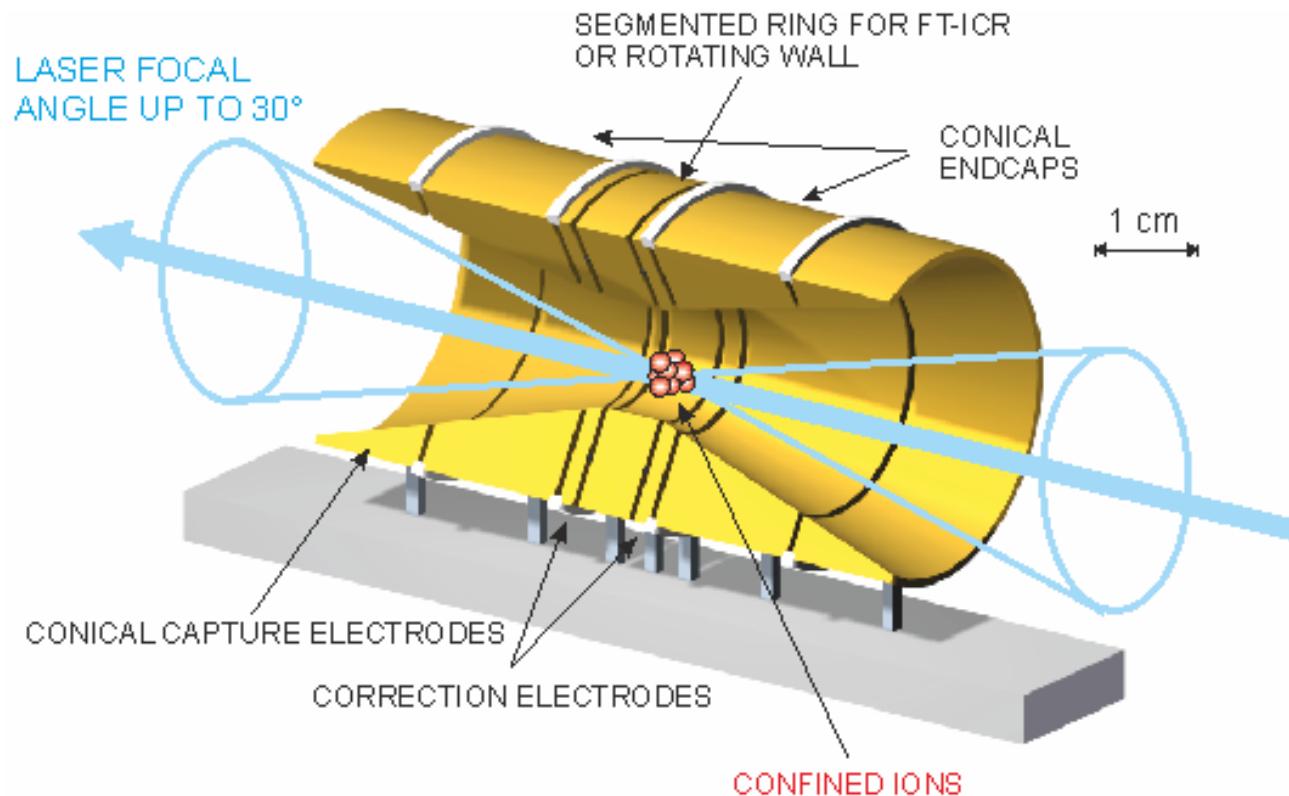


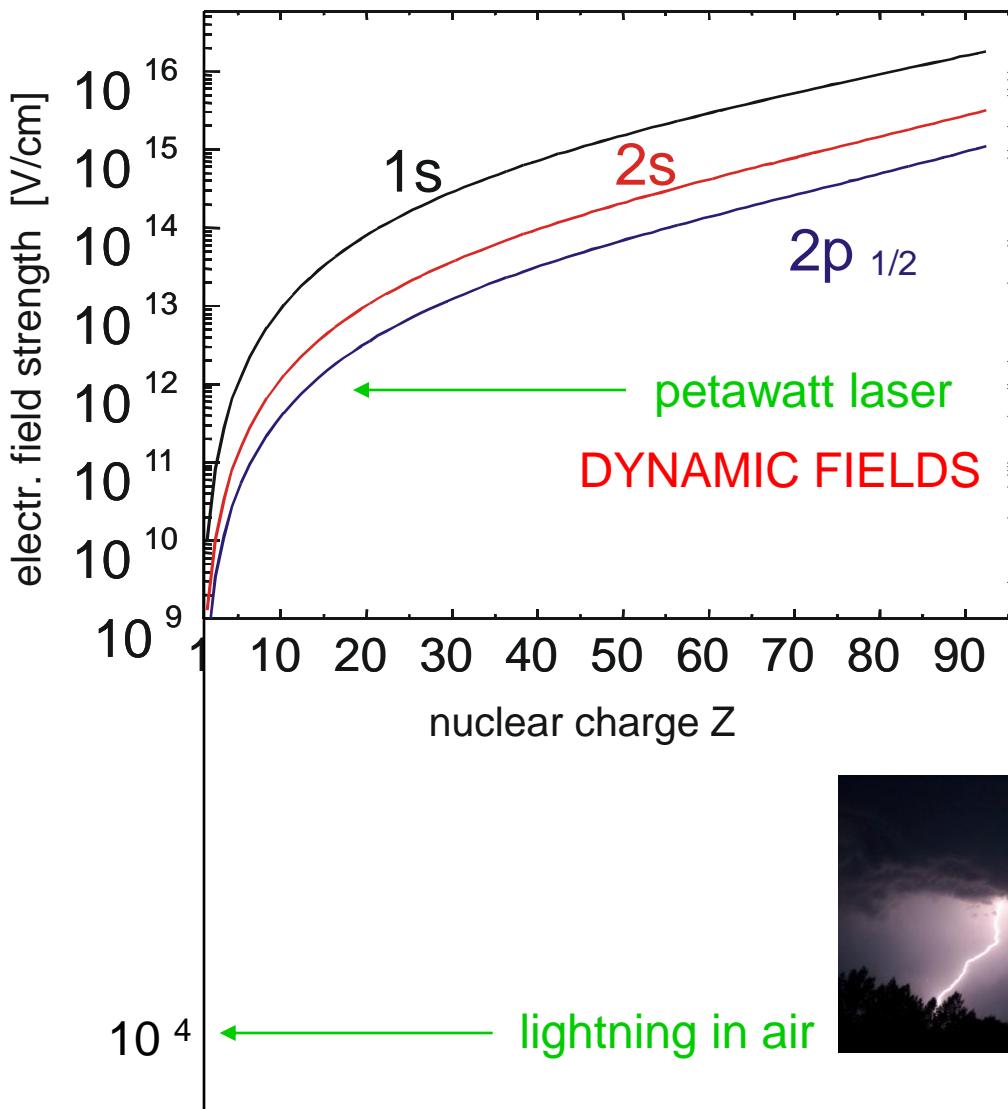
# TRAP-ASSISTED STUDIES @ A HIGH-INTENSITY LASER



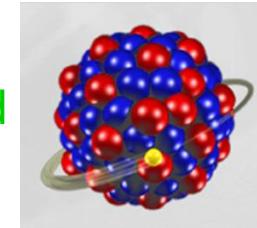
seit 1558

Helmholtz-Institut Jena  
Schiller-Universität Jena

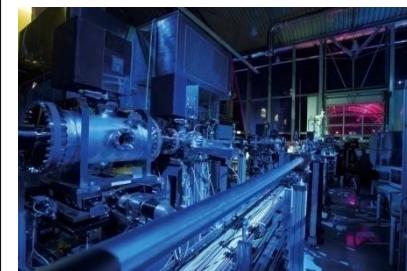
# extreme fields



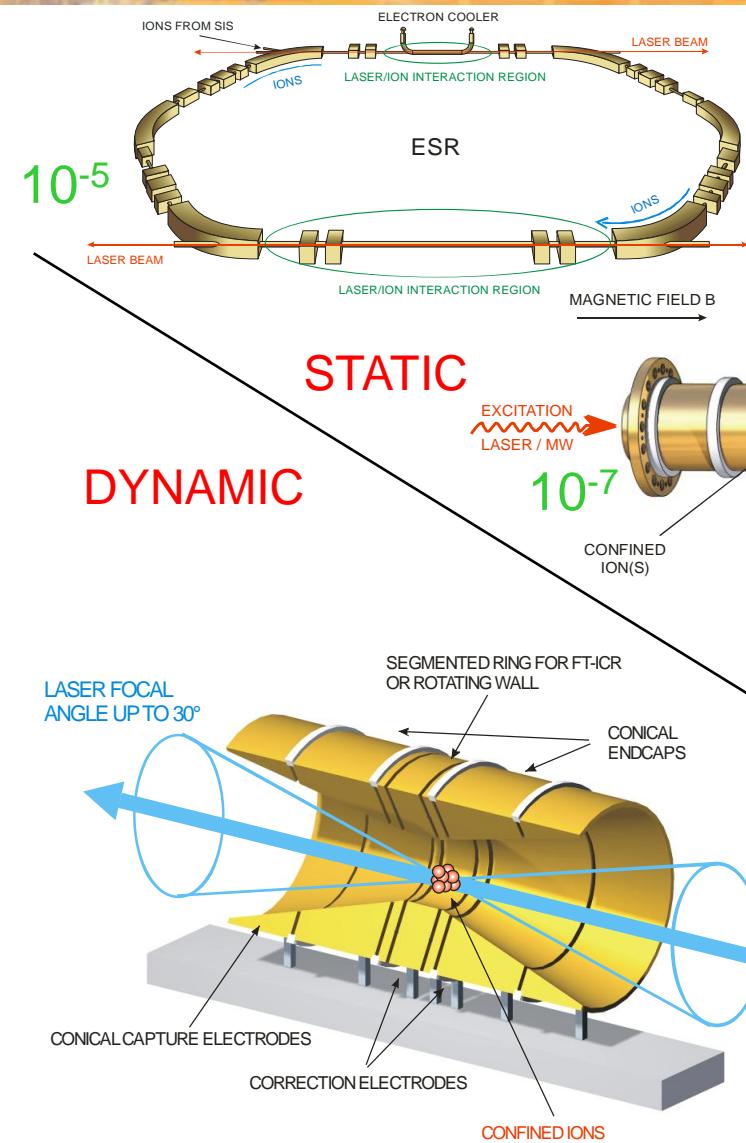
highly charged ions



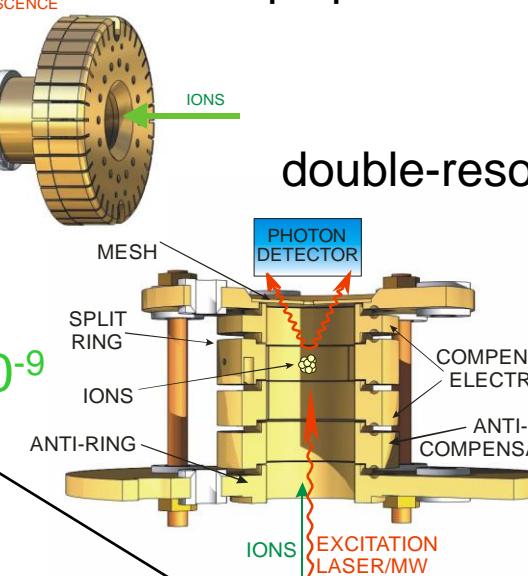
STATIC FIELDS



# our extreme field physics in traps



in-ring experiments (e.g. LIBELLE)  
(or EBIT)



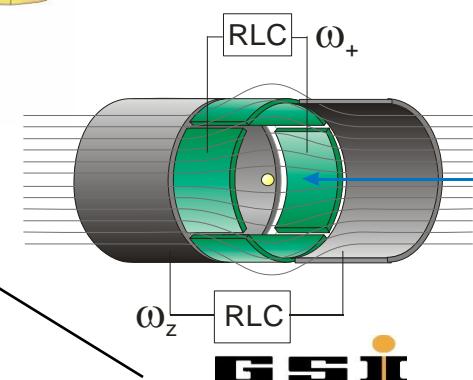
in-trap spectroscopy (SPECTRAP)

double-resonance spectroscopy  
(g-factor)

PHELIX  
POLARIS  
FLASH  
JETI

$10^{-9}$

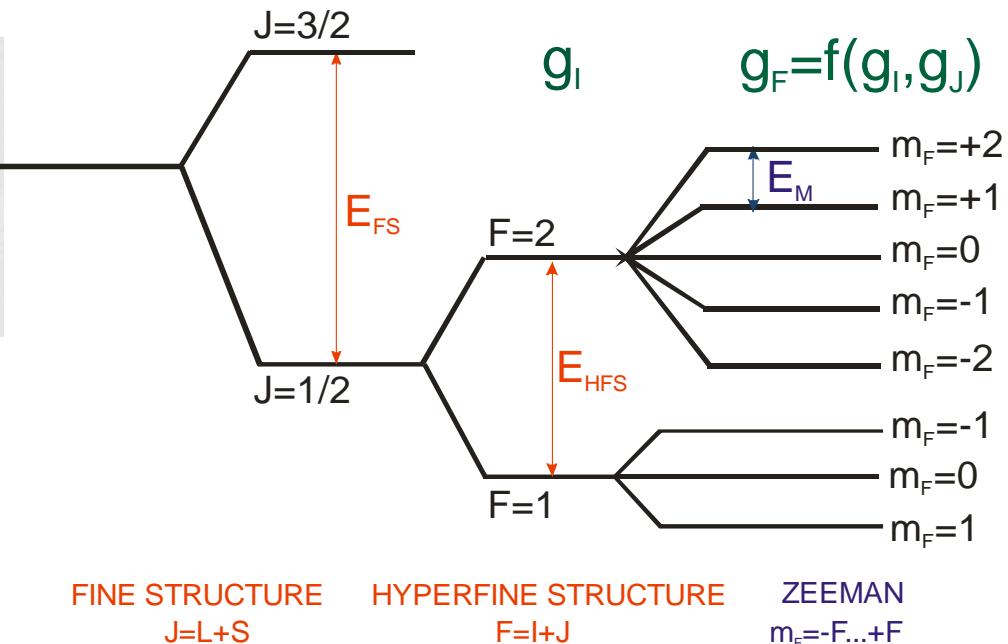
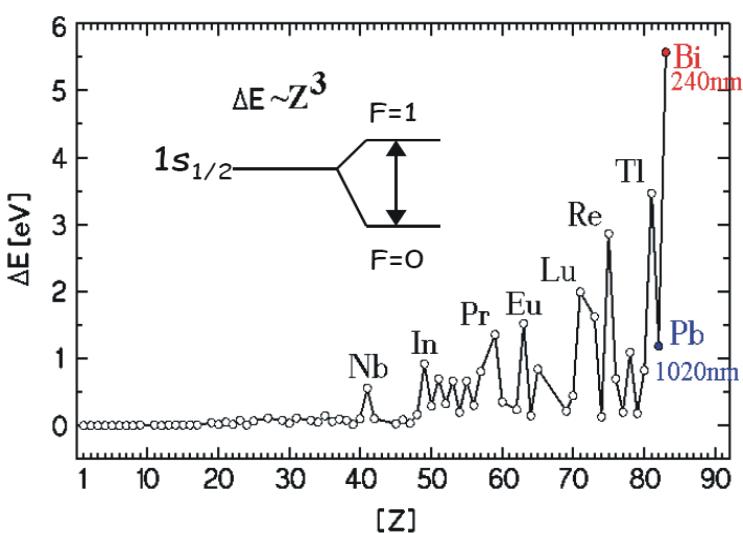
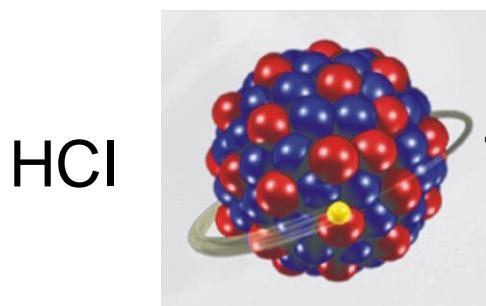
blind spectroscopy



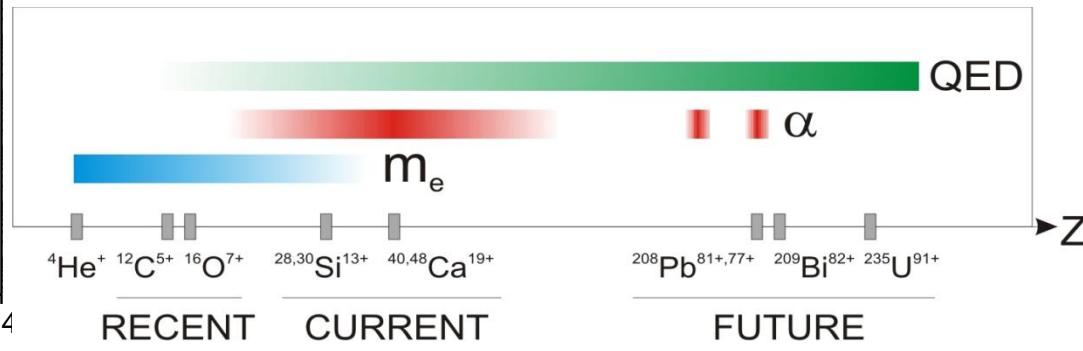
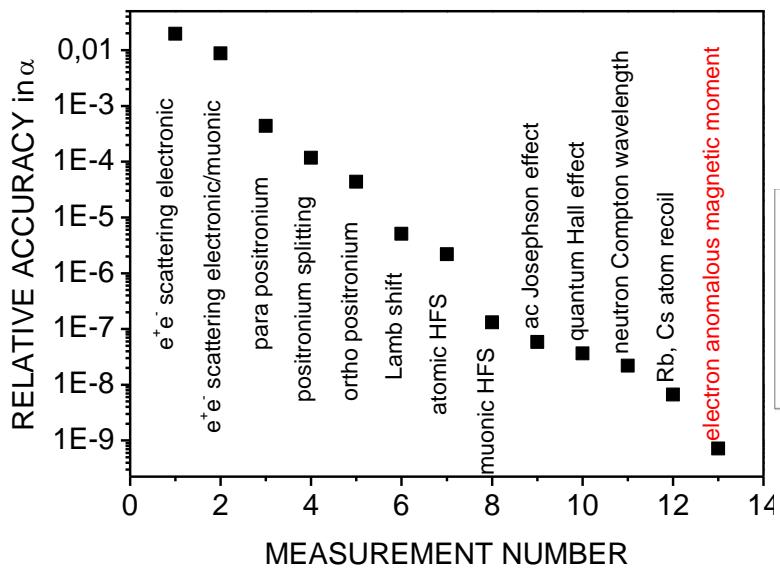
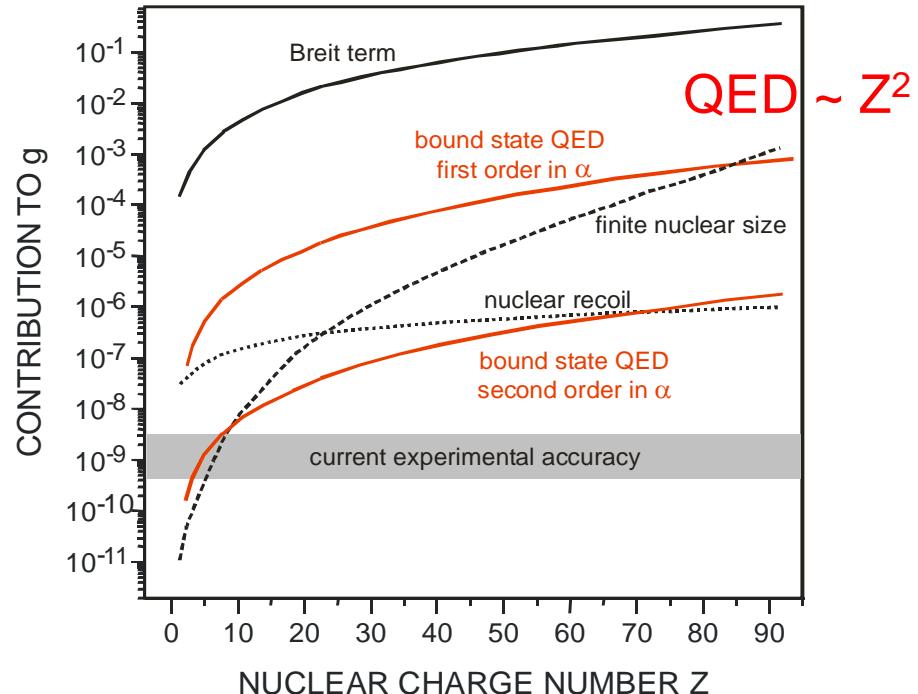
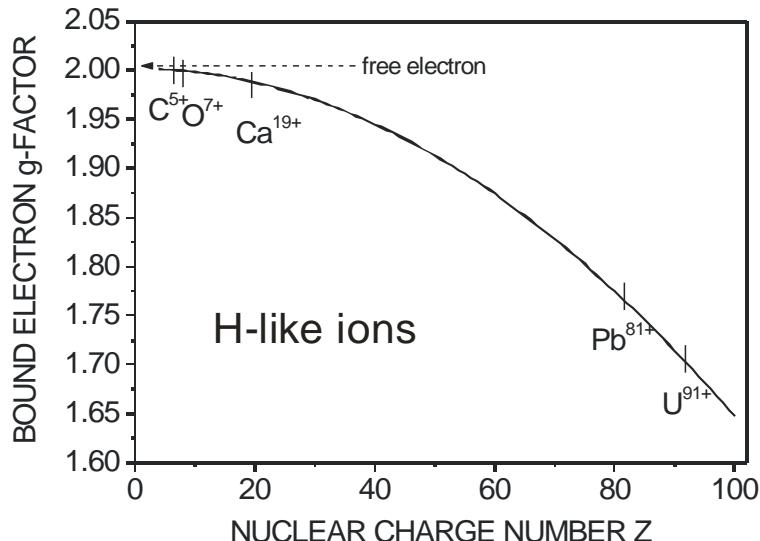
# precision spectroscopy of forbidden transitions in HCl (optical laser / MW)

- Transition energy (frequency)
- Fine-structure (multiplet) splitting
- Hyperfine splitting
- Lamb shift

$$\begin{aligned} Z^2 \\ \alpha^2 Z^4 \\ (m_e/m_p)\alpha^2 Z^3 \\ \alpha^5 Z^4 \end{aligned}$$

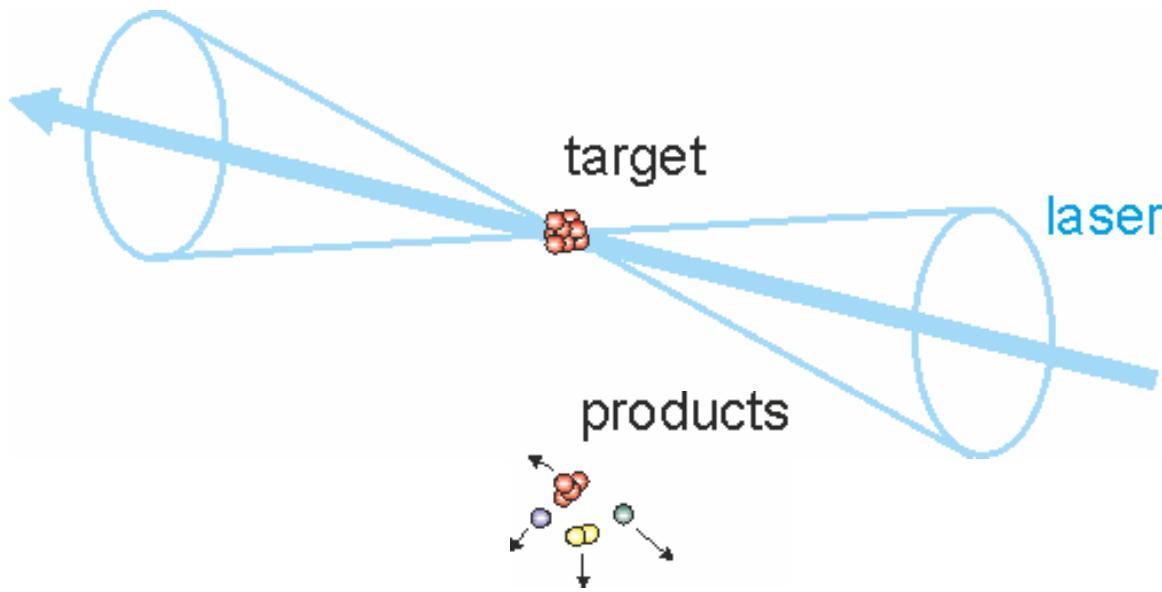


# background: QED and fundamental constants



# laser-induced reaction / analysis

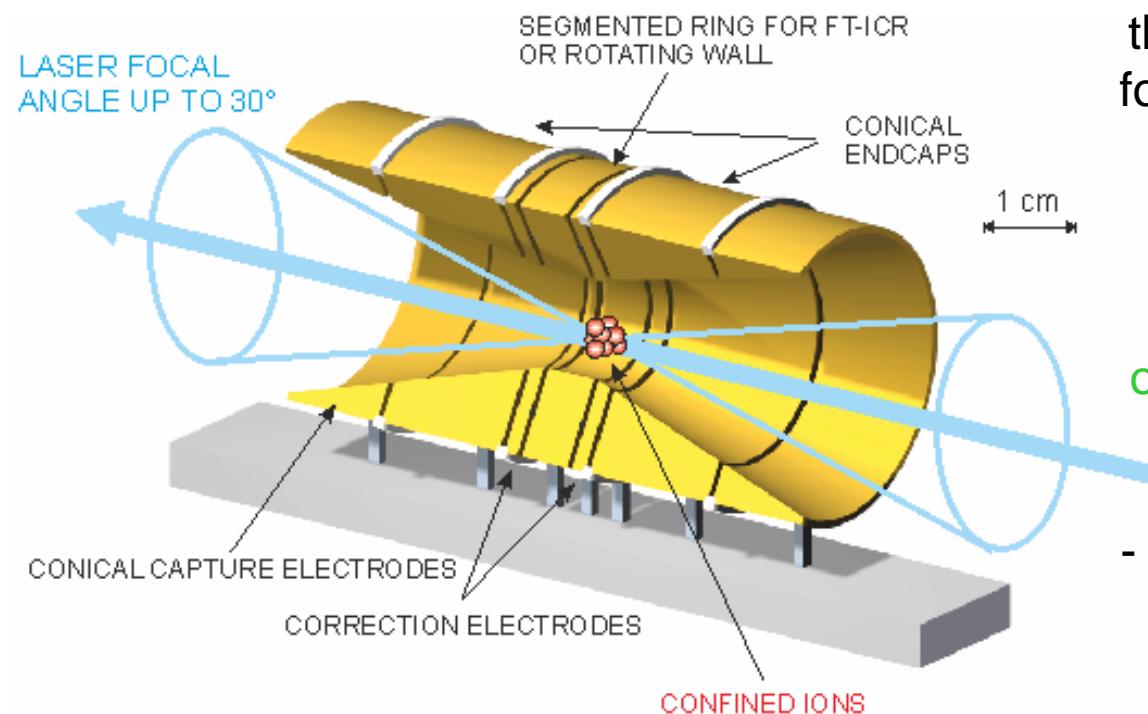
applications may be different but tools can be the same



typical problems: what exactly was the target?  
what exactly are the products?

for charged particles both can be controlled by use of a Penning trap

# use of a Penning trap



non-linear photoionisation  
giant resonances  
strong-field effects

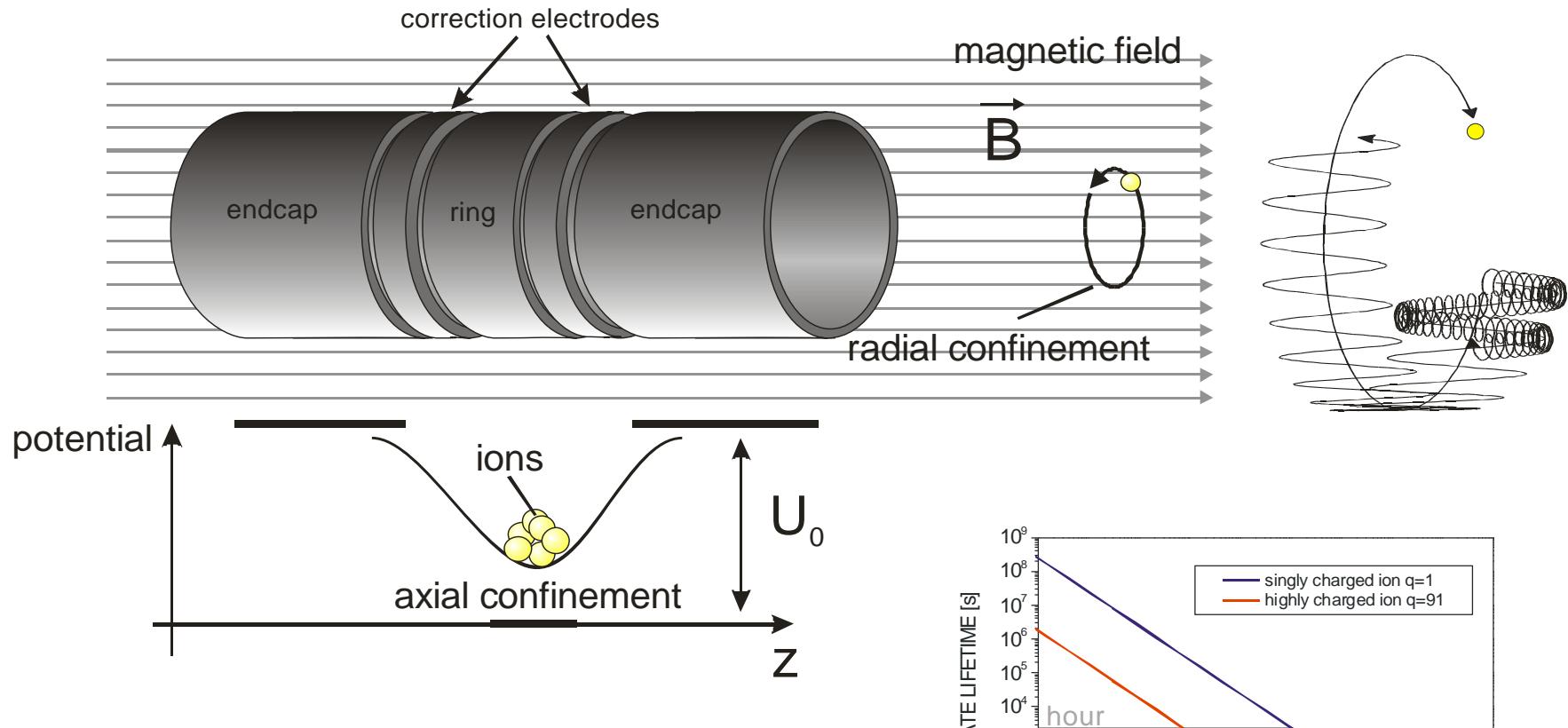
„transparent“ experiment

the Penning trap as a **universal tool** for the study of reactions induced by a high-intensity laser

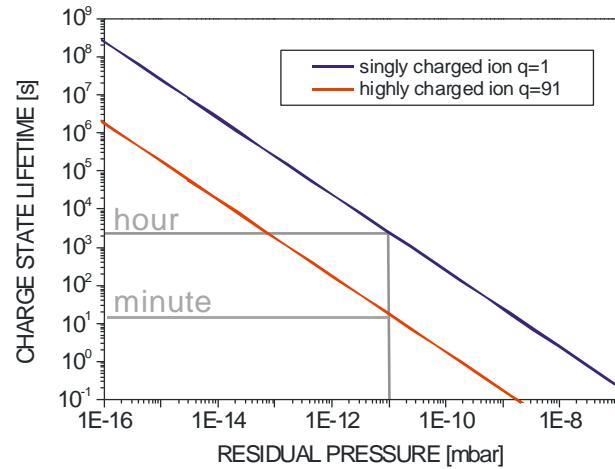
main motivation: non-destructive real-time analysis of multiphoton ionization processes with high resolution

- good ion localization and extended storage times (minutes)
  - well-controlled and clean environment for various reactions
- numerous manipulation techniques
  - non-destructive detection

# Penning trap principles



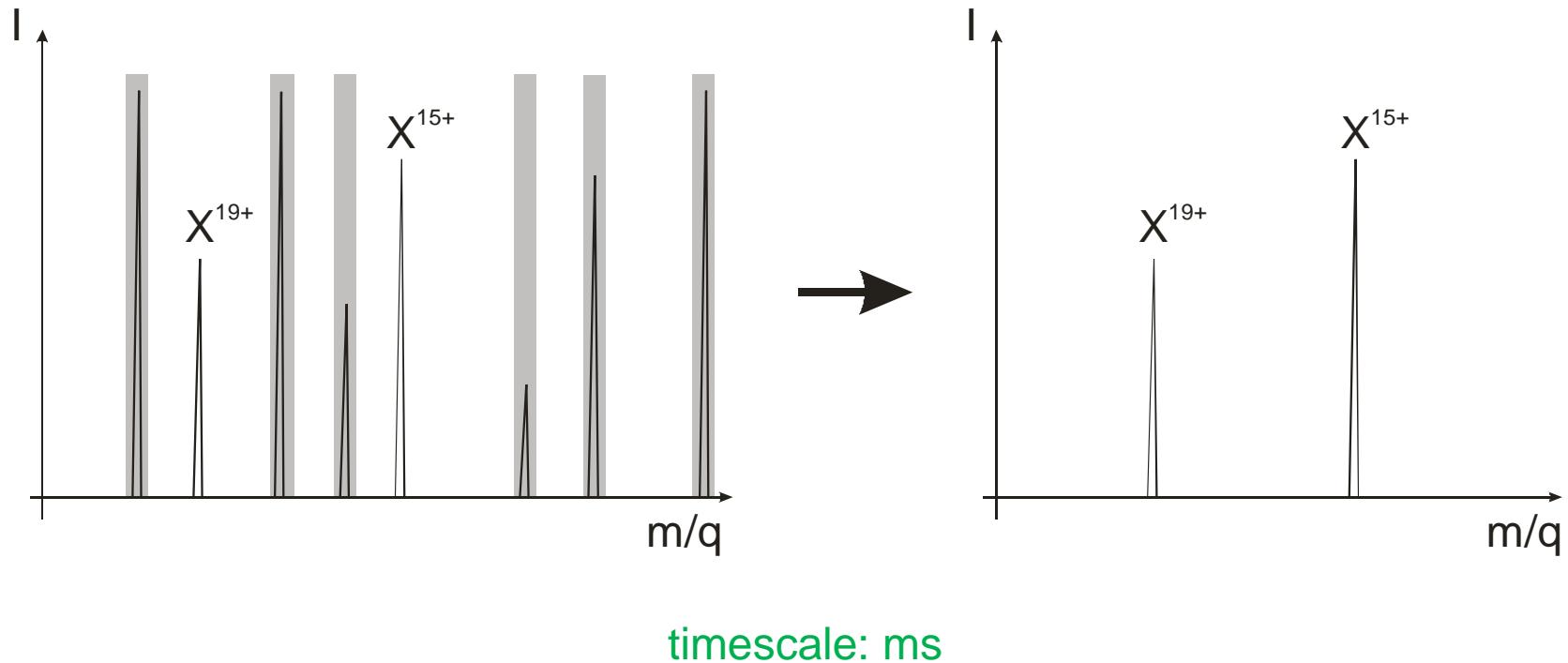
typical parameters:  $B = 6 \text{ T}$   
 $U_0 = 100 \text{ V}$   
 $d = 20 \text{ mm}$   
 $\omega_-, \omega_z, \omega_+ = 10 \text{ kHz}, 100 \text{ kHz}, 10 \text{ MHz}$



# techniques (1) : ion selection

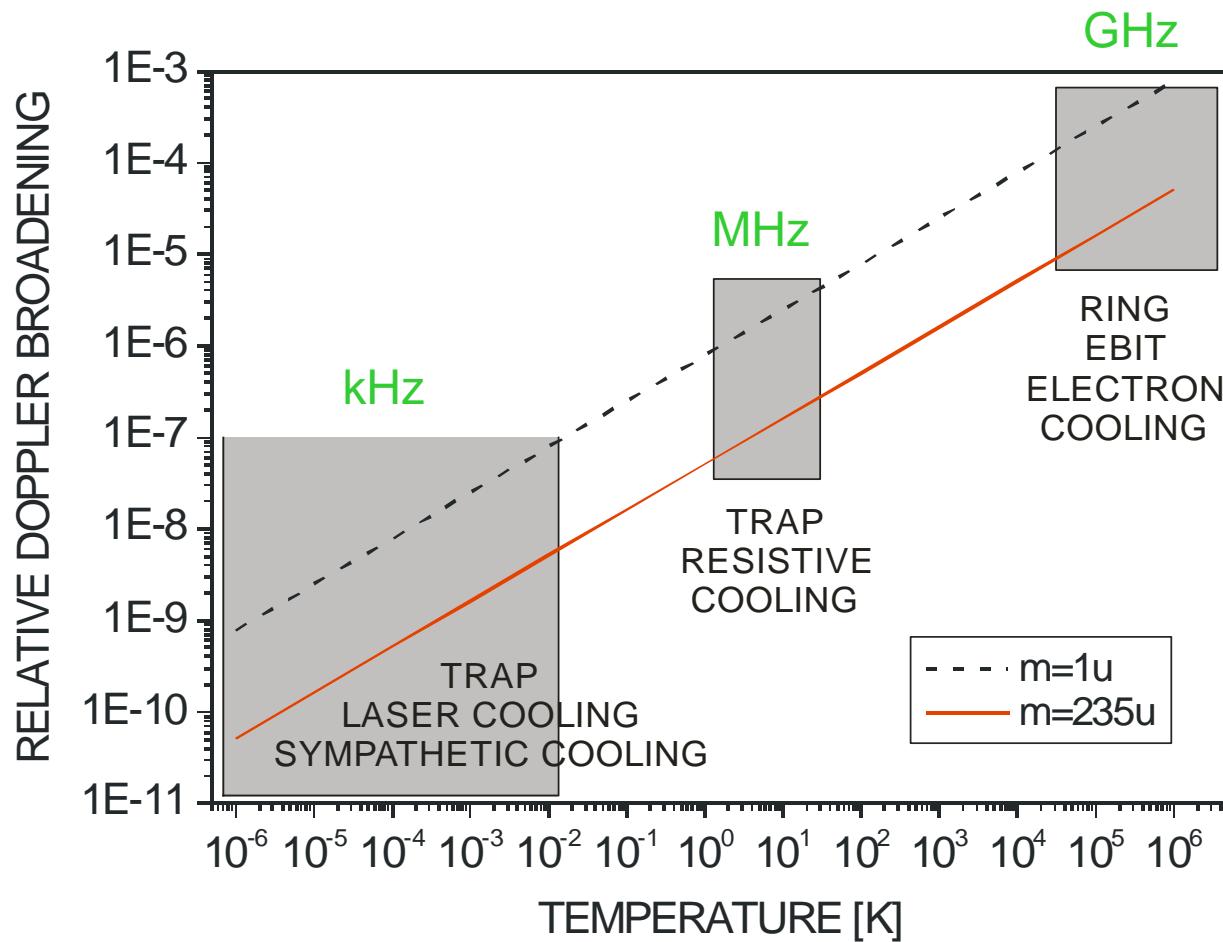
selection of **any** mass to charge ratio **combination**  
by resonant **ejection of all unwanted ions**

**SWIFT** (stored waveform inverse Fourier transform)

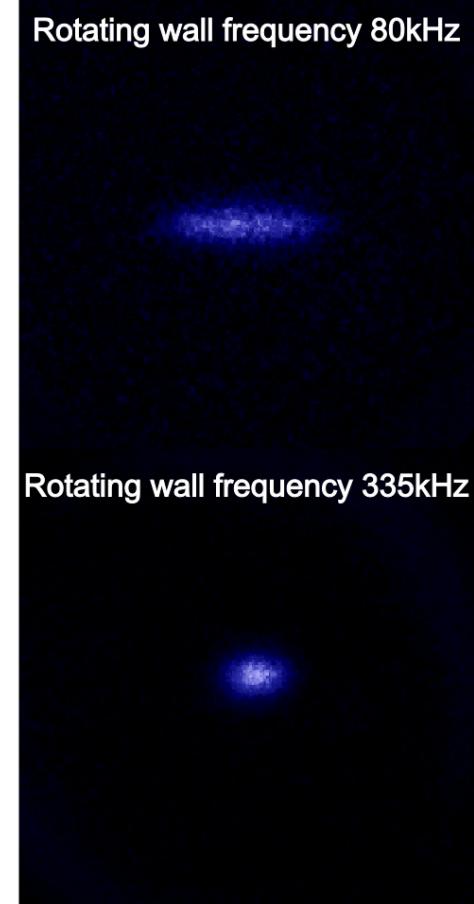
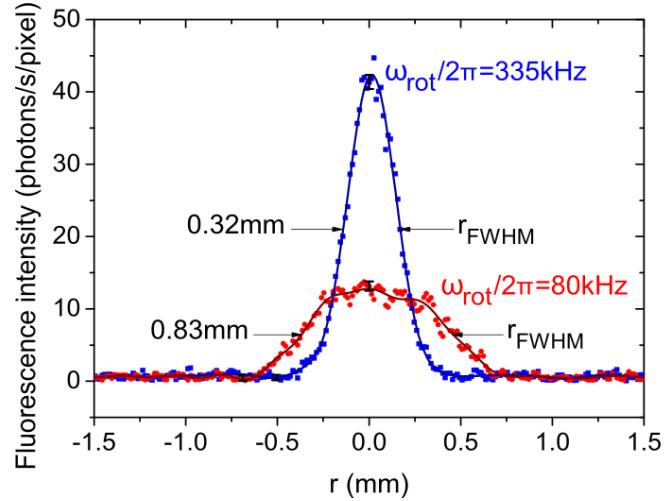
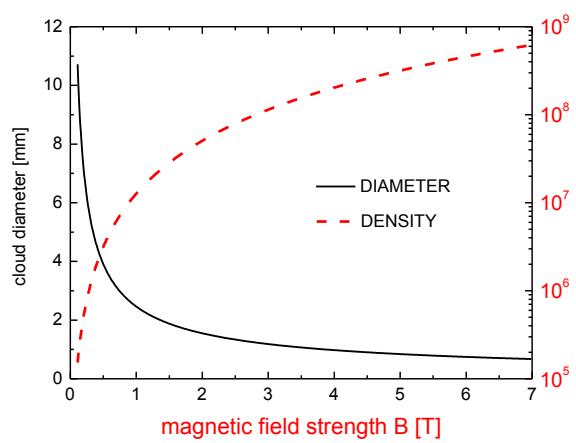
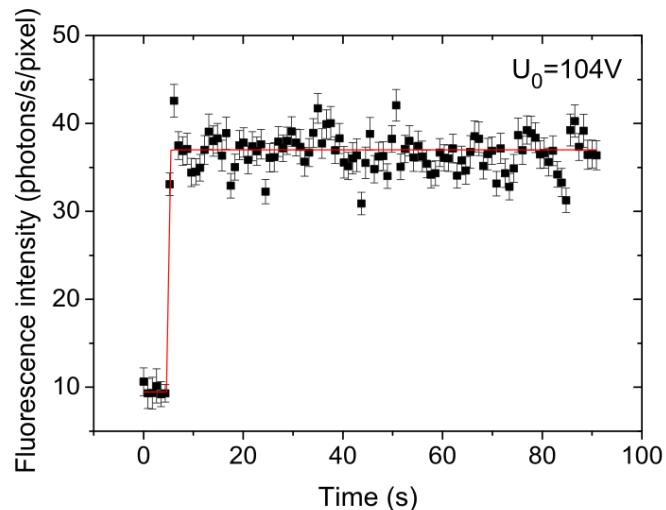
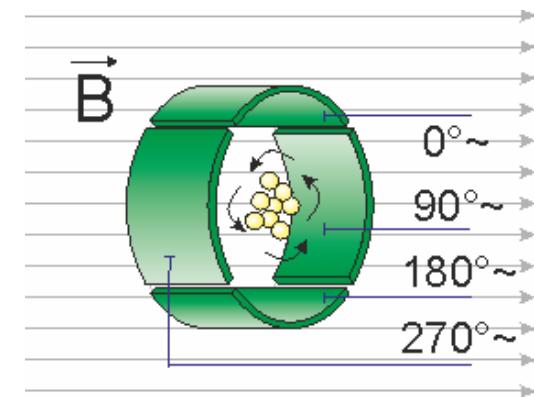


# techniques (2) : ion cooling

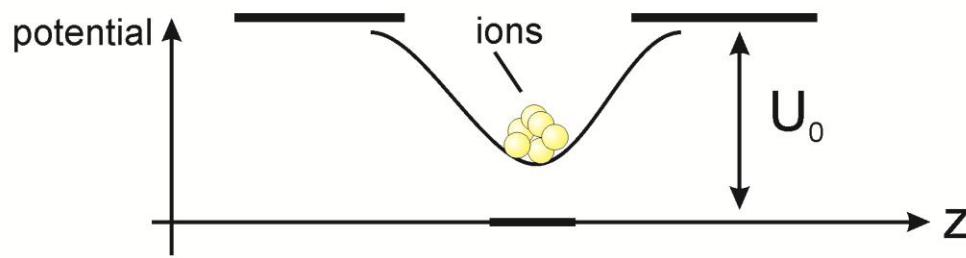
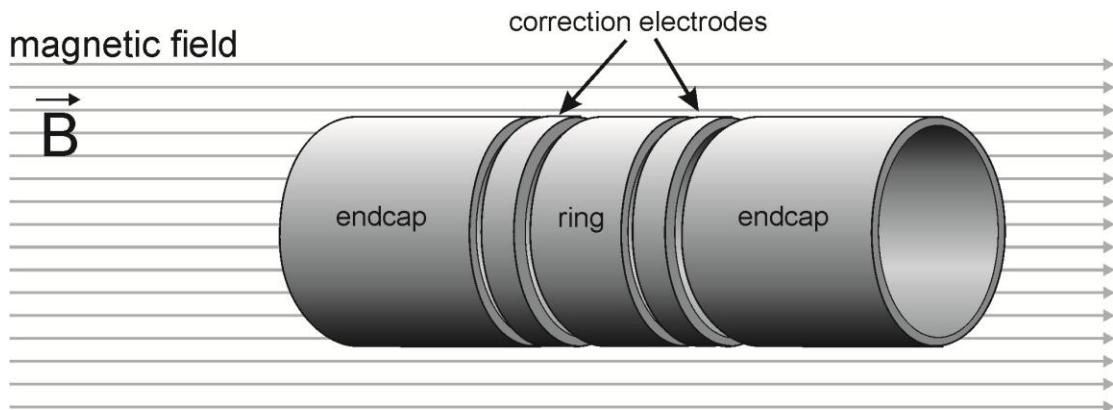
relative Doppler broadening of optical transitions



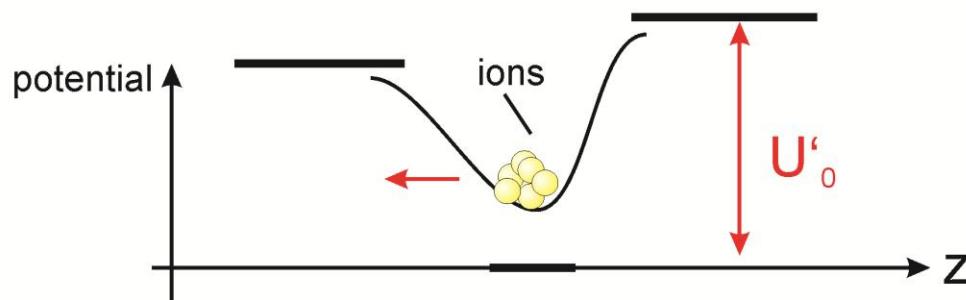
# techniques (3) : compression by rotating wall



# techniques (4) : ion positioning



ion position control  
relative to laser focus  
by potential asymmetry  
ms,  $\mu\text{m}$



shot-to-shot scan of  
the effective laser  
field strength

# combination: experimental cycle (example)

0 s

- ion capture from external source or in-trap production
- possible ion cooling to 4 K or below
- ion selection („clean target“) e.g. by „SWIFT“
- ion centering / positioning and compression by „rotating wall“

PREPARATION

5 s

@ 10 Hz

- laser interaction, e.g. multiphoton ionisation
- non-destructive detection (charge state evolution)
- product ions still in trap for further measurements...

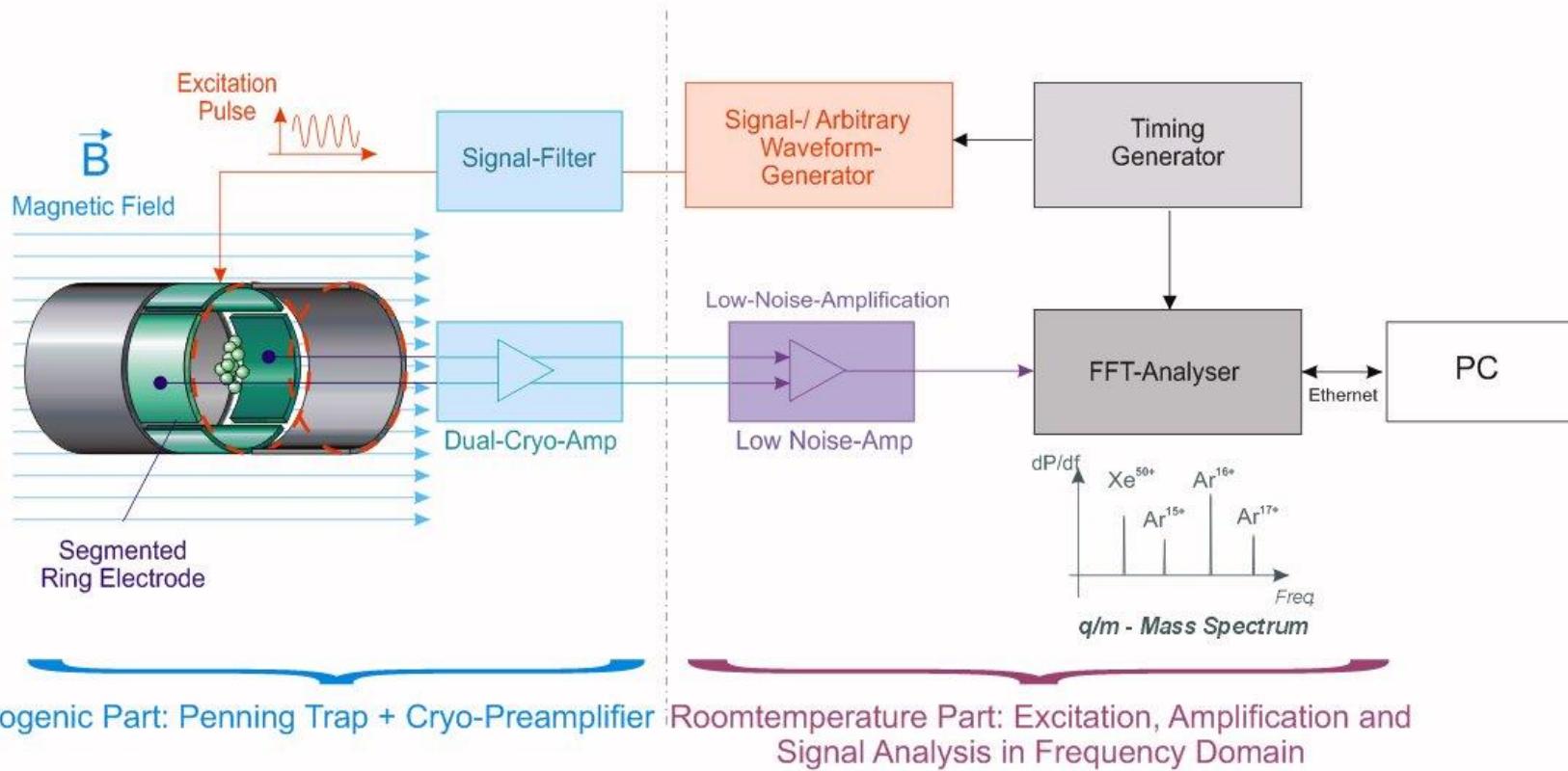
INVESTIGATION

# FTICR mass spectrometry

limit: of order 1000 charges (broad band), single ion (resonant)

10 Hz

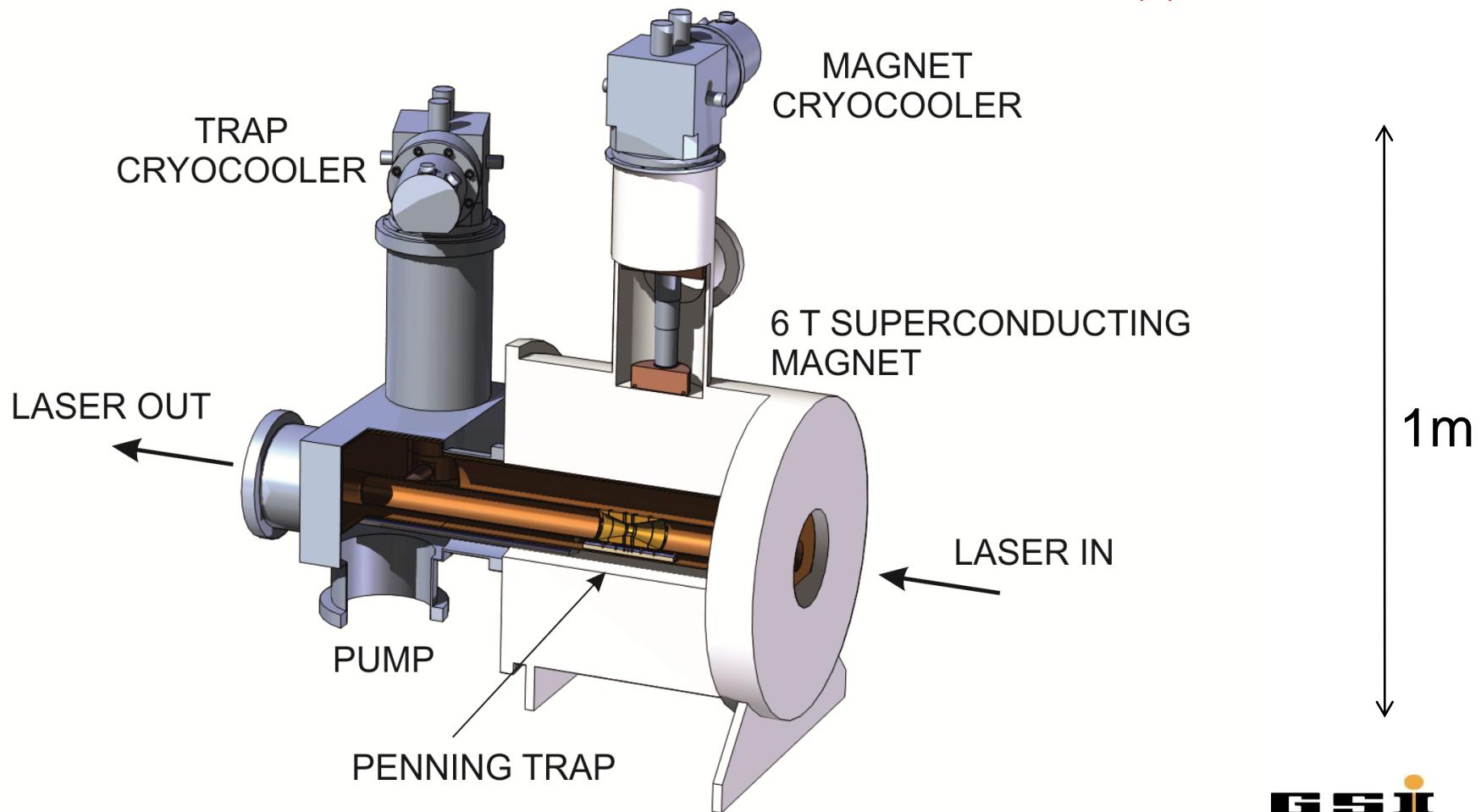
seconds



non-destructive detection

# setup schematic

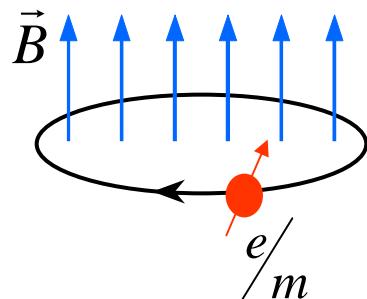
Penning trap in a dry superconducting magnet:  
high operation stability, high resolution, yet easy transport and flexible use  
@ PHELIX, FLASH (II), JETI, POLARIS,...



# Determination of the Proton g-Factor

$$\omega_c = \frac{e}{m_p} B$$

Cyclotron frequency



$$\omega_c = \sqrt{\omega_+^2 + \omega_-^2 + \omega_z^2}$$

$$\omega_+ \approx 2\pi \cdot 29 \text{ MHz}$$

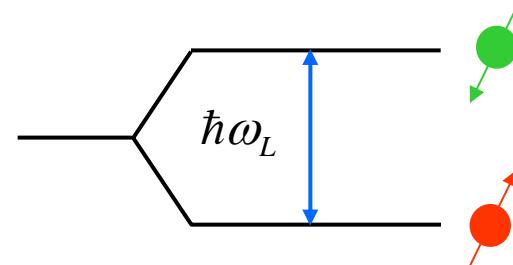
$$\omega_z \approx 2\pi \cdot 690 \text{ kHz}$$

$$\omega_- \approx 2\pi \cdot 8.5 \text{ kHz}$$

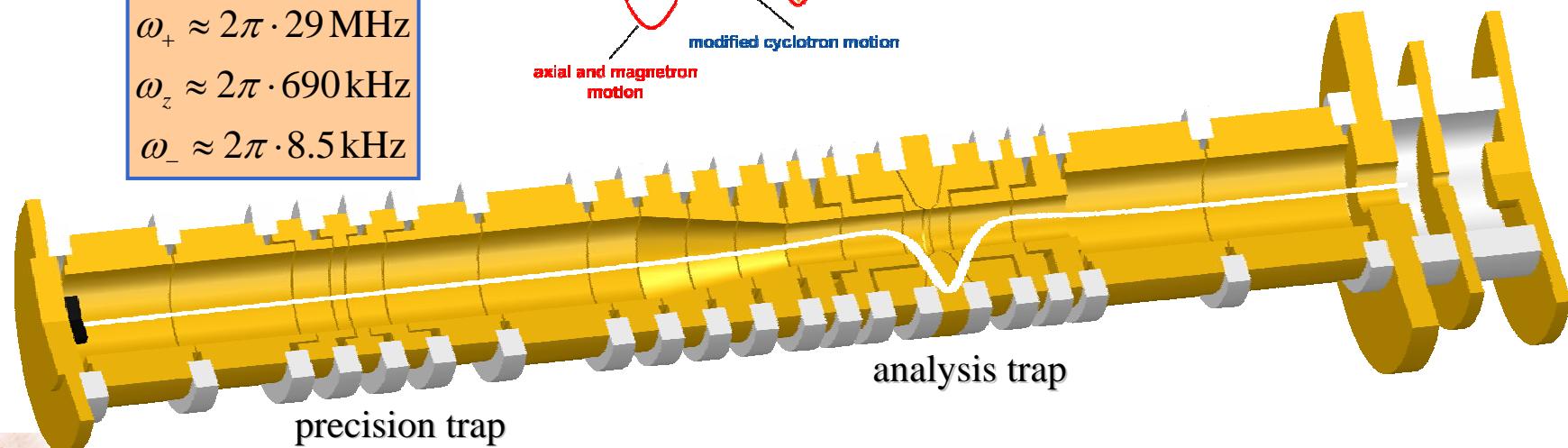
$$g = 2 \frac{\omega_L}{\omega_c}$$

$$\omega_L = g \frac{e}{2m_p} B$$

Larmor frequency



$$\omega'_z (\uparrow) - \omega'_z (\downarrow) = \Delta \omega_z$$



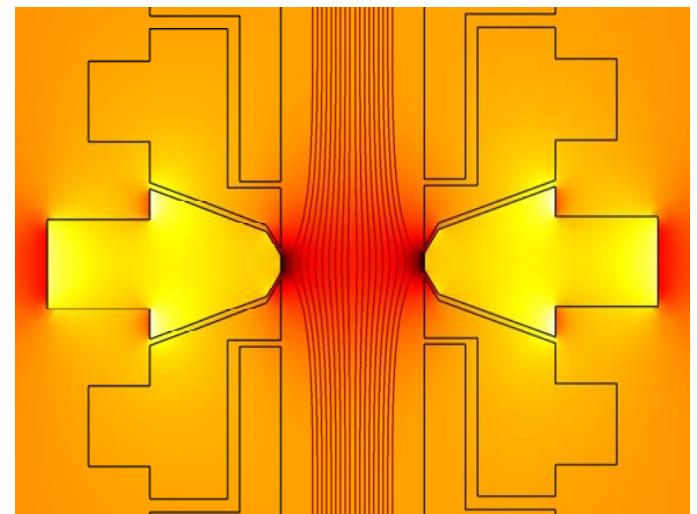
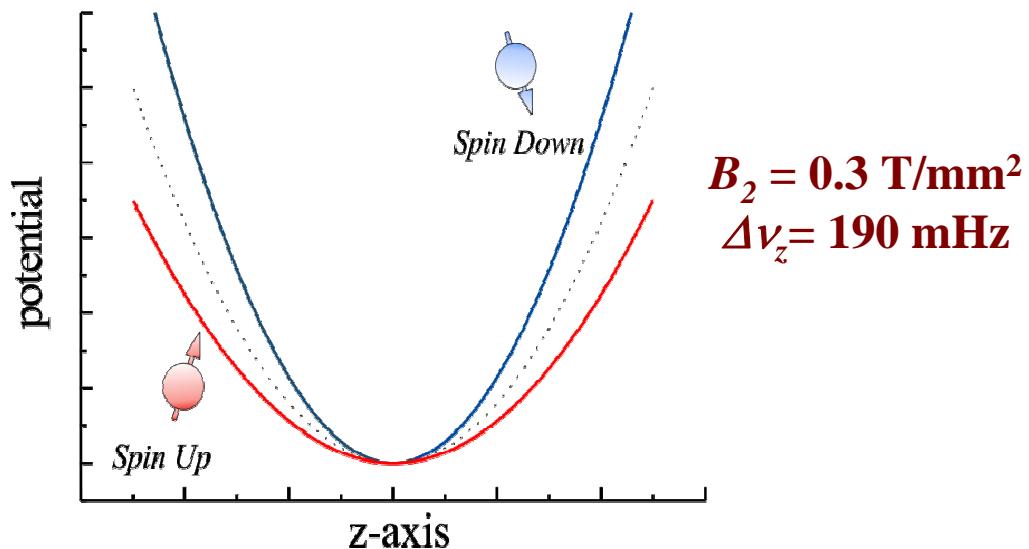
# A Single Trapped Proton and the Continuous Stern-Gerlach Effect

axial frequency shift  
due to spinflip:

$$\Delta\nu_z \approx \frac{1}{2\pi^2} \frac{\mu_z B_2}{m v_z}$$



Proton measurement is 10 000 times harder compared to electron g-2 measurement.

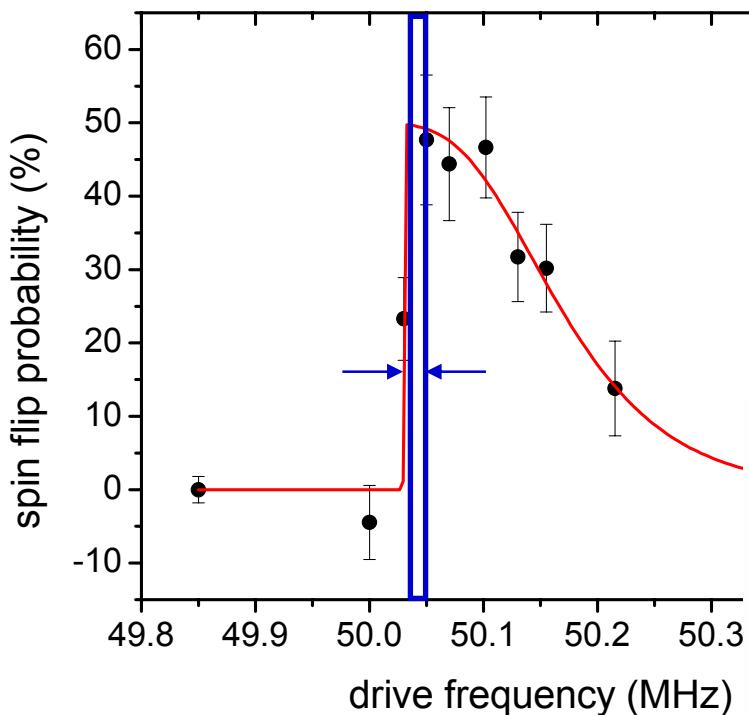


# First Larmor Resonance Curve of a Single Proton in the Analysis Trap

- ✓ Axial frequency fluctuation reduced
- ✓ Larmor resonance curve measured

$$\frac{\Delta g}{g} = 6 \cdot 10^{-4}$$

$$g = 2 \frac{\nu_L}{\nu_c}$$



PRL 106, 253001 (2011)

Selected for a Viewpoint in Physics  
PHYSICAL REVIEW LETTERS

week ending  
24 JUNE 2011

## Observation of Spin Flips with a Single Trapped Proton

S. Ulmer,<sup>1,2,3</sup> C. C. Rodegheri,<sup>1,2</sup> K. Blaum,<sup>1,3</sup> H. Kracke,<sup>2,4</sup> A. Mooser,<sup>2,4</sup> W. Quint,<sup>3,5</sup> and J. Walz<sup>2,4</sup>

<sup>1</sup>Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, D-69117 Heidelberg, Germany

<sup>2</sup>Institut für Physik, Johannes Gutenberg-Universität Mainz, D-55099 Mainz, Germany

<sup>3</sup>Ruprecht Karls-Universität Heidelberg, D-69047 Heidelberg, Germany

<sup>4</sup>Helmholtz Institut Mainz, D-55099 Mainz, Germany

<sup>5</sup>GSI—Helmholtzzentrum für Schwerionenforschung, D-64291 Darmstadt, Germany

(Received 28 February 2011; published 20 June 2011)

Radio-frequency induced spin transitions of one individual proton are observed. The spin quantum jumps are detected via the continuous Stern-Gerlach effect, which is used in an experiment with a single proton stored in a cryogenic Penning trap. This is an important milestone towards a direct high-precision measurement of the magnetic moment of the proton and a new test of the matter-antimatter symmetry in the baryon sector.

# The Ion g-Factor Team in Mainz (before spinflip)



Birgit  
Schabinger

Sven  
Sturm

Anke  
Wagner

# The Ion g-Factor Team (after spinflip)



# Bound Electron Magnetic Moment Measurement on Hydrogen-like Silicon $^{28}\text{Si}^{13+}$

PRL 107, 023002 (2011)

PHYSICAL REVIEW LETTERS

week ending  
8 JULY 2011

## $\mathfrak{g}$ $g$ Factor of Hydrogenlike $^{28}\text{Si}^{13+}$

S. Sturm,<sup>1,2</sup> A. Wagner,<sup>1</sup> B. Schabinger,<sup>1,2</sup> J. Zatorski,<sup>1</sup> Z. Harman,<sup>1,3</sup> W. Quint,<sup>4</sup> G. Werth,<sup>2</sup> C. H. Keitel,<sup>1</sup> and K. Blaum<sup>1</sup>

<sup>1</sup>Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, 69117 Heidelberg, Germany

<sup>2</sup>Institut für Physik, Johannes Gutenberg-Universität, 55099 Mainz, Germany

<sup>3</sup>ExtreMe Matter Institute EMMI, Planckstraße 1, 64291 Darmstadt, Germany

<sup>4</sup>GSI Helmholtzzentrum für Schwerionenforschung GmbH, Planckstraße 1, 64291 Darmstadt, Germany

(Received 6 May 2011; published 7 July 2011)

We determined the experimental value of the  $g$  factor of the electron bound in hydrogenlike  $^{28}\text{Si}^{13+}$  by using a single ion confined in a cylindrical Penning trap. From the ratio of the ion's cyclotron frequency and the induced spin flip frequency, we obtain  $g = 1.995\,348\,958\,7(5)(3)(8)$ . It is in excellent agreement with the state-of-the-art theoretical value of  $1.995\,348\,958\,0(17)$ , which includes QED contributions up to the two-loop level of the order of  $(Z\alpha)^2$  and  $(Z\alpha)^4$  and represents a stringent test of bound-state quantum electrodynamics calculations.

