

# Is $Z_{cs}(3982)$ a molecular partner of $Z_c(3900)$ and $Z_c(4020)$ ?

Vadim Baru

Institut für Theoretische Physik II, Ruhr-Universität Bochum Germany

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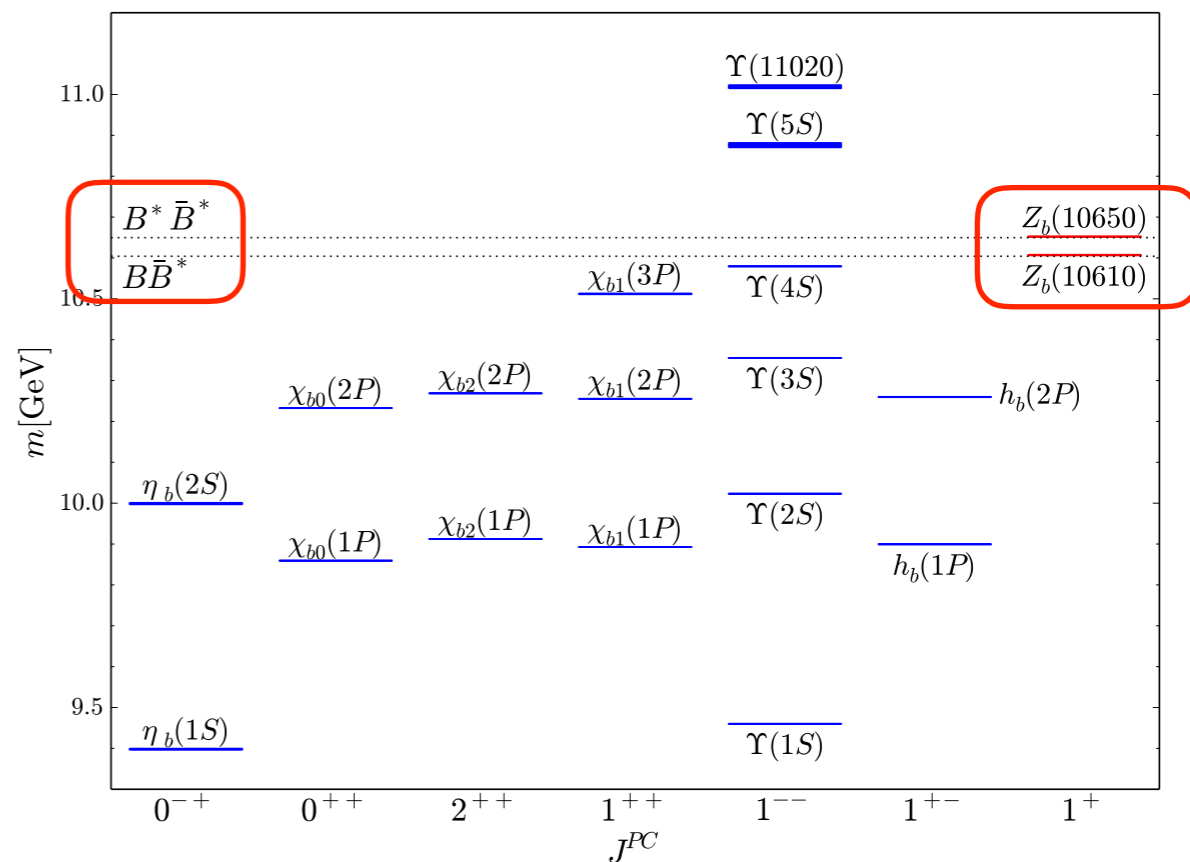
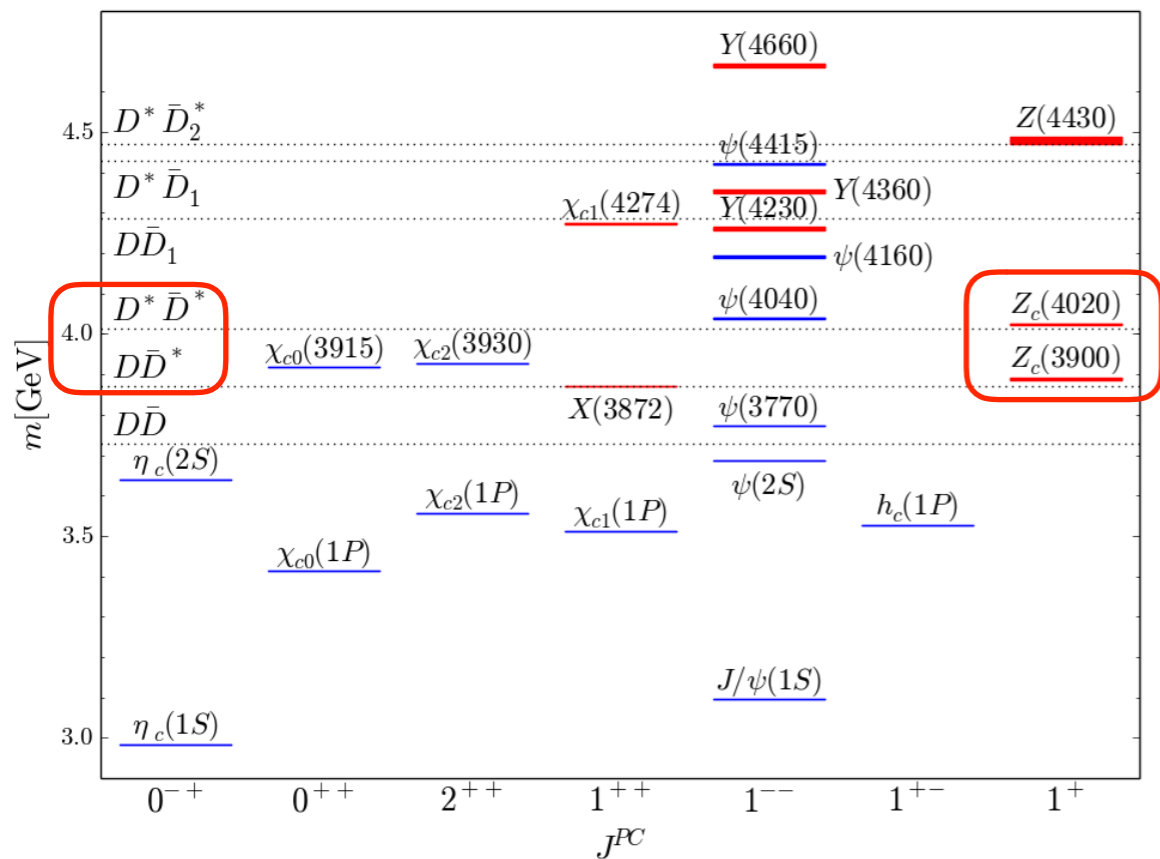
DARMSTADT

in collaboration with

E. Epelbaum, A.A. Filin, C. Hanhart, and A.V. Nefediev

Key Ref: PRD 105, 034014 (2022)

# “Tetraquarks” in quarkonium spectrum

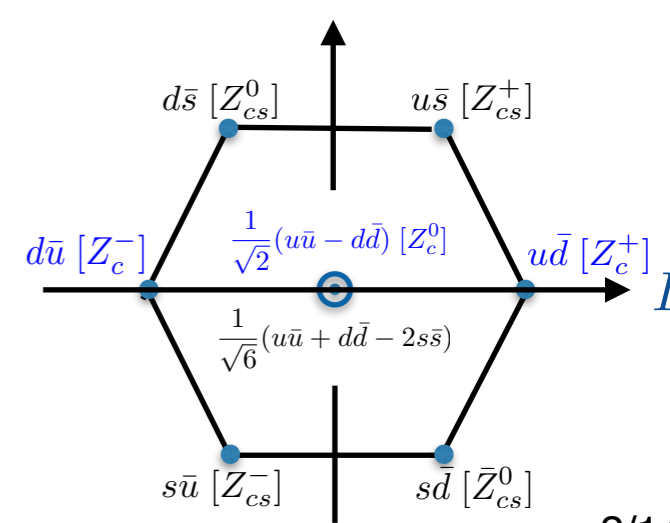


- Twin I=1 nearthreshold states  $Z_b(10610)/Z_b(10650)$  by Belle and  $Z_c(3900)/Z_c(4020)$  by BES III

- clearly exotic: seen in charged modes  $\pi^\pm h_b(mP), \pi^\pm \Upsilon(nS)$  Belle  $\implies b\bar{b}u\bar{d}, \dots$
- $J^{PC} = 1^{+-}$   $\pi^\pm J/\psi, \pi^\pm \Psi(2S)$  BES III  $\implies c\bar{c}u\bar{d}, \dots$

- Natural candidates for hadronic molecules

$\implies$  Flavour SU(3) and heavy-quark spin (HQSS) partners must exist

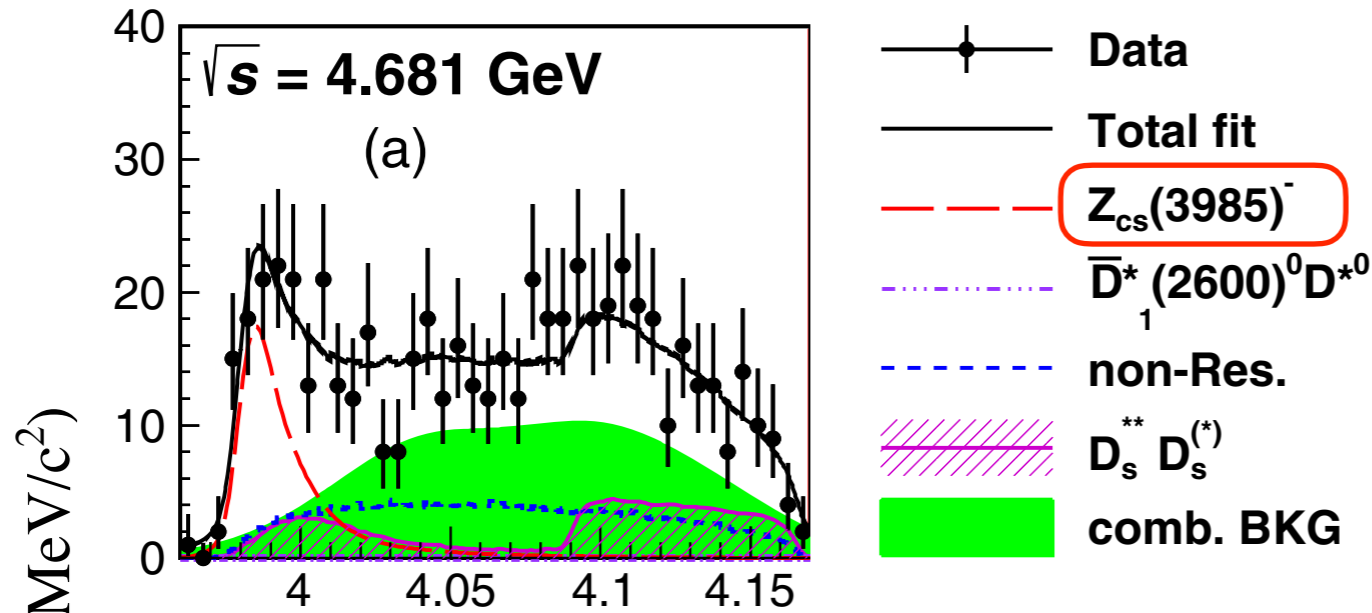


our works (2016-2022), Hidalgo-Duque et al.(2013), Guo et al (2015), Meng et al 2020, ...

# $Z_{cs}(3982) — c\bar{c}s\bar{u}$ state by BESIII

PRL 126 (2021), 102001

$$e^+e^- \rightarrow K^+ (D_s^- D^{*0} + D_s^{*-} D^0)$$

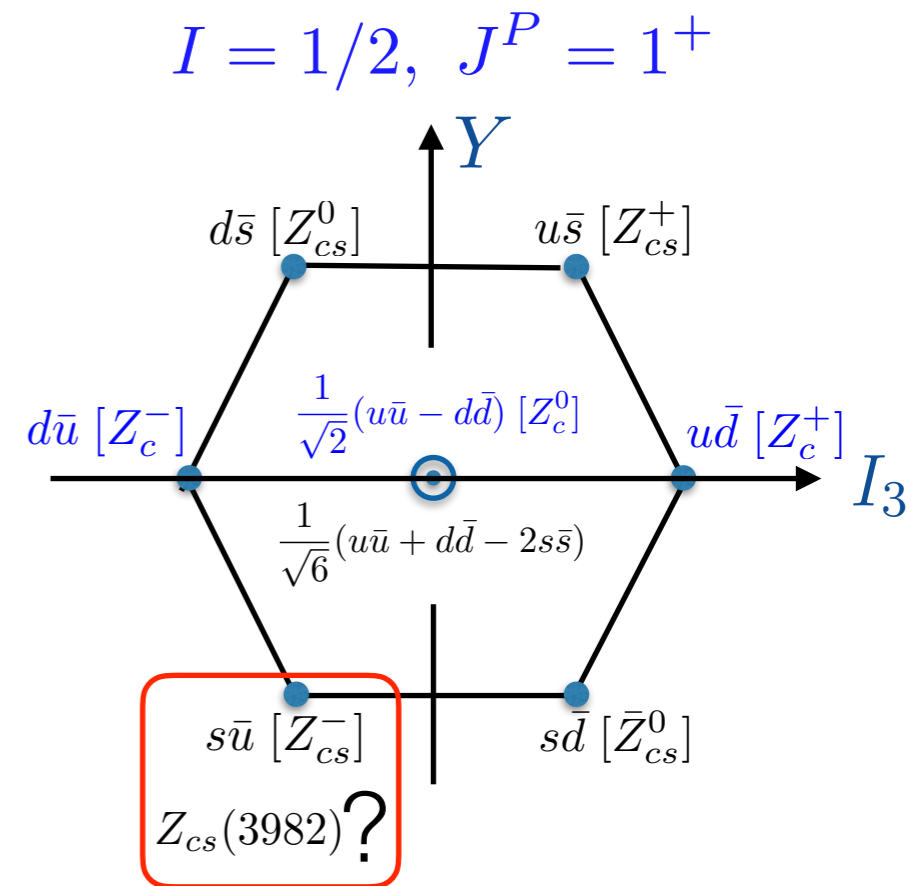
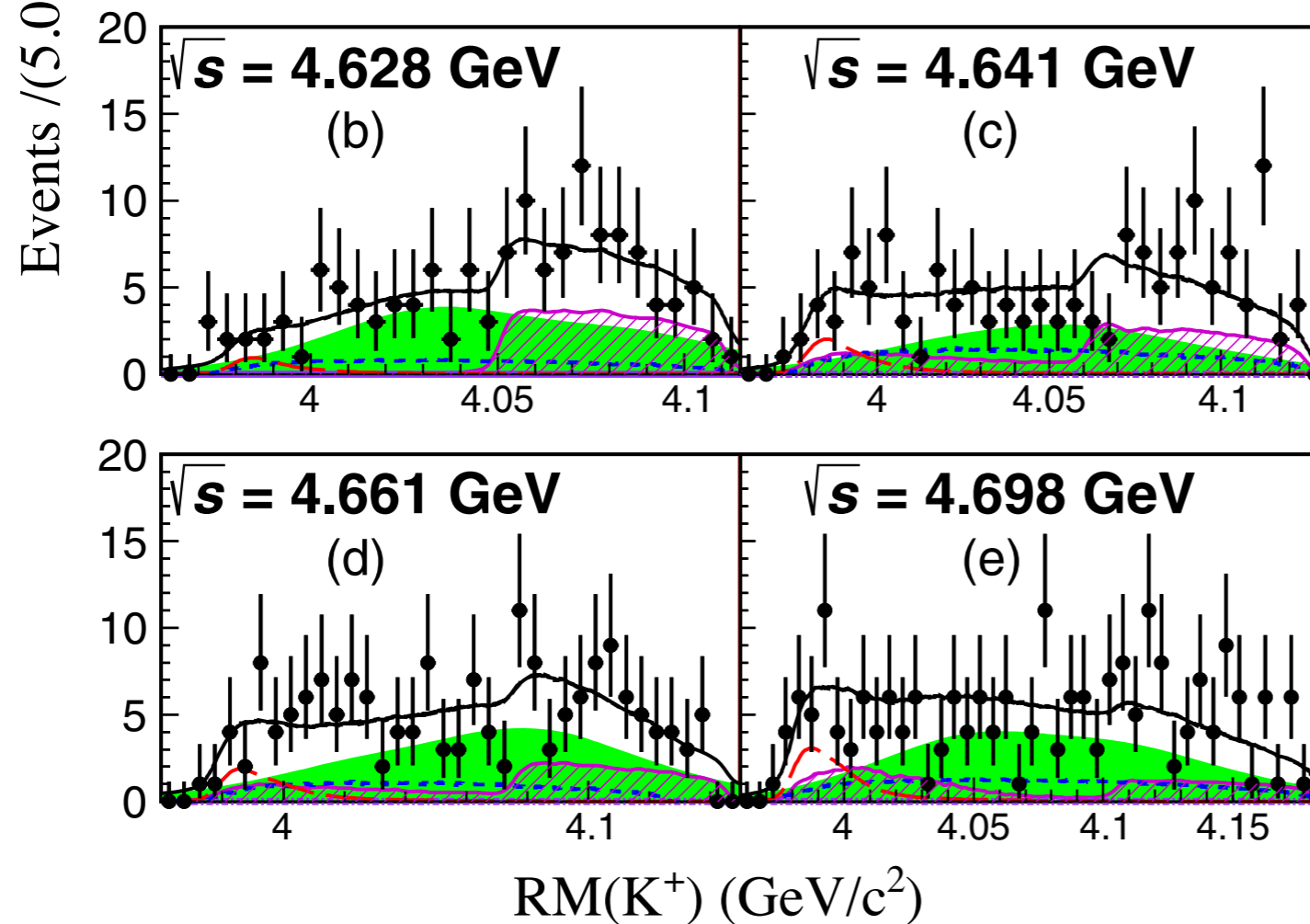


Breit-Wigner type fit:

$$M = 3982.5_{-2.6}^{+1.8} \pm 2.1 \text{ MeV}$$

$$\Gamma = 12.8_{-4.4}^{+5.3} \pm 3.0 \text{ MeV}$$

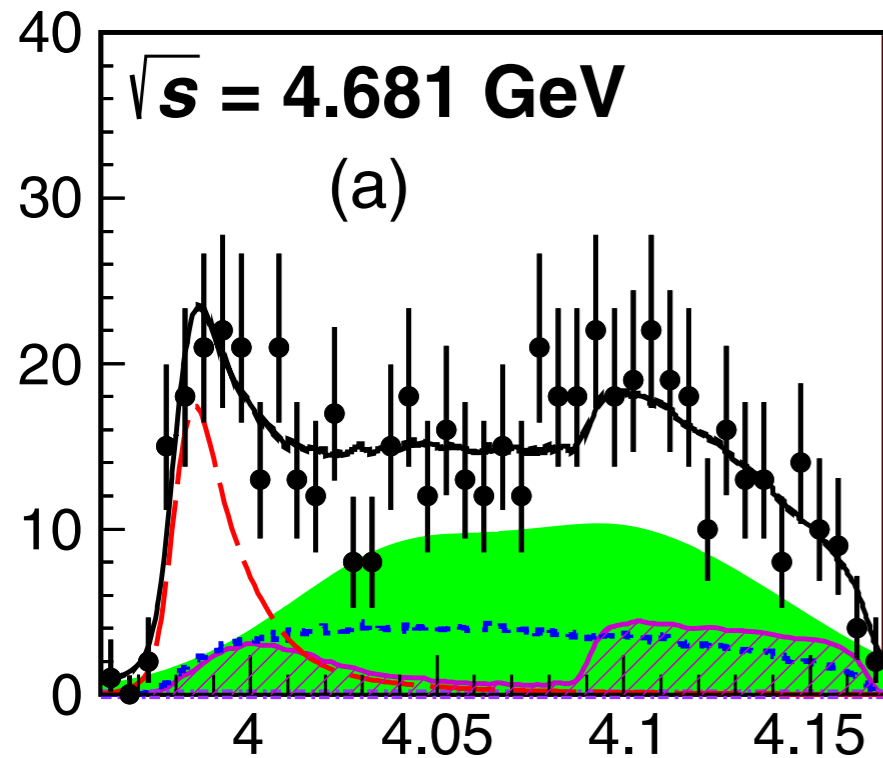
Near  $D_s^- D^{*0} + D_s^{*-} D^0$  thr.



Appealing possibility that  $Z_{cs}$  is a strange molecular partner of  $Z_c(3900)$

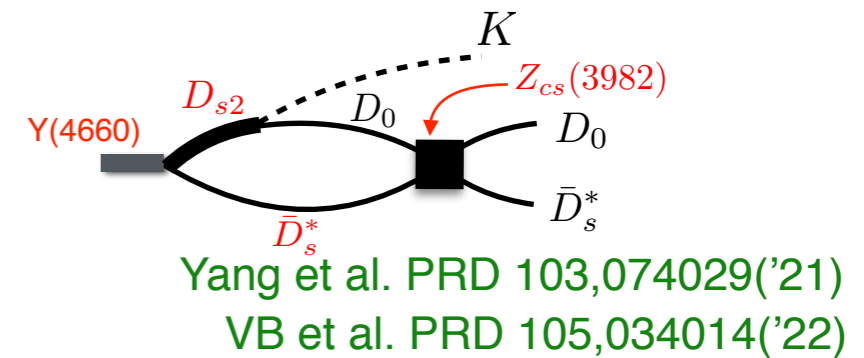
Properties of  $Z_{cs}$ : A coupled-channel EFT analysis of the line shape is needed!

# From $Z_{cs}(3982)$ to molecular nature of vector mesons



## BESIII

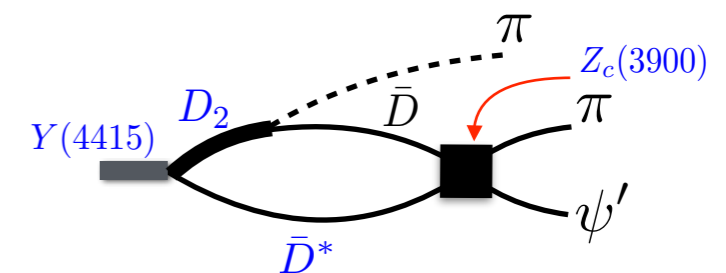
- BESIII data are measured in the energy range 4.628-4.698  
 $\implies$  Excitation of  $J^{PC}=1^{--}$   $Y(4660)$   $\Gamma = 62_{-7}^{+9}$  MeV
- $D_{s2}(2573)$  and  $D_{s1}(2536)$  are important parts of the background



- Triangle singularity from  $D_{s2}D_s^*D$  is near  $\sqrt{s} = 4.681$  GeV

$\implies Y(4660)$  should have sizeable coupling to  $D_{s2}D_s^* \implies D_{s2}D_s^*$  molecule!

- Complete analogy to  $Z_c$ : BESIII observation of  $Z_c$  at 4415 MeV suggests  $Y(4415)$  is a  $D_2D^*$  molecule



Similar argument for  $Y(4230)$ : Wang et al. PRL111,132003('13), Clevén et al. PRD90,074039('14)

If  $Z_{cs}$  is a strange molecular partner of  $Z_c(3900)$   
 $\implies Y(4660)$  is likely to be a strange partner of  $Y(4415)$

# Contact SU(3) flavour and HQSS Lagrangian

- $H\bar{H} \rightarrow H\bar{H}$  LO contact Lagrangian

$$\mathcal{L} = -\frac{C_{00}}{4} \text{Tr}[\bar{H}_a^\dagger H_a^\dagger H_b \bar{H}_b] - \frac{C_{01}}{4} \text{Tr}[\bar{H}_a^\dagger \sigma^i H_a^\dagger H_b \sigma^i \bar{H}_b] \\ - \frac{C_{10}}{4} \text{Tr}[\bar{H}_a^\dagger \lambda_{aa'}^A H_{a'}^\dagger H_b \lambda_{bb'}^A \bar{H}_{b'}] - \frac{C_{11}}{4} \text{Tr}[\bar{H}_a^\dagger \lambda_{aa'}^A \sigma^i H_{a'}^\dagger H_b \lambda_{bb'}^A \sigma^i \bar{H}_{b'}]$$

Al Fiky et al. PLB640('06)

Grinstein et al., NPB380('92);

- SU(3) flavour symmetry:  $H_a \sim (Q\bar{u}, Q\bar{d}, Q\bar{s}) \sim (D^0, D^+, D_s^+)$

- HQSS:  $H_a = P_a + V_a^i \sigma^i$   $\bar{H}_a = \bar{P}_a - \bar{V}_a^i \sigma^i$

$$P = D, D_s$$

$$V = D^*, D_s^*$$

# Contact SU(3) flavour and HQSS Lagrangian

- $H\bar{H} \rightarrow H\bar{H}$  LO contact Lagrangian

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- HQSS:  $H_a = P_a + V_a^i \sigma^i$   $\bar{H}_a = \bar{P}_a - \bar{V}_a^i \sigma^i$   $P = D, D_s$   $V = D^*, D_s^*$

- 4 Param's in general, but only two are relevant for (I=1/2)  $Z_{cs}$  and (I=1)  $Z_c$

$$C_d = \frac{1}{8}(C_{11} + C_{10}), \quad C_f = \frac{1}{8}(C_{11} - C_{10})$$

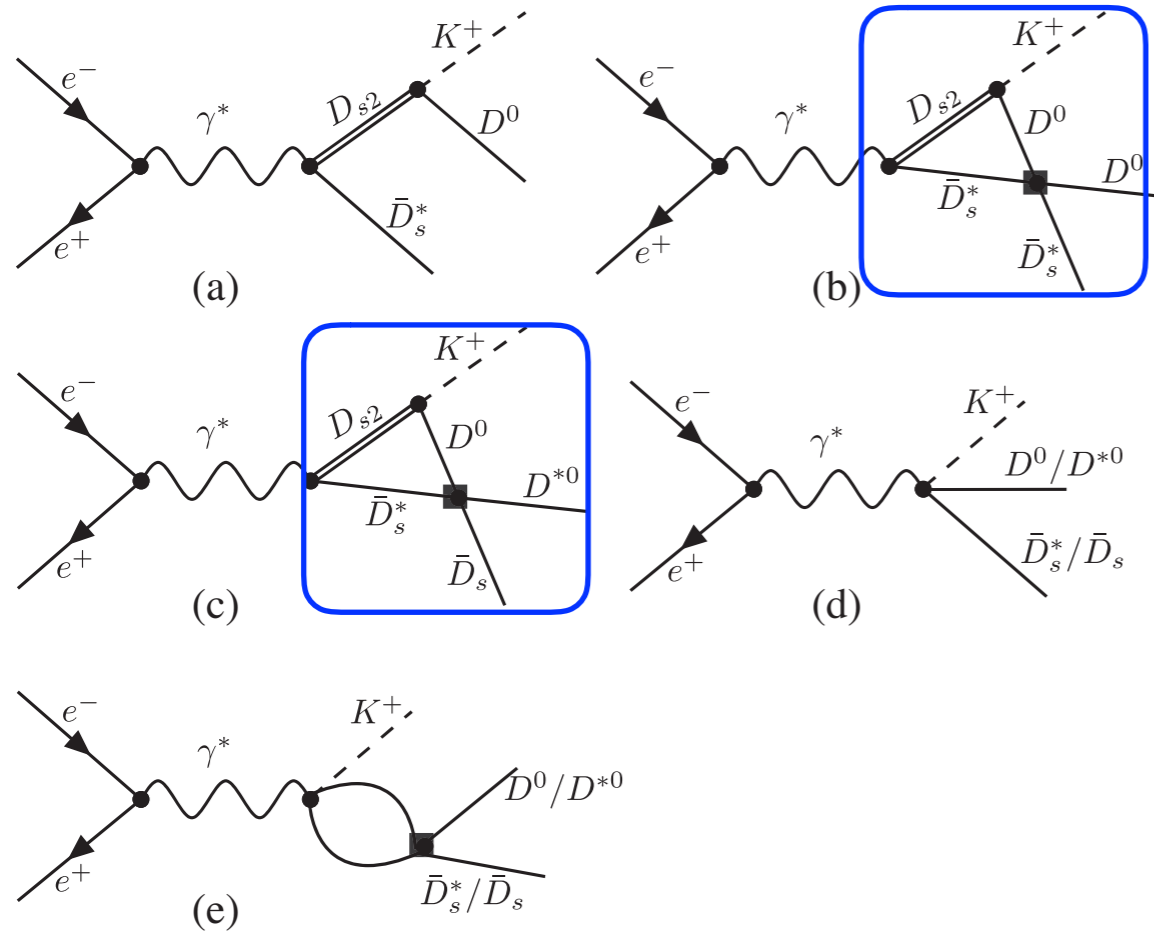
- Resulting effective potentials to be used in the dynamical equations

$$Z_{cs} \quad 1^+ : \{ \bar{D}_s D^*, D \bar{D}_s^*, D^* \bar{D}_s^* \}$$

$$V^{\text{CT}}[1^+] = \begin{pmatrix} C_d + \frac{1}{2}C_f & \frac{1}{2}C_f & -\frac{1}{\sqrt{2}}C_f \\ \frac{1}{2}C_f & C_d + \frac{1}{2}C_f & \frac{1}{\sqrt{2}}C_f \\ -\frac{1}{\sqrt{2}}C_f & \frac{1}{\sqrt{2}}C_f & C_d \end{pmatrix}$$

$$Z_c \quad 1^{+-} : \left\{ \frac{D\bar{D}^* - D^*\bar{D}}{\sqrt{2}}, D^*\bar{D}^* \right\}$$

$$V^{\text{CT}}[1^{+-}] = \begin{pmatrix} C_d & C_f \\ C_f & C_d \end{pmatrix}$$



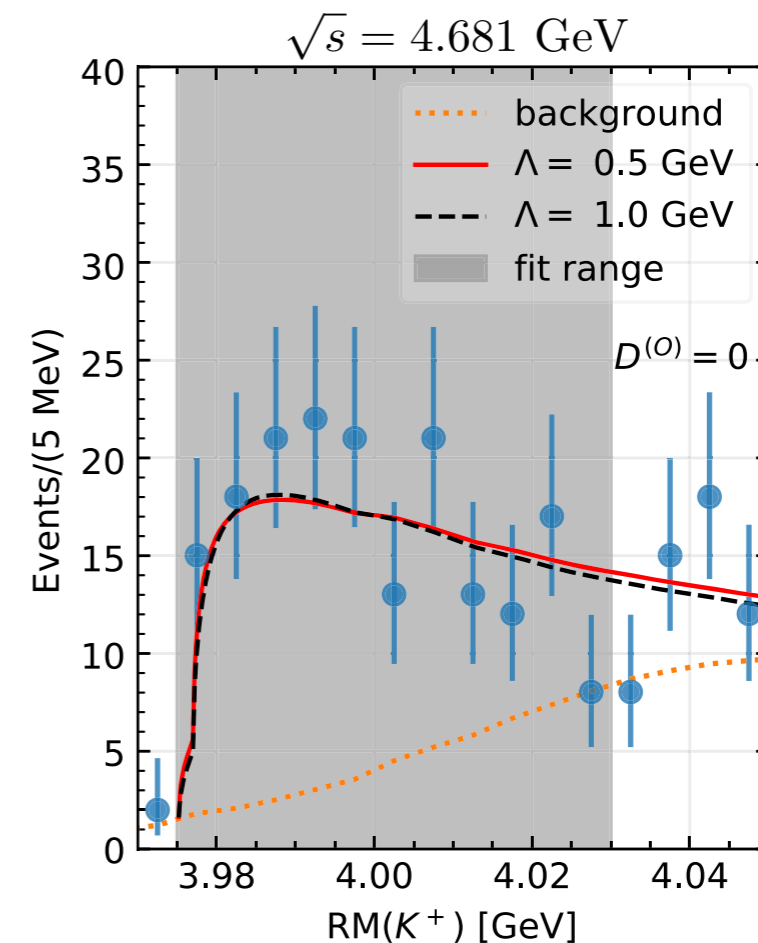
## Production

- Triangle singularity from  $D_{s2}D_s^*D_0$  is nearby
- Pointlike production

- Contact EFT neglecting coupled-ch. to  $D_s^*D^*$  and  $D^*D^*$
- Limited energy range in fits, 50 MeV above  $D_sD^*/D_s^*D$

⇒ Extracted poles are consistent with that  $Z_{cs}$  is the SU(3) flavour partner of  $Z_c(3900)$

– Supported by a combined  $Z_c(3900)$  and  $Z_{cs}$  analysis  
by Du et al. PRD105, 074018('22)]

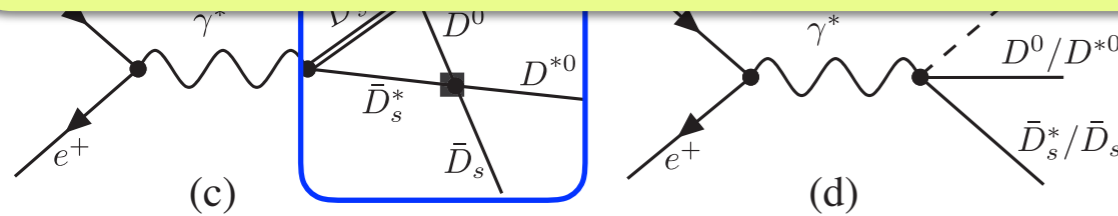


$$Z_{cs} \quad 1^+ : \{ \bar{D}_s D^*, D \bar{D}_s^*, \cancel{D^* \bar{D}_s^*} \}$$

$$Z_c \quad 1^{+-} : \left\{ \frac{D \bar{D}^* - D^* \bar{D}}{\sqrt{2}}, \cancel{D^* \bar{D}^*} \right\}$$

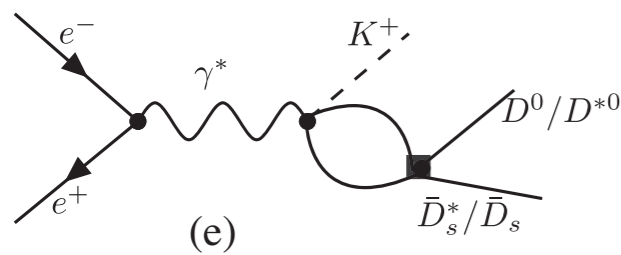
$$V^{\text{CT}}[1^+] = \begin{pmatrix} C_d + \frac{1}{2}C_f & \frac{1}{2}C_f & -\frac{1}{\sqrt{2}}C_f \\ \frac{1}{2}C_f & C_d + \frac{1}{2}C_f & \frac{1}{\sqrt{2}}C_f \\ -\frac{1}{\sqrt{2}}C_f & \frac{1}{\sqrt{2}}C_f & C_d \end{pmatrix}$$

$$V^{\text{CT}}[1^{+-}] = \begin{pmatrix} C_d & C_f \\ C_f & C_d \end{pmatrix}$$



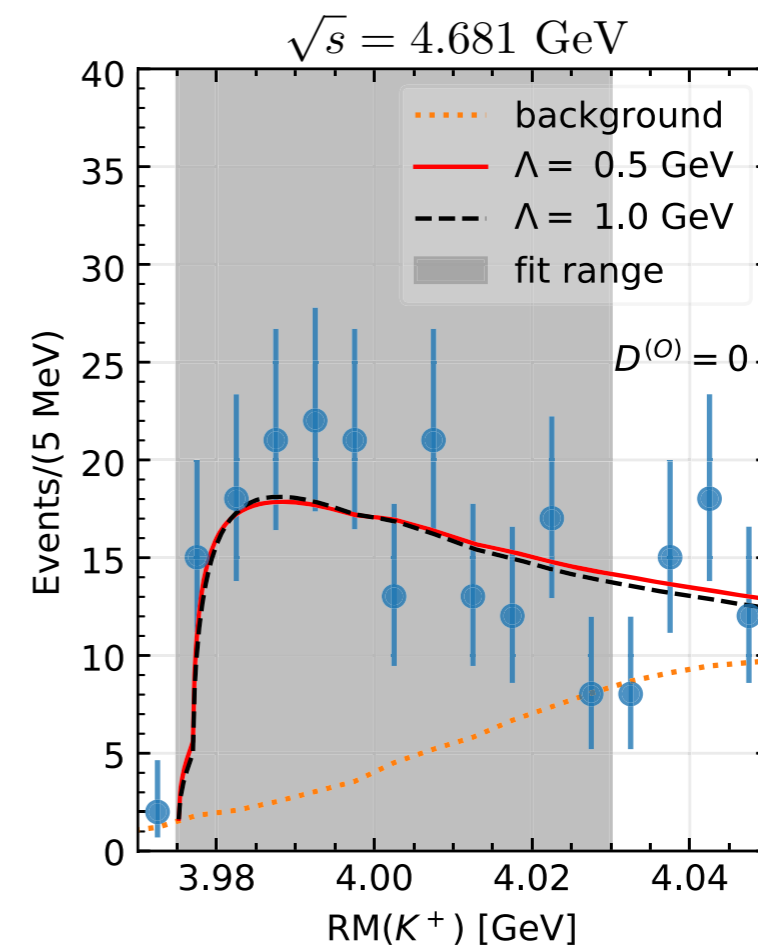
—Contact EFT neglecting coupled-ch. to  $D_s^* D^*$  and  $D^* D^*$

—Limited energy range in fits, 50 MeV above  $D_s D^*/D_s^* D$



⇒ Extracted poles are consistent with that  $Z_{cs}$  is the SU(3) flavour partner of  $Z_c(3900)$

—Supported by a combined  $Z_c(3900)$  and  $Z_{cs}$  analysis by Du et al. PRD105, 074018('22)





# Goals of our study

- Include all coupled channels, i.e.  $\bar{D}_s D^*$ ,  $D \bar{D}_s^*$ ,  $D^* \bar{D}_s^*$
- Complete analysis in the whole energy range covered by the BES III data  
 $\implies$  reliable prediction for a possible spin partner state near  $D^* \bar{D}_s^*$

- Include missing production mechanisms:

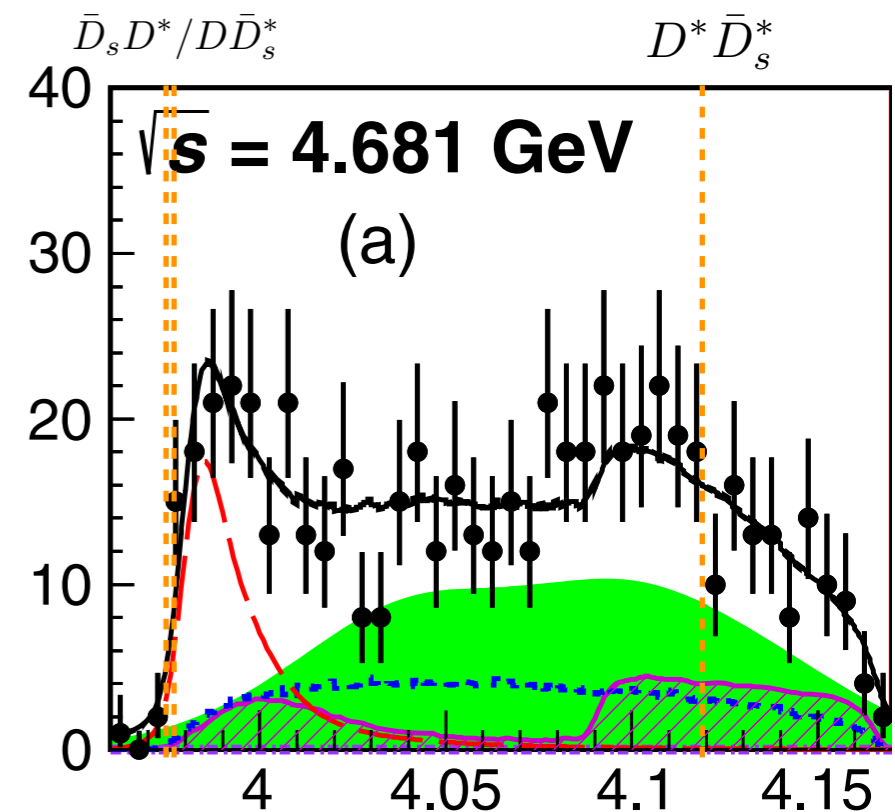
— the role of  $D_{s1}$

— triangle singularities from  $D_{s2} D_s^* D$ ,  $D_{s2} D_s^* D^*$  and  $D_{s1} D_s^* D^*$

All in the range of BESIII data!

Not covered by our analysis

- — possible SU(3) breaking due to lightest pseudoscalar GB octet
- — combined analysis of  $Z_c(3900)/Z_c(4020)$  and  $Z_{cs}(3982)$  data



# Coupled-channel production $e^+e^- \rightarrow Y(4660) \rightarrow K^+(D_s^- D^{*0} + D_s^{*-} D^0)$

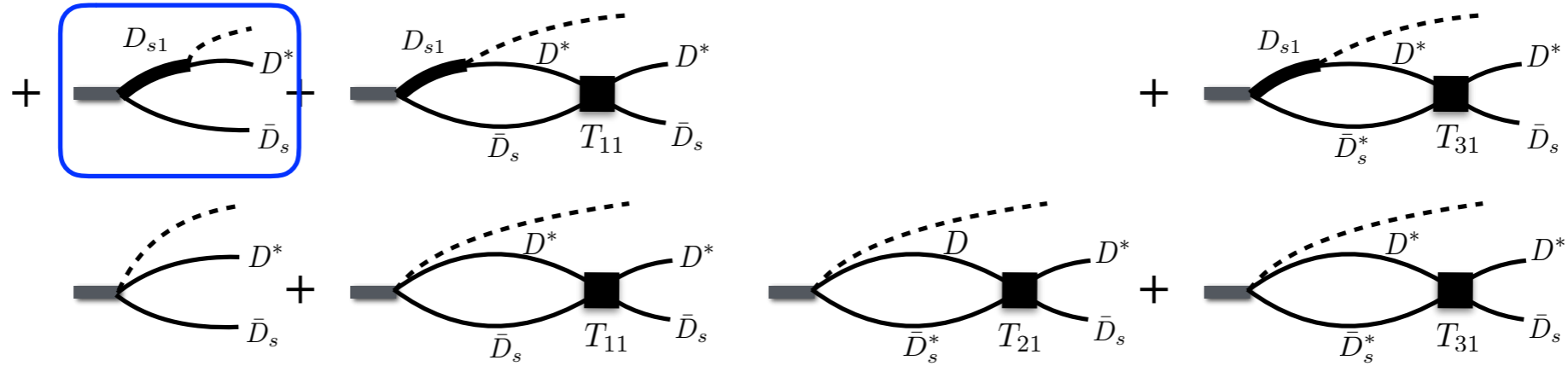
$$\begin{aligned}
 M_{Y \rightarrow K D^* \bar{D}_s} = & \text{Diagram 1} + \text{Diagram 2} \\
 & + \text{Diagram 3} + \text{Diagram 4} + \text{Diagram 5} + \text{Diagram 6} \\
 & + \text{Diagram 7} + \text{Diagram 8} + \text{Diagram 9} + \text{Diagram 10} \\
 M_{Y \rightarrow K D \bar{D}_s^*} = & \text{Diagram 11} + \text{Diagram 12} + \text{Diagram 13} + \text{Diagram 14} \\
 & + \text{Diagram 15} + \text{Diagram 16} + \text{Diagram 17} + \text{Diagram 18} \\
 & + \text{Diagram 19} + \text{Diagram 20} + \text{Diagram 21} + \text{Diagram 22}
 \end{aligned}$$

—  $Z_{cs}(3982)$  is a pole in the coupled-channel  $\{\bar{D}_s D^*, D \bar{D}_s^*, D^* \bar{D}_s^*\}$  scattering amplitude

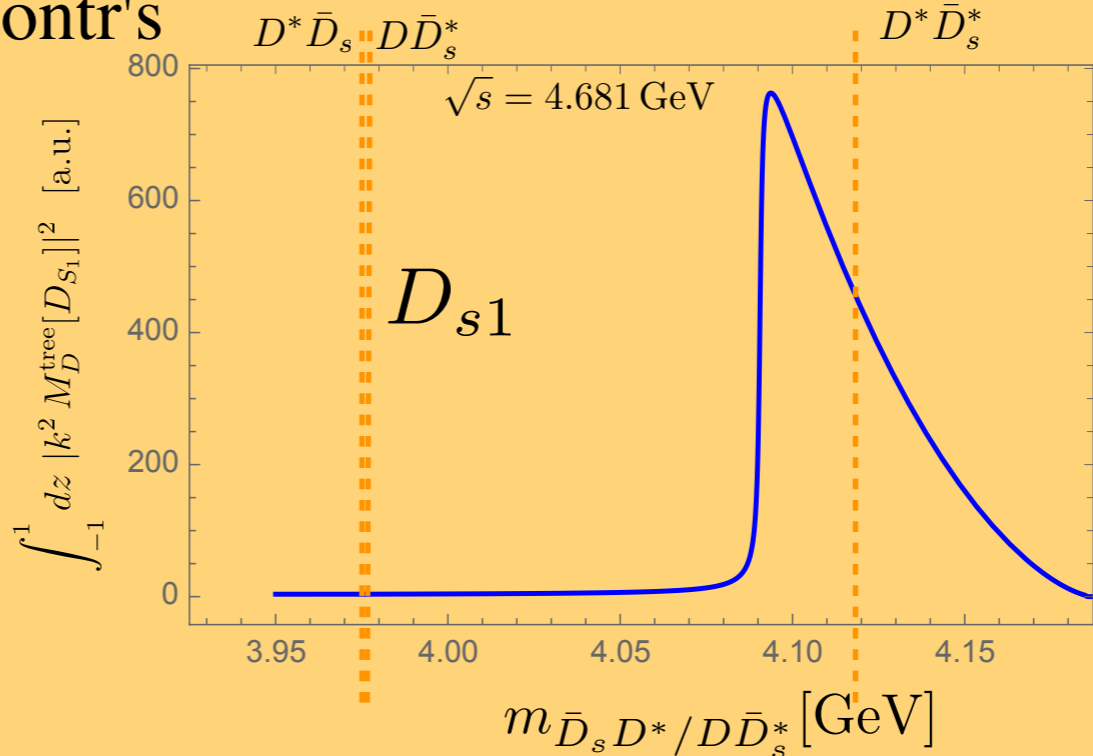
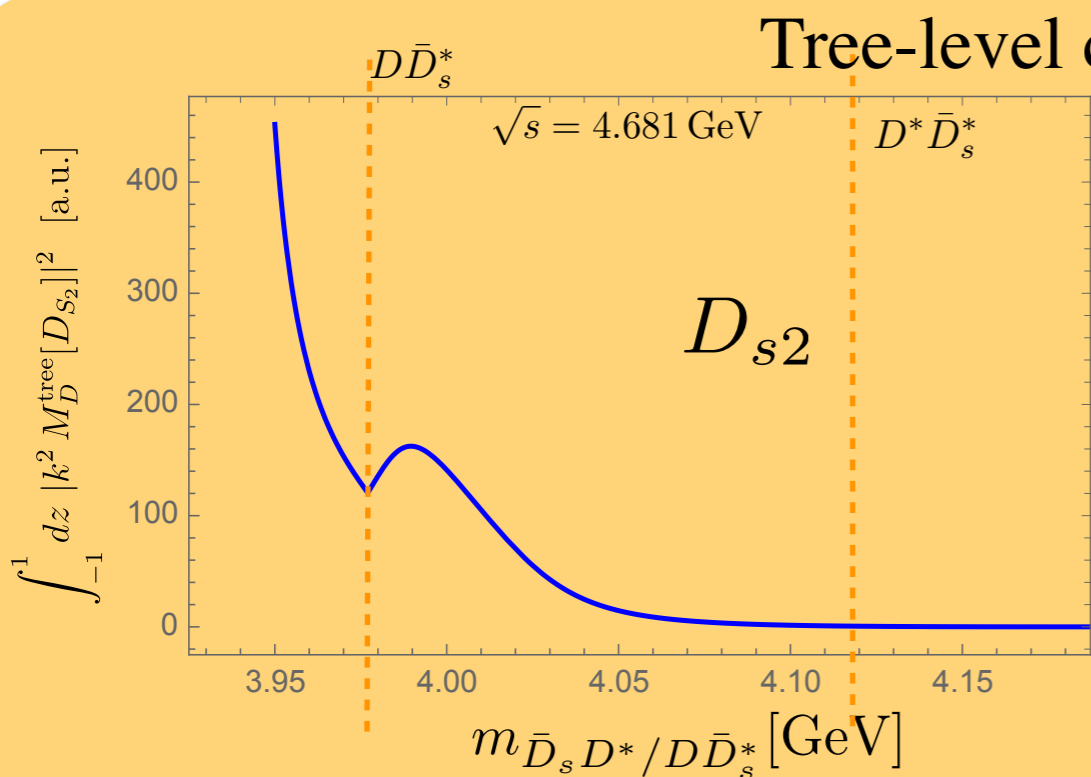
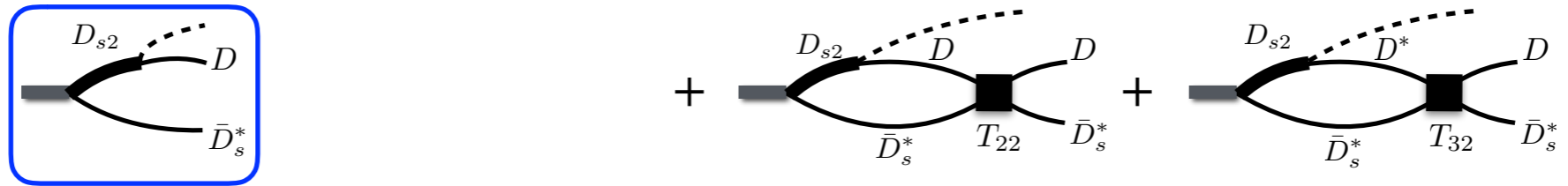
$$T_{\alpha\beta}(\sqrt{s}, p, p') = V_{\alpha\beta}(p, p') - \sum_{\gamma} \int \frac{d^3 q}{(2\pi)^3} V_{\alpha\gamma}(p, q) G_{\gamma}(\sqrt{s}, q) T_{\gamma\beta}(\sqrt{s}, q, p')$$

# Coupled-channel production $e^+e^- \rightarrow Y(4660) \rightarrow K^+(D_s^- D^{*0} + D_s^{*-} D^0)$

$$M_{Y \rightarrow KD^* \bar{D}_s} =$$



$$M_{Y \rightarrow KDD\bar{D}_s^*} =$$



Amplitude

# Coupled-channel production $e^+e^- \rightarrow Y(4660) \rightarrow K^+(D_s^- D^{*0} + D_s^{*-} D^0)$

$$\begin{aligned}
 M_{Y \rightarrow K D^* \bar{D}_s} = & \text{[Diagram 1]} + \text{[Diagram 2]} \\
 & + \text{[Diagram 3]} + \text{[Diagram 4]} + \text{[Diagram 5]} + \text{[Diagram 6]} \\
 & + \text{[Diagram 7]} + \text{[Diagram 8]} + \text{[Diagram 9]} + \text{[Diagram 10]} \\
 M_{Y \rightarrow K D \bar{D}_s^*} = & \text{[Diagram 11]} + \text{[Diagram 12]} + \text{[Diagram 13]} + \text{[Diagram 14]} \\
 & + \text{[Diagram 15]} + \text{[Diagram 16]} + \text{[Diagram 17]} + \text{[Diagram 18]} \\
 & + \text{[Diagram 19]} + \text{[Diagram 20]} + \text{[Diagram 21]} + \text{[Diagram 22]}
 \end{aligned}$$

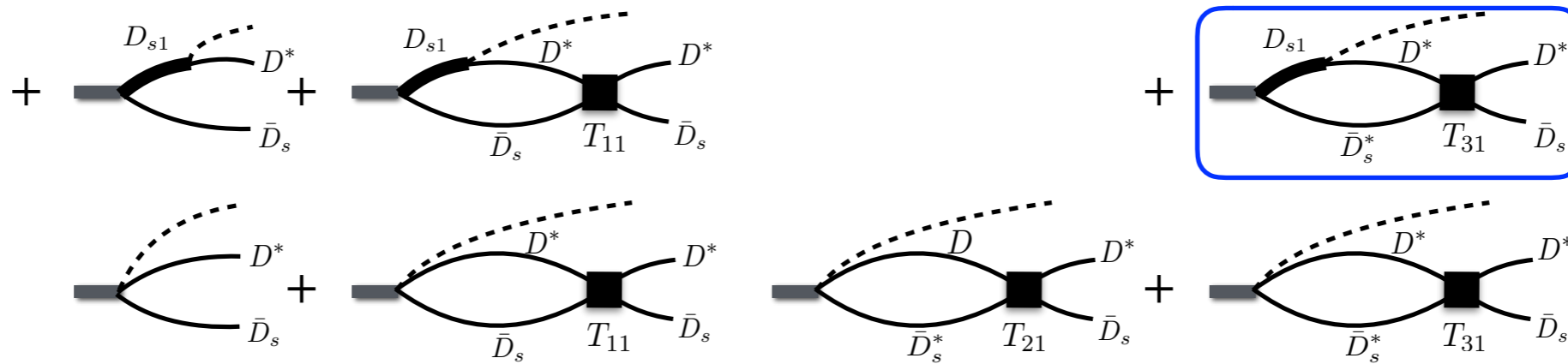
The diagrams represent various scattering amplitudes  $T_{ij}$  in the coupled-channel formalism. Each diagram shows a central interaction vertex (black square) connected to external lines representing particles:  $D_s$  (solid),  $D_s^*$  (dashed),  $D$  (solid), and  $\bar{D}_s$  (solid). The diagrams are arranged in a grid, with the first row showing the production of  $D^* \bar{D}_s$  and the second row showing the production of  $D \bar{D}_s^*$ . The vertices are labeled  $T_{21}$ ,  $T_{31}$ ,  $T_{11}$ ,  $T_{22}$ ,  $T_{32}$ ,  $T_{12}$ , and  $T_{32}$ .

—  $Z_{cs}(3982)$  is a pole in the coupled-channel  $\{\bar{D}_s D^*, D \bar{D}_s^*, D^* \bar{D}_s^*\}$  scattering amplitude

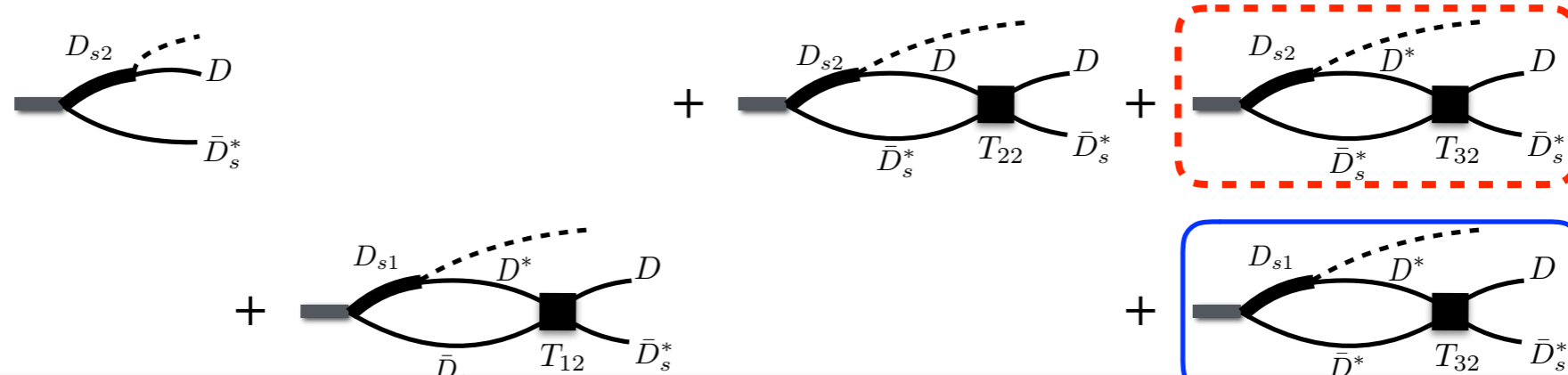
$$T_{\alpha\beta}(\sqrt{s}, p, p') = V_{\alpha\beta}(p, p') - \sum_{\gamma} \int \frac{d^3 q}{(2\pi)^3} V_{\alpha\gamma}(p, q) G_{\gamma}(\sqrt{s}, q) T_{\gamma\beta}(\sqrt{s}, q, p')$$

# Coupled-channel production $e^+e^- \rightarrow Y(4660) \rightarrow K^+(D_s^- D^{*0} + D_s^{*-} D^0)$

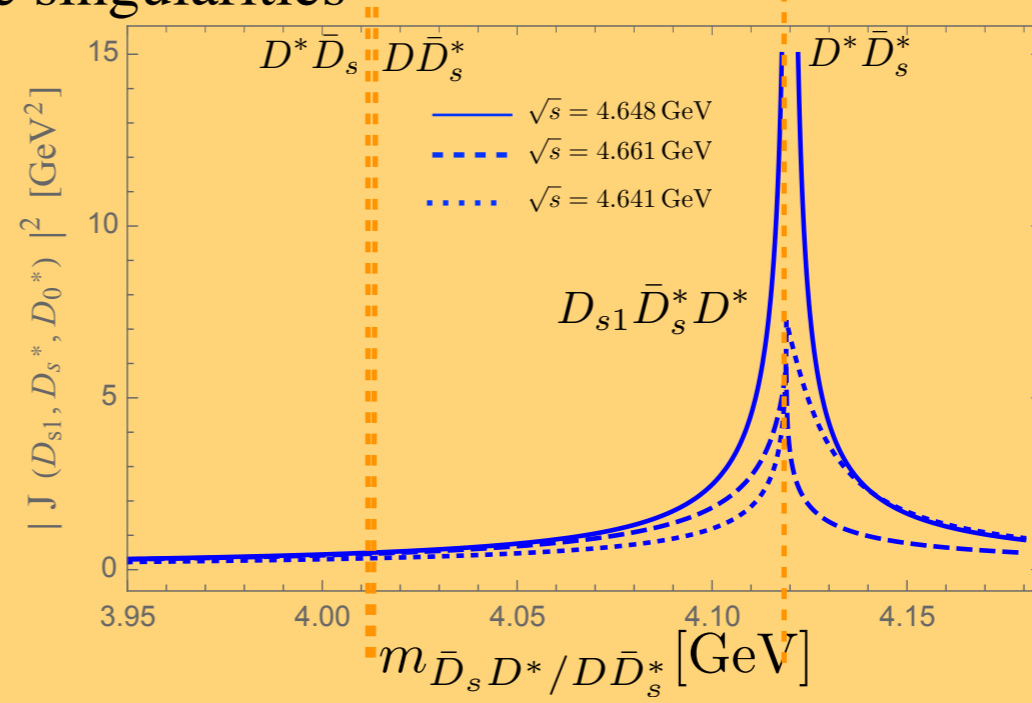
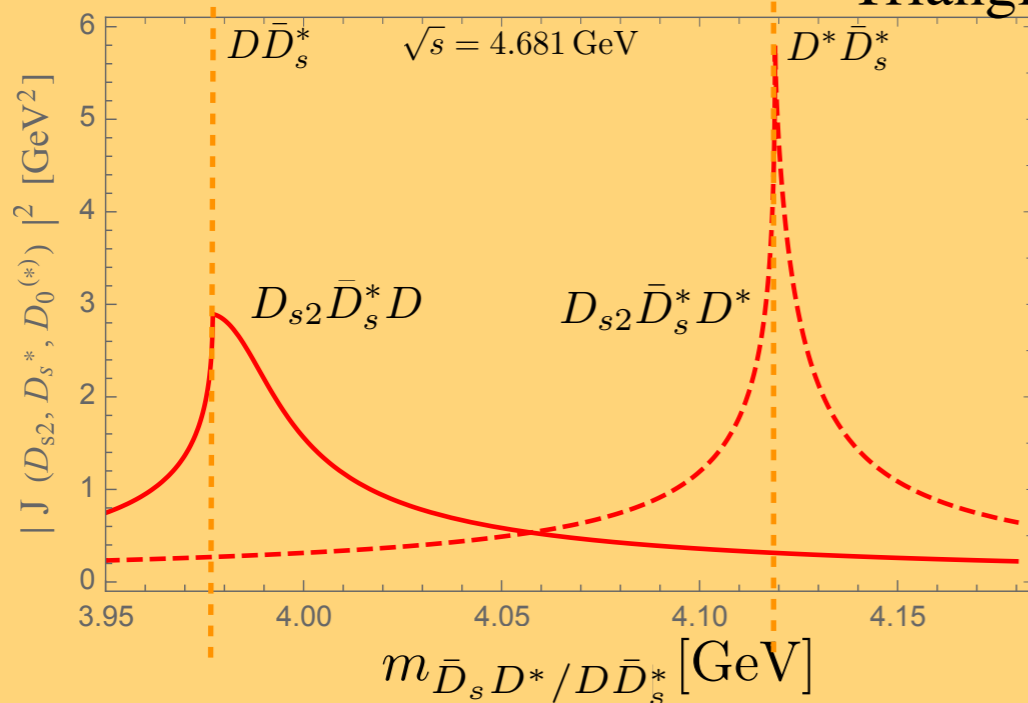
$$M_{Y \rightarrow KD^* \bar{D}_s} =$$



$$M_{Y \rightarrow KDD_s^*} =$$



## Triangle singularities



Amplitude

# From amplitudes to observables

- Differential cross section

$$\frac{d\sigma}{dm_{23}} = \mathcal{N} \left| \frac{2m_Y}{s - m_Y^2 + im_Y\Gamma_Y} \right|^2 \sum_{\alpha=1}^2 \frac{k q^{(23)}[\alpha]}{s} \int_{-1}^1 \frac{dz}{2} \left[ k^4 |M_D[\alpha]|^2 + |M_S[\alpha]|^2 \right] \quad \begin{array}{l} \alpha = 1 \equiv \bar{D}_s D^* \\ \alpha = 2 \equiv \bar{D}_s^* D \end{array}$$

- Number of events

$$\frac{dN}{dm_{23}} = \frac{d\sigma}{dm_{23}} \bar{\epsilon} \mathcal{L}_{\text{int}} f_{\text{corr}}$$

$\bar{\epsilon}$  — efficiency,

$\mathcal{L}_{\text{int}}$  — integrated luminosity,

$f_{\text{corr}}$  — radiative & vacuum polarisation

- Maximum likelihood

$$-2 \log \mathcal{L} = 2 \sum_i \left( \mu_i - n_i + n_i \log \frac{n_i}{\mu_i} \right)$$

- Combined fit of 5 distributions with 5 parameters:

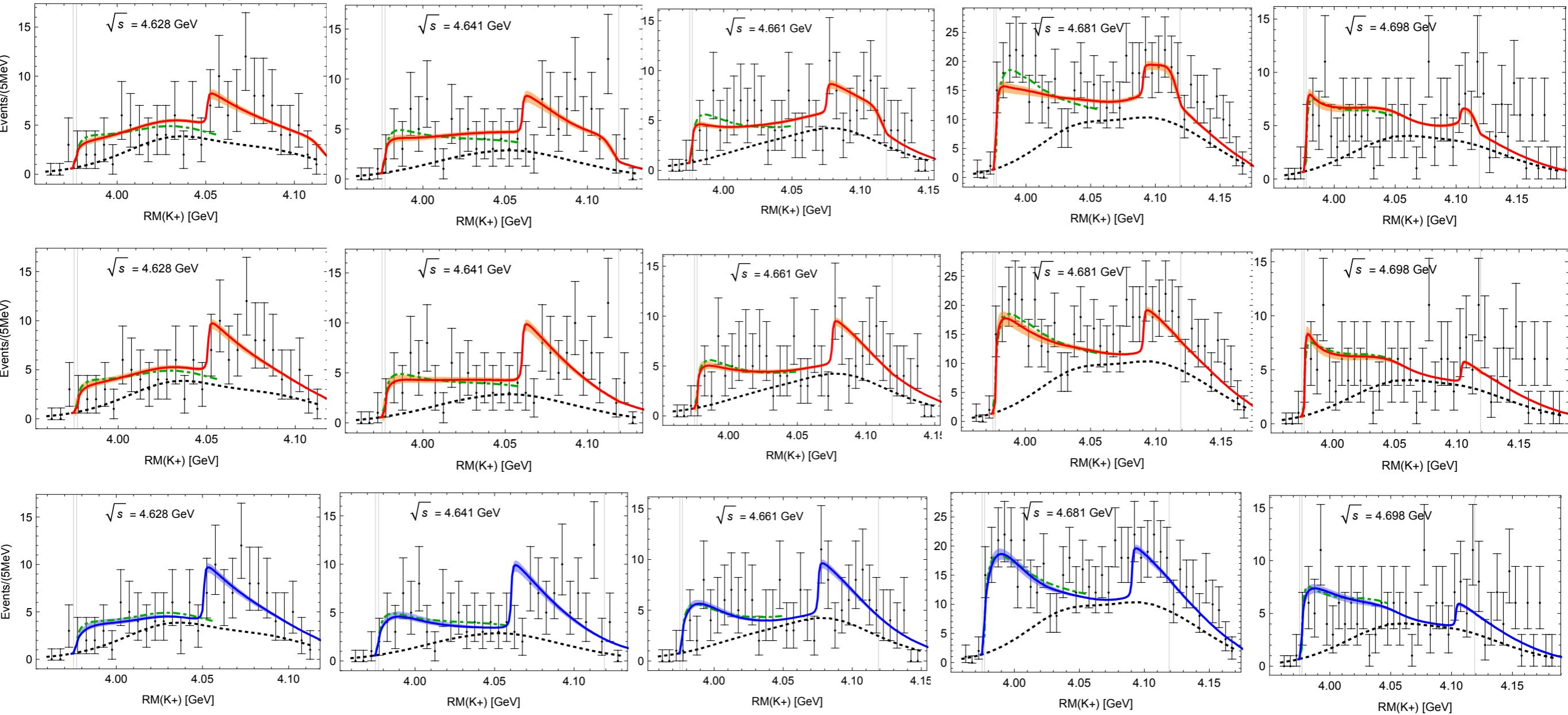
- 2 LO contact terms  $C_d$  and  $C_f$
- 2 ratios of production couplings
- Overall Norm

# Two classes of Solutions

Class 1: minor coupled-ch. effects—fits 1 and 1'

Class 2: large coupled-ch. effects—fit 2

similar to Yang et al. PRD103,074029(21)



Upper, middle, lower row for fit 1, fit 1', fit 2, respectively

5 parameter fits

Fit	$C_d, \text{fm}^2$	$C_f, \text{fm}^2$	$g_{D_{s1}}/g_{D_{s2}}$	$g/g_{D_{s2}}$	$\mathcal{N}, 10^{-2} \frac{\text{pb}}{\text{GeV}}$	$-2 \log \mathcal{L}$
fit 1	$-0.51 \pm 0.02$	$0.18 \pm 0.02$	$0.26 \pm 0.02$	$-2.5 \pm 0.3$	$0.46 \pm 0.05$	138
fit 1'	$-0.24 \pm 0.05$	$-0.1 \pm 0.05$	$0.37 \pm 0.03$	$-2.8 \pm 0.6$	$0.35 \pm 0.04$	144
fit 2	0.50 (fixed)	$-1.04 \pm 0.01$	$-0.44 \pm 0.03$	$-6.5 \pm 2.5$	$0.28 \pm 0.03$	146

← Global minimum

# Lessons from this study

## Pole positions

$J^{P(C)}$	State	Threshold, MeV	RS	Poles fit 1	RS	Poles fit 2
$1^+$	$Z_{cs}(3982)$	$\bar{D}_s D^* / \bar{D}_s^* D$ 3975.2/3977.0	(+ + +)	$3942 \pm 11$	(+ + +)	$3954 \pm 2$
$1^+$	$Z_{cs}(3982)$	$\bar{D}_s D^* / \bar{D}_s^* D$ 3975.2/3977.0	(- - +)	$3971 \pm 2$	(- - +)	$3959 \pm 7 - (47 \pm 16)i$
$1^+$	$Z'_{cs}$	$\bar{D}_s^* D^*$ 4119.1	(- - +)	$4115 \pm 2 - (10 \pm 2)i$		No state/not spin partner
$1^{+-}$	$Z_c(3900)$	$(D\bar{D}^*, -)$ 3871.7	(++)	$3841 \pm 11$	(-+)	$3864 \pm 7 - (58 \pm 13)i$
$1^{+-}$	$Z_c(4020)$	$\bar{D}^* D^*$ 4013.7	(-+)	$4009 \pm 18 - (9 \pm 2)i$		Not spin partner
$1^{++}$	$W_{c1}$	$(D\bar{D}^*, +)$ 3871.7	(-)	$3864 \pm 2$	(+)	$3852 \pm 2$
$2^{++}$	$W_{c2}$	$\bar{D}^* D^*$ 4013.7	(-)	$4009 \pm 2$	(+)	$3990 \pm 2$

fit 1 (1'): — bound (virtual) state, amplified by the  $D_{s2}D_s^*D$  triangle singularity

$Z_{cs}(3982)$

⇒ strong enhancement near the  $D_s^*D$  threshold

— SU(3) partner of  $Z_c(3900)$

— has a HQSS partner near  $D_s^*D^*$  threshold  $Z'_{cs}$  — SU(3) partner of  $Z_c(4020)$



# Lessons from this study

## Pole positions

$J^{P(C)}$	State	Threshold, MeV	RS	Poles fit 1	RS	Poles fit 2
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$1^{++}$	$W_{c1}$	$(D\bar{D}^*, +)$ 3871.7	(-)	$3864 \pm 2$	(+)	$3852 \pm 2$
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fit 1 (1'): — bound (virtual) state, amplified by the  $D_{s2}D_s^*D$  triangle singularity

$Z_{cs}(3982)$

$\implies$  strong enhancement near the  $D_s^*D$  threshold

— SU(3) partner of  $Z_c(3900)$

— has a HQSS partner near  $D_s^*D^*$  threshold  $Z'_{cs}$  — SU(3) partner of  $Z_c(4020)$

fit 2: — may be quite broad resonance state, amplified by the  $D_{s2}D_s^*D$  triangle singularity

$Z_{cs}(3982)$

$\implies$  strong enhancement near the  $D_s^*D$  threshold

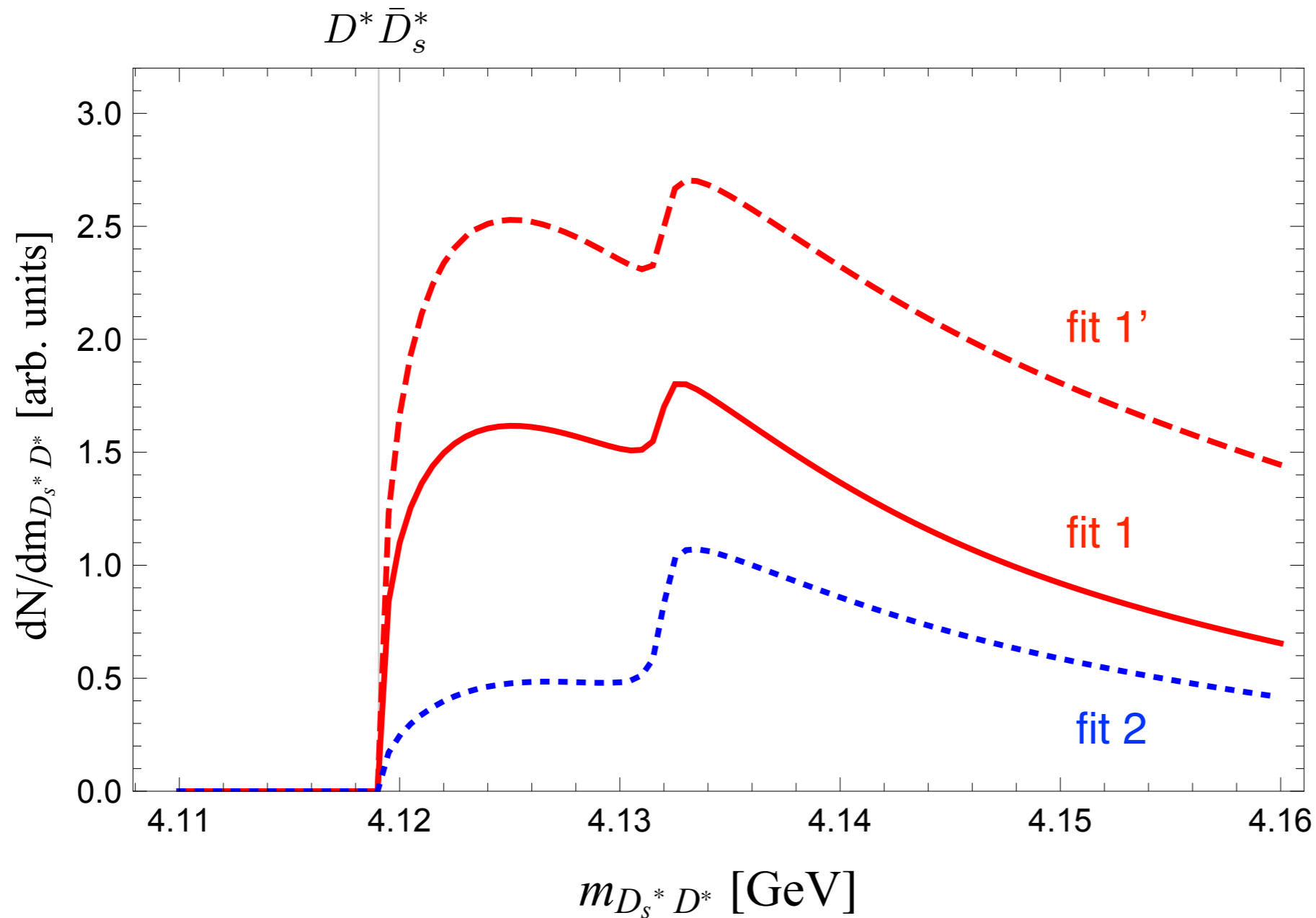
— may still be an SU(3) partner of  $Z_c(3900)$

— No HQSS partners near  $D_s^*D^*$  and  $D^*D^*$  thresholds

$\implies Z_c(4020)$  must have a different origin

# How to discriminate?

Predictions for  $e^+e^- \rightarrow K^+ D_s^{*-} D^{*0}$



- Strong enhancement near  $D_s^* D^*$  threshold in fits 1 (1'), because of  $Z_{cs}'$  pole
- No  $Z_{cs}'$  pole in fit 2  $\implies$  only smooth phase space
- Structure near 4.133 is due to  $D_{s1}$

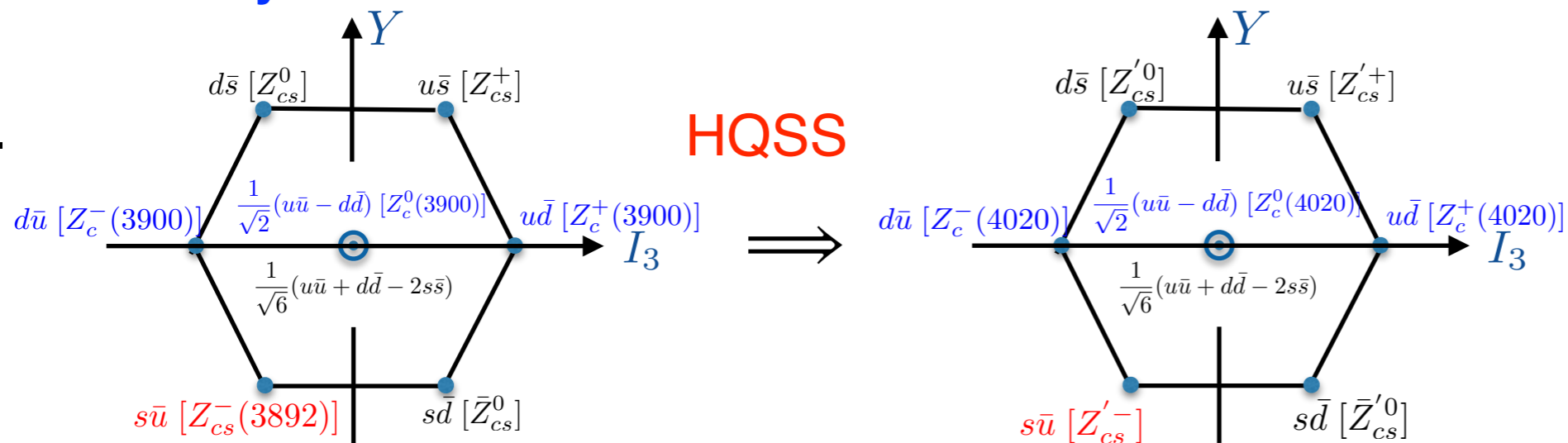
# Conclusions

• **BES III lineshapes**  $e^+e^- \rightarrow K^+(D_s^- D^{*0} + D_s^{*-} D^0)$  are analysed :

- an EFT approach with LO contact interactions
- **SU3 and HQSS**  $\Rightarrow$  multiplets of particles
- Various production mechanisms including triangle singularities

• **Two very different scenarios:**

1.



All states: bound/virtual

consistent with Yang et al.('21), Du et al.('22)

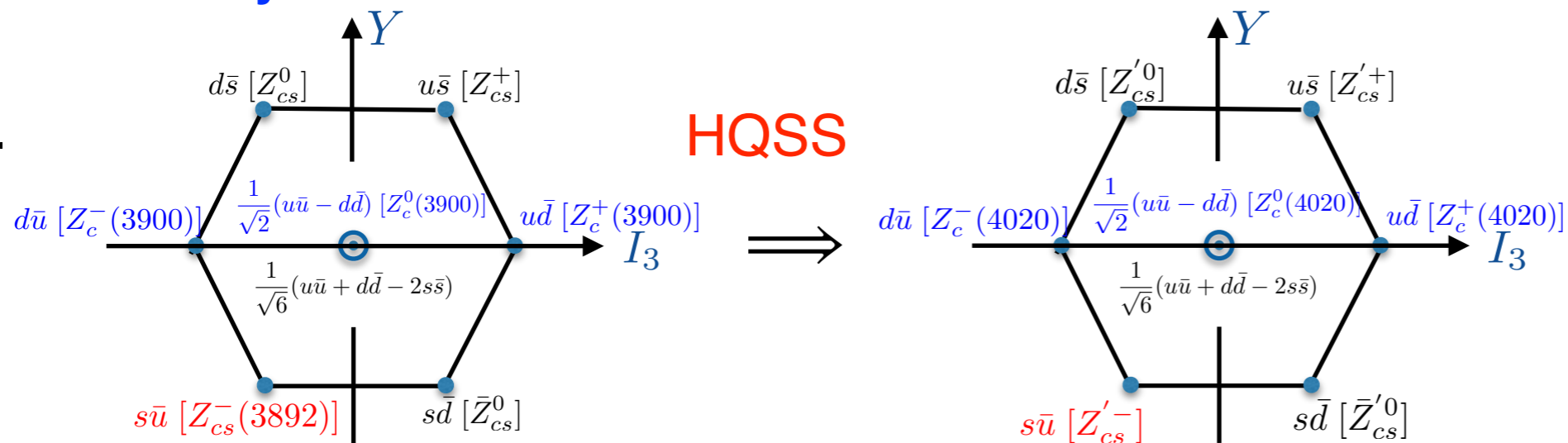
2.  $Z_{cs}(3982)$  — may be a strange partner of  $Z_c(3900)$ , if both are resonances but no HQSS partners

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2.  $Z_{cs}(3982)$  — may be a strange partner of  $Z_c(3900)$ , if both are resonances but no HQSS partners

$\Rightarrow$  Precise data in  $D_s D^* + D_s^* D$  and  $D_s^* D^*$  needed!

$\Rightarrow$  Role of pseudosclar GB octet to be understood

$\Rightarrow$  Possible connection of the  $Z_{cs}(3982)$  with the structures by LHCb

Ortega et al. ('21)  
Meng et al. (21')

Thanks for your attention!

