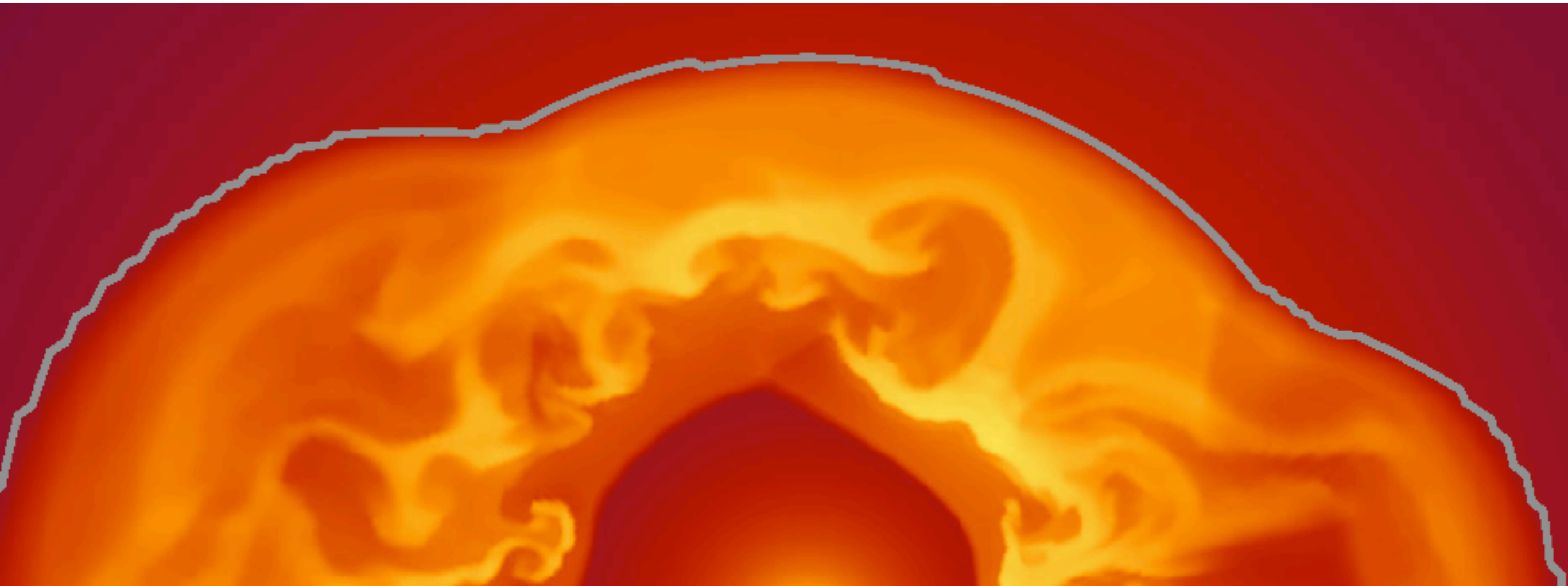
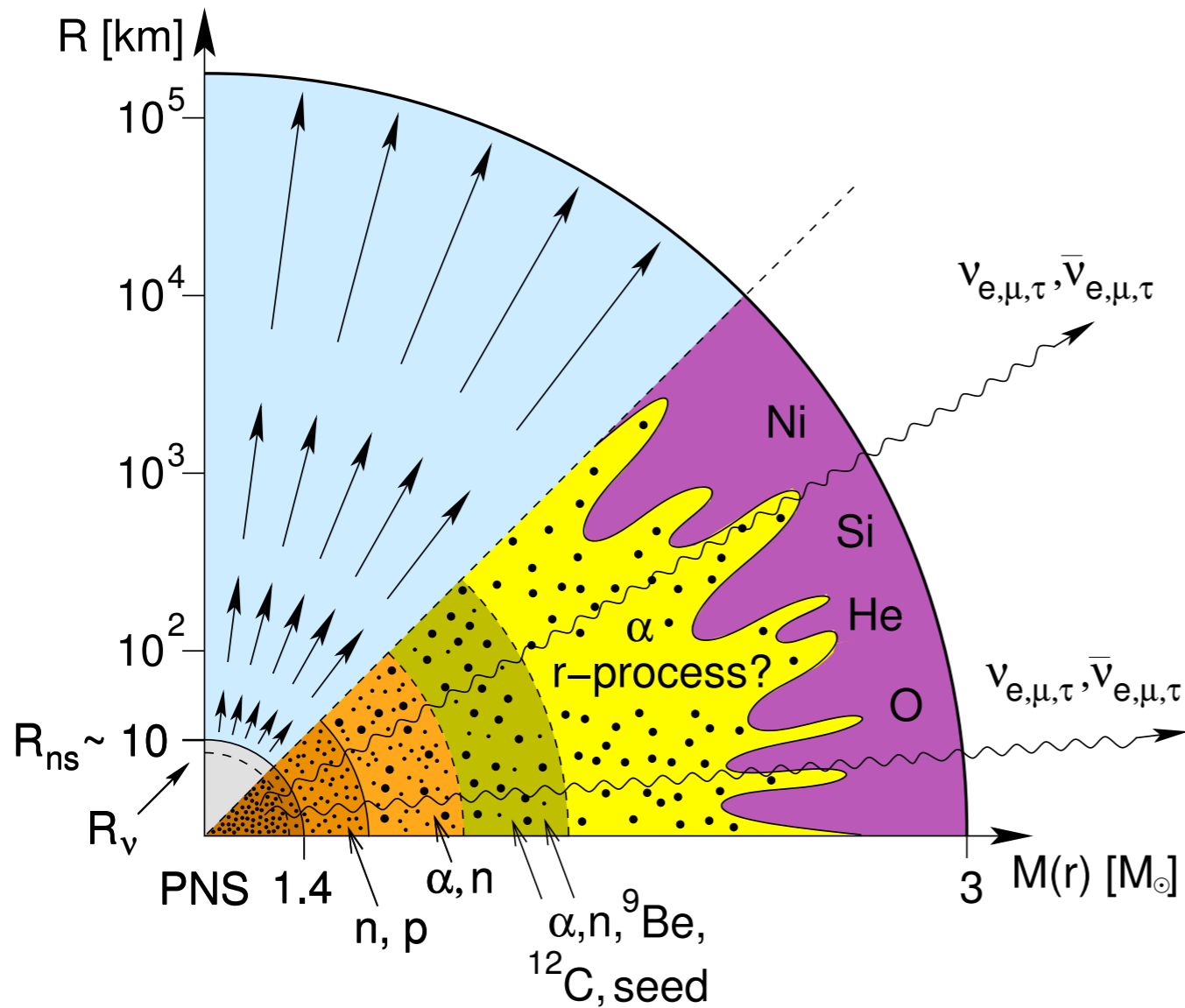


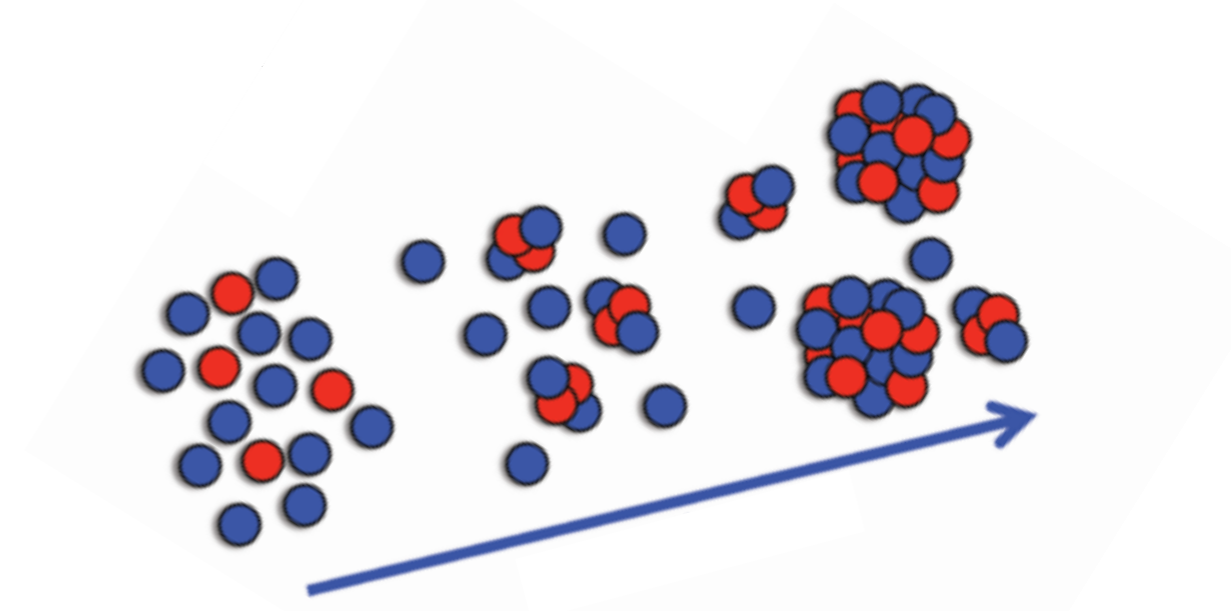
# Lighter element primary process in neutrino-driven winds



# Neutrino-driven winds



neutrons and protons form alpha particles  
alpha particles recombine into seed nuclei

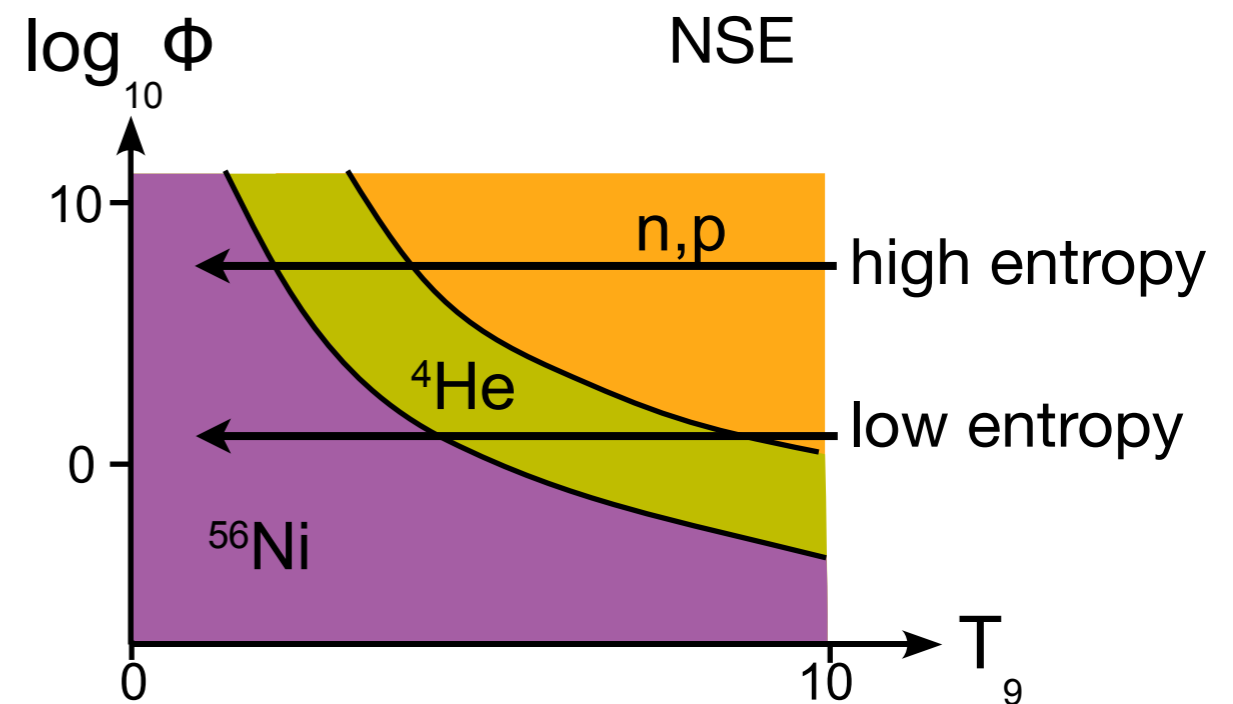
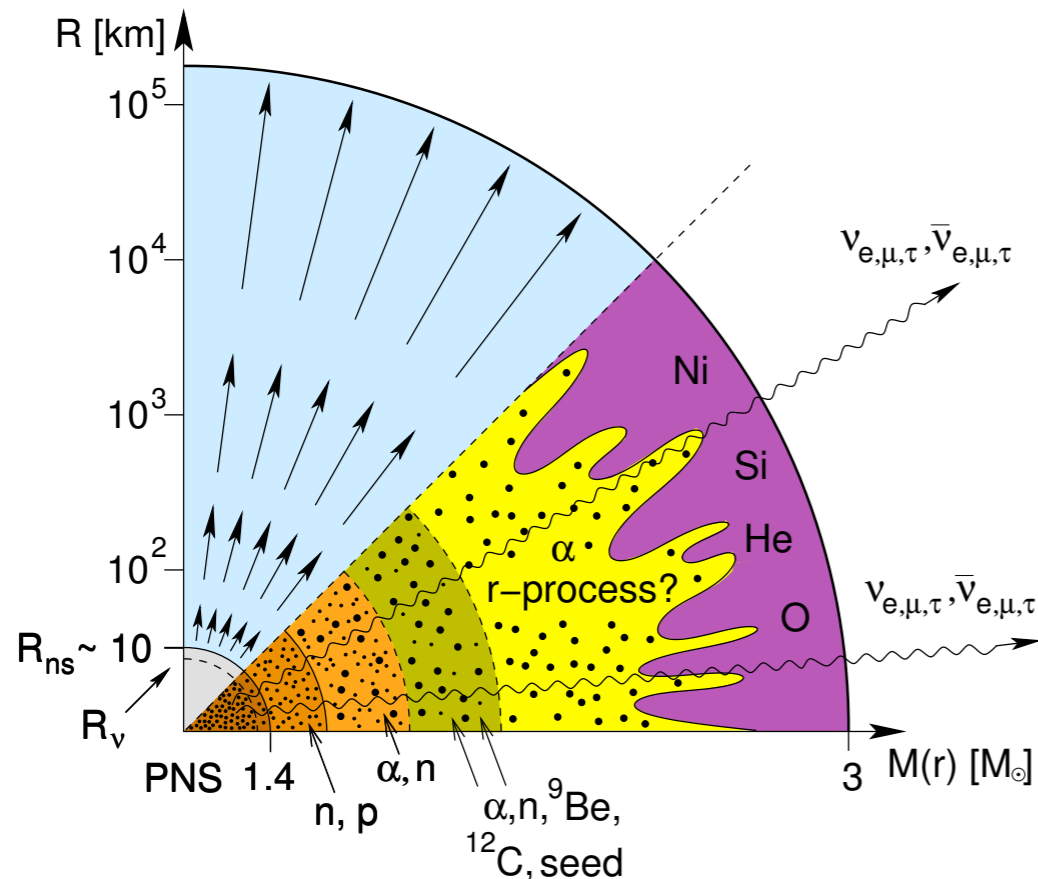


NSE	→	charged particle reactions / $\alpha$ -process	→	r-process weak r-process vp-process
$T = 10 - 8 \text{ GK}$		$8 - 2 \text{ GK}$		$T < 3 \text{ GK}$

# Neutrino-driven wind parameters

r-process  $\Rightarrow$  high neutron-to-seed ratio ( $Y_n/Y_{\text{seed}} \sim 100$ )

- Short **expansion time scale** to inhibit  $\alpha$ -process and formation of seed nuclei
- High **entropy** is equivalent to high photon-to-baryon ratio: photons dissociate seed nuclei into nucleons
- **Electron fraction**:  $Y_e < 0.5$



Entropy per baryon in relativistic gas:  
 $s \propto (kT^3) / (\rho N_A) \Rightarrow s = 10/\Phi$

Photon-to-baryon ratio:  
 $\Phi = n_\gamma / (\rho N_A) \propto (kT^3) / (\rho N_A)$

# Wind and r-process

Meyer et al. 1992 and Woosley et al. 1994:  
r-process: high entropy and low  $Y_e$

Witti et al., Takahashi et al. 1994 needed factor  
5.5 increased in entropy

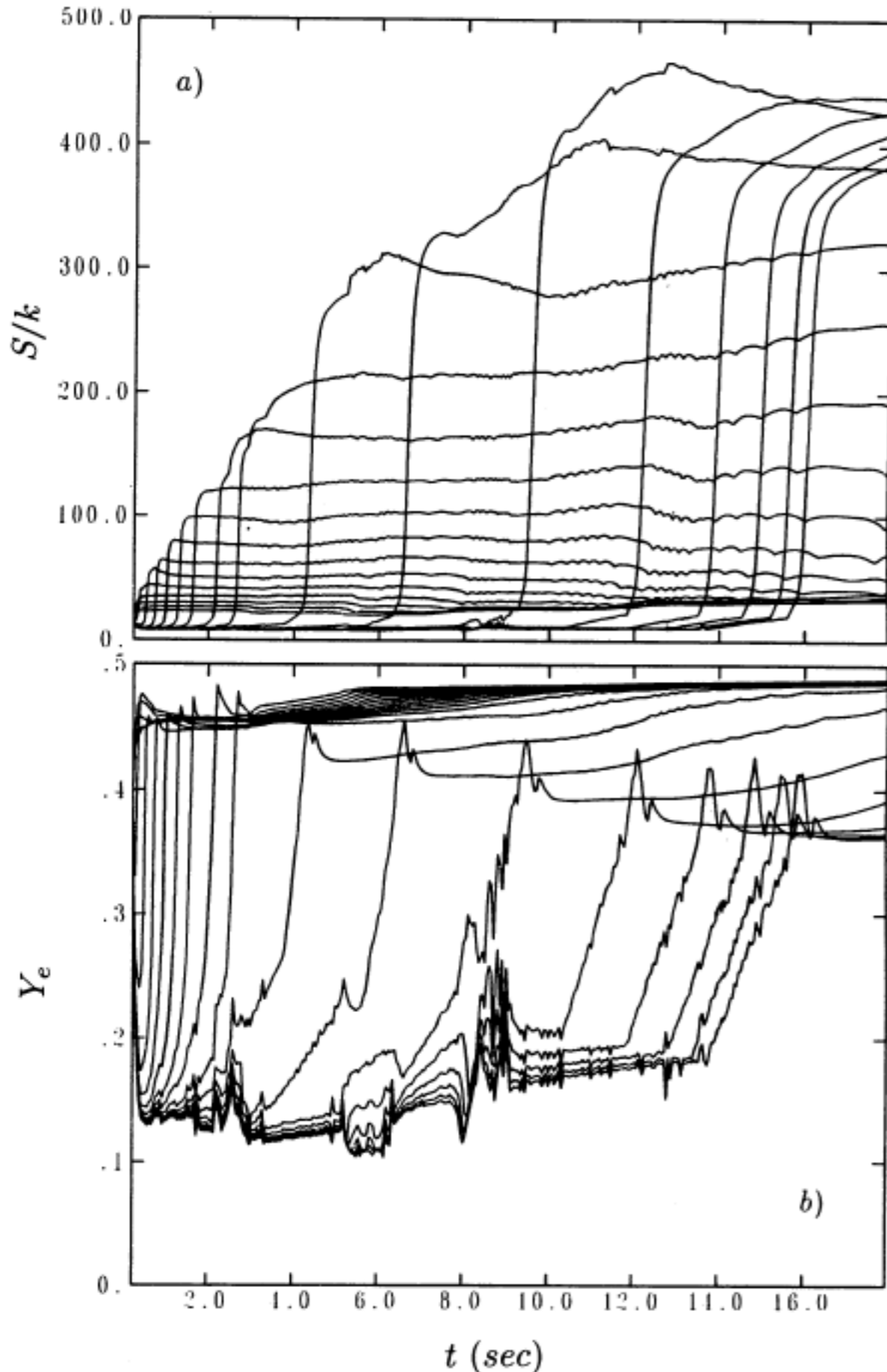
Qian & Woosley 1996: analytic model

$$\dot{M} \propto L_\nu^{5/3} \epsilon_\nu^{10/3} R_{ns}^{5/3} M_{ns}^{-2},$$

$$s \propto L_\nu^{-1/6} \epsilon_\nu^{-1/3} R_{ns}^{-2/3} M_{ns},$$

$$\tau \propto L_\nu^{-1} \epsilon_\nu^{-2} R_{ns} M_{ns}.$$

Thompson, Otsuki, Wanajo, ... (2000-...) parametric steady state winds



# Electron fraction

depends on accuracy of supernova neutrino transport and on details of neutrino interactions in outer layers of neutron star.

$$Y_e \approx \left[ 1 + \frac{L_{\bar{\nu}_e} (\epsilon_{\bar{\nu}_e} - 2\Delta + 1.2\Delta^2/\epsilon_{\bar{\nu}_e})}{L_{\nu_e} (\epsilon_{\nu_e} + 2\Delta + 1.2\Delta^2/\epsilon_{\nu_e})} \right]^{-1} \quad \text{Qian \& Woosley 1996}$$

( $\Delta = m_n - m_p$ )

The neutrino energies are determined by the position (temperature) where neutrinos decouple from matter: **neutrinosphere**

Raffelt 2001



radius

Neutrino sphere ( $T_{NS}$ )

# Electron fraction

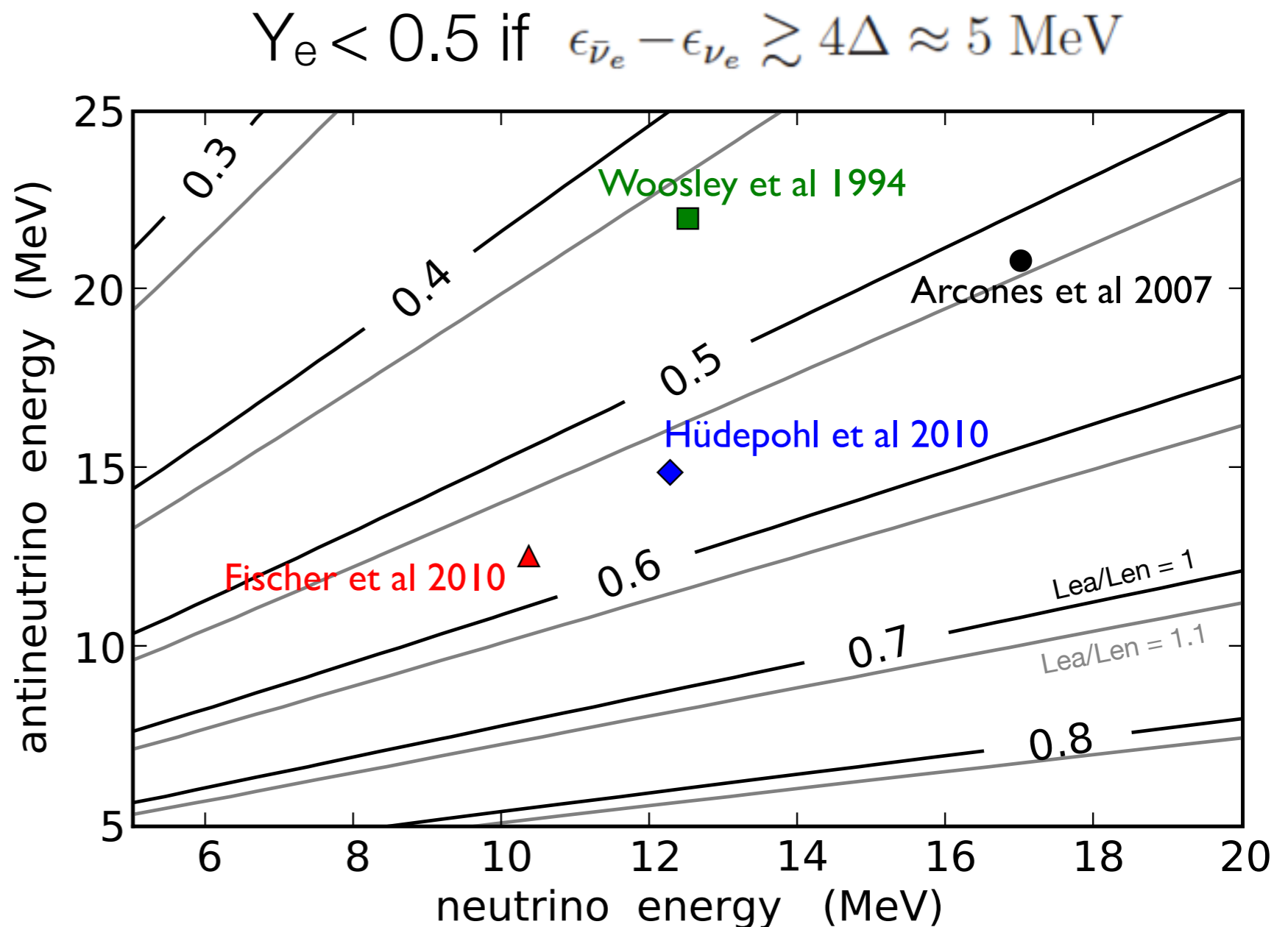
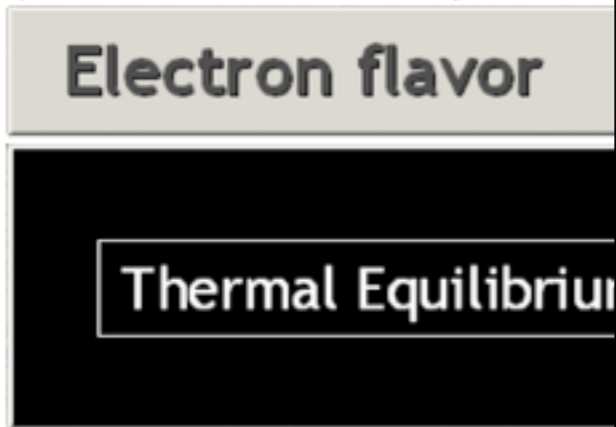
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( $\Delta = m_n - m_p$ )

The neutrino energies are decouple from matter: n

Raffelt 2001



Charged-current weak interaction processes in hot and dense matter and its impact on the spectra of neutrinos emitted from proto-neutron star cooling

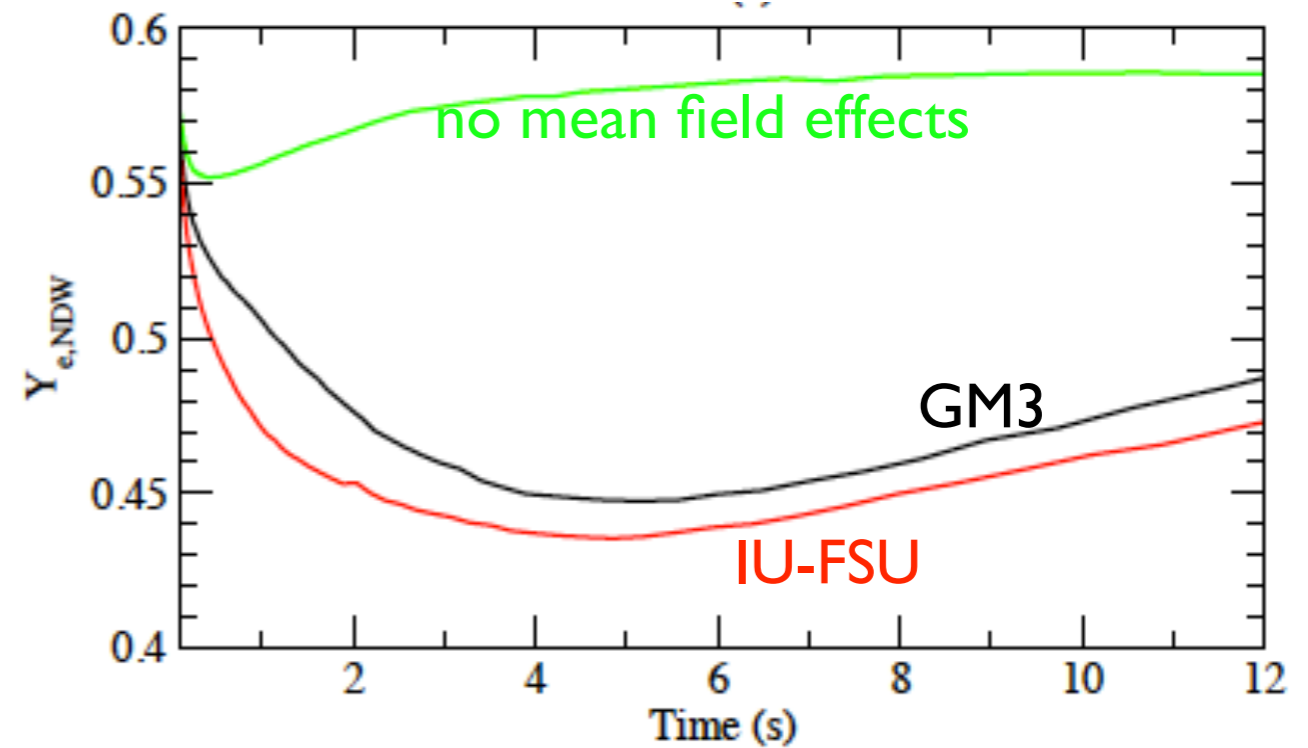
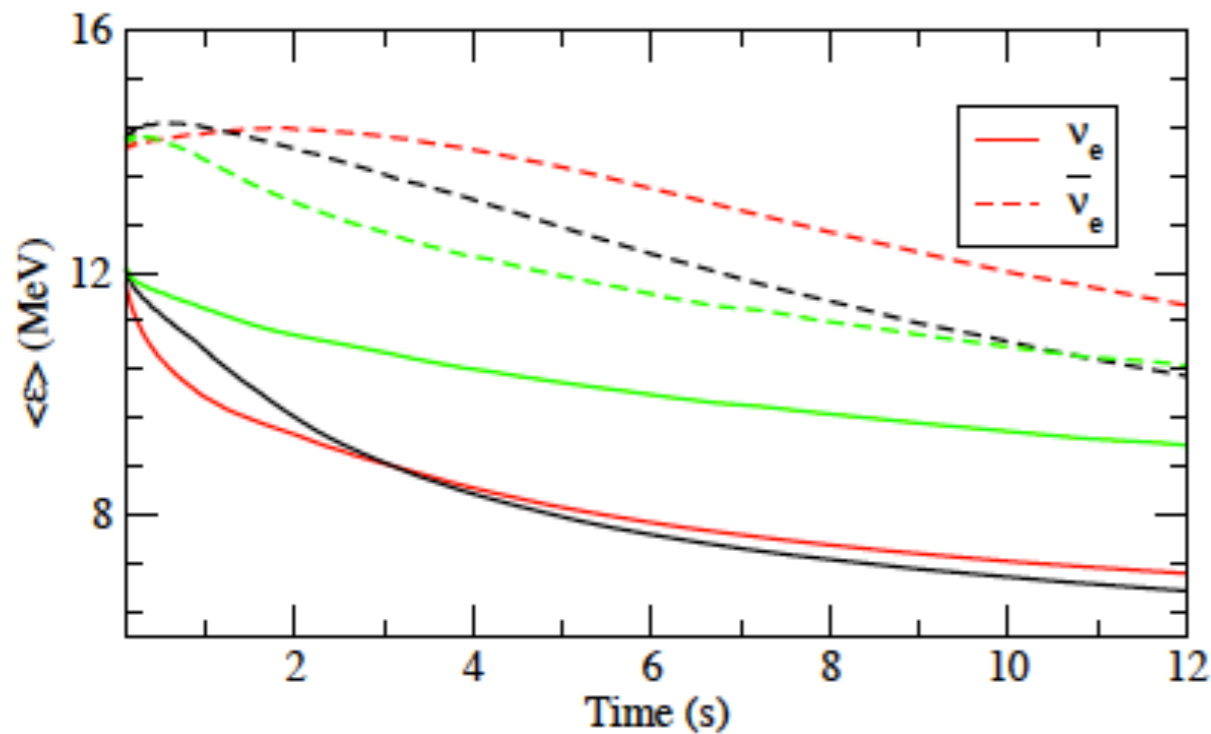
G. Martínez-Pinedo,<sup>1,2</sup> T. Fischer,<sup>2,1</sup> A. Lohs,<sup>1</sup> and L. Huther<sup>1</sup>

A NEW CODE FOR PROTO-NEUTRON STAR EVOLUTION

L. F. ROBERTS<sup>†</sup>

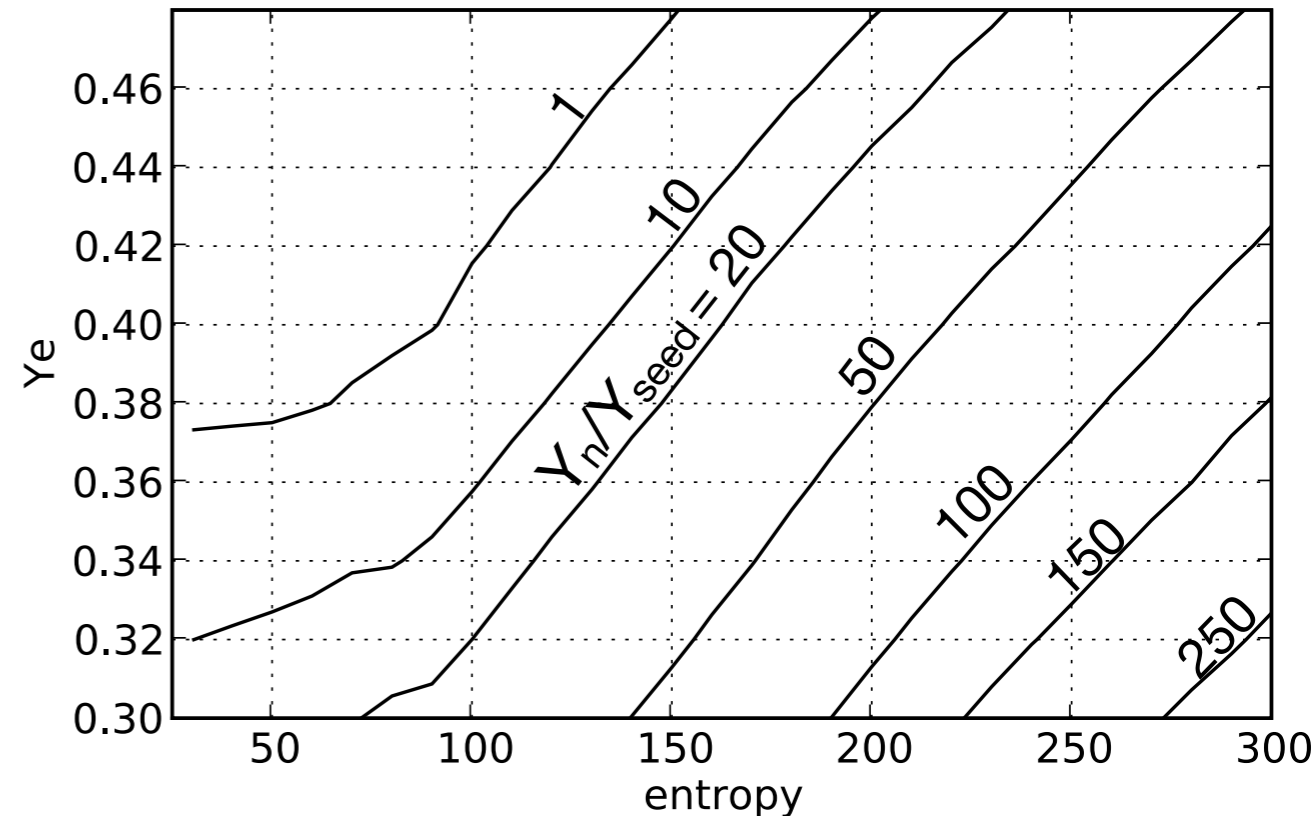
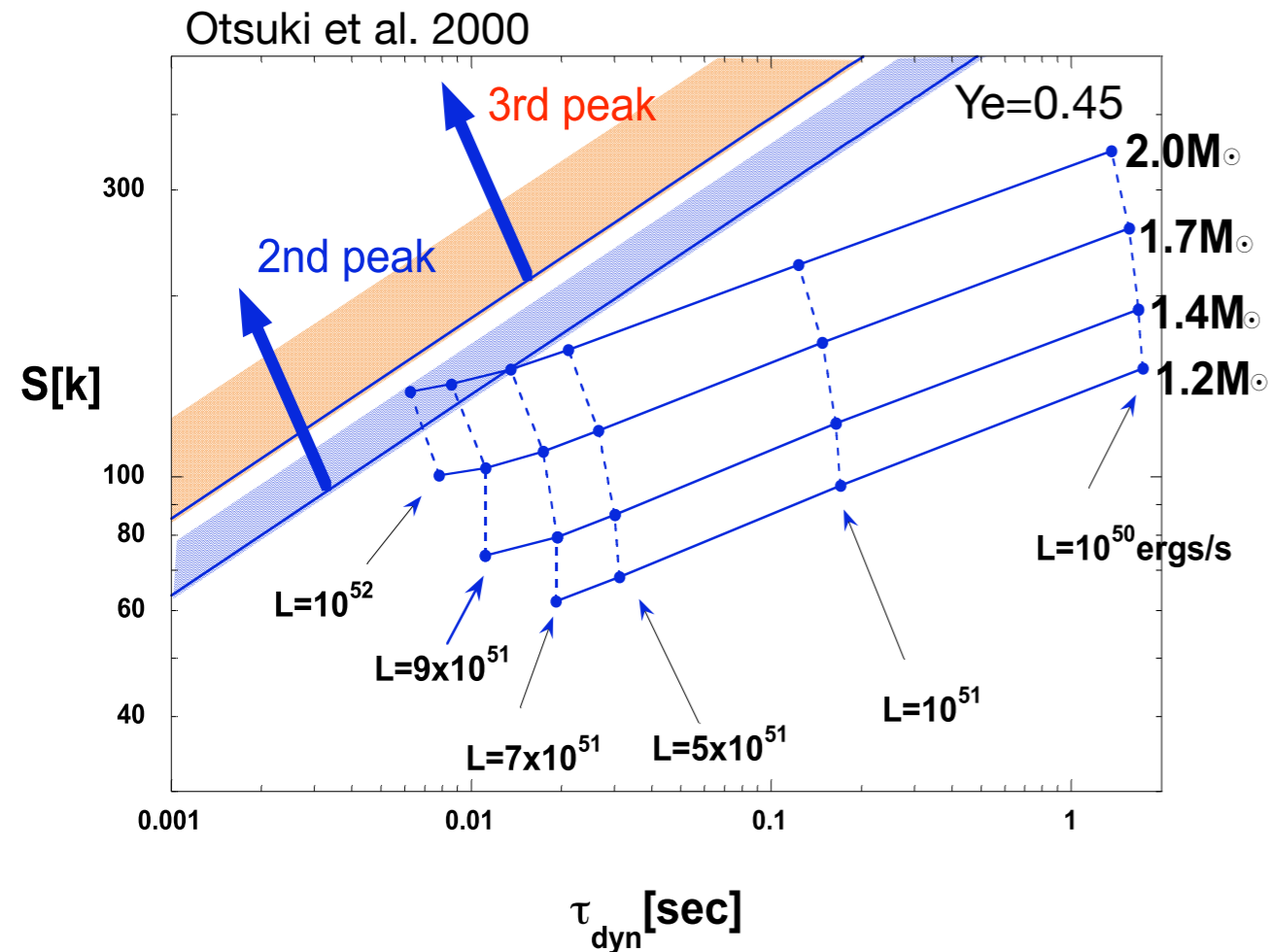
Medium modification of the charged current neutrino opacity and its implications

L. F. Roberts<sup>1</sup> and Sanjay Reddy<sup>2</sup>



# Wind parameters and r-process

Necessary conditions identified by steady-state models (e.g., Otsuki et al. 2000, Thompson et al. 2001)



Conditions are not realized in recent simulations

(Arcones et al. 2007, Fischer et al. 2010, Hüdepohl et al. 2010, Roberts et al. 2010, Arcones & Janka 2011)

$$S_{\text{wind}} = 50 - 120 \text{ k}_B/\text{nuc}$$

$$\tau = \text{few ms}$$

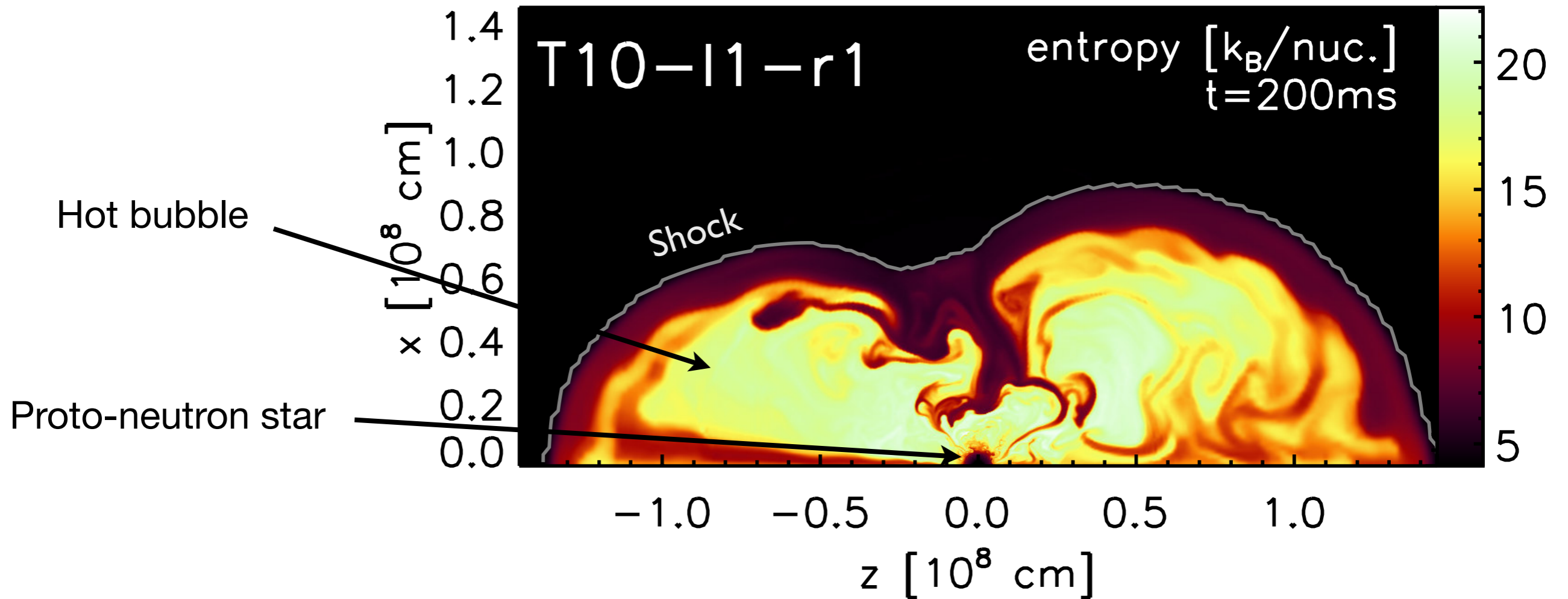
$$Y_e > 0.5?$$

**Additional ingredients:** wind termination, extra energy source, rotation and magnetic fields, neutrino oscillations

Review: Arcones & Thielemann (arxiv: 1207.2527)



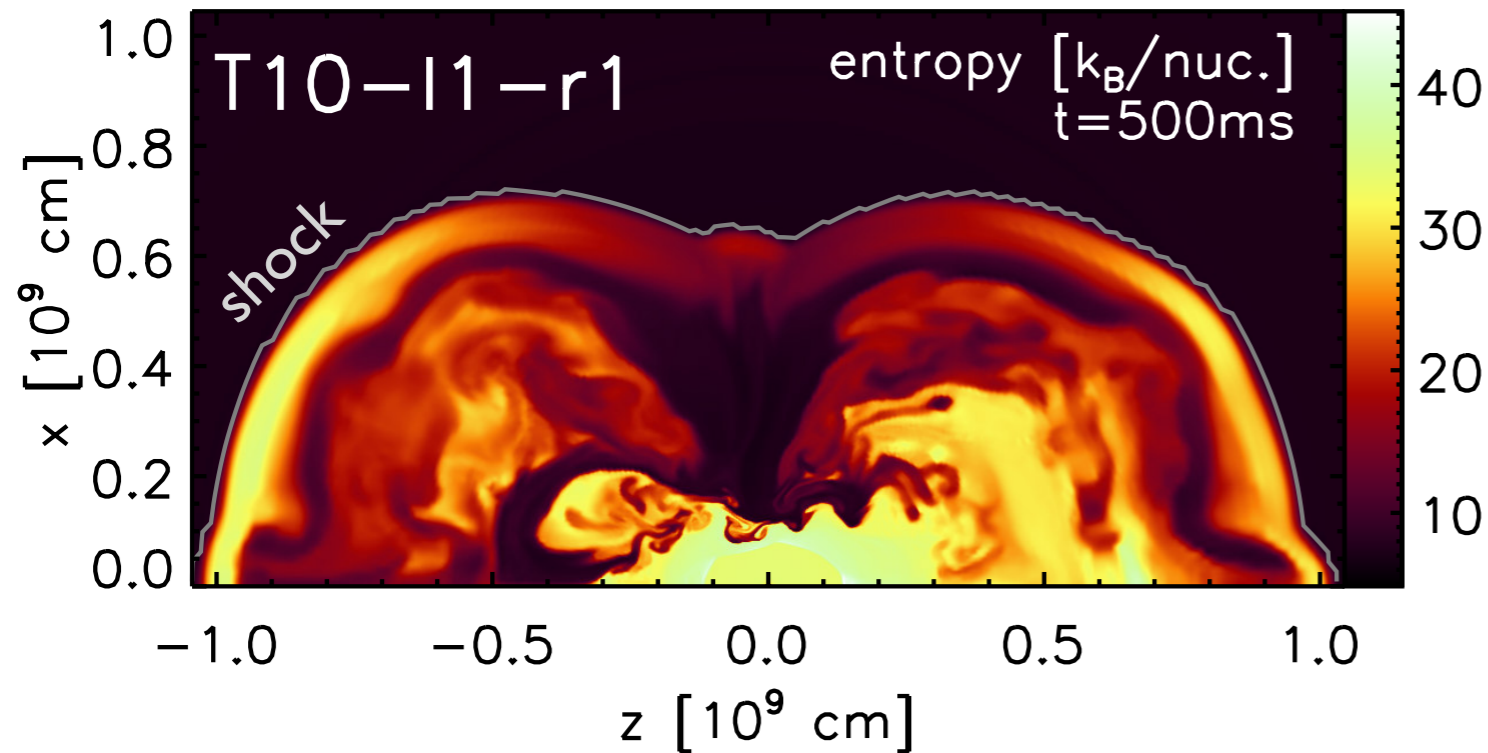
# Core-collapse supernova simulations



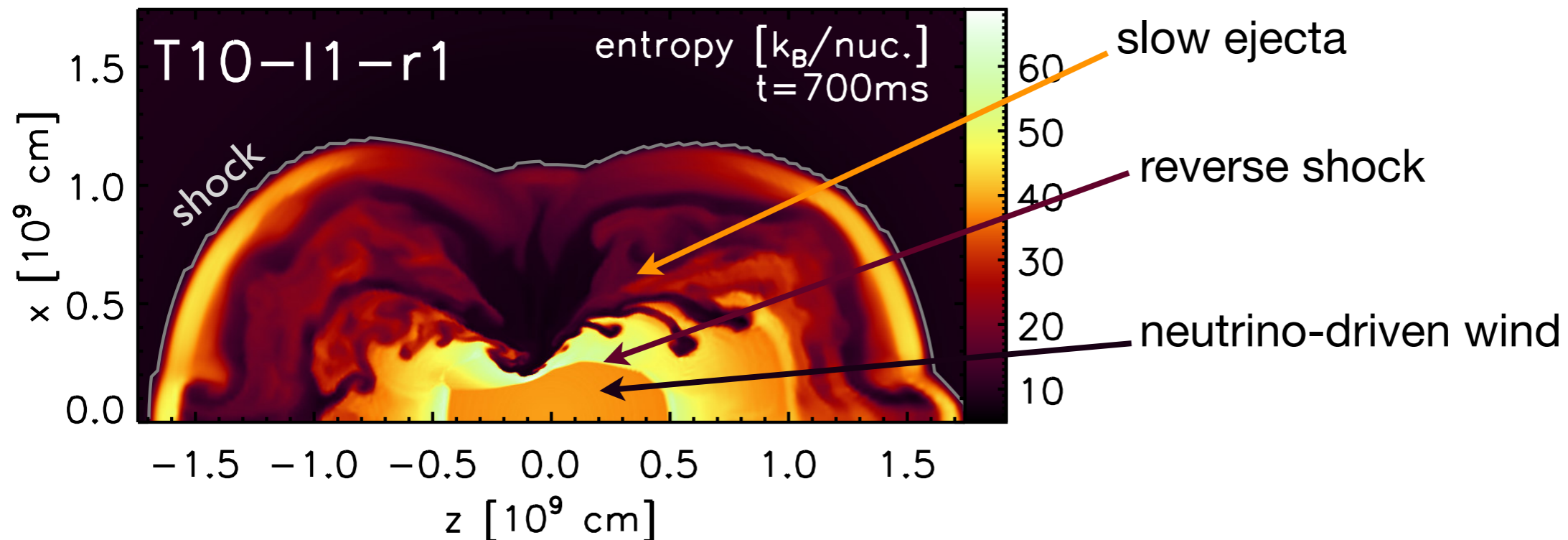
Long-time hydrodynamical simulations:

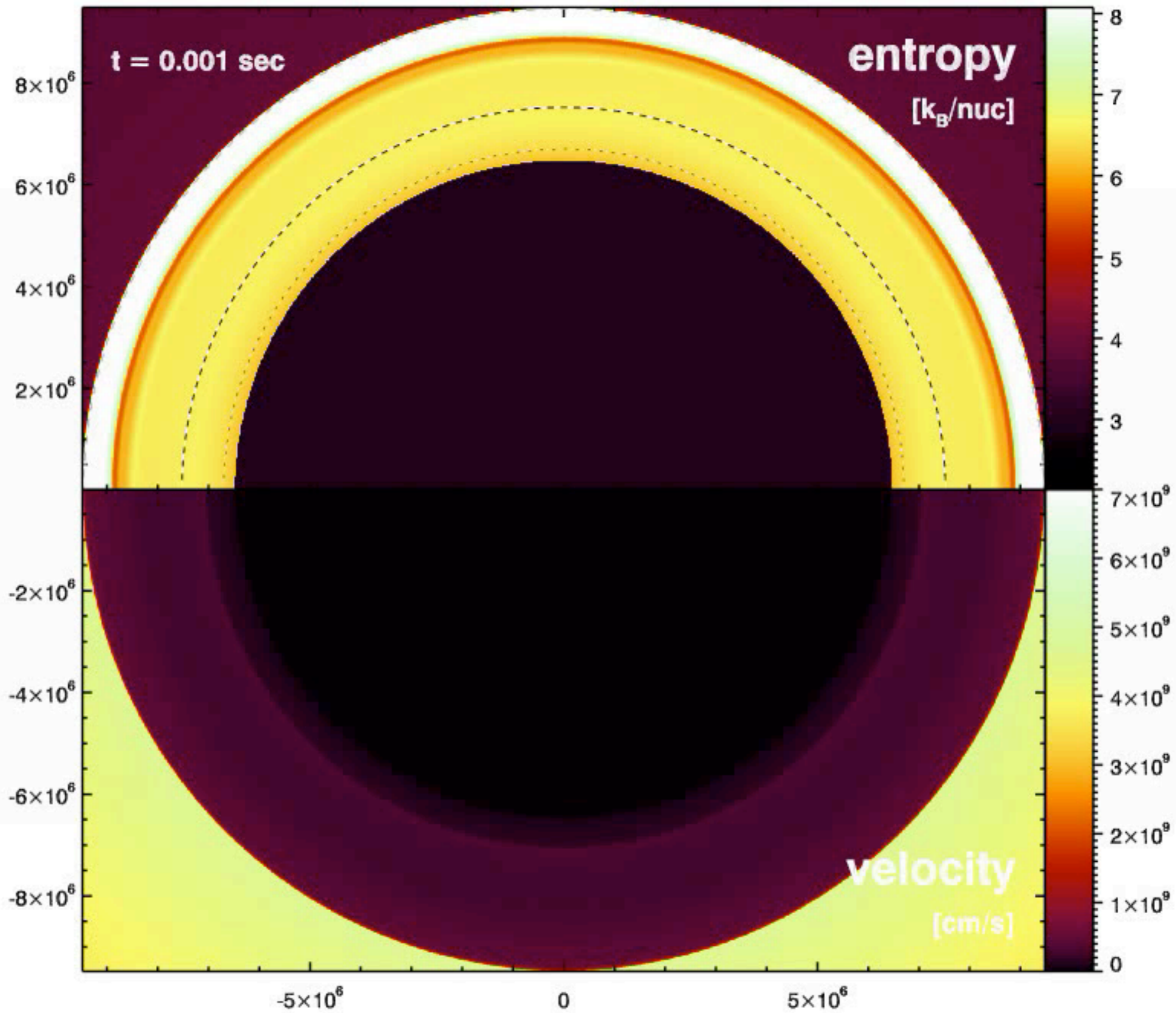
- ejecta evolution from  $\sim 5\text{ms}$  after bounce to  $\sim 3\text{s}$  in 2D (Arcones & Janka 2011) and  $\sim 10\text{s}$  in 1D (Arcones et al. 2007)
- explosion triggered by neutrinos
- detailed study of nucleosynthesis-relevant conditions

# Neutrino-driven wind in 2D



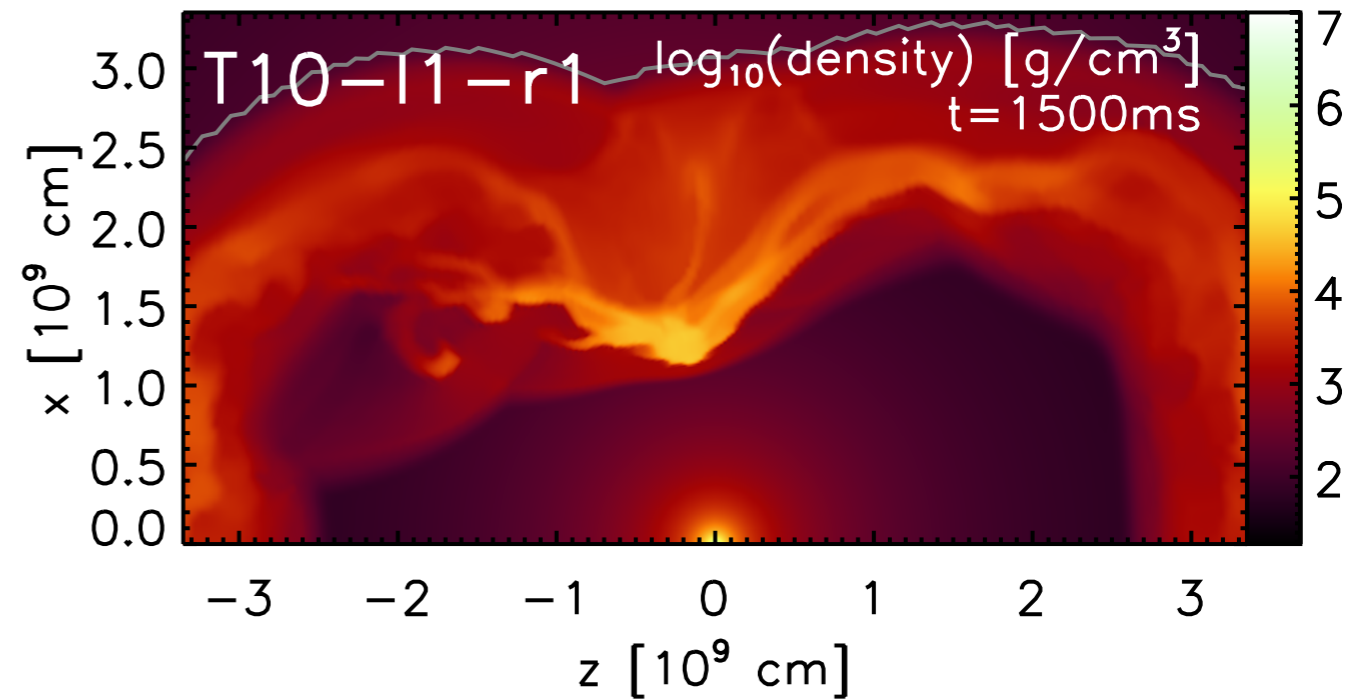
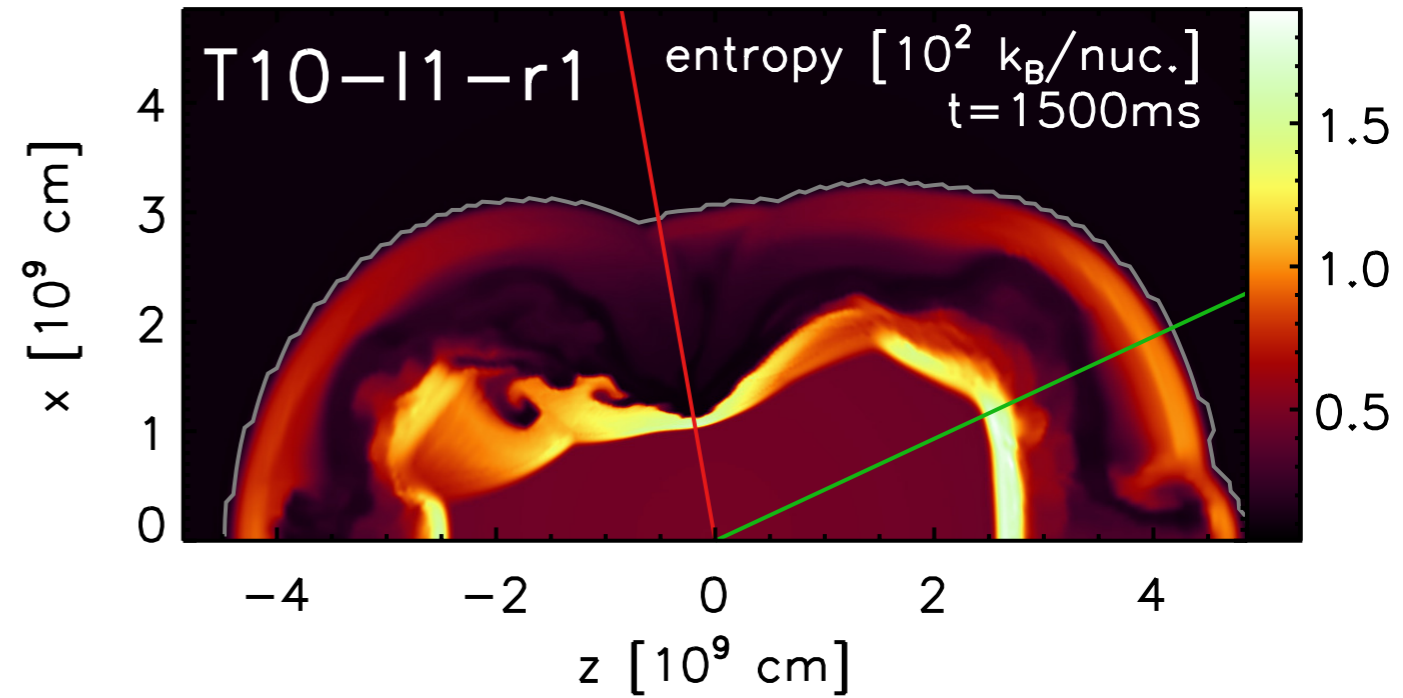
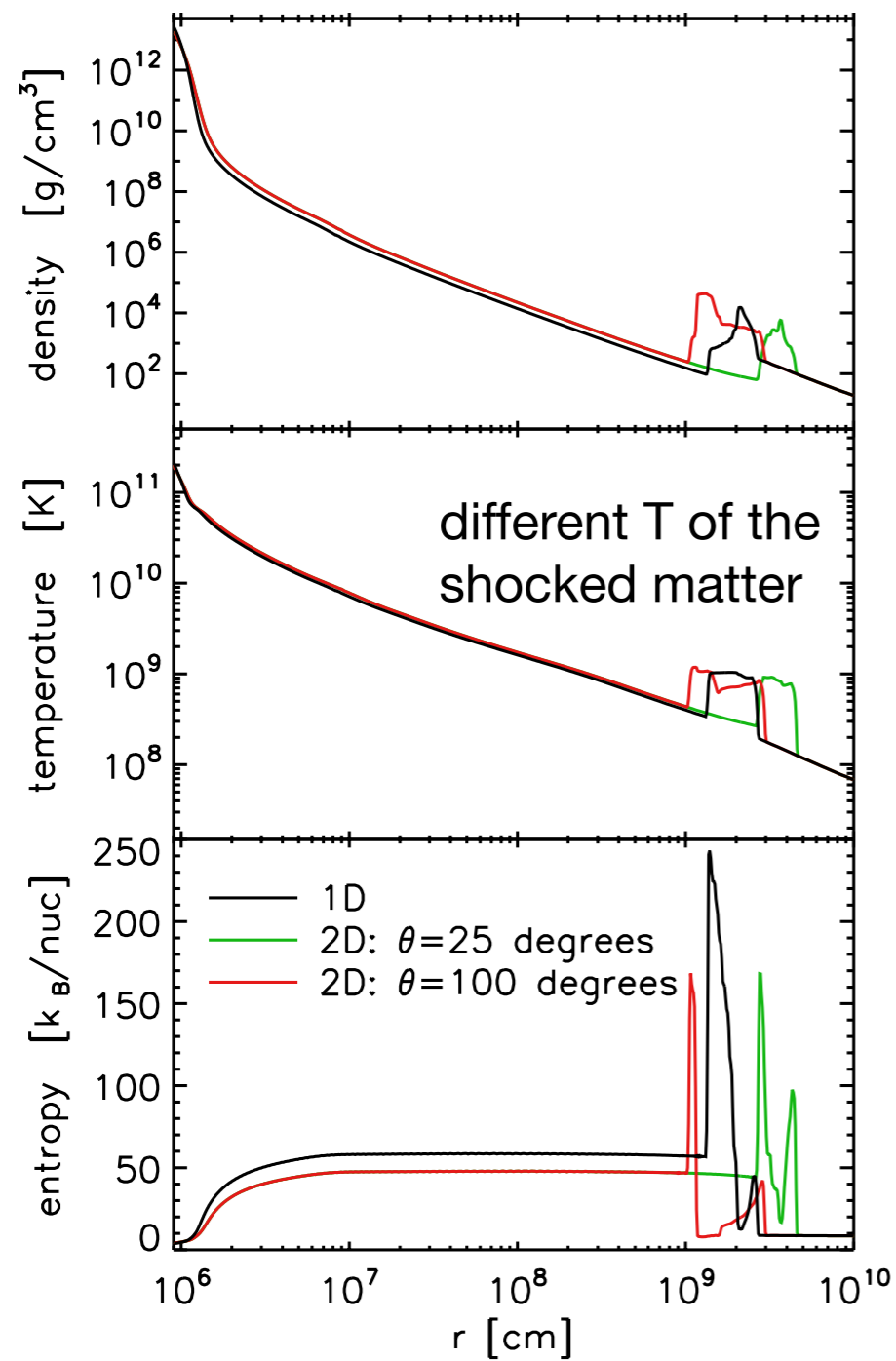
Supersonic neutrino-driven wind collides with slow supernova ejecta: reverse shock





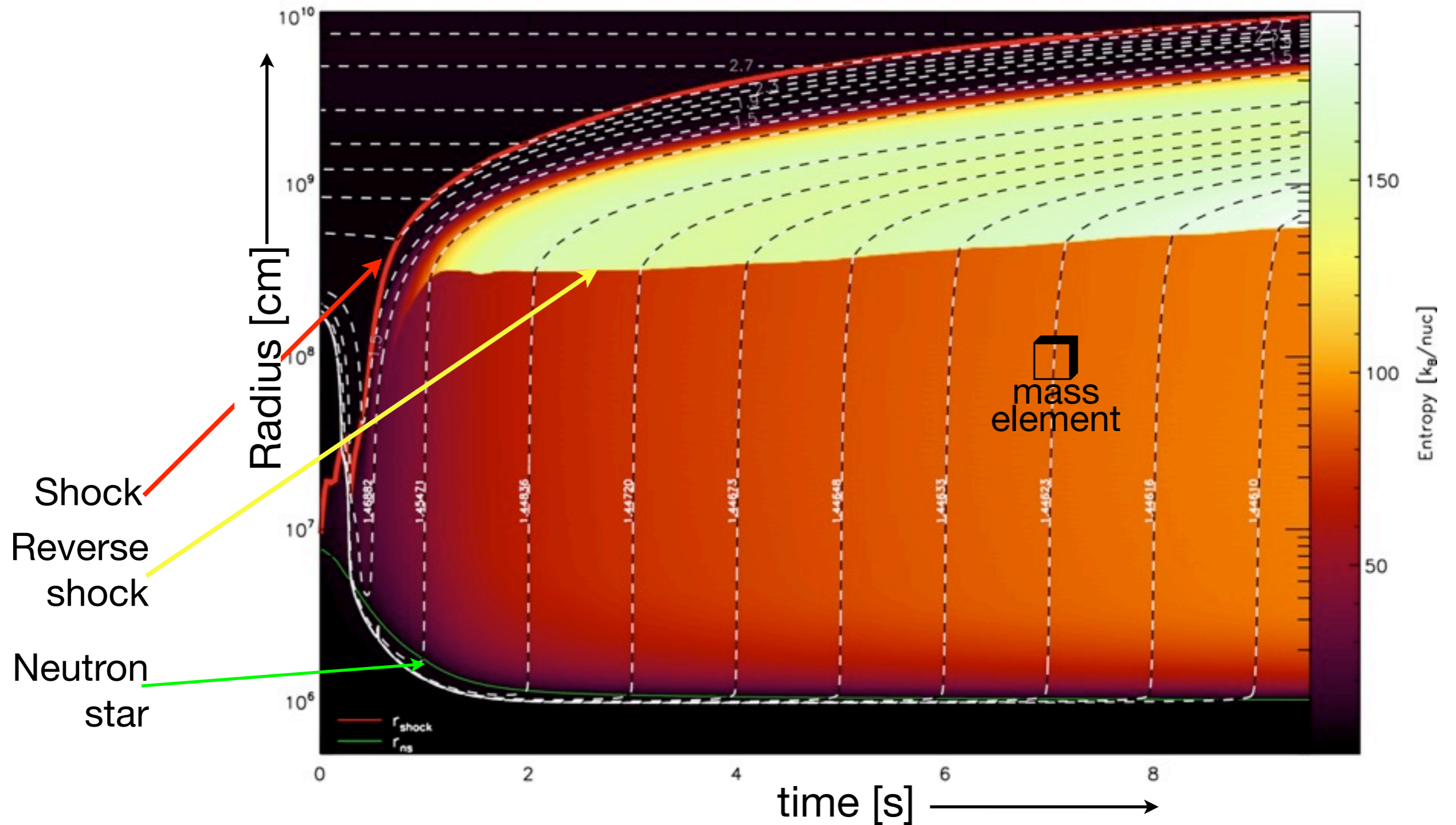
# Neutrino-driven wind in 2D and 1D

## Spherically symmetric wind



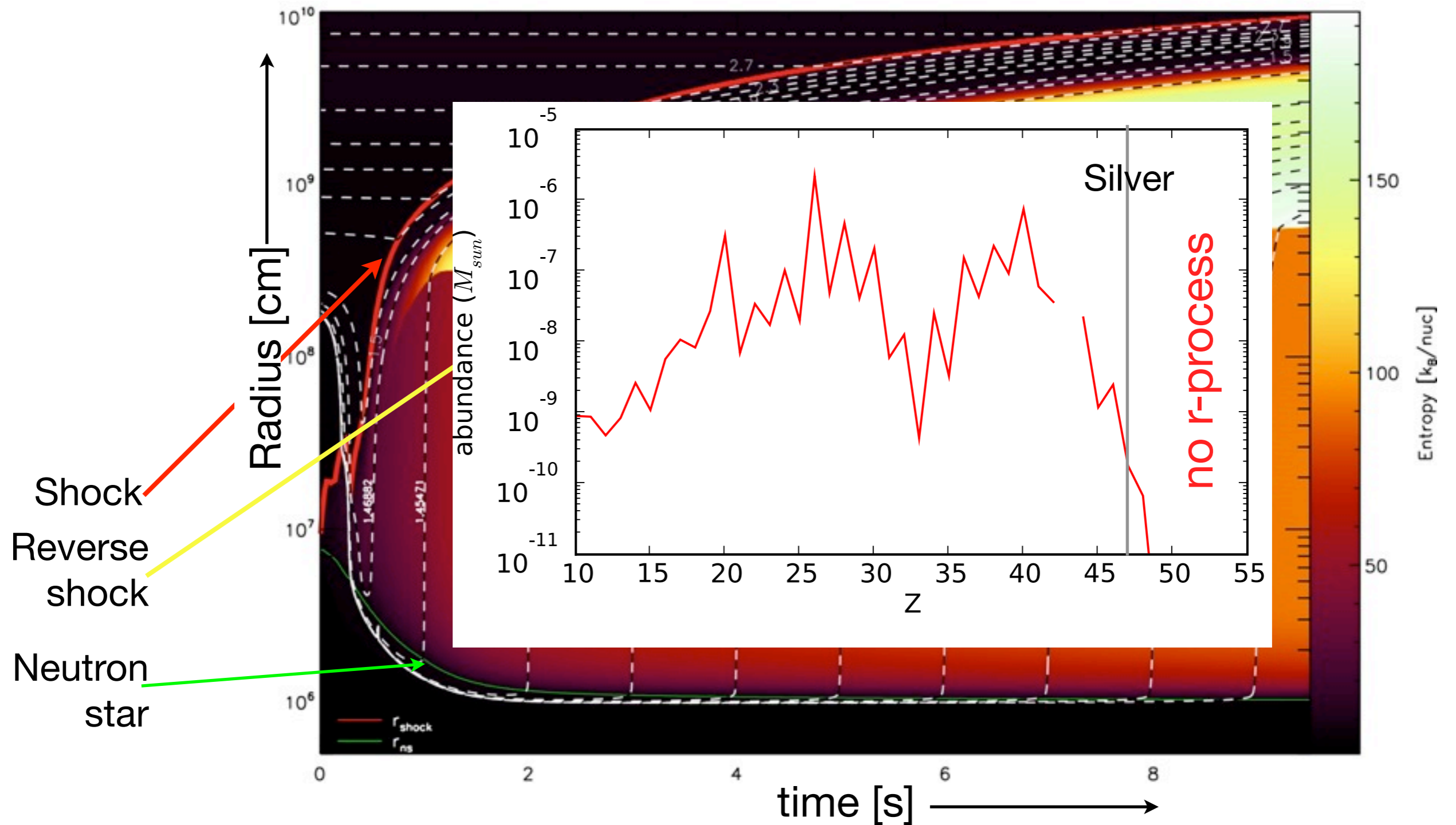
# 1D simulations for nucleosynthesis studies

Arcones et al 2007



# 1D simulations for nucleosynthesis studies

Arcones et al 2007



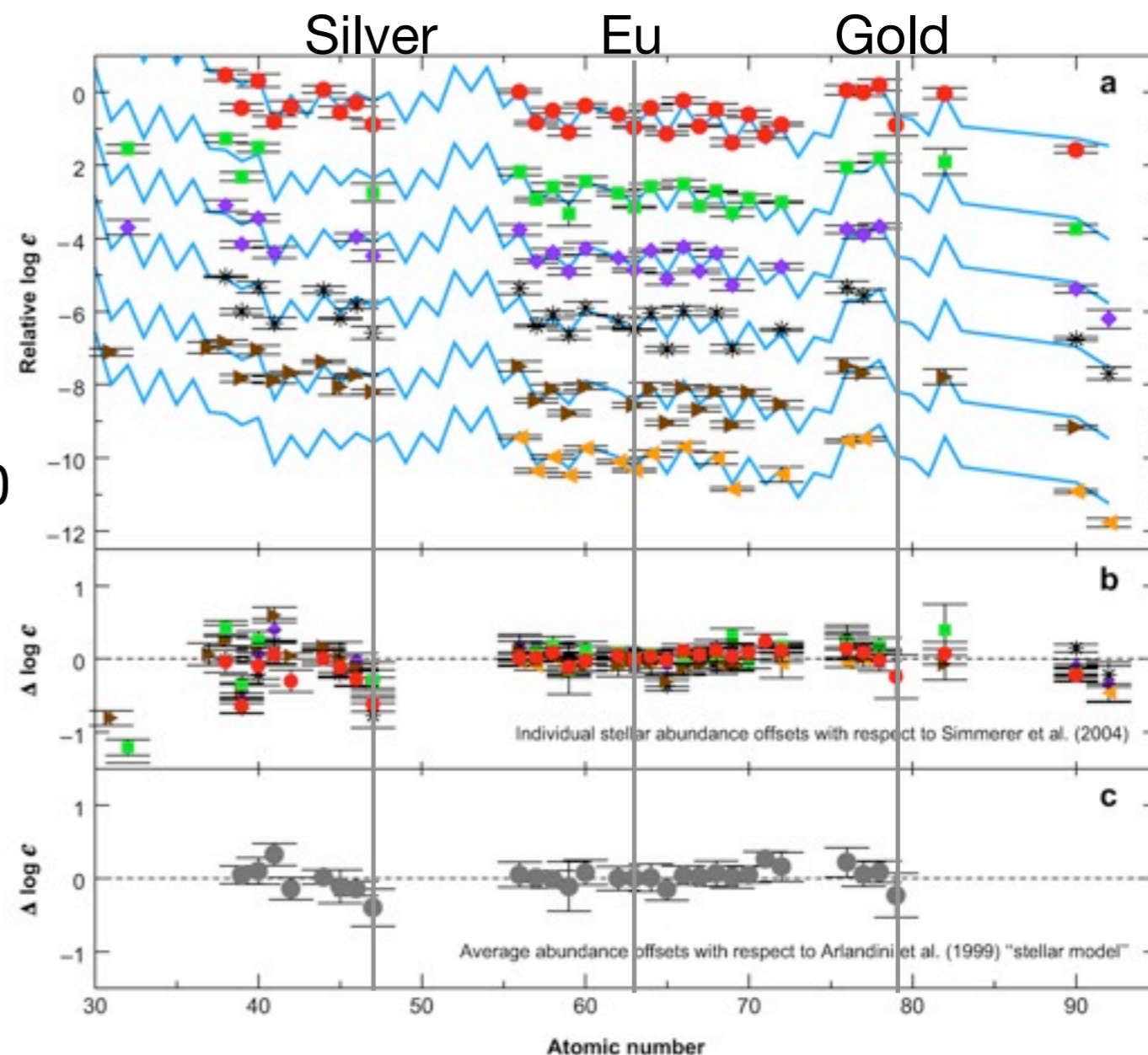
# r-process in ultra metal-poor stars

Abundances of r-process elements in:

- ultra metal-poor stars and
- r-process solar system:  $N_{\text{solar}} - N_{\text{s}}$

Robust r-process for  $56 < Z < 83$

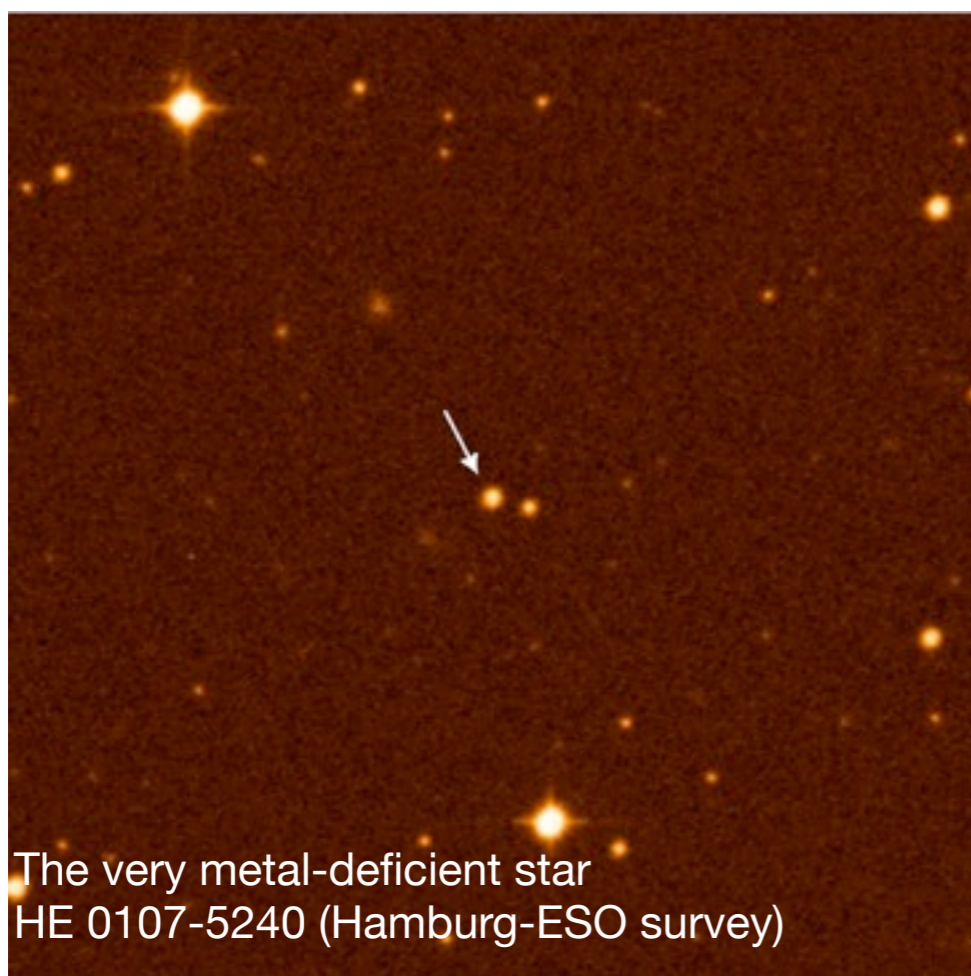
Scatter for lighter heavy elements,  $Z \sim 40$



- CS 22892-052: Sneden et al. (2003)
- HD 115444: Westin et al. (2000)
- ◆ BD+17°324817: Cowan et al. (2002)
- \* CS 31082-001: Hill et al. (2002)
- ▲ HD 221170: Ivans et al. (2006)
- ▲ HE 1523-0901: Frebel et al. (2007)

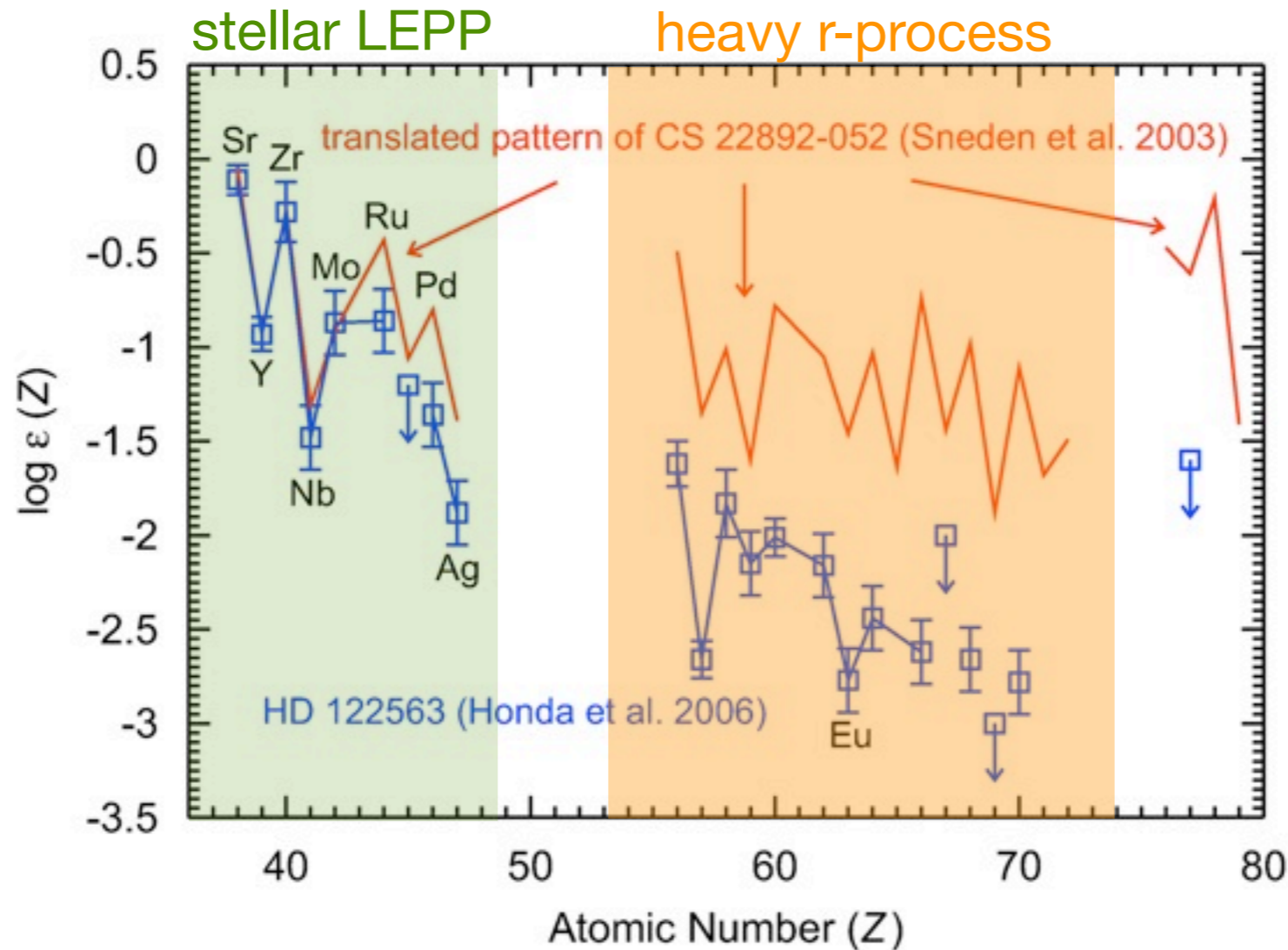
$$\log(\epsilon(E)) = \log(N_E/N_H) + 12$$

Sneden, Cowan, Gallino 2008



# LEPP: Lighter Element Primary Process

Ultra metal-poor stars with **high** and **low** enrichment of heavy r-process nuclei suggest: two components or sites (Qian & Wasserburg):



Travaglio et al. 2004: solar = r-process + s-process + **solar LEPP**

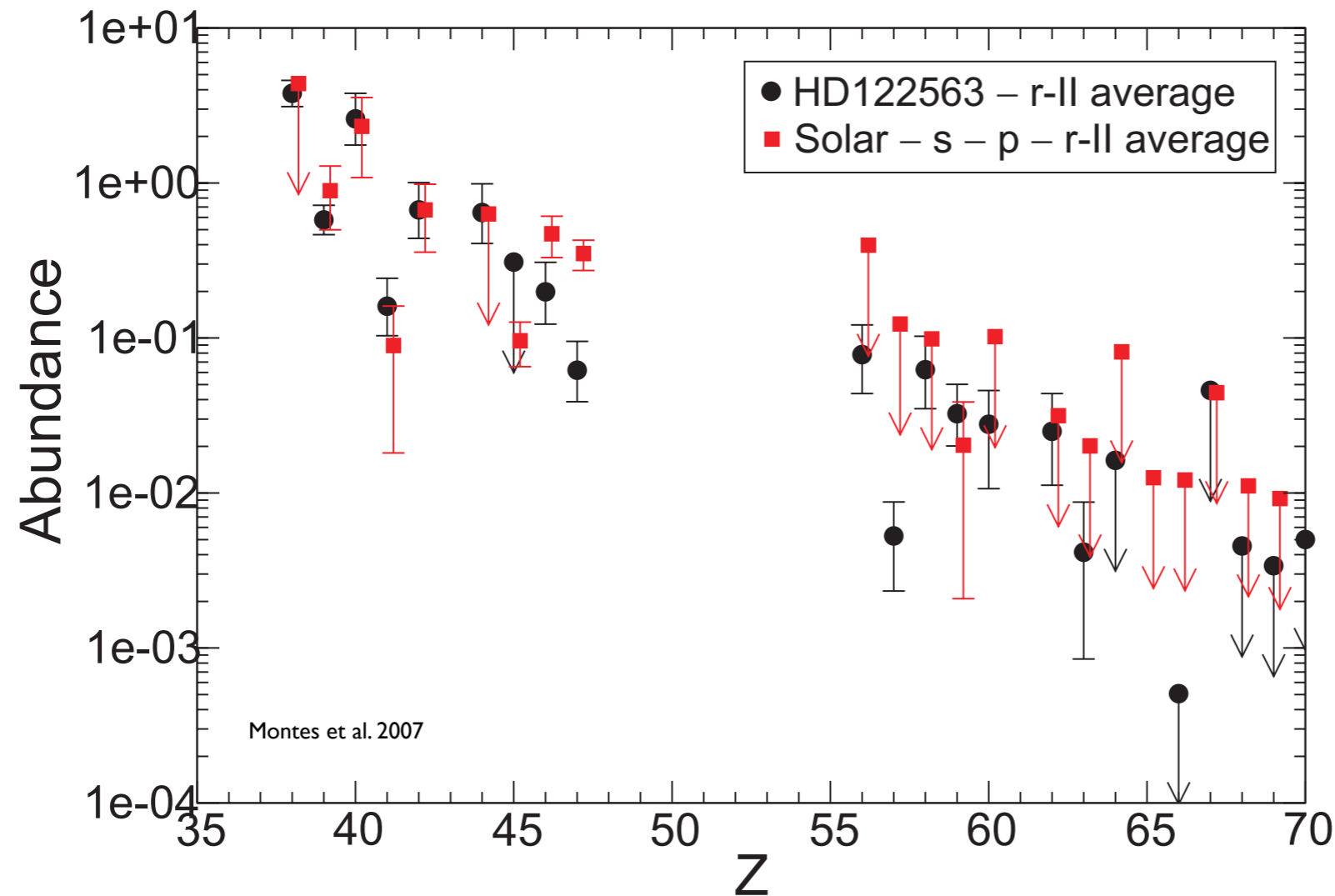
LEPP contributes 20-30% of solar Sr-Y-Zr and explains under-productions of "s-only" isotopes from  $^{96}\text{Mo}$  to  $^{130}\text{Xe}$

Montes et al. 2007: solar LEPP ~ stellar LEPP → unique?



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Montes et al. 2007: solar LEPP ~ stellar LEPP → unique?

# LEPP in neutrino-driven winds

Integrated abundances for different progenitors

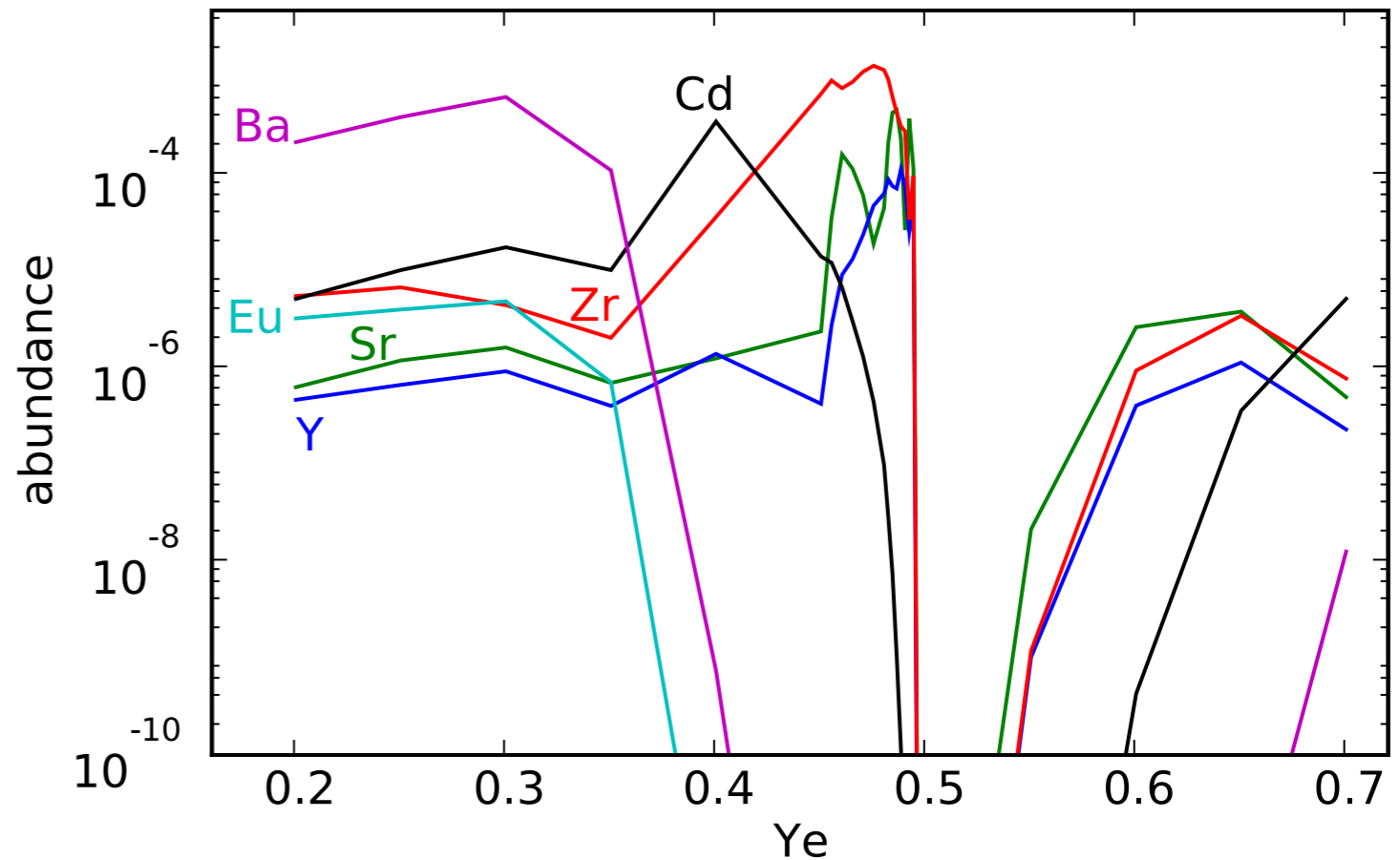
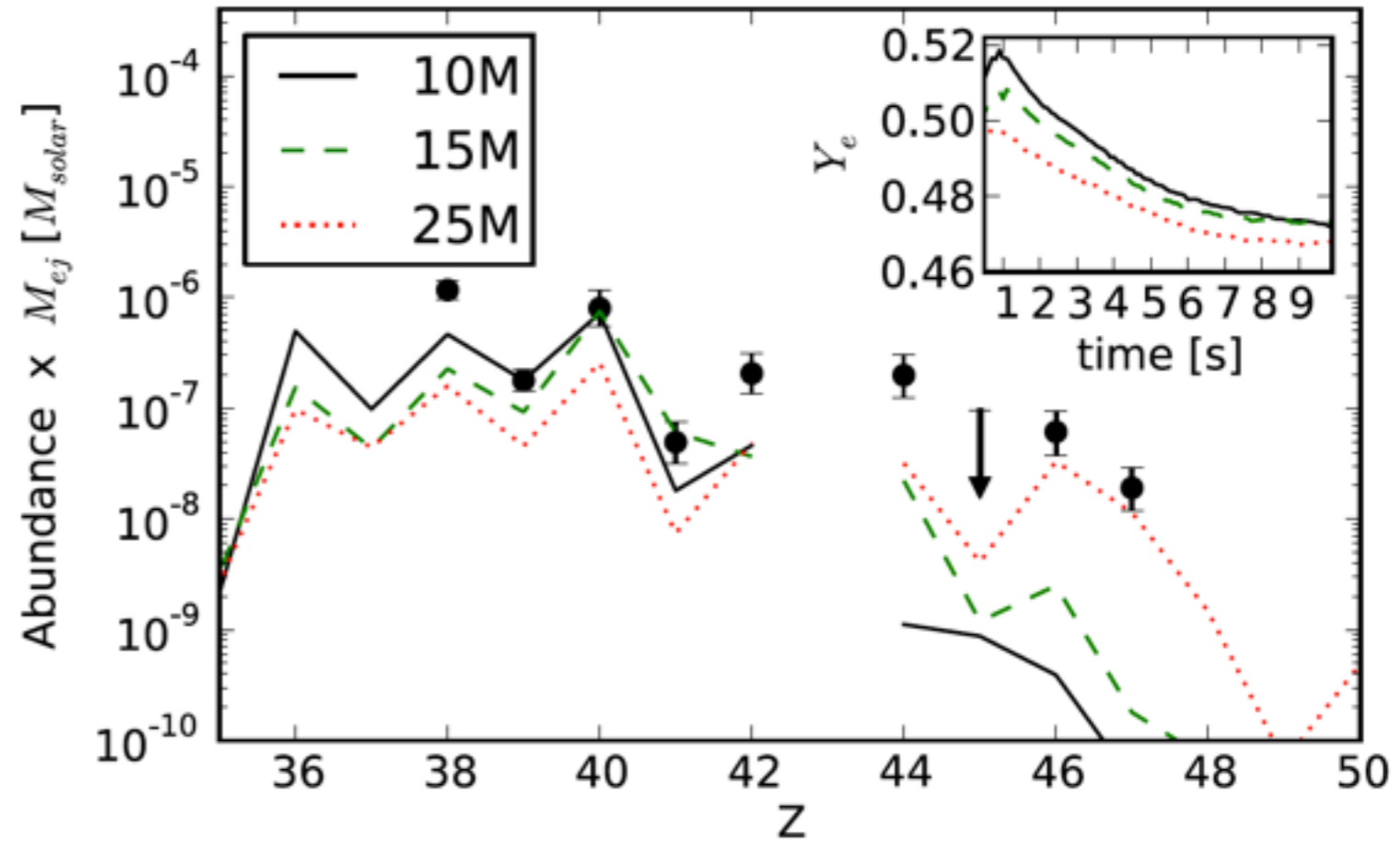
Massive progenitors:  
higher entropy  $\Rightarrow$  heavier nuclei

Simplified neutrino transport:  
approximated  $Y_e$

Impact of  $Y_e$  on wind  
nucleosynthesis:

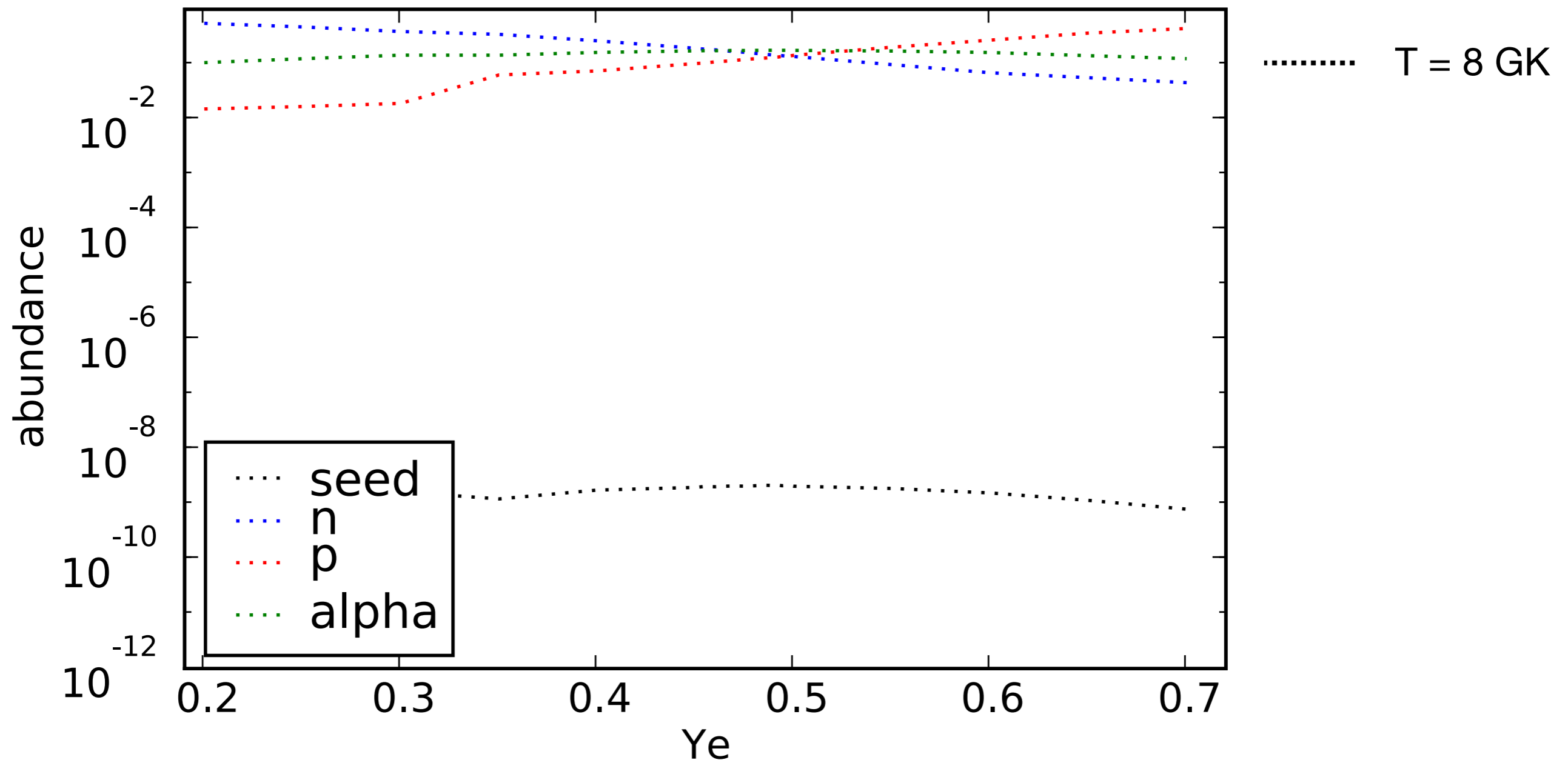
- r-process only for extreme low  $Y_e$
- LEPP in neutron- and proton-rich conditions

(Arcones & Montes, 2011)



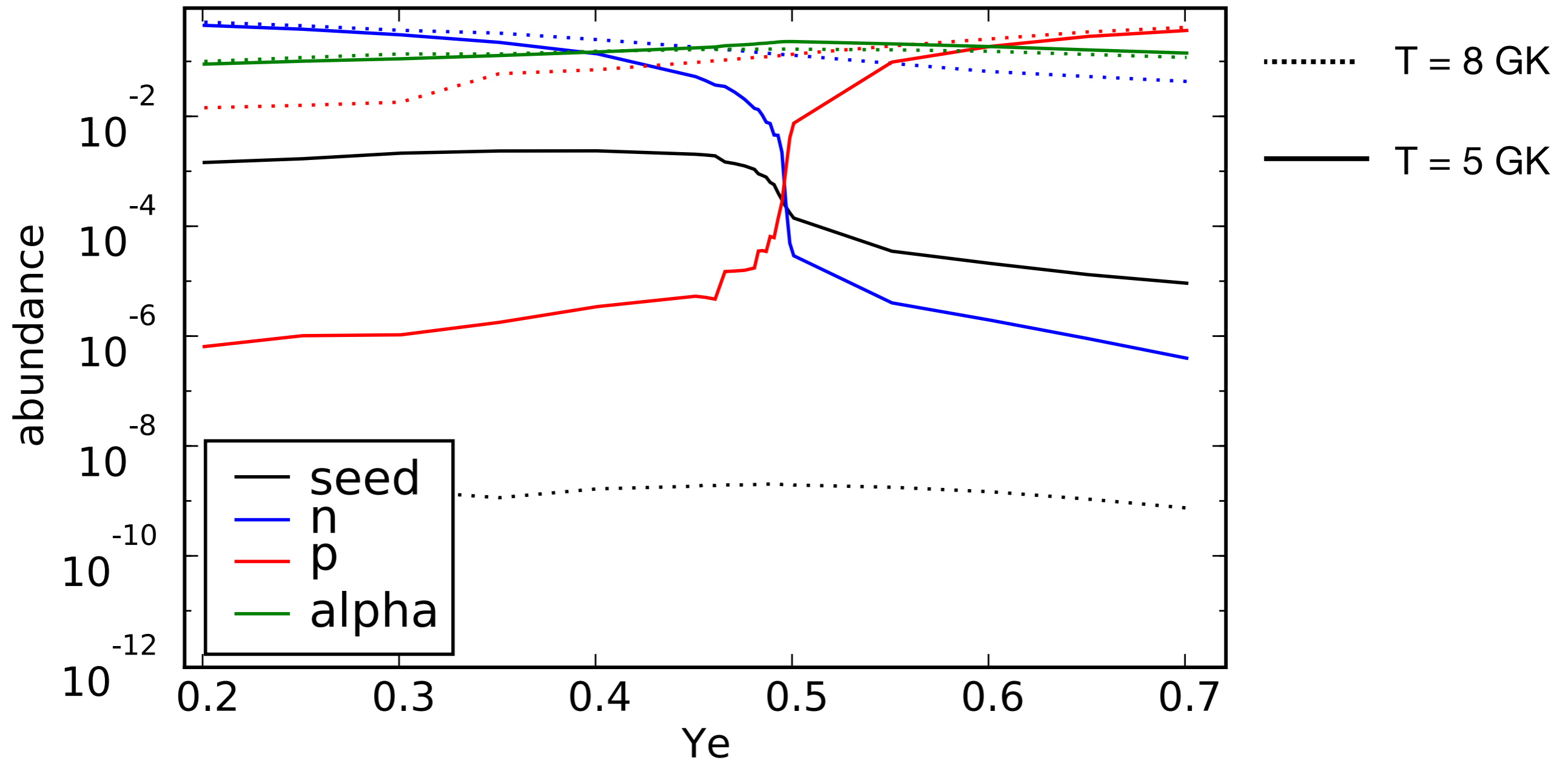
# Wind nucleosynthesis and $Y_e$

Initial composition is given by NSE, at high temperatures only n, p and alphas.



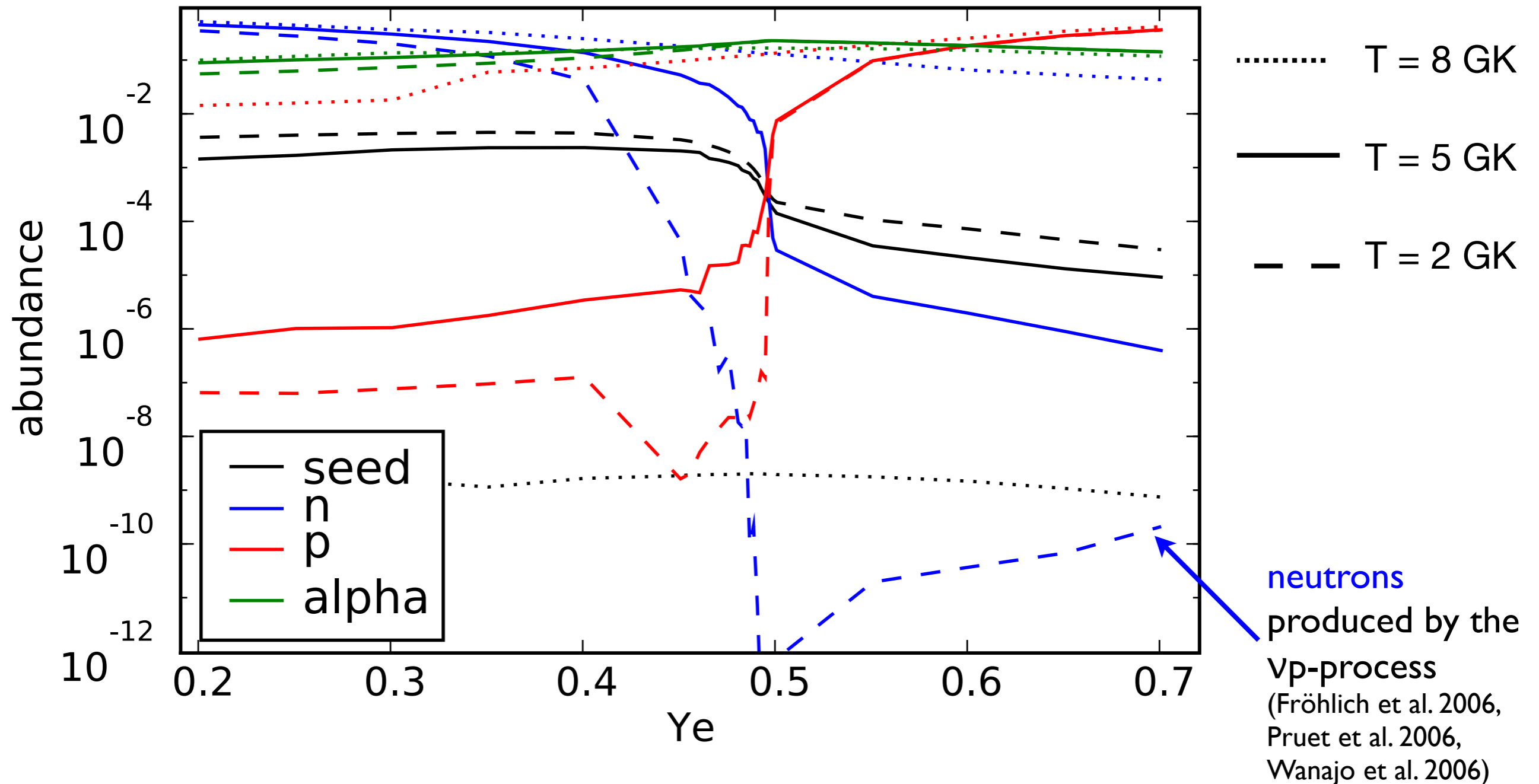
# Wind nucleosynthesis and $Y_e$

Alpha particles recombine forming seed nuclei.



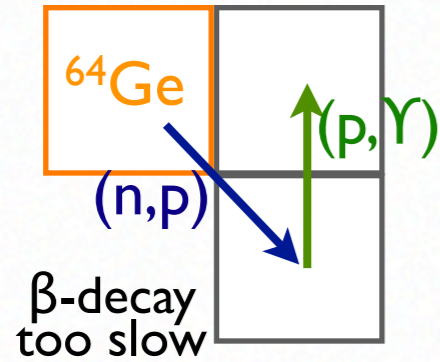
# Wind nucleosynthesis and $Y_e$

At freeze-out neutron- and proton-to-seed ratio determine production of heavy elements.



# vp-process

Z

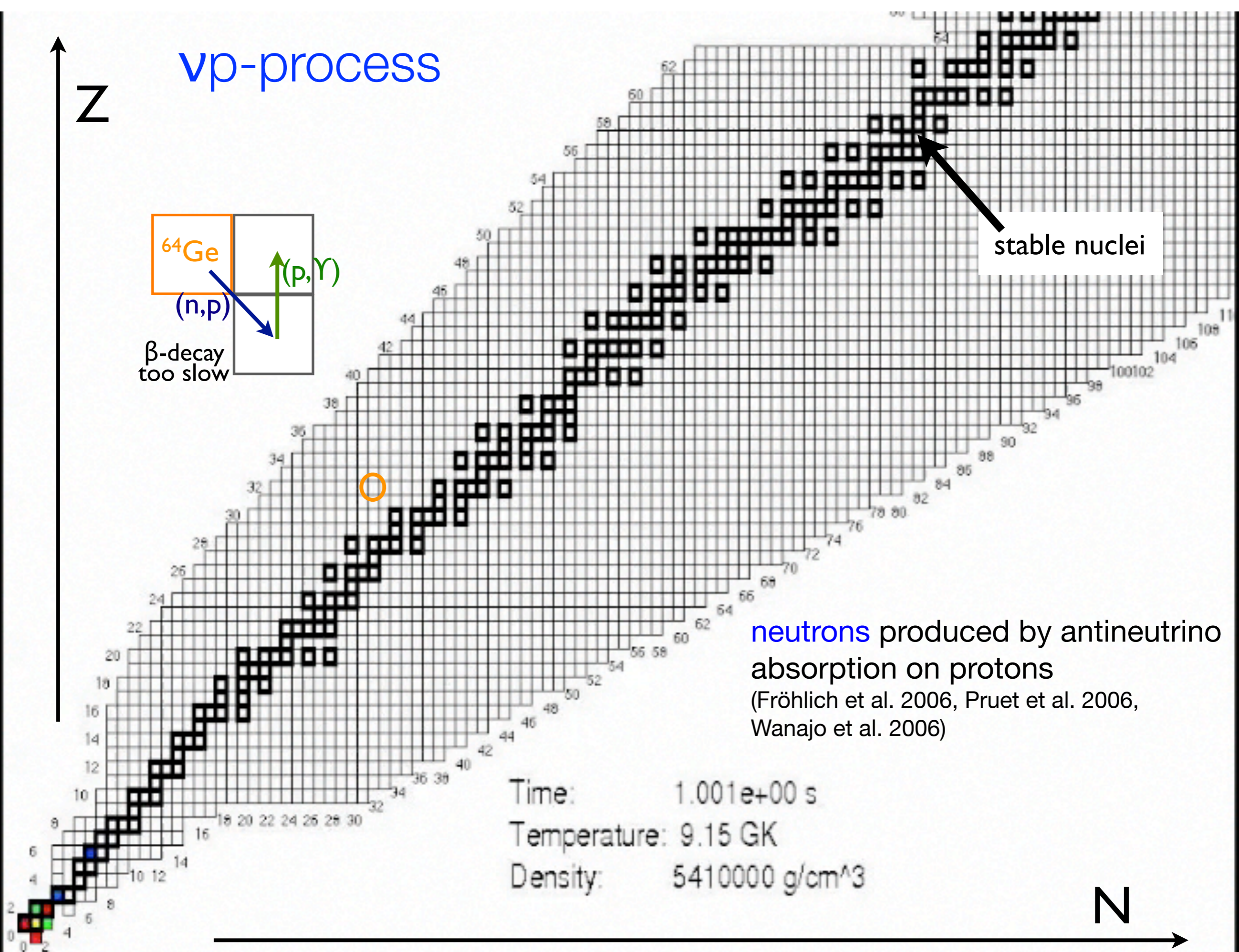


stable nuclei

neutrons produced by antineutrino absorption on protons  
(Fröhlich et al. 2006, Pruet et al. 2006, Wanajo et al. 2006)

Time: 1.001e+00 s  
Temperature: 9.15 GK  
Density: 5410000 g/cm<sup>3</sup>

N



# vp-process

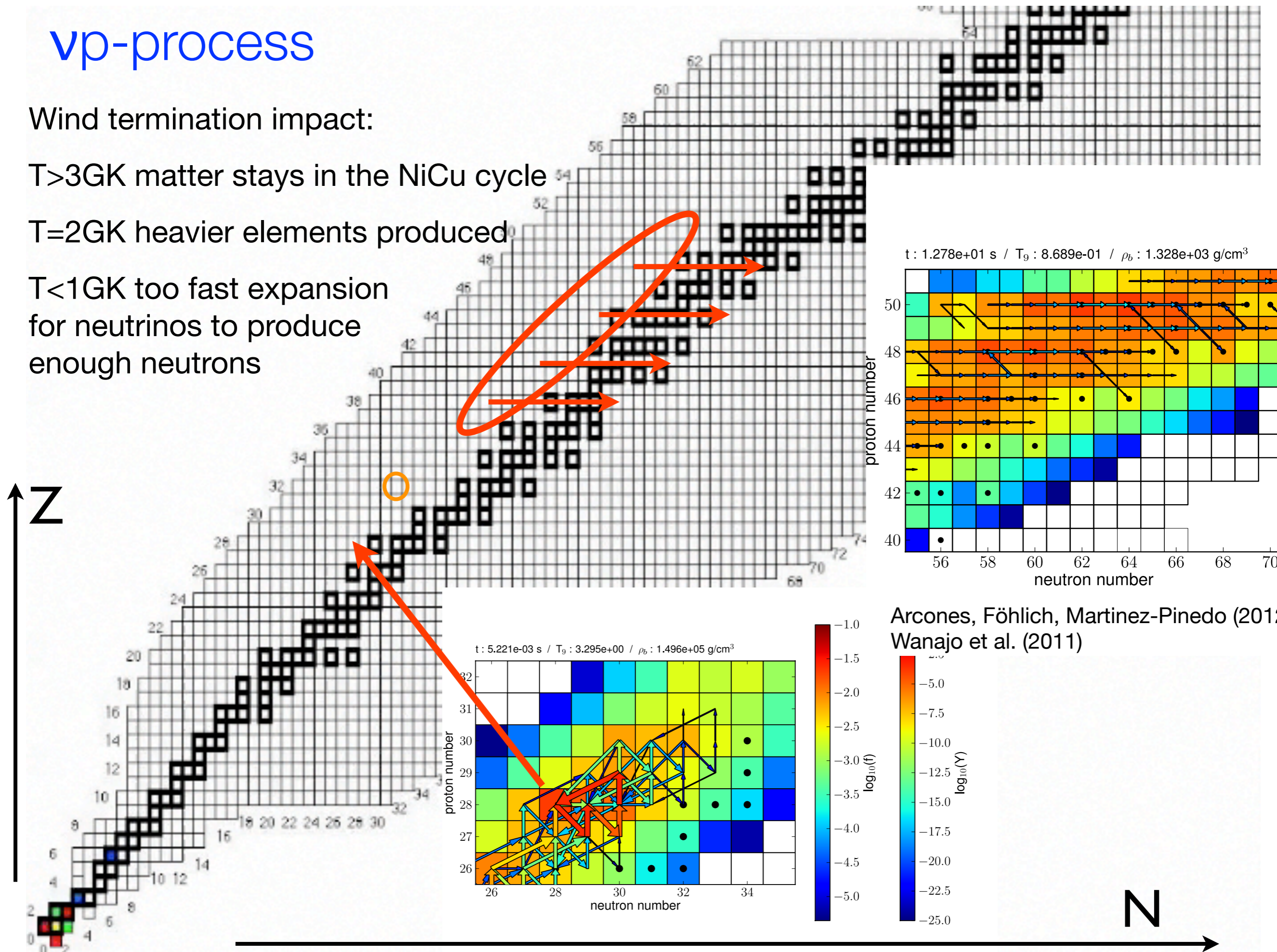
Wind termination impact:

$T > 3\text{GK}$  matter stays in the NiCu cycle

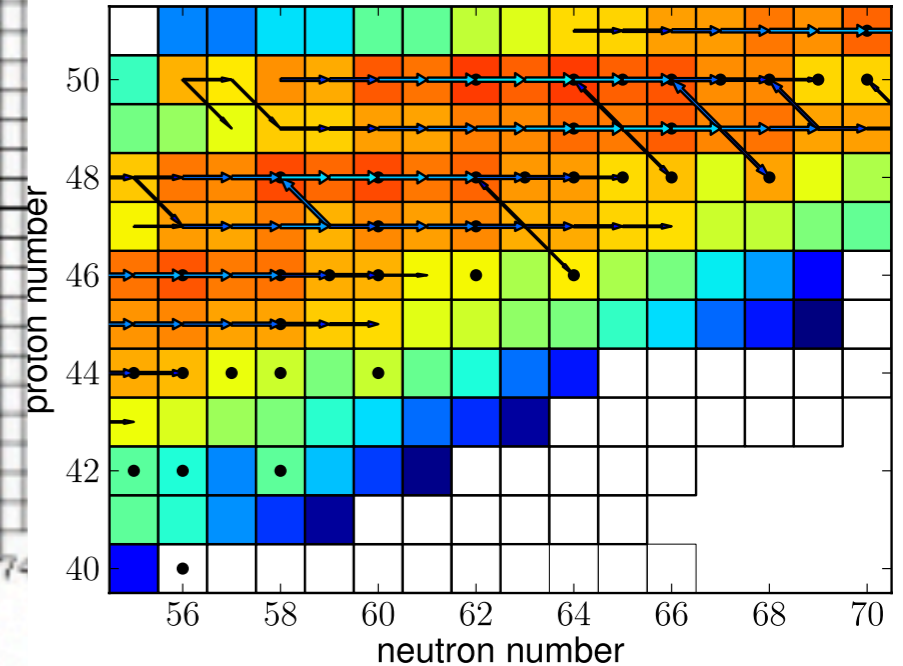
$T = 2\text{GK}$  heavier elements produced

$T < 1\text{GK}$  too fast expansion  
for neutrinos to produce  
enough neutrons

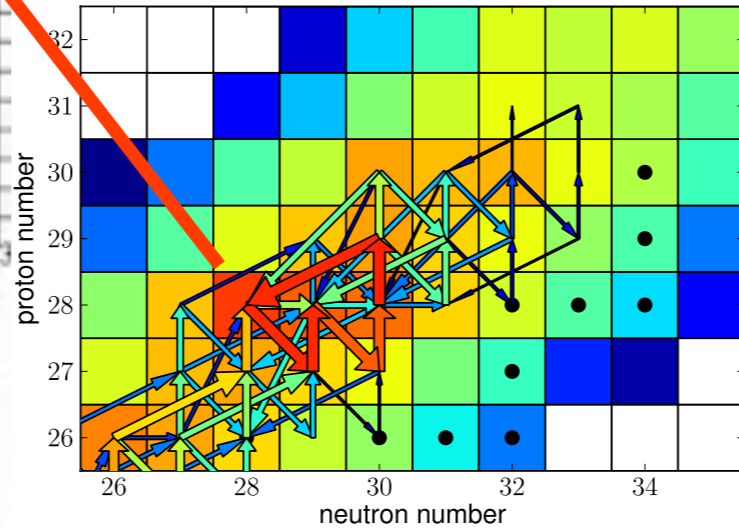
Z



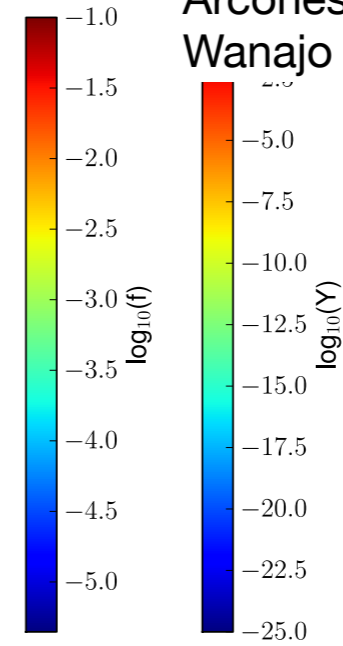
$t : 1.278\text{e}+01 \text{ s} / T_9 : 8.689\text{e}-01 / \rho_b : 1.328\text{e}+03 \text{ g/cm}^3$



$t : 5.221\text{e}-03 \text{ s} / T_9 : 3.295\text{e}+00 / \rho_b : 1.496\text{e}+05 \text{ g/cm}^3$



Arcones, Föhlich, Martinez-Pinedo (2012)  
Wanajo et al. (2011)



N

# vp-process

Wind termination impact:

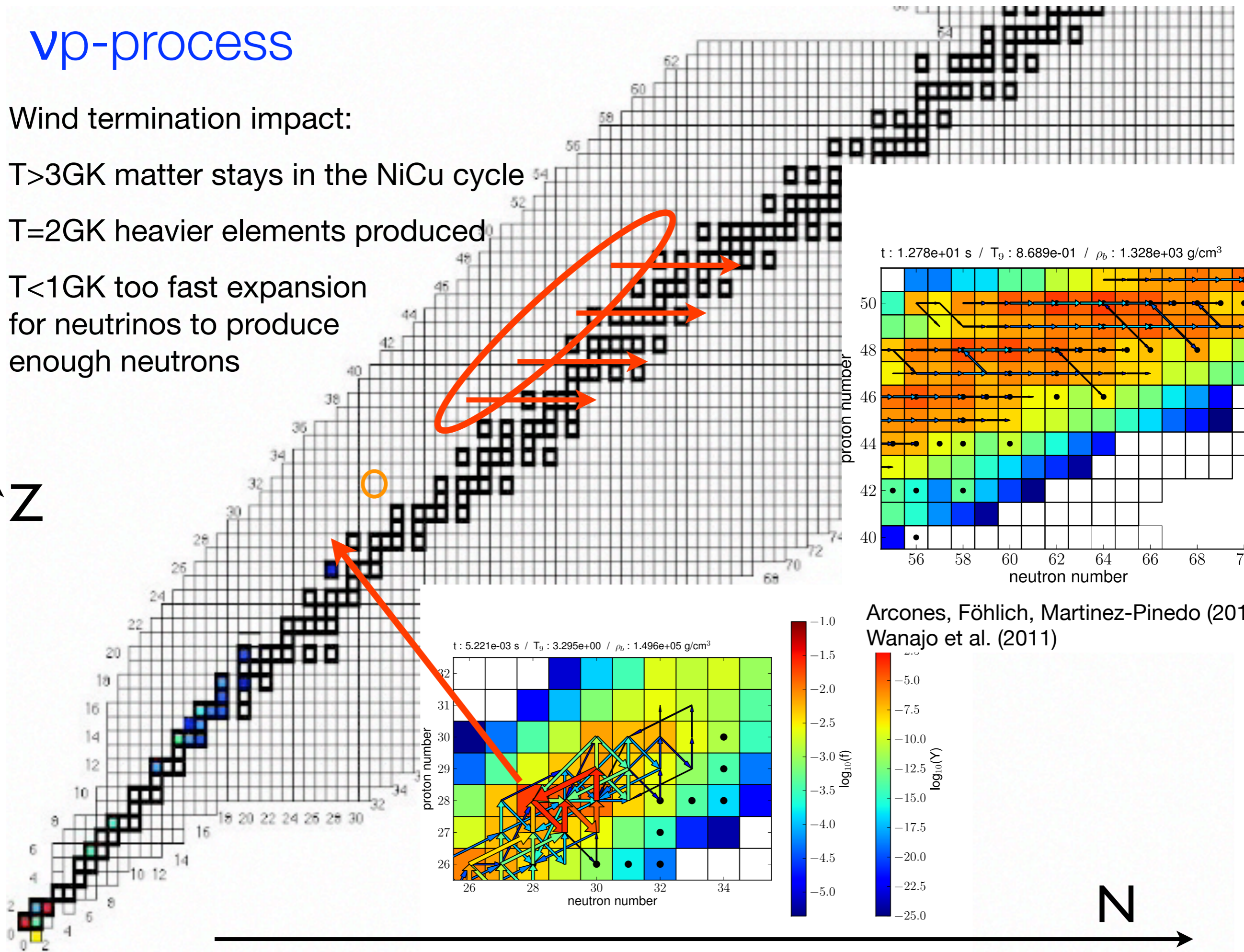
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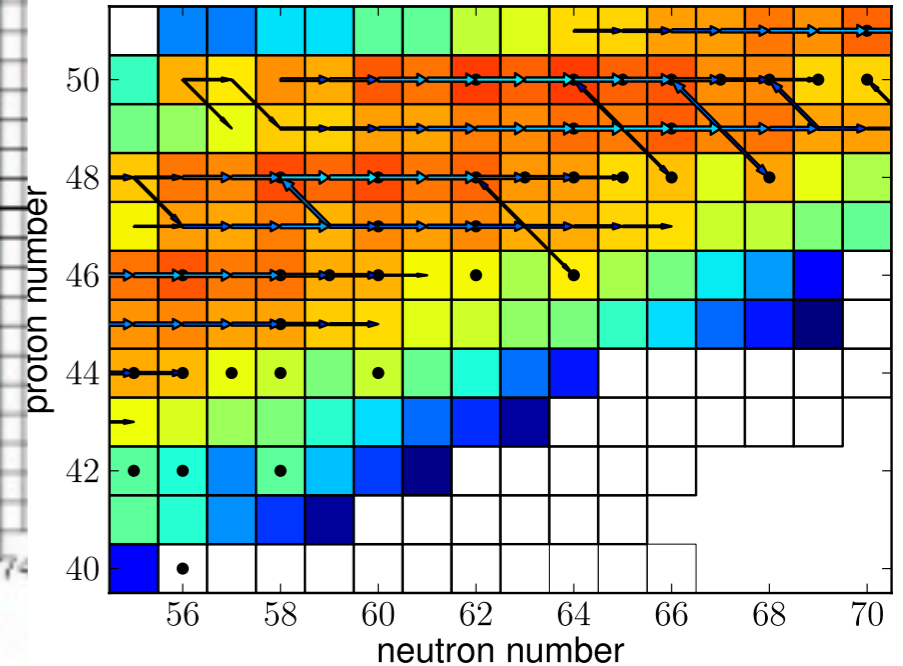
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**Z**  
↑

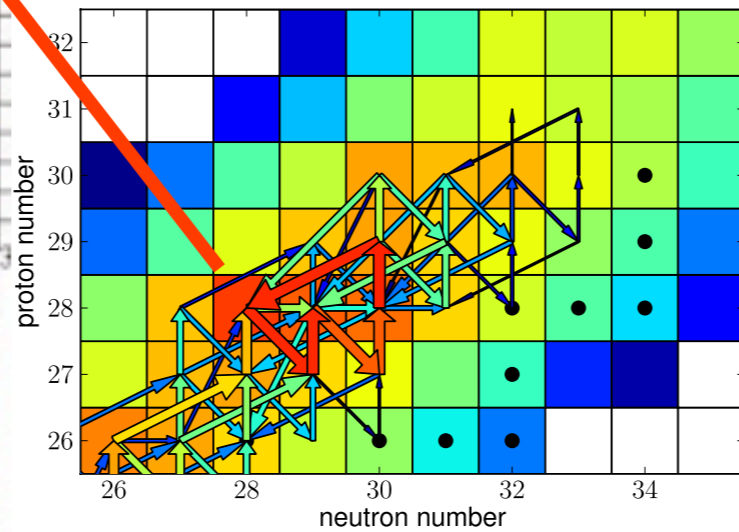
→ **N**



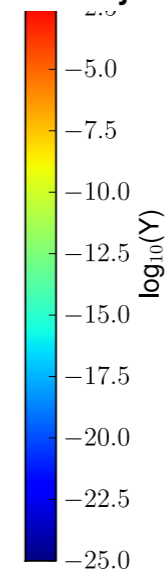
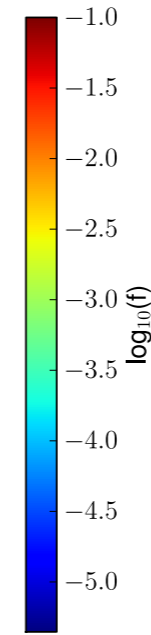
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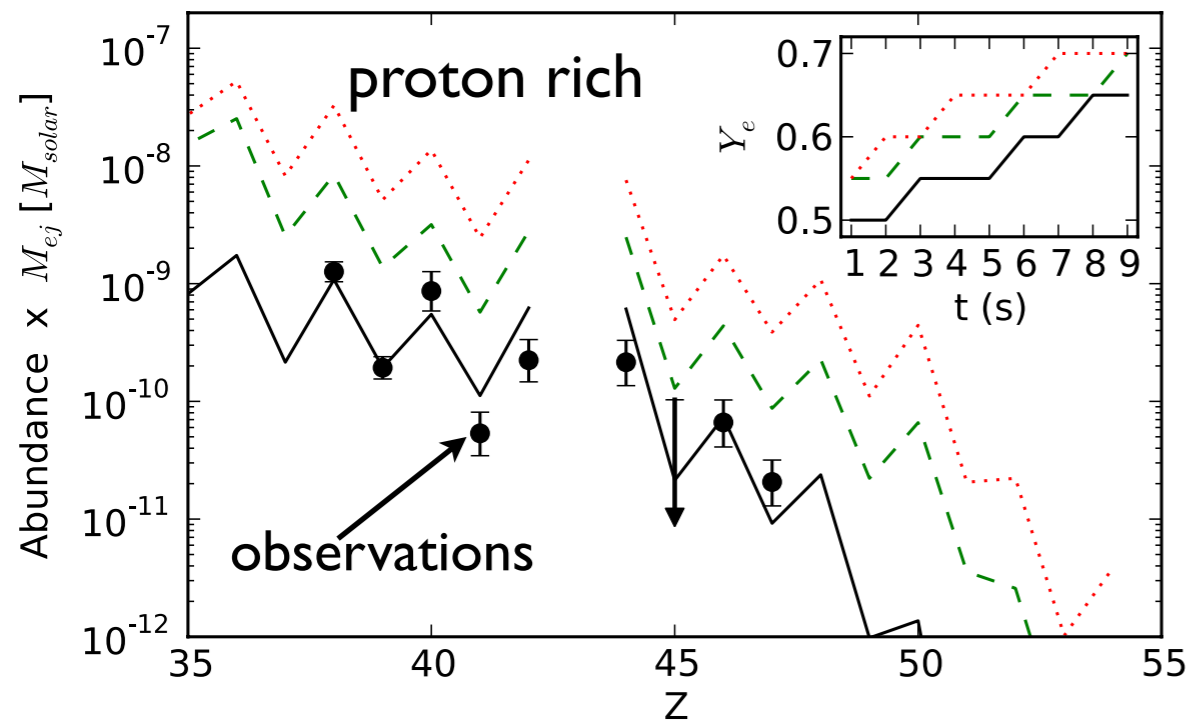


# Lighter heavy elements in neutrino-driven winds

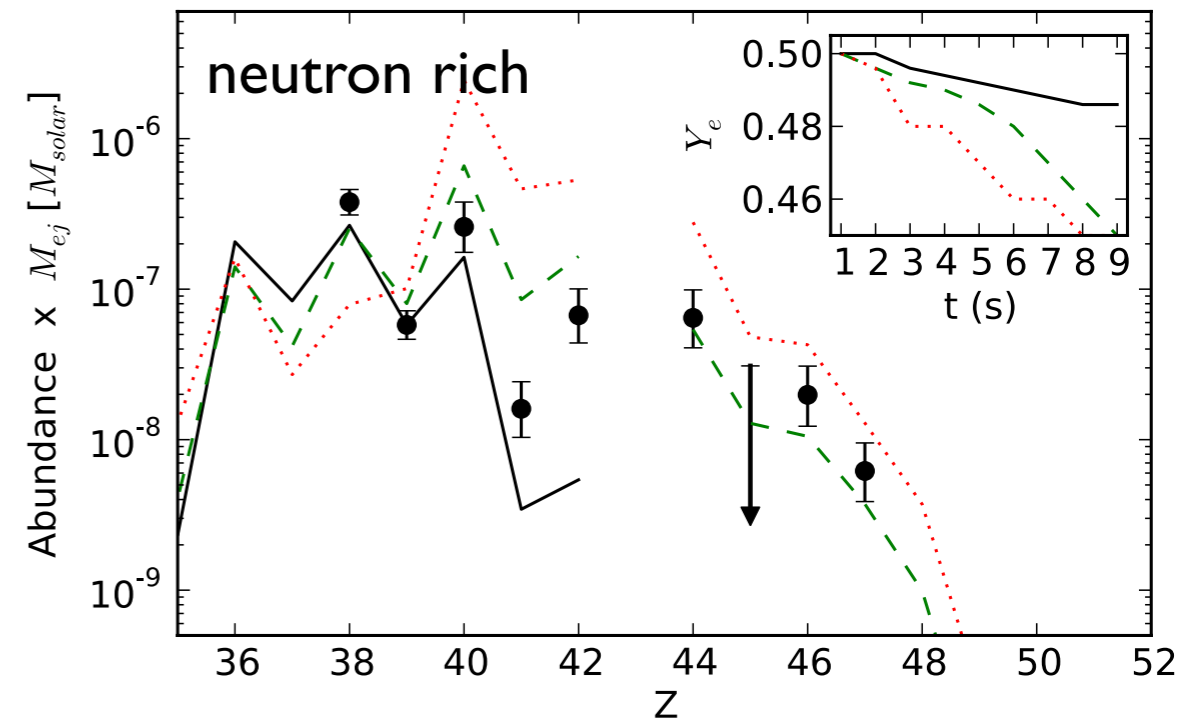
Can the LEPP pattern be produced based on neutrino-driven wind simulations?

Which nuclear process is the LEPP? Charged-particle reactions (Qian & Wasserburg 2001)

vp-process



weak r-process



Observation pattern can be reproduced!

Production of p-nuclei

Overproduction at  $A=90$ , magic neutron number  $N=50$  (Hoffman et al. 1996) suggests: only a fraction of neutron-rich ejecta

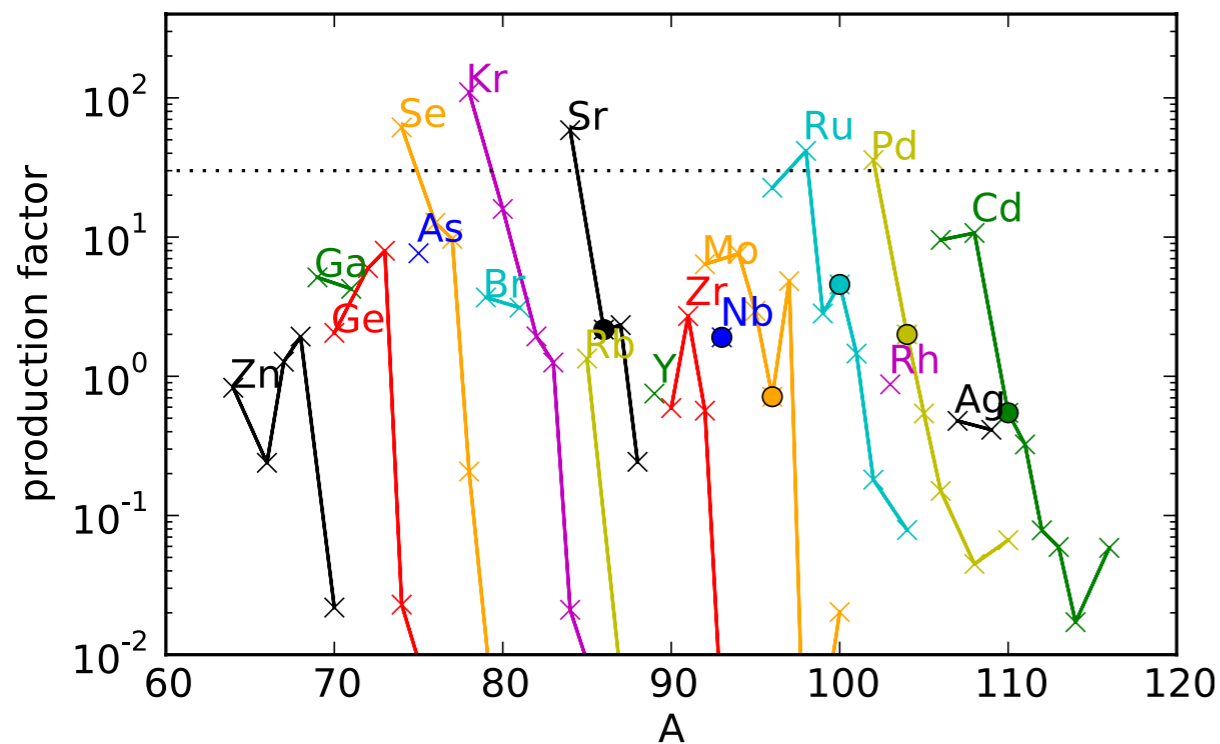
(Arcones & Montes, 2011)

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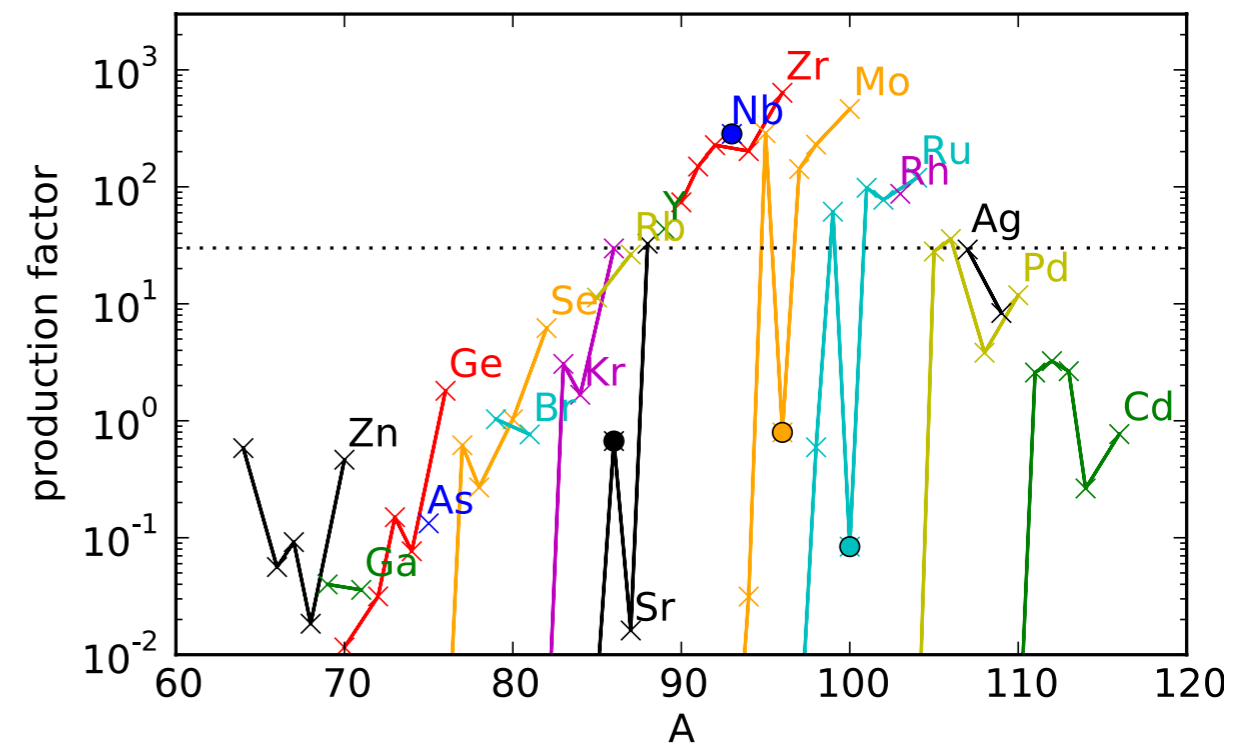
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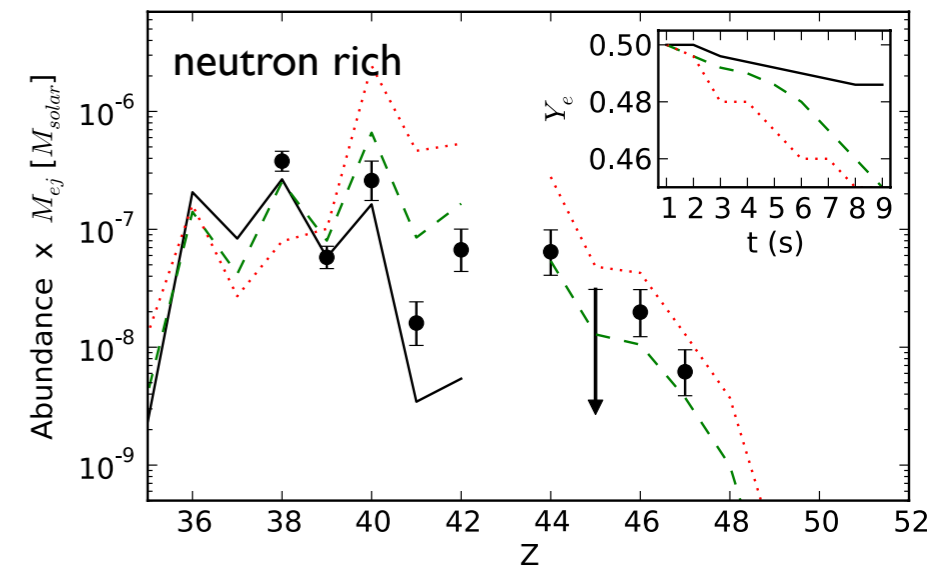
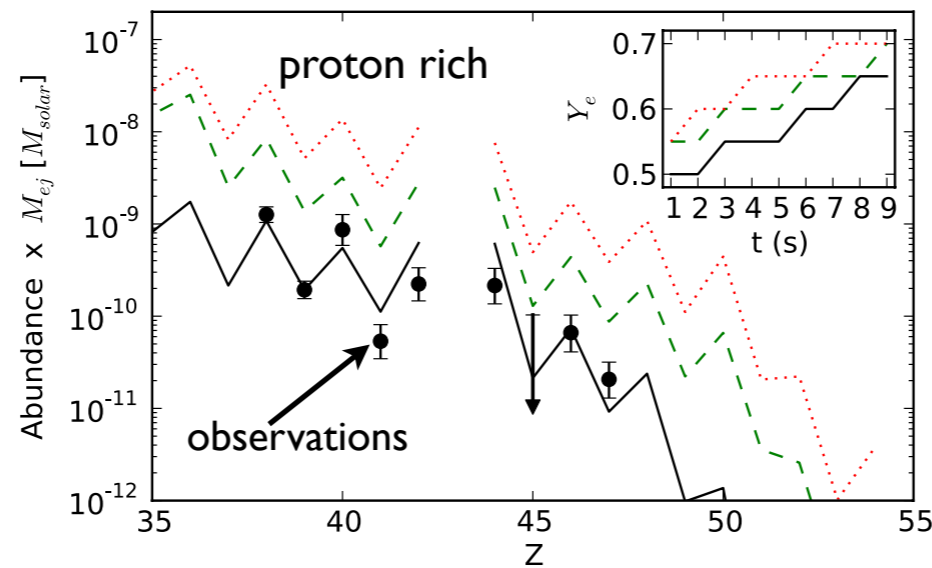


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(Arcones & Montes, 2011)

# Conclusion

LEPP pattern can be produced based on neutrino-driven wind simulations



LEPP = charged-particle reactions + vp-process  
weak r-process

Observations and better constraints on  $Y_e$  are required

Other possible LEPP sites: super-AGB stars at low Z (Herwig et al. 2011); fast rotating massive stars (Frischknecht et al. 2011)