



PHOTOIONIZATION DYNAMICS WITH ATTO-SECOND PULSE TRAINS

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Dynamics in real time

To capture a moving object we need...

Camera with long exposure time



an exposure time /shutter faster than the motion !

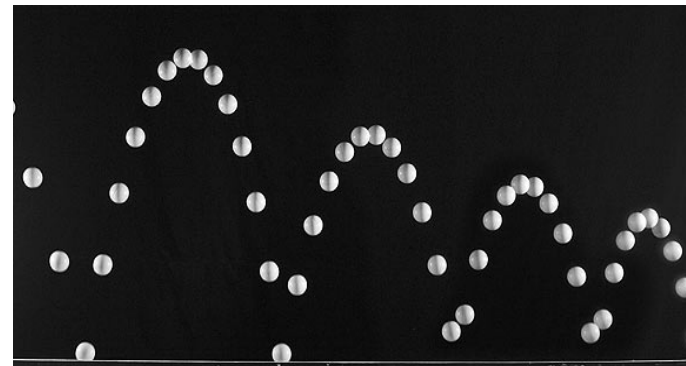
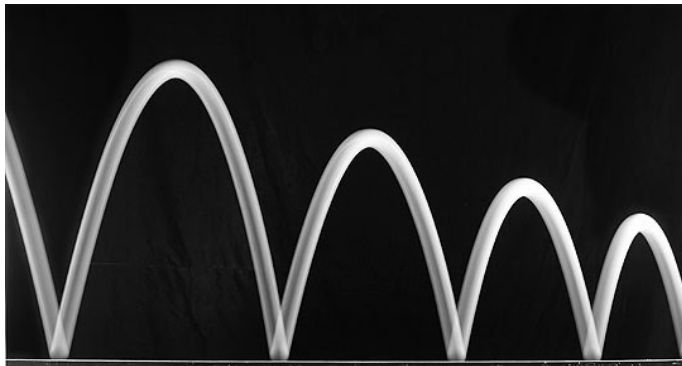


Sequential approach



Dynamics in real time

Camera with long exposure time and a strobe light



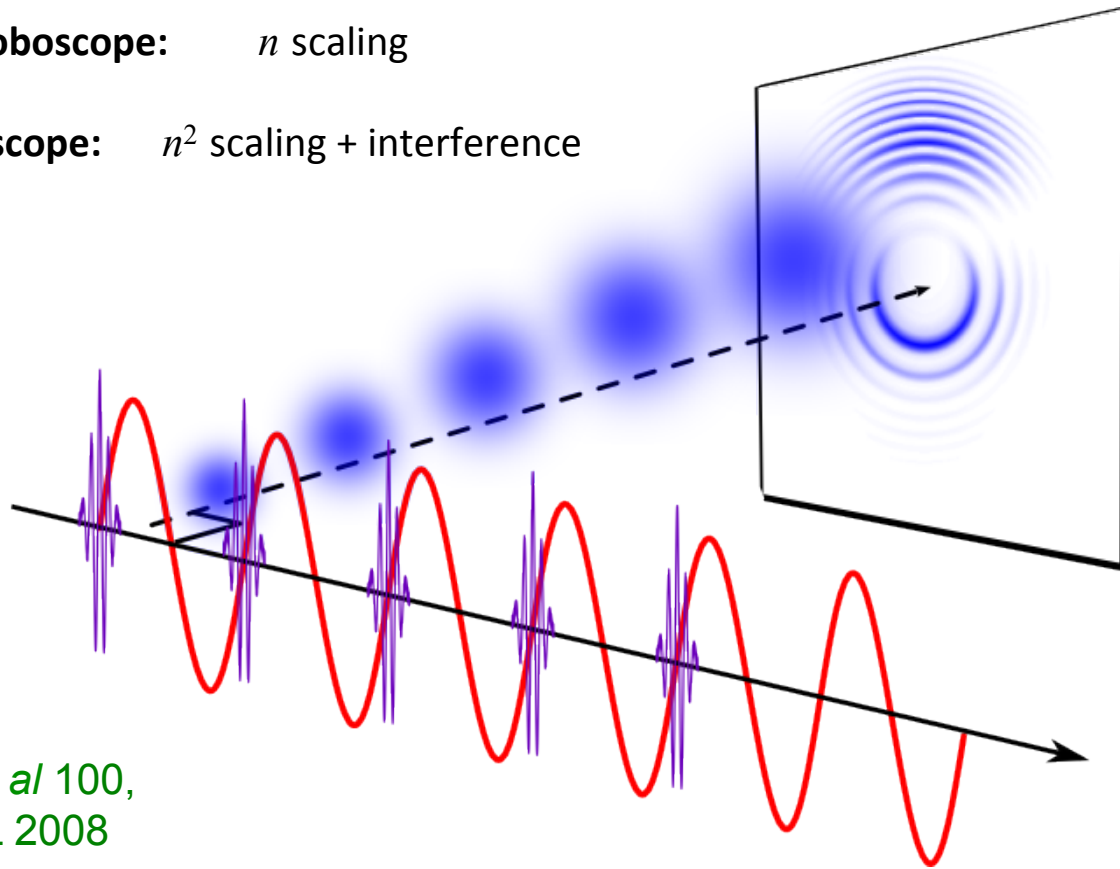
Stroboscope approach



The quantum stroboscope with Attosecond Pulse Trains (APT)

Conventional stroboscope: n scaling

Quantum stroboscope: n^2 scaling + interference



Mauritsson *et al* 100,
073003 PRL 2008



Applications

Atto-second Pulse Trains

Weak IR field

Strong IR Field

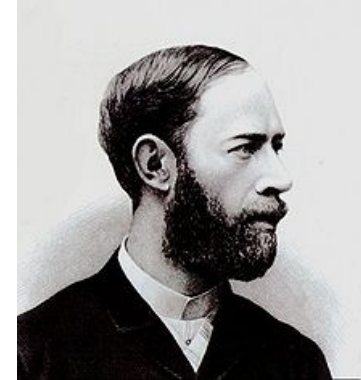
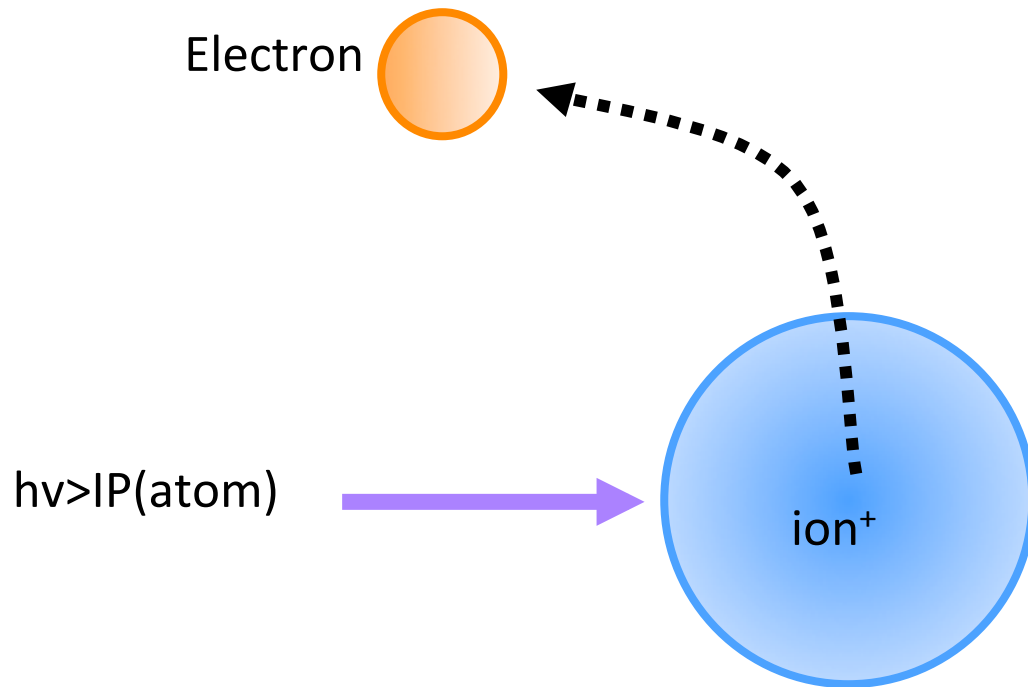


Interference in the
frequency domain

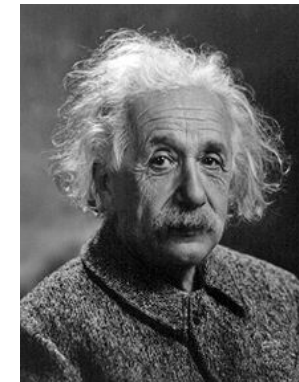
Interference in the
time domain



Photo-ionization



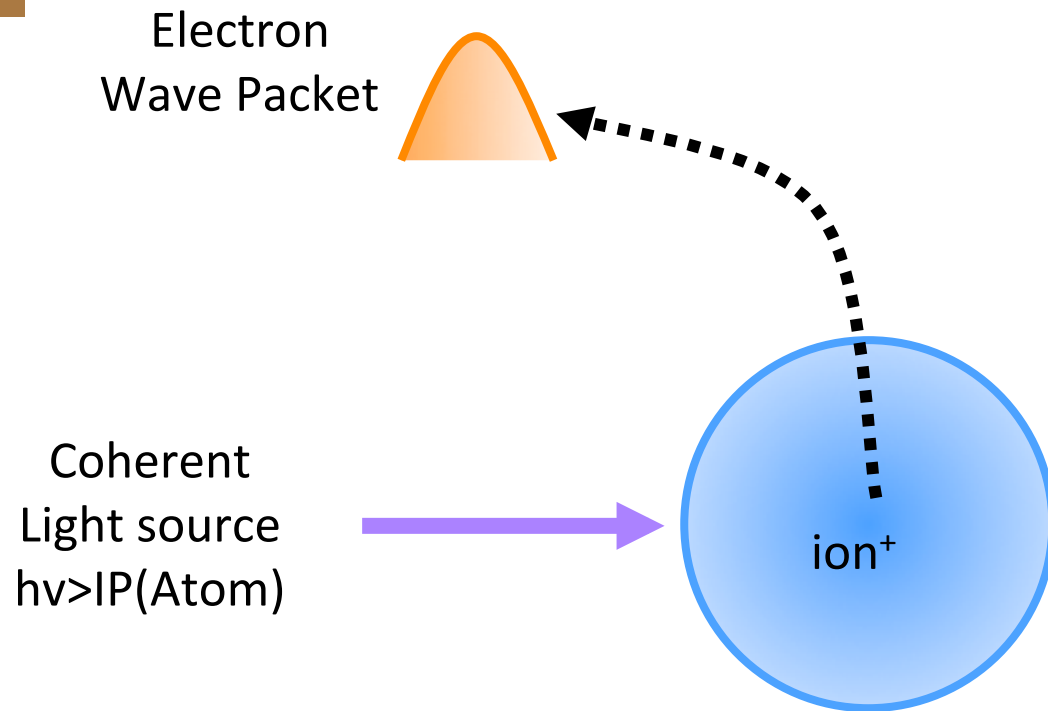
H.R Hertz - 1887



A. Einstein - 1905



Photo-ionization with APTs



Atto-second regime

$$\Delta t < 1 \text{ fs} \longleftrightarrow \Delta E > 2 \text{ eV}$$



Wave-particle duality



Photo-ionization with APTs

Electron



Can we "read-out" properties
of **one**-photon ionization
in the Electron Wave Packet ?

Ligh
 $h\nu >$



Wave-particle duality



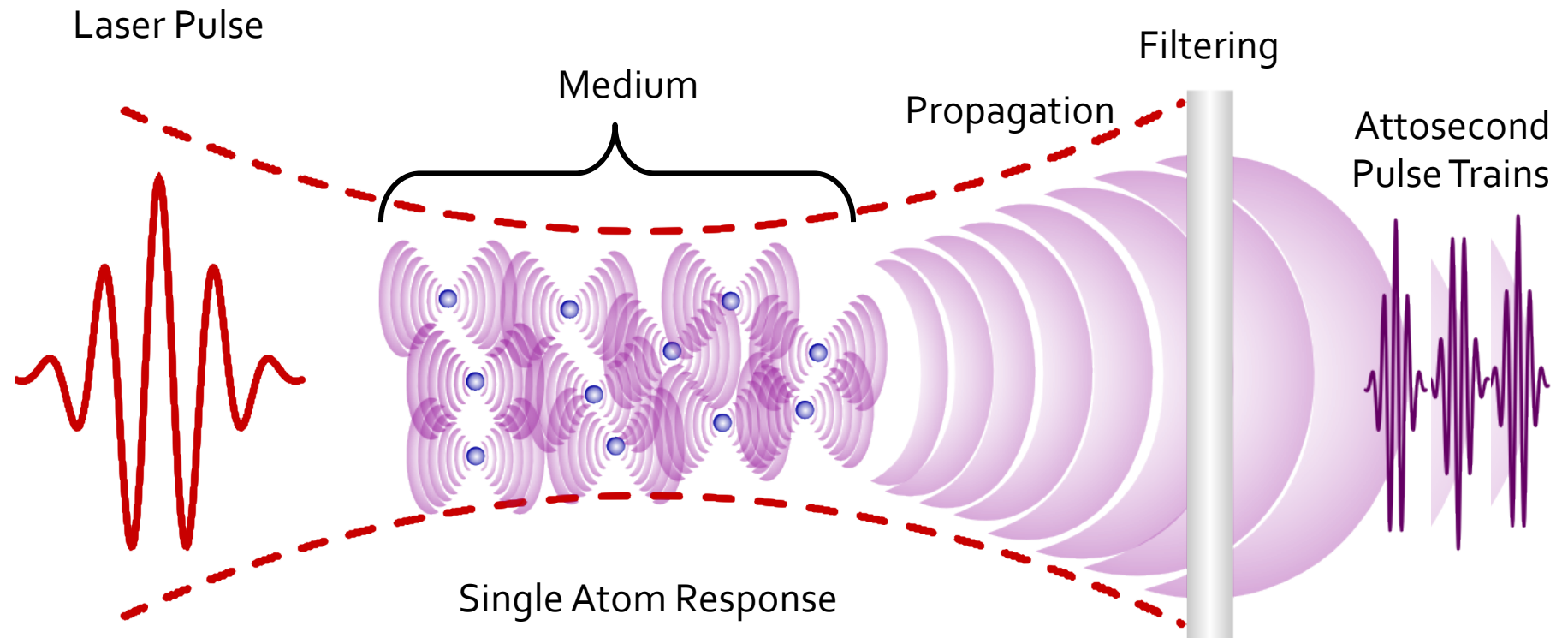


TOOLS

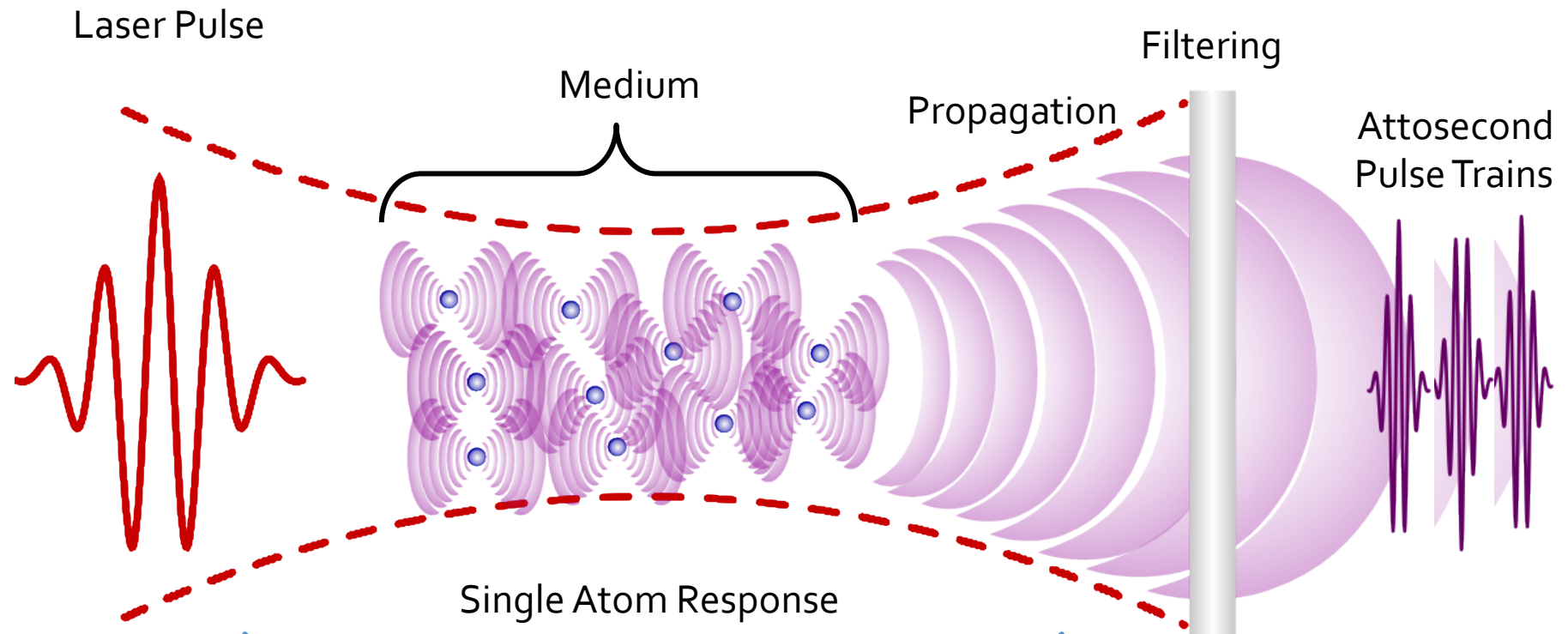
Workshop of the Extreme Matter Institute, Darmstadt, 2011



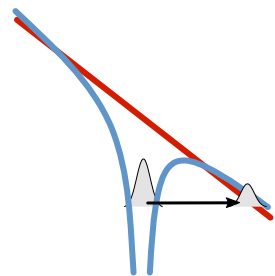
Attosecond Pulse Trains



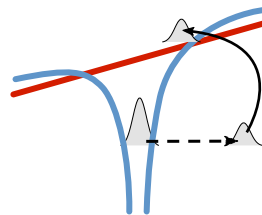
Attosecond Pulse Trains



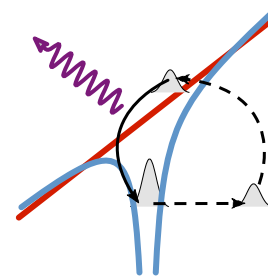
Single Atom Response



1. Tunneling



2. Acceleration

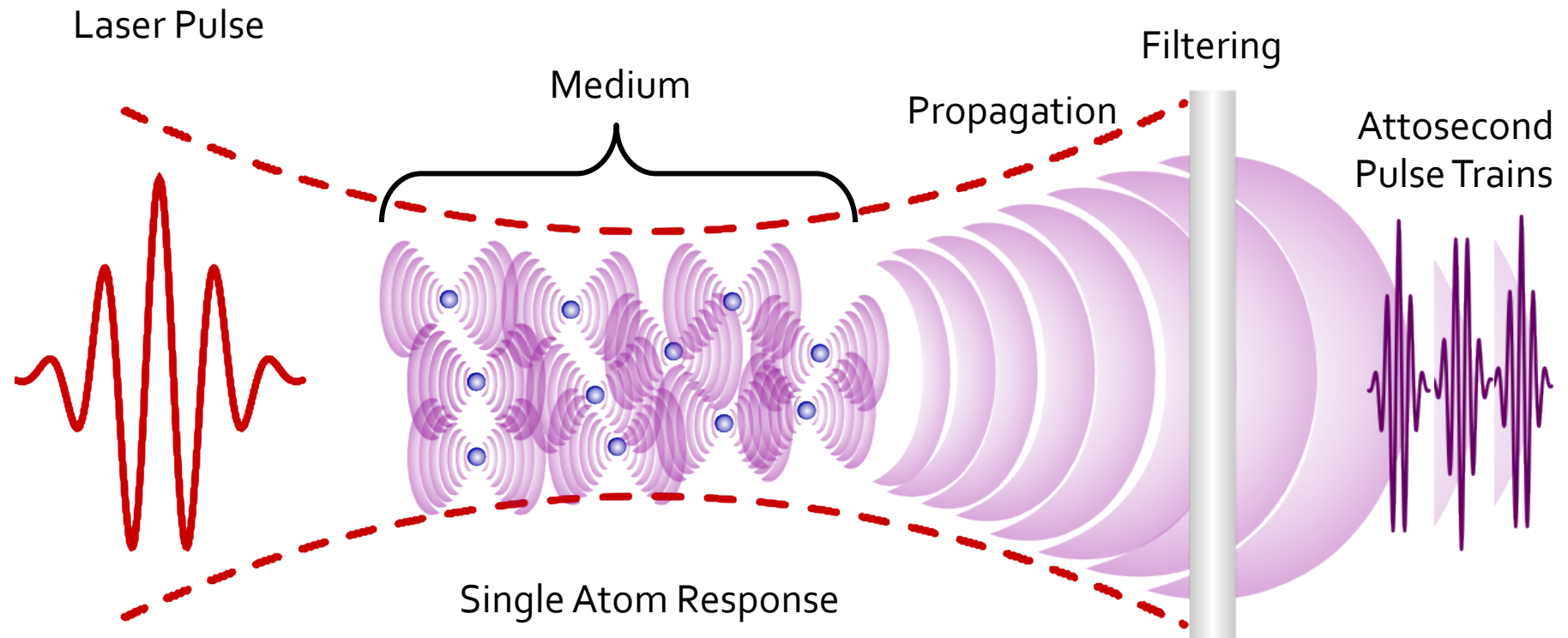


3. Recombination

P. Corkum, PRL
71, 1994 (1993)



Attosecond Pulse Trains



$$x(t) \propto i \int_{-\infty}^t d\tau \int d\vec{p} d^* (\vec{p} + e\vec{A}(t)) e^{-iS(\vec{p}, t, \tau)/\hbar} E(t - \tau) d(\vec{p} + e\vec{A}(t - \tau))$$

Lewenstein, 49,
2117, PRA 1994

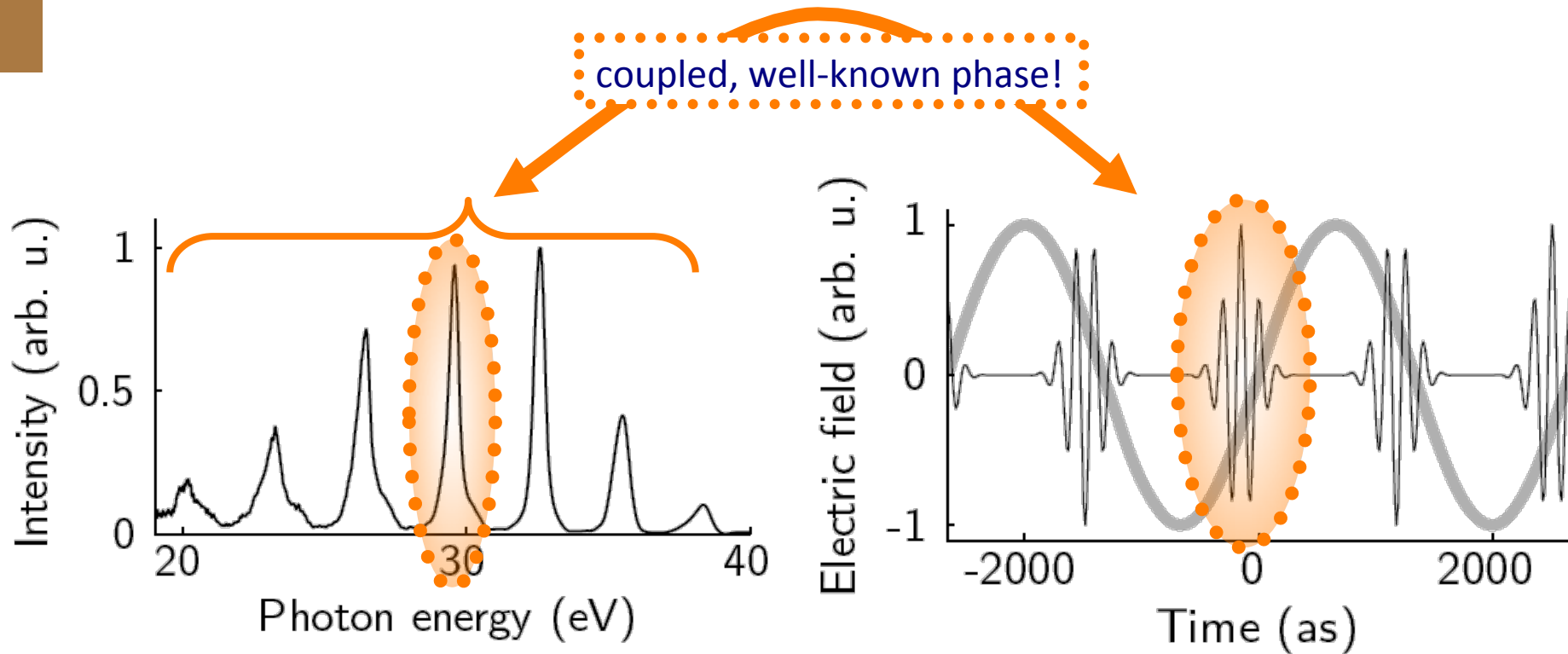
Recombination

Acceleration in
the continuum

Transition to the
continuum



Attosecond Pulse Trains

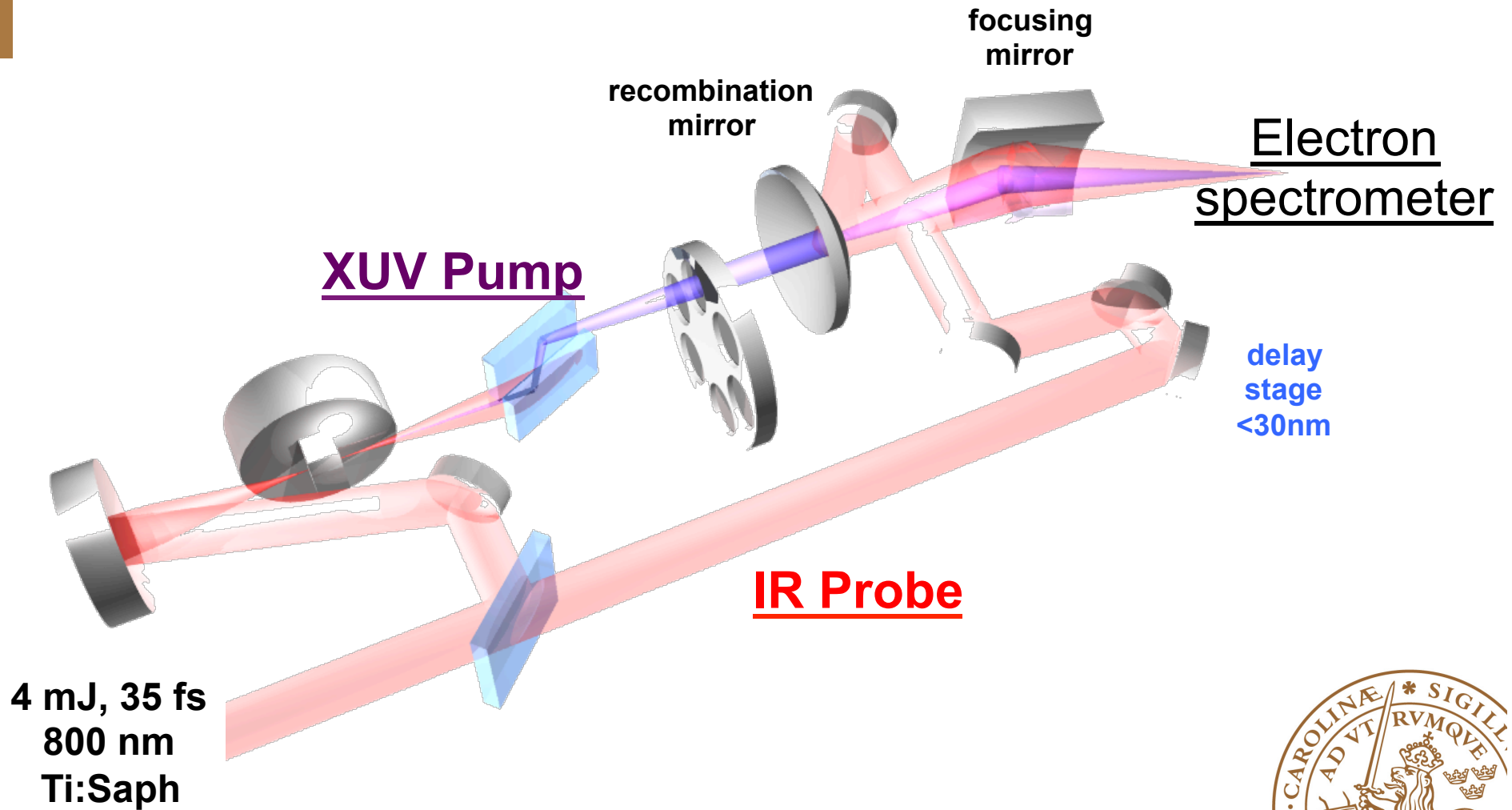


100 meV bandwidth/harmonic!

260 as duration/pulse!



Pump-probe setup





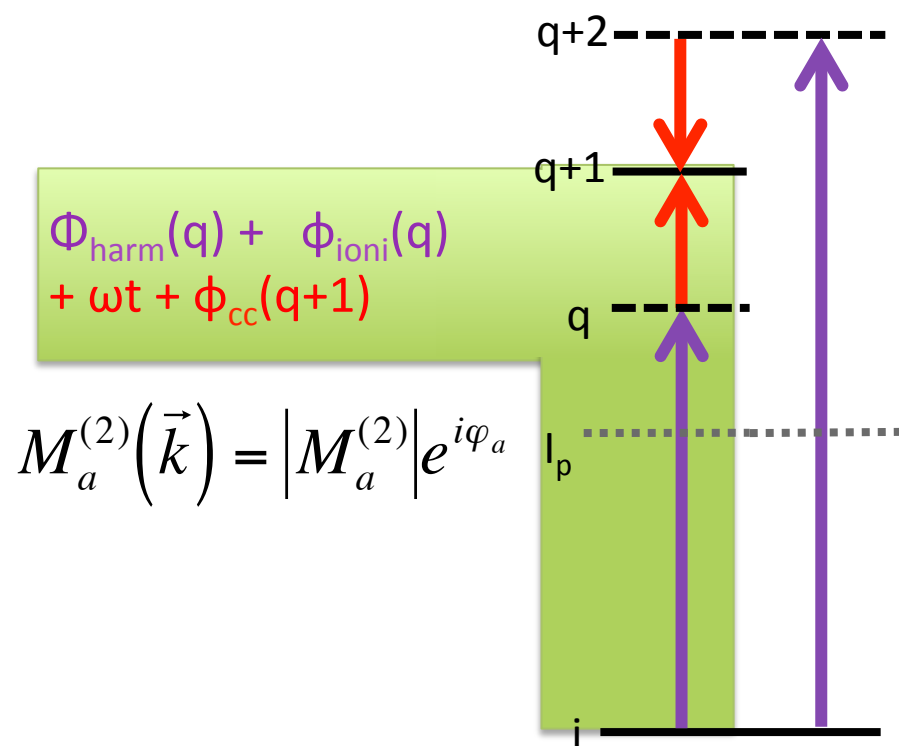
INTERFEROMETRY IN FREQUENCY DOMAIN

Workshop of the Extreme Matter Institute, Darmstadt, 2011



Interferometric measurement

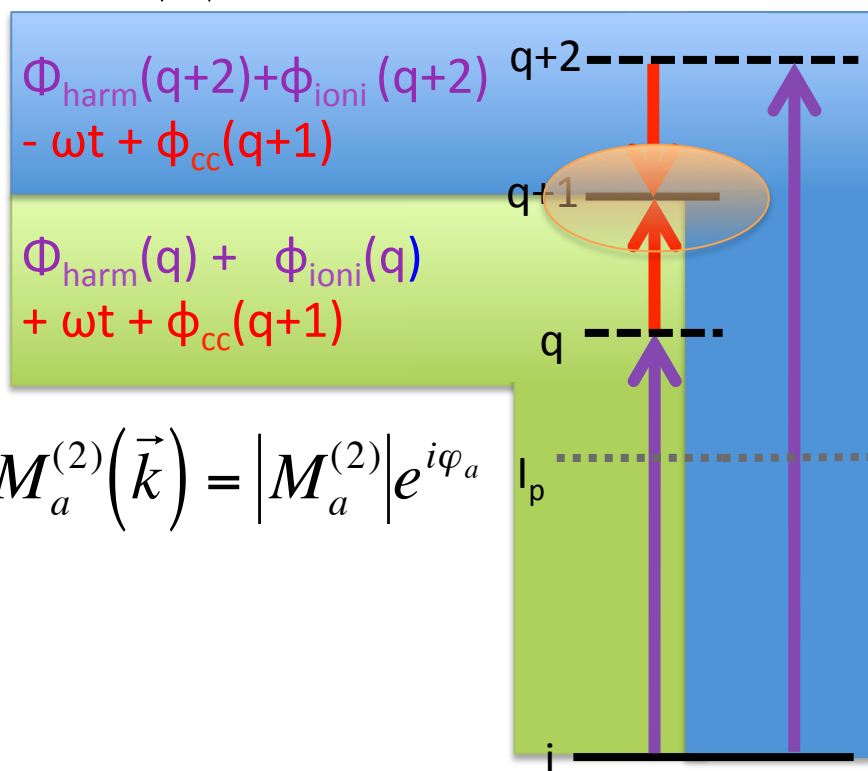
Ionization with two harmonics in the presence
of a weak infrared field



Interferometric measurement

Ionization with two harmonics in the presence of a weak infrared field

$$M_e^{(2)}(\vec{k}) = |M_e^{(2)}| e^{i\varphi_e}$$



$$M_a^{(2)}(\vec{k}) = |M_a^{(2)}| e^{i\varphi_a}$$

Interference between two quantum paths

$$S_{q+1} = |M_a^{(2)}(\vec{k}) + M_e^{(2)}(\vec{k})|^2$$

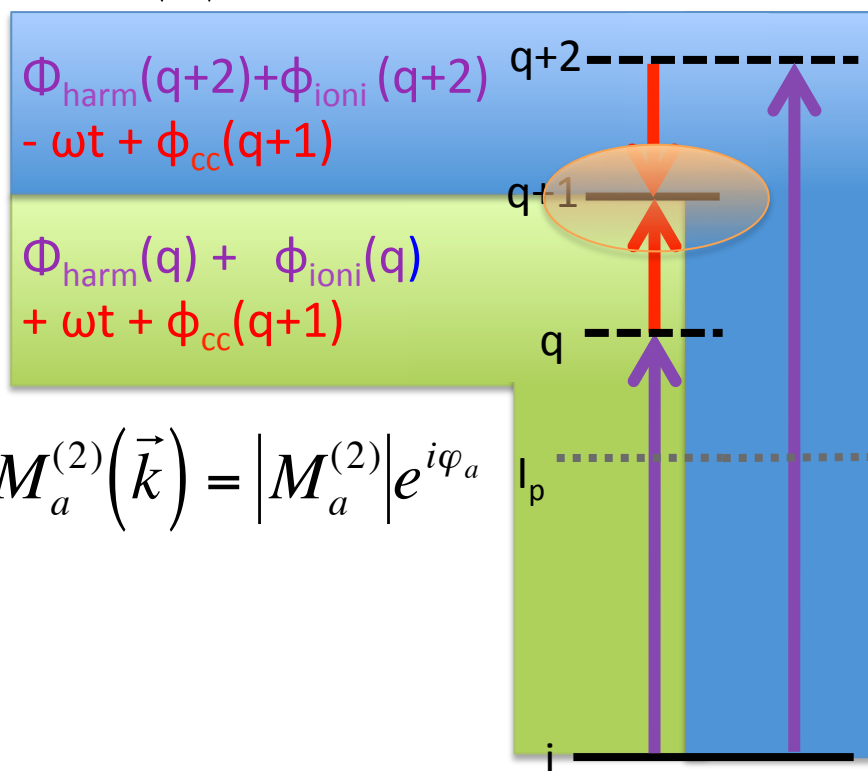
$$S_{q+1} \propto \cos(2\omega t - \Delta\Phi_{harm} - \Delta\Phi_{ioni})$$



Interferometric measurement

Ionization with two harmonics in the presence of a weak infrared field

$$M_e^{(2)}(\vec{k}) = |M_e^{(2)}| e^{i\varphi_e}$$



$$M_a^{(2)}(\vec{k}) = |M_a^{(2)}| e^{i\varphi_a}$$

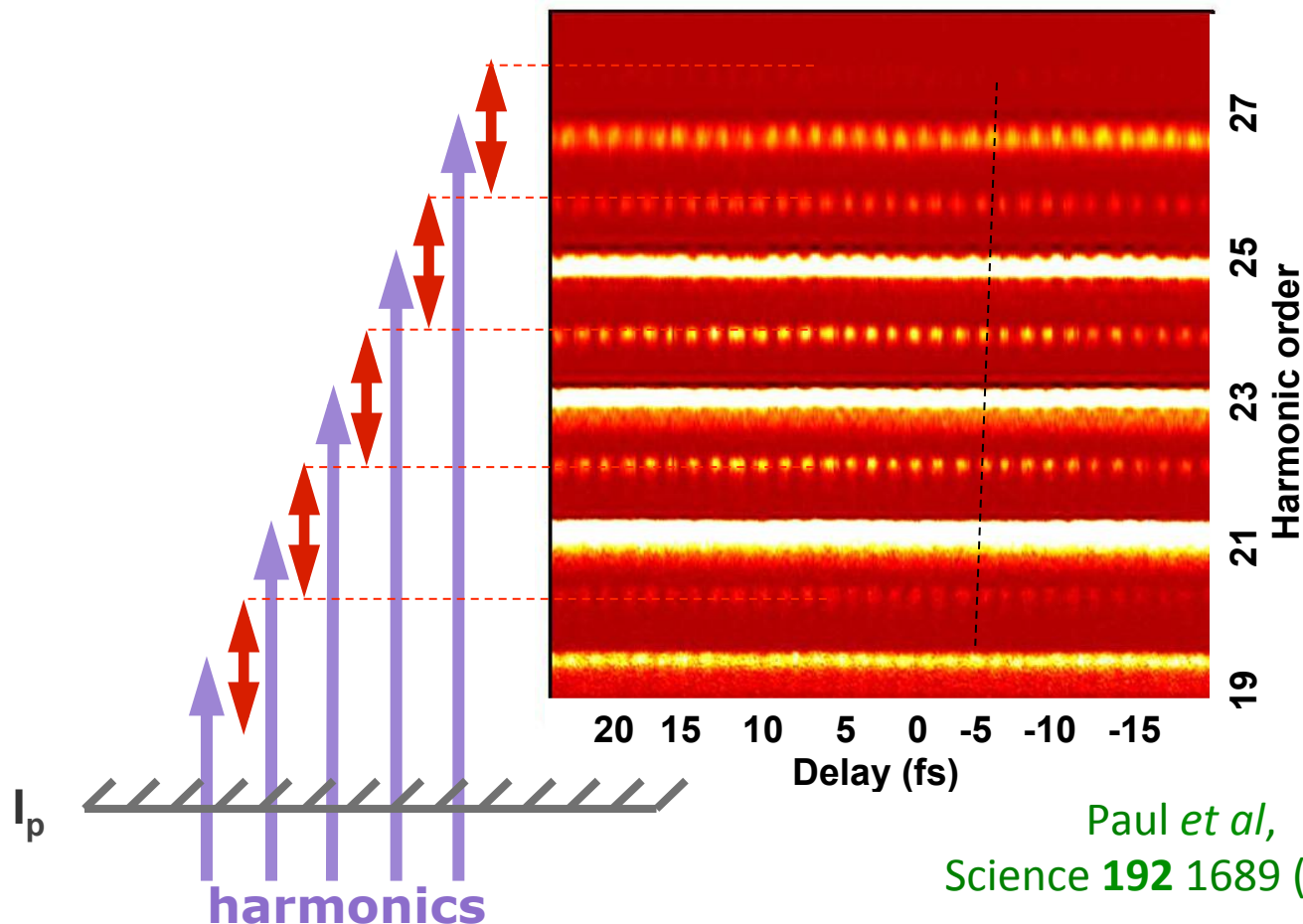
Interference between two quantum paths

$$S_{q+1} = |M_a^{(2)}(\vec{k}) + M_e^{(2)}(\vec{k})|^2$$

$$S_{q+1} \propto \cos(2\omega t - \Delta\Phi_{harm} - \Delta\Phi_{ioni}^{\approx 0})$$



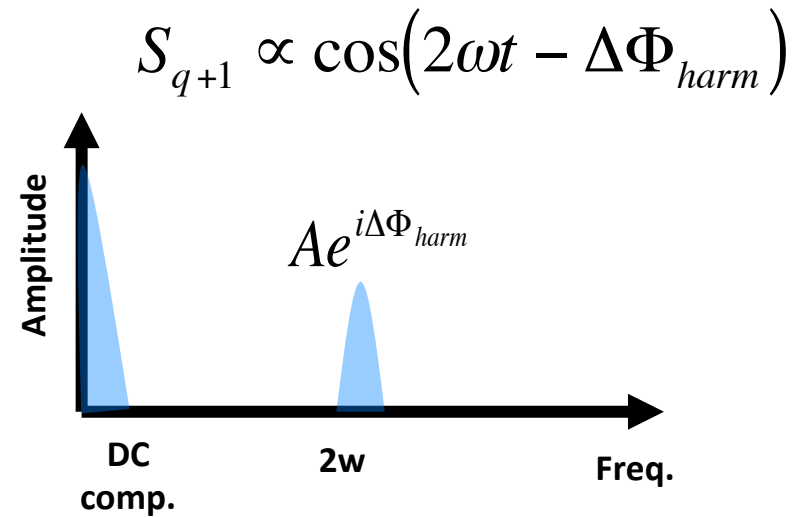
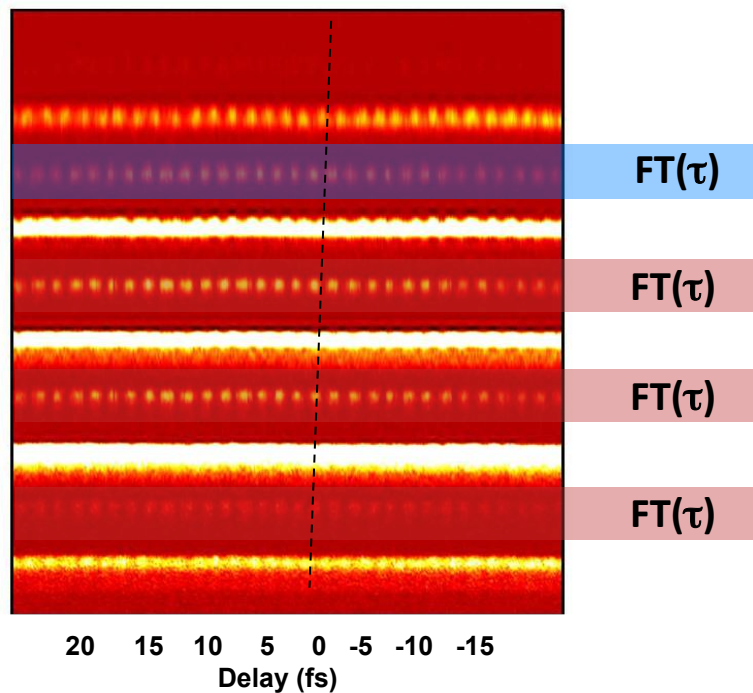
Reconstruction of Attosecond Beating by Interference of Two-photon Transition



Paul *et al*,
Science **192** 1689 (2001)



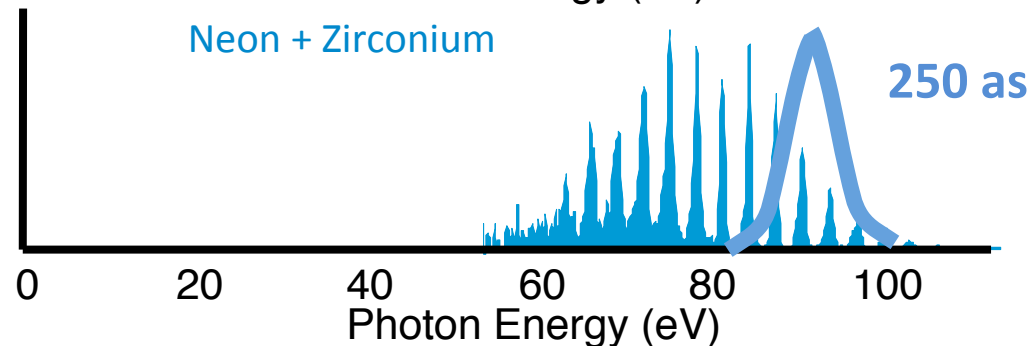
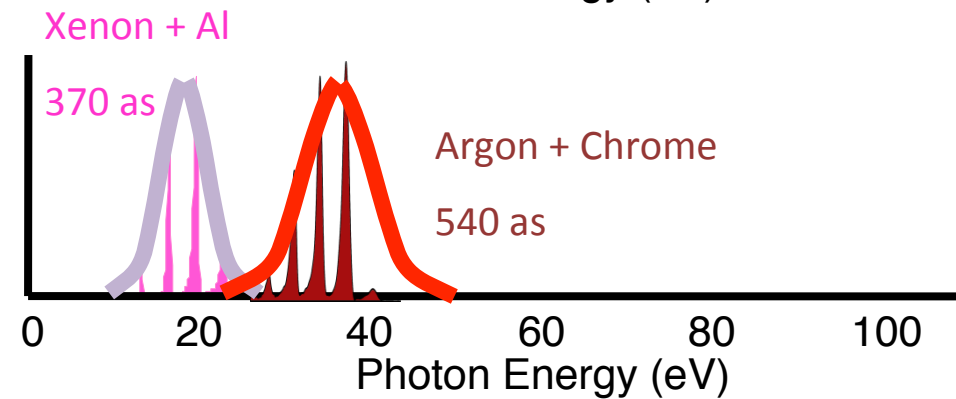
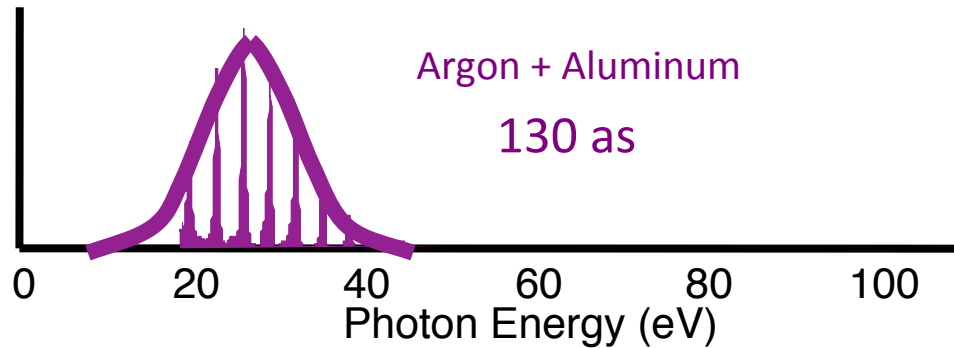
Reconstruction of Attosecond Beating by Interference of Two-photon Transition



Phase determination down to 20 as !



Reconstruction of Attosecond Beating by Interference of Two-photon Transition



...
López-Martens *et al*,
Phys Rev Lett **94** 3301 (2005)

...
E. Gustafsson *et al*,
Optics Letters **32**, 1353 (2007)

...



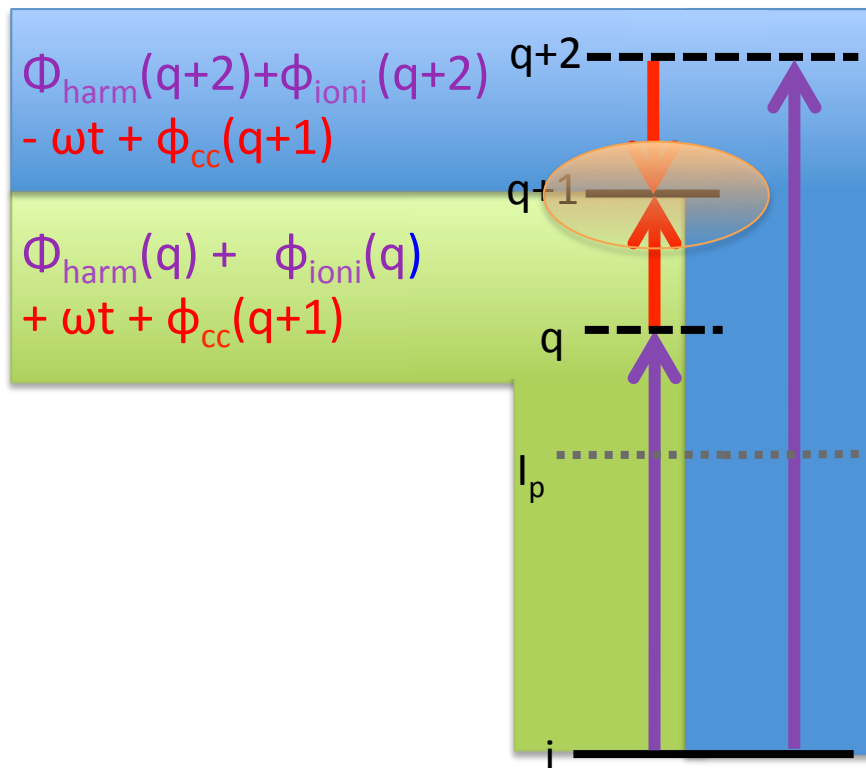


Resonant photo-excitation – atomic case



Interferometric measurement

Ionization with two harmonics in the presence of a weak infrared field



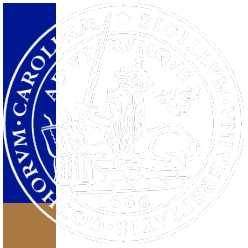
Interference between two quantum paths

$$S_{q+1} \sim \cos(2\omega t - \Delta\phi_{\text{harm}} - \Delta\phi_{\text{ioni}})$$

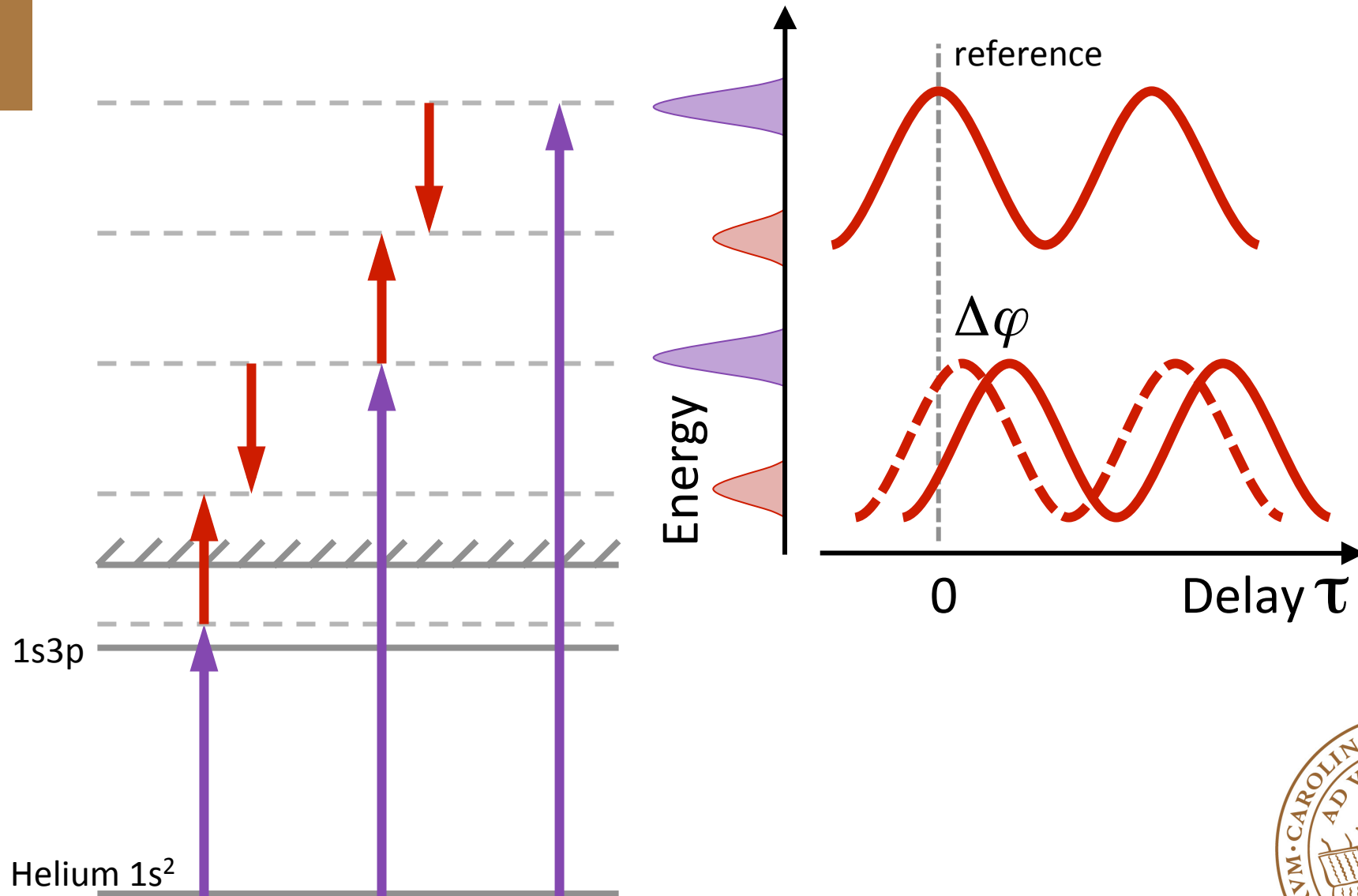
Idea

“Two photon Resonant the photoionization pathway will introduce a phase jump of π ”



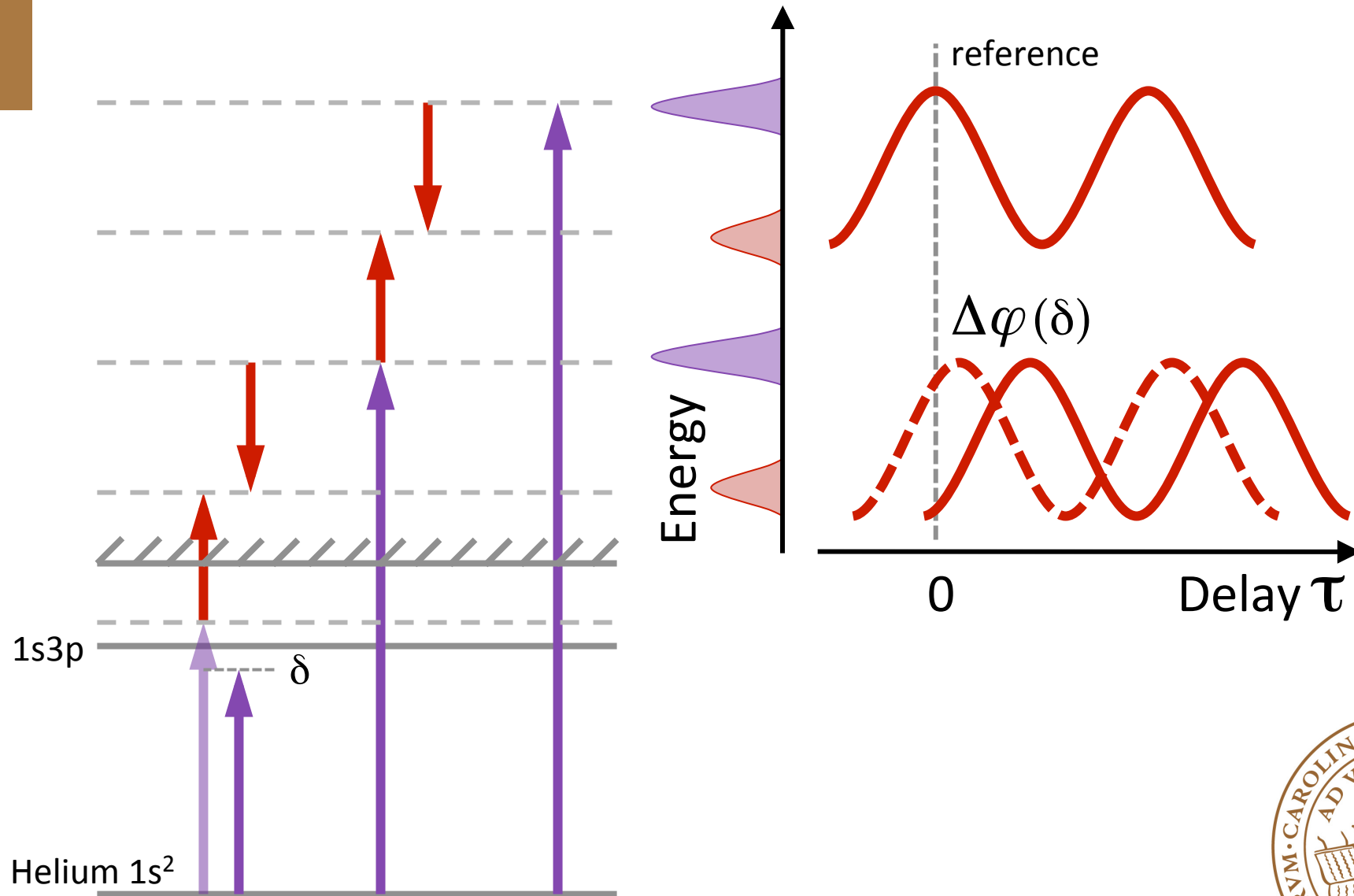


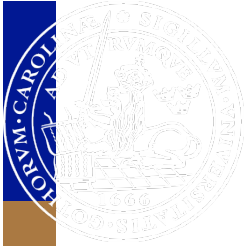
Two-Color Two-Photon Ionization





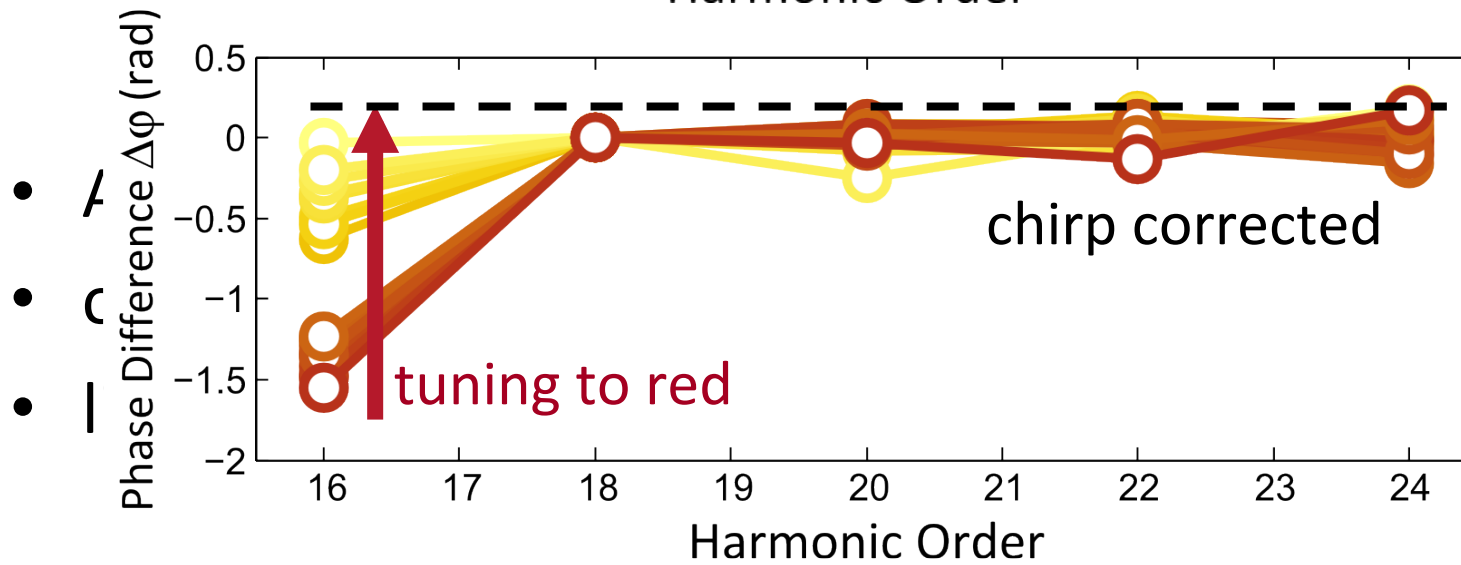
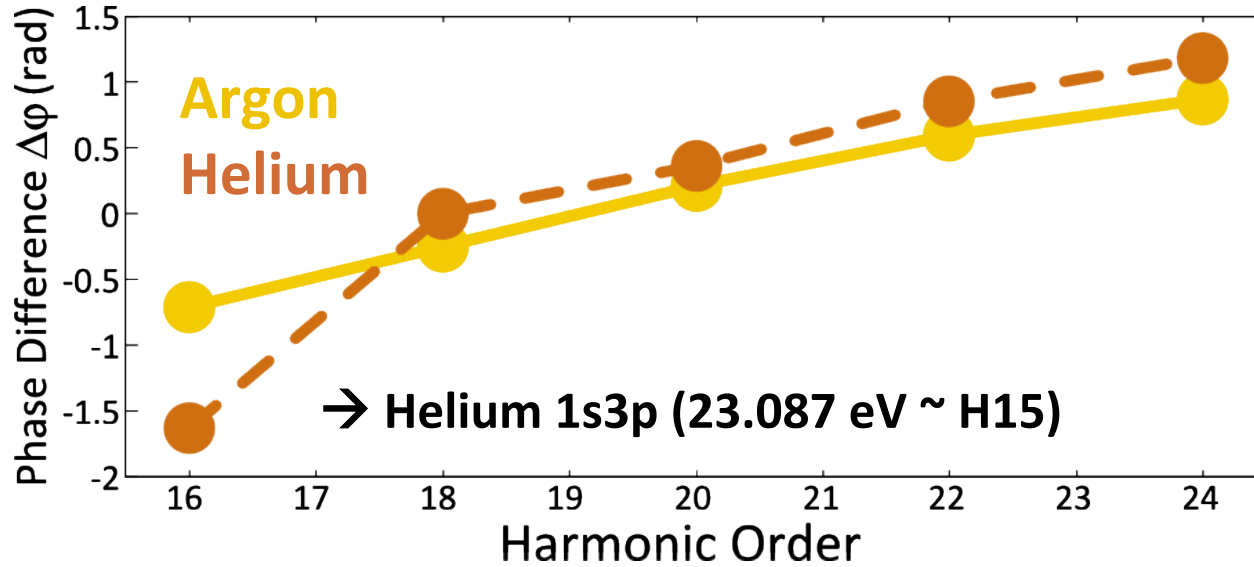
Two-Color Two-Photon Ionization





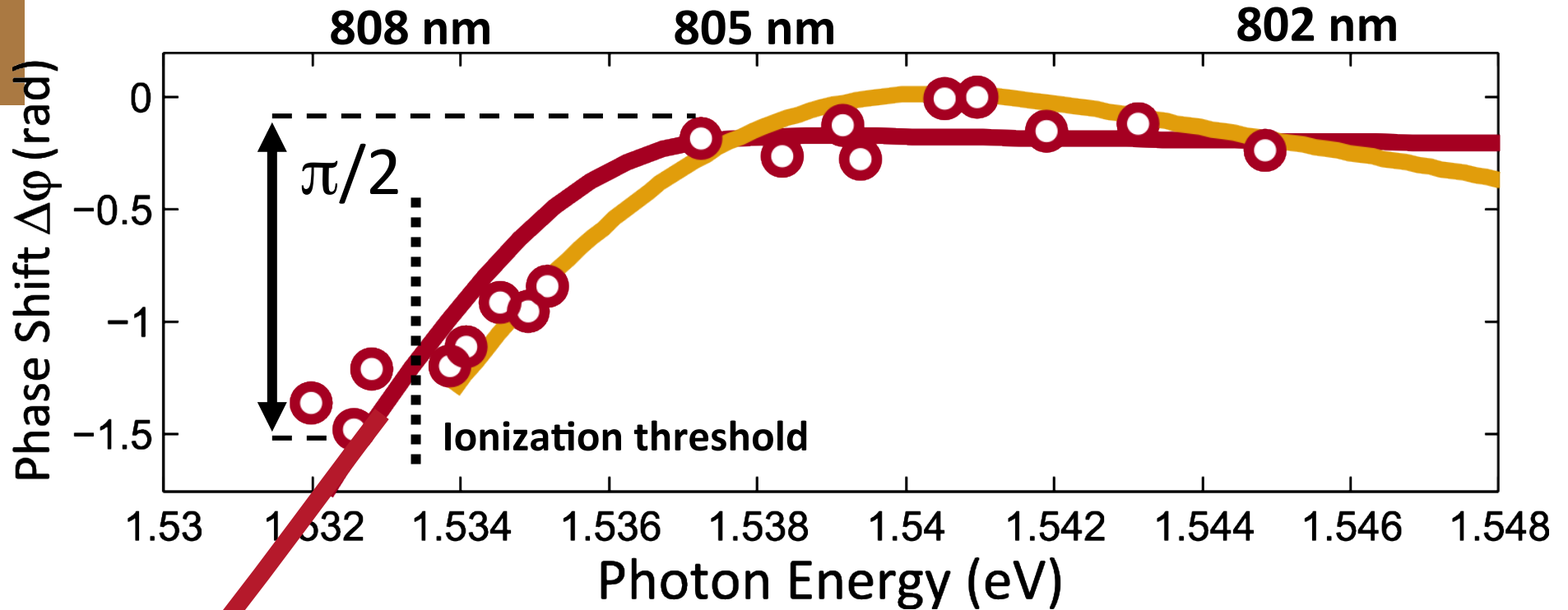
Harmonic $\Delta\varphi$: Argon vs Helium

Swodoba *et al*, PRL **104**, 103003 (2010)





Wavelength Dependence

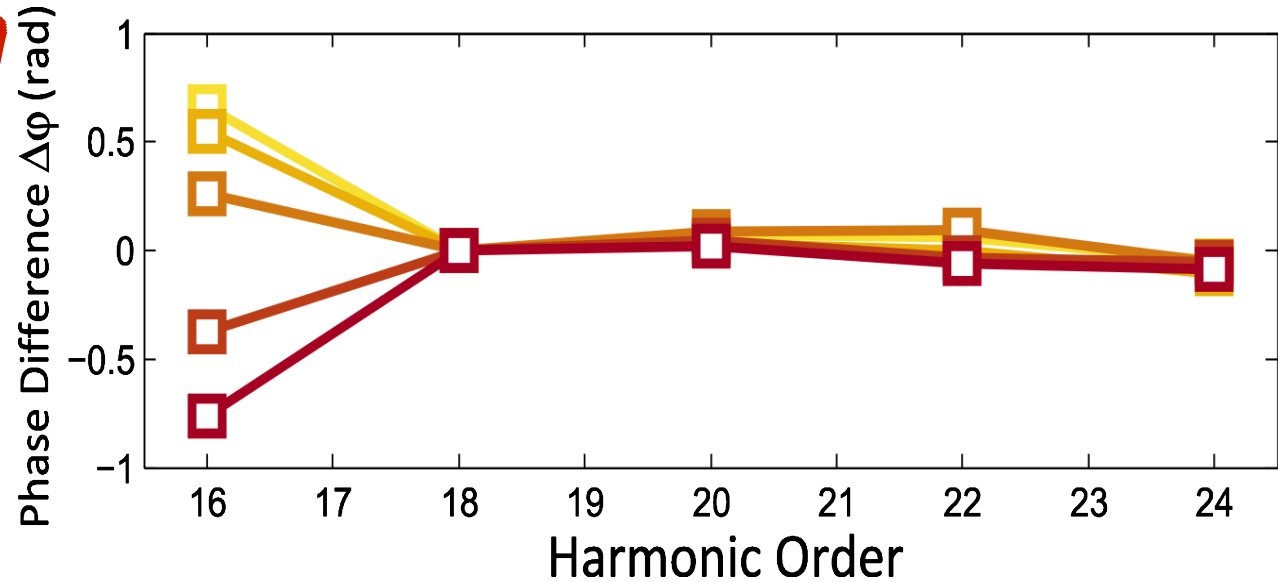
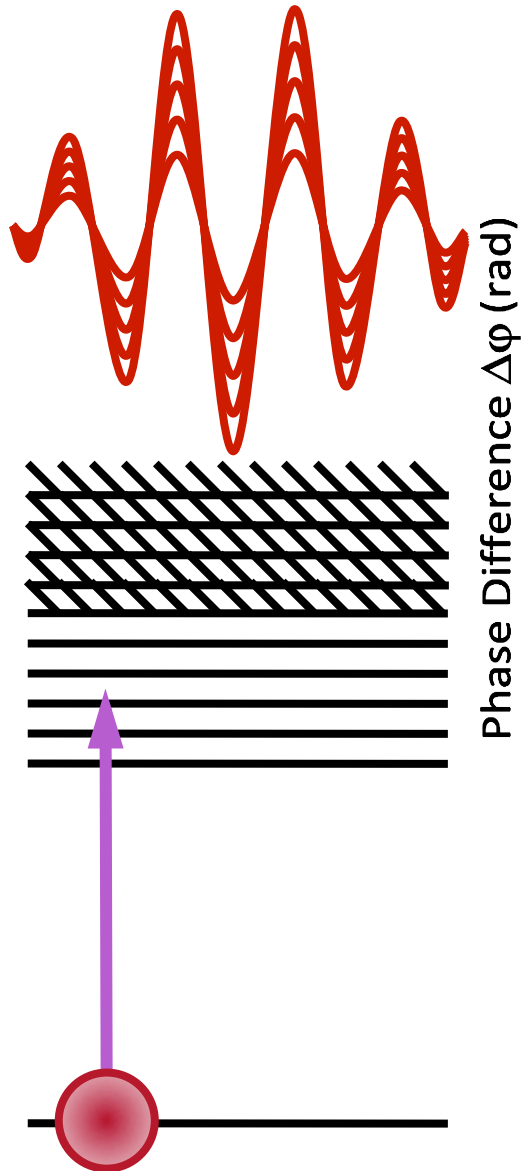


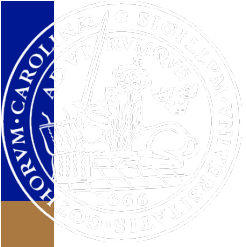
- Perturbation theory model (Gaussian pulses)
- TDSE simulation (@ $1e11$ W/cm²)



Intensity Dependence

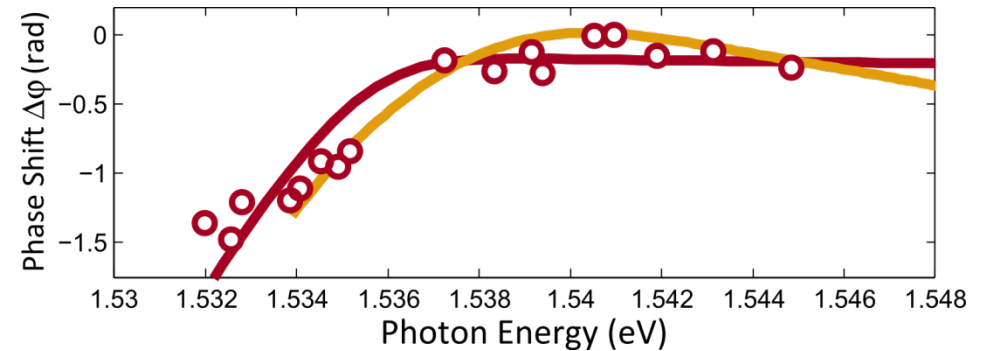
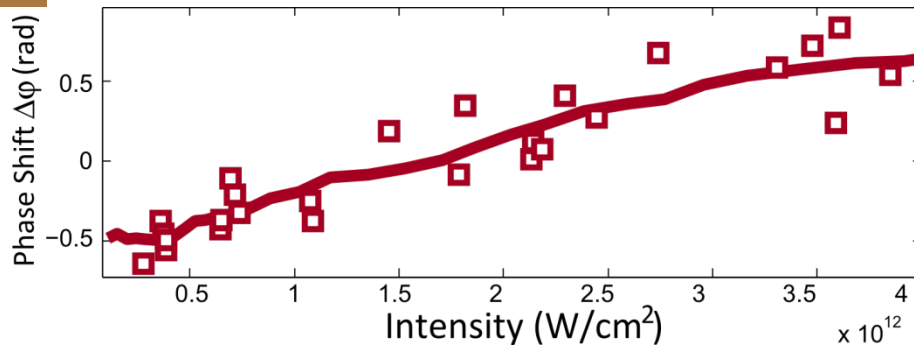
Move the state instead of the photon frequency





Measure the AC Stark Shift

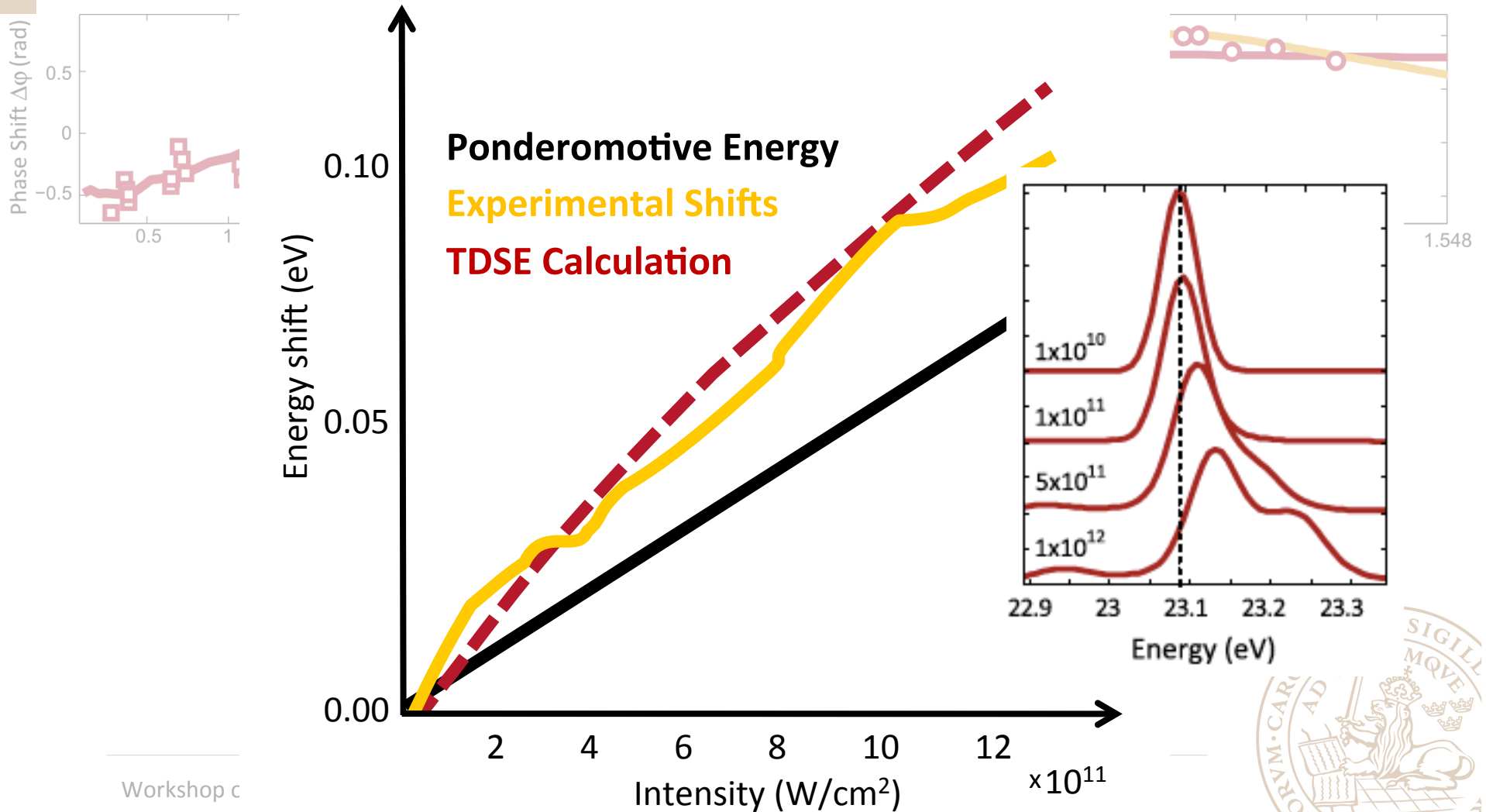
- combine intensity & wavelength dependence of $\Delta\phi$:





Measure the AC Stark Shift

- combine intensity & wavelength dependence of $\Delta\varphi$:



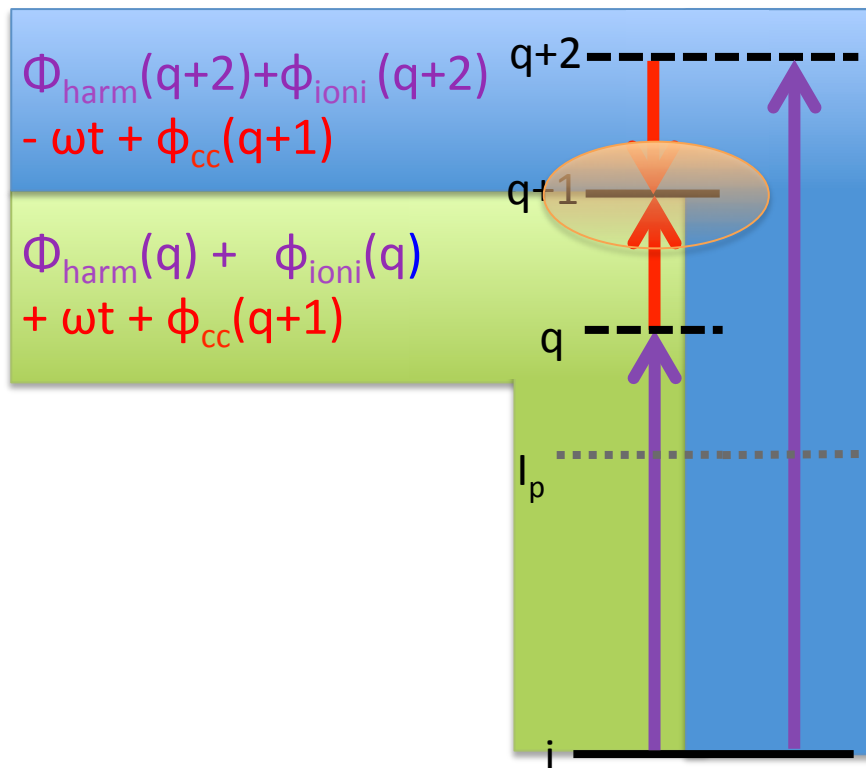


Emission time in photo-ionization



Interferometric measurement

Ionization with two harmonics in the presence of a weak infrared field



Interference between two quantum paths

$$S_{q+1} \sim \cos(2\omega t - \Delta\phi_{\text{harm}} - \Delta\phi_{\text{ioni}})$$

$$\Delta\phi_{\text{ioni}} = \phi_{\text{ioni}}(q+2) - \phi_{\text{ioni}}(q)$$

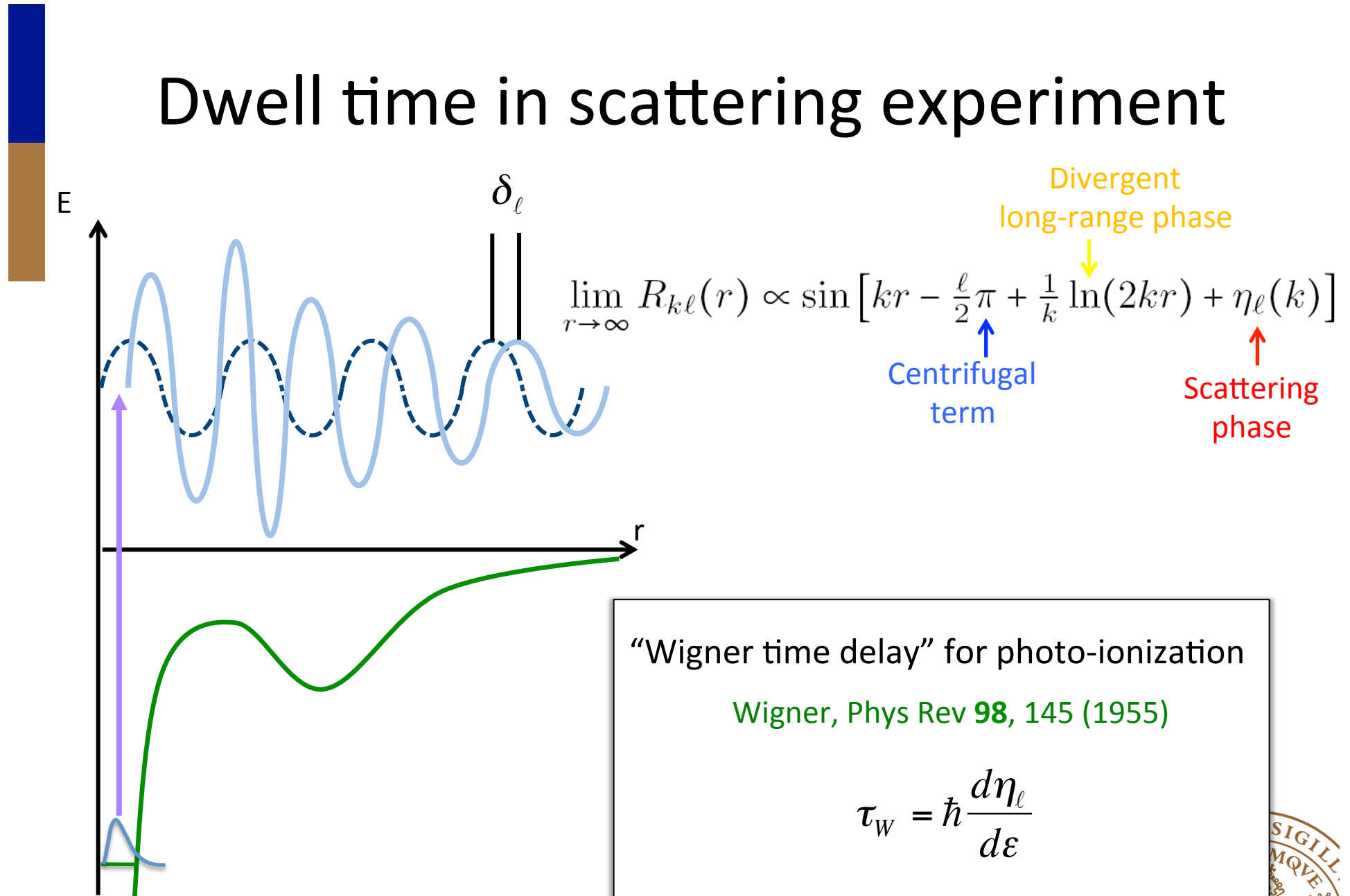
$$\sim 2\omega_0 * d\phi_{\text{ioni}}(q+1)/d\omega$$

$$\sim 2\omega_0 * GD$$

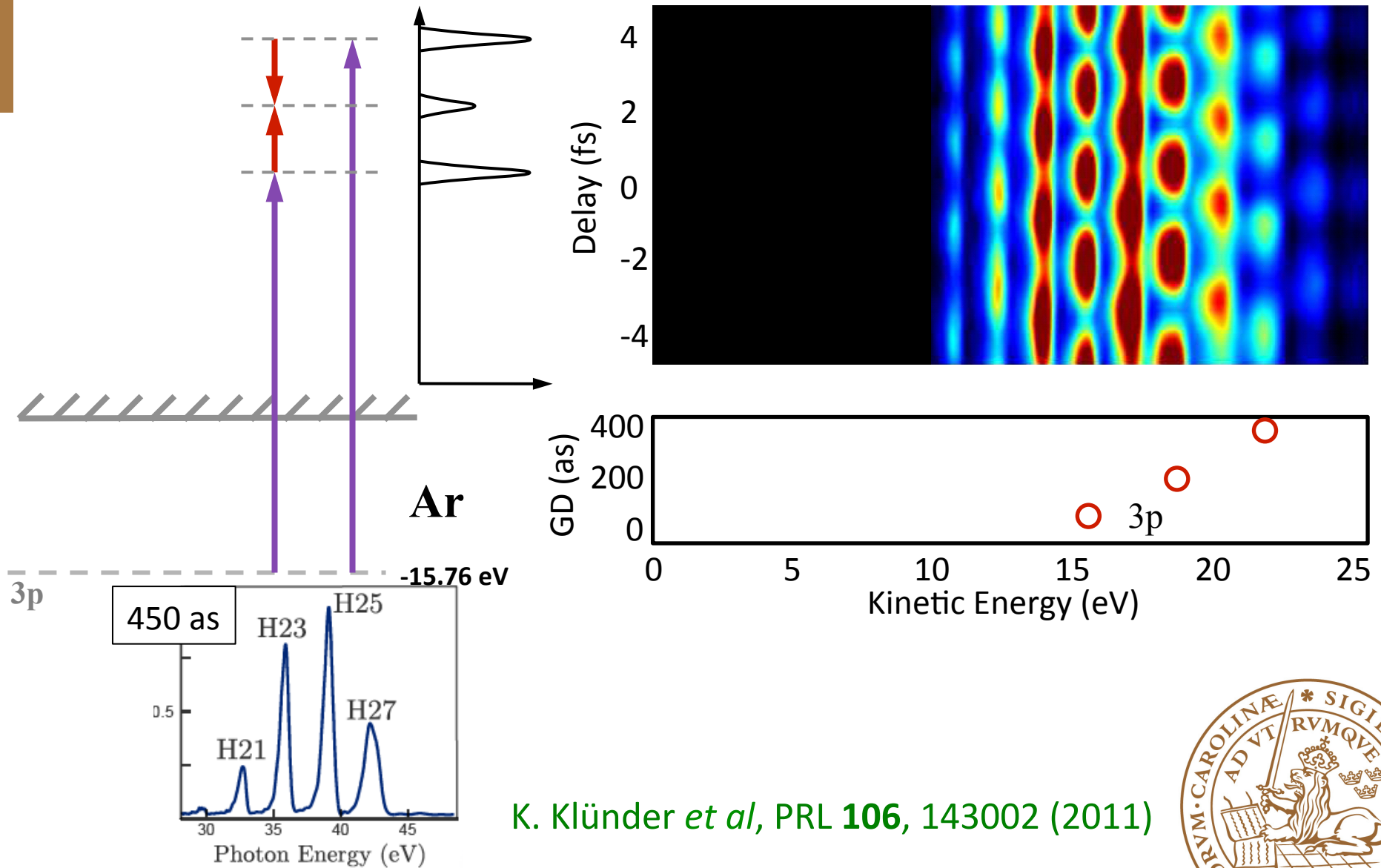
(GD: Group Delay)



Dwell time in scattering experiment



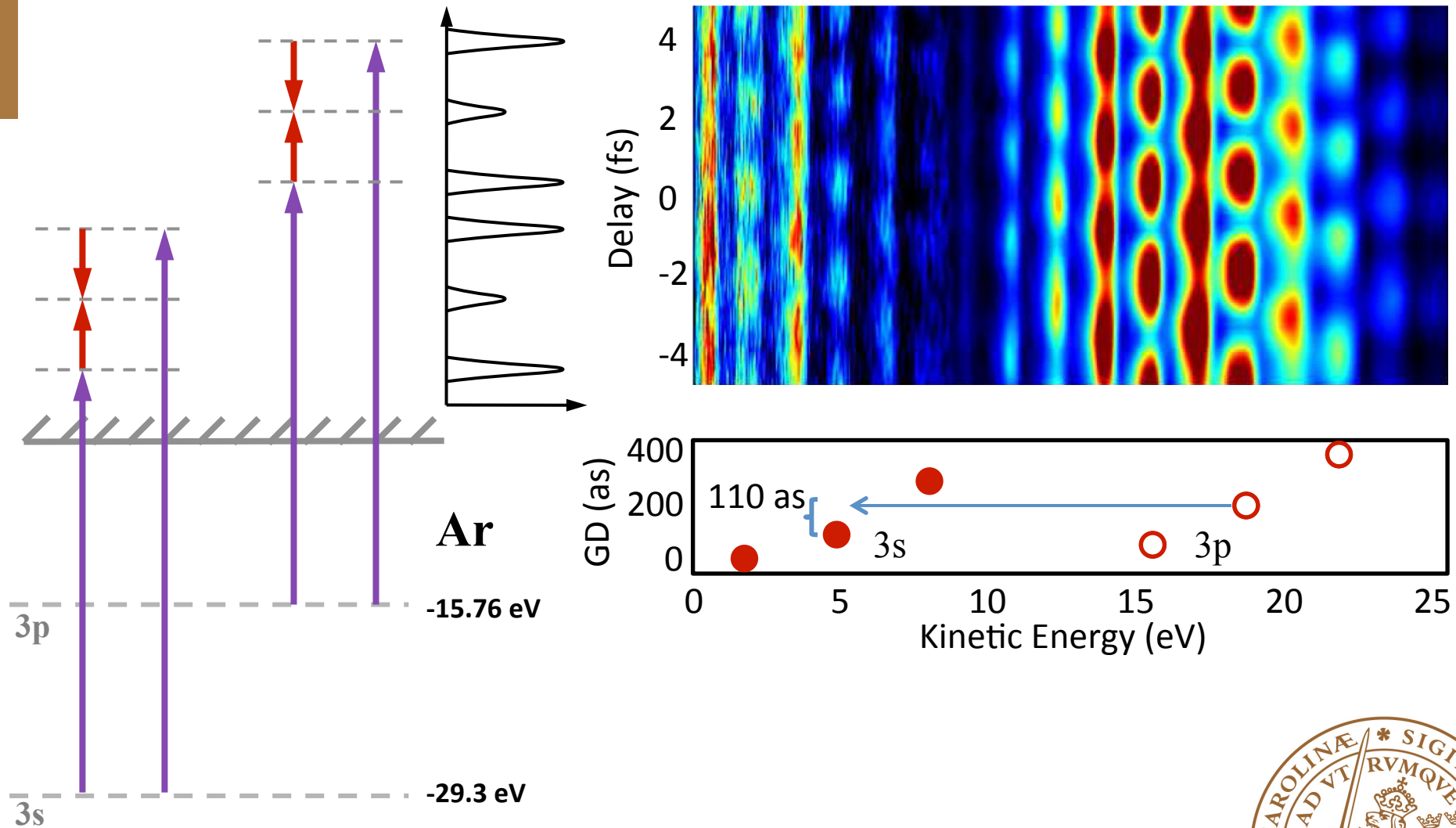
Ionization from the 3s and 3p shells



K. Klünder *et al*, PRL **106**, 143002 (2011)



Ionization from the 3s and 3p shells

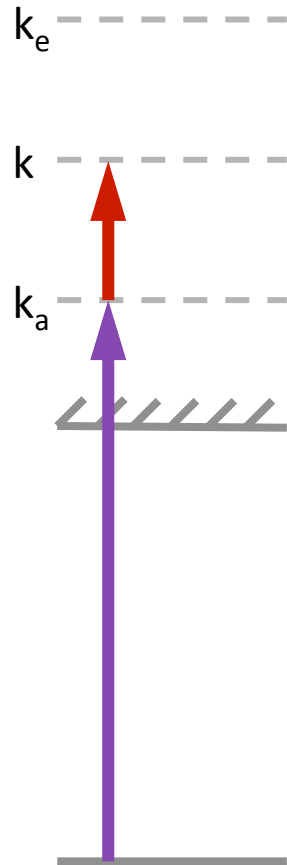


K. Klünder *et al*, PRL **106**, 143002 (2011)



Interpretation

Perturbation theory - Independent particle model



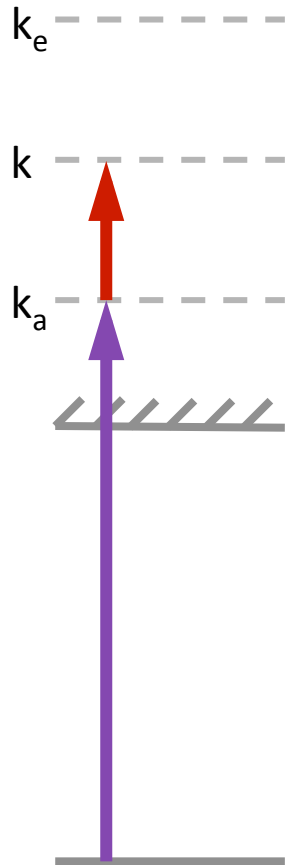
All intermediate states, n , must be considered

$$M_a^{(2)}(\vec{k}) = -iE_L E_H \lim_{\epsilon \rightarrow 0^+} \sum_n \frac{\langle \varphi_{\vec{k}} | \vec{E} \cdot \vec{r} | \varphi_n \rangle \langle \varphi_n | \vec{E} \cdot \vec{r} | \varphi_i \rangle}{\epsilon_i + \omega_H - \epsilon_n + i\epsilon}$$



Interpretation

Perturbation theory - Independent particle model



All intermediate states, n , must be considered

$$M_a^{(2)}(\vec{k}) = -iE_L E_H \lim_{\epsilon \rightarrow 0^+} \sum_n \frac{\langle \varphi_{\vec{k}} | \vec{\epsilon} \cdot \vec{r} | \varphi_n \rangle \langle \varphi_n | \vec{\epsilon} \cdot \vec{r} | \varphi_i \rangle}{\epsilon_i + \omega_H - \epsilon_n + i\epsilon}$$

Partial wave expansion

$$M_a^{(2)}(\vec{k}) = -iE_L E_H \sum_{l=0,2} C_{l0} Y_{l0}(\hat{k}) e^{i\eta_l(k)} T_a^{(2)}(k)$$

Scattering phase
of the final state

Radial part

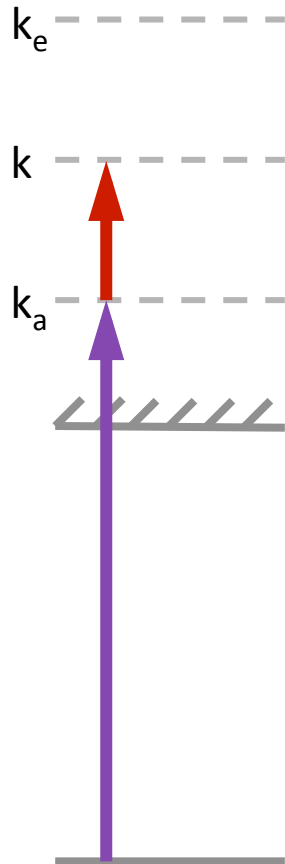
Define the perturbed wavefunction ρ

$$T_a^{(2)}(k) = \sum_n \frac{\langle R_{kl} | r | R_{n1} \rangle \langle R_{n1} | r | R_{i0} \rangle}{\epsilon_i + \omega_H - \epsilon_n + i\epsilon} = \langle R_{kl} | r | \rho_{k_a 1} \rangle$$



Interpretation

Perturbation theory - Independent particle model



$$M_a^{(2)}(\vec{k}) = -iE_L E_H \lim_{\epsilon \rightarrow 0^+} \sum_n \frac{\langle \varphi_{\vec{k}} | \vec{\epsilon} \cdot \vec{r} | \varphi_n \rangle \langle \varphi_n | \vec{\epsilon} \cdot \vec{r} | \varphi_i \rangle}{\epsilon_i + \omega_H - \epsilon_n + i\epsilon}$$



Using asymptotic development
 R_{kl} and ρ_{ka1}

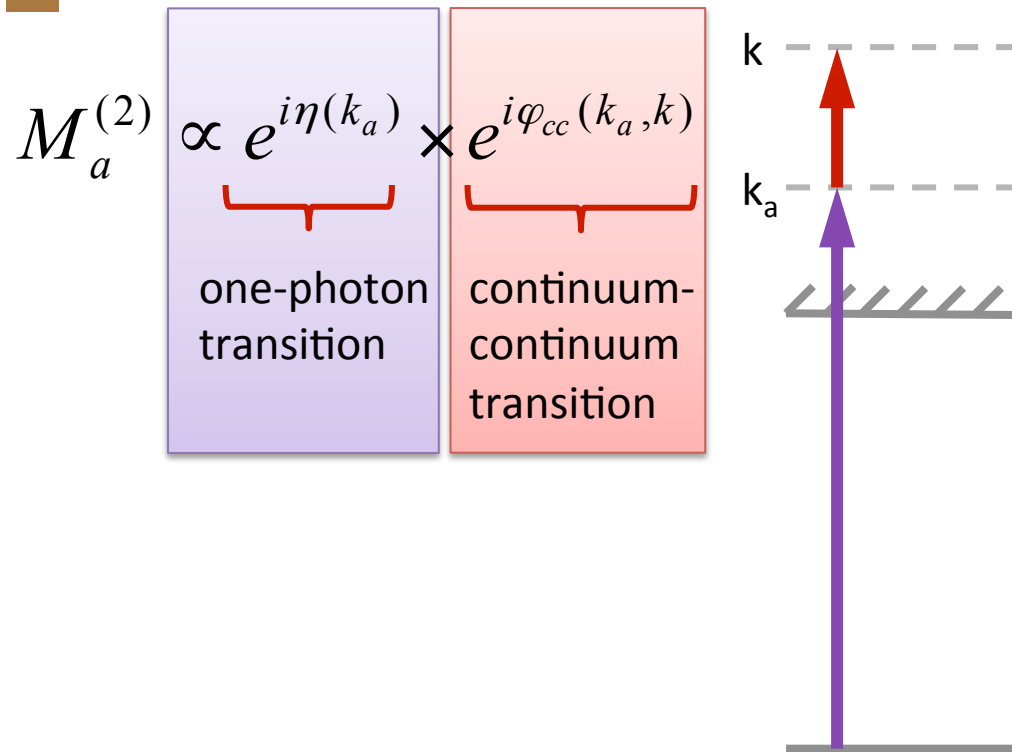
$$M_a^{(2)}(k) \propto \underbrace{e^{i\eta_1(k_a)}}_{(I)} \times \underbrace{\left(\frac{i}{k_a - k} \right)^{iz} \frac{(2k_a)^{\frac{i}{k_a}}}{(2k)^{\frac{i}{k}}} \Gamma(2 + iz)}_{(II)}$$

- (I) Intermediate scattering phase
- (II) Continuum-Continuum phase



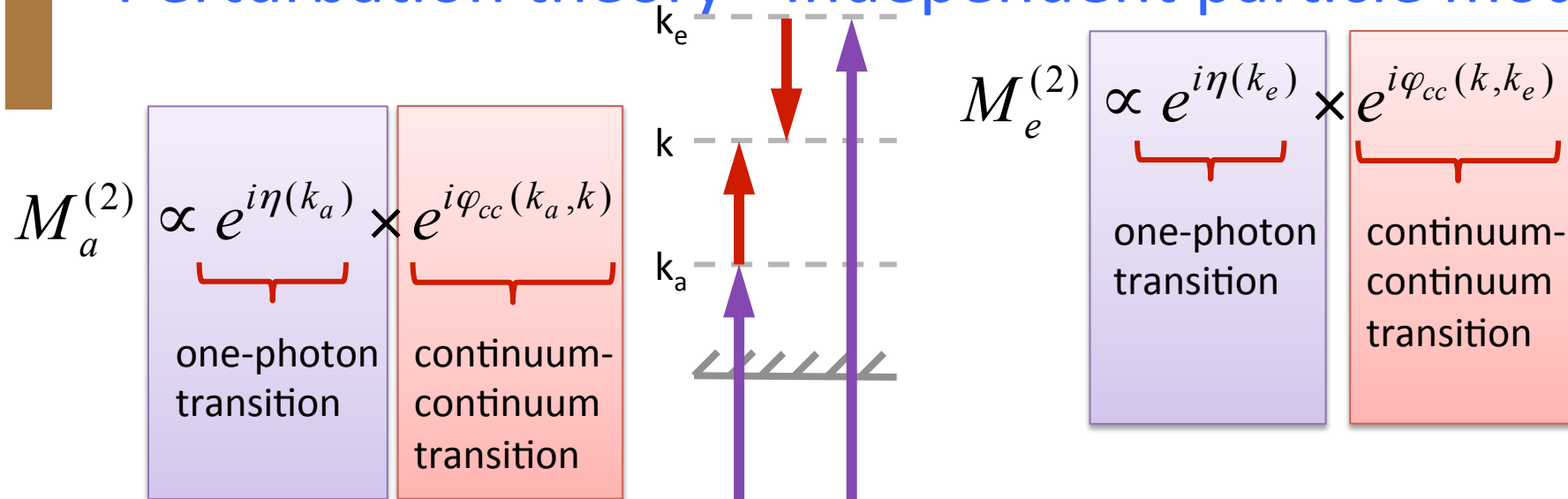
Interpretation

Perturbation theory - Independent particle model



Interpretation

Perturbation theory - Independent particle model

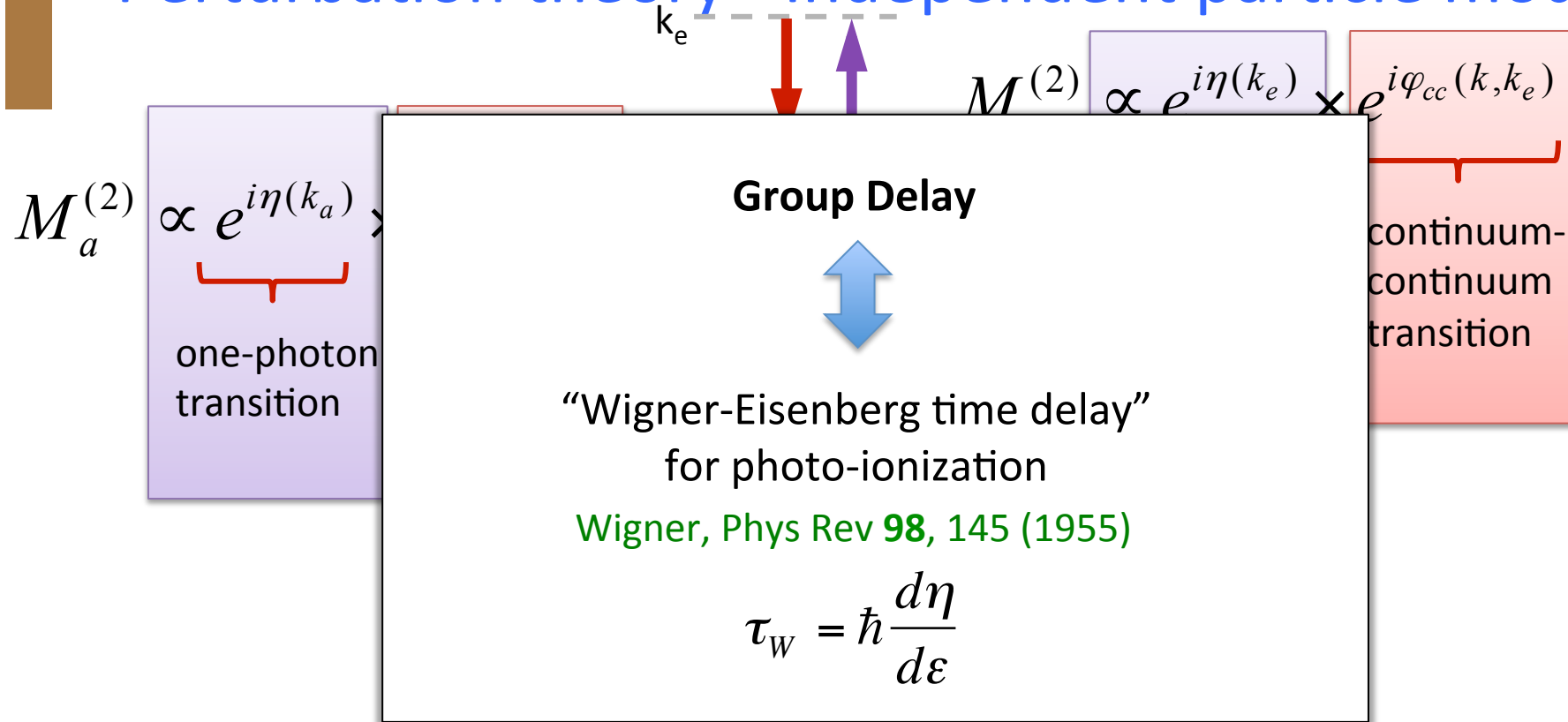


$\Rightarrow GD = \frac{\eta(k_e) - \eta(k_a)}{2\omega} + \frac{\varphi_{cc}(k, k_e) - \varphi_{cc}(k_a, k)}{2\omega}$



Interpretation

Perturbation theory - Independent particle model



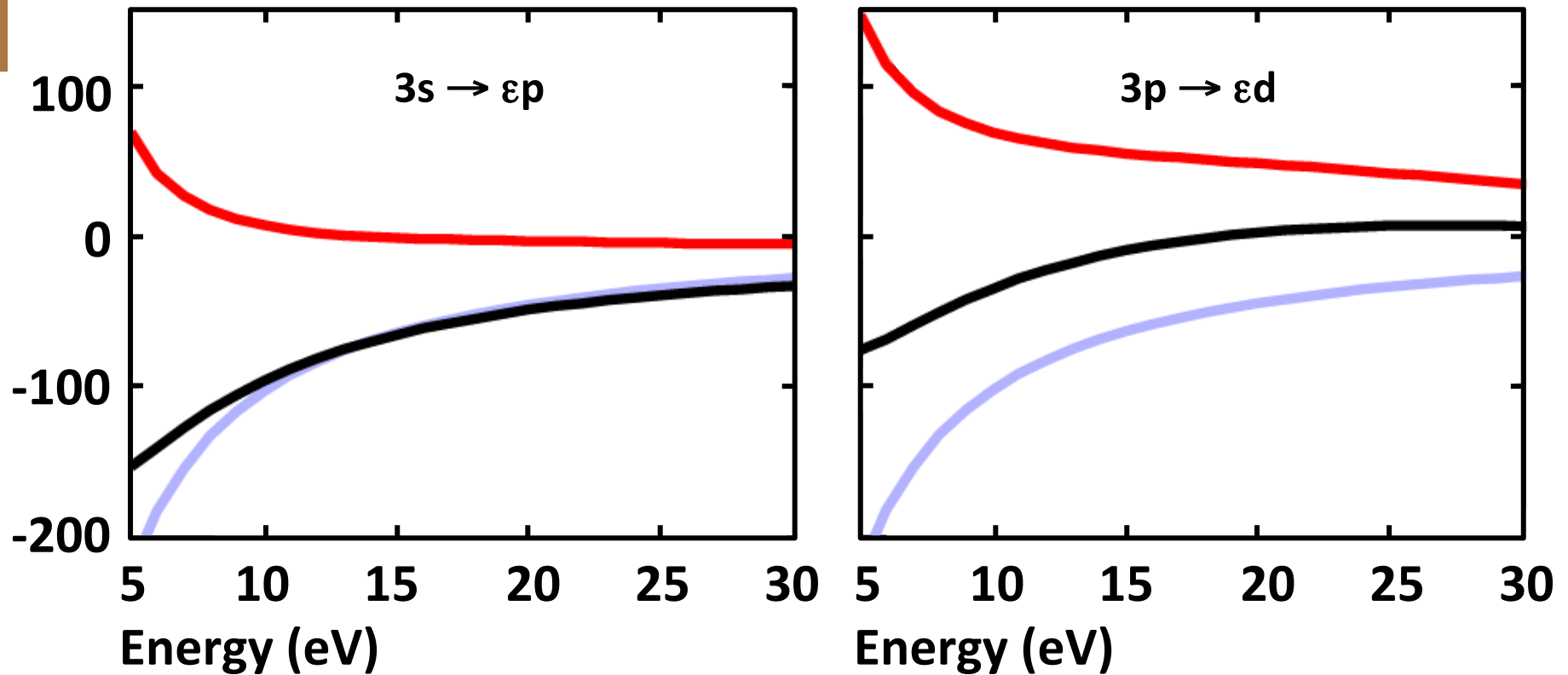
$$GD = \frac{\eta(k_e) - \eta(k_a)}{2\omega} + \frac{\varphi_{cc}(k, k_e) - \varphi_{cc}(k_a, k)}{2\omega}$$

Wigner time delay Continuum-continuum delay



Calculated Group Delay

Independent particle model



HF Wigner time delay (1-photon transition)

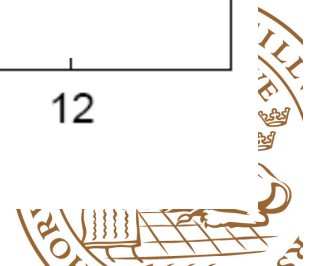
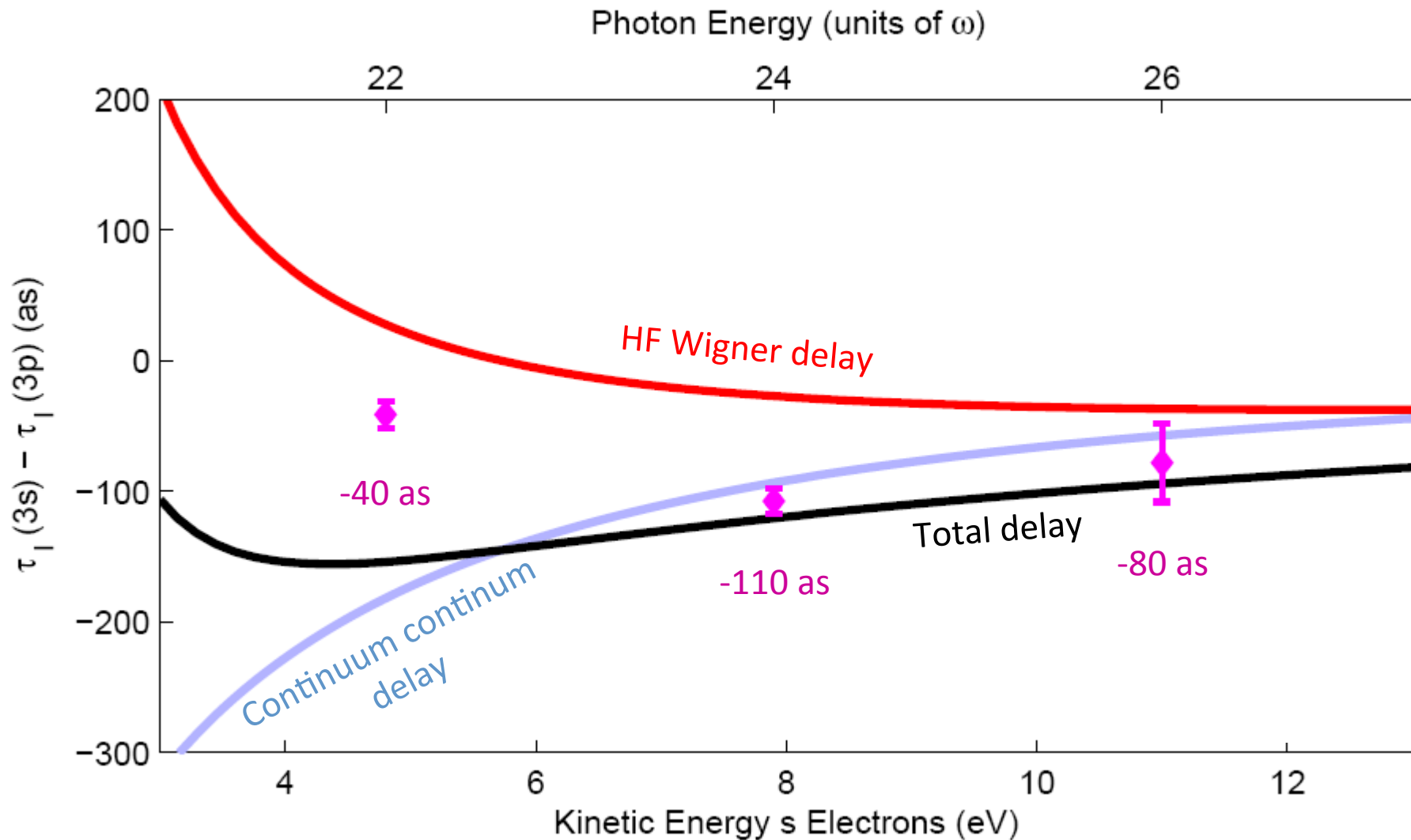
Continuum-continuum delay

Total

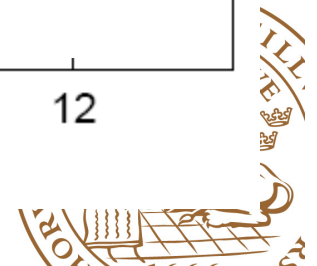
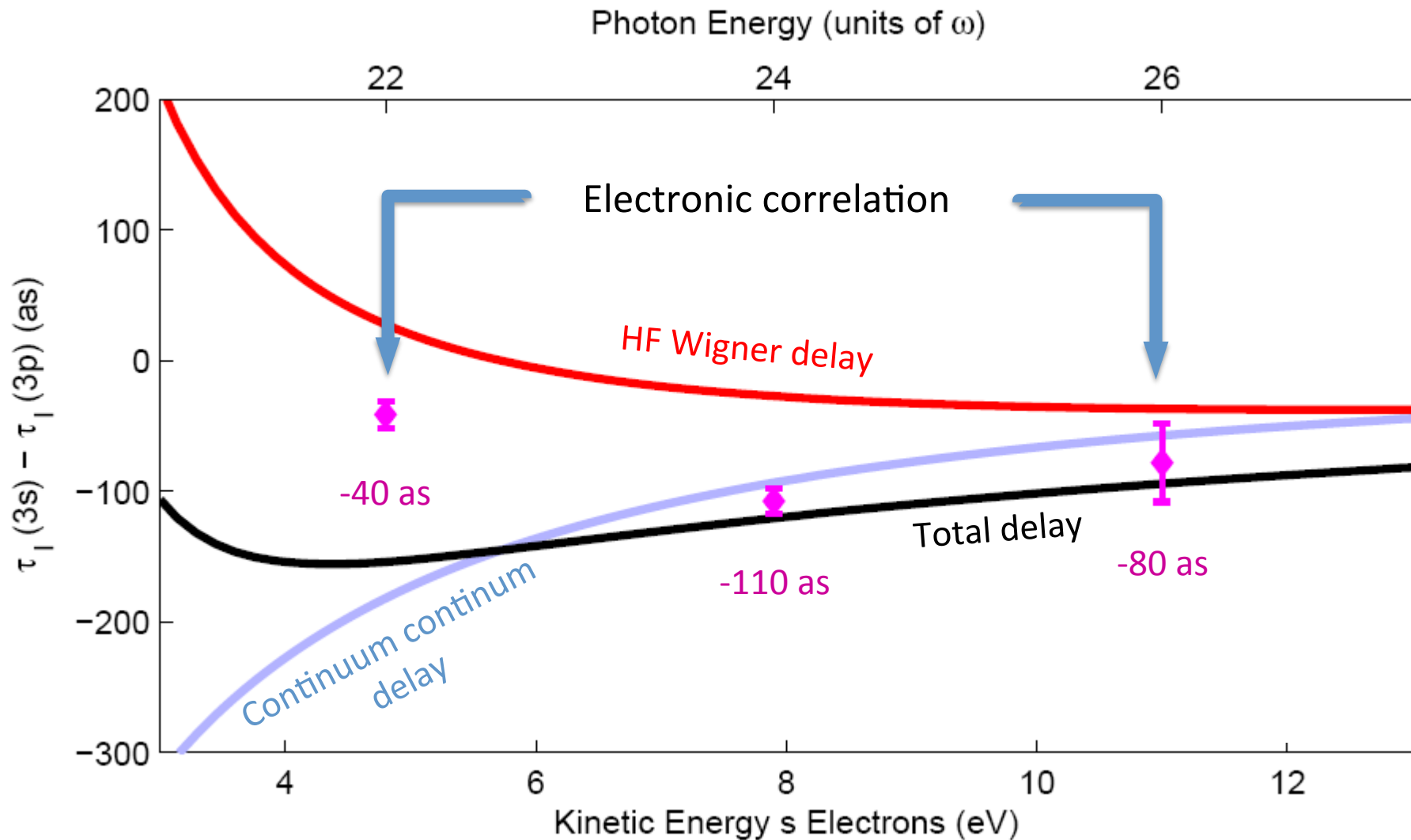
Kennedy and Manson,
Phys. Rev. A 5, 227 (1972)



Comparison with HF

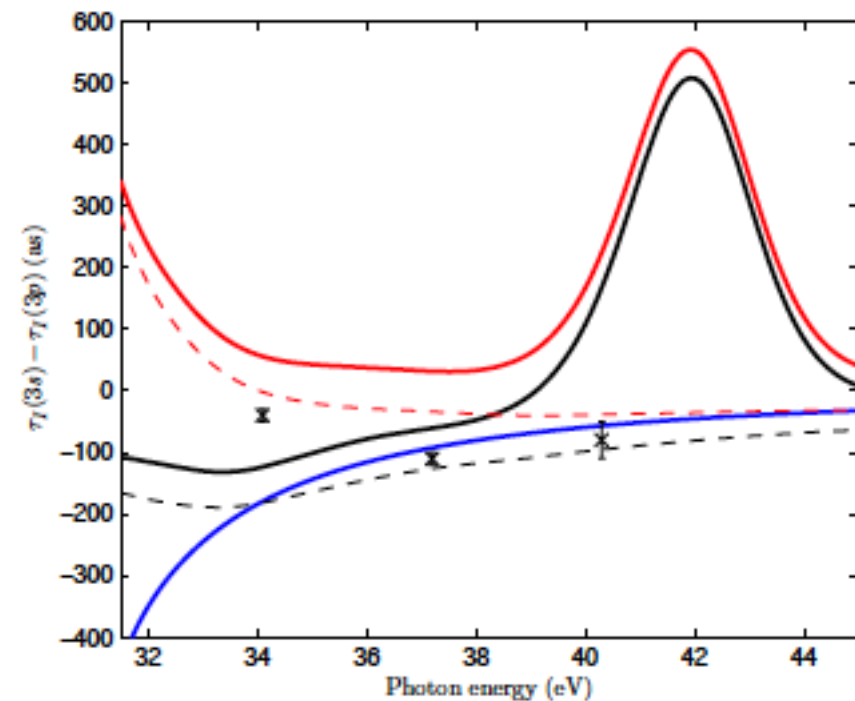
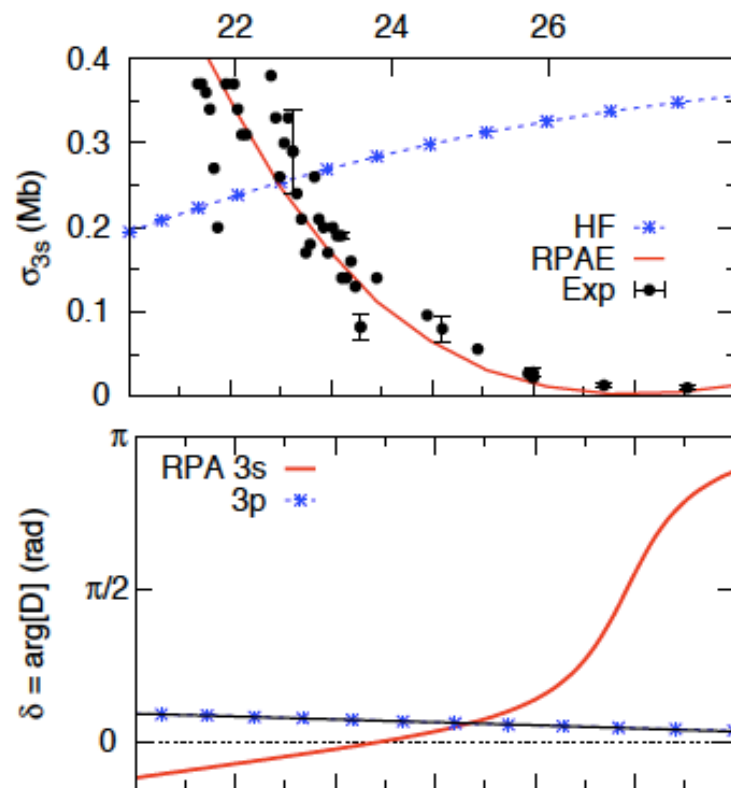


Comparison with HF



Comparison with RPAE

”Cooper minimum” in the 3s channel



With the courtesy of A. Kheifets



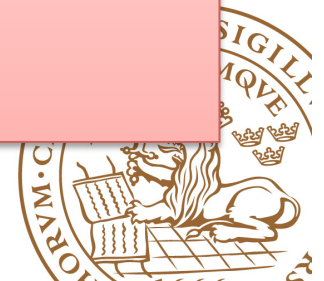
Conclusion

- Interferometric methods enables to read out the Group Delay (GD) of the Electronic Wave Paquet

- The GD encodes information on:
One - photon transition phase
- The measurement introduces a contribution to the GD
Continuum - continuum transition phase



New insight into electronic correlation



Acknowledgement

SWEDEN

M. Dahlström
K. Klünder

D. Guénot
M. Swoboda
T. Fordell
M. Miranda
J. Mauritsson
P. Johnsson
A. L'Huillier

FRANCE

J. Caillat
R. Taïeb
A. Maquet

USA

C. Buth
K.J. Schafer

