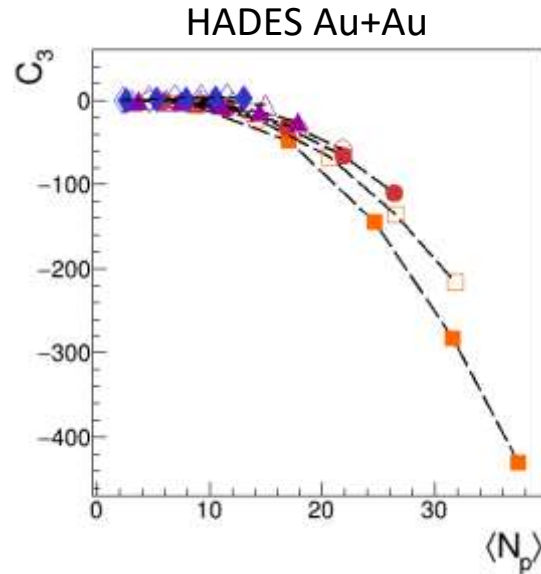
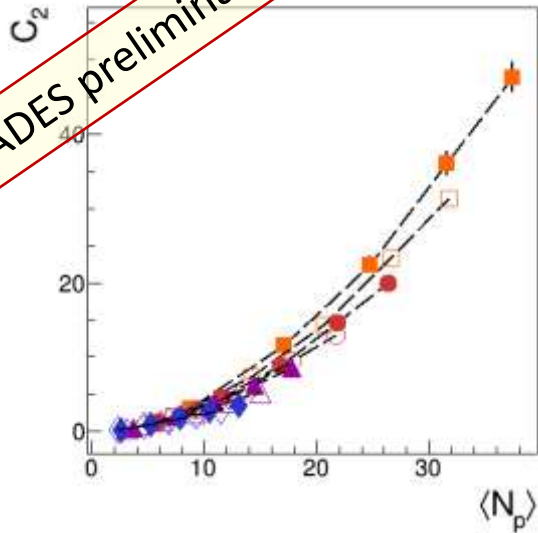
HADES preliminary



The n-particle correlators are very robust!

Re-analysis of the proton fluctuations: proton correlators (Bayesian N_{part})

HADES preliminary



The n-particle correlators are very robust!!