

FRAGMENTATION IN QUARKONIUM HADROPRODUCTION



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Geoffrey T. Bodwin, HSC, U-Rae Kim, Jungil Lee, PRL113, 022001 (2014)

Geoffrey T. Bodwin, HSC, U-Rae Kim, Jungil Lee, Yan-Qing Ma, Kuang-Ta Chao, PRD 93, 034041 (2016)

Production at Large p_T

- ▶ Large p_T ($p_T \gg m_H$) cross section of a hadron H is given in QCD factorization by

$$\frac{d\sigma_H}{dp_T^2} = \sum_{i=g,q,\bar{q}} \frac{d\sigma_i}{dp_T^2} \otimes D_{i \rightarrow H}(z, \mu) \quad (\text{LP: } \sim 1/p_T^4)$$

Leading-power fragmentation

$$+ \sum_n \frac{d\sigma_{Q\bar{Q}(n)}}{dp_T^2} \otimes D_{Q\bar{Q}(n) \rightarrow H}(z, \zeta_1, \zeta_2, \mu) \quad (\text{NLP: } \sim 1/p_T^6)$$

Next-to-leading-power fragmentation

$$+ O(1/p_T^8)$$

J.C.Collins and D.E.Soper, NPB194, 445 (1982)
 Z.-B. Kang, J.-W. Qiu, G. Sterman, PRL108, 102002 (2012)
 S. Fleming, A. K. Leibovich, T. Mehen, I. Z. Rothstein, PRD86, 094012 (2012)
 Y.-Q. Ma, J.-W. Qiu, G. Sterman, H. Zhang, PRL113, 142002 (2014)

Fragmentation Production at Large p_T

- ▶ Leading-power fragmentation

$$\frac{d\sigma_H}{dp_T^2} \Big|_{\text{LP}} = \sum_{i=g,q,\bar{q}} \frac{d\sigma_i}{dp_T^2}(z) \otimes D_{i \rightarrow H}(z, \mu)$$

Parton cross section Fragmentation function

The diagram illustrates the decomposition of the total fragmentation cross section. A large bracket on the right side of the equation groups the term $\frac{d\sigma_i}{dp_T^2}(z) \otimes D_{i \rightarrow H}(z, \mu)$. Two arrows point to this bracket: one from the text "Parton cross section" above it, and another from the text "Fragmentation function" below it.

- ▶ Compute parton cross sections in perturbative QCD

$$\sigma_{AB \rightarrow i} = \sum_{j,k=g,q,\bar{q}} \int_0^1 dx_1 dx_2 f_{j/A}(x_1) f_{k/B}(x_2) \hat{\sigma}_{jk \rightarrow i}$$

in collinear factorization

Available through NLO in α_s

The diagram shows the integral expression for the parton cross section. An upward arrow points from the text "Available through NLO in α_s " to the term $\hat{\sigma}_{jk \rightarrow i}$ in the integral, indicating that this term is available at next-to-leading order in the strong coupling constant α_s .

- ▶ Determine fragmentation functions by perturbatively matching to NRQCD

$$D_{i \rightarrow H}(z, \mu) = \sum_N d_{i \rightarrow Q\bar{Q}(N)}(z, \mu) \langle \mathcal{O}^H(N) \rangle$$

Analytically available through order α_s^2 ,
order α_s^3 numerically available for some channels

NLP Corrections

- If needed, compute power-suppressed NLP corrections by either matching to fixed-order calculations,

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$$\frac{d\sigma_H}{dp_T^2} \Big|_{\text{NLP}} = \frac{d\sigma_H}{dp_T^2} \Big|_{\text{fixed order}} - \frac{d\sigma_H}{dp_T^2} \Big|_{\text{LP}}$$

- or compute NLP fragmentation contributions

Y.-Q. Ma, J.-W. Qiu, G. Sterman, H. Zhang, PRL113, 142002 (2014)

$$\frac{d\sigma_H}{dp_T^2} = \sum_{i=g,q,\bar{q}} \frac{d\sigma_i}{dp_T^2} \otimes D_{i \rightarrow H}(z, \mu) \quad (\sim 1/p_T^4)$$

Leading-power fragmentation

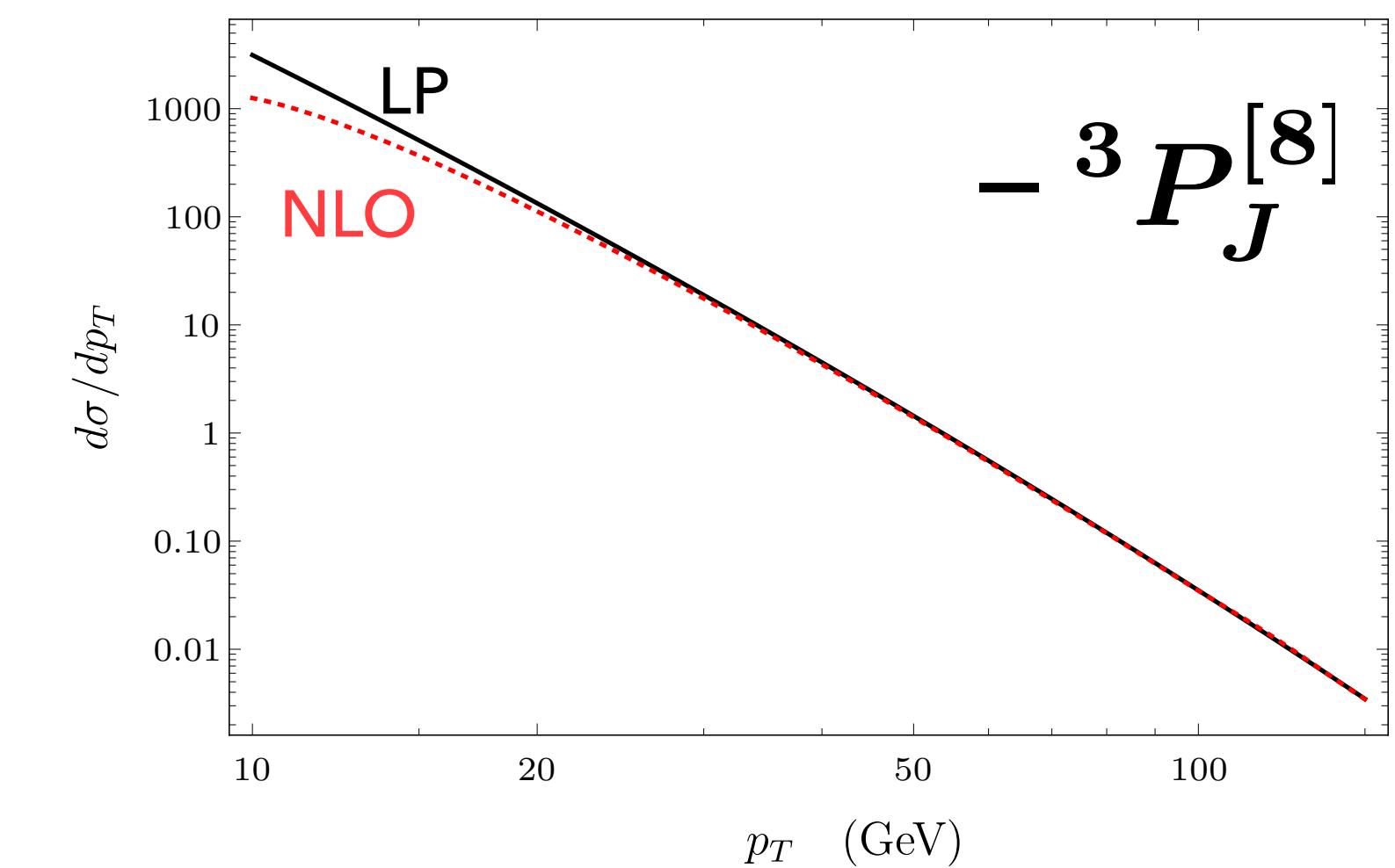
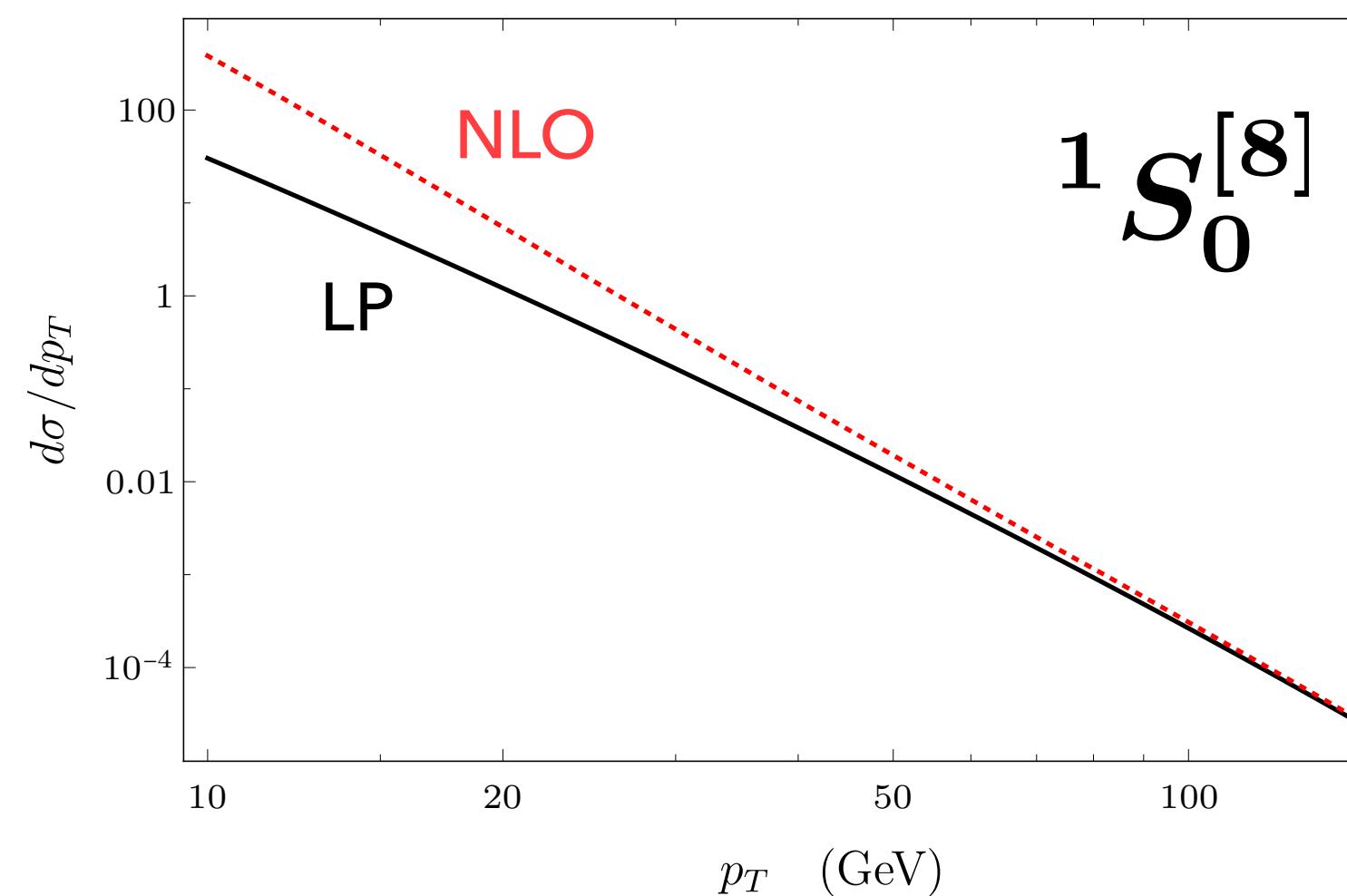
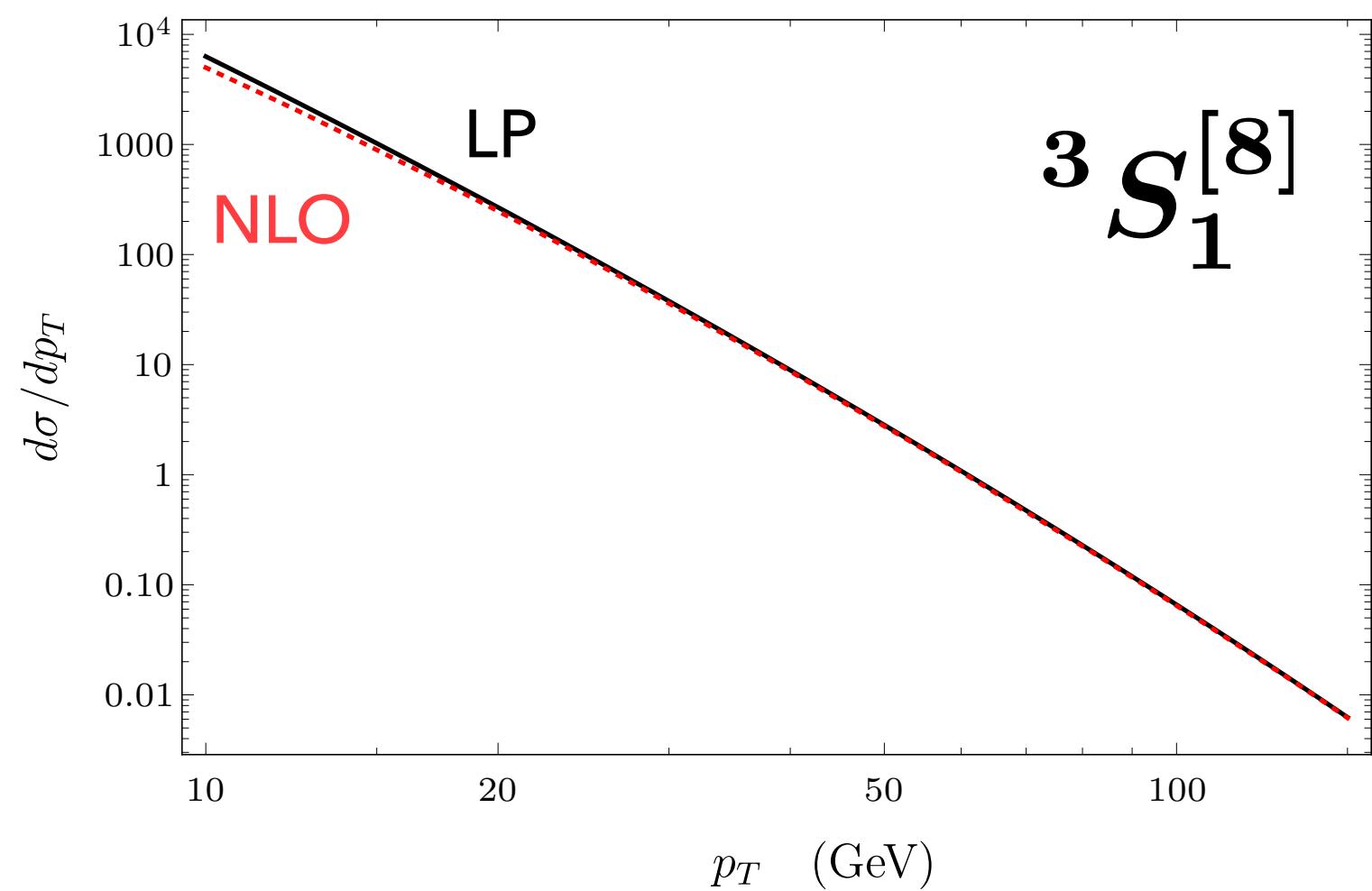
$$+ \sum_n \frac{d\sigma_{Q\bar{Q}(n)}}{dp_T^2} \otimes D_{Q\bar{Q}(n) \rightarrow H}(z, \zeta_1, \zeta_2, \mu) \quad (\sim 1/p_T^6)$$

Next-to-leading-power fragmentation

$$+ O(1/p_T^8)$$

Matching Fragmentation and Fixed-order Calculations

- ▶ Comparison of LP fragmentation and NLO fixed-order calculation



- ▶ Difference between LP and NLO gives the power-suppressed corrections

Resummation of DGLAP Logarithms

- ▶ Logarithms of p_T/m can be resummed by implementing DGLAP evolution.

Leading logarithm:

$$\frac{d}{d \log \mu_f^2} \begin{pmatrix} D_S(\mu_f) \\ D_g(\mu_f) \end{pmatrix} = \frac{\alpha_s(\mu_f)}{2\pi} \begin{pmatrix} P_{qq} & 2n_f P_{gq} \\ P_{qg} & P_{gg} \end{pmatrix} \otimes \begin{pmatrix} D_S(\mu_f) \\ D_g(\mu_f) \end{pmatrix}$$

- ▶ DGLAP evolution for fragmentation available through NLL accuracy
- ▶ NRQCD results for fragmentation functions singular at $z=1$, requires care.

Numerical procedure for distributions singular at $z=1$:

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 Bodwin, HSC, Kim, Lee, Ma, Chao, PRD 93, 034041 (2016)

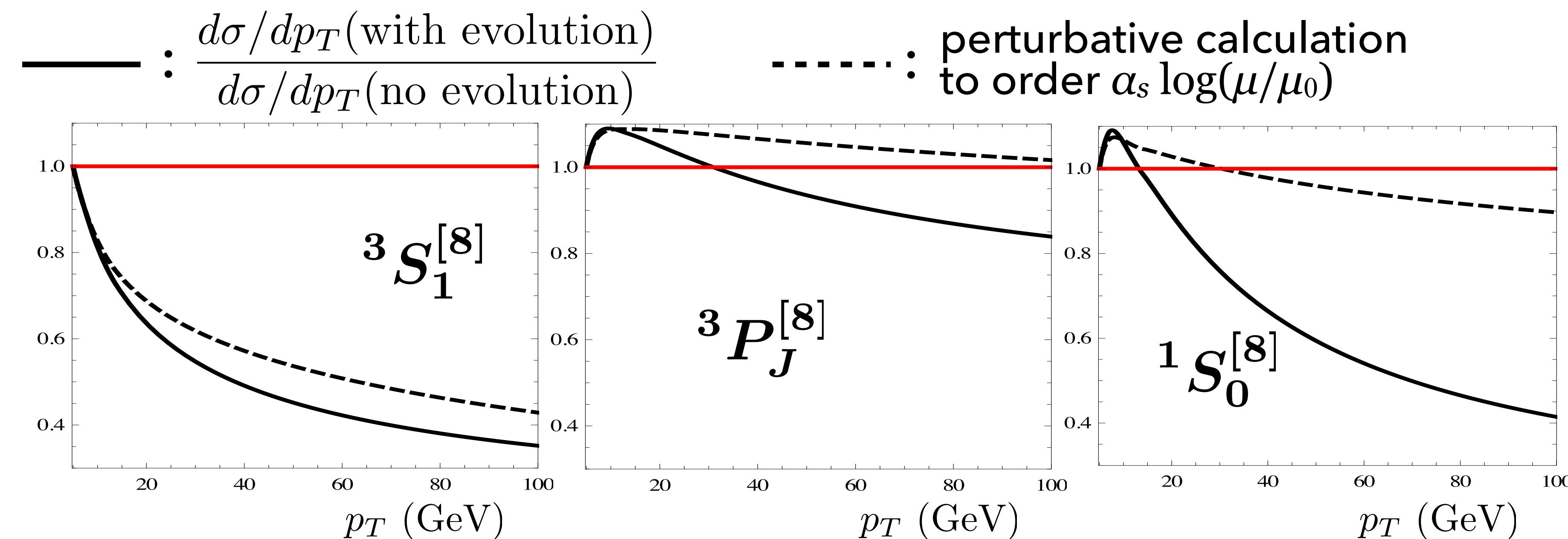
$$\begin{aligned} \int_0^1 dz \hat{\sigma}(z) D(z) &= \int_0^{1-\epsilon} dz \hat{\sigma}(z) D(z) + \int_{1-\epsilon}^1 dz \hat{\sigma}(z) D(z) \\ &\approx \int_0^{1-\epsilon} dz \hat{\sigma}(z) D(z) + \hat{\sigma}(z=1) \int_{1-\epsilon}^1 dz z^N D(z) \end{aligned}$$

$$\int_{1-\epsilon}^1 dz z^N D(z) = \int_0^1 dz z^N D(z) - \int_0^{1-\epsilon} dz z^N D(z)$$

↑ ↑
Well defined in Mellin space **Well behaved numerically**

Resummation of DGLAP Logarithms

- ▶ Modification of p_T -shape by evolution from $\mu_0=3$ GeV to $\mu=p_T$:
Resummation modifies the shape differently for each channel.



- ▶ When μ is not too large, numerical results agree well with perturbative calculation of DGLAP evolution to order $\alpha_s \log(\mu/\mu_0)$

Quarkonium Production in LP Fragmentation

- ▶ To obtain $\frac{d\sigma}{dp_T dy}$,
Computing time on a laptop:
- ▶ Compute NLO parton cross sections
Aversa, Chiappetta, Greco, Guillet, NPB 327, 105 (1989)
 ~5 min per p_T and y
- ▶ Compute DGLAP evolution of fragmentation functions (LL, NLL)
 ~5-10 min per p_T and y
 for 10 polarized &
 unpolarized channels
- ▶ Fortran and Mathematica versions
- ▶ PDFs can be swapped
- ▶ To change initial state, swap parton cross sections
e.g. use photoproduction parton cross sections
to compute photoproduction
Fontannaz, Guillet, Heinrich, EPJC 21, 303 (2001),
EPJC 22, 303 (2001), EPJC26, 209 (2002), EPJC34, 191 (2004)