

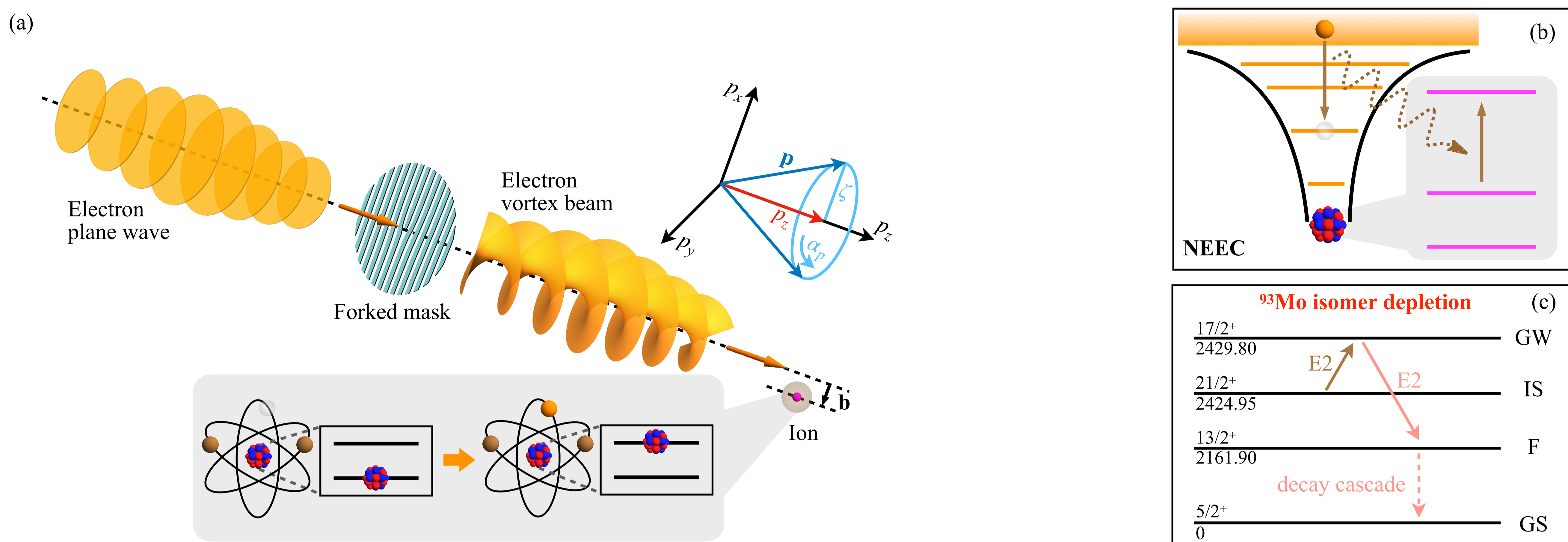
# Nuclear excitation by electron capture with electron vortex beams for isomer depletion

Yuanbin Wu<sup>1</sup>, Simone Gargiulo<sup>2</sup>, Fabrizio Carbone<sup>2</sup>, Christoph H. Keitel<sup>1</sup>, and Adriana Pálffy<sup>1,3</sup>

1. Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, D-69117 Heidelberg, Germany

2. Institute of Physics, Laboratory for Ultrafast Microscopy and Electron Scattering, École Polytechnique Fédérale de Lausanne, Station 6, Lausanne 1015, Switzerland

3. Department of Physics, Friedrich-Alexander-Universität Erlangen-Nürnberg, D-91058 Erlangen, Germany



## NEEC — fundamental process at the nuclear-atomic interface

First proposed theoretically in 1976

Phys. Lett. B **62**, 393 (1976)

First experimental observation claimed in 2018

Nature **554**, 216 (2018)

Population mechanisms of excited nuclear levels

Atomic vacancy effects on nuclear lifetime

Dense astrophysical plasmas

Isomer depletion

### NEEC with electron vortex beams

Key points for NEEC:

Vacancies of atomic levels

Free electrons

Shaping electron wave functions to manipulate nuclei?

#### Vortex beam

$$\psi_s(\mathbf{r}) = \int \frac{d^2\mathbf{p}_\perp}{(2\pi)^2} a_{\zeta m}(\mathbf{p}_\perp) u_{\mathbf{p}s} e^{i\mathbf{p}\cdot\mathbf{r}}, \quad a_{\zeta m}(\mathbf{p}_\perp) = (-i)^m e^{im\alpha_p} \delta(|\mathbf{p}_\perp| - \zeta) / \zeta$$

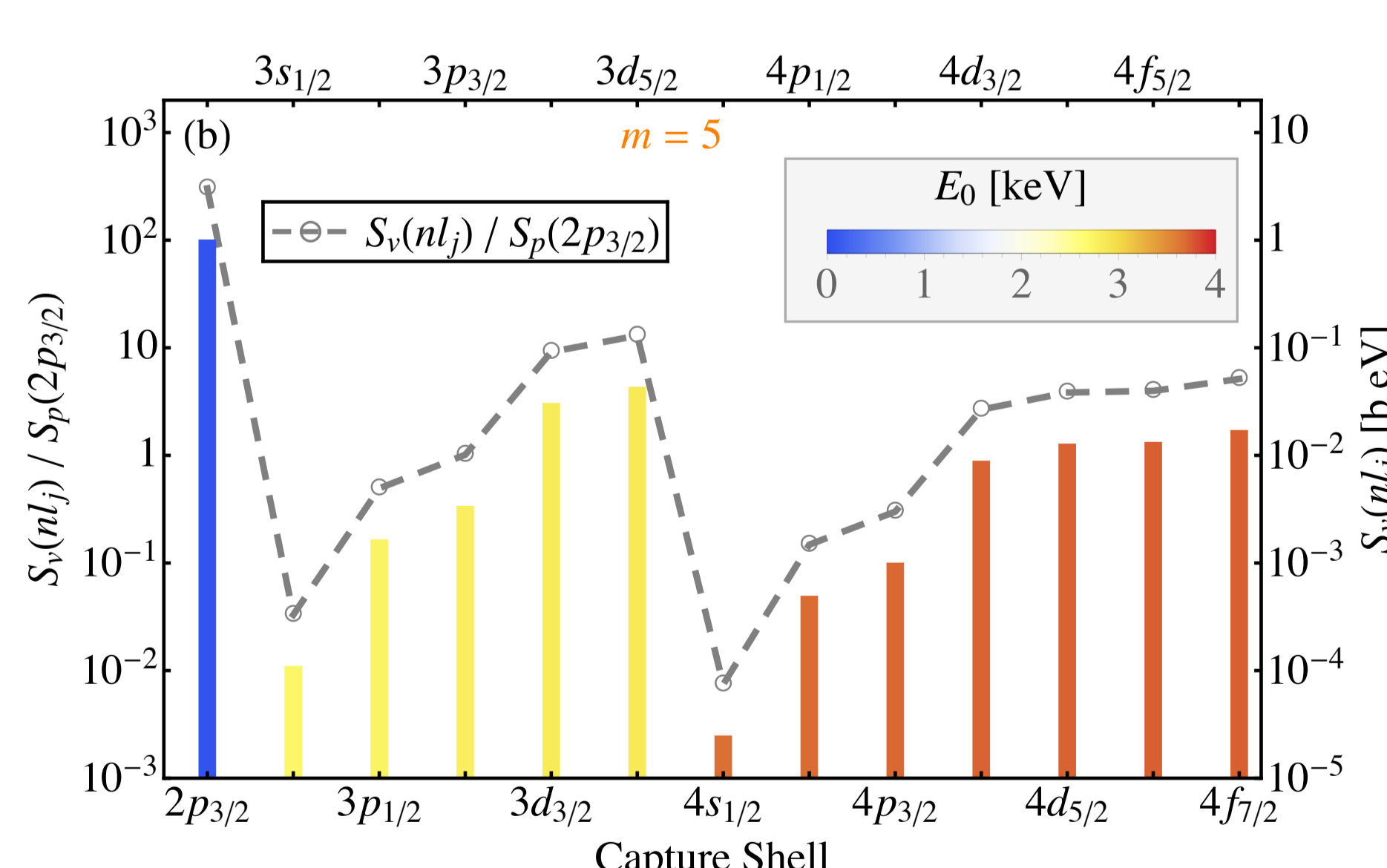
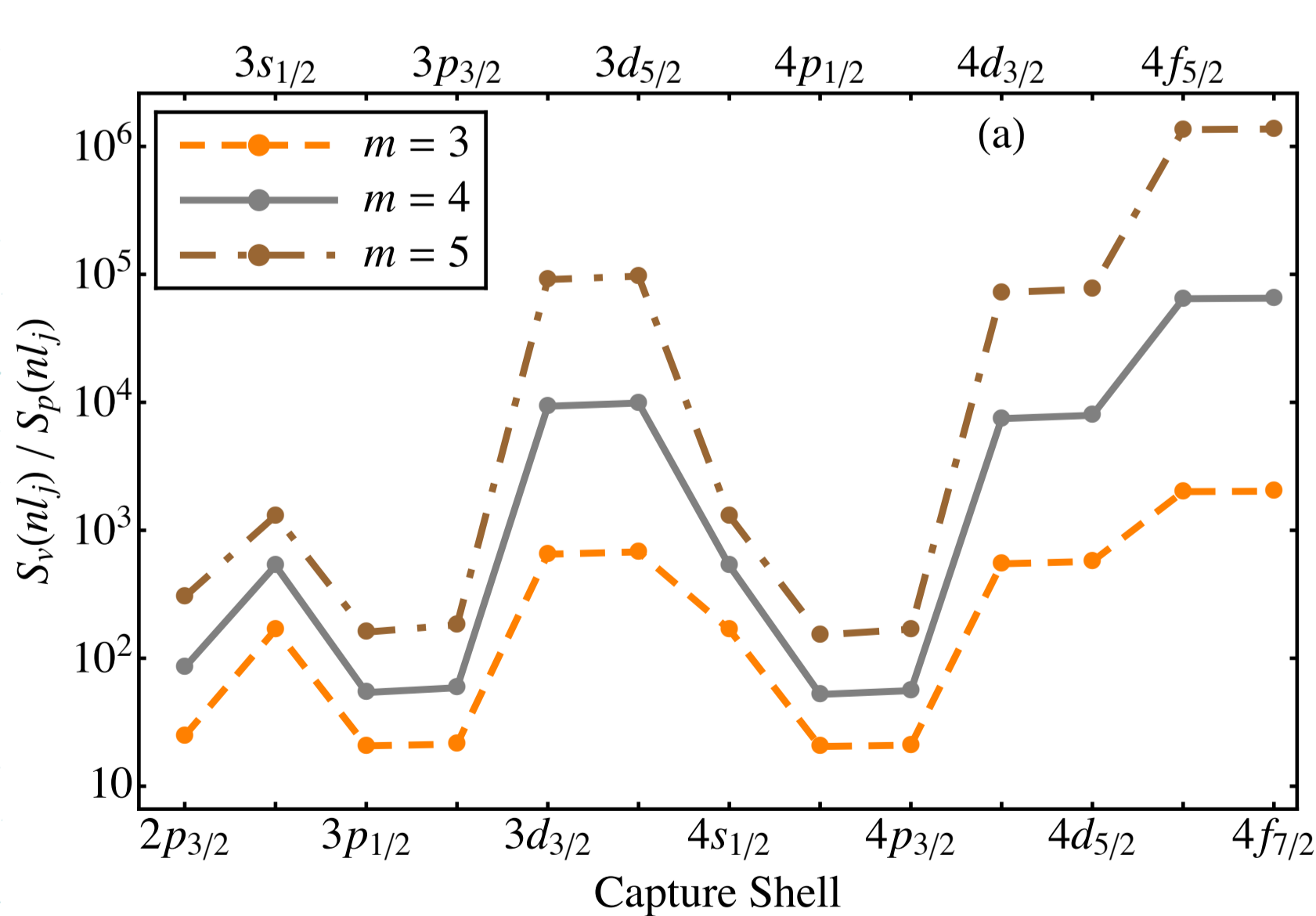
### Theoretical formalism

$$\sigma_{nec}^{i \rightarrow g}(E) = \frac{4\pi^2}{pJ_z} Y_{nec}^{i \rightarrow g} \mathcal{L}(E - E_0), \quad Y_{nec}^{i \rightarrow g} \propto |\langle \Psi_g^N | \langle \Psi_g^e | H_N | \Psi_i^e, \psi_s \rangle | \Psi_i^N \rangle|^2$$

$$Y_{nec}^{i \rightarrow g} = \frac{b^2}{4\pi} \int_0^{2\pi} \int_0^{2\pi} \frac{d\alpha_p}{2\pi} \frac{d\alpha_k}{2\pi} e^{im(\alpha_p - \alpha_k)} \mathcal{Y}_{nec}^{i \rightarrow g}(\mathbf{p}, \mathbf{k}) F_1(2; u)$$

$$\mathcal{Y}_{nec}^{i \rightarrow g}(\mathbf{p}, \mathbf{k}) = \frac{16\pi^3(2J_g + 1)}{(2J_i + 1)(2L + 1)^2} \mathcal{B} \uparrow(\lambda L) \rho_i \sum_{\kappa, m_i} \frac{\mathcal{Y}_b}{2l + 1} Y_{lm_i}^*(\theta_k, \varphi_k) Y_{lm_i}(\theta_p, \varphi_p)$$

### <sup>93m</sup>Mo isomer depletion E2 transition



### <sup>152m</sup>Eu isomer depletion M1 transition

<i>nlj</i>	<i>E<sub>d</sub></i> [keV]	<i>S<sub>p</sub></i> [b eV]	<i>S<sub>v</sub></i> [b eV] <i>m</i> = 3	<i>S<sub>v</sub></i> [b eV] <i>m</i> = 5
2s <sub>1/2</sub>	5.20	8.05 × 10 <sup>-4</sup>	1.14 × 10 <sup>-3</sup>	1.14 × 10 <sup>-3</sup>
2p <sub>1/2</sub>	5.19	7.85 × 10 <sup>-5</sup>	1.35 × 10 <sup>-3</sup>	3.34 × 10 <sup>-3</sup>
2p <sub>3/2</sub>	6.02	1.25 × 10 <sup>-5</sup>	4.21 × 10 <sup>-4</sup>	7.61 × 10 <sup>-3</sup>

The choice of impact parameter **b** is crucial

$$\zeta = p_z; \quad \zeta b = 1$$

- Introduce the theory for NEEC with an electron vortex beam
- 2 orders of magnitude enhancement for NEEC cross section for <sup>93m</sup>Mo isomer depletion

- 6 orders of magnitude enhancement for higher shells
- Control nuclear excitations by shaping electron wave functions

#### Reference:

Y. Wu, S. Gargiulo, F. Carbone, C. H. Keitel, and A. Pálffy, *Phys. Rev. Lett.* **128**, 162501 (2022).

