

Observation of ω contribution to $\chi_{c1}(3872) \rightarrow \pi^+ \pi^- J/\psi$ decays at LHCb

Excellence Cluster ORIGINS, Munich, Germany

on behalf of the LHCb Collaboration

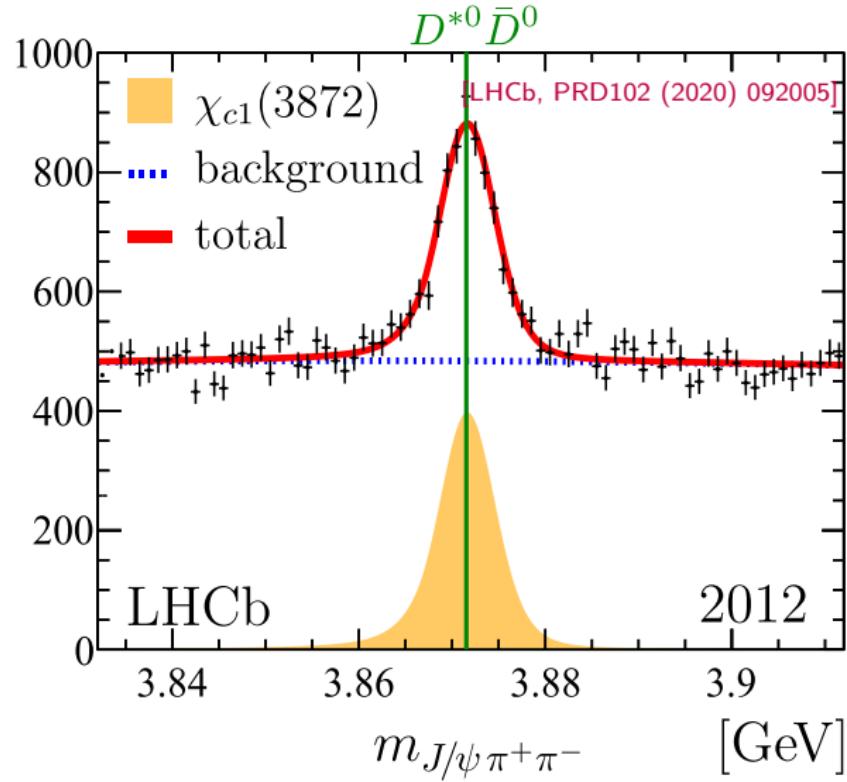
August 30th, 2022

QWG, Darmstadt, Germany

The first bird of the exotic era: $\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-$

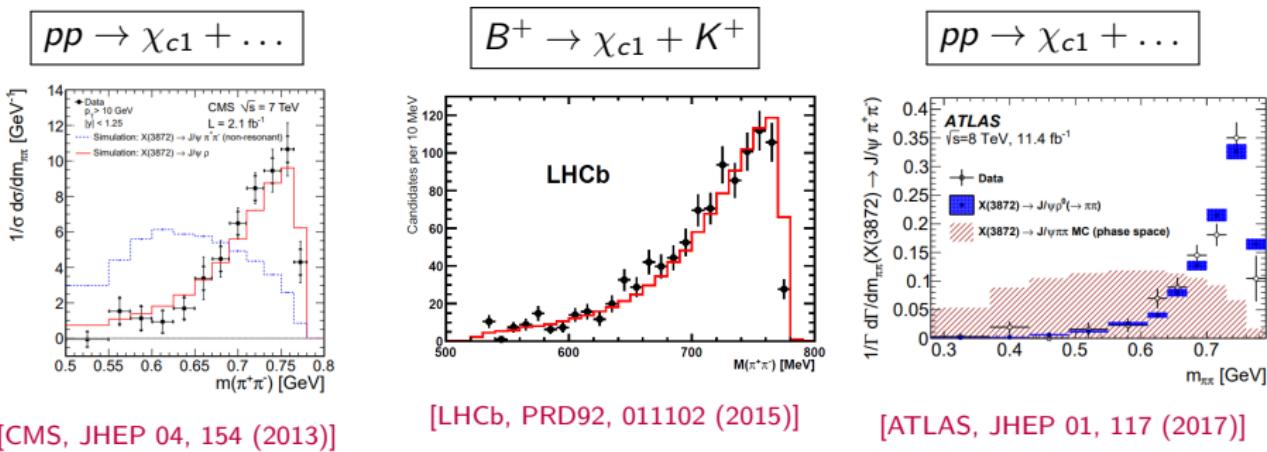
First seen by Belle [PRL91 (2003), 262001]

- Charmonium state ($c\bar{c}$), $\chi_{c1}(2P)$ is expected* 70 MeV above
- The peak right at the $D^{*0}\bar{D}^0$ threshold
- Large isospin violation in the strong decays to $J/\psi \pi^+ \pi^-$ [LHCb 2022, 2204.12597]



Isospin violation in decays of $\chi_{c1}(3872)$

- $\chi_{c1}(3872)$ is mostly isosinglet, no $\chi_{c1} \rightarrow J/\psi \pi^\pm \pi^0$
 [BaBar PRD71, 031501 (2005); Belle PRD84, 052004 (2011)]
 - ▶ $\chi_{c1} \rightarrow J/\psi \omega$ conserves isospin
 - ▶ $\chi_{c1} \rightarrow J/\psi \rho$ violates isospin
- However, ($\pi\pi$) in the decay is in isovector model:
 - ▶ No $\chi_{c1}(3872) \rightarrow \pi^0 \pi^0 J/\psi$ observed
 - ▶ $\pi^+ \pi^-$ distribution is consistent with ρ .



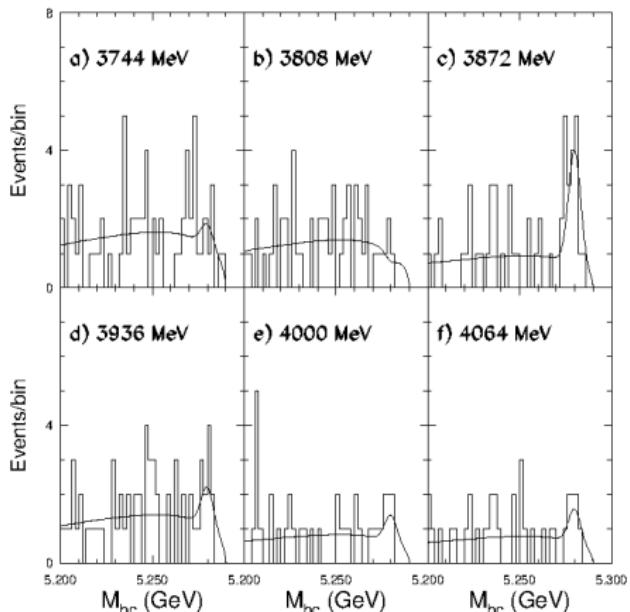
[CMS, JHEP 04, 154 (2013)]

[LHCb, PRD92, 011102 (2015)]

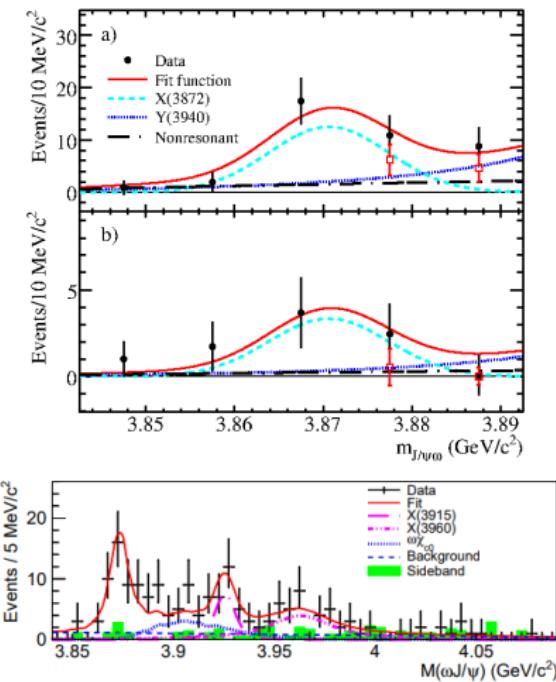
[ATLAS, JHEP 01, 117 (2017)]

Observation of $\chi_{c1} \rightarrow \omega J/\psi$ using $B^{0,+} \rightarrow \chi_{c1} K^{0,+}$

[Belle, 0505037 (2005)]



[BaBar, 1005.5190 (1010)]



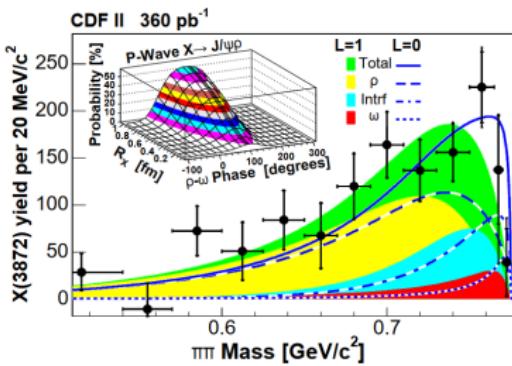
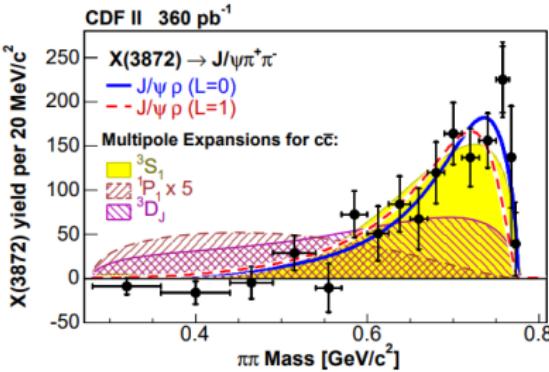
Averaged value of the ratio:

$$\frac{\mathcal{B}(\chi_{c1}(3872) \rightarrow \omega J/\psi)}{\mathcal{B}(\chi_{c1}(3872) \rightarrow \pi^+ \pi^- J/\psi)} = 1.4 \pm 0.3$$

Attempts to find ω in $\pi\pi$ by CDS II

[CDF II, PRL 96, 102002 (2006)]

$$pp \rightarrow \chi_{c1} + \dots$$



- Prompt production

- $N_{ev} = 1260 \pm 130$

- Check of quantum numbers:

- Consistent with S -wave, $J^{PC} = 1^{++}$

Contribution of ω :

- $\mathcal{M} = BW_\rho + A_\omega e^{i\phi_\omega} BW_\omega$

- $\phi_\omega = 95^\circ = \arctan \Gamma_\rho / (2\Delta m)$
[Goldhaber (1969)]

- A_ω is fixed to match $J/\psi 3\pi$ rate
[Belle, 0505037 (2005)]

- gives worse fit, p-value: 55% \rightarrow 19%

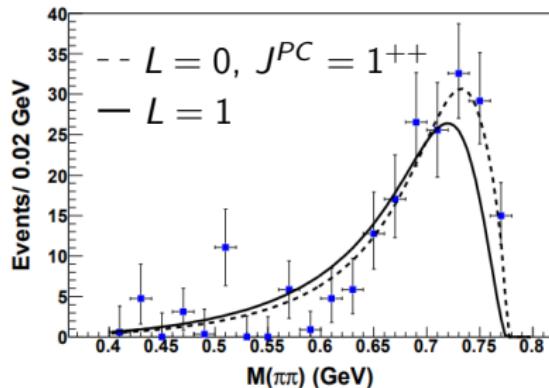
- No evidence of ω , however,
- contribution of ω on the level of < 23% is not excluded

Note: sum of BW does not respect unitarity

Attempts to find ω in $\pi\pi$ by Belle

[Belle, PRD 84, 052004 (2011)]

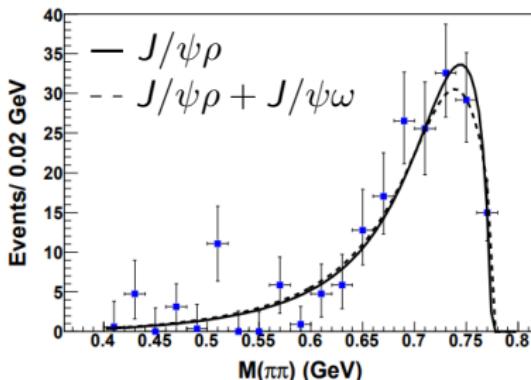
$$B^+ \rightarrow \chi_{c1} + K^+$$



- clean B decays gives similar sensitivity as in the CDF II analysis
- $N_{ev} = 159 \pm 15$
- Check of quantum numbers:
 - ▶ S -wave $J^{PC} = 1^{++}$ fits better

Contribution of ω :

- $\mathcal{M} = BW_\rho + A_\omega e^{i\phi_\omega} BW_\omega$
- $\phi_\omega = 95^\circ$ is fixed
- A_ω is floated.
- Significance of ω is 1.3σ .
- with the contribution of $(12 \pm 10)\%$



Note: sum of BW does not respect unitarity

New LHCb analysis searching for ω

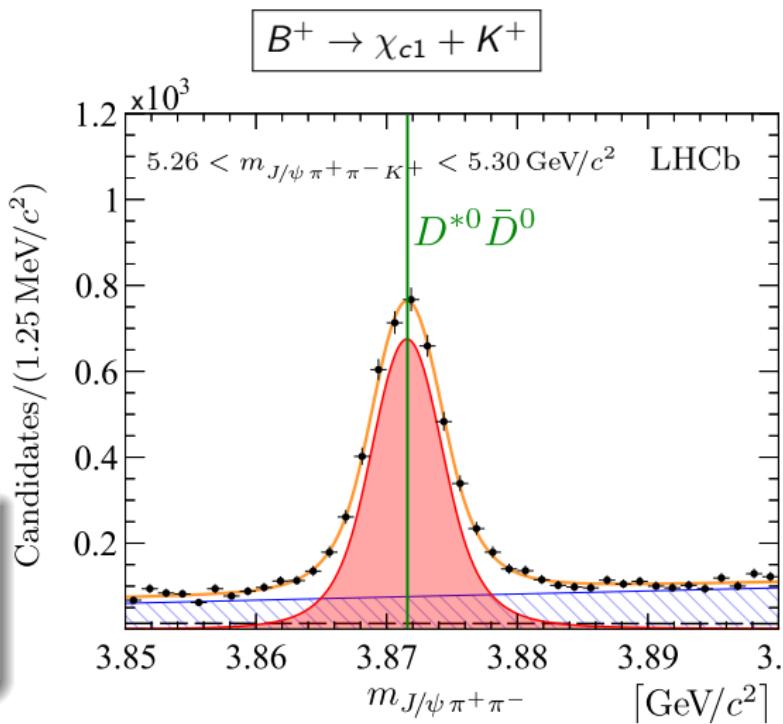
[LHCb, JHEP 08 (2020) 123]

Details on the selection:

- displaced B -vertex
- $N_{ev} = 6788 \pm 117$ of $B^+ \rightarrow \chi_{c1} K^+$
- $\times 43$ events of Belle
- $\times 6$ better statistical errors than in Belle and CDF II

Extraction of χ_{c1} yield in bins of $\pi\pi$ spectrum:

- 2d unbinned fit of $m_{J/\psi\pi\pi} \times m_{\pi\pi}$



$J/\psi\pi^+\pi^-$ spectrum from [LHCb, JHEP 08 (2020) 123]

Attempt to fit the spectrum by $\rho \rightarrow \pi\pi$ amplitude

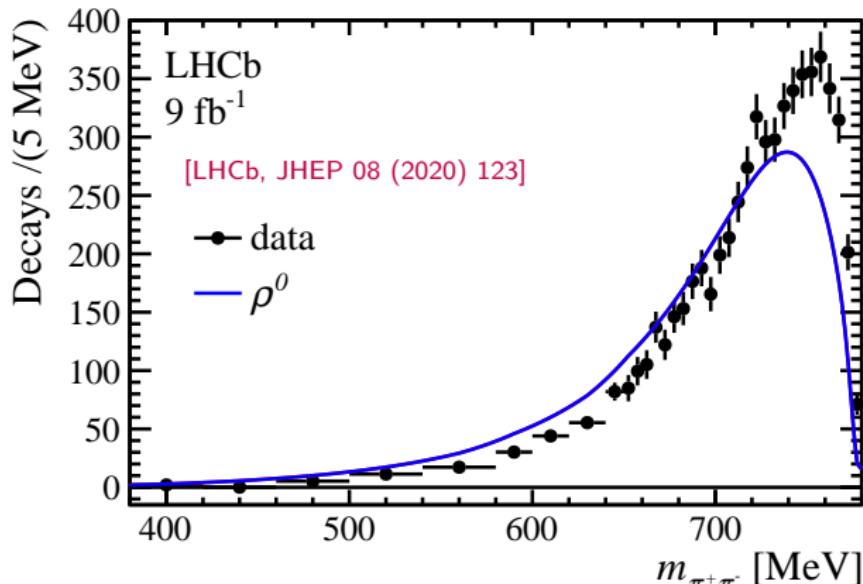
ρ lineshape is well known:

- combined analysis of $\pi\pi$ P -wave, GKPY [Garcia Martin et al, PRD83, 074004 (2011)]
- $e^+e^- \rightarrow \pi^+\pi^-$ well described by Gounaris-Sakurai [BaBar, PRD86, 032013 (2012)]
- P -wave Breit-Wigner with Blatt-Weisskopf barrier factor, $R = 1.45 \text{ GeV}^{-1}$.

$$\frac{dN}{dm_{\rho\pi}} = p_\pi p_{J/\psi} |\mathcal{A}|^2 .$$

$\chi_{c1}/\text{ndf} = 366.6/34$:

- Can ρ shape be different in the decay?
- $\omega \rightarrow \pi^+\pi^-$
- Tails of higher ρ states?



Coupled-channel model of $\pi\pi/3\pi$ scattering

[LHCb, JHEP 08 (2020) 123]

- Two-channel K -matrix approach, $T = [1 - iK\rho]^{-1}K$ with $\pi\pi \leftrightarrow \rho\pi$

$$K = \frac{1}{m_\rho^2 - s} \begin{pmatrix} g_{\rho \rightarrow 2\pi}^2 & 0 \\ 0 & 0 \end{pmatrix} + \frac{1}{m_\omega^2 - s} \begin{pmatrix} g_{\omega \rightarrow 2\pi}^2 & g_{\omega \rightarrow 2\pi} g_{\omega \rightarrow 3\pi} \\ g_{\omega \rightarrow 2\pi} g_{\omega \rightarrow 3\pi} & g_{\omega \rightarrow 3\pi}^2 \end{pmatrix},$$

- ρ is diag. loop matrix,
 - ▶ $\rho_{2\pi}$: $\pi\pi$ P-wave,
 - ▶ $\rho_{3\pi}$: $\rho\pi$ P-wave with $\rho \rightarrow \pi\pi$.
- The parameters of K are fixed from PDG:

$$\begin{aligned} g_{\rho \rightarrow 2\pi}^2 &= m_\rho \Gamma_\rho / \rho_{2\pi}(m_\rho^2), \\ g_{\omega \rightarrow 3\pi}^2 &= m_\omega \Gamma_\omega \mathcal{B}(\omega \rightarrow 3\pi) / \rho_{3\pi}(m_\omega^2), & \mathcal{B}(\omega \rightarrow 3\pi) &= (89.3 \pm 0.6)\% \\ g_{\omega \rightarrow 2\pi}^2 &= m_\omega \Gamma_\omega \mathcal{B}(\omega \rightarrow \pi^+ \pi^-) / \rho_{2\pi}(m_\omega^2), & \mathcal{B}(\omega \rightarrow 2\pi) &= (1.53 \pm 0.06)\% \end{aligned}$$

- Use Q -vector approach with the production parameters α – subject to fit,

$$A_{2\pi} = (\alpha_{2\pi} T_{2\pi,2\pi} + \alpha_{3\pi} T_{3\pi,2\pi}) \sqrt{B_1} \quad P\text{-wave amplitude}$$

Parametrization of the production

[LHCb, JHEP 08 (2020) 123]

Intuitive understanding in LO

Can simplify by neglecting the second-order terms, $g_{\omega \rightarrow 2\pi}^2/g_{\rho \rightarrow 2\pi}^2 \sim 0.0009$

$$\hat{T}_{2\pi,2\pi} \approx g_{\rho \rightarrow 2\pi}^2 \text{BW}_\rho, \quad \text{BW}_\rho = 1/(m_\rho^2 - s - ig_{\rho \rightarrow 2\pi}^2 \rho_{2\pi})$$

$$\hat{T}_{2\pi,3\pi} \approx g_{\omega \rightarrow 2\pi} g_{\omega \rightarrow 3\pi} (m_\rho^2 - s) \text{BW}_\rho \text{BW}_\omega.$$

Parametrization of the production, $\hat{A}_{2\pi} = \alpha_{2\pi} T_{2\pi,2\pi} + \alpha_{3\pi} T_{3\pi,2\pi}$:

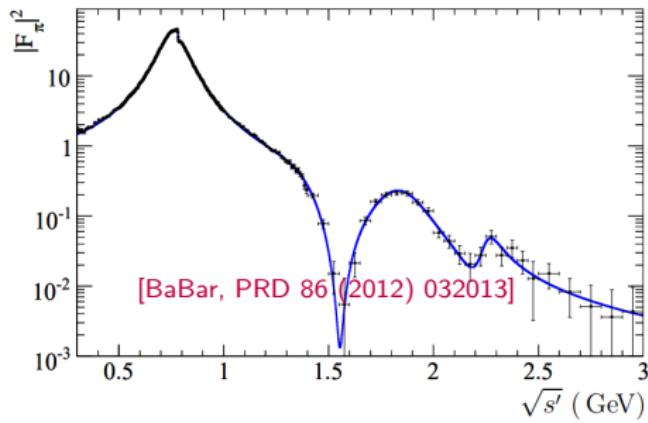
- $\alpha_{2\pi}$ is the 1st-order real polynomial
- $\alpha_{3\pi} = A_\omega / (m_\rho^2 - s)$ kills the artificial K-matrix zero

$$\text{LO} \quad \Rightarrow \quad \hat{A}_{2\pi} = \text{BW}_\rho (1 + k \text{BW}_\omega)$$

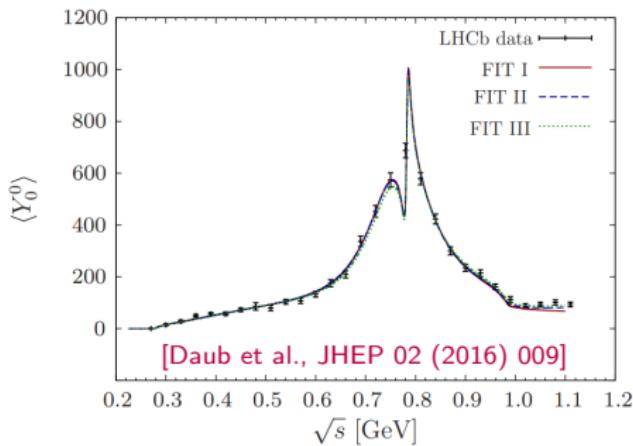
For the fit, the amplitude with **no-approximation** is used.

$\rho - \omega$ interference in pion form-factor

$$e^+ e^- \rightarrow \pi^+ \pi^-$$



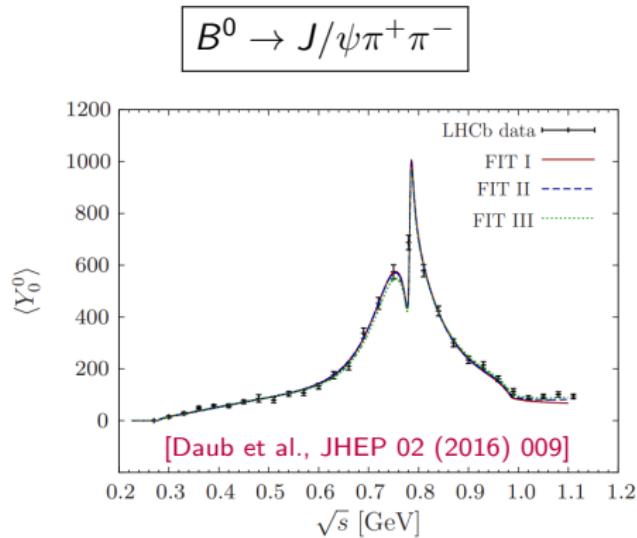
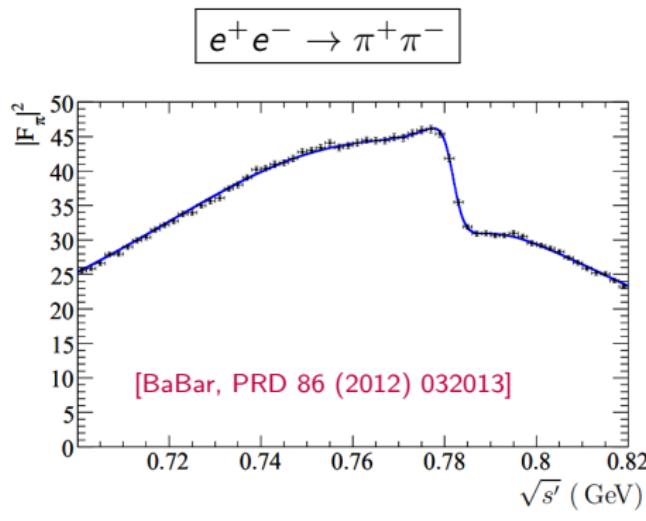
$$B^0 \rightarrow J/\psi \pi^+ \pi^-$$



A recipe for accounting for $\rho - \omega$ interference,
LO in isospin violation, [Gardner-O'Connell (1998)]

$$\text{BW}_\rho \rightarrow \text{BW}_\rho \left(1 + \frac{k_{\text{e.m.}} s}{m_\omega^2 - s - im_\omega \Gamma_\omega} \right)$$

$\rho - \omega$ interference in pion form-factor

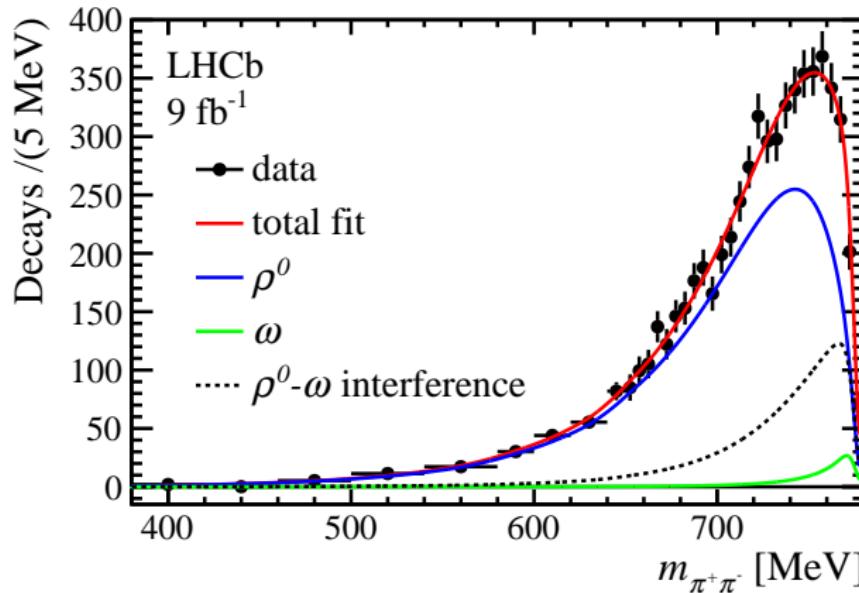


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The default fit

[LHCb, JHEP 08 (2020) 123]



- 3 parameter fit
 - $\chi^2/\text{ndf} = 24.7/32$
 - $R_\omega^{\text{all}} = 0.214 \pm 0.023$
due to the large interference
- !! Pure contribution
of ρ is
 $(78.6 \pm 2.3)\%$

using $\omega \rightarrow 2\pi$ (this) : $R_\omega = (1.93 \pm 0.44)\%$

using $\omega \rightarrow 3\pi$ (PDG) :
$$\frac{\mathcal{B}(\chi_{c1}(3872) \rightarrow \omega J/\psi) \mathcal{B}(\omega \rightarrow 2\pi)}{\mathcal{B}(\chi_{c1}(3872) \rightarrow \pi^+\pi^- J/\psi)} = (2.1 \pm 0.5)\%$$

Systematic studies

[LHCb, JHEP 08 (2020) 123]

Variations of the model

Fit type	χ^2/NDof	p-value	$\mathcal{R}_\omega^{\text{all}}$	\mathcal{R}_ω^0	$\mathcal{R}_{\omega/p}^0$	n_σ
Default	24.7/32	0.82	0.214 ± 0.023	0.019 ± 0.004	0.025 ± 0.006	8.1σ
$P_2 \neq 0$	24.6/31	0.78	0.206 ± 0.035	0.018 ± 0.006	0.023 ± 0.009	5.5σ
Gaussian $\chi_{c1}(3872)$	20.0/32	0.95	0.194 ± 0.024	0.016 ± 0.004	0.020 ± 0.006	7.3σ
cubic $\epsilon(m_{\pi^+\pi^-})$	24.5/32	0.83	0.221 ± 0.023	0.021 ± 0.005	0.027 ± 0.007	8.1σ
had.ID corrections	24.6/32	0.82	0.214 ± 0.023	0.019 ± 0.004	0.025 ± 0.006	8.1σ
BDT selection	24.6/32	0.82	0.207 ± 0.022	0.018 ± 0.004	0.023 ± 0.006	7.9σ
$\sigma(m_{\pi^+\pi^-}) \times 1.0$	26.6/32	0.74	0.213 ± 0.023	0.019 ± 0.004	0.025 ± 0.006	8.1σ
$\sigma(m_{\pi^+\pi^-}) \times 1.14$	22.6/32	0.89	0.215 ± 0.023	0.020 ± 0.004	0.026 ± 0.006	8.1σ
$m_{\pi^+\pi^-} < 775$ MeV	18.0/31	0.97	0.196 ± 0.024	0.016 ± 0.004	0.021 ± 0.006	7.1σ
$\cos \theta_X < 0$	26.9/32	0.72	0.211 ± 0.035	0.019 ± 0.007	0.024 ± 0.010	5.2σ
$\cos \theta_X > 0$	42.2/32	0.11	0.217 ± 0.030	0.021 ± 0.006	0.027 ± 0.009	4.2σ
NR prod. of 2π	24.7/32	0.82	0.214 ± 0.022	0.019 ± 0.004	0.025 ± 0.006	8.1σ
D -wave free	24.5/31	0.79	0.210 ± 0.029	0.017 ± 0.005	0.021 ± 0.007	7.8σ
D -wave fixed at 4%	24.5/32	0.82	0.208 ± 0.023	0.018 ± 0.004	0.023 ± 0.006	7.9σ
ρ'	25.1/32	0.80	0.209 ± 0.023	0.018 ± 0.004	0.024 ± 0.006	8.1σ
$R_{\text{prod}} = 0$ GeV $^{-1}$	24.7/32	0.82	0.209 ± 0.023	0.019 ± 0.004	0.024 ± 0.006	7.9σ
$R_{\text{prod}} = 30$ GeV $^{-1}$	24.6/32	0.82	0.229 ± 0.022	0.021 ± 0.004	0.028 ± 0.006	8.7σ
$R = 1.3$ GeV $^{-1}$	24.7/32	0.82	0.216 ± 0.022	0.020 ± 0.004	0.026 ± 0.006	8.2σ
$R = 1.6$ GeV $^{-1}$	24.7/32	0.82	0.212 ± 0.023	0.019 ± 0.004	0.025 ± 0.006	8.0σ
GS model	24.8/32	0.81	0.221 ± 0.024	0.021 ± 0.005	0.028 ± 0.007	7.8σ
Summary			$0.214 \pm 0.023 \pm 0.020$	$0.019 \pm 0.004 \pm 0.003$	$0.025 \pm 0.006 \pm 0.005$	$> 7.1\sigma$

Testing the Gounaris-Sakurai shape

[LHCb, JHEP 08 (2020) 123]

[Barkov et al., NPB 256 (1985) 365-384], [BaBar, PRD86, 032013 (2012)]

$$\mathcal{A}_{2\pi} = p_\pi(s) \left\{ \text{BW}_\rho^{\text{GS}}(1 + A_\omega e^{i\phi_\omega} \text{BW}_\omega) + A_{\rho'} e^{i\phi_{\rho'}} \text{BW}_{\rho'}^{\text{GS}} \right\}$$

$\rho - \omega$ is equivalent to the coupled-channel, ρ' is added by sum of BW.

$$\text{BW}_\rho^{\text{GS}} = \frac{1}{m_\rho^2 - s - im_\rho \tilde{\Gamma}_\rho(s)}$$

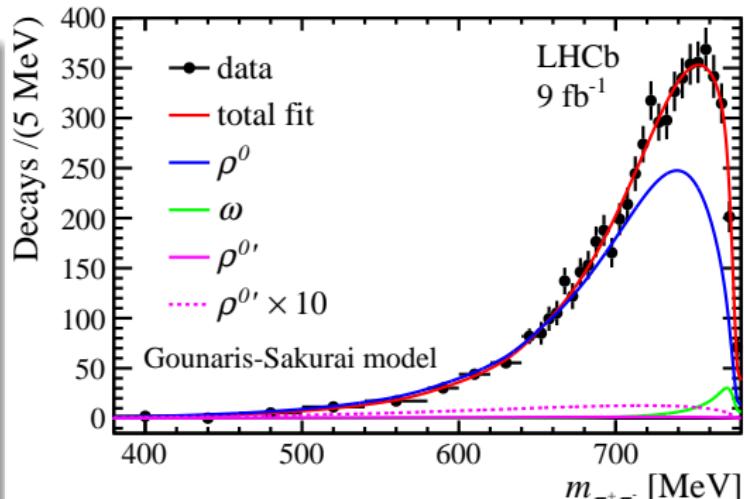
with $i\tilde{\Gamma}_\rho(s)$ calculated from using dispersion relations ()

- P-wave

$$\text{Im}(i\tilde{\Gamma}_\rho(s)) = \Gamma \frac{m_\rho}{\sqrt{s}} \left[\frac{p(s)}{p(m^2)} \right]^3$$

- Twice-subtracted DR

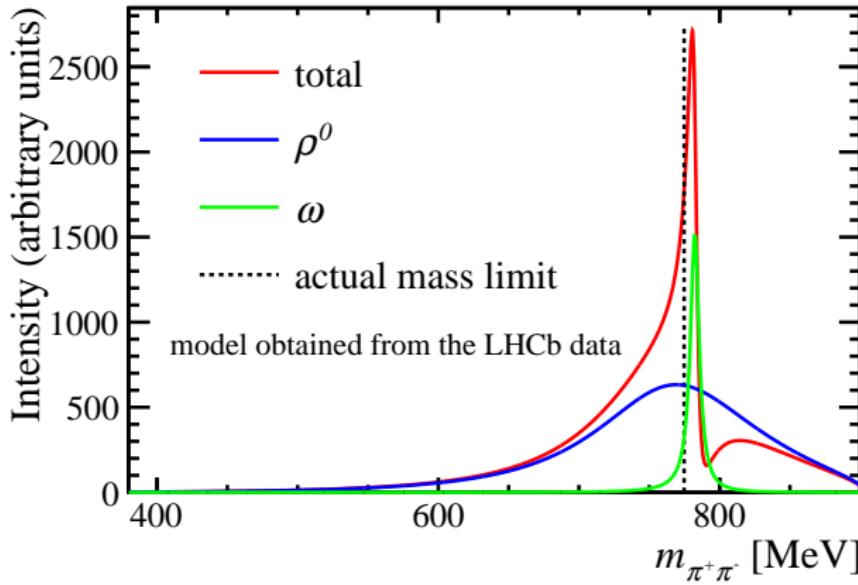
- Two subt.constants are fixed at $\sqrt{s} = m_\rho$,



- $\chi^2/\text{ndf} = 24.8/32$
- Compatible results

Size of the isospin violation

[LHCb, JHEP 08 (2020) 123]



- Extend the upper limit of ph.sp: $m_{\chi_{c1}} \rightarrow 4 \text{ GeV}$
- Ratio of the integrals:

$$R_{\omega/\rho} = 0.18 \pm 0.05(\text{stat})$$

The estimate for couplings:

$$\frac{g_{\chi_{c1} \rightarrow \rho J/\psi}}{g_{\chi_{c1} \rightarrow \omega J/\psi}} = \sqrt{\frac{\mathcal{B}(\omega \rightarrow 2\pi)}{\mathcal{B}(\rho \rightarrow 2\pi)}} \frac{1}{R_{\omega/\rho}} = 0.29 \pm 0.04$$

Can be compared to $g_{\psi(2S) \rightarrow \pi^0 J/\psi} / g_{\psi(2S) \rightarrow \eta J/\psi} = 0.045 \pm 0.001$

Summary

- Significant, $> 7\sigma$ contribution of ω is found in $\chi_{c1} \rightarrow J/\psi \pi^+ \pi^-$,
- Pure contribution of ρ : $R_\rho = (78.6 \pm 2.3 \pm 2.0)\%$,
- $R_\omega = (1.9 \pm 0.4 \pm 0.3)\%$
- Consistent with the direct measurements using $\omega \rightarrow 3\pi$

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$$g_{\chi_{c1} \rightarrow \rho J/\psi} / g_{\chi_{c1} \rightarrow \omega J/\psi} \approx 0.3$$

The isospin violation is larger by an order of magnitude than for ordinary charmonium state

- Possible in molecular picture due to the proximity to $D^{*0}\bar{D}^0$ rather than $D^{*+}\bar{D}^-$ [Törnqvist(2003, 2004), Voloshin(2004), Swanson(2004), Suzuki(2005), ...]
- Also possible in compact tetraquark picture[Terasaki,(2007), Maiani(2018,2020)]

Thank you for the attention

How well we know the $\pi\pi$ P-wave

Three shapes are consistent:

- GKY: combined analysis with th. constrains
- GS (P-wave CM, 2 subtracted)
- P-wave BW with Blatt-Weisskopf barrier factor, $R = 1.5 \text{ GeV}^{-1}$.

$$B_1(p) = p^2 \frac{1}{1 + (pR)^2}.$$

- CDF-II and Belle used $R = 1.5 \text{ GeV}^{-1}$
- PDG(1983): $R = 5.3^{+0.9}_{-0.7} \text{ GeV}^{-1}$

