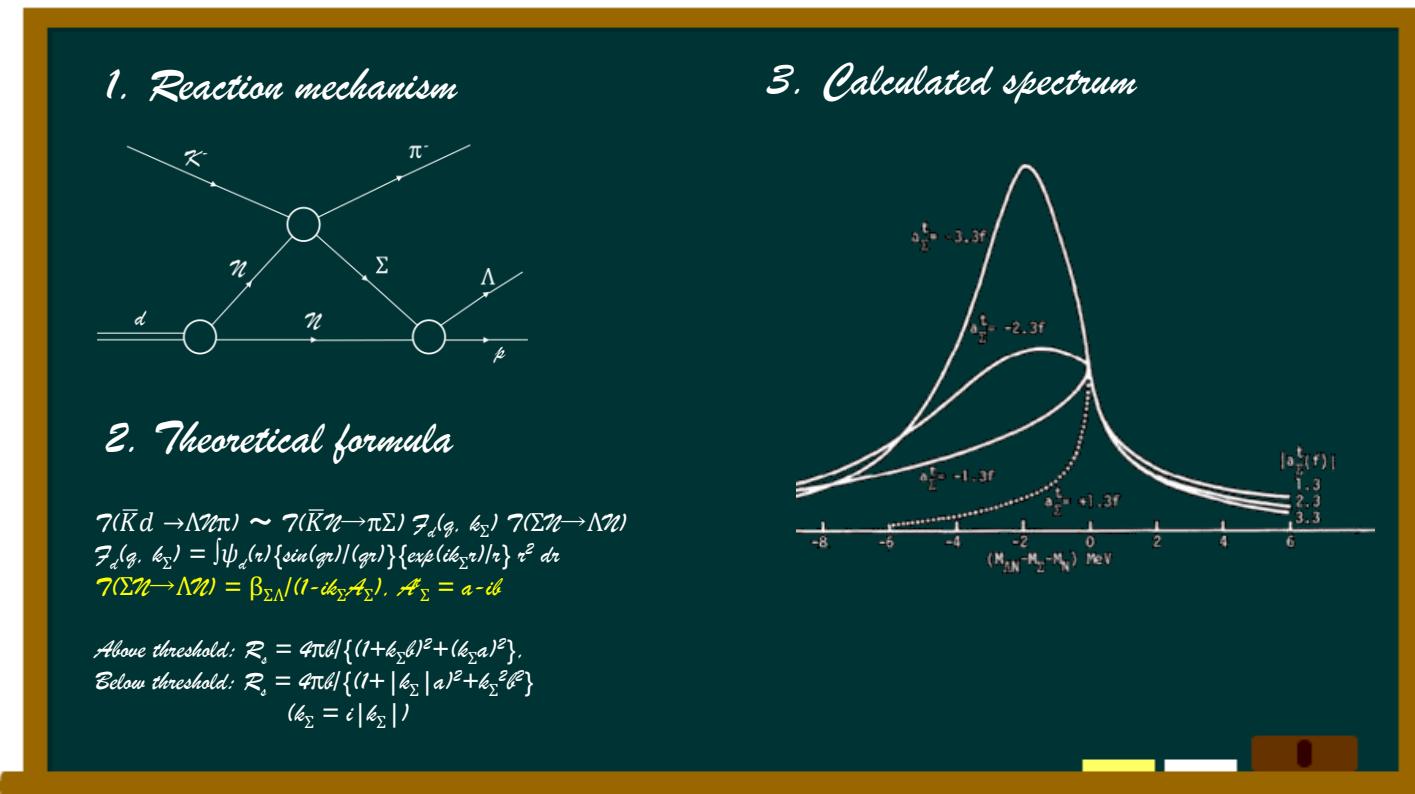


# J-PARC P90: HIGH RESOLUTION SPECTROSCOPY OF THE $\Sigma N$ CUSP BY USING $d(K^-, \pi^-)$ REACTION

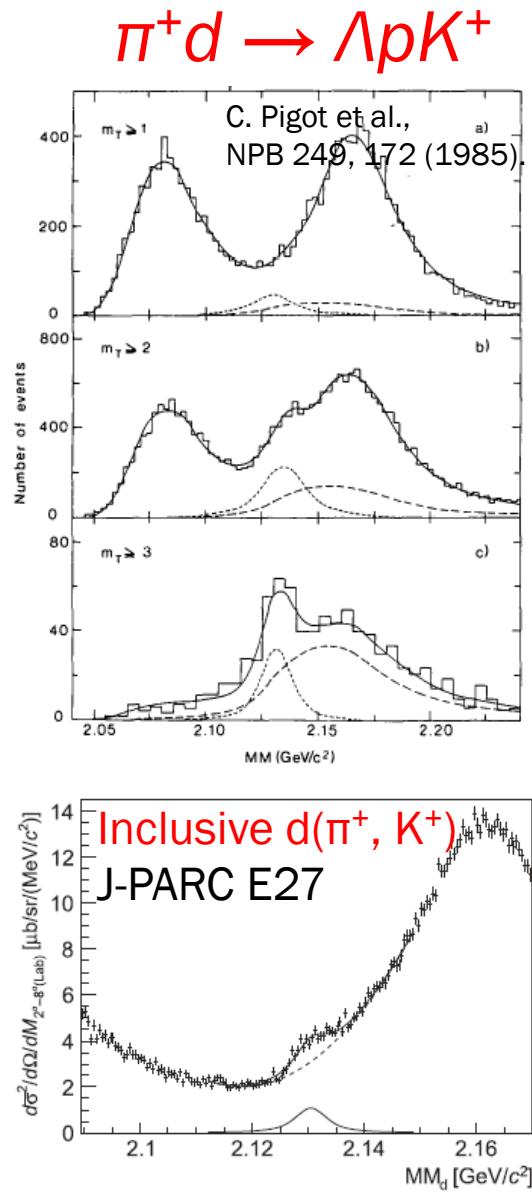
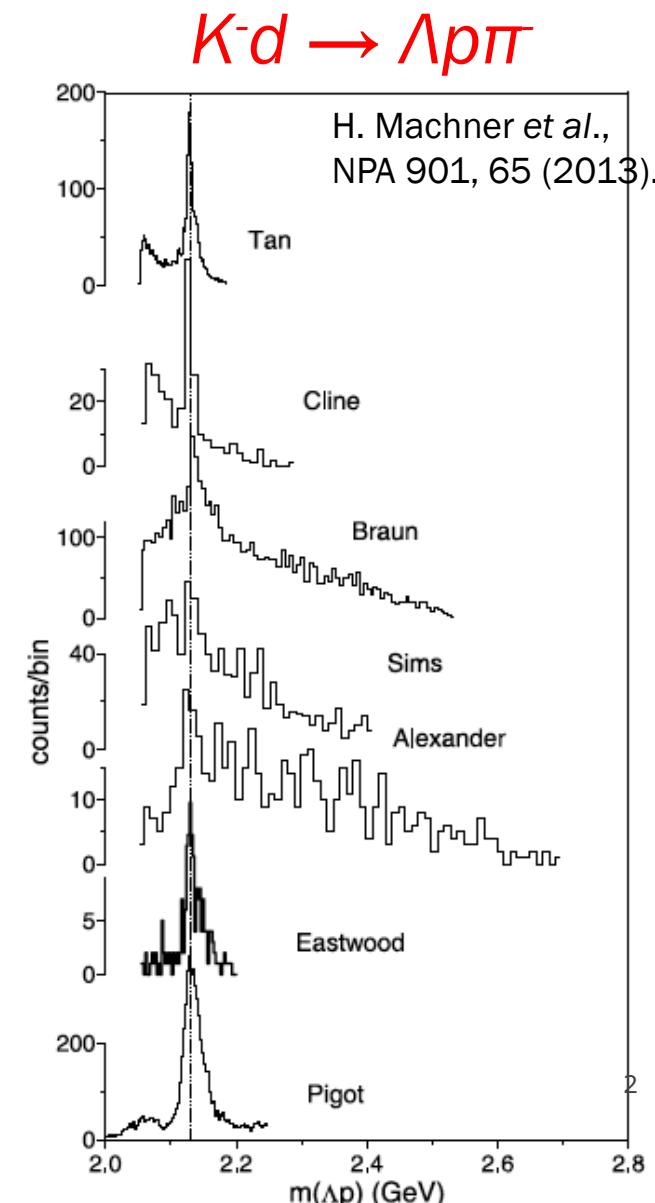
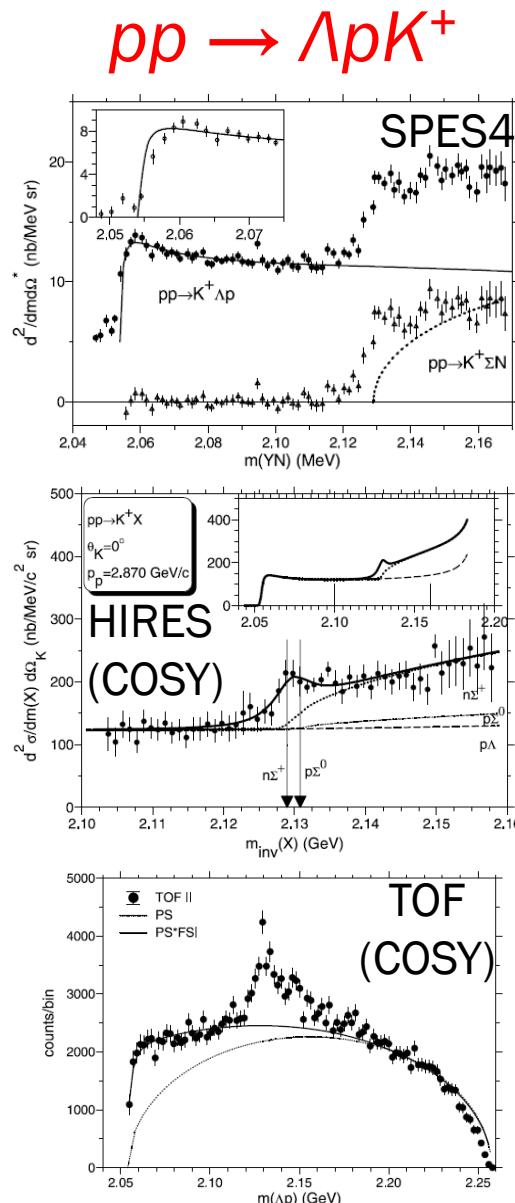
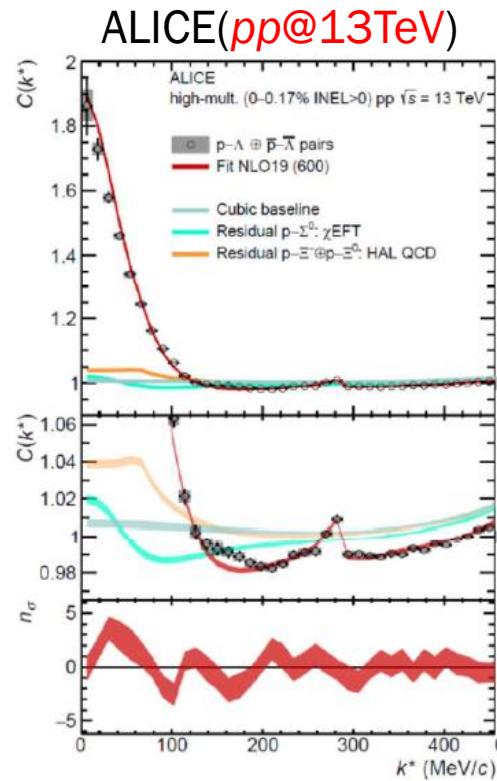
*YUDAI ICHIKAWA (JAEA)*

JOINT THEIA-STRONG2020 AND JAEA/MAINZ REIMEI WEB-SEMINAR 2022/01/26



# “ $\Sigma$ N CUSP”

*Clear enhancement  
around  $\Sigma$ N threshold  
(~2.13 GeV/c<sup>2</sup>)*



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## PURPOSE OF THE PROPOSED EXPERIMENT

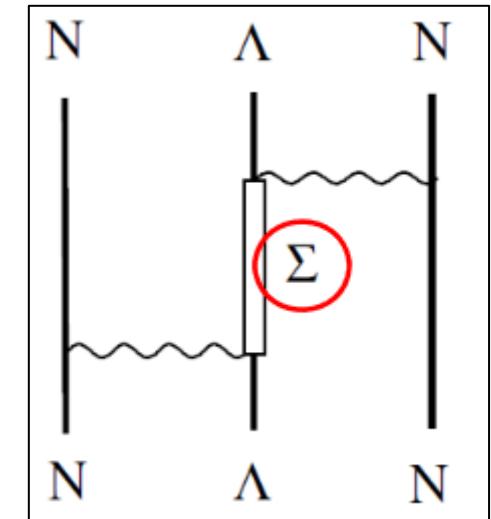
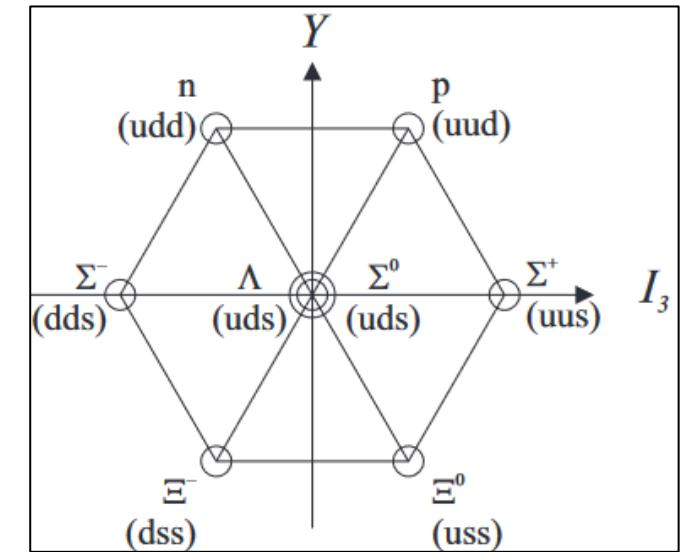
- Deduce the  $\Sigma N$  scattering length of  $(T, S) = (1/2, 1)$  channel by fitting “ $\Sigma N$  cusp” spectrum shape observed in the missing mass of the  $d(K^-, \pi^-)$  reaction.
  - Unstable bound state ( $\Sigma N$  dibaryon)? or Virtual state?

The key of this experiment is **the excellent missing-mass resolution** thanks to the **S-2S** spectrometer (used in E70) and high statistics. We will be able to achieve the best resolution of **0.4 MeV** in  $\sigma$ , which is **two times better** than the past experiment (HIRES at COSY).

# $\Sigma N$ INTERACTION

- $\Sigma N$  interaction is one of the key to understand  $B_8 B_8$  interaction in SU(3)
  - Relation with E40 ( $\Sigma N$  scattering experiment)
    - E40:  $\Sigma N$  scattering ( $p_\Sigma > 470 \text{ MeV}/c$ )  $\rightarrow$  Short range interaction  
\* $\Sigma N$  scattering experiment in lower momentum is difficult
    - P90: “ $\Sigma N$  cusp”  $\rightarrow$   $\Sigma N$  scattering length (0 energy interaction)
      - $\Sigma N$  scattering length:  $A_\Sigma = a + ib$
      - **a(real part)**  $\rightarrow$  Important for the  $\Sigma$ -hypernuclei
      - **b(imaginary part)**  $\rightarrow$   $\Lambda N - \Sigma N$  coupling strength
- Complementary
- ↓
- Important for the  $\Lambda$ -hypernuclei

Octet baryon



# THRESHOLD CUSP

Cusp structure can be expressed by the scattering length (for  $B'C'$ ),  $\underline{A} = \underline{a} + i\underline{b}$

- $B'C' \rightarrow BC$  amplitude

- $f_{B'C', BC} \sim \frac{\sqrt{b}}{1-ikA}$ , **Pole position:**  $k \sim -\frac{i}{A}$

- (Two body scattering amplitude)

- $f = \frac{1}{k} \sum_{l=0}^{\infty} (2l+1) e^{i\delta_l} \sin \delta_l P_l(\cos \theta)$
- $\rightarrow (s\text{-wave}) \quad f = \frac{1}{k} e^{i\delta_0} \sin \delta_0 = \frac{1}{k \cot \delta - ik} \rightarrow \frac{a}{1-ika}$
- $k \cot \delta = \frac{1}{a} - \left( \frac{r_{eff}}{2} \right) k^2 + \dots$

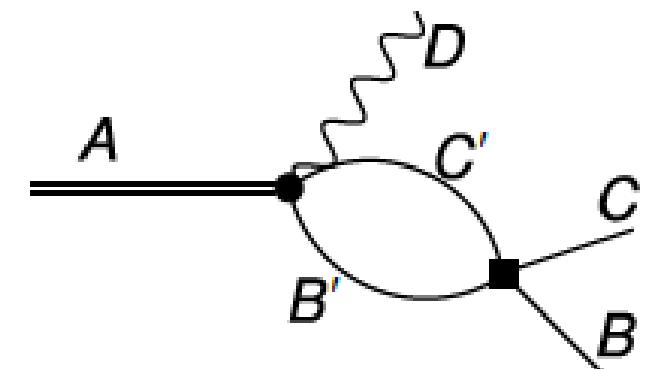
- Reaction rate ( $R$ ):  $\frac{d\sigma^2}{d\Omega dE} \propto |f_{B'C', BC}|^2$

- Above threshold:  $R = \frac{4\pi b}{\{(1+kb)^2 + (ka)^2\}}$

- Below threshold:  $R = \frac{4\pi b}{\{(1+\kappa a)^2 + (\kappa b)^2\}}$ ,  $k = i\kappa$  (due to analytic continuation)

Reduced mass  
 $\underline{\mu} = \underline{m}_{B'} \underline{m}_{C'}/(\underline{m}_{B'} + \underline{m}_{C'})$

$k$ (relative momentum for  $B'C'$ )  $\sim \sqrt{2\mu E}$

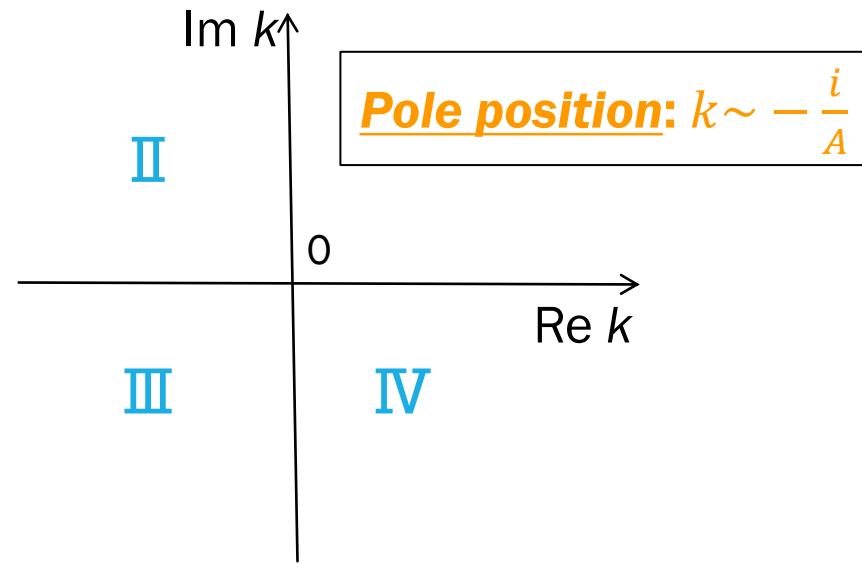
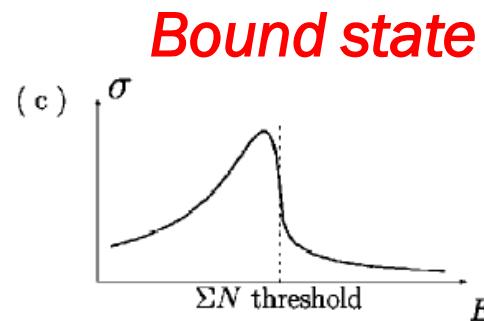
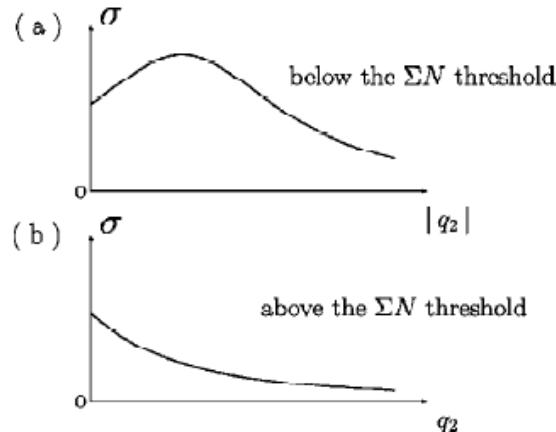


For the “ $\Sigma N$  cusp”,  
 $B'=\Sigma$ ,  $C'=N$ ,  $B=\Lambda$ ,  $C=N$

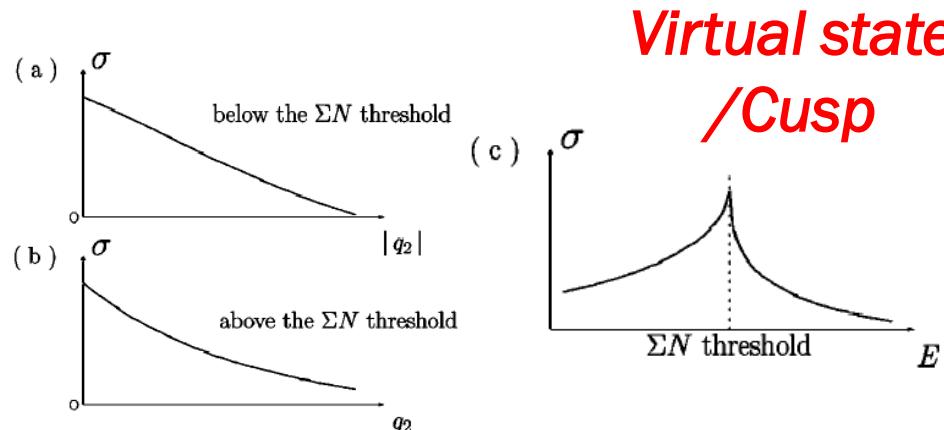
# POLE POSITION vs CROSS SECTION ( $d\sigma/dE$ )

K. Miyagawa and H. Yamamura, PRC 60, 024003 (1999).

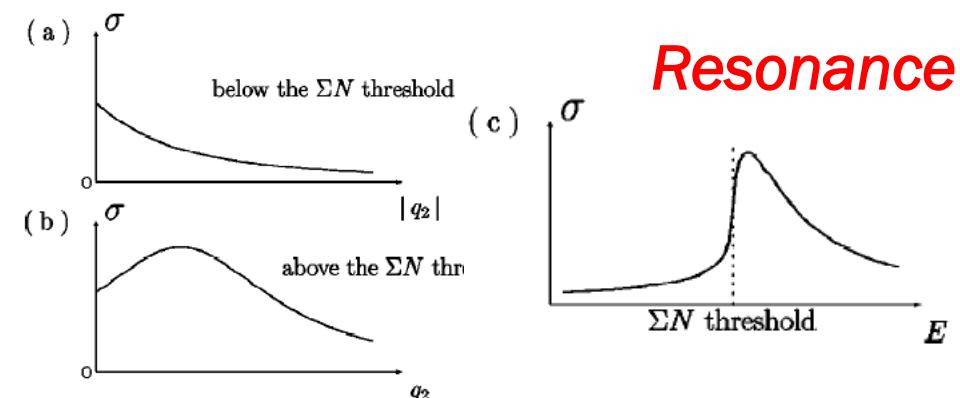
- Pole in (II) quad



- Pole in (III) quad



- Pole in (IV) quad



# THRESHOLD CUSP

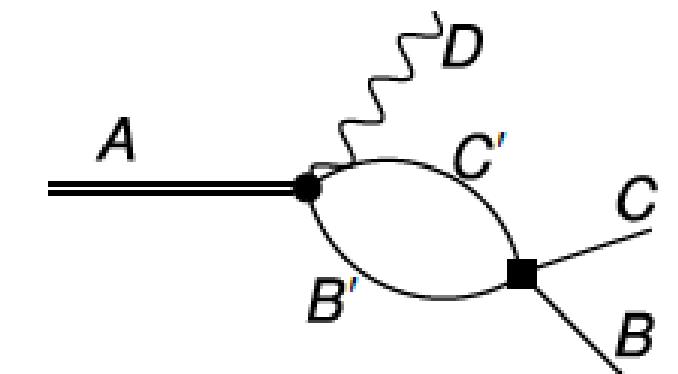
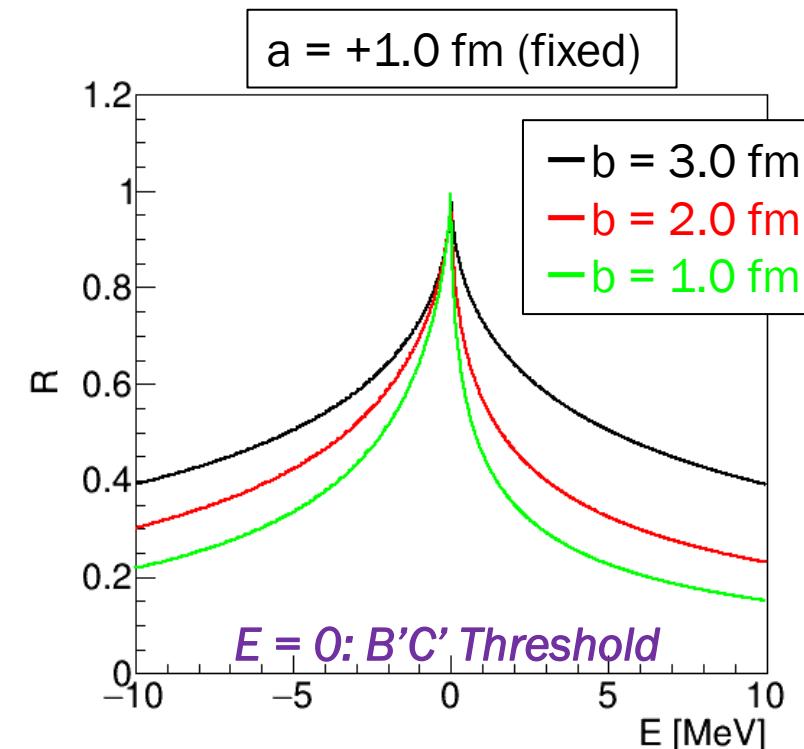
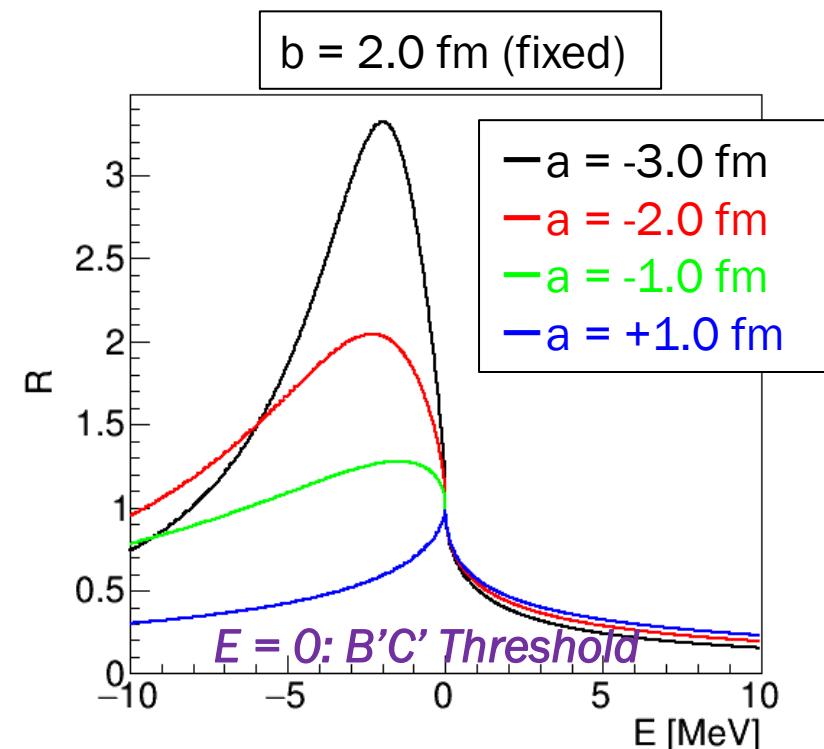
Cusp structure can be expressed by the scattering length (for  $B'C'$ ),  $\mathbf{A} = \mathbf{a} + i\mathbf{b}$

- Above threshold:**  $R = \frac{4\pi b}{\{(1+kb)^2+(ka)^2\}} \sim 1 - 2kb + O(k^2)$

- Below threshold:**  $R = \frac{4\pi b}{\{(1+\kappa a)^2+(\kappa b)^2\}} \sim 1 - 2\kappa a + O(\kappa^2), k = i\kappa$

Reduced mass  
 $\mu = m_{B'}m_{C'}/(m_{B'}+m_{C'})$

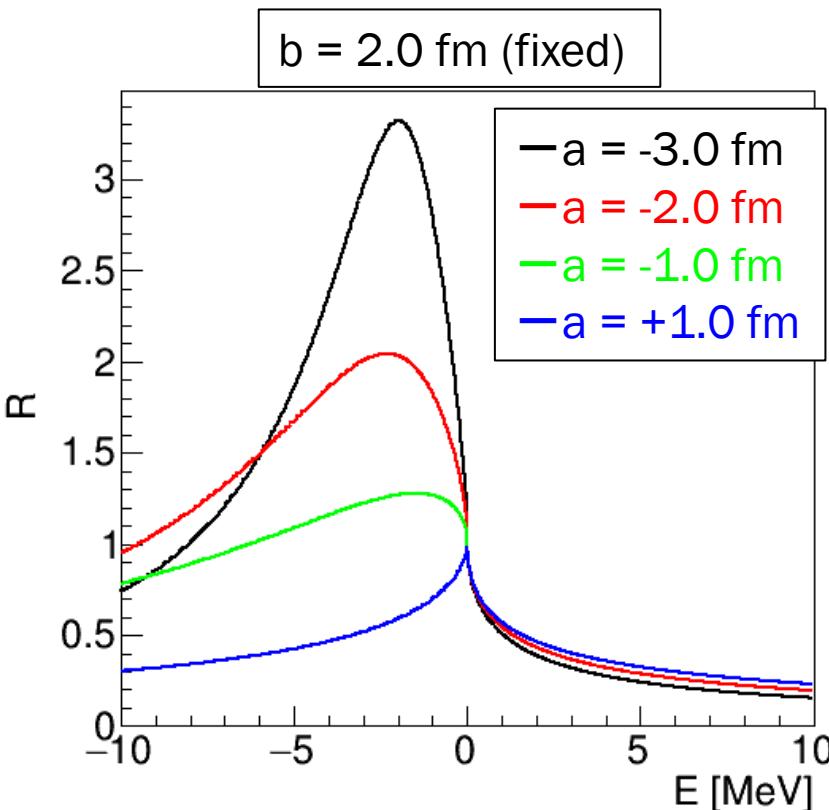
$k(\text{relative momentum for } B'C') \sim \sqrt{2\mu E}$



For the “ $\Sigma N$  cusp”,  
 $B'=\Sigma, C'=N, B=\Lambda, C=N$

# “ $\Sigma N$ CUSP”

- “ $\Sigma N$  cusp” is measured by  $K^-d \rightarrow \pi^-\Lambda p$  reaction etc..
  - T:  $T = 1/2$  ( $\Lambda p$  final state)
  - S:  ${}^3S_1$  is favored, D-target; observed in forward angles



**Above threshold:**  $R = \frac{4\pi b}{\{(1+kb)^2+(ka)^2\}} \sim 1 - 2kb + O(k^2)$

**Below threshold:**  $R = \frac{4\pi b}{\{(1+\kappa a)^2+(\kappa b)^2\}} \sim 1 - 2\kappa a + O(\kappa^2), k = i\kappa$   
( $k \sim \sqrt{2\mu E}$ )

“ $\Sigma N$  Cusp” can be expressed by the  $\Sigma N$  scattering length ( $A_\Sigma = a + ib$ ) of the  $(T, S) = (1/2, {}^3S_1)$  channel!!

# “ $\Sigma N$ cusp” $K^-d \rightarrow \pi^- \Lambda p$ reaction

- Amplitude of the elementary reaction
- $f(\theta) \propto T(\bar{K}d \rightarrow \Lambda N \pi) \sim T(\bar{K}N \rightarrow \pi \Sigma) F_d(\vec{Q}_\Sigma, k_\Sigma) T(\Sigma N \rightarrow \Lambda N)$

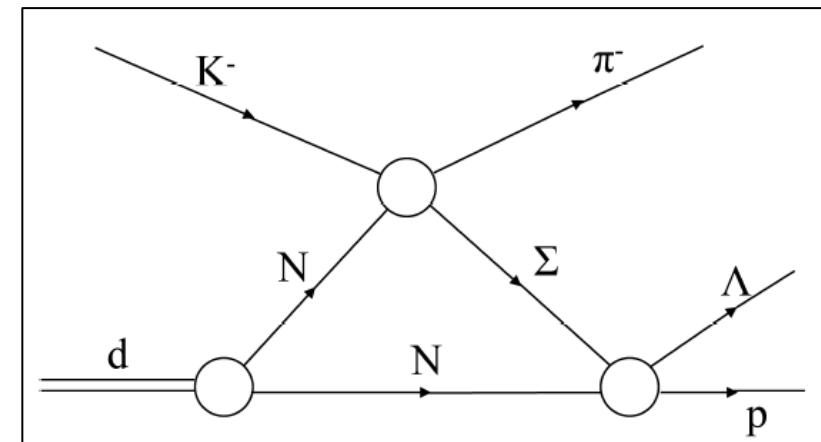
**Deuteron factor**

$$F_d(Q_\Sigma, k_\Sigma) = \int \frac{e^{ik_\Sigma r}}{r} e^{i\vec{Q}_\Sigma \cdot \vec{r}} \psi_d(r) d^3 r$$

$$\vec{Q}_\Sigma = \vec{q}m_N / (m_N + m_\Sigma)$$

$\psi_d(r)$  : deuteron wave function

R.H. Dalitz, Nucl. Phys. A354, 101 (1981).



**Amplitude of  $\Sigma N \rightarrow \Lambda N$  reaction (Important term)**

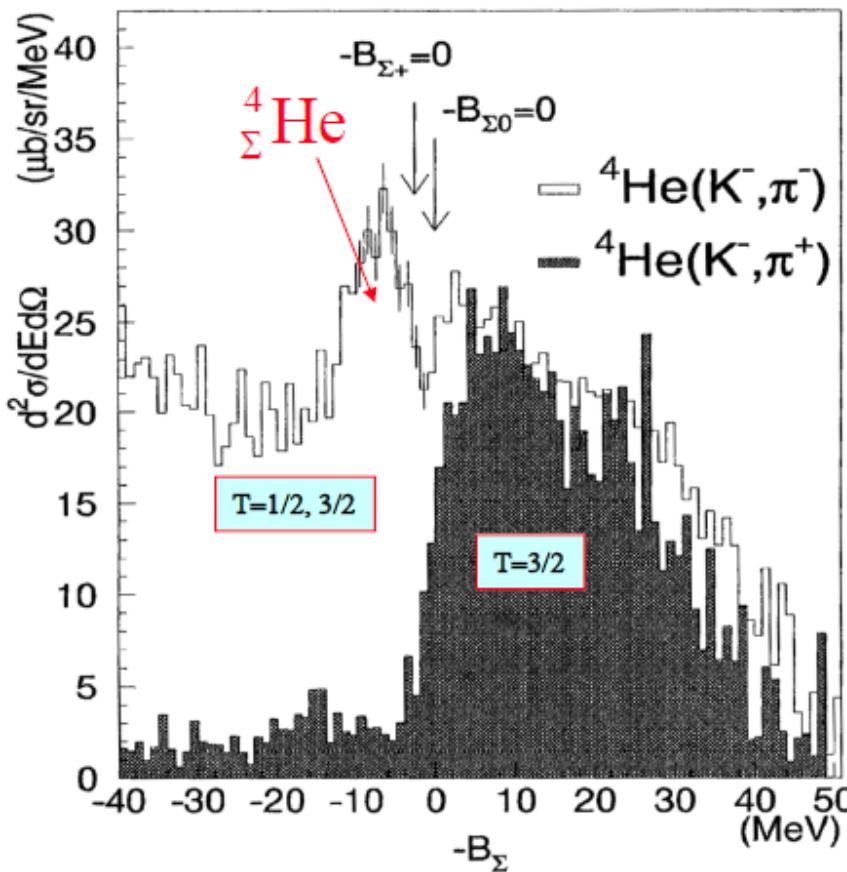
$$T(\Sigma N \rightarrow \Lambda N) \propto \frac{\sqrt{b}}{1 - ikA}$$

**Above threshold:**  $R = \frac{4\pi b}{\{(1+kb)^2 + (ka)^2\}} \sim 1 - 2kb + O(k^2)$

**Below threshold:**  $R = \frac{4\pi b}{\{(1+\kappa a)^2 + (\kappa b)^2\}} \sim 1 - 2\kappa a + O(\kappa^2)$ ,  $k = ik$   
 $(k \sim \sqrt{2\mu E})$

**“ $\Sigma N$  Cusp” can be expressed by the  $\Sigma N$  scattering length ( $A_\Sigma = a + ib$ ) of the  $(T, S) = (1/2, {}^3S_1)$  channel!!**

# IMPORTANCE OF $(T, S) = (1/2, 1)$ CHANNEL IN ${}^4\Sigma$ He



${}^4\text{He}(\text{K}^-, \pi^+)$ :  $\frac{3}{2}$  0  
 ${}^4\text{He}(\text{K}^-, \pi^-)$ :  $\frac{1}{2}$  0

$\Sigma N$  interaction has strong T and S dependence.

$V_{T=3/2, S=1}$  and  $V_{T=1/2, S=0}$  are expected to be repulsive due to quark Pauli-blocking effect.  
 → No  $\Sigma$ -hypernuclei in large A system.

$\Sigma N$  cusp channel:  $V_{T=1/2, S=1}$   
 (expected to be attractive potential and origin for the  ${}^4\Sigma$ He bound state)

$\Sigma N N N$  (4body) system

$T$		$S$		$\bar{V}_{\Sigma N}$ ( $V_{TS}$ )	
$\frac{3}{2}$	0	$\frac{5}{18} V_{\frac{3}{2}0}$	$+\frac{1}{2} V_{\frac{3}{2}1}$	$+\frac{2}{9} V_{\frac{1}{2}0}$	
$\frac{1}{2}$	0	$\frac{4}{9} V_{\frac{3}{2}0}$		$+\frac{1}{18} V_{\frac{1}{2}0}$	$+\frac{1}{2} V_{\frac{1}{2}1}$

## $\Sigma N(T=1/2, {}^3S_1)$ SCATTERING LENGTH (THEORY), $A_\Sigma = a + ib$

Model	J04	J04c	J-A	NSC 97f	NSC 89	ND	NF	NB
a [fm]	3.83	3.63	-2.37	-1.03	2.54	2.06	-1.29	-3.0
b [fm]	3.01	3.09	3.74	2.41	0.26	4.64	3.02	1.8
Model	chiral EFT (NLO13)					chiral EFT (NLO19)		
$\Lambda$ [MeV]	500	550	600	650	500	550	600	650
a [fm]	-2.61	-2.44	-2.27	-2.06	-0.95	-0.98	-2.29	-1.95
b [fm]	2.89	3.11	3.29	3.59	4.77	4.59	3.39	3.38

$a > 0$ : Attractive

$a < 0$ : Bound state

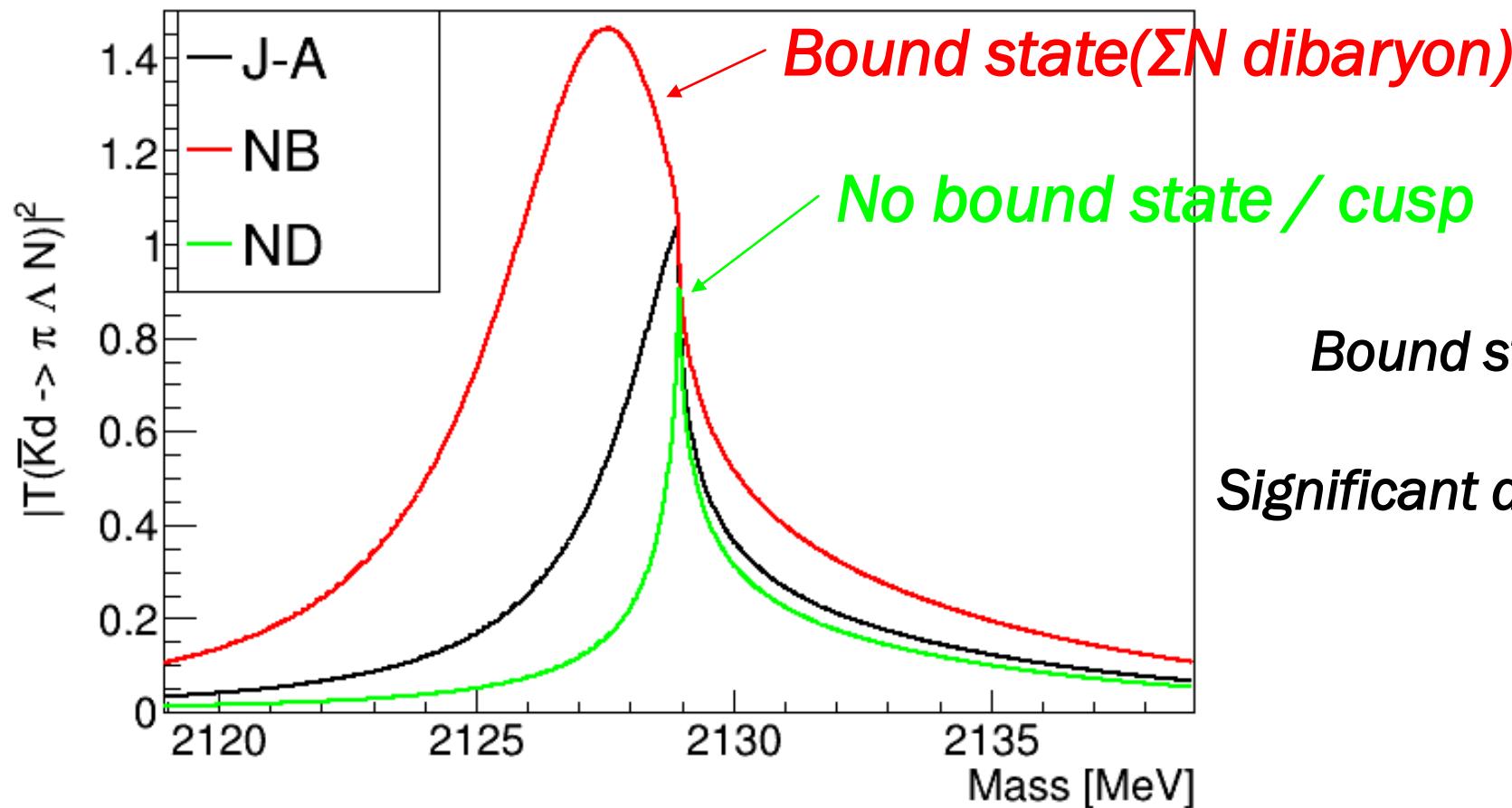
Large ambiguity!!

# $\Sigma N$ ( $T=1/2, {}^3S_1$ ) SCATTERING LENGTH (THEORY)

**Shallow bound state/cusp:** e.g. J-A ( $A_\Sigma = -2.37 + i3.74$  fm)

**Deeply bound state ( $\sim$ BW):** e.g. NB ( $A_\Sigma = -3.00 + i1.8$  fm)

**No bound state/cusp:** e.g. ND ( $A_\Sigma = 2.06 + i4.64$  fm)



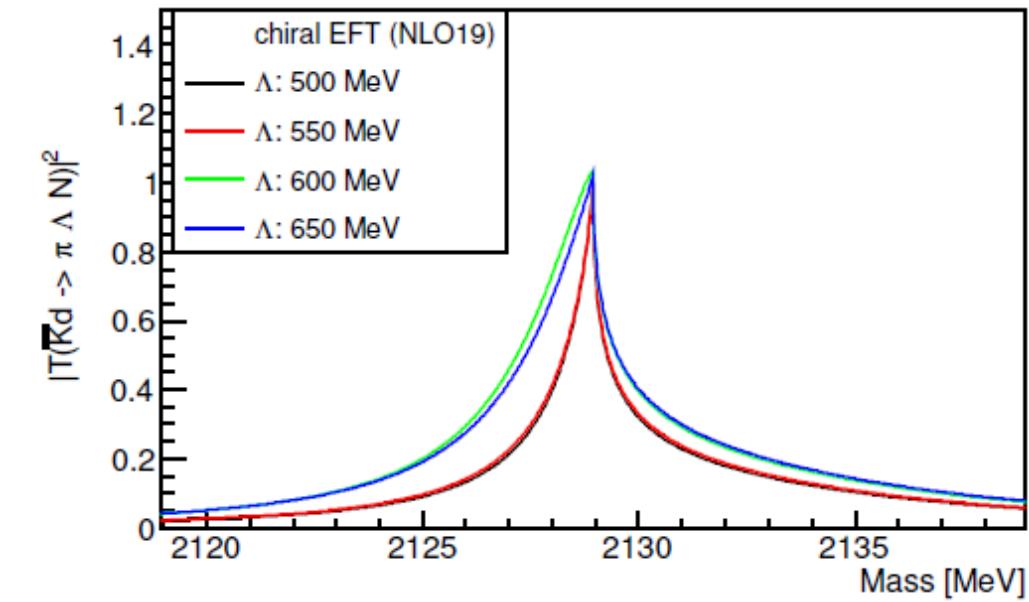
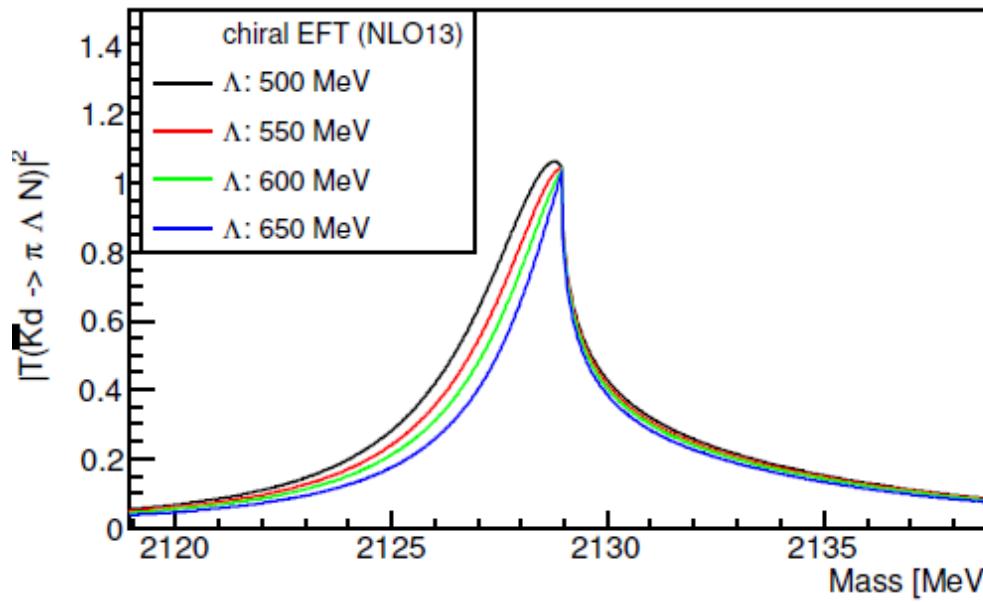
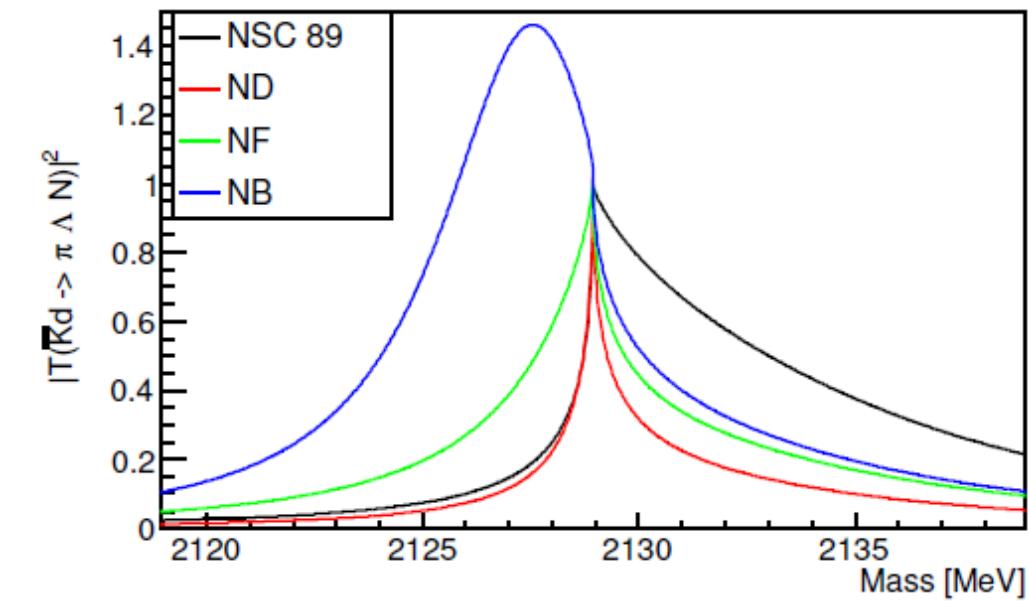
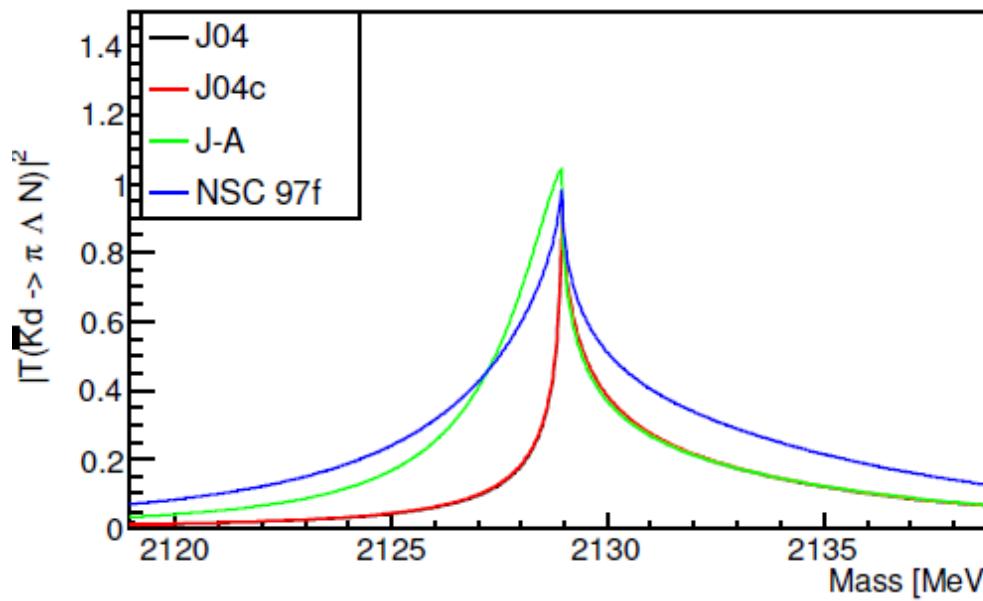
Pole position:  $k \sim -i/A_\Sigma$

Bound state (deuteron like dibaryon)

or not?

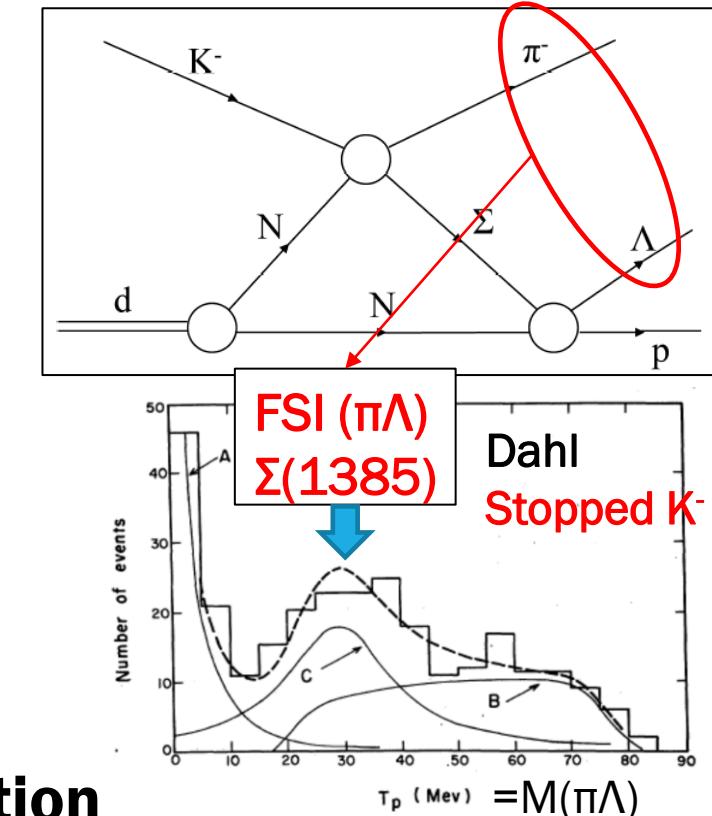
Significant difference in the spectrum shape

# $\Sigma$ (T=1/2, $^3S_1$ ) SCATTERING LENGTH (THEORY)



# REQUIREMENTS FOR THE “ $\Sigma N$ CUSP” EXPERIMENT

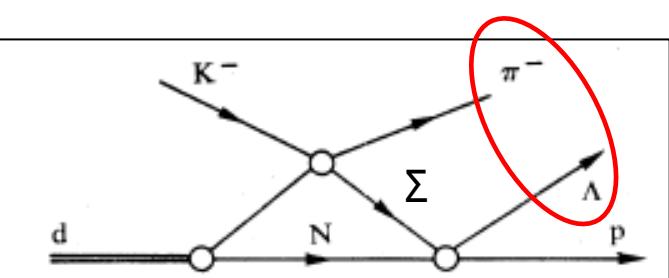
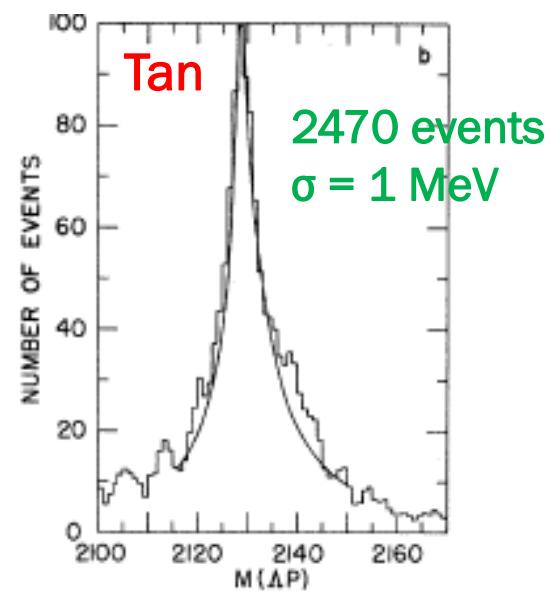
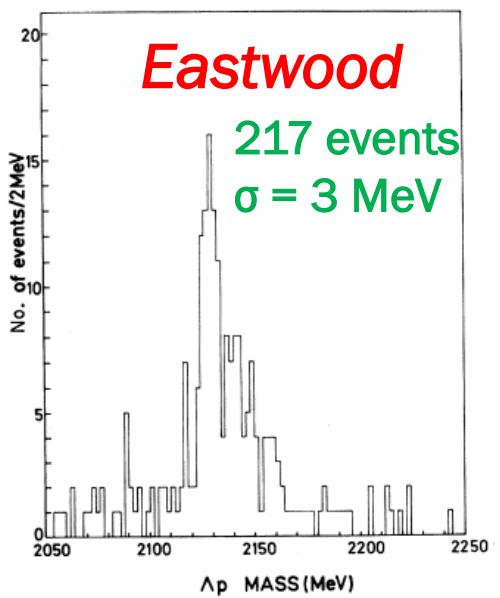
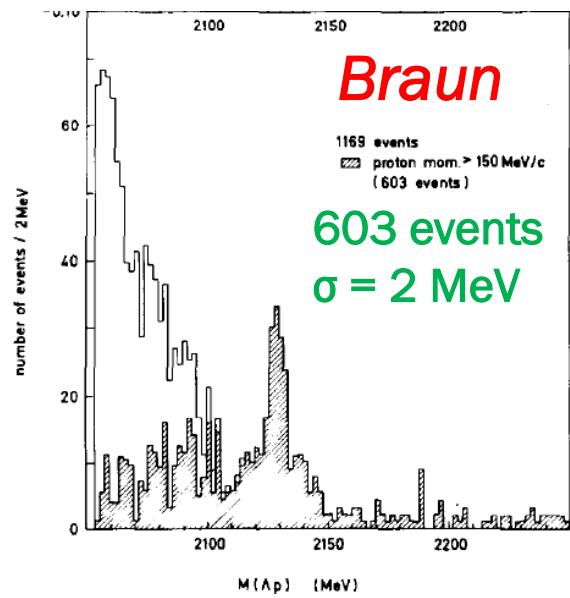
- Good energy resolution ( $\sigma < 1$  MeV)
- High statistics ( $> 10^4$  events)
- Good Signal / Noise (S/N) ratio
- Avoid FSI except for the  $\Sigma N$ 
  - $\times$  Stopped  $K^-d \rightarrow \Lambda p \pi^-$  reaction ( $p_\pi \sim p_\Lambda \sim p_p$ )  
FSI:  $\pi\Lambda$ ,  $\pi p$ ,  $YN$  ( $YN$  FSI =  $\Sigma N$  cusp signal)
  - $\circ$  In-flight  $K^-d \rightarrow \Lambda p \pi^-$  reaction ( $p_K \sim p_\pi \gg p_\Lambda \sim p_p$ )  
FSI:  $YN$  ( $YN$  FSI =  $\Sigma N$  cusp signal),  $\circ$  impulse approximation
- Decompose  $^1S_0$  and  $^3S_1$  contribution  
( $K^-d \rightarrow \Lambda p \pi^-$  reaction: extract only  $^3S_1$  contribution by D-target property)



*There was no experiment to satisfy these requirements!!*

## Bubble chamber

	Reaction	Comments	Statistics	Resolution
Braun	Inflight $d(K^-, \pi^-) \Lambda p$ 680 – 840 MeV/c	Low statistic, worse resolution	603 events ( $\cos \theta > 0.9$ , momcut)	2 MeV
Eastwood	Inflight $d(K^-, \pi^-) \Lambda p$ 1450, 1650 MeV/c	Low statistic worse resolution	217 events ( $\cos \theta > 0.9$ , momcut)	3 MeV
Tan	stopped $d(K^-, \pi^-) \Lambda p$	Large FSI	2470 events	1 MeV
Pigot	Inflight $d(K^-, \pi^-), d(\pi^+, K^+)$	Poor resolution	Uncertain	9.1 MeV ( $d(K^-, \pi^-)$ 1.4 GeV/c)
$pp \rightarrow \Lambda p K^+$ (COSY etc)	$pp \rightarrow \Lambda p K^+$	$^1S_0 + ^3S_1$ admixture Worse SN	High	0.8 MeV
ALICE	pp (Femtoscopy)	$^1S_0 + ^3S_1$ admixture	High	No description
J-PARC E27	$d(\pi^+, K^+)$ (Inclusive)	Worse SN (inclusive)	High	1.4 MeV



**Stopped  $K^-$  reaction**

- Multiple  $K^-$  scattering
- $\Sigma \Lambda, \pi p$  ( $p_\pi \sim p_\Lambda \sim p_p$ )

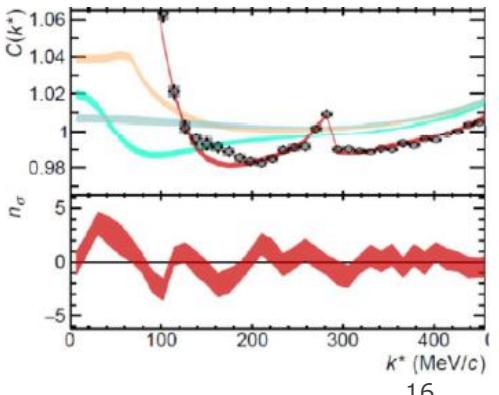
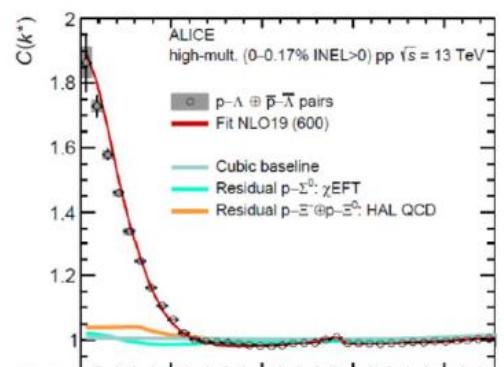
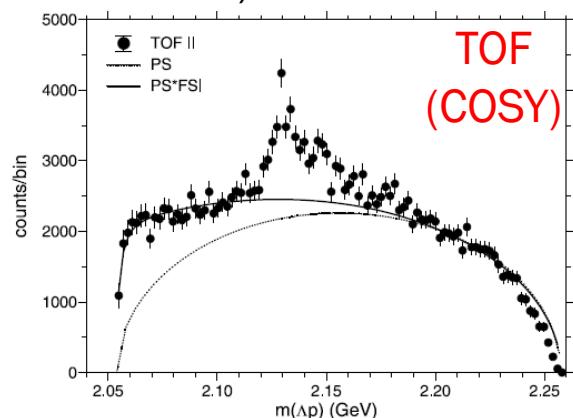
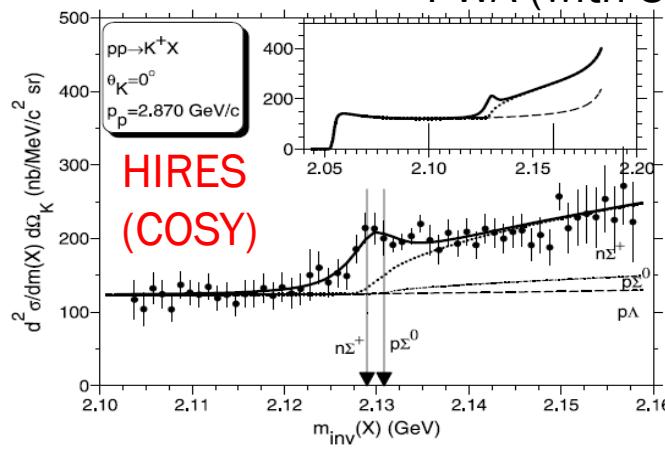
[YN FSI = Signal]

## Bubble chamber

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$pp \rightarrow \Lambda p K^+$ (COSY etc)	$pp \rightarrow \Lambda p K^+$	${}^1S_0 + {}^3S_1$ admixture Worse SN	High	0.8 MeV
ALICE	pp (Femtoscopy)	${}^1S_0 + {}^3S_1$ admixture	High	No description
J-PARC E27	$d(\pi^+, K^+)$ (Inclusive)	Worse SN (inclusive)	High	1.4 MeV

$pp \rightarrow \Lambda p K^+$ : Good resolution, Worse SN,  ${}^1S_0 + {}^3S_1$  mixed,  
Complicated reaction mechanism (via  $N^*$ ,  $\Delta^*$ )  
→PWA (with Spin observable)

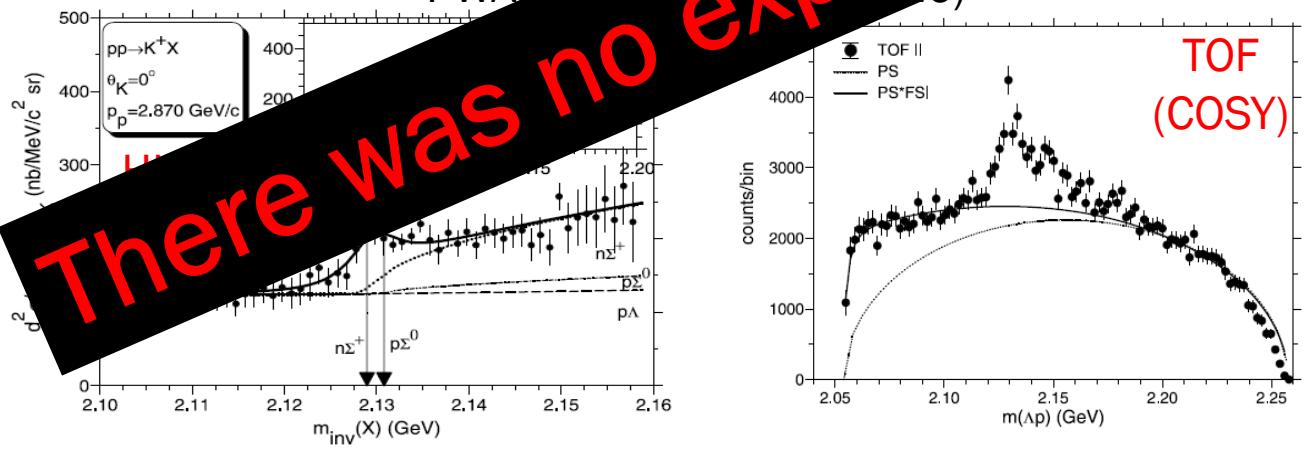
ALICE(pp@13TeV, Femtoscopy)  
 ${}^1S_0 + {}^3S_1$  mixed → Spin observable



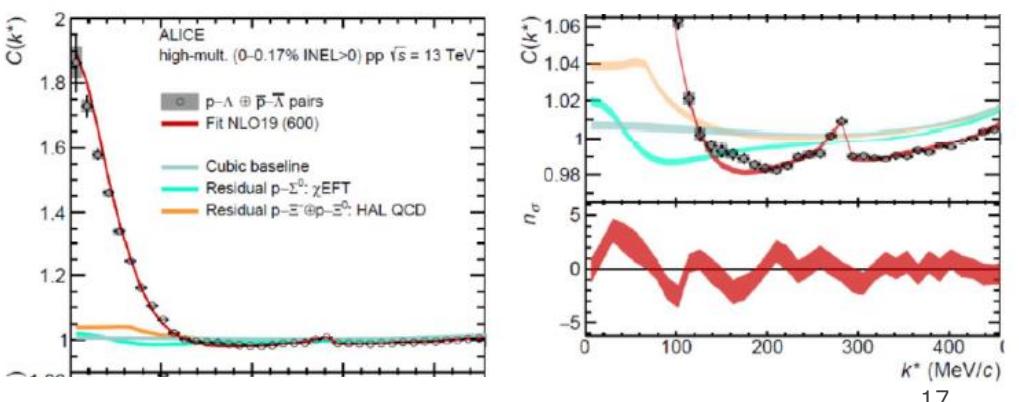
## Bubble chamber

	Reaction	Comments	Statistics	Resolution
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Eastwood	Inflight $d(K^-, \pi^-) \Lambda p$ 1450, 1650 MeV/c	Low statistic worse resolution	217 events ( $\cos \theta > 0.9$ , momcut)	3 MeV
Tan	stopped $d(K^-, \pi^-) \Lambda p$	Large FSI	2470 events	
Pigot	Inflight $d(K^-, \pi^-), d(\pi^+, K^+)$	Poor resolution	Unpublished ( $d(K^-, \pi^-)$ 1.1 MeV $d(\pi^+, K^+)$ 1.4 GeV/c)	
$pp \rightarrow \Lambda p K^+$ (COSY etc)	$pp \rightarrow \Lambda p K^+$	${}^1S_0 + {}^3S_1$ admixt. Worse resolution		0.8 MeV
ALICE	$pp$ (Femtoscopy)	${}^1S_0 + {}^3S_1$ mixed	High	No description
J-PARC E27	$d(\pi^+, K^+)$ (Inclusive)	Inclusive	High	1.4 MeV

$pp \rightarrow \Lambda p K^+$ : Good resolution, Worse resolution  
Complicated reaction, difficult to understand,  
 $\rightarrow$  PWA(will be done by ALICE)



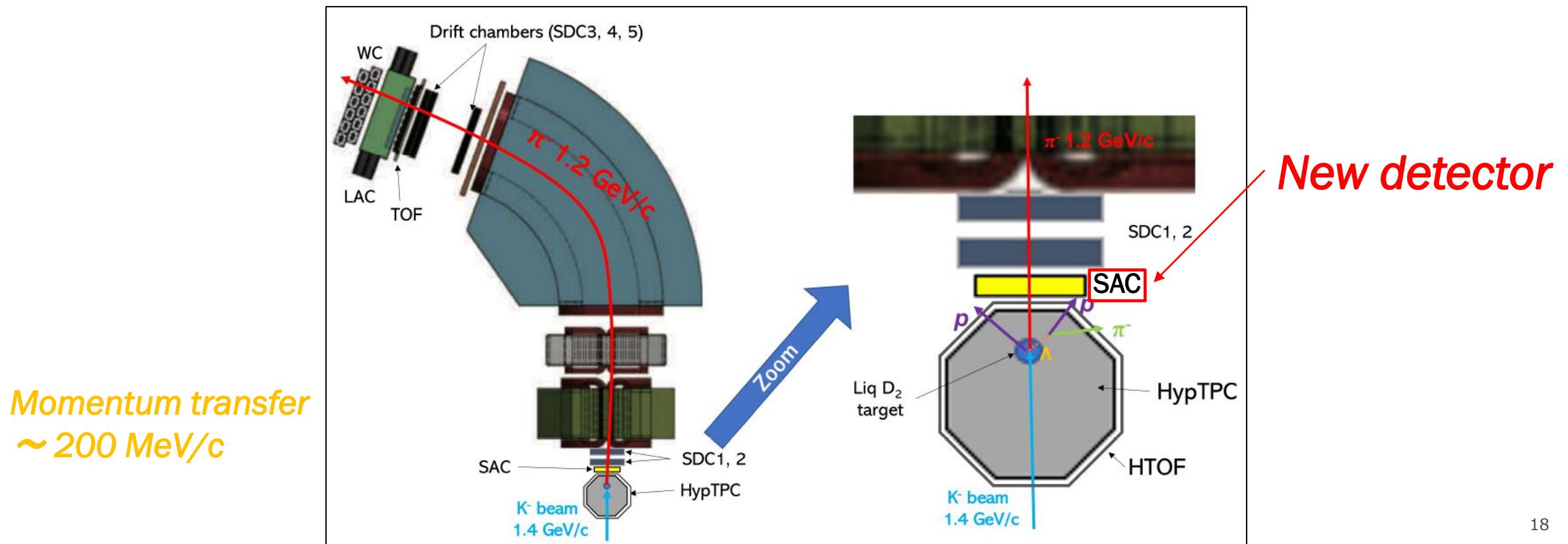
ALICE(pp@13TeV, Femtoscopy)  
 ${}^1S_0 + {}^3S_1$  mixed  $\rightarrow$  Spin observable



There was no experiment to satisfy the requirements!!

## SET UP

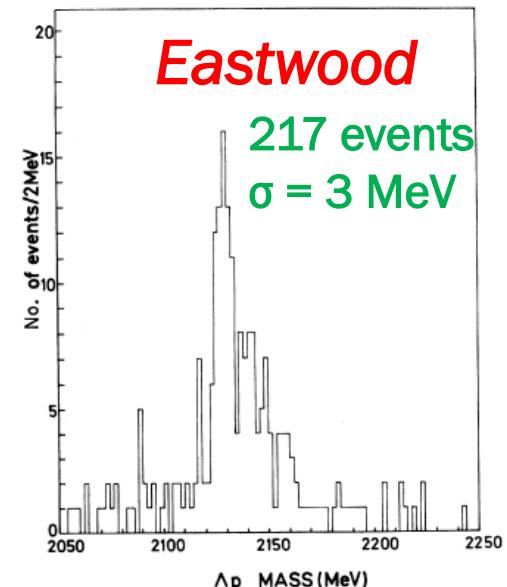
- Reaction:  $K^-d \rightarrow \Lambda p \pi^-$  at  $1.4 \text{ GeV}/c$
- S-2S(developed for E70):  $\pi^-$  measurements → measurement of missing mass spectrum
  - Good mass resolution:  $\Delta M \sim 0.4 \text{ MeV} (\sigma)$ ,  $(\Delta p/p)_{K18} = 3.3 \times 10^{-4}$ (FWHM),  $\Delta p/p_{(S-2S)} = 6.0 \times 10^{-4}$ (FWHM))
- HypTPC(developed for E42): Final state ( $\Lambda p$ ) restriction and background suppression



# YIELD ESTIMATION

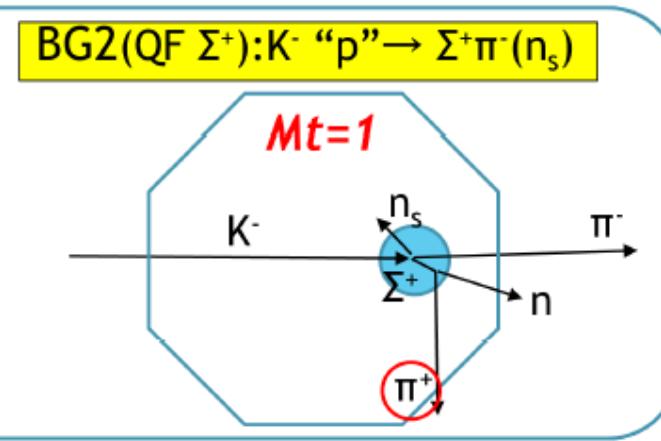
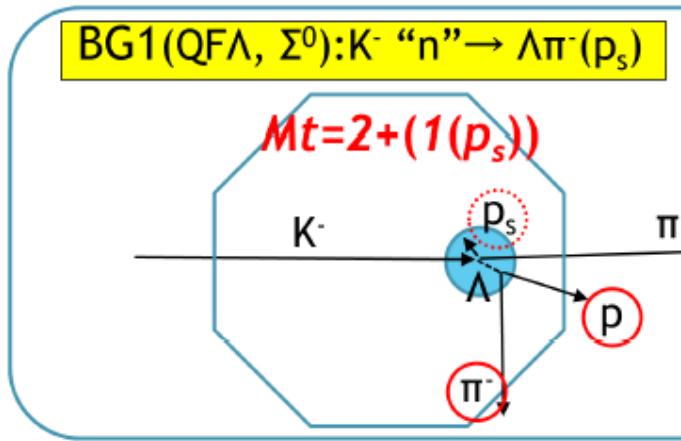
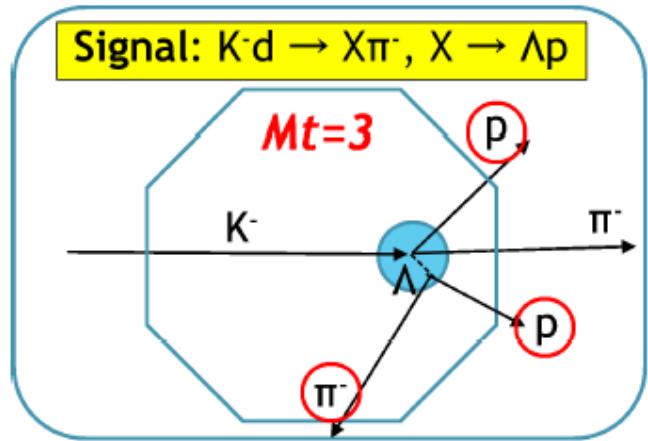
$$N = \left( \frac{d\sigma}{d\Omega} \right) \times d\Omega_{S2S} \times \left( \frac{N_{beam} \times N_A \times (\rho x)}{A} \right) \times \epsilon.$$

- $d\sigma/d\Omega = 127 \text{ nb/sr}$  in Lab (D. Eastwood et al., Phys. Rev. D 3 (1971) 2603.)
- $d\Omega = 50 \text{ msr}$
- $N_A = 6.02 \times 10^{23}$ ;
- $\rho x = 0.54 \text{ g/cm}$  by taking into account the beam size,  $\sigma_x = 23 \text{ mm}$
- $N_{beam} = 0.5M \text{ K}^-/\text{sec}$ , 90%Acc eff
- $A = 2$ .
- $\epsilon = 0.5$  (including decay factor and DAQ eff)

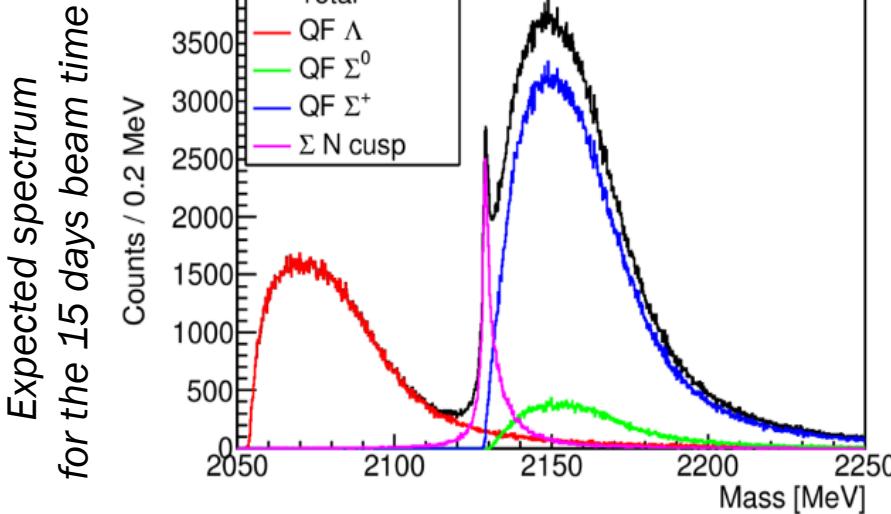


*Inclusive:  $7.6 \times 10^4$  events for 15 days beam time*  
*Exclusive:  $1.4 \times 10^4$  events for 15 days beam time*

# QF BACKGROUND SUPPRESSION BY HYPTPC

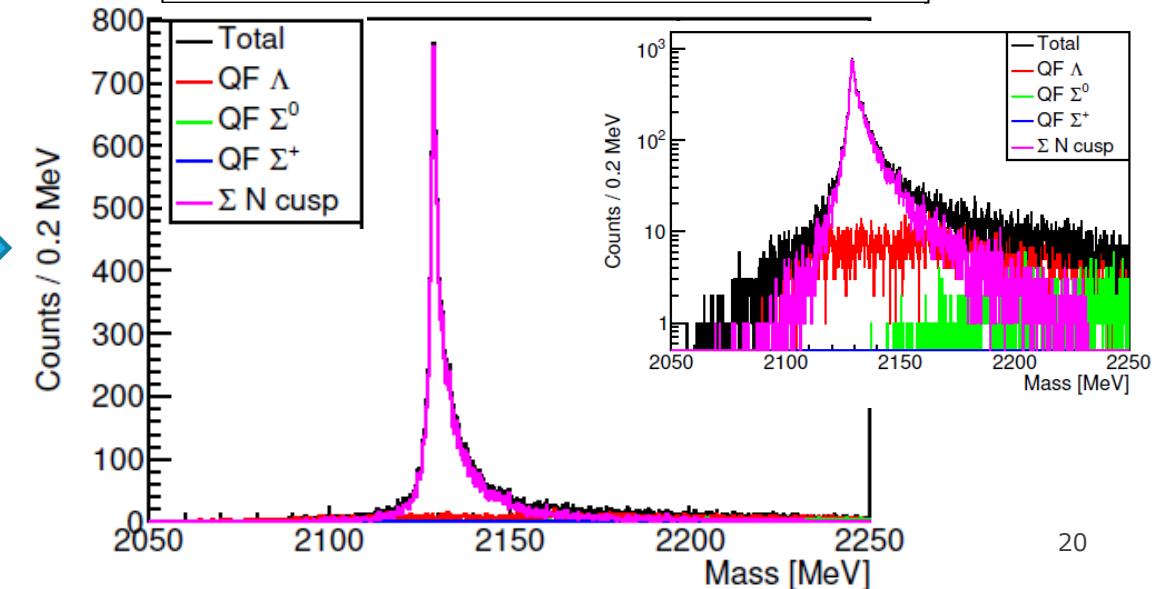


Simulated inclusive spectrum  $d(K^-, \pi^-)$

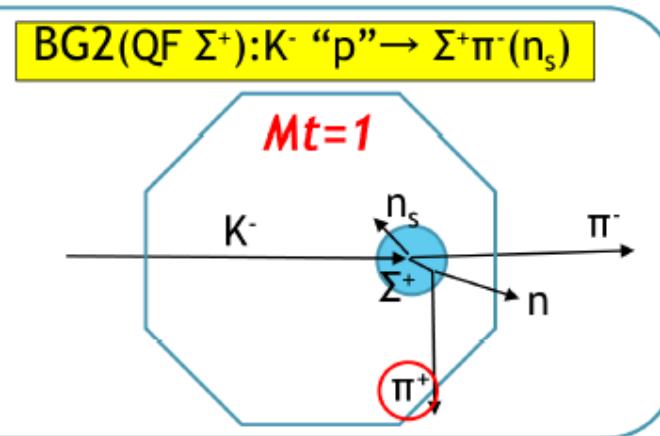
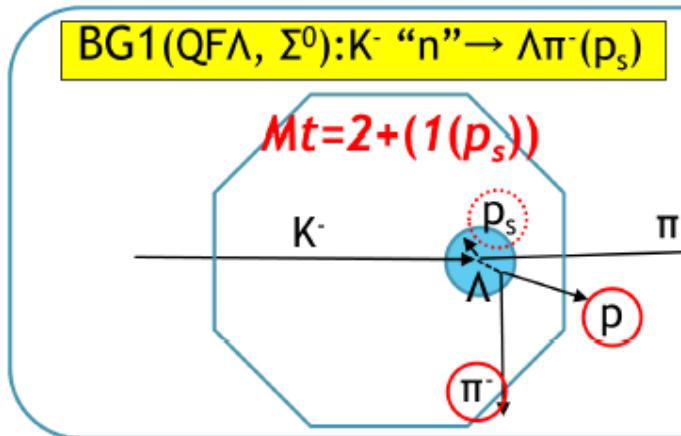
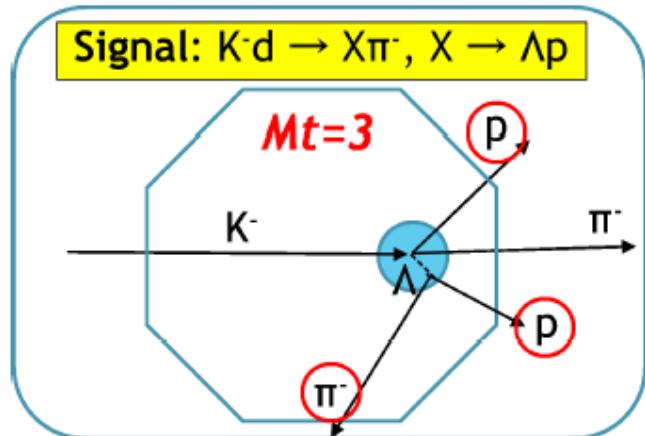


**Multiplicity = 3**

**Multiplicity = 3 without ( $K^-, \pi^-$ )**

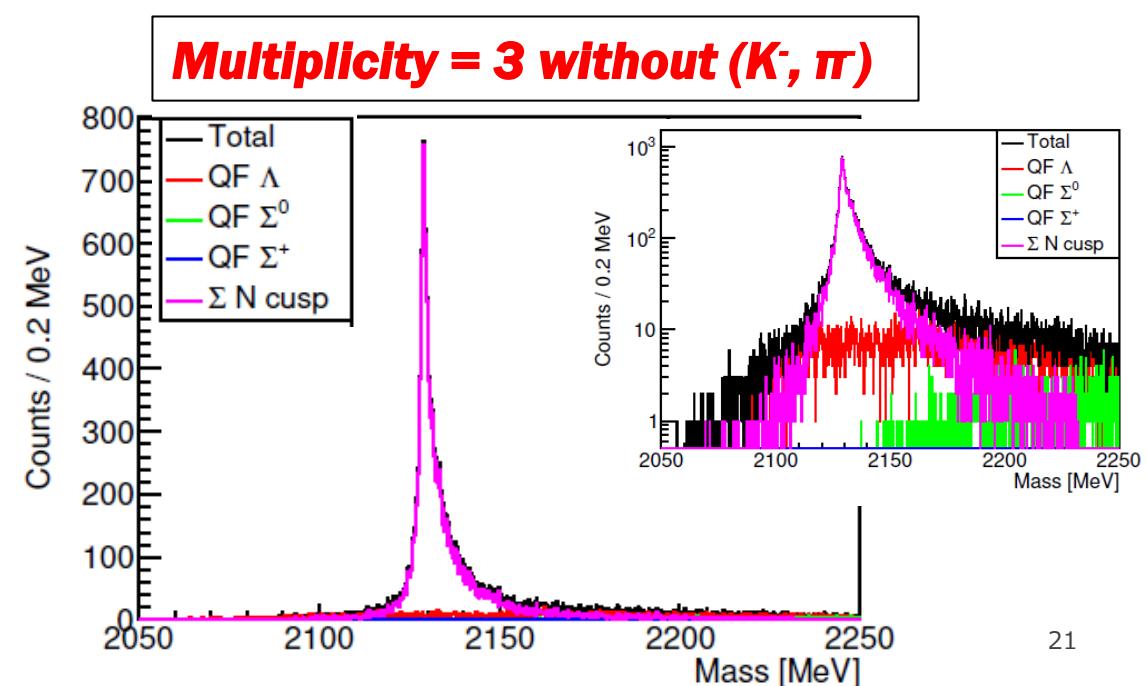
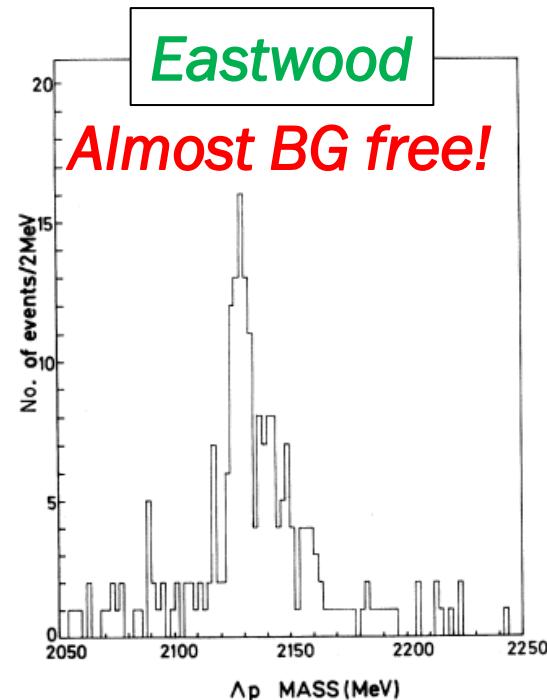


# QF BACKGROUND SUPPRESSION BY HYPTPC



$K^-d \rightarrow \Lambda p \pi^-$   
@1.45 and 1.65 GeV/c  
(Bubble chamber)

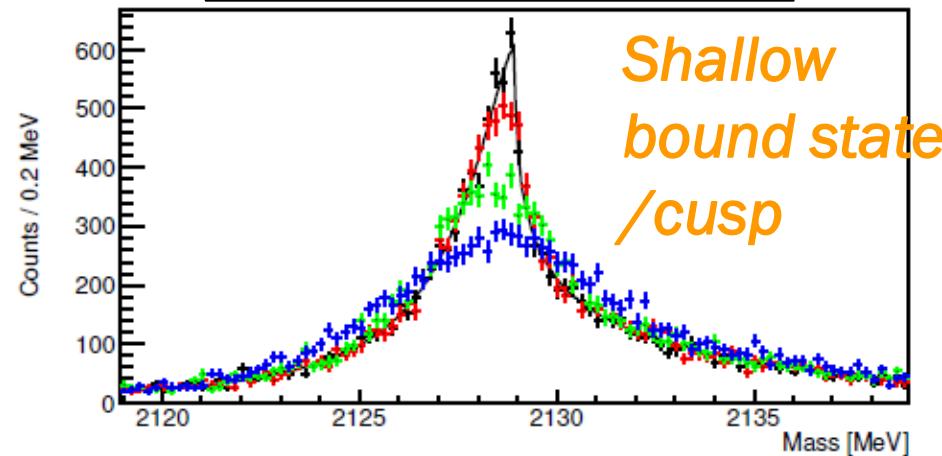
$\cos\theta_{CM} > 0.9$   
 $p_{proton} > 150$  MeV/c



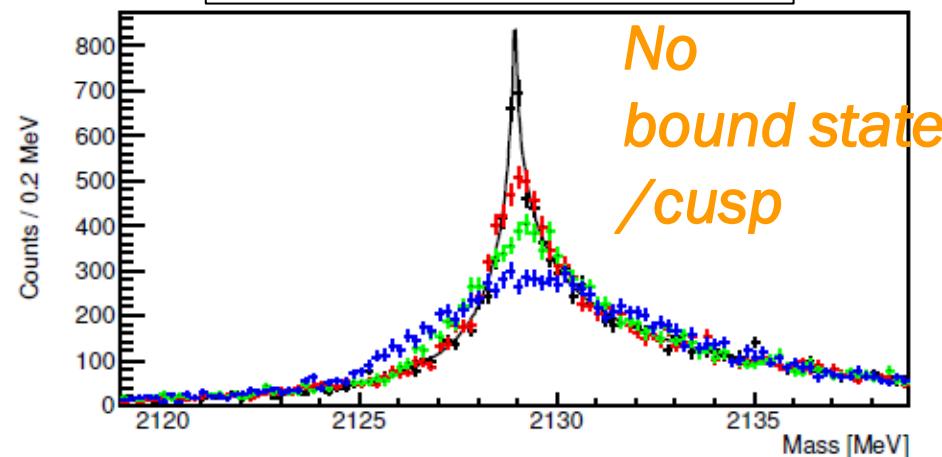
# EXPECTED SPECTRA (RESOLUTION EFFECT)

*Good energy resolution is necessary to discuss the cusp shape!!*

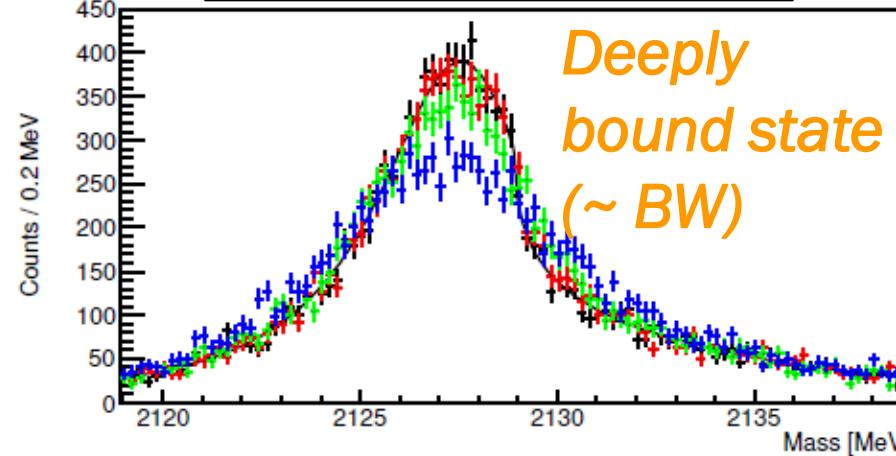
J-A ( $A_\Sigma = -2.37 + i3.74 \text{ fm}$ )



ND ( $A_\Sigma = 2.06 + i4.64 \text{ fm}$ )



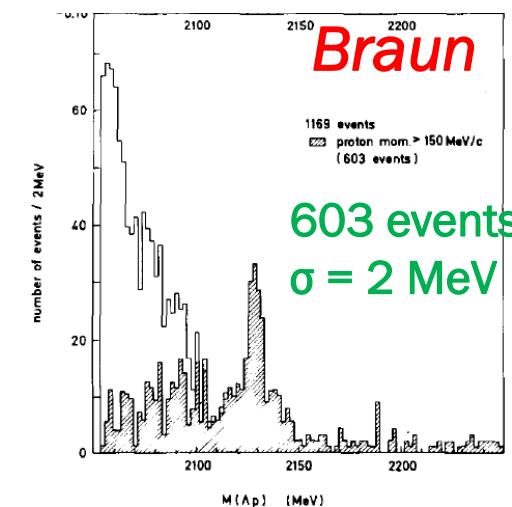
NB ( $A_\Sigma = -3.00 + i1.8 \text{ fm}$ )



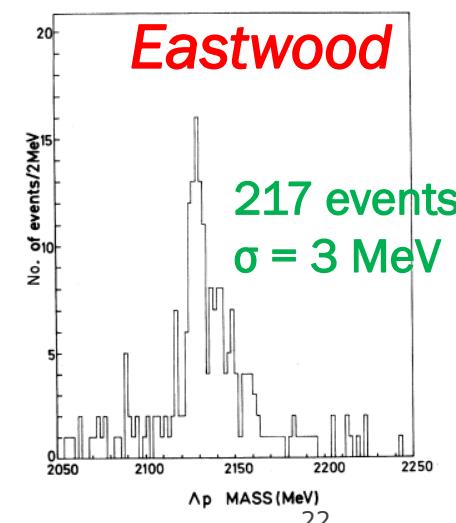
$1.4 \times 10^4$   
Events  
(P90)

- Ideal
- $\Delta M = 0 \text{ MeV}$
- $\Delta M = 0.4 \text{ MeV} (\text{P90})$
- $\Delta M = 1 \text{ MeV}$
- $\Delta M = 2 \text{ MeV}$  (Braun)

Braun



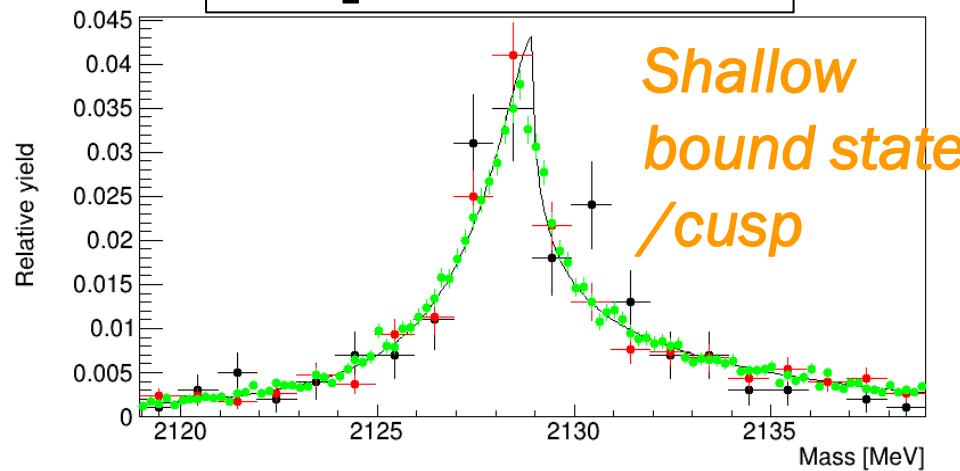
Eastwood



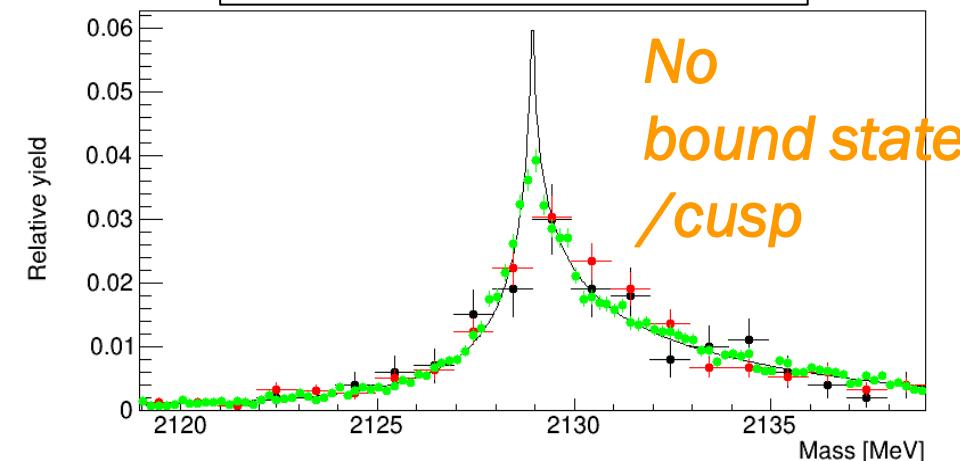
# EXPECTED SPECTRA (STATISTICAL EFFECT)

**>10<sup>4</sup> statistics is necessary!!**

**J-A ( $A_\Sigma = -2.37 + i3.74 \text{ fm}$ )**

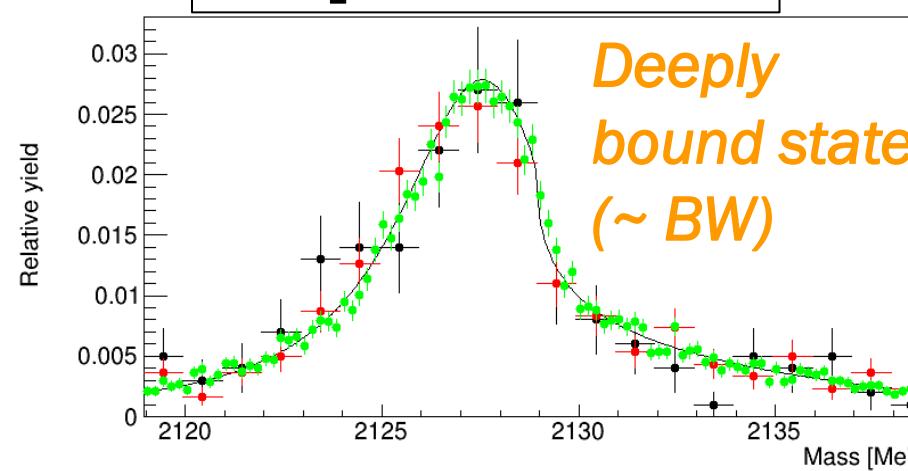


**ND ( $A_\Sigma = 2.06 + i4.64 \text{ fm}$ )**



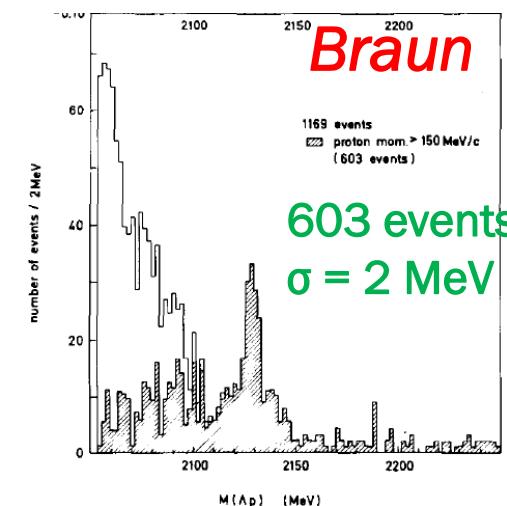
$$\Delta M = 0.4 \text{ MeV} \text{ (P90)}$$

**NB ( $A_\Sigma = -3.00 + i1.8 \text{ fm}$ )**

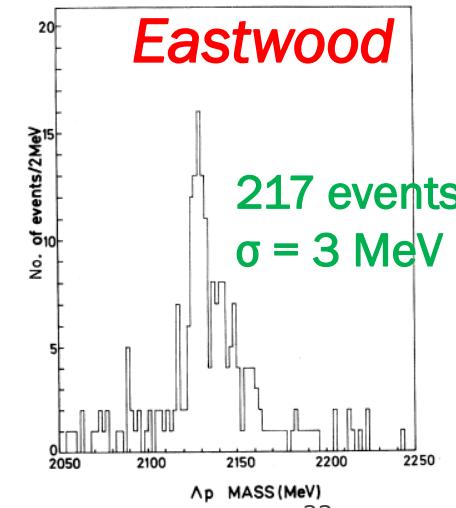


- Ideal (w/o resolution)
- 200 events (Eastwood)
- 600 events (Braun)
- 14000 events (P90)

**Braun**



**Eastwood**



217 events  
 $\sigma = 3 \text{ MeV}$

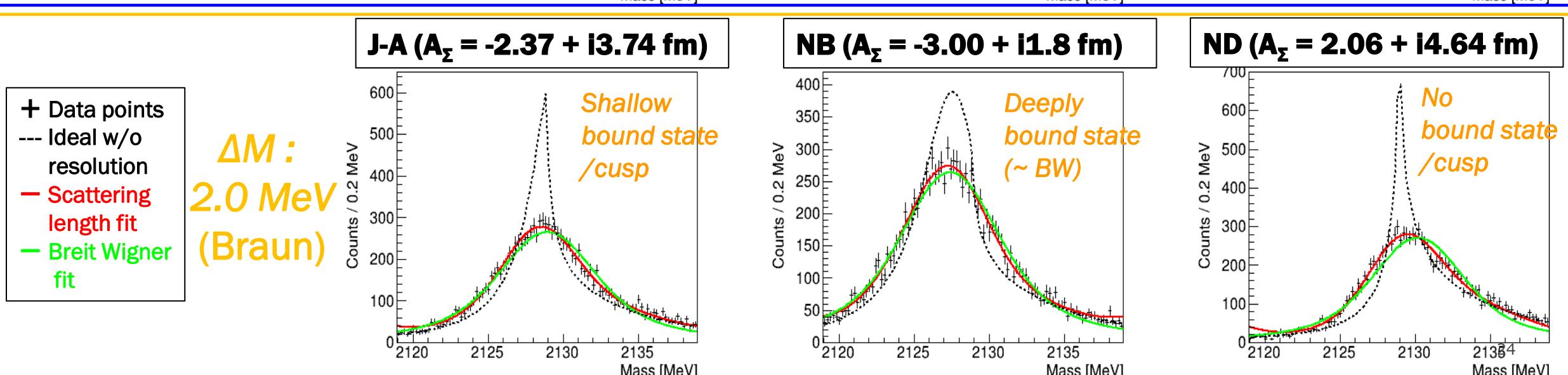
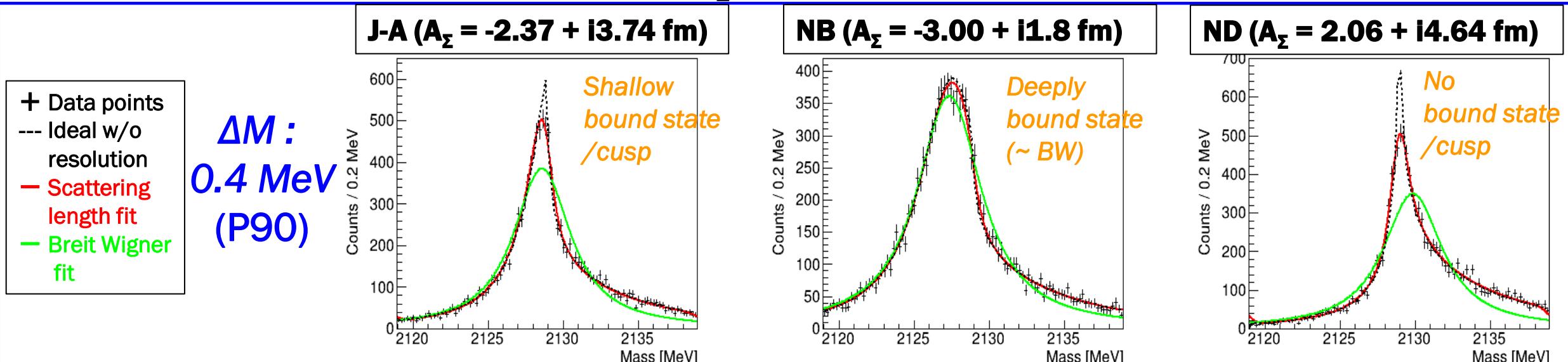
**Braun**

603 events  
 $\sigma = 2 \text{ MeV}$

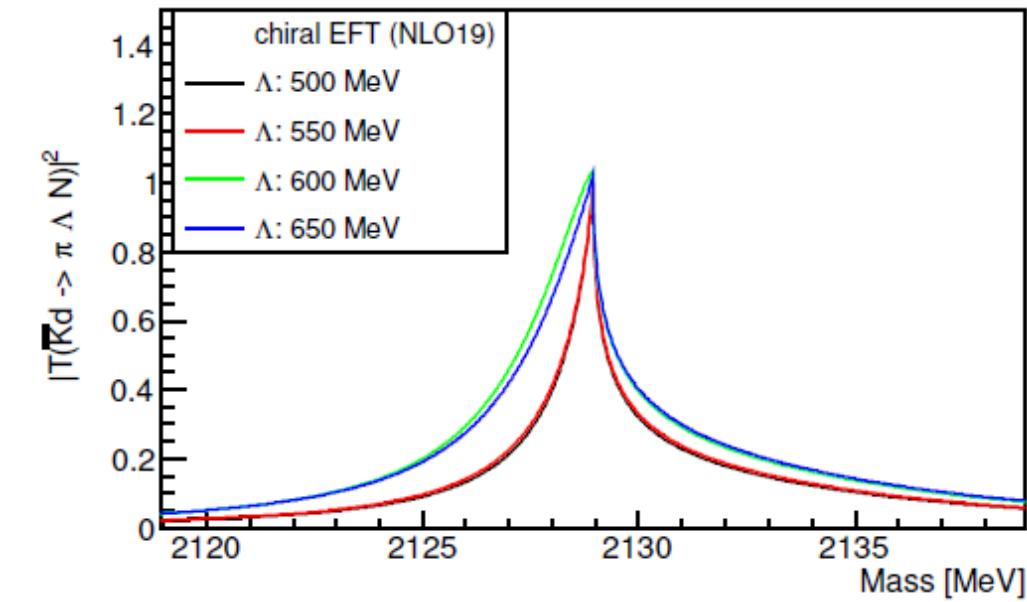
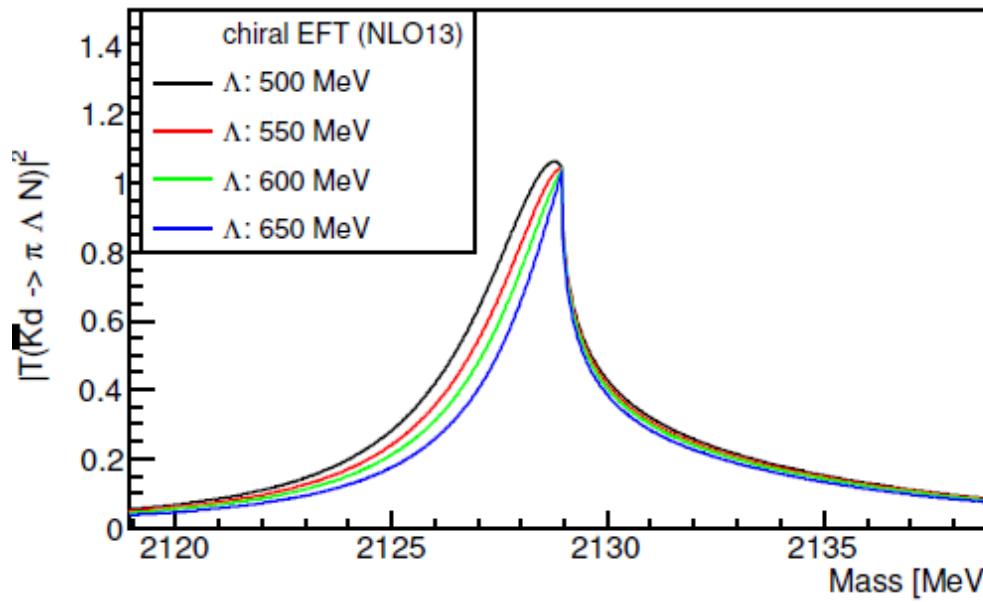
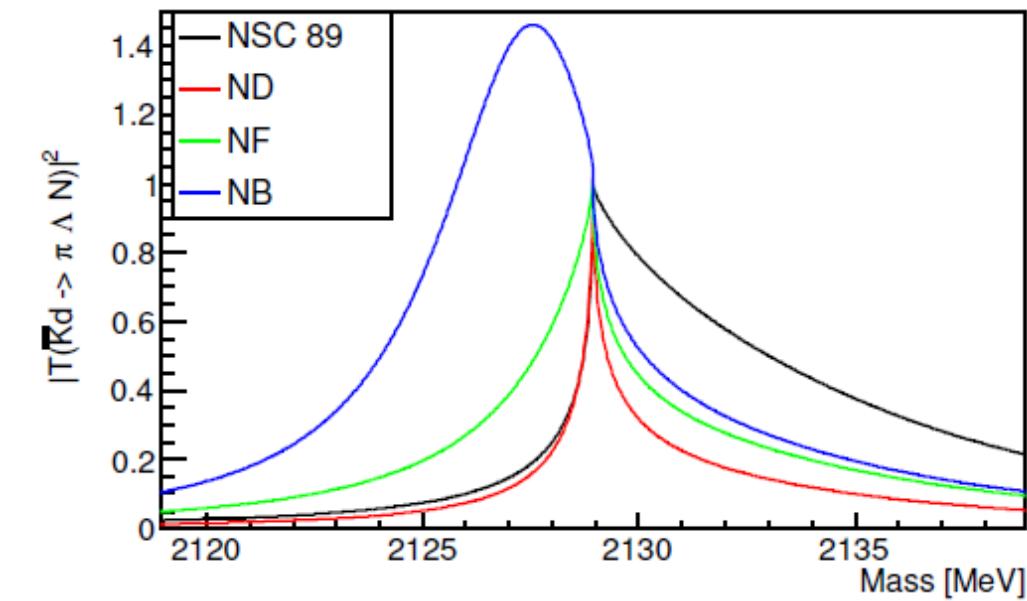
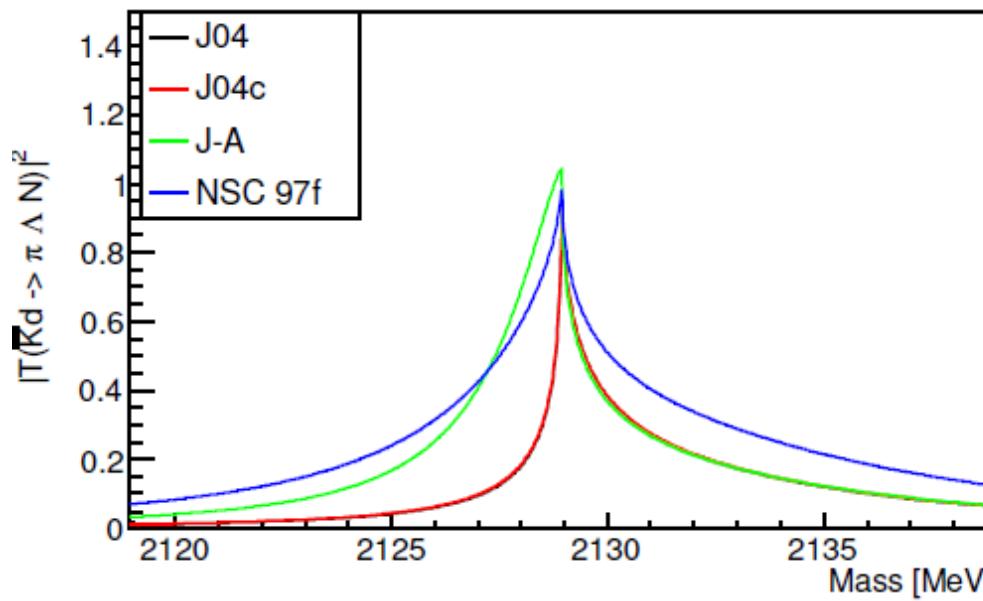
# SCATTERING LENGTH FIT VS BREIT-WIGNER FIT

Significant difference between scattering length fit and Breit-Wigner fit in  $\Delta M = 0.4 \text{ MeV}!!$

Statistical error for the scattering length ( $A_\Sigma = a + ib$ ) determination is  $< 0.3 \text{ fm}$ .

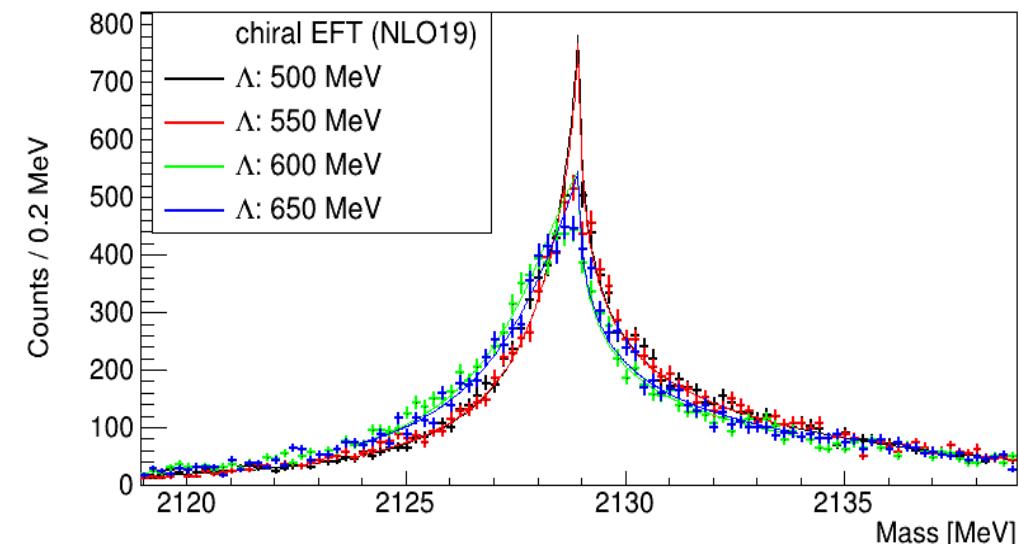
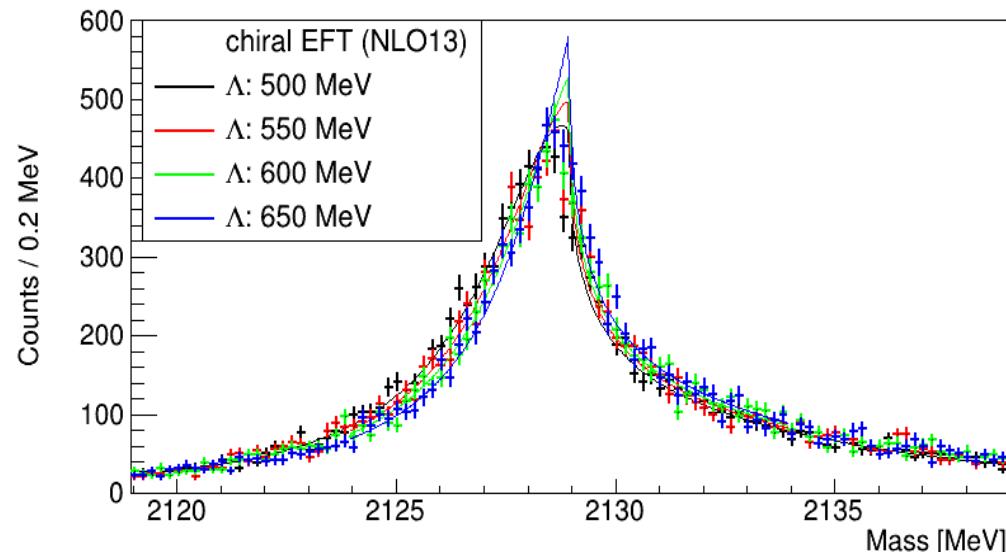
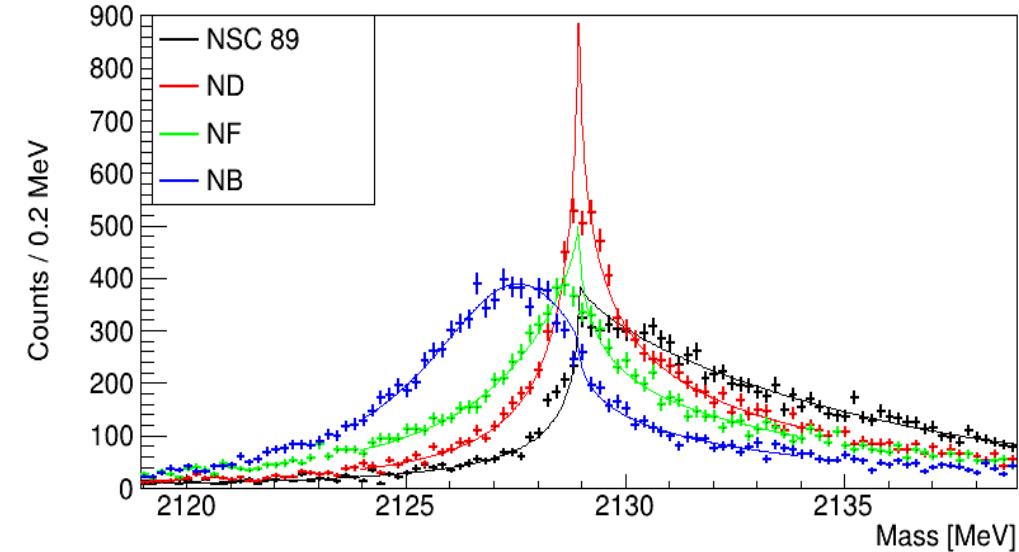
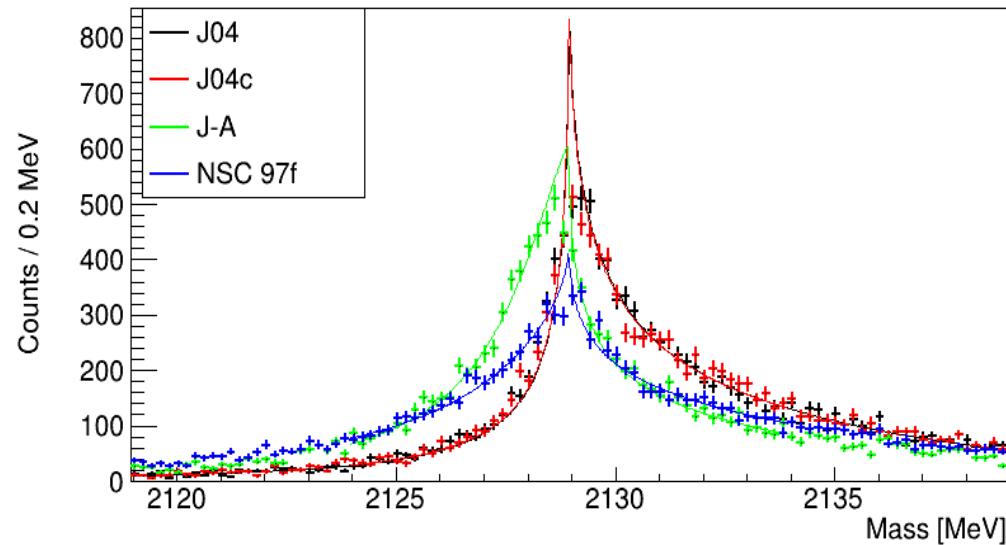


# $\Sigma$ (T=1/2, $^3S_1$ ) SCATTERING LENGTH (THEORY)



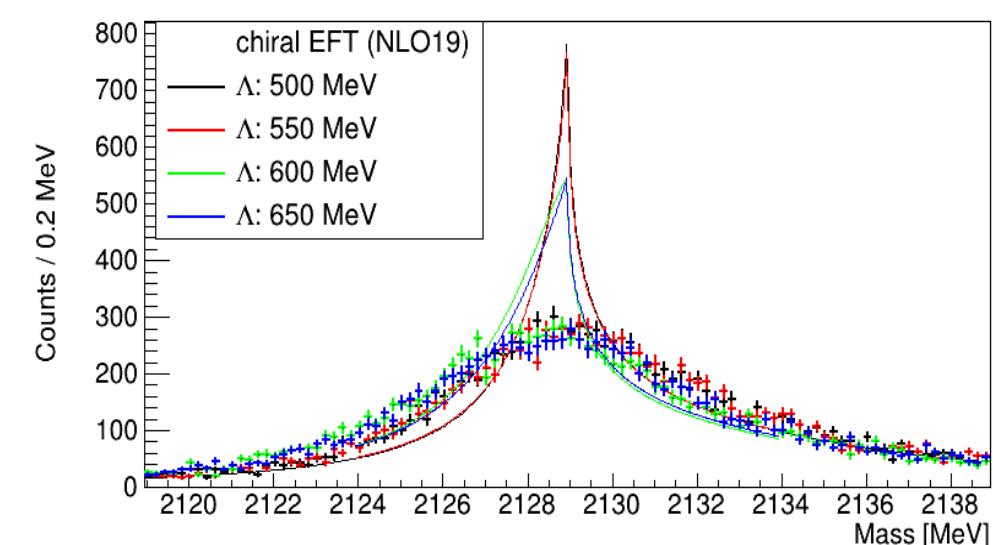
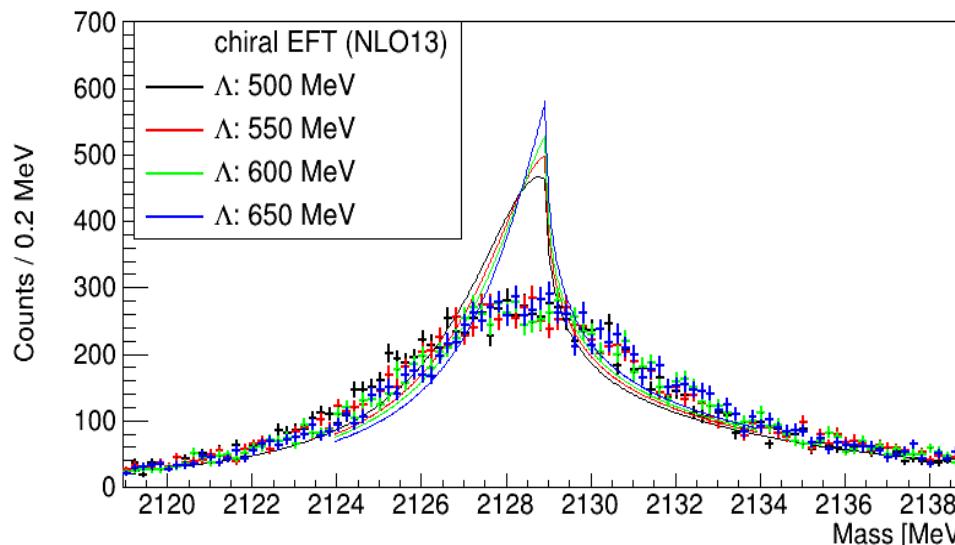
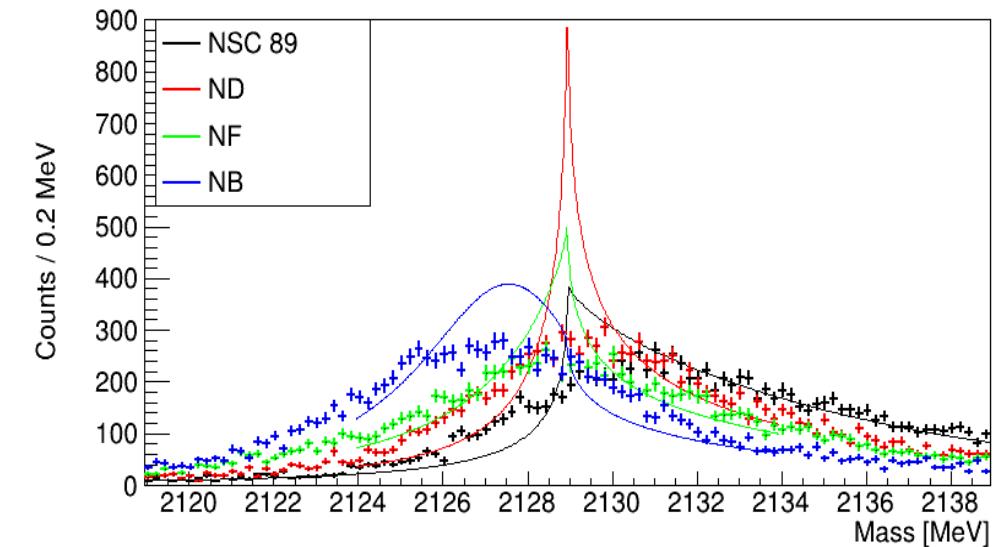
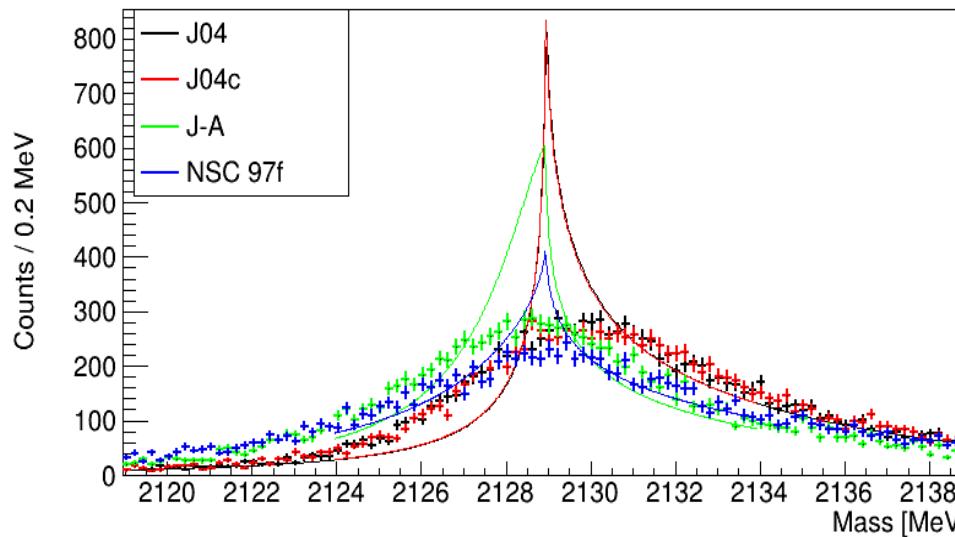
# EXPECTED SPECTRA WITH P90 QUALITY

$\Delta M = 0.4 \text{ MeV}, 1.4 \times 10^4 \text{ events}$



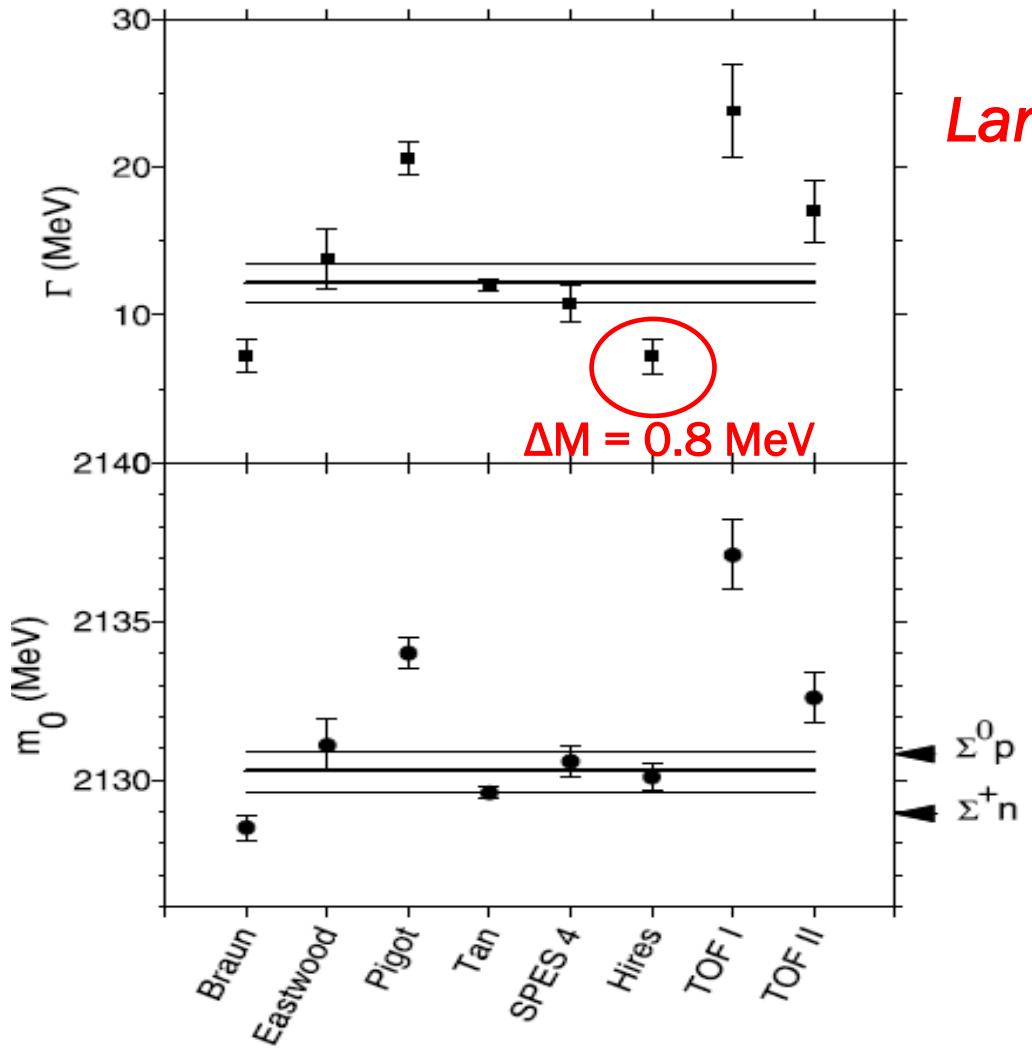
# EXPECTED SPECTRA WITH WORSE RESOLUTION

$\Delta M = 2 \text{ MeV}, 1.4 \times 10^4 \text{ events}$



# COMPARISON OF $M_0$ , $\Gamma$ (1 BW FIT)

H. Machner et al., NPA 901, 65 (2013).

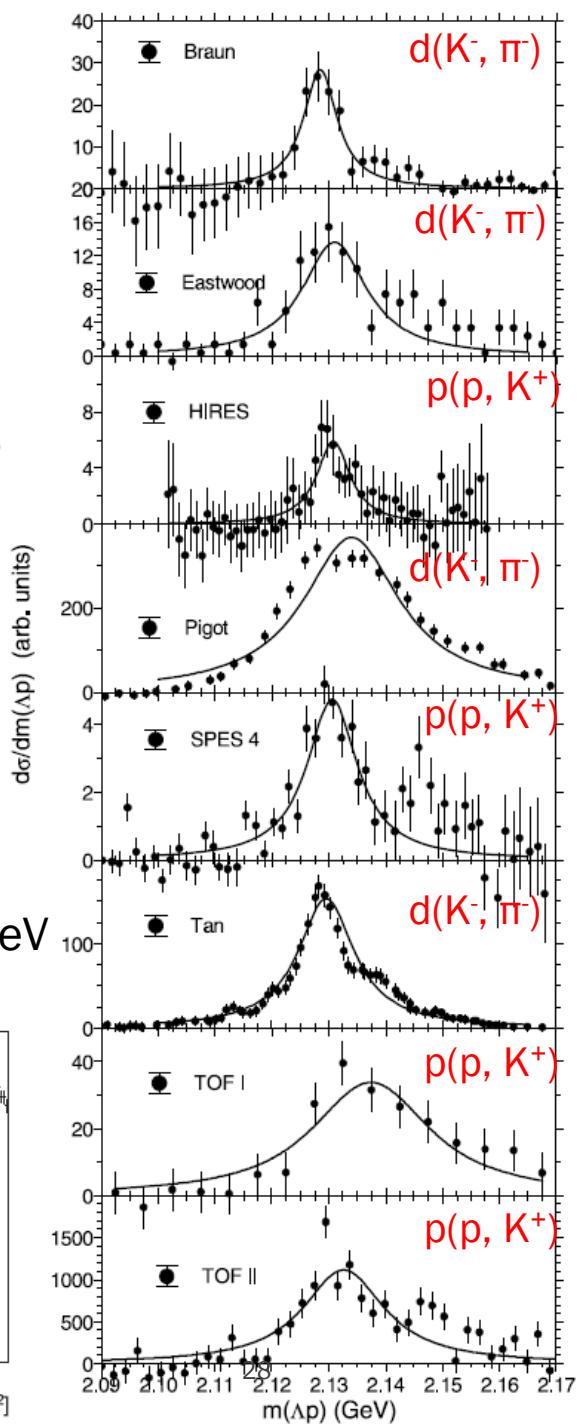
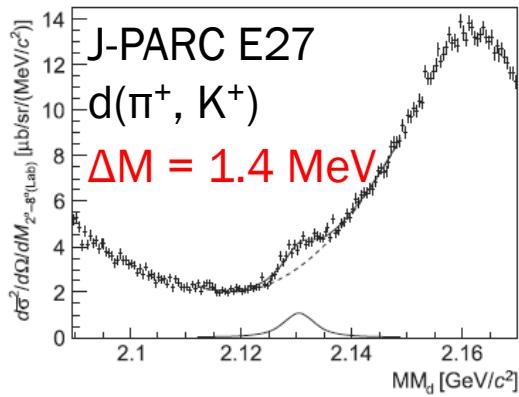


*Large width may be come from  
the worse resolution.*



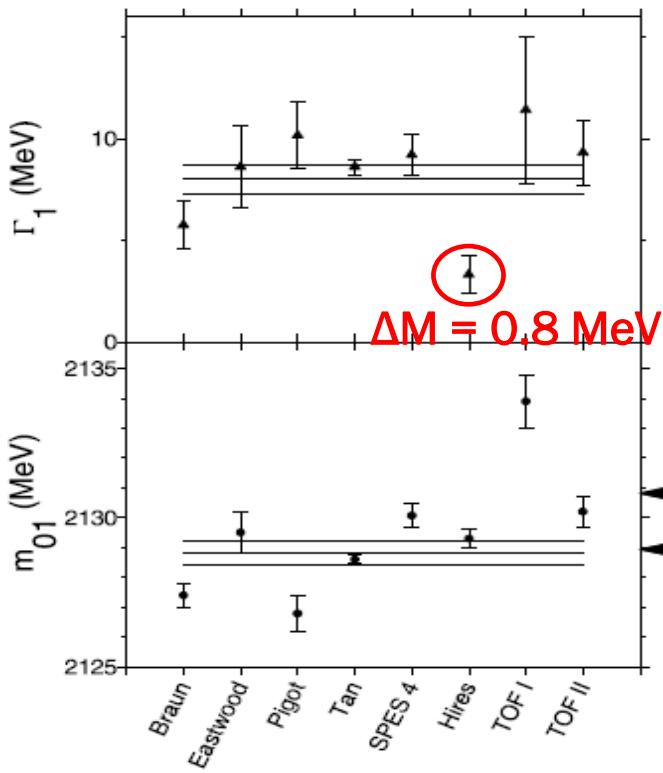
*P90 is essential!!*

E27:  $M = 2130.5 \pm 0.4(\text{stat.}) \pm 0.9(\text{syst.})$  MeV  
 $\Gamma = 5.3^{+1.4}_{-1.2}(\text{stat.})^{+0.6}_{-0.3}(\text{syst.})$  MeV



# COMPARISON OF $M_0$ , $\Gamma$ (2 BW FIT)

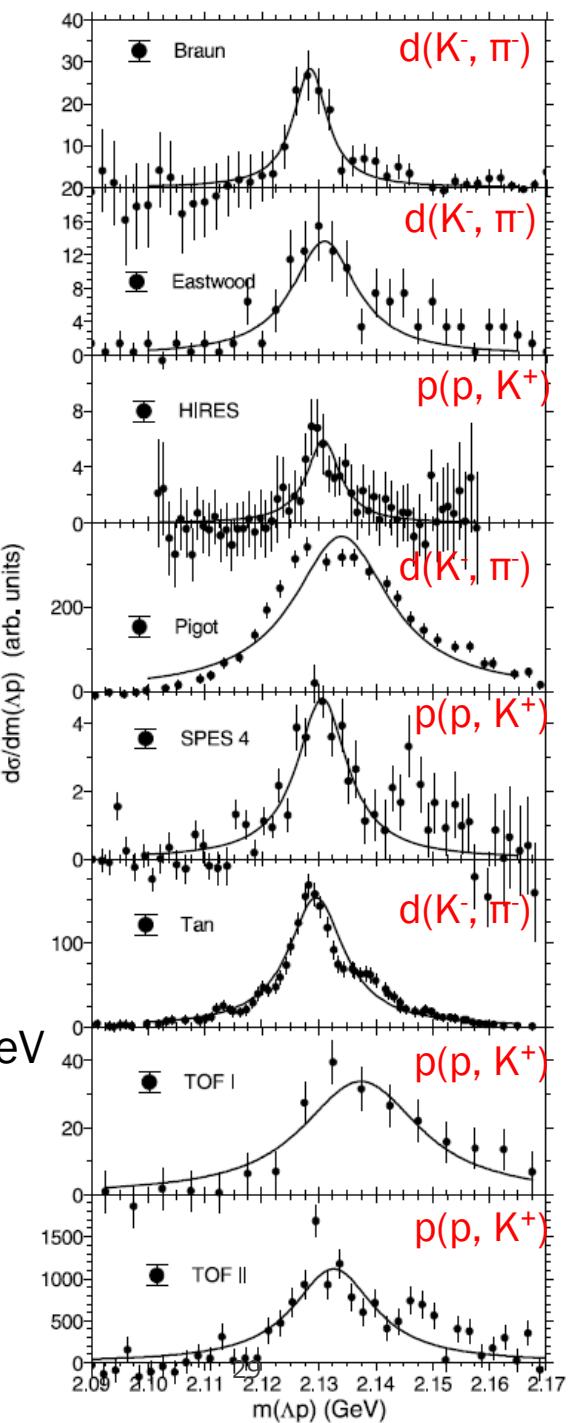
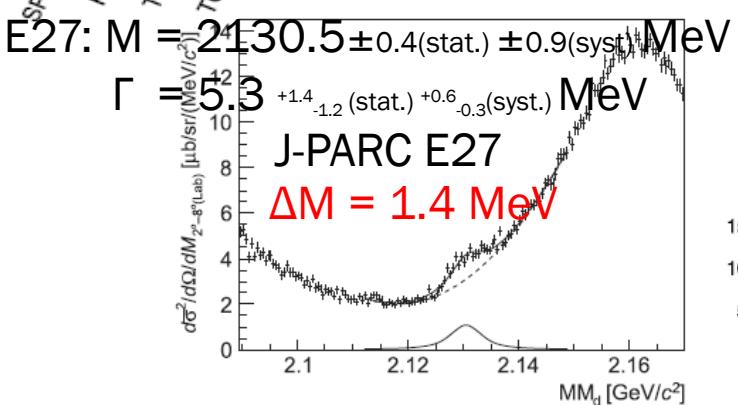
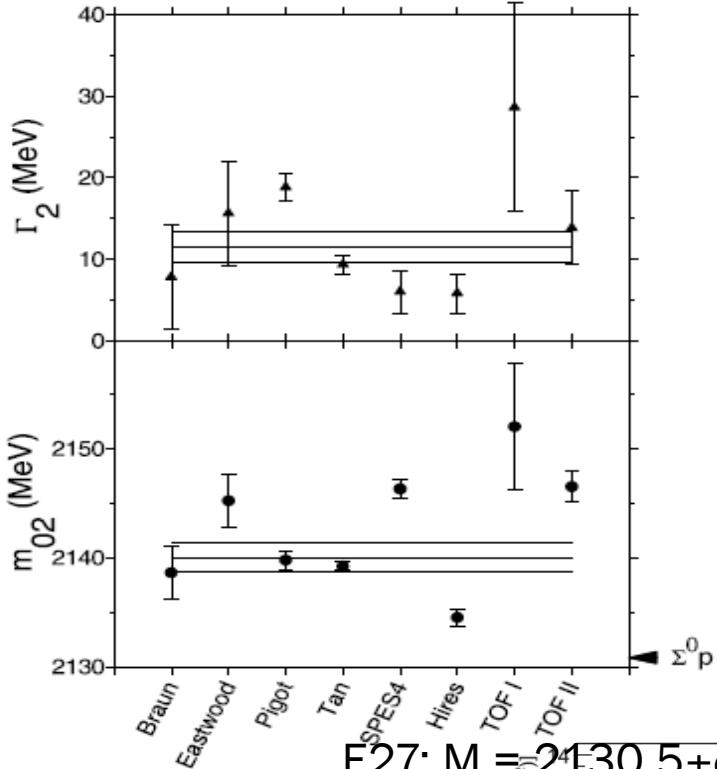
H. Machner et al., NPA 901, 65 (2013).



*Large width may be come from  
the worse resolution.*



*P90 is essential!!*



## SUMMARY

- $\Sigma N$  interaction is the important key of the  $B_8B_8$  interaction and  $(\Lambda, \Sigma)$  hypernuclei.
  - “ $\Sigma N$  cusp” can be expressed by the  $\Sigma N$  interaction (scattering length).
    - There are a lot of past experiments to measure the “ $\Sigma N$  cusp”. However, the origin of the “ $\Sigma N$  cusp” remains unclear yet.  $\Sigma N$  dibaryon or not?
    - Inflight  $d(K^-, \pi^-)$  reaction has advantage to dedicate ( $T=1/2, {}^3S_1$ ) channel.
  - P90 will investigate the nature of “ $\Sigma N$  cusp” with the world’s best quality.
    - K1.8 Beam line, S-2S for  $\pi^-$  measurement, and HypTPC for BG suppression.
    - $1.4 \times 10^4$   $\Sigma N$  cusp events are expected in **15 days** beam time.
    - **0.4 MeV** ( $\sigma$ ) mass resolution will be achieved, 2 times better than past exp.
    - We can deduce scattering length with the statistical error less than 0.3 fm.
- We will be ready to install by the end of JFY 2023. [SAC, TPC stands]