



# About new data and BESIII individual determination of time-like form factors

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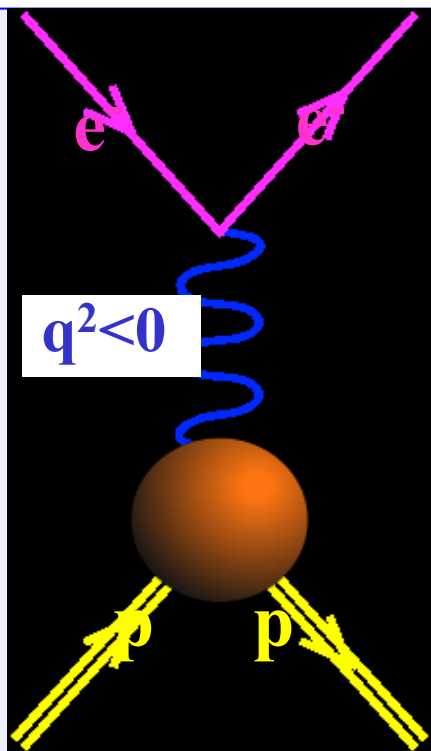
*INFN e Università di Brescia, Italia*

*ArXiv: nucl-th / 2012.145656, to appear in Phys Rev. C*

*PANDA Coll. Meeting, EM Session, March 9, 2021*



# Proton Charge and Magnetic Distributions



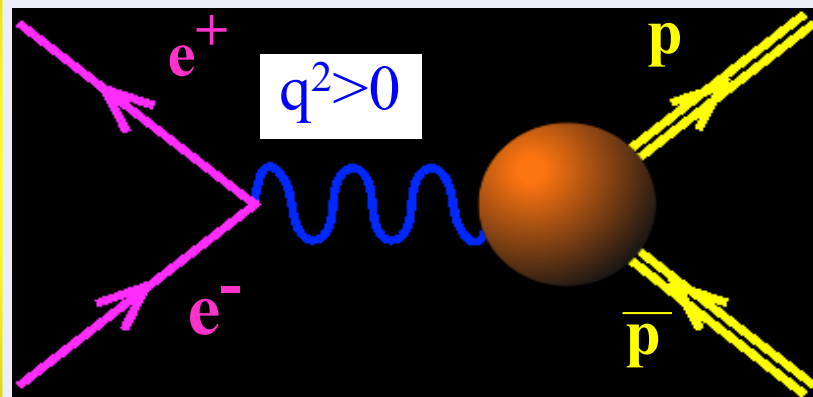
$$G_E(0) = 1$$

$$G_M(0) = \mu_p$$

*Space-like  
FFs are real*

*Asymptotics*

- QCD
- analyticity



*Time-Like  
FFs are complex*

*Unphysical region*  
 $p + \bar{p} \leftrightarrow e^+ e^- + \pi^0$

$$e + p \rightarrow e + p$$

$$0 \quad q^2 = 4m_p^2$$

$G_E = G_M$

$$p + \bar{p} \leftrightarrow e^+ + e^-$$

$q^2$



# CMS Cross section of $e^+e^- \rightarrow p\bar{p}$

$$\frac{d\sigma_{e^+e^- \rightarrow p\bar{p}}(s, \theta)}{d\Omega} = \frac{\alpha^2 \beta \mathcal{C}(\beta)}{4s} \left[ (1 + \cos^2(\theta)) |G_M(s)|^2 + \frac{1}{\tau} \sin^2(\theta) |G_E(s)|^2 \right]$$

$$\tau = s/(4m_p^2)$$

*A. Zichichi, S. Berman, N. Cabibbo, and R. Gatto, Nuovo Cim. 24, 170 (1962).*

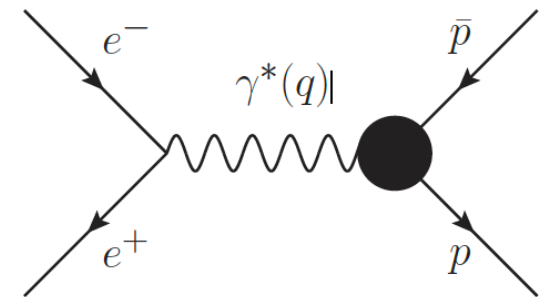
$$\mathcal{C}(\beta) = \frac{y(\beta)}{1 - e^{-y(\beta)}}$$

$$y(\beta) = \frac{\pi\alpha}{\beta} \sqrt{1 - \beta^2}$$

- Coulomb factor
  - *effective within  $< 1$  MeV*
  - *insures finite cross section at threshold*
  - *not present in the time reverse reaction*

$\beta = \sqrt{1 - 1/\tau}$  is the proton velocity

Assumes one-photon exchange :  
 $\cos^2 \theta$  – even

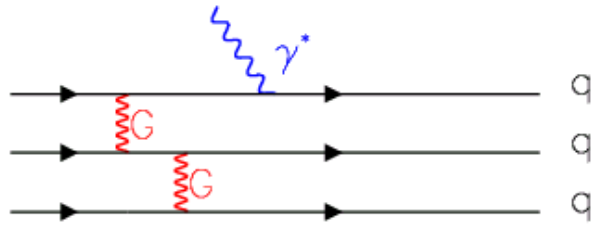


# Oscillations

- Recent and precise data on the proton time-like form factors show a **systematic sinusoidal modulation** in the near-threshold region.
- The relevant variable is **the momentum  $p$**  associated to the **relative motion of the final hadrons**.
- The periodicity and the simple shape of the oscillations point **to a unique interference mechanism**, which occurs when the hadrons are separated by about 1 fm.
- The hadronic matter is distributed in non-trivial way.
- The oscillation period corresponds to **hadronic-scale**
  - **scaling-violating parameter**
  - **origin ?**



# The Time-like Region

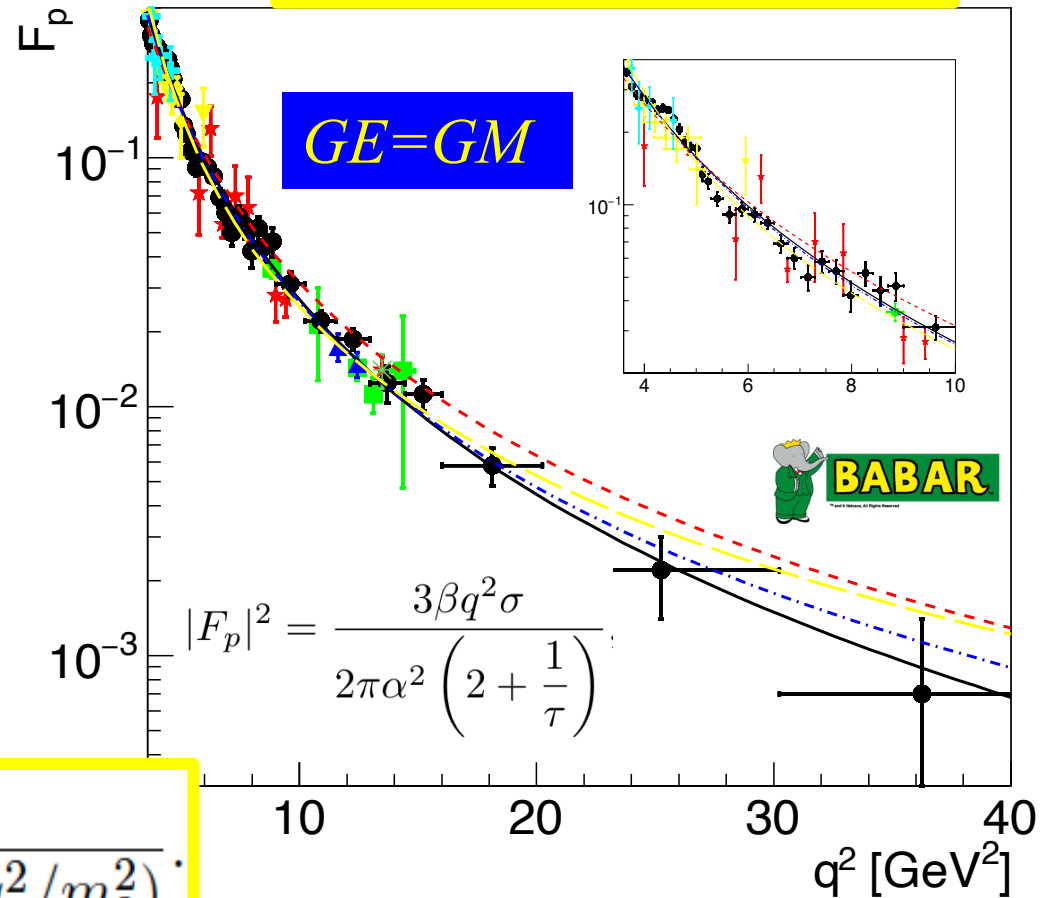
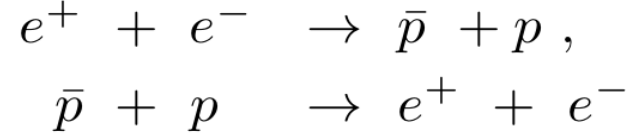


Expected QCD scaling  $(q^2)^2$

$$\frac{A}{(q^2)^2 [\log^2(q^2/\Lambda^2) + \pi^2]}$$

$$\frac{A}{(1 + q^2/m_a^2) [1 - q^2/0.71]^2}$$

$$|F_{T3}(q^2)| = \frac{A}{(1 - q^2/m_1^2)(2 - q^2/m_2^2)}$$

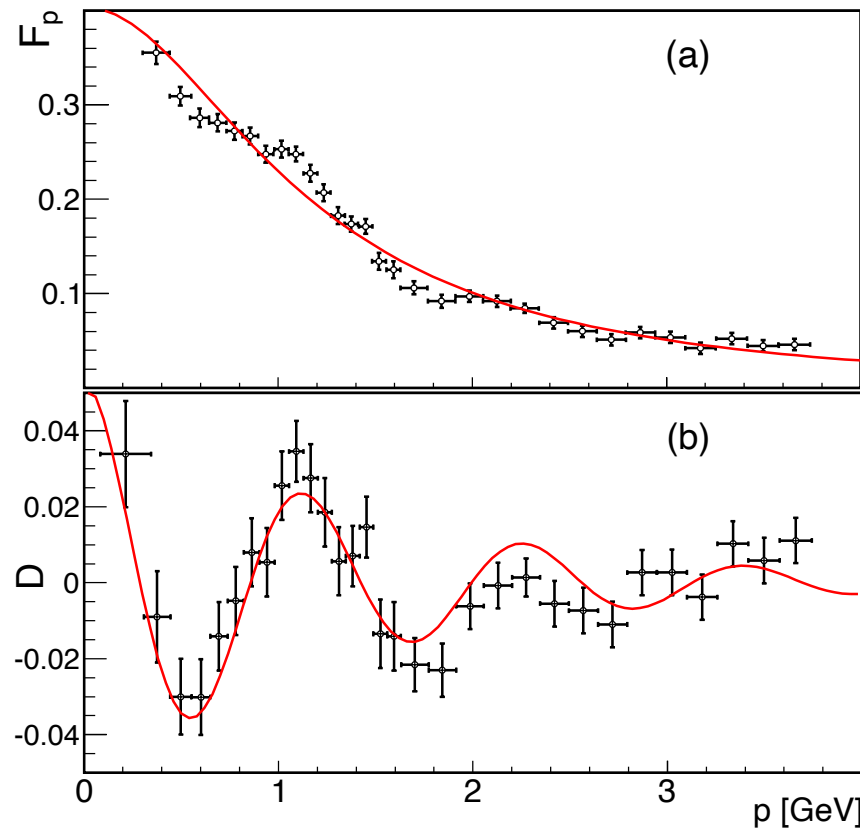


*A. Bianconi, E. T-G. Phys. Rev. Lett. 114,232301 (2015)*



# Oscillations : regular pattern in $P_{Lab}$

The relevant variable is  $p_{Lab}$  associated to the relative motion of the final hadrons.



$$F_{osc}(p) \equiv A \exp(-Bp) \cos(Cp + D).$$

$A \pm \Delta A$	$B \pm \Delta B$	$C \pm \Delta C$	$D \pm \Delta D$	$\chi^2/n.d.f$
	$[GeV]^{-1}$	$[GeV]^{-1}$		
$0.05 \pm 0.01$	$0.7 \pm 0.2$	$5.5 \pm 0.2$	$0.03 \pm 0.3$	1.2

**A:** Small perturbation    **B:** damping  
**C:**  $r < 1\text{fm}$             **D=0:** maximum at  $p=0$

*Simple oscillatory behaviour*  
*Small number of coherent sources*

*A. Bianconi, E. T-G. Phys. Rev. Lett. 114,232301 (2015)*



# Fourier Transform

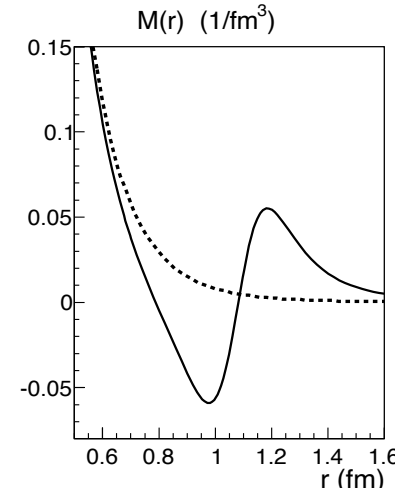
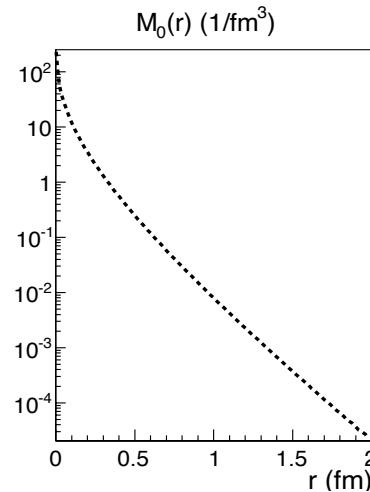
*A. Bianconi, E. T-G.,  
Phys. Rev. Lett. 114, 232301 (2015)*

$$F_0(p) \equiv \int d^3\vec{r} \exp(i\vec{p} \cdot \vec{r}) M_0(r)$$

$$F(p) = F_0(p) + F_{osc}(p) \equiv \int d^3\vec{r} \exp(i\vec{p} \cdot \vec{r}) M(r).$$

$$F_0 = \frac{A}{(1 + q^2/m_a^2) [1 - q^2/0.71]^2},$$

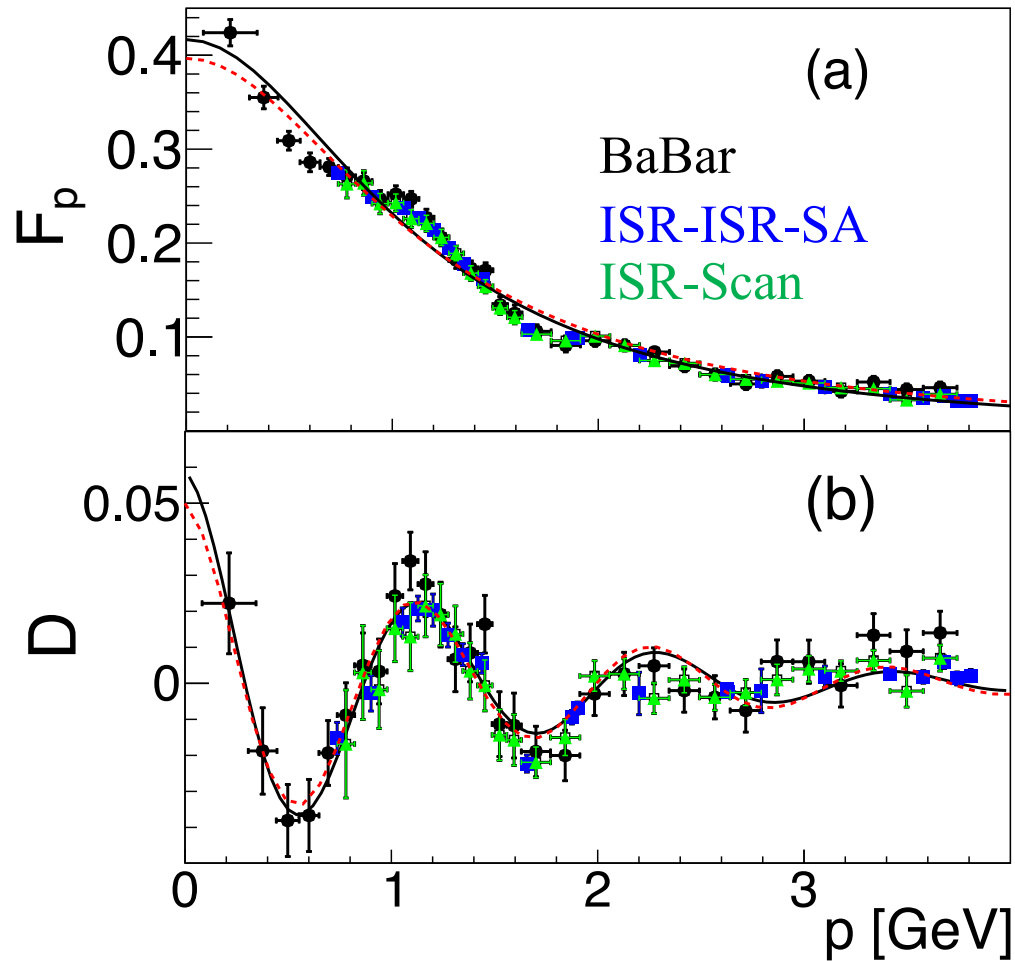
$$F_{osc}(p) \equiv A \exp(-Bp) \cos(Cp + D).$$



- *Rescattering processes*
- *Large imaginary part*
- *Related to the time evolution of the charge density?*  
(E.A. Kuraev, E. T-G., A. Dbeyssi, PLB712 (2012) 240)
- *Consequences for the SL region?*
- *Data from BESIII, expected from PANDA*



# Confirmation of regular oscillations



$$F_p^{\text{fit}}(s) = F_{3p}(s) + F_{\text{osc}}(p(s))$$

$$F_{3p}(s) = \frac{F_0}{\left(1 + \frac{s}{m_a^2}\right) \left(1 - \frac{s}{m_0^2}\right)^2},$$

$$F_{\text{osc}}(p(s)) = Ae^{-Bp} \cos(Cp + D).$$

$$s = 2m_p \left( m_p + \sqrt{p^2 + m_p^2} \right),$$

$$p = \sqrt{s \left( \frac{s}{4m_p^2} - 1 \right)}.$$





# Monopole Fit

Ref.	Exp.	N	$F_0$	$m_a^2$ (GeV <sup>2</sup> )
[3, 5, 12]	BaBar	85	$7.7 \pm 0.3$	$15 \pm 1$
[3–6]	BaBar, BESIII-ISR, BESIII-SC	107	$8.9 \pm 0.2$	$8.8 \pm 0.6$

# Damped Oscillatory Function

Ref.	Data set	$A \pm \Delta A$	$B \pm \Delta B$ (GeV <sup>-1</sup> )	$C \pm \Delta C$ (GeV <sup>-1</sup> )	$D \pm \Delta D$	$\chi^2/n.d.f$
[3, 5, 12]	BaBar	$0.05 \pm 0.01$	$0.59 \pm 0.2$	$5.6 \pm 0.1$	$0.2 \pm 0.2$	$57/(55-4) = 1.1$
[3–6]	BESIII-ISR, SC, BaBar	$0.07 \pm 0.01$	$0.93 \pm 0.09$	$5.9 \pm 0.1$	$0.1 \pm 0.2$	$227/(107-4) = 2.2$

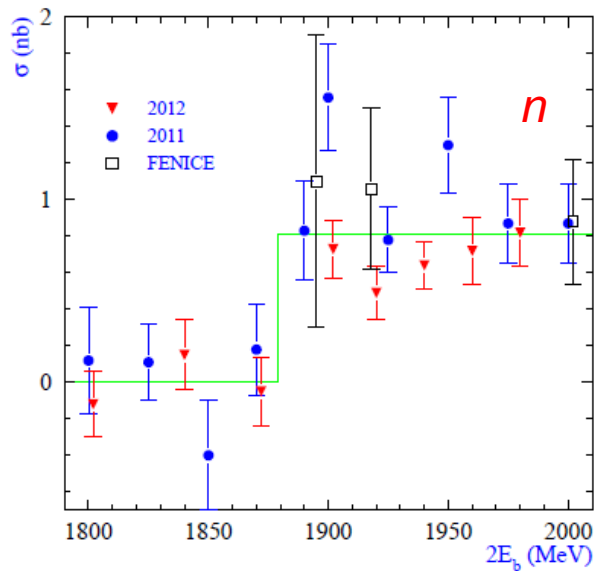
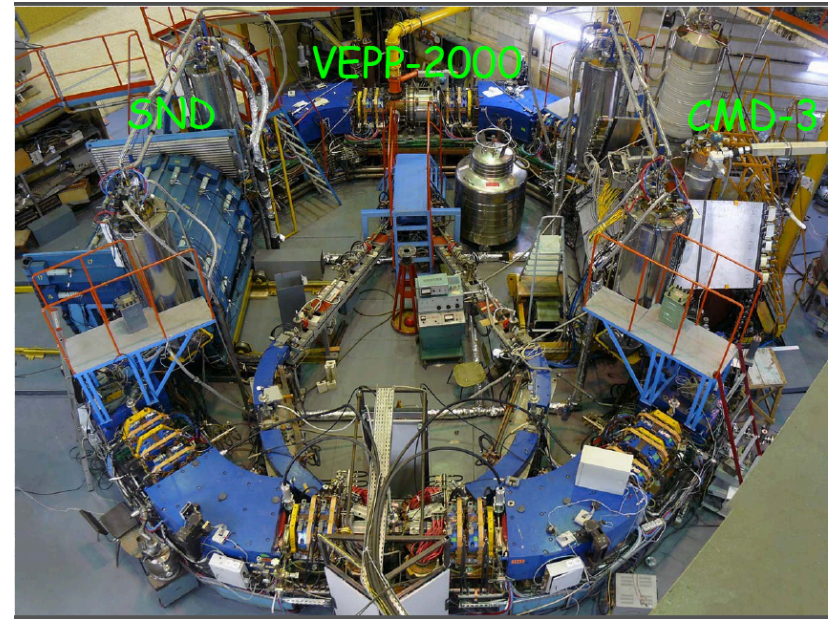
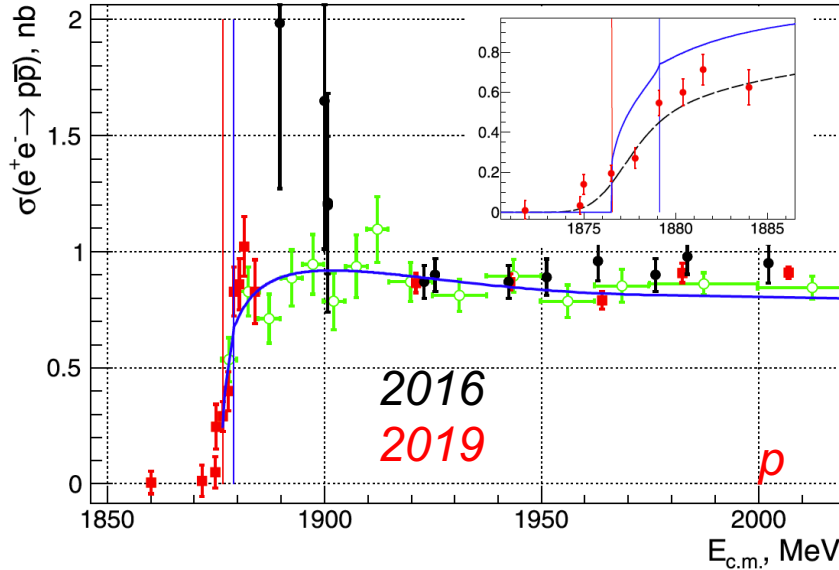
# 6-parameter Fit

Ref.	$F_0$	$m_a^2$ (GeV <sup>2</sup> )	$A$	$B$ (GeV <sup>-1</sup> )	$C$ (GeV <sup>-1</sup> )	$D$	$\frac{\chi^2}{n.d.f.}$
[3–6, 11]	$9.7 \pm 0.3$	$7.1 \pm 0.5$	$0.073 \pm 0.007$	$1.05 \pm 0.07$	$5.51 \pm 0.09$	$0.04 \pm 0.1$	$\frac{278}{118 - 6} = 2.5$



# Threshold physics

VEPPII, Novosibirsk



- The beam energy is measured at 0.1 MeV precision (back scattering laser light system).
- The energy spread due to radiation and energy resolution is small enough to differentiate the proton and neutron thresholds.

$$\sigma_{\text{Born}}(E_{\text{c.m.}}) = A + B \left[ 1 - \exp\left(-\frac{(E_{\text{c.m.}} - E_{\text{thr}})}{\sigma_{\text{thr}}}\right) \right],$$

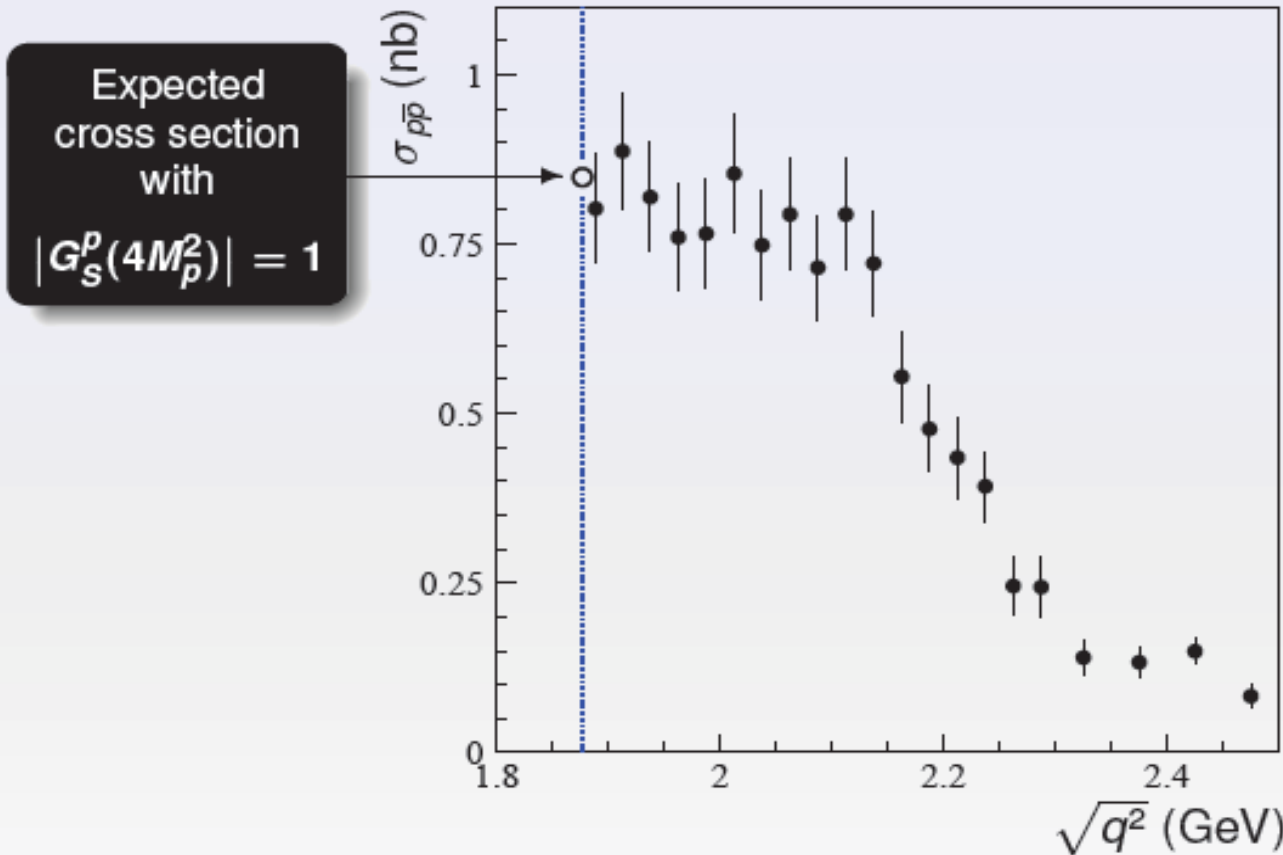


# Point-like form factors?

**BABAR:  $e^+e^- \rightarrow p\bar{p}$**

EPJA39, 315

*S. Pacetti*  
*EPJA 39, 315 (2009)*



At the threshold

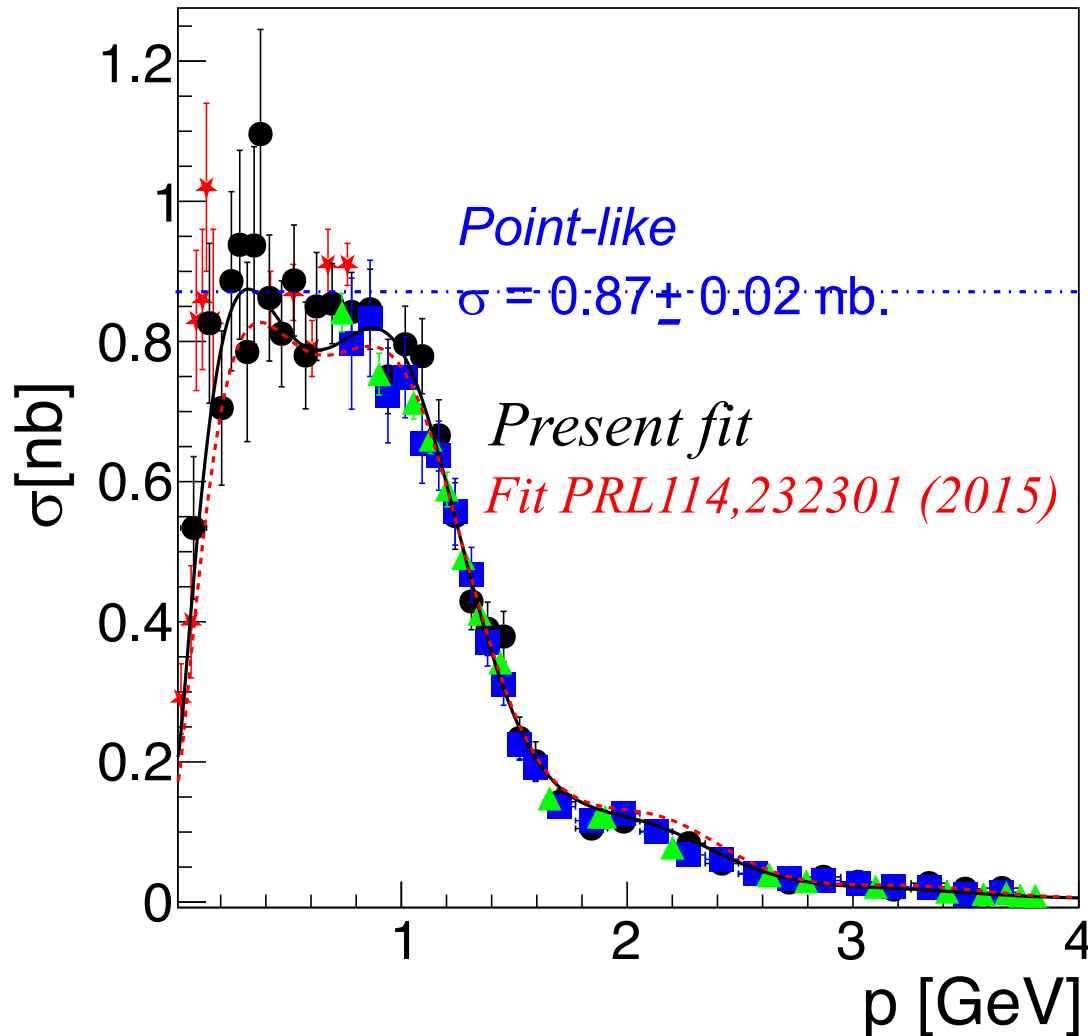
$$\sigma_{p\bar{p}}(4M_p^2) = \frac{\pi^2 \alpha^3}{2M_p^2} \frac{\cancel{\beta_p}}{\cancel{\beta_p}} |G_S^p(4M_p^2)|^2$$

$$\sigma_{p\bar{p}}(4M_p^2) = 0.85 |G_S^p(4M_p^2)|^2 \text{ nb}$$



$|G_S^p(4M_p^2)| \equiv 1$   
as pointlike fermion pairs!

# Cross section from $e^+e^- \rightarrow p\bar{p}$



Novosibirsk 38pt  
 $1.9 < 2E < 4.5$   
*PLB794,64 (2019)*

BaBar 85pt  
 $1.9 < 2E < 4.5$   
*PRD87,092005 (2013)*

ISR-ISR-SA 30pt  
 $2 < 2E < 3.6$   
*PRD99,092002 (2019)*

ISR-Scan 22pt  
 $2 < 2E < 3.1$   
*PRL124,042001 (2020)*



# Total Cross Section from $e^+e^- \rightarrow \bar{p}p$

$$\sigma_{e^+e^- \rightarrow \bar{p}p}(s) = \frac{4\pi\alpha^2\beta\mathcal{C}(\beta)}{3s} \left( |G_M(s)|^2 + \frac{1}{2\tau} |G_E(s)|^2 \right)$$

- Effective FF:  $\sigma_{\text{Tot}} \sim F_p^2$

$$F_p(s)^2 = \frac{2\tau |G_M(s)|^2 + |G_E(s)|^2}{2\tau + 1}$$

- Equivalent to:

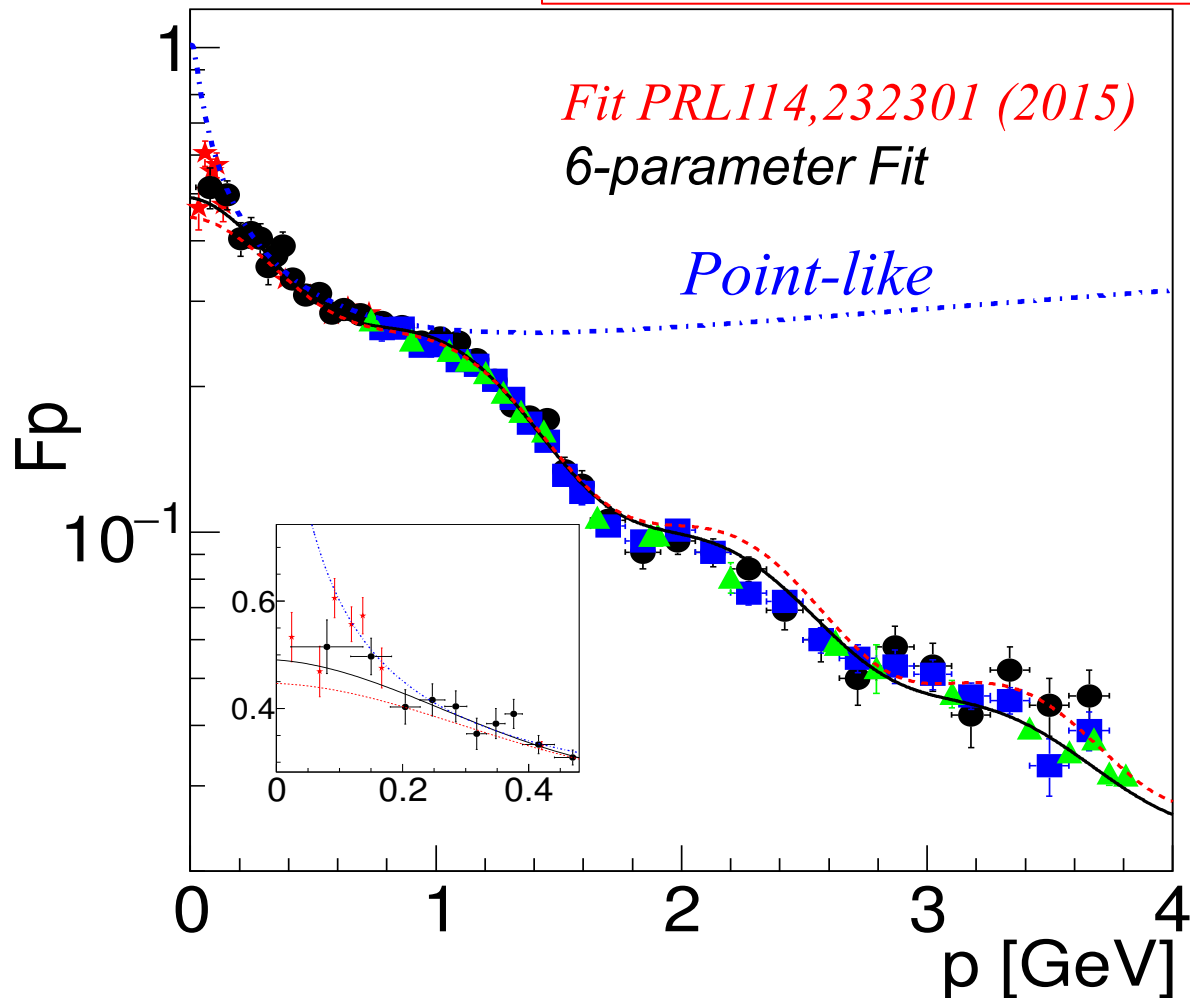
$$|G_E(s)| = |G_M(s)| \equiv F_p(s)$$

*Strictly valid at threshold, where only one amplitude is present*



# Generalized Form Factor

$$F_p(s)^2 = \frac{2\tau |G_M(s)|^2 + |G_E(s)|^2}{2\tau + 1}$$

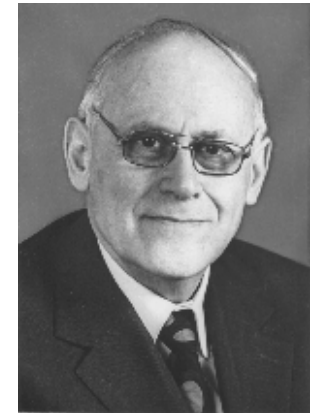
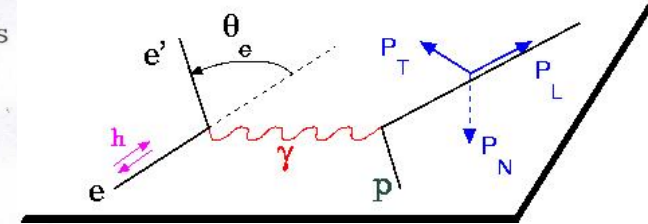


## POLARIZATION PHENOMENA IN ELECTRON SCATTERING BY PROTONS IN THE HIGH-ENERGY REGION

Academician A. I. Akhiezer\* and M. P. Rekaló

Physicotechnical Institute, Academy of Sciences of the Ukrainian SSR  
Translated from Doklady Akademii Nauk SSSR, Vol. 180, No. 5,  
pp. 1081-1083, June, 1968  
Original article submitted February 26, 1967

M.P. Rekaló  
(1938-2004)



A.I. Akhiezer  
(1911-2000)

$$s_2 \frac{d\sigma}{d\Omega_R} = 4p_2 \frac{(\mathbf{s} \cdot \mathbf{q})}{1 + \tau} \Gamma(\theta, \epsilon_1) \left[ \tau G_M (G_M + G_E) - \frac{1}{4\epsilon_1} G_M (G_E - \tau G_M) \right],$$

The polarization induces a term in the cross section proportional to  $G_E G_M$   
*Polarized beam and target or polarized beam and recoil proton polarization*

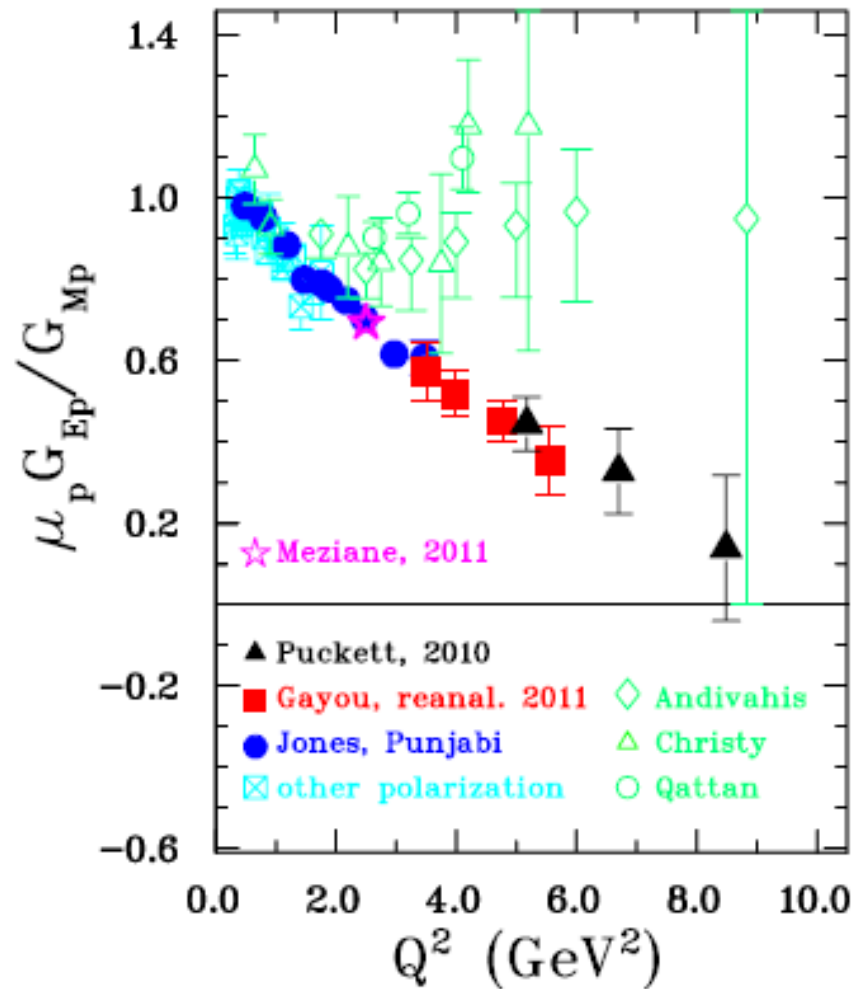


# Polarization experiments @ JLab

*A.I. Akhiezer and M.P. Rekalo, 1967*

## GEp collaboration

- 1) "standard" **dipole function** for the nucleon magnetic FFs  $G_{Mp}$  and  $G_{Mn}$
- 2) **linear deviation** from the dipole function for the electric proton FF  $G_{Ep}$
- 3) **QCD scaling** not reached
- 3) **Zero crossing** of  $G_{Ep}$ ?
- 4) **contradiction between polarized and unpolarized measurements**



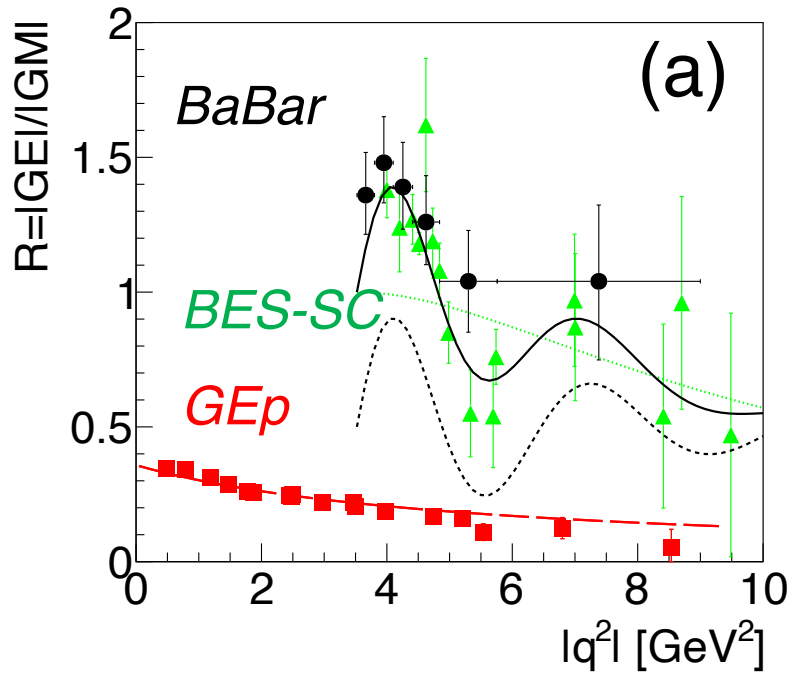
*JLab-GEp Collaboration*

*J.R. Puckett et al, PRC 96, 055203 (2017)*





# Form Factor Ratio $R=|GE|/|GM|$

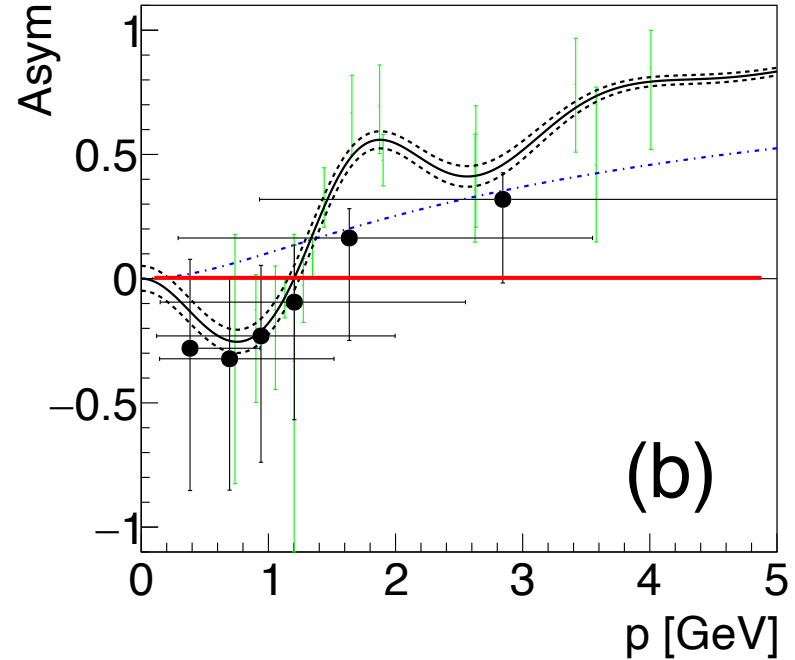
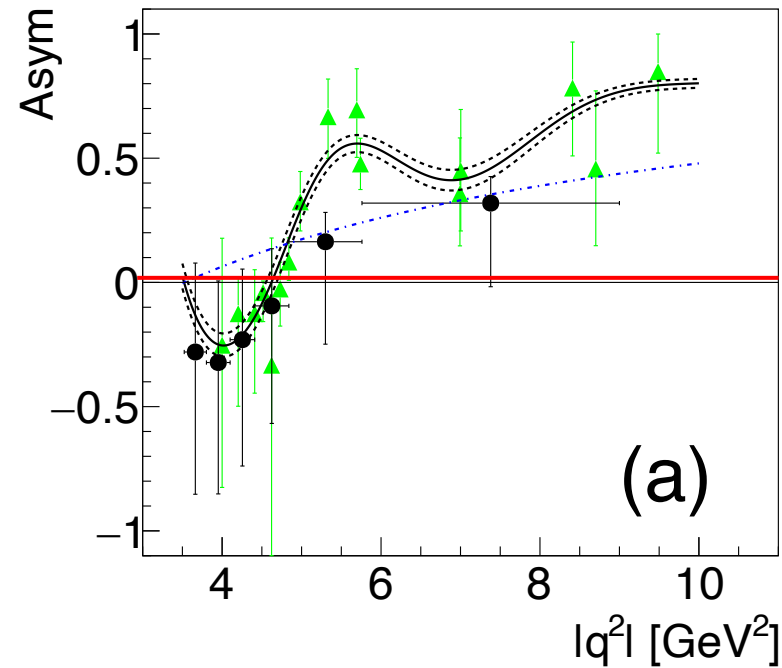


- Precise data from BESIII
- Dip at  $|q^2| \sim 5.8$  GeV<sup>2</sup>
- Comparison with SL (Jlab-GEp data)
- Oscillations on top of a monopole: from GE or GM?

$$F_R(\omega(s)) = \frac{1}{1 + \omega^2/r_0} [1 + r_1 e^{-r_2 \omega} \sin(r_3 \omega)], \quad \omega = \sqrt{s} - 2m_p,$$



# Angular Asymmetry



$$\frac{d\sigma_{e^+e^- \rightarrow \bar{p}p}}{d\Omega}(s, \theta) = \sigma_0(s) |1 + \mathcal{A}(s) \cos^2(\theta)|$$

$$\sigma_0(s) = \frac{\alpha^2 \beta \mathcal{C}(\beta)}{4s} \left( |G_M(s)|^2 + \frac{1}{\tau} |G_E(s)|^2 \right)$$

$$\mathcal{A}(s) = \frac{\tau |G_M(s)|^2 - |G_E(s)|^2}{\tau |G_M(s)|^2 + |G_E(s)|^2} = \frac{\tau - R(s)^2}{\tau + R(s)^2}$$

$$q^2 = (4.60 \pm 0.07) \text{ GeV}^2$$

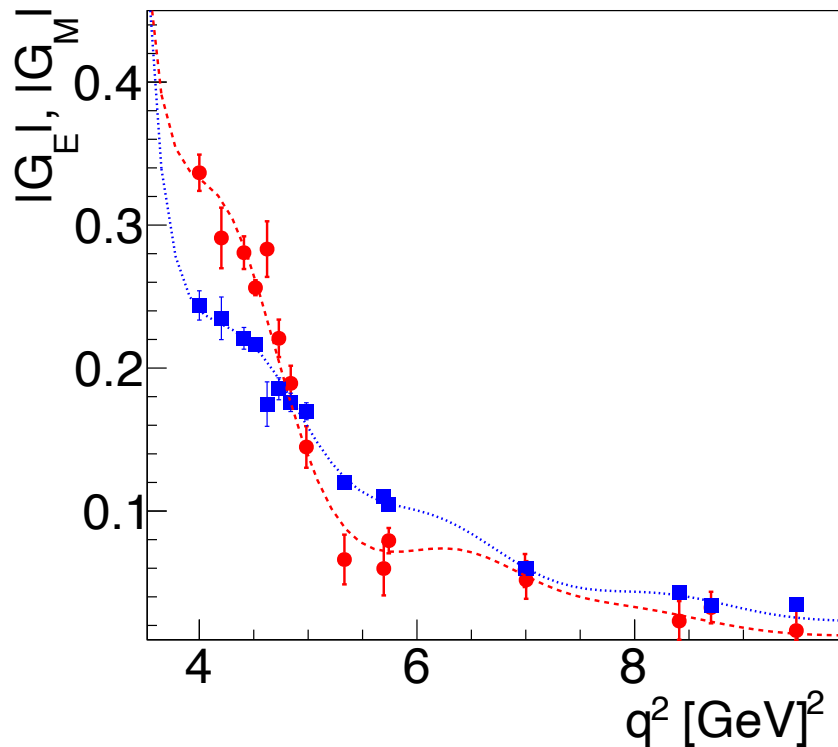
$$p = (1.20 \pm 0.04) \text{ GeV}$$

*Zero of the angular asymmetry*



# Sachs form factors: $|G_E|$ , $|G_M|$

From the fit on  $F_p$  and the fit on  $R$ ,  
the Sachs FFs (moduli) can be reconstructed



$$|G_E(s)| = F_p(s) \sqrt{\frac{1 + 2\tau}{R^2(s) + 2\tau/R^2(s)}}$$
$$|G_M(s)| = F_p(s) \sqrt{\frac{1 + 2\tau}{R^2(s) + 2\tau}}$$

Threshold constrain  $R=1$  for  $\tau=1$

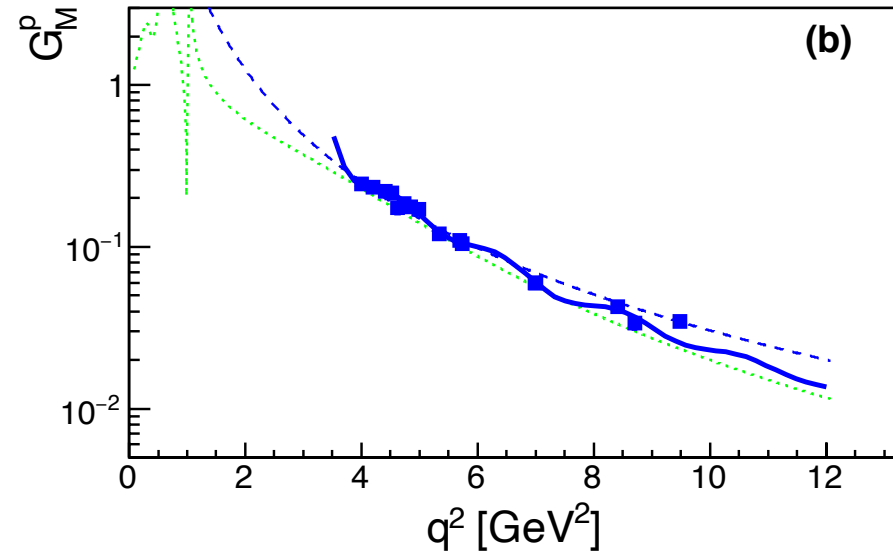
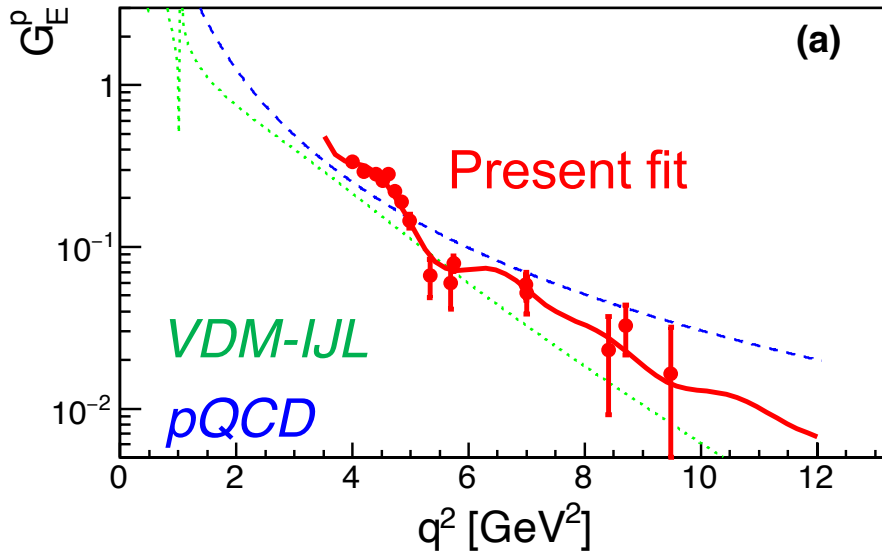
The fit gives :

$$|G_E| = |G_M| = 0.48$$



# Models

Parametrizations have been determined by fitting  $F_p$



$|G_E|$ : more pronounced oscillations  
faster  $q^2$ -decrease

Threshold constrain  $R=1$  for  $\tau=1$

The fit gives :

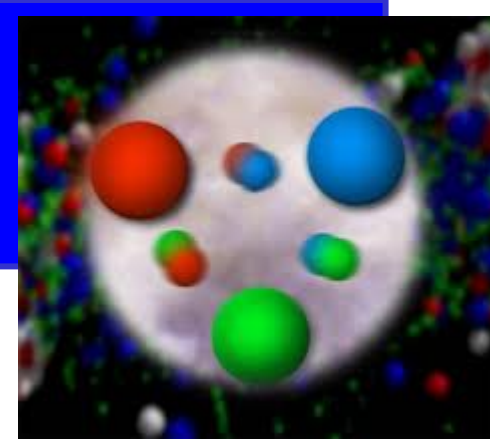
$$|G_E| = |G_M| = 0.48$$

pQCD : 0.34

VDM-IJL : 0.29



# The nucleon



3 valence quarks and  
a neutral sea of  $q\bar{q}$  pairs

antisymmetric state of  
colored quarks

$$|p\rangle \sim \epsilon_{ijk} |u^i u^j d^k\rangle$$
$$|n\rangle \sim \epsilon_{ijk} |u^i d^j d^k\rangle$$

## Main assumption

Does not hold in the spatial center of the nucleon: the center of the nucleon *is electrically neutral*, due to the strong gluonic field

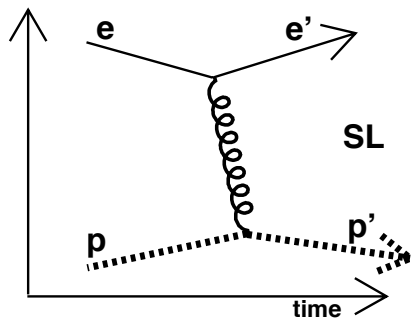
*E.A. Kuraev, E. T-G, A. Dbeyssi, Phys.Lett. B712 (2012) 240*



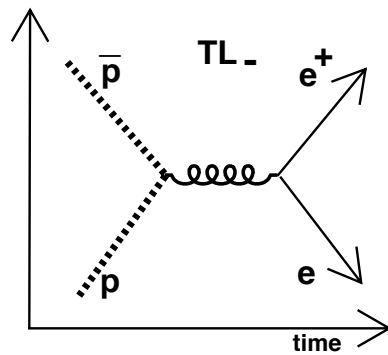
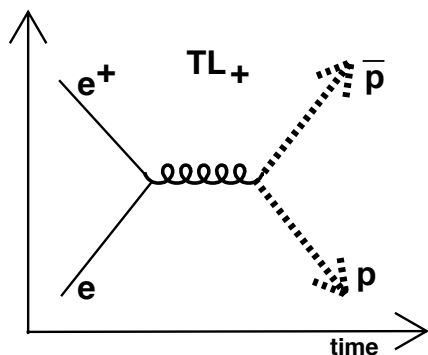
# Definition of TL-SL Form Factors

$$F(q^2) = \int_{\mathcal{D}} d^4x e^{iq_\mu x^\mu} \rho(x), \quad q_\mu x^\mu = q_0 t - \vec{q} \cdot \vec{x}$$

$\rho(x) = \rho(\vec{x}, t)$  space-time distribution of the electric charge in the space-time volume  $\mathcal{D}$ .



SL photon 'sees' a charge density



*TL photon can NOT test a space distribution*

*How to connect and understand the amplitudes?*



# $\rho(x)$ in the space-like region

and in the Breit frame or at small  $x$ :

density $\rho(r)$	Form factor $F(q^2)$	r.m.s. $\langle r_c^2 \rangle$	comments
$\delta$	1	0	pointlike
$e^{-ar}$	$\frac{a^4}{(q^2 + a^2)^2}$	$\frac{12}{a^2}$	dipole
$\frac{e^{-ar}}{r}$	$\frac{a^2}{q^2 + a^2}$	$\frac{6}{a^2}$	monopole
$\frac{e^{-ar^2}}{r^2}$	$e^{-q^2/(4a^2)}$	$\frac{1}{2a}$	gaussian
$\rho_0$ for $x \leq R$ 0 for $r \geq R$	$\frac{3(\sin X - X \cos X)}{X^3}$ $X = qR$	$\frac{3}{5}R^2$	square well



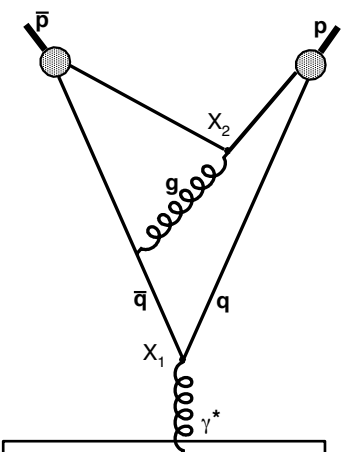
# Photon-Charge coupling

$$\rho(\vec{x})$$

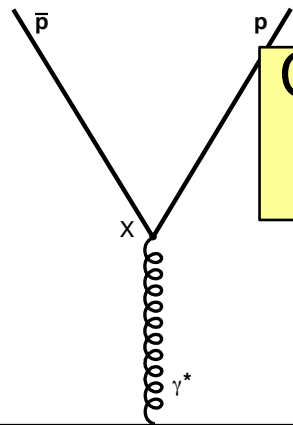
Fourier transform of a stationary charge and current distribution

$$R(t)$$

Amplitude for creating *charge-anticharge pairs* at time  $t$



*Resolved*



*Unresolved*

Charge distribution: distribution in time of  $\gamma^* \rightarrow$  *charge-anticharge vertices*

The simplest picture: qq pair + compact di-quark

*representation*

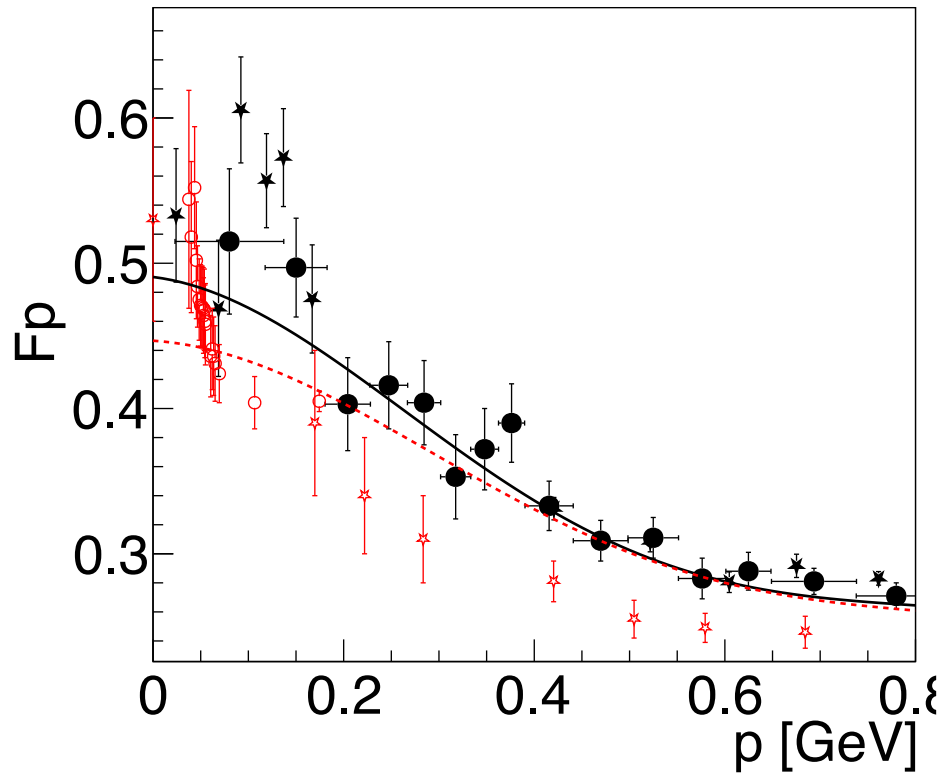


# Conclusions

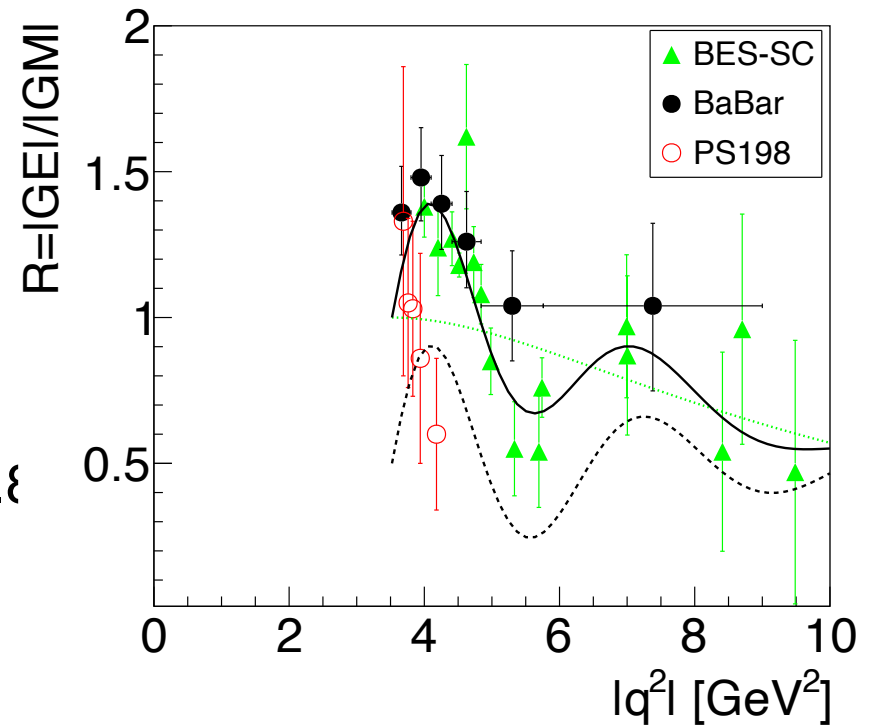
- New, precise data on time-like form factors and *first determination of individual FFs ( $|G_E|$  and  $|G_M|$ )*
- New data on *FFs ratio* and *angular asymmetry*
- *FFs ratio*: damped oscillations around **a monopole decrease**
- Oscillations more pronounced in  $|G_E|$ .
- Origin of oscillatory phenomena ? Di-quark as a necessary step towards hadron creation?
- Time reversal holds for ‘elastic’ annihilation channels: *are the two channels really equivalent* in the whole  $q^2$ -region?



# Data...



LEAR: antiproton  
BaBar, Novosibirsk  
BESIII





*Thank you for your attention*