

Hyperfine spectroscopy of simple systems in external magnetic and electric fields

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Simple systems?

- “Simple” enough to allow for high accuracy theoretical description
- Accessible to high precision spectroscopy experiments

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- Long-lived,
- dominant electromagnetic interactions

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- 2-body:

Ordinary: hydrogen-like atoms and ions

Exotic: positronium, muonium, μ^-p

Simple systems?

- Long-lived,
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- 2-body:
Ordinary: hydrogen-like atoms and ions
Exotic: positronium, muonium, μ^-p
- 3-body: molecular ions H_2^+ , HD^+ ,
exotic: antiprotonic helium, $p\mu^-p$

Precision HD⁺ spectroscopy

- High accuracy goals (Duesseldorf team):

10^{-10} for the determination of the electron-to-proton and electron-to-deuteron mass ratio

10^{-16} for testing the time variability of fundamental constants

Precision HD⁺ spectroscopy

- Theoretical uncertainties: below 10 kHz

This requires

- Higher orders relativistic and QED effects
- Hyperfine structure

but also

Corrections due to external magnetic and electric fields (constant or oscillating)

External fields

- Magnetic field of the Earth ($\sim 0.5\text{G}$)
- Incompletely shielded fields (?)

but also:

- Trap fields
- Bias fields
- ...

Hyperfine structure of HD⁺

- Classification:

L: orbital momentum

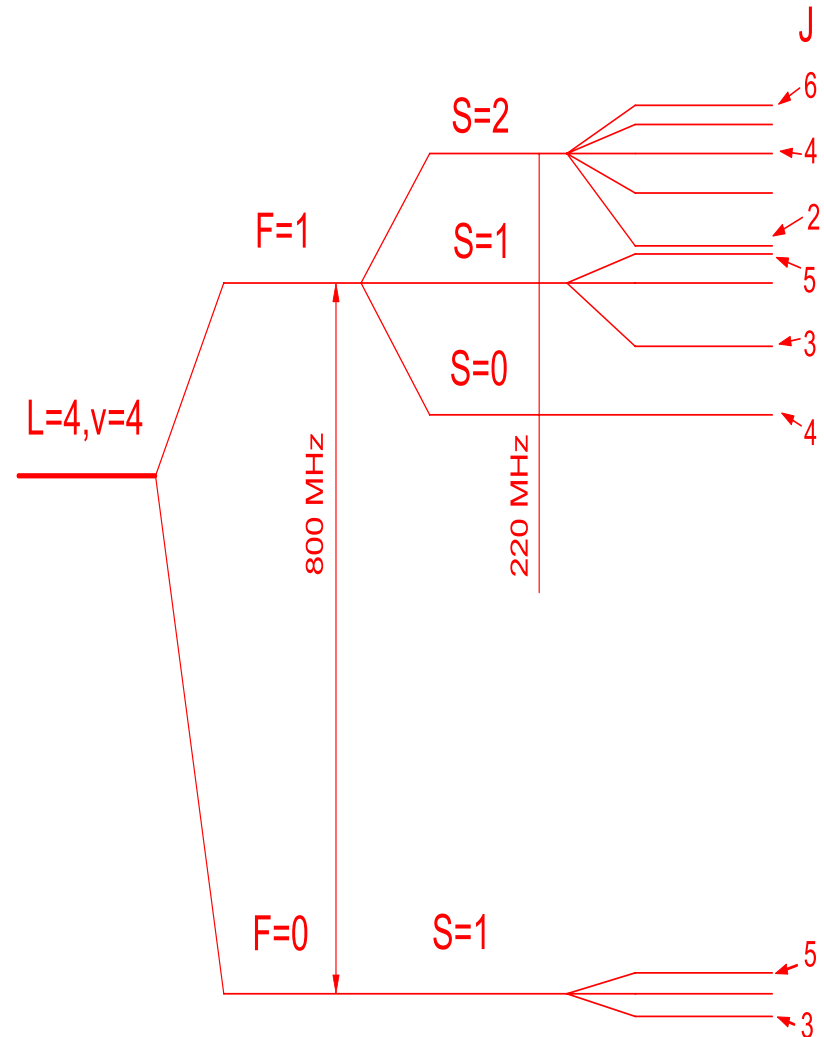
V: vibrational q.n.

$$F = S_p + S_e$$

$$S = F + S_d$$

$$J = S + L$$

$$(J_z)$$



Effective Hamiltonian

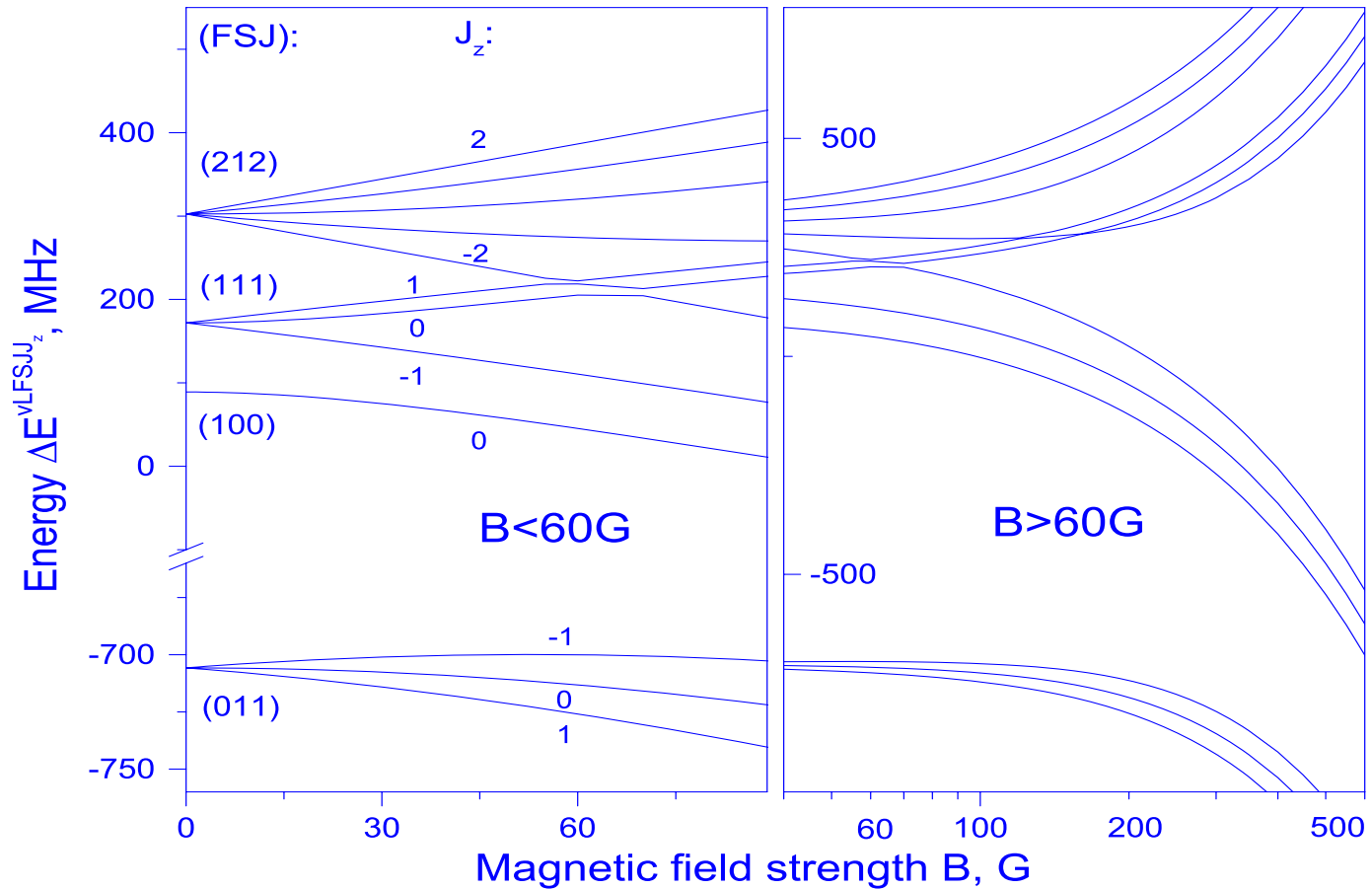
$$H^{\text{hfs}} = E_1(\mathbf{S}_p \cdot \mathbf{S}_e) + E_2(\mathbf{S}_d \cdot \mathbf{S}_e) + E_3(\mathbf{S}_d \cdot \mathbf{S}_p) + \\ + E_4(\mathbf{S}_d \cdot \mathbf{L}) + E_5(\mathbf{S}_p \cdot \mathbf{L}) + E_6(\mathbf{S}_e \cdot \mathbf{L}) + \dots$$

$$H = H^{\text{hfs}} + H^{\text{mag}}$$

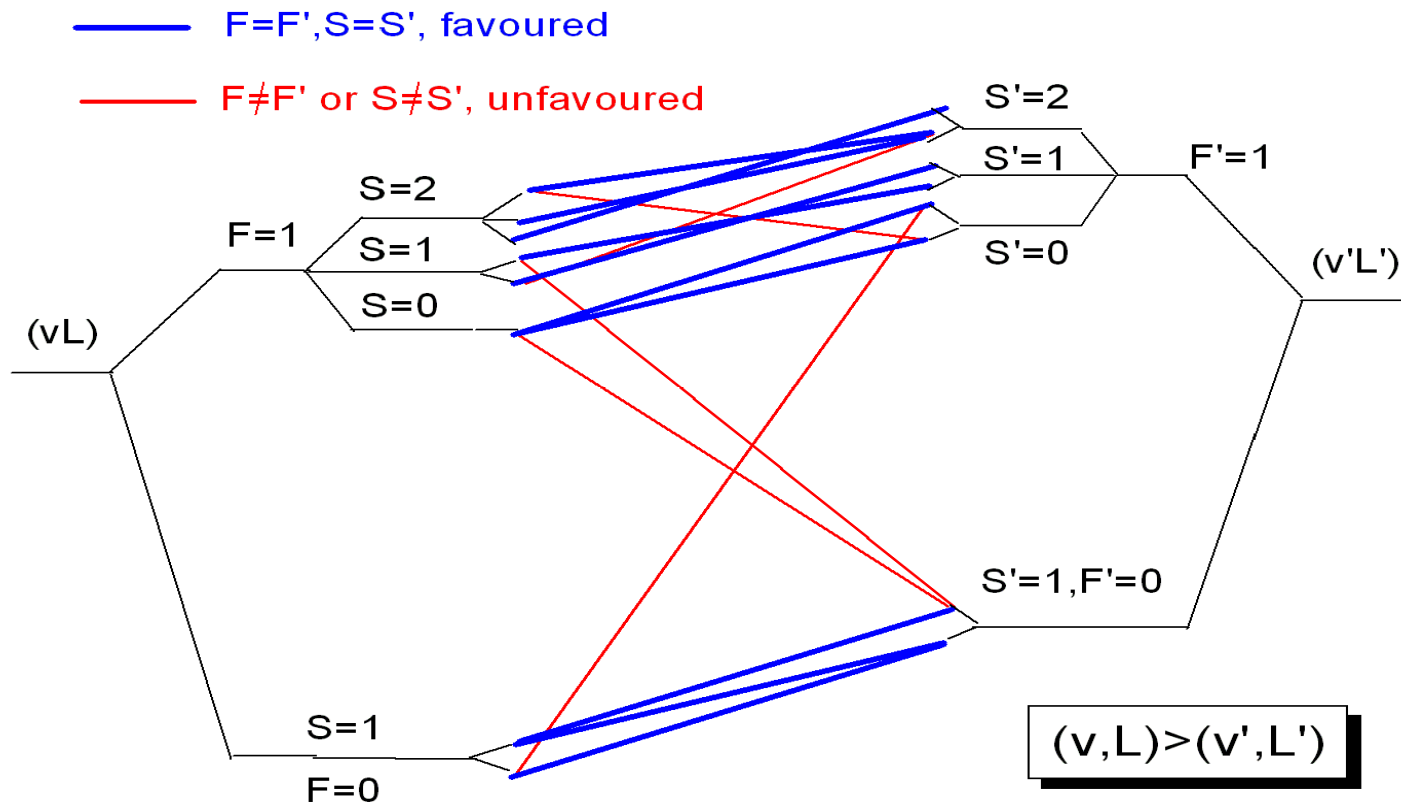
$$H^{\text{mag}} =$$

$$E_{10}(\mathbf{L} \cdot \mathbf{B}) + E_{11}(\mathbf{S}_p \cdot \mathbf{B}) + E_{12}(\mathbf{S}_d \cdot \mathbf{B}) + E_{13}(\mathbf{S}_e \cdot \mathbf{B})$$

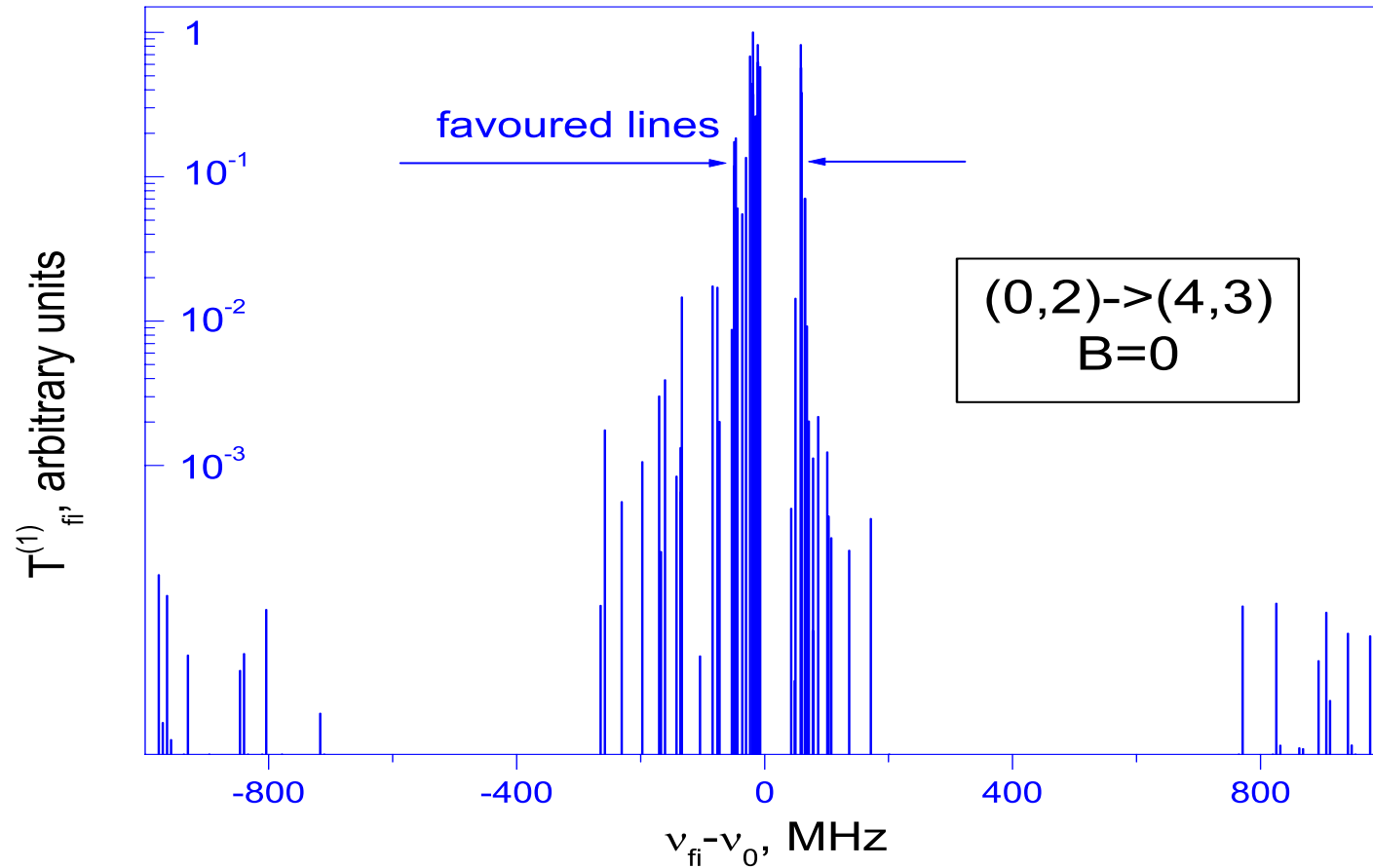
Zeeman splitting ($L=0, v=0$)



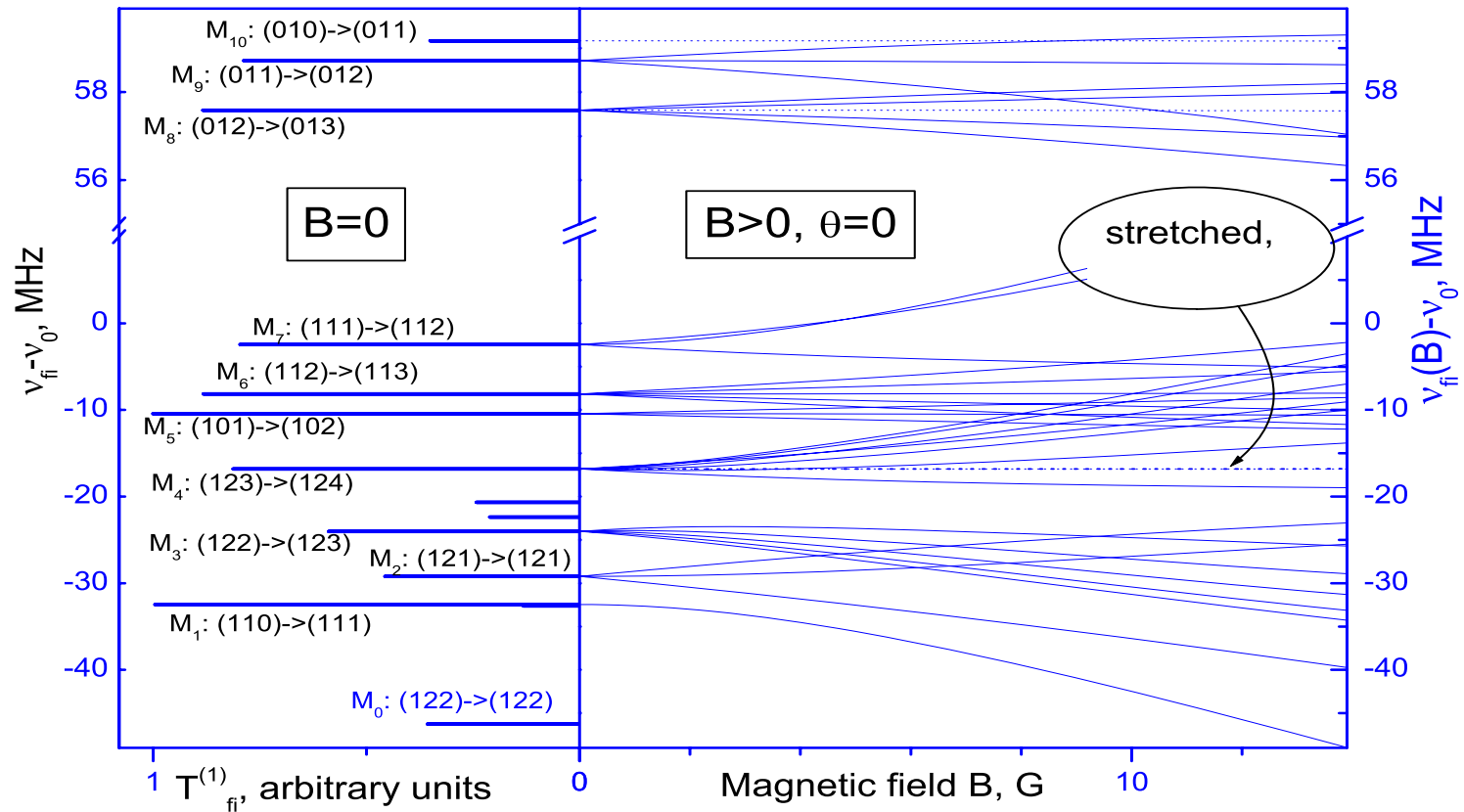
E1 transition spectra



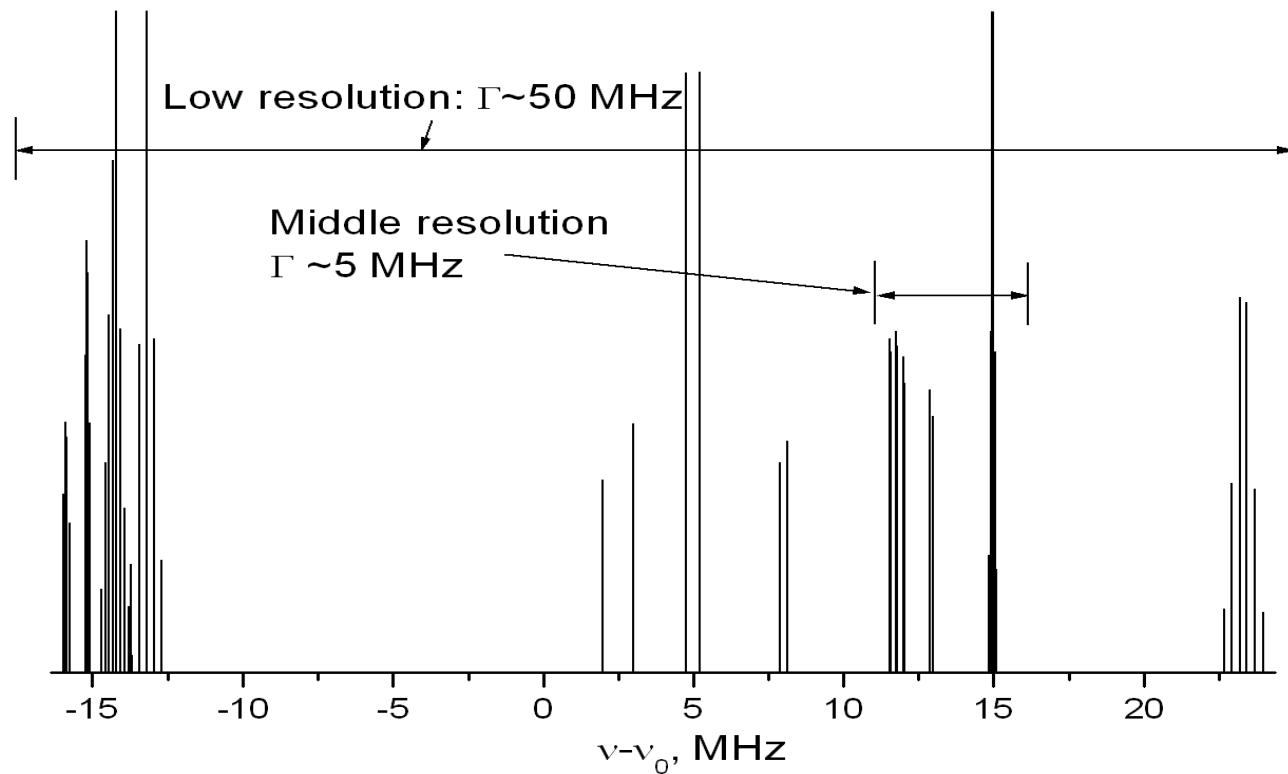
HFS of transition spectra



Zeeman structure of $(0,1) \rightarrow (4,2)$



Observable Zeeman effects



Observable Zeeman effects

Dependence on the spectral resolution.

Low resolution >50 MHz: **line broadening**

Middle resolution: also **shift of the h.f. lines:**

(F, S, J)	(1, 2, 3)		(1, 2, 1)		(1, 1, 2)		(0, 1, 3)		(0, 1, 2)		(0, 1, 1)	
$\bar{\nu} - \nu_0, \text{MHz}$	-23.481		-14.257		-6.902		58.251		58.861		59.426	
$\theta =$	0	$\pi/2$	0	$\pi/2$	0	$\pi/2$	0	$\pi/2$	0	$\pi/2$	0	$\pi/2$
$\delta\bar{\nu}(0.5 \text{ G}), \text{kHz}$	5.1	5.1	-1.5	-3.0	-1.4	-2.7	-0.1	0.1	-0.2	-0.1	0.0	-0.1
$\delta\bar{\nu}(1.0 \text{ G}), \text{kHz}$	16.9	17.0	-5.1	-9.8	-4.8	-8.7	-0.3	0.4	-0.5	-0.1	0.0	-0.5

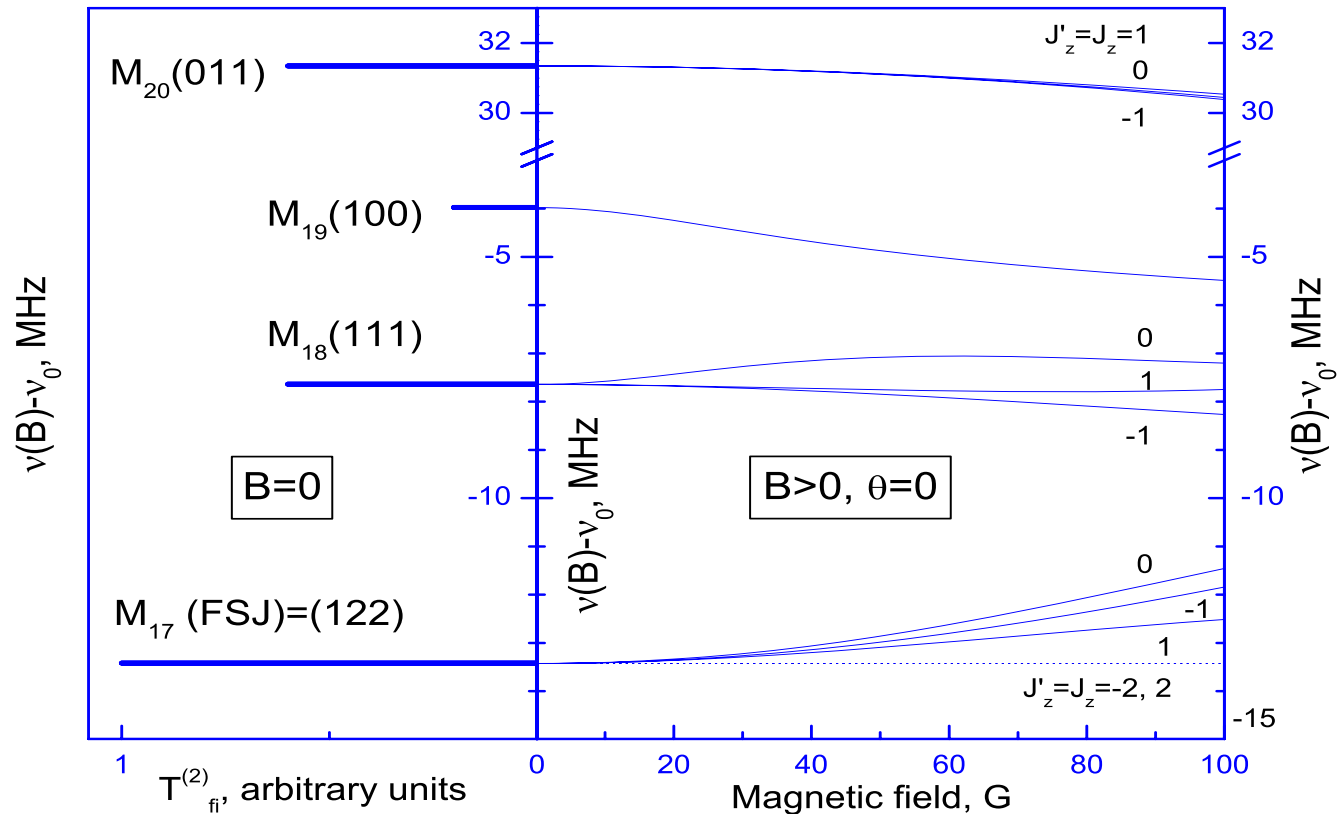
High resolution spectroscopy

- Individual Zeeman components resolved
- Search for transitions, insensitive to B

$$f(B)-f(0) = t.J_z.B+(q+r.J_z^2).B^2$$

1. **Transition between stretched lines:** no quadratic dependence on B
2. **More appropriate choices:**
Zeeman shift below 40 Hz/G²

2-gamma transition $(00) \rightarrow (20)$



Stark effect in HD^+

$$H = H^{\text{hfs}} + H^{\text{mag}} + H^{\text{el}}$$

$$H^{\text{el}} = -E \cdot d + q \cdot Q + \dots$$

The $E \cdot d$ term – in second order of P.T.

The $q \cdot Q$ term – in first order of P.T.

Polarizabilities of HD^+

Static Stark effect in HD⁺

- In absence of magnetic fields and HFS:

$$\Delta f_i = f_i(E) - f_i(0) = -\frac{1}{2}\alpha_i^{(D)}\mathbf{E}^2 + \frac{1}{2}\alpha^{(Q)} \cdot \mathbf{q}$$

$$\mathbf{E} = -\nabla V; \quad \mathbf{q}_{mn} = \nabla_m \nabla_n V$$

$$\alpha^{(D)} = -2 \sum_k \frac{|\langle k | \mathbf{d} | i \rangle|^2}{f_i(0) - f_k(0)} \quad \alpha_{mn}^{(Q)} = \langle i | \mathbf{Q}_{nm} | i \rangle$$

Convergence: $v'=v$, $L'=L\pm 1$: ~99%; el. excitations: ~1%

Moss et al, 2002, Koelemeij, 2011, Bakalov et al., 2011

Static Stark effect in HD⁺

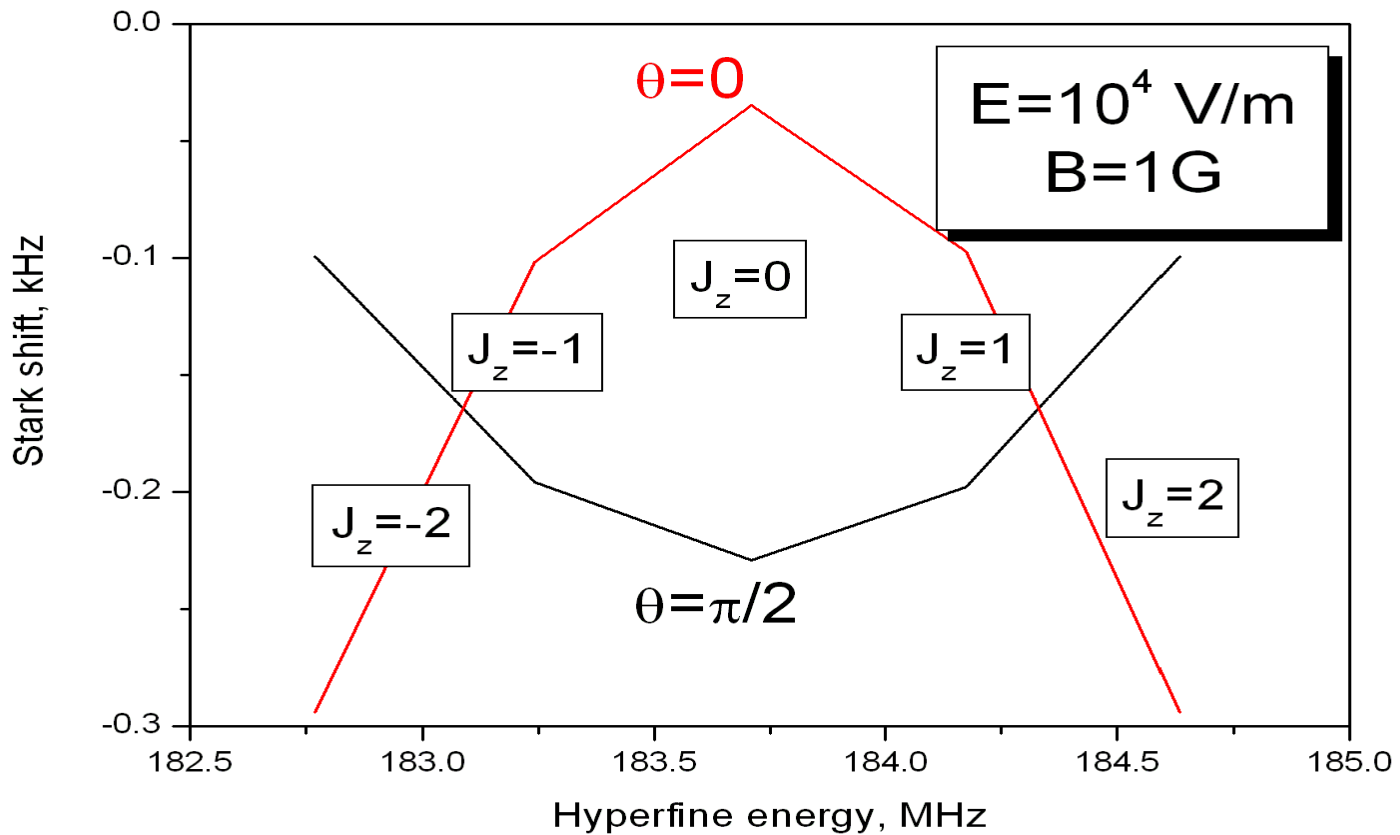
- Dipole polarizabilities $a_{vLM}^{(D)}$:
 - increase fast with v
 - decrease slowly with LTypically $0.3-3 \times 10^{-8}$ kHz/(V/m)²
and up to 10^{-6} kHz/(V/m)² for $v > 6$
- Quadrupole polarizabilities $a_{vLM}^{(Q)}$:
 - same behavior with v and LTypically: 1-2 orders of magnitude smaller

Stark and Zeeman effects with HFS

$$f_i(E) - f_i(0) = -\mathbf{E}^2 \sum_k (\sin^2 \theta (A_{-1}^{(ki)2} + A_1^{(ki)2}) + \cos^2 \theta A_0^{(ki)2}) / (f_i(0) - f_k(0))$$

$$A_q^{(ki)} = \langle k || \mathbf{d} || i \rangle \sum_{FSJ'J} \sqrt{2J+1} \begin{Bmatrix} J' & 1 & J \\ L & S & L' \end{Bmatrix} (-1)^{J'+L'+S} \beta_{FSJ'}^k \beta_{FSJ}^i C_{JJ_z, 1q}^{J' J'_z}$$

Stark + Zeeman effects

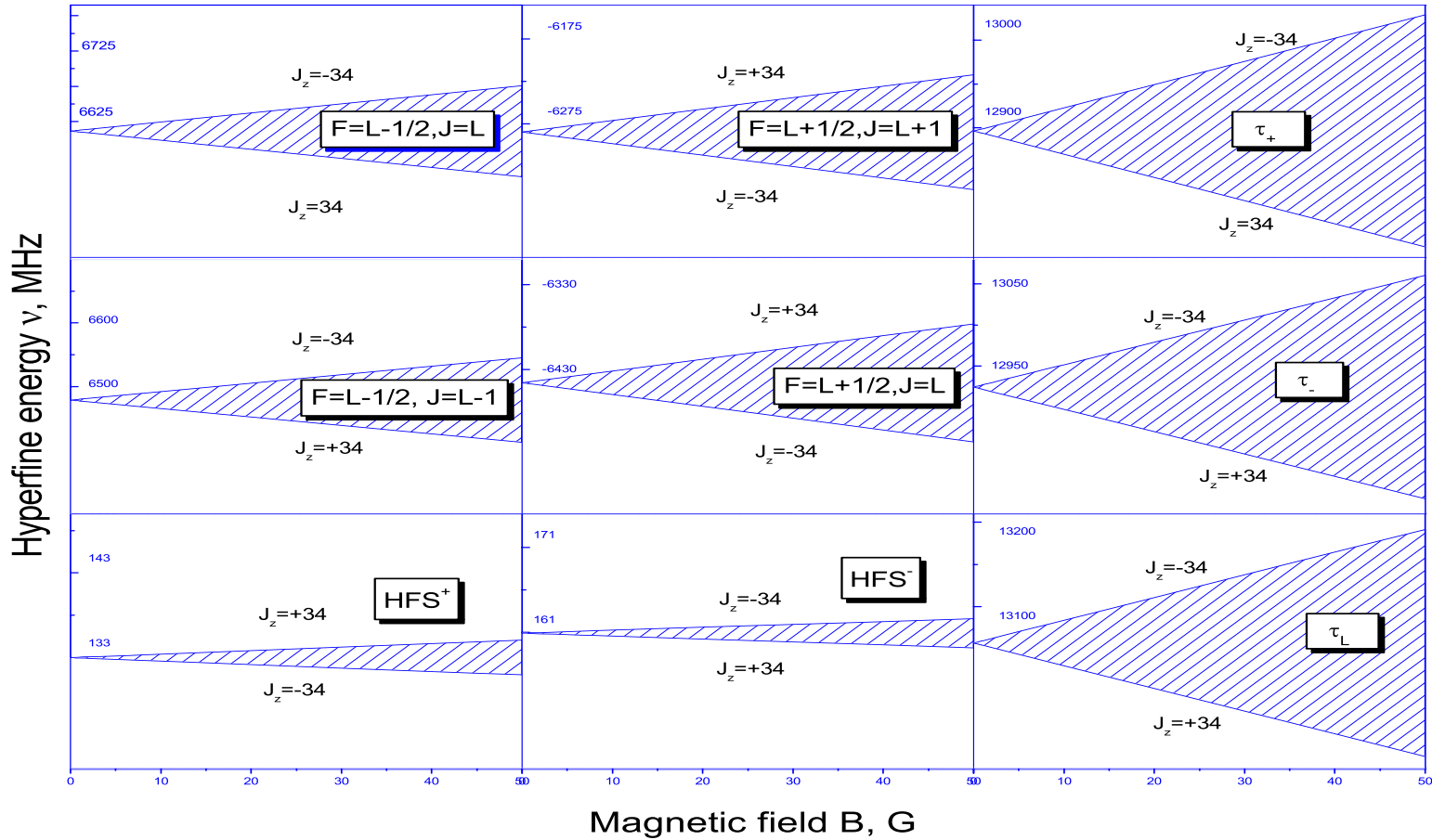


Stark + Zeeman effects

- The Stark shift depends on $|J_z|$, while the Zeeman depends on J_z :
- The Stark shift of the weighed mean of a h.f. multiplet differs from the simple form.
- Deformation of the hyperfine structure

Effects of up to a few kHz for $E \sim 10^8$ V/m

Zeeman shift in antiprotonic He



Zeeman shift in antiprotonic He

- The response of antiprotonic He levels to external magnetic fields – weaker, but

$$J_z \gg 1$$

- Quadratic effects – need to be studied,
expected to be stronger in ^3He
- Possible shift of the weighed mean of the multiplets of unresolved Zeeman lines.

Next steps to do

- **Dynamic Stark** (and Zeeman) effects
incl. **Black Body Radiation** shift