

# **Study of laser matter interaction in an intensity regime relevant for shock ignition**

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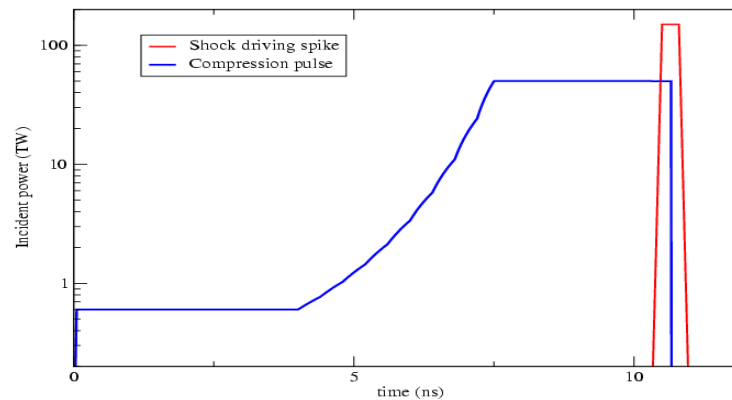
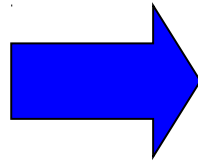
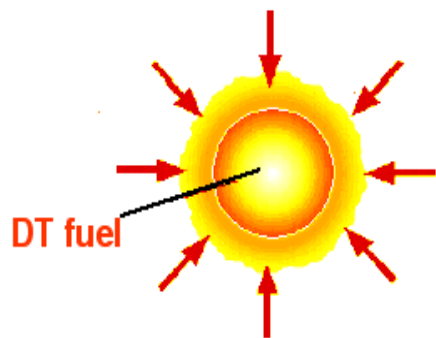
**LULI, Ecole Polytechnique, Paris, France**

*Experiment performed with the support of LASERLAB and HiPER*

# *The concept of shock ignition*

1: “normal” compression of a thermonuclear DT pellet with ns laser beams at  $I \approx 10^{14}$  W/cm<sup>2</sup> (we are able to compress!)

2: a high-intensity pulse ( $I \approx 10^{15}$  -  $10^{16}$  W/cm<sup>2</sup>) generates a strong shock ( $P \approx$  several 100 Mbar) which heats the central spot and creates the conditions for ignition



Timing is essential

# ***Goals of PALS experiment***

We wanted to investigate:

- 1) The effect of laser-plasma instabilities at  $I \approx 10^{16}$  W/cm<sup>2</sup>. Do they develop? How much light do they reflect? Do they create many hot electrons and at what energy?
- 2) Are we really able to couple the high-intensity laser beam to the payload through an extended plasma corona? Are we able to create a strong shock?

# ***PALS experiment***

## **Iodine Laser**

**=1.3  $\mu\text{m}$   $\tau = 300 \text{ ps}$   $E = 1500 \text{ J}$**

**$3\omega$   $\lambda = 0.44 \mu\text{m}$   $E \leq 500 \text{ J}$**

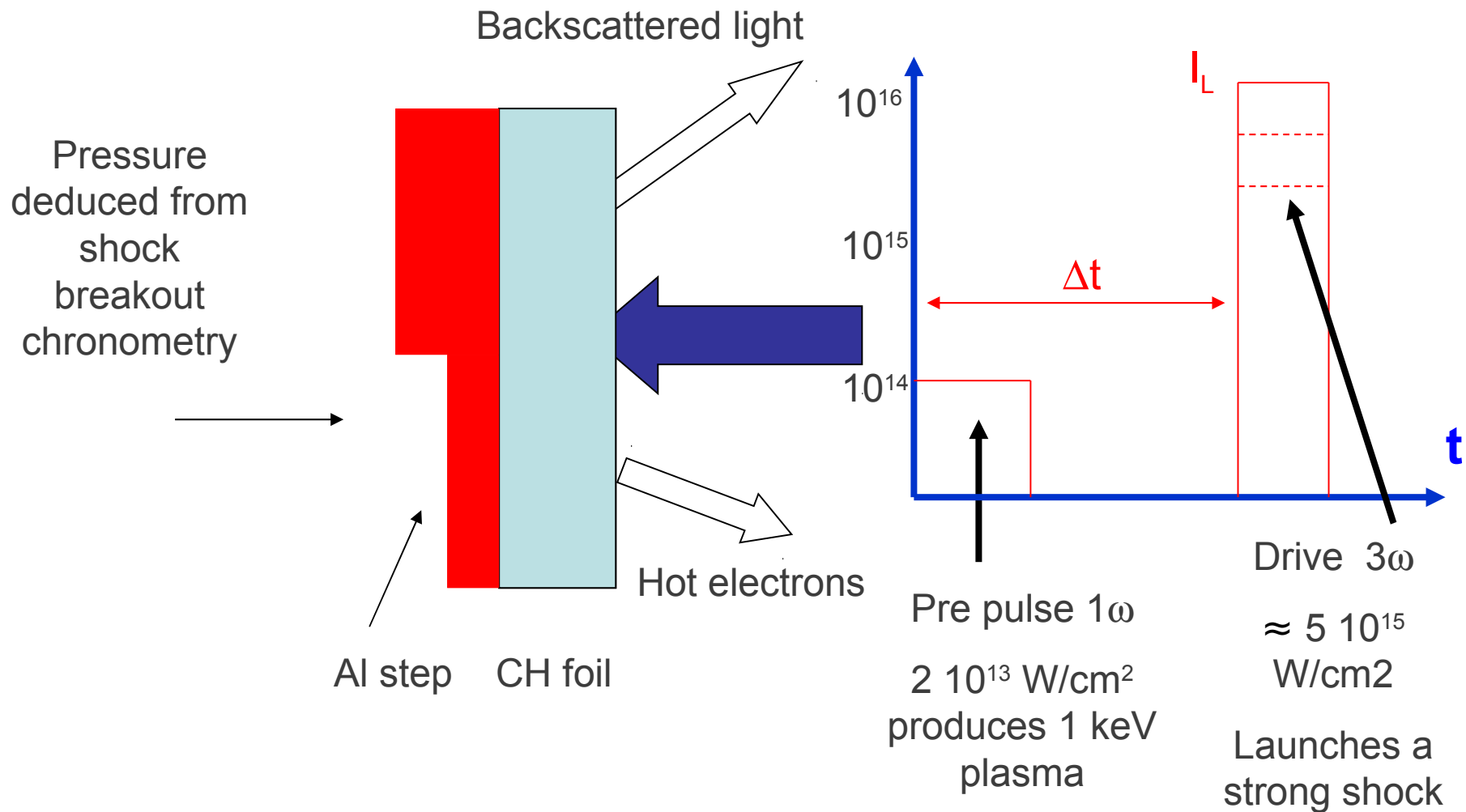
Laser Interaction Chamber at the  
PALS Laboratory in Prague



“Full” use of PALS laser facility:  
Main beam at  $3\omega$   
Auxiliary beam creating  
extended plasma  
XRL beam for diagnostics

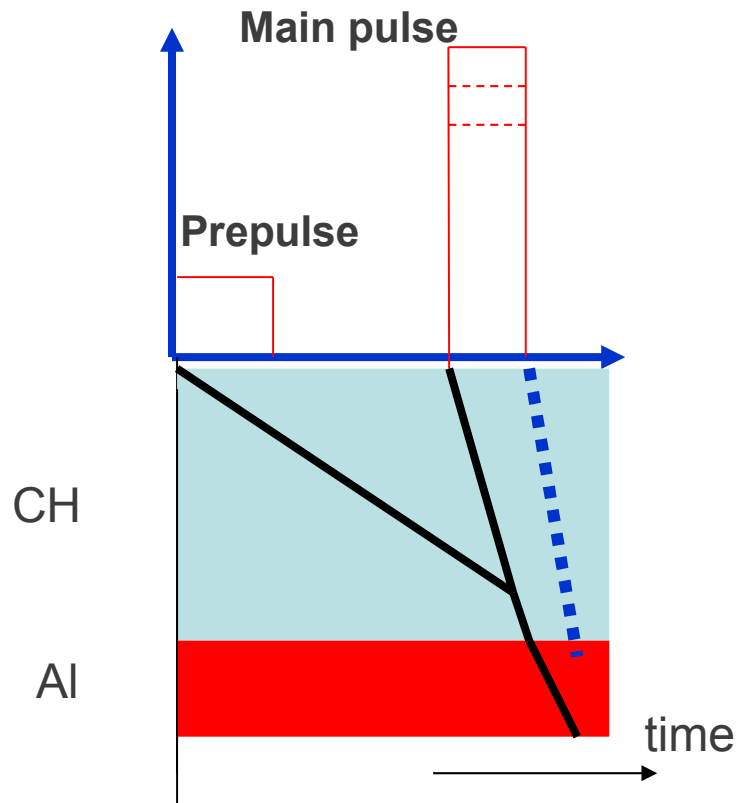


# Sketch of expt. set-up



The CH layer simulates the low-Z material of a pellet shell.  
The Al layer is a standard material for shock measurements

# Constraints on pulse and target



- Prepulse: produces 1keV, 1 mm plasma
- Drive : long enough to support pressure
- CH thickness : 2<sup>nd</sup> shock must catch up with 1st shock in CH
- Al thickness : trade off between
  - THIN Al : measurement must occur before the end of the main pulse (before the rarefaction wave launched at the end of the pulse catches up with the shock) **[actually this could NOT be realized in our conditions due to the short laser pulse...]**
  - THICK Al increases the time resolution

Final target configuration:

25  $\mu\text{m}$  CH layer (containing Cl to allow for X-ray spectroscopy)

25  $\mu\text{m}$  Al layer plus 10  $\mu\text{m}$  Al step (on some targets)

Targets produced by Ch.Spindloe and coworkers at the RAL Target Prep Lab

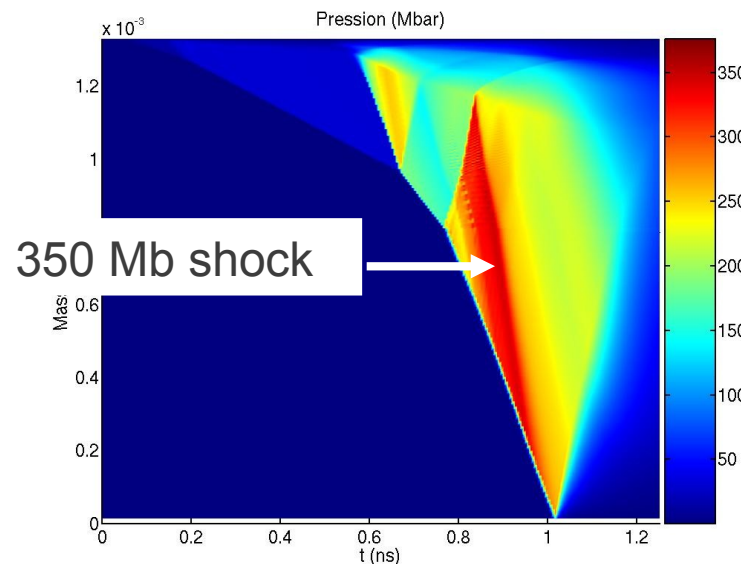
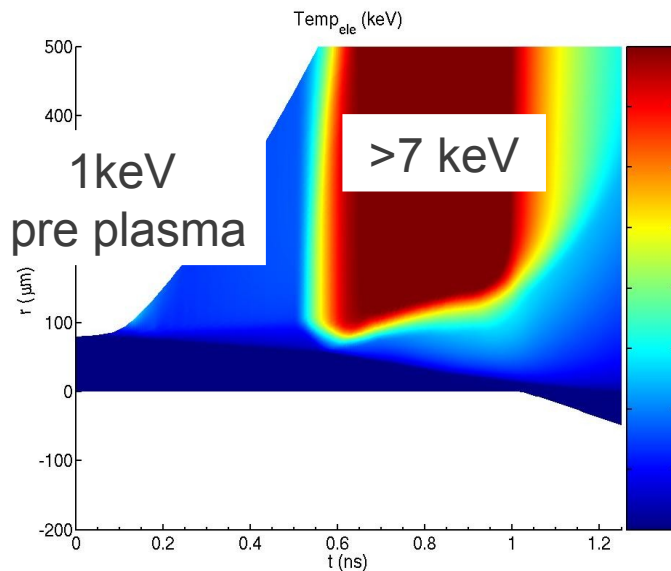
# Target and Laser parameters

Prepulse -  $E = 30 \text{ J}$ ,  $\lambda = 1.3 \text{ }\mu\text{m}$ ,  $\Phi = 1 \text{ mm}$ ,  $\tau = 300 \text{ ps} \Rightarrow I = 1.2 \cdot 10^{13} \text{ W/cm}^2$

Temperature  $T_e (eV) = 10^{-6} \left( I (W/cm^2) \lambda^2 (mm) \right)^{2/3} \approx 750 eV$

Main Pulse -  $E = 250 \text{ J}$ ,  $\lambda = 438 \text{ nm}$ ,  $\Phi = 100 \text{ }\mu\text{m}$ ,  $\tau = 300 \text{ ps} \Rightarrow I = 10^{16} \text{ W/cm}^2$

Pressure  $P (MBar) = 8.6 \left( \frac{I (W/cm^2)}{10^{14} \lambda (mm)} \right)^{2/3} \left( \frac{A}{2Z} \right)^{1/3} = 320 \text{ MBar}$



Absorption on the critical surface



# ***Diagnositics***

## **Phase 1**

### **Creation and characterization of preplasma with:**

- 1) X-ray deflectometry, using the PALS X-ray laser to obtain the density profile (Pals group)
- 2) X-ray spectroscopy, to obtain plasma temperature (Milan group)
- 3) X-ray pin-hole cameras (Pals group)

## **Phase 2**

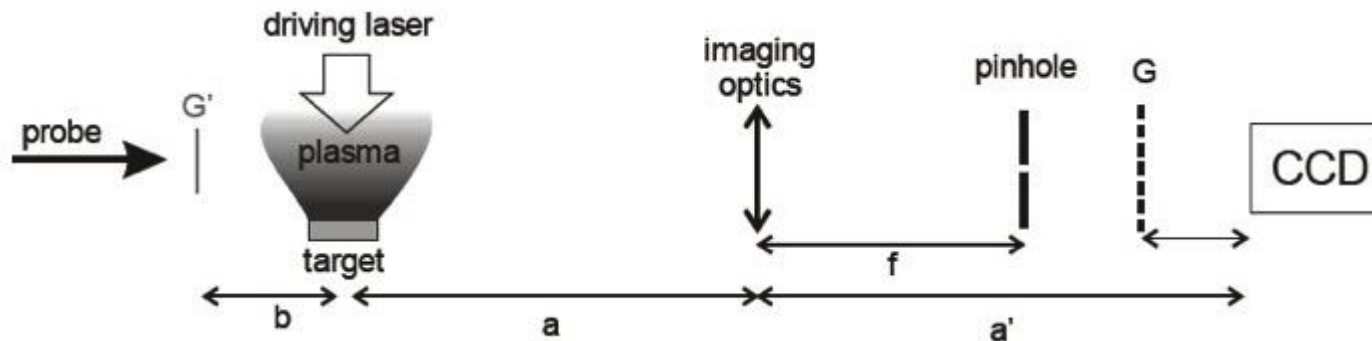
### **Characterization of shock formation and laser-plasma interaction with:**

The interaction of the main pulse has been studied using:

- 1) EEPHC diagnostic (energy encoded X-ray pin-hole camera) to measure plasma extension and characterize its emission (Pisa group).
- 2) Shock chronometry (measuring the self emission from the target rear side with a streak camera) (Milan group)
- 3) Optical imaging, spectroscopy, and calorimetry of back reflected radiation to evaluate the onset and amount of back reflected light from parametric instabilities (SRS, SBS, TPD) (Pisa group)

# Density measurement - XRL deflectometry

The technique used for this measurement is based on the deformation of Talbot pattern of 2D grating caused by gradients of index of refraction (electron density) of plasma

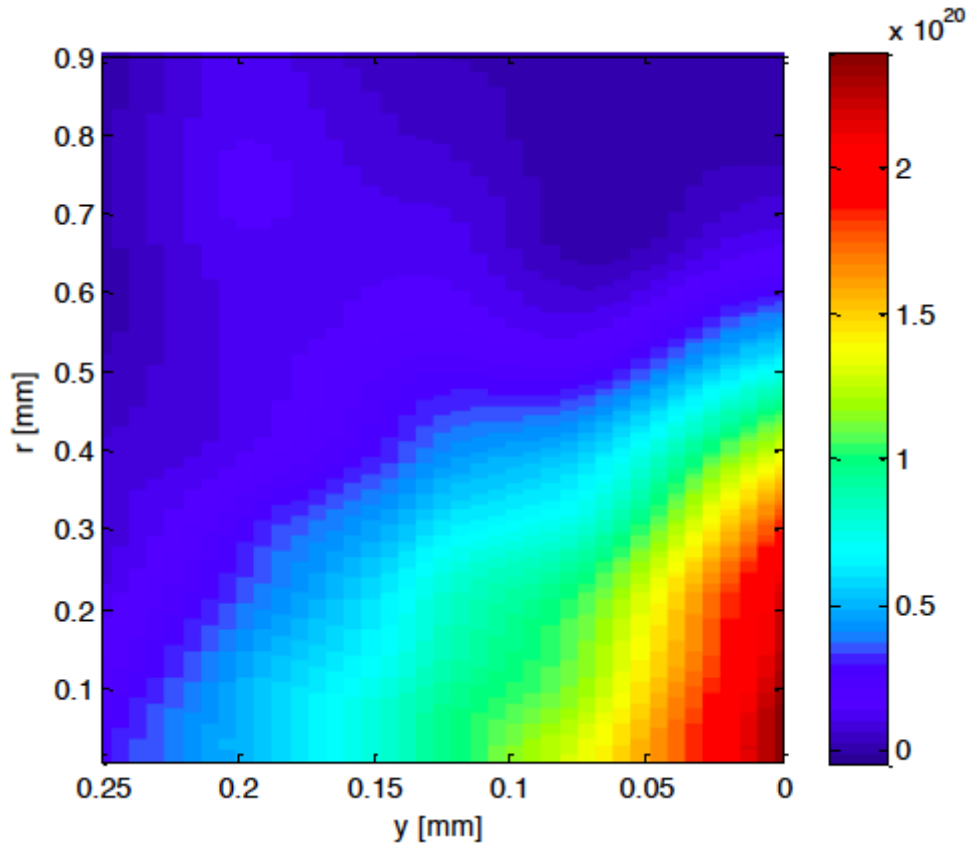


Ne-like zinc X-ray laser emitting at 21.2 nm, operated in single pass with 150-ps pulses of 200  $\mu$ J. Mo-Si multilayered spherical mirror with  $f=250$ mm used to image the plasma on back-illuminated X-ray CCD with  $M = 8.2$ . A pinhole of 0.5mm diameter was put to the image of the XRL source (2500 mm from the imaging mirror), to reduce the signal of plasma self-emission. The 100  $\mu$ m period laser-drilled grid made of 5  $\mu$ m thick steel was at 1275 mm from CCD.

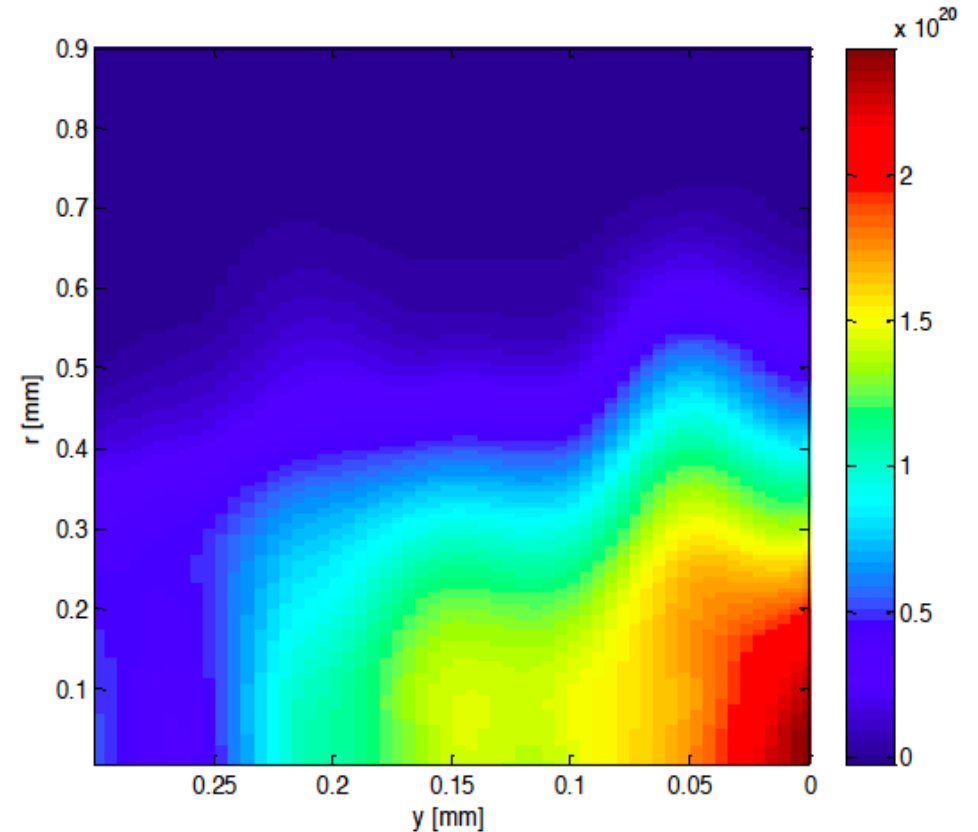
J. Nejdli, M. Kozlova, Plasma density-gradient measurement using x-ray laser wave-front distortion, Proc. SPIE Vol. 7451, 745117 (2009)

# Density profiles of pre-plasma

2D density profiles at different times after irradiation



**0.3 ns**



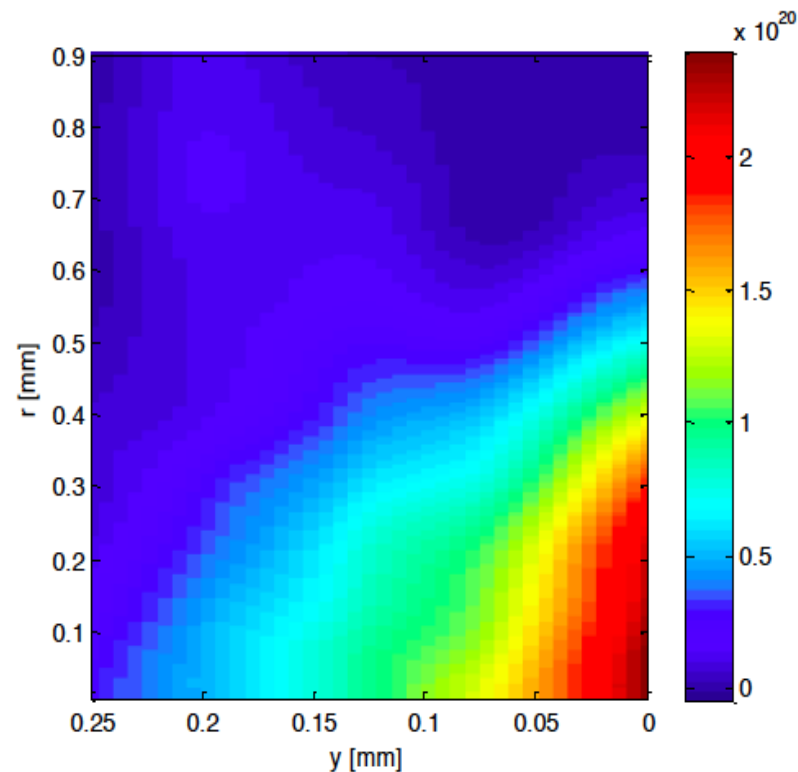
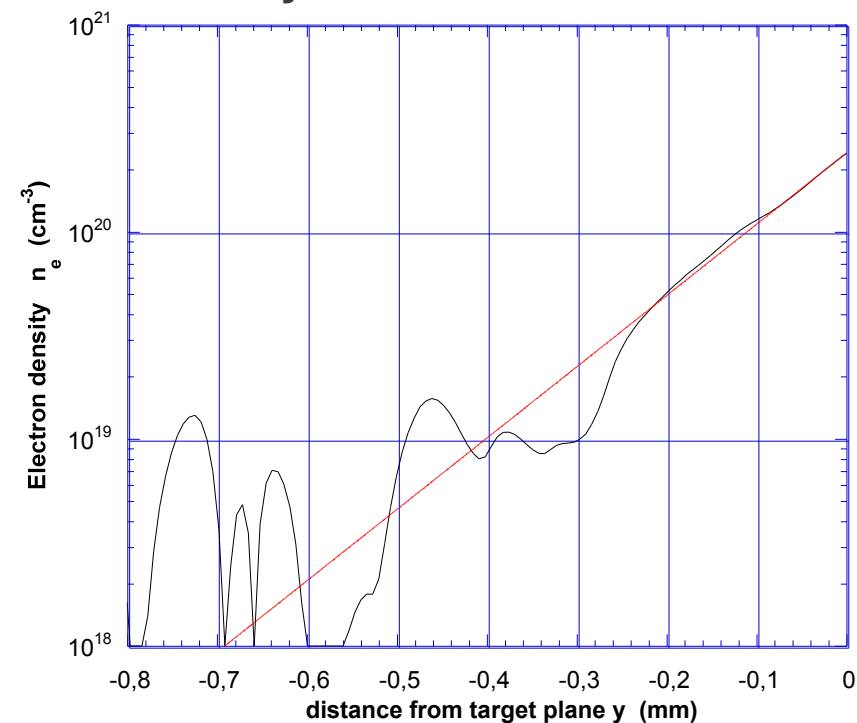
**0.9 ns**

Probing of CH plasma by XRL beam 0.3ns after driving pulse arrival; a) reference pattern without plasma, b) signal recorded when probing the plasma, c) field of pattern deformation

**Cylindrical symmetry allows resolving the electron density profile via inverse Abel transform**

# Density profiles of pre-plasma

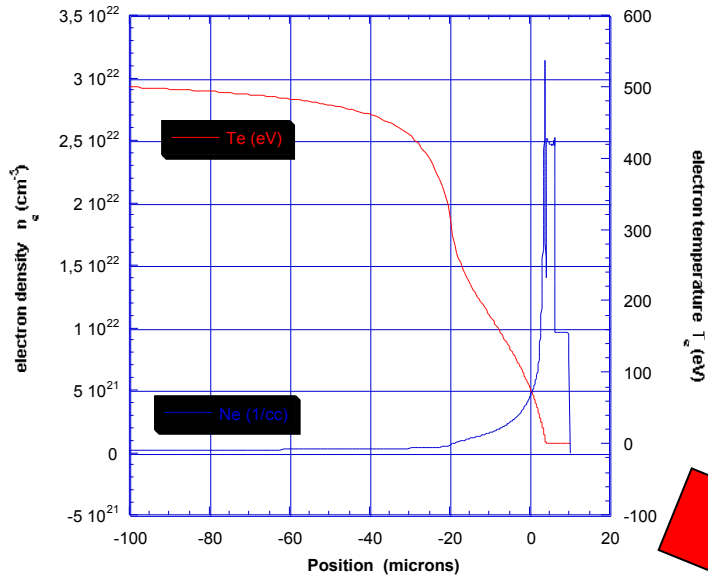
Along the axis, plasma expansion is practically 1D and the profile is exponential (as expected) but with some “bumps” in density



**Max accessible density  $\approx 2.4 \times 10^{20} \text{ cm}^{-3}$**

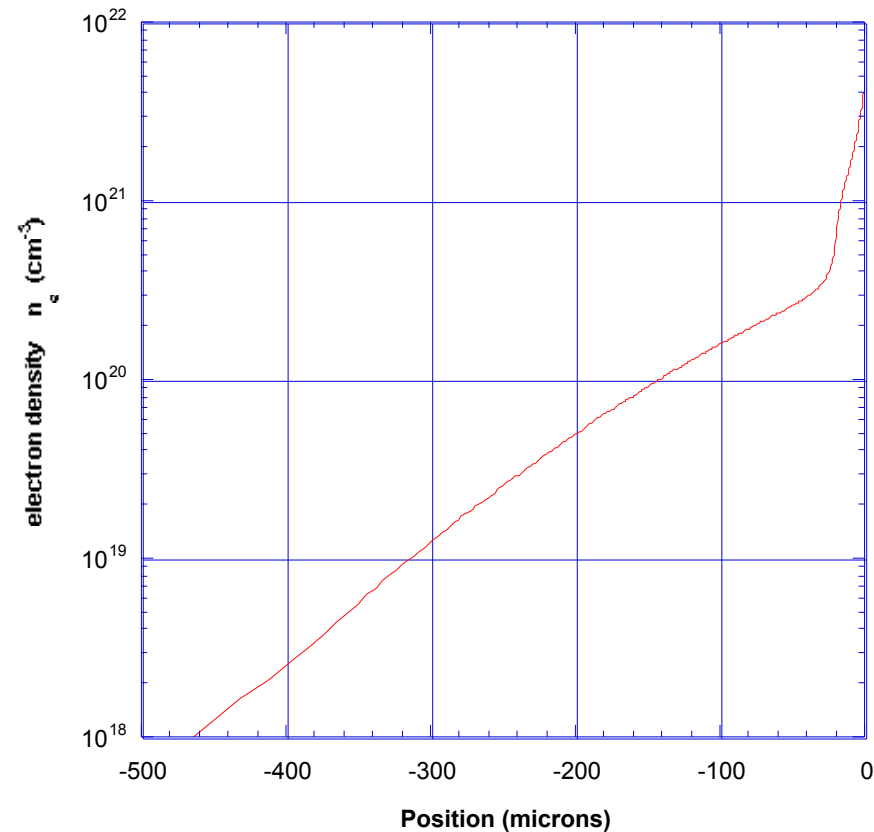
# Density profiles of pre-plasma

“Classical” exponential profile is well reproduced by 1D hydro simulations performed with the code MULTI



There is a coincidence in the exponential behavior between experimental and simulated data if we take in account of a prepulse of laser system. Hydro simulations allow to recover the plasma temperature (between 300 and 500 eV depending on laser intensity).

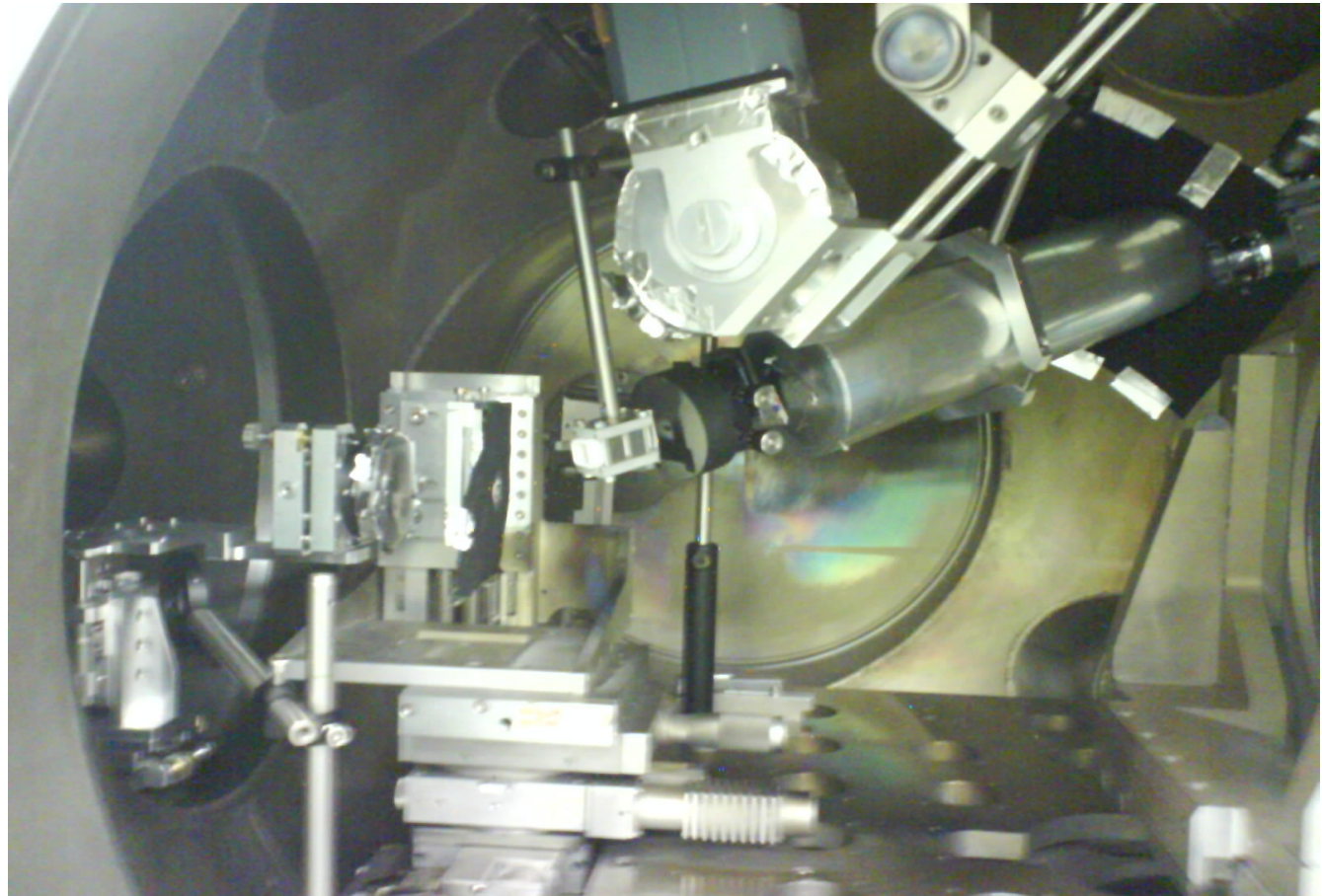
Exponential below a few  $10^{20} \text{ cm}^{-3}$



# *keV Spectrometer*

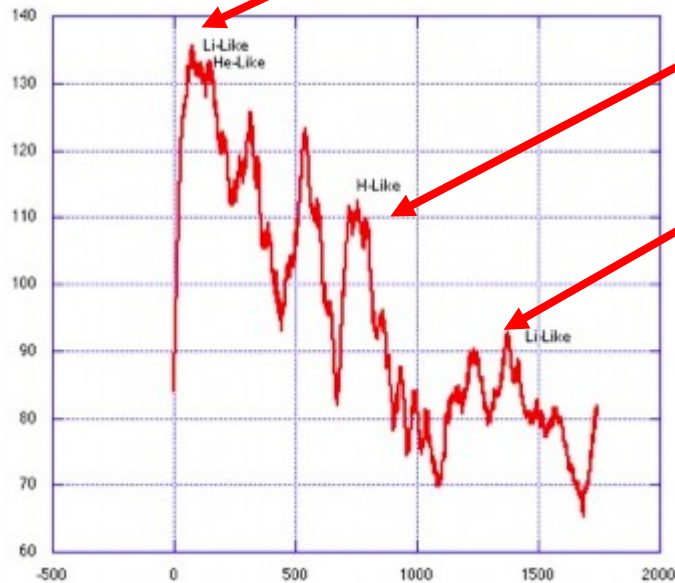
- ADP (ammonium dihydrogen phosphate) with  $2d \approx 10.659 \text{ \AA}$
- Placed at  $\approx 20 \text{ cm}$  from the source and with a Bragg angle  $\approx 19^\circ$

- Observed range from 2600 to 3600 eV



# X-ray spectra from pre-plasma

Poor spectral resolution due to large source size ( $\approx 1\text{mm}$ ) implying a line width  $\Delta E \approx 40\text{ eV}$



## Observed Li-, He- and H-like Cl ions

Cl	Li	3.8383	3p-1s
Cl	Li	3.8509	3p-1s
Cl	Li	3.8565	3p-1s
Cl	Li	3.862	3p-1s
Cl	Li	3.8701	3p-1s
Cl	Li	3.883	2s3d-1s2p
Cl	H	4.18518	2p-1s
Cl	H	4.1852	2p-1s
Cl	H	4.18533	2p-1s
Cl	H	4.19055	2p-1s
Cl	H	4.1906	2p-1s
Cl	H	4.19069	2p-1s
Cl	H	4.19074	2p-1s
Cl	He	4.4438	2p-1s
Cl	Li	4.4585	2p-1s
Cl	Li	4.463	2p-1s
Cl	He	4.4669	2p-1s
Cl	Li	4.4837	2p-1s
Cl	Li	4.4925	2p-1s
Cl	Li	4.497	2p-1s
Cl	Li	4.5215	2p-1s

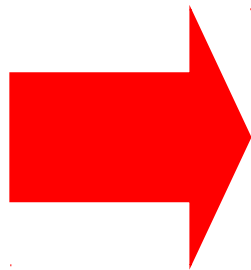
Due to low resolution and presence of non-identified lines (from impurities?) this spectrum only allowed to estimate the order of magnitude for plasma temperature of the order of 500 - 1000 eV

# ***Preplasma characterization***

Plasma with  $n_e > 10^{20} \text{ cm}^{-3}$  extends over 200  $\mu\text{m}$  (perpendicularly to target surface) and over 800  $\mu\text{m}$  radially (expected spot size  $\approx 1 \text{ mm}$ )

Hydro compatible with plasma temperature of 300-500 eV

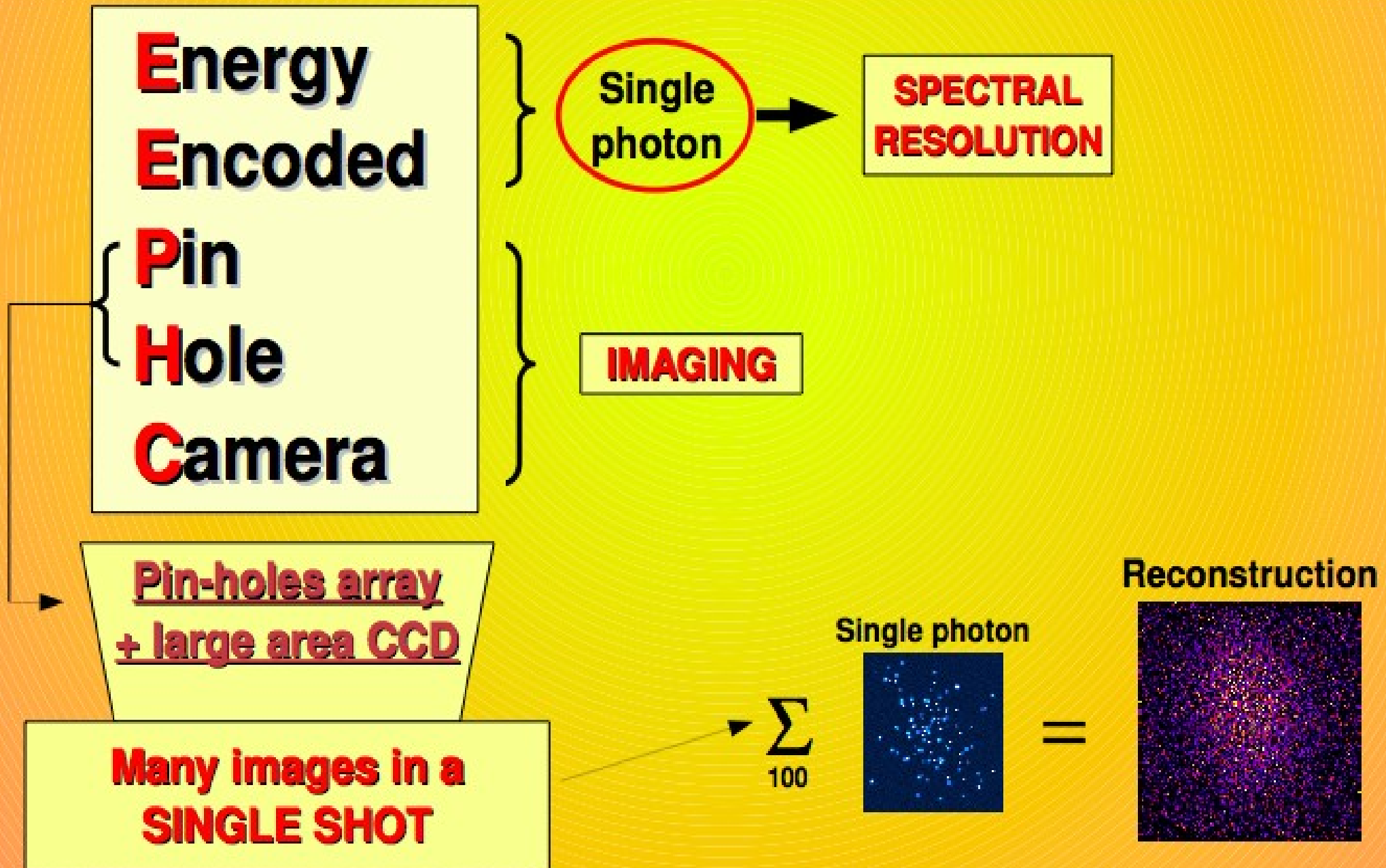
(Size confirmed by X-ray PHC images)



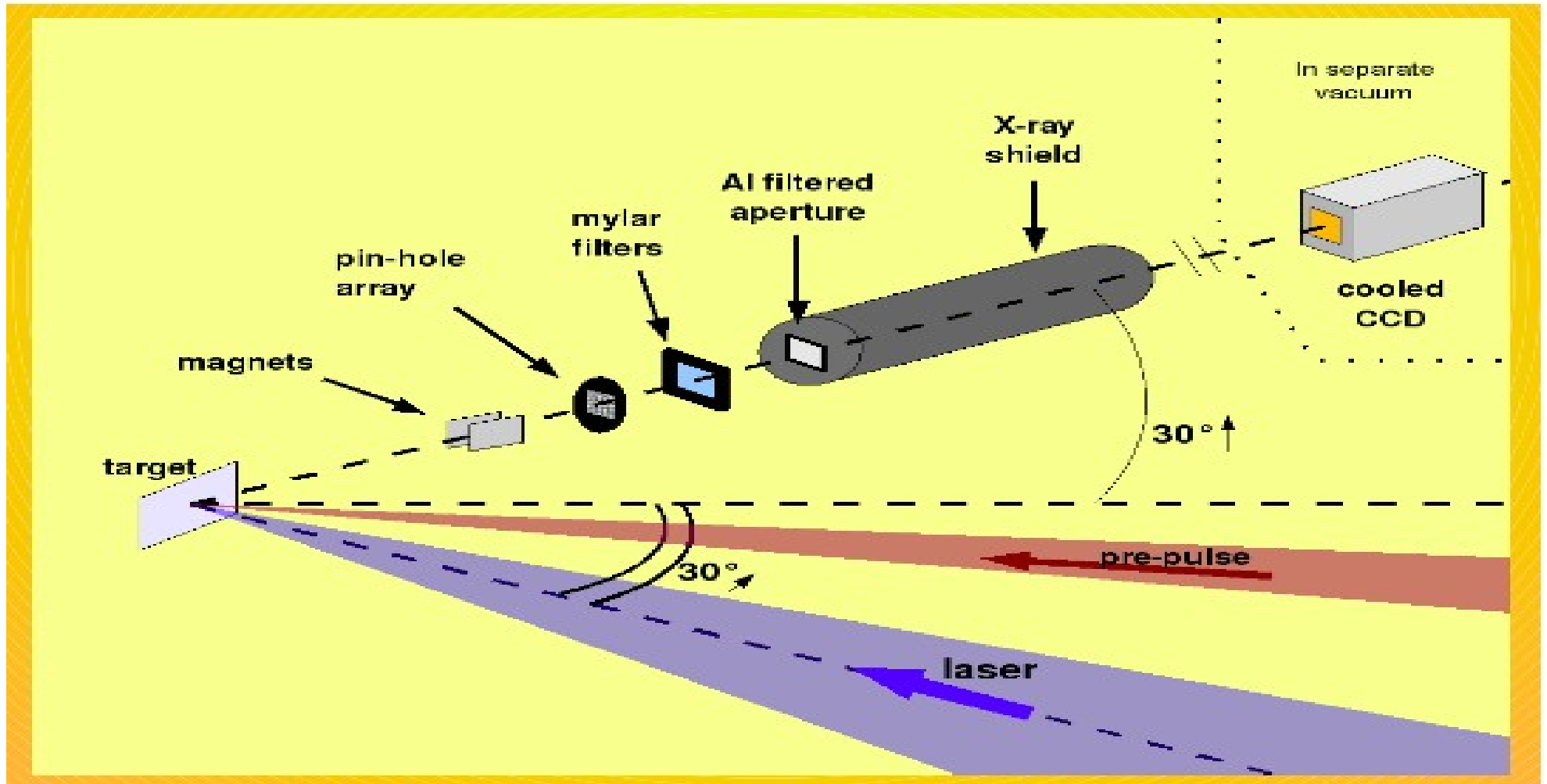
***Shock Ignition on  
PALS: Phase 2***



# *Energy encoded X-ray PHC*



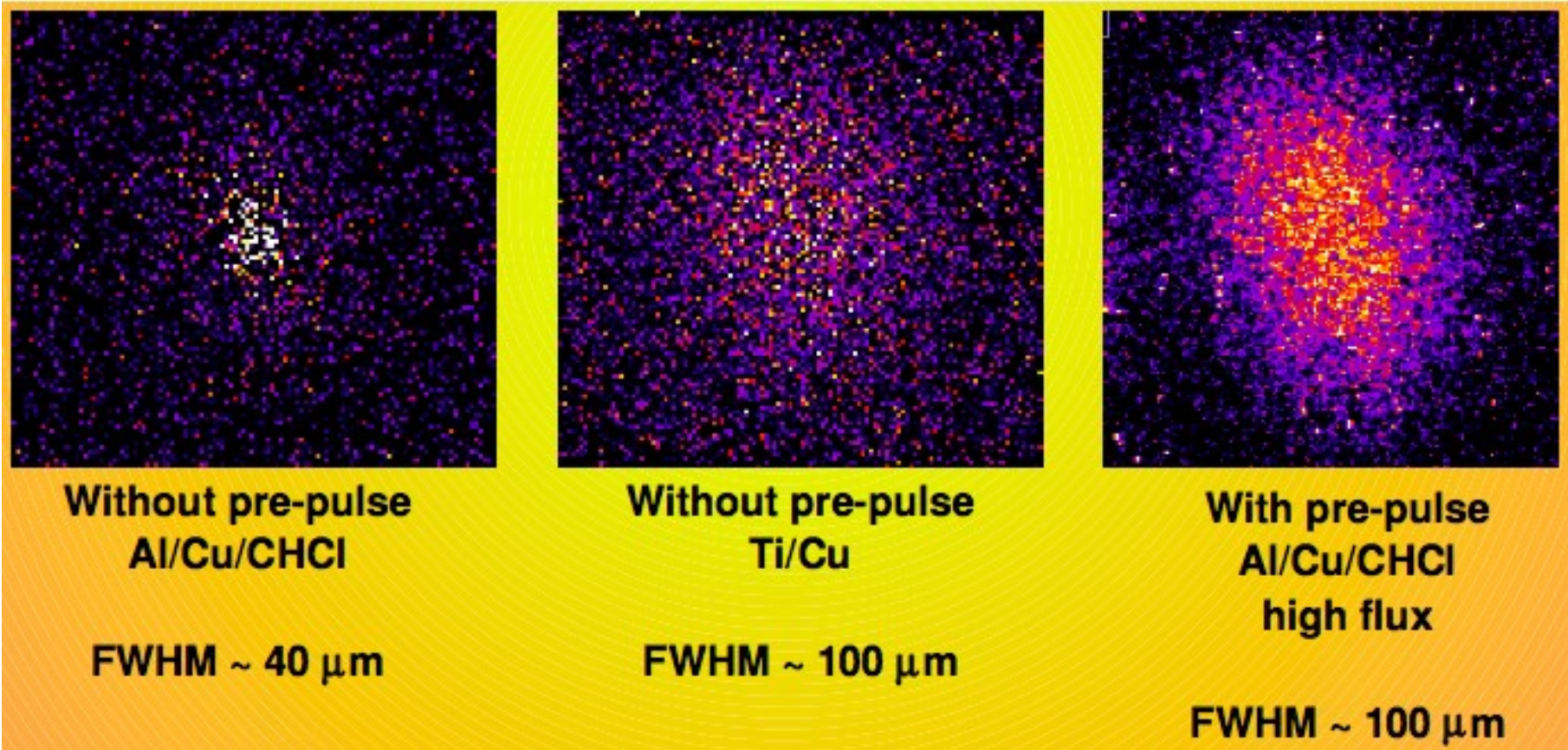
# EEPHC Set-up



When using such a diagnostics we added a thin (1  $\mu\text{m}$ ) Cu layer between Al and CH. We also used Ti/Cu targets

No exploitable results without main beam

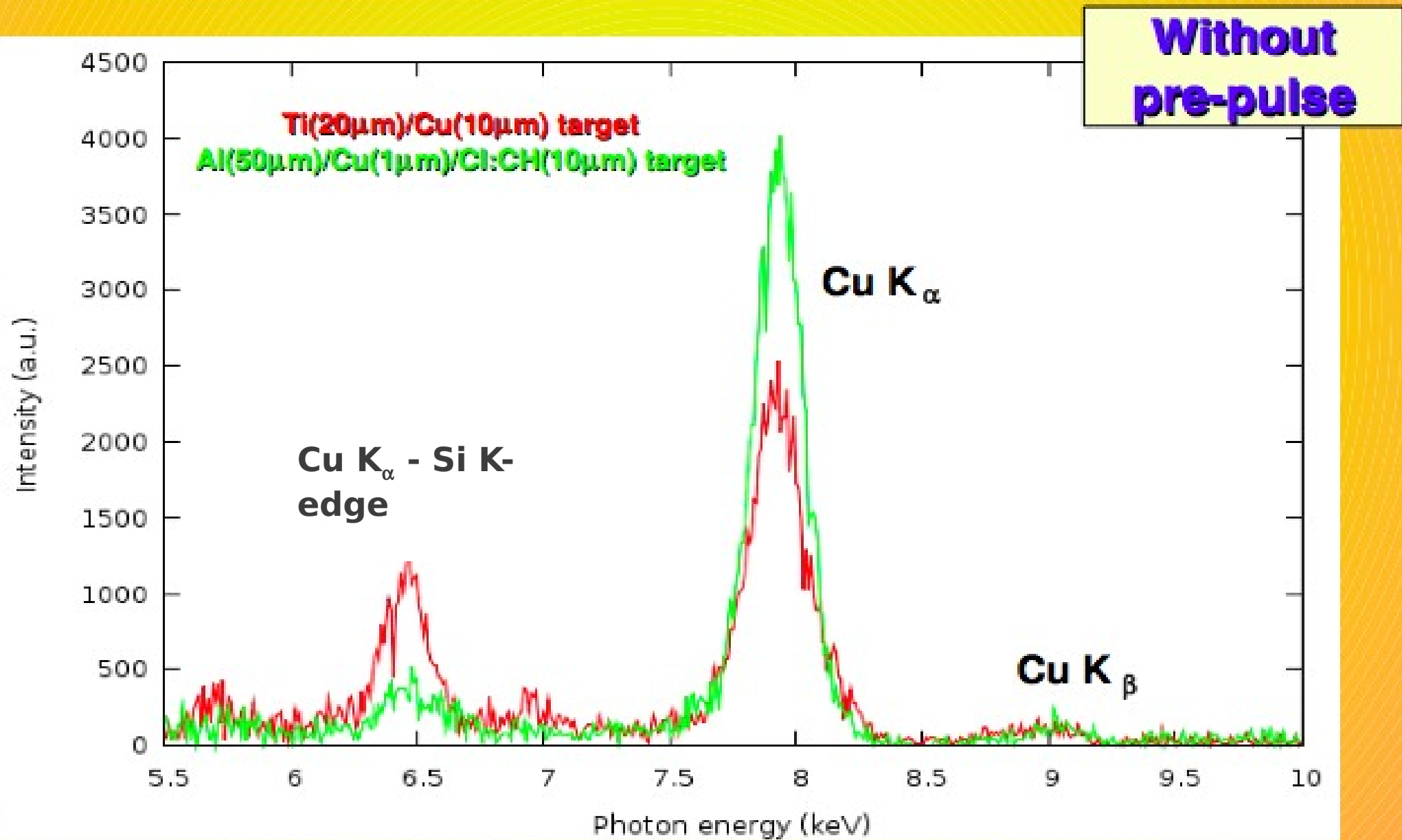
# *EEPHC results: imaging*



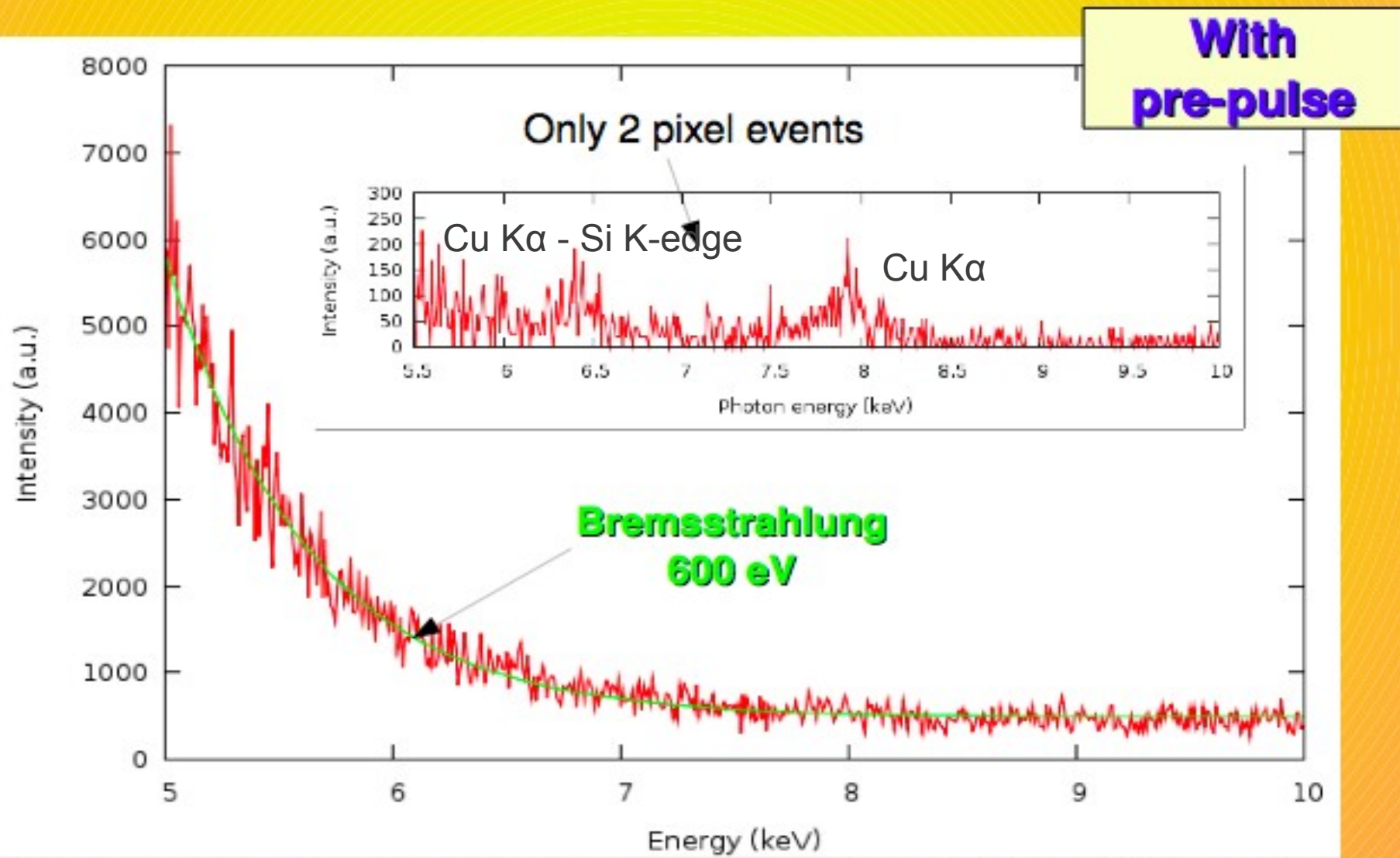
Focal spot size 100 μm (Range of 50 keV electrons in Cu 7 μm, in CH 40 μm)

Work in progress to retrieve hot electron energy. At  $10^{16}$  W/cm<sup>2</sup> one expects  $T_e = 100 \text{ keV} (I_{17} \lambda^2)^{1/3} \leq 30 \text{ keV}$  but higher energies may be produced by SRS

# ***EEPHC results: imaging***

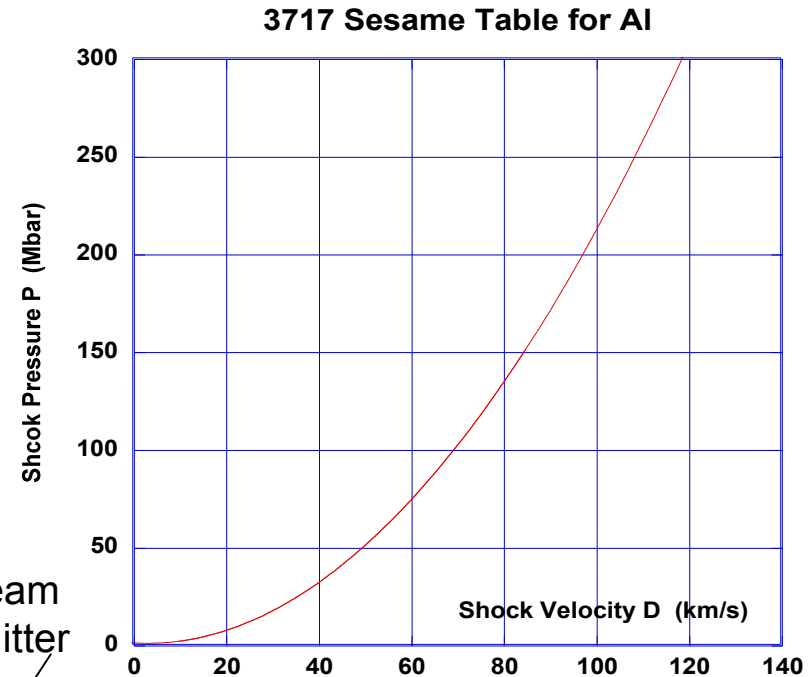
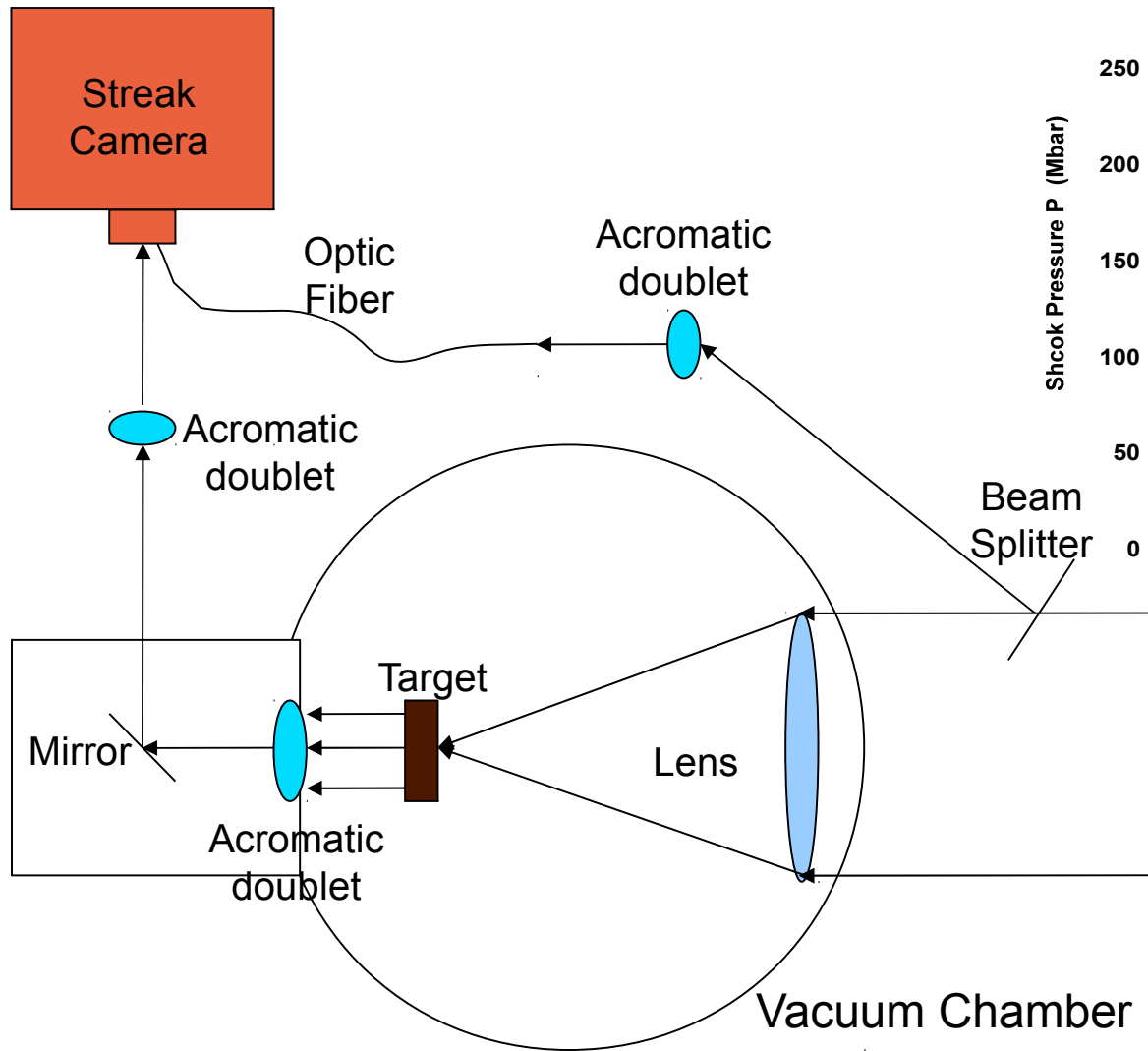


# ***EEPHC results: spectroscopy***



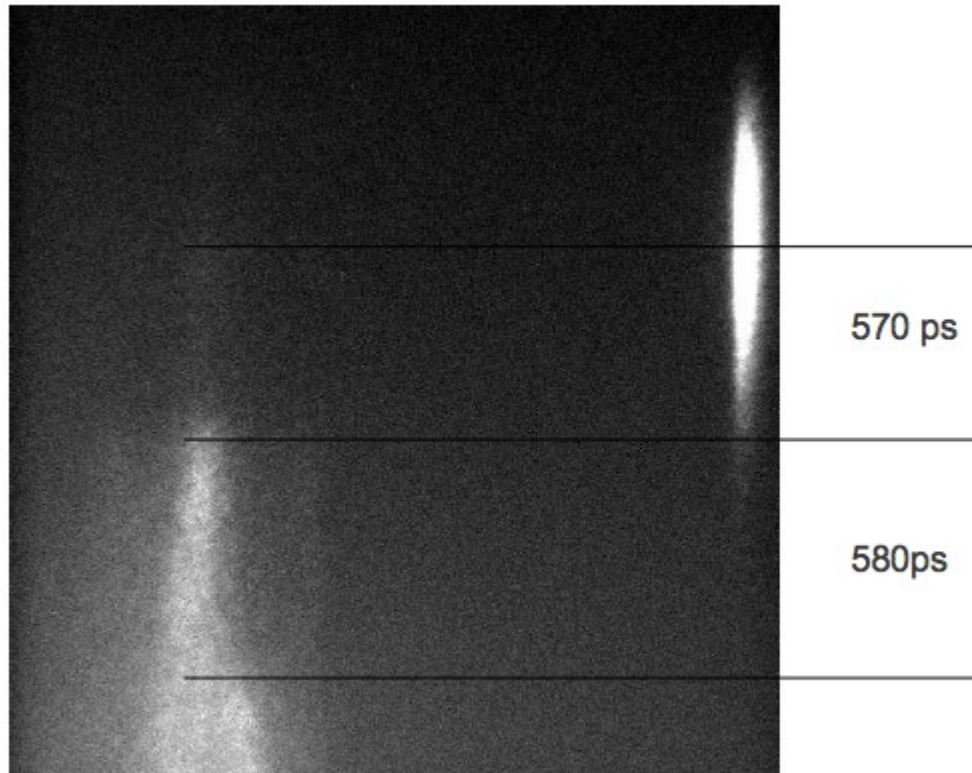
**Is this representative of the true plasma temperature? Again detailed analysis in progress...**

# Shock chronometry: set-up



**The measurement of shock velocity provides the value of shock pressure using an EOS**

# Shock chronometry



**2 diagnostics per shot:**

**$\Delta t$  between irradiation and first shock breakout**

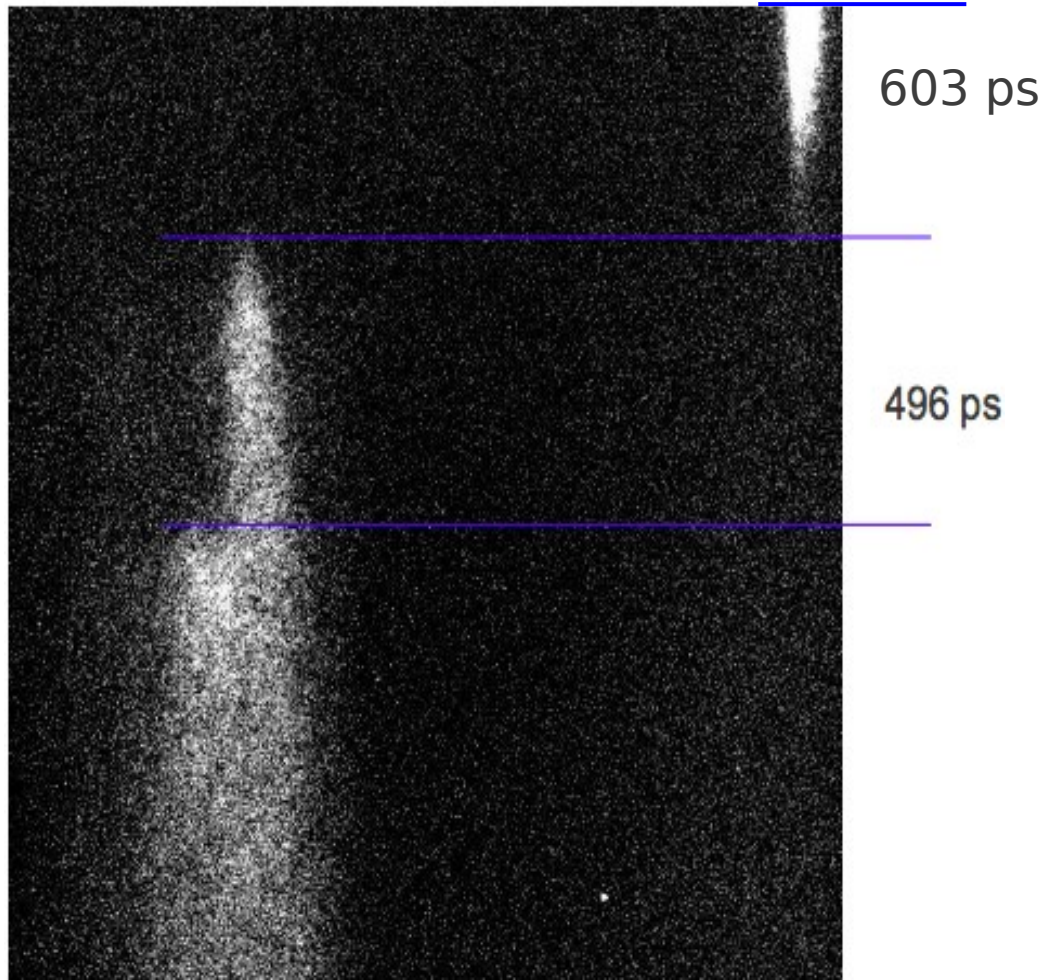
**$\Delta t$  between breakout on base and on step**

**Stepped target with  $E(3\omega) = 260$  J**

**Main only (shot 040110\_03)**

**$D = 17.3$  km/s  $\Rightarrow$   $P = 4.4$  MBar**

# Shock chronometry



Stepped target with  
 $E(3\omega) = 245 \text{ J}$   
 $E \text{ pre-pulse} = 29 \text{ J}$ ,  
delay prepulse  $500 \text{ ps}$   
 **$D = 20.2 \text{ km/s}$**   
 **$\Rightarrow P = 6.3 \text{ MBar}$**



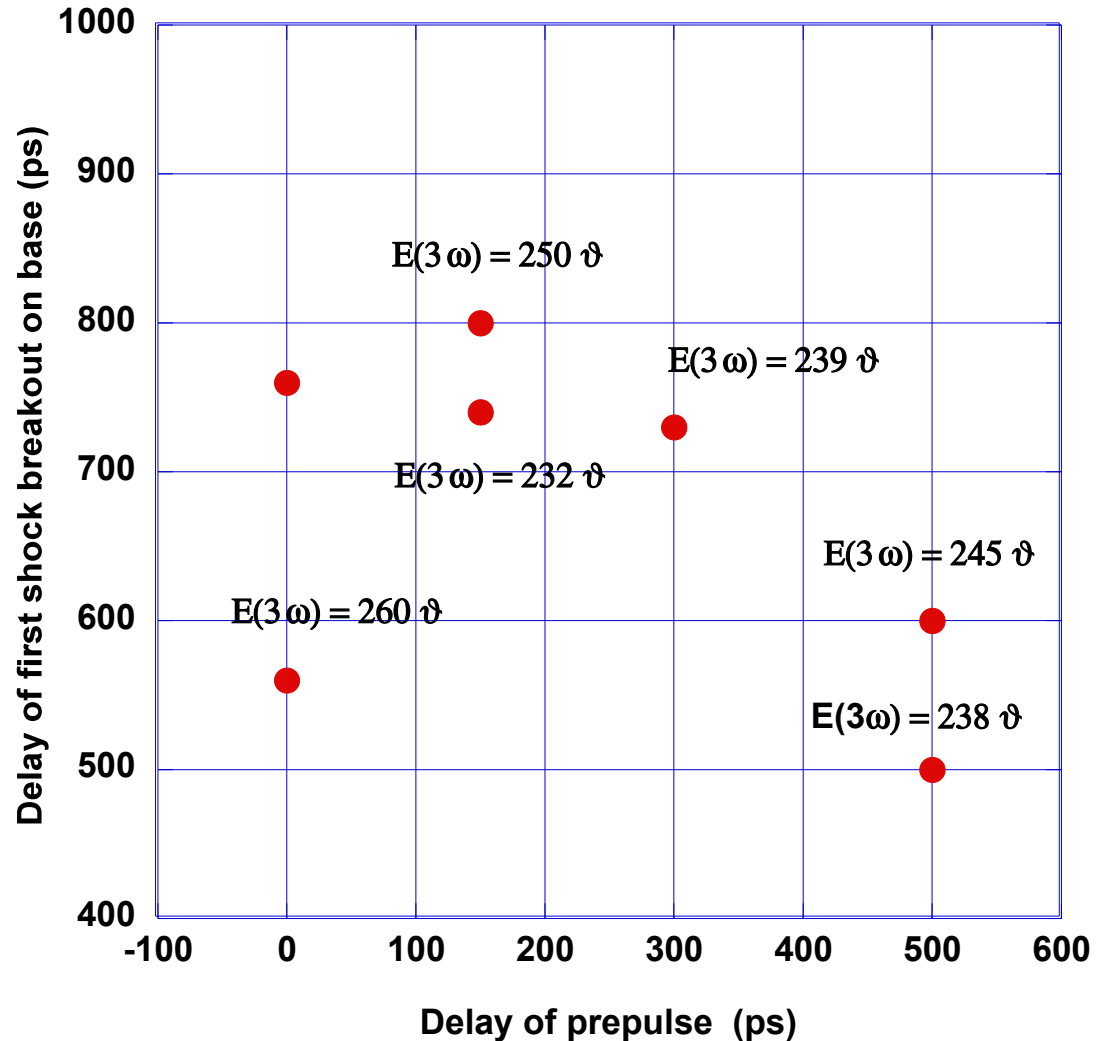
# Shock chronometry: results

	240 J	120 J	60 J
150 ps Flat	040110_05 740 ps		
150 ps Step	040110_06 804 ps		
300 ps Flat	040110_07 1023 ps		040210_13 856 ps
300 ps Step	040110_08 657 ps 040110_12 730 ps		
500 ps Flat	040110_10 498 ps	040210_06 764 ps	
500 ps Step	040210_05 603 ps		

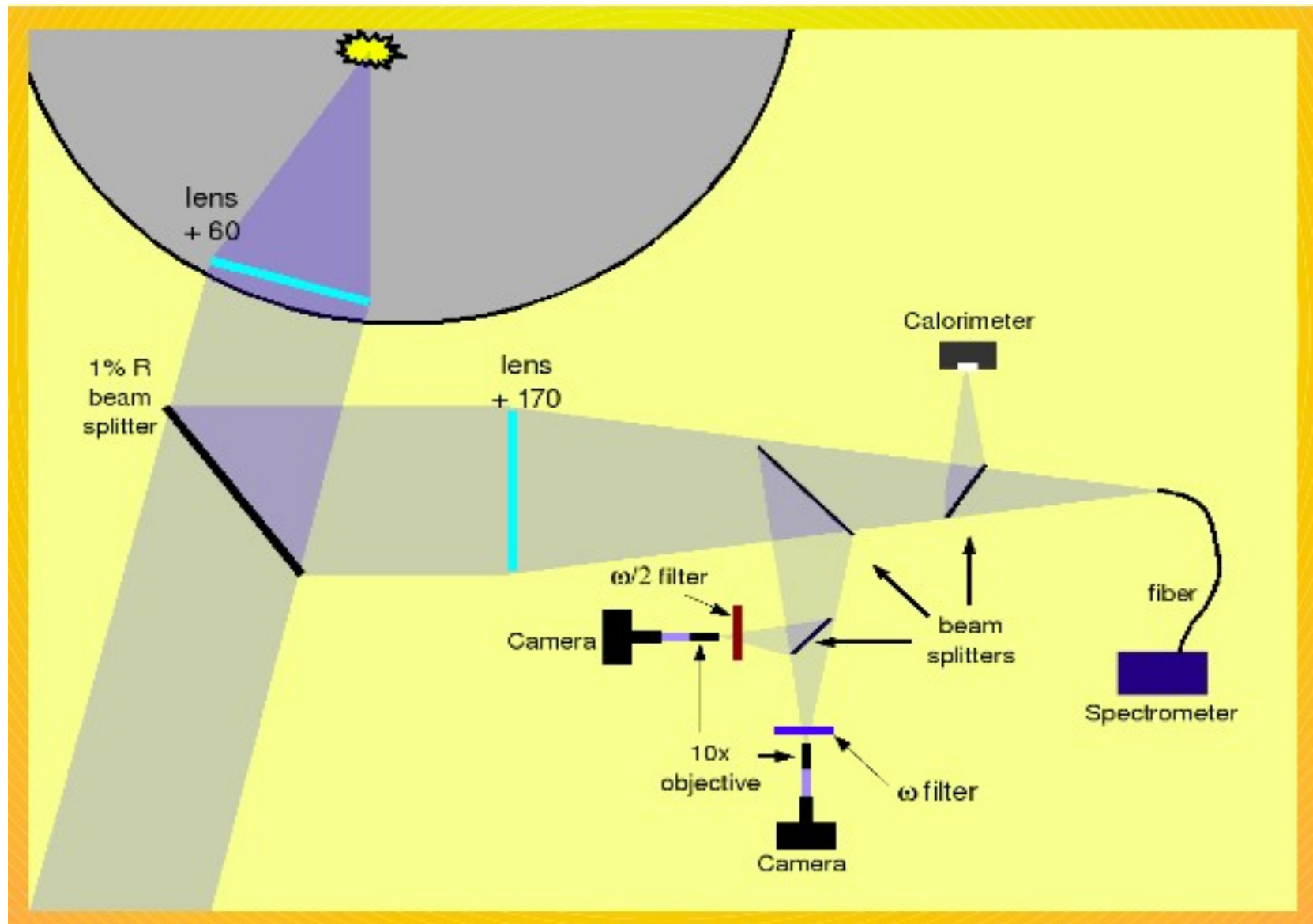
Same “bell shape” behavior of ion measurements.

Value for  $\Delta t = 500$  ps are every close to those without prepulse.

The plasma is dispersed in 500 ps? (it does not seem to be supported by XRL deflectometry data nor by SRS data)



# Back scattering diagnostics

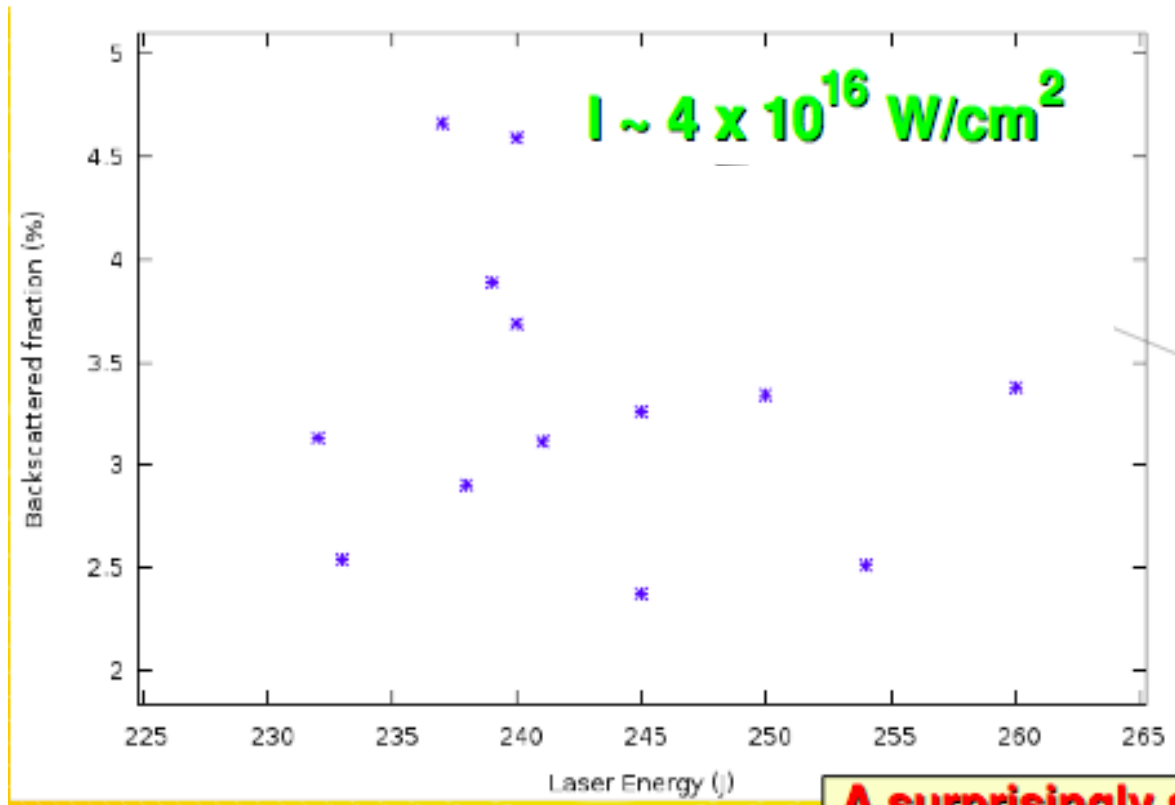


# Back scattering: calorimetry

Results at the Omega facility (Usa, 2009) give 33% back reflection at  $I \approx 8 \cdot 10^{15}$  W/cm<sup>2</sup>



**A surprisingly small fraction of light is backscattered in our experimental conditions ( $I \approx 10^{16}$  W/cm<sup>2</sup>,  $\lambda = 0.44 \mu\text{m}$ )**



Simulations (Klimo et al., PPCF, 52, 2010) also predict large SRS

**A surprisingly s**

# Back scattering: spectroscopy

## • SBS

$$\omega_{i.a.} \ll \omega_{e.m.} \Rightarrow \omega_{scatt} \sim \omega_L$$

$$c_{i.a.}(1\text{ keV}) \sim 3 \times 10^5 \text{ m/s}$$

In detail

$$\left( \frac{\Delta\lambda}{\lambda} \right)_{SBS} \sim \frac{c_{i.a.}}{c} \approx \left( \frac{\Delta\lambda}{\lambda} \right)_{res}$$

Can't distinguish SBS signal from scattered laser light

## • TPD

$$\omega_L \rightarrow \frac{\omega_L}{2} + \frac{\omega_L}{2}$$

└─▶ +  $\omega_L^{back}$  ─▶ Typical emission @  $\frac{3}{2}\omega_L$

No experimental evidence of TPD

TPD cut by window transmission?

# Back scattering: spectroscopy

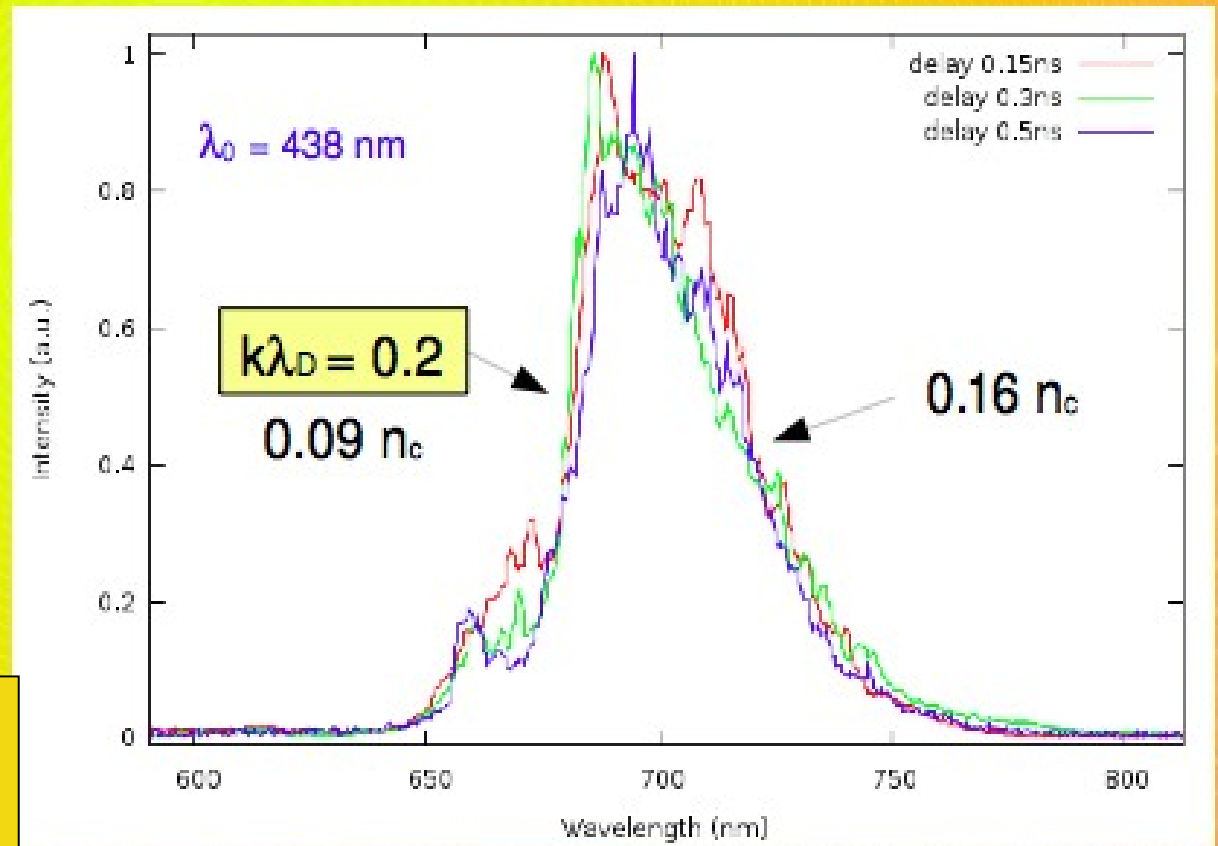
- SRS**

- Emission between  $\omega$  ( $n \sim 0$ ) and  $\omega/2$  ( $n = n_c/4$ )

$$\lambda_{SRS} = \lambda_0 [1 - (n/n_c)^{1/2} (1 + 3k^2 \lambda_D^2)^{1/2}]^{-1}$$

- Blue cut-off due to Landau damping

No SRS emission from  $n_c/4$  layer. Depletion of laser beam due to delocalised collisional absorption?



$k\lambda_D$

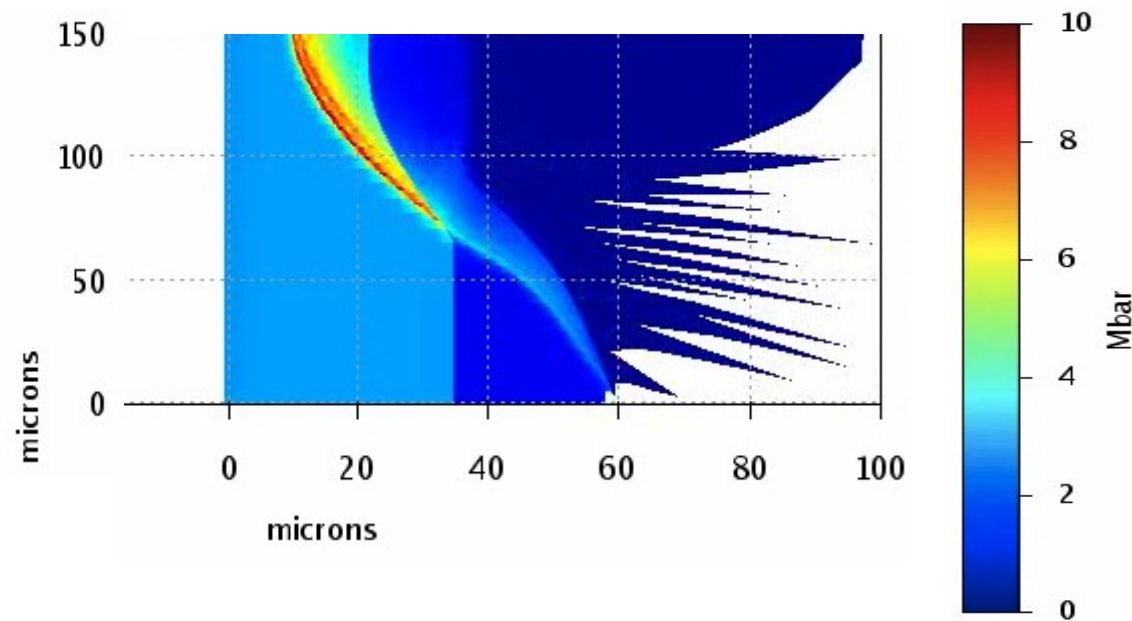
$n \rightarrow n_c/4$

**Laser-plasma coupling seems to occur at densities lower than critical**

# 2D Hydro simulations

The measured values of P at target rear corresponds to higher values at front. Indeed shock pressure undergoes a rapid decrease due to:

- 1) 2D effects during shock propagation
- 2) Relaxation waves from front side when laser turns off



**2D Simulation for  $I \approx 2 \times 10^{15} \text{ W/cm}^2$  and focal spot diameter  $\approx 100 \mu\text{m}$ .  
Breakout time = 1 ns - 0.35 ns (time of laser max in simulation) = 650 ps  
[Tommaso Vinci, code DUED]**

**Final pressure  $\approx 10 \text{ Mbar}$ , Initial Pressure  $\approx 90 \text{ Mbar}$ , in agreement with laser intensity used in simulation**

# 2D Hydro simulations

The expectation  $\approx 300$  Mbar is based on the classical model of laser absorption at the critical density

In our case the extended plasma corona likely implies delocalised absorption. Assuming absorption takes place at density  $\approx 0.16$  critical (as suggested by SRS spectra) we may expect  $\approx 160$  Mbar

However in order to reproduce the shock breakout time we need a pressure on front side  $P \approx 90$  Mbar and to get this we must either

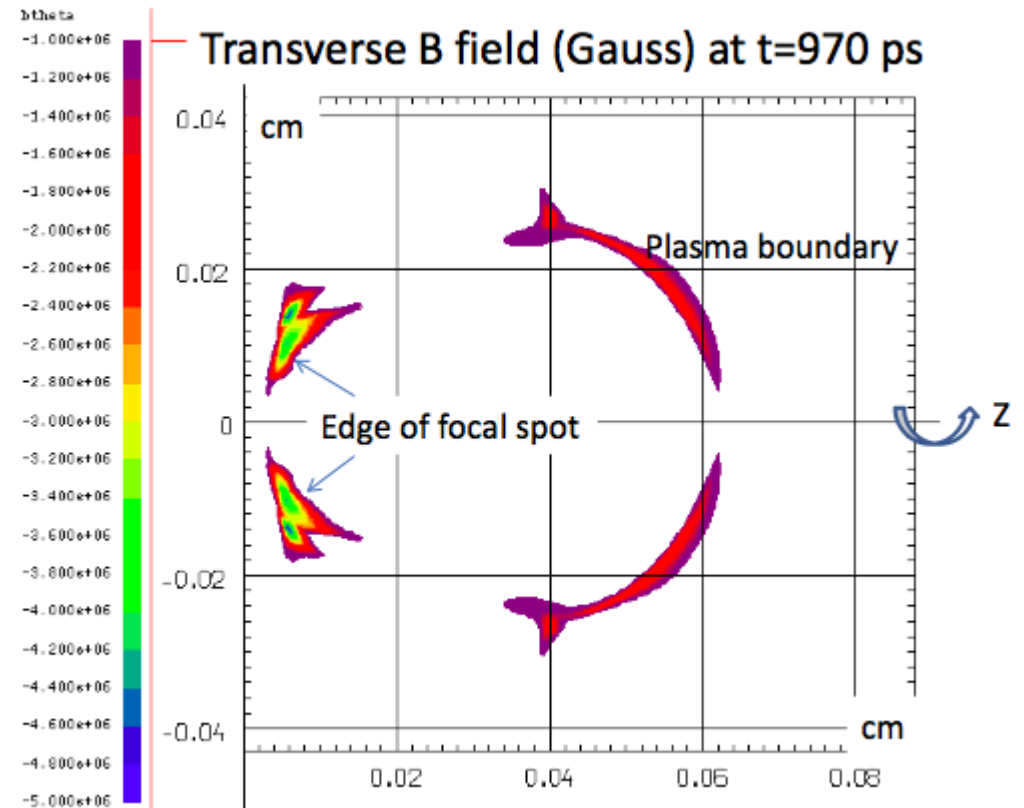
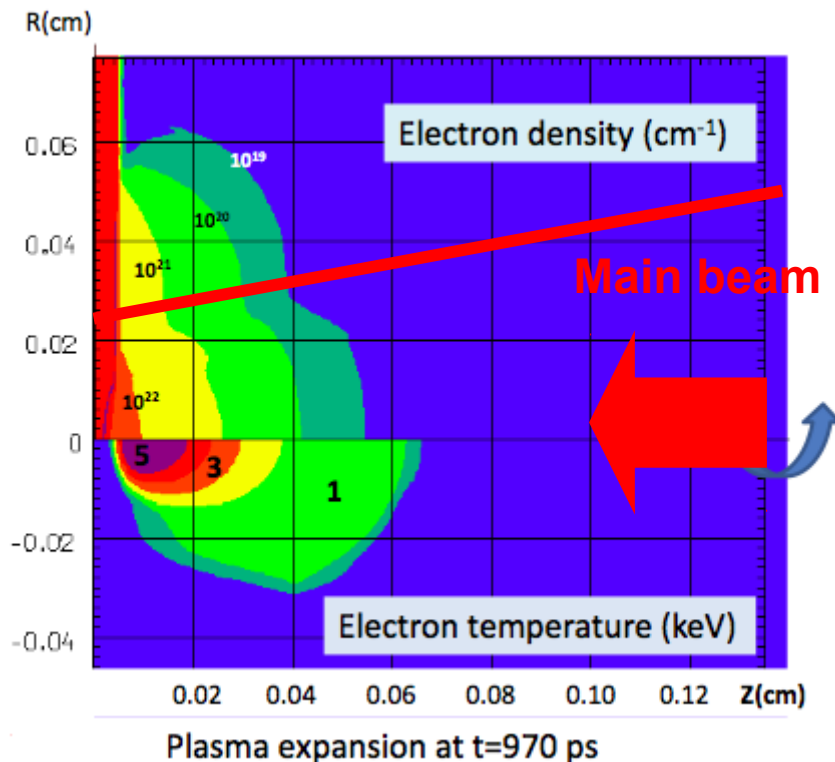
- 1) Reduce the laser intensity to  $I \approx 2 \times 10^{15}$  W/cm<sup>2</sup> with the “normal” flux limiter  $f = 0.06$
- 2) Use the nominal value 2D Simulation for  $I \approx 10^{16}$  W/cm<sup>2</sup> with a reduced flux limiter  $f = 0.01$

The flux limiter is a simple way to describe the physics of the transport of energy by electrons inside the target. Reducing  $f$  implies a different transport mechanism and an inhibition of thermal transport

Is there any possible evidence of such inhibition in our experimental situation?

# Large magnetic fields

Very large magnetic fields (up to 5 Mgauss) are created due to the  $\nabla n \times \nabla T$  mechanism in our experimental conditions (simulations by Guy Schurtz, CELIA, using CHIC)

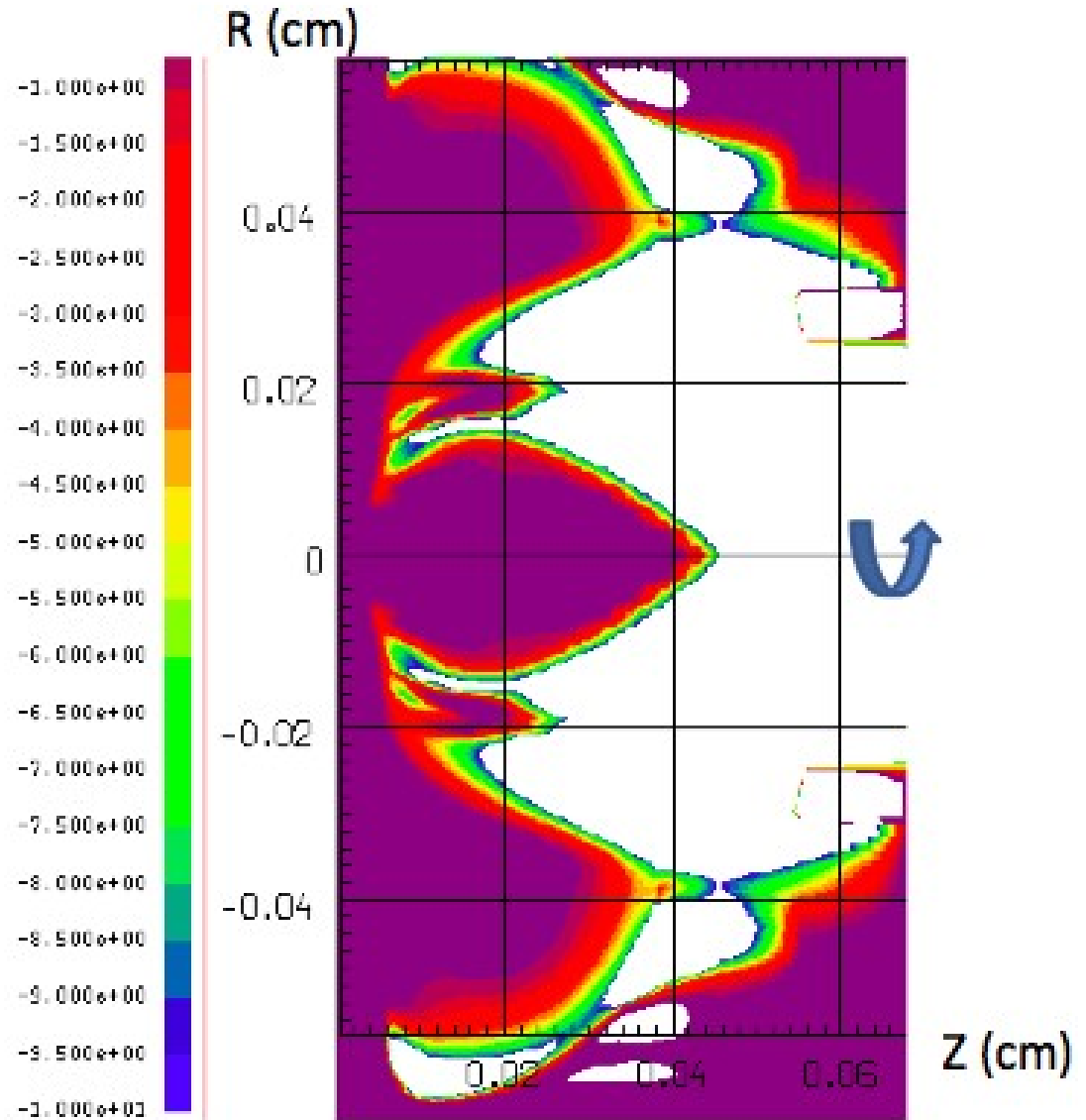




# Large magnetic fields

These imply a large value of  $\omega\tau$ , the Hall parameter, where  $\omega = eB/mc$  is the electron cyclotron frequency and  $\tau$  the collision time.

Therefore the motion of electrons is largely magnetized and the transport of heat towards the inside of the target may be inhibited



HALL PARAMETER  $\omega\tau$

# Conclusions

- Analysis of results still in progress
- The preplasma and the interaction of the main beam have been characterised using several diagnostics.
- We are able to couple a strong which is initially strong (100 Mbar) corresponding to an effective laser intensity  $\approx 2 \times 10^{15}$  W/cm<sup>2</sup>. However we were expecting  $\approx 10^{16}$  W/cm<sup>2</sup> and  $\approx 300$  Mbar. SRS spectra seem to suggest that laser couples to plasma at lower densities bringing to lower effectiveness of shock generation
- It seems that light lost via PI is surprisingly low (however we need an accurate calibration of calorimetric data due to spectral response of the apparatus). This could be in agreement with low laser wavelength, relatively low effective intensity and coupling at lower density
- Some evidence of hot electron generation by Cu K<sub>α</sub> emission (energy  $\approx 50$  keV). Need of more detailed characterization.
- SRS spectra are independent on prepulse delay, at least for  $\Delta t \leq 500$  ps
- “Bell shaped” behavior from several diagnostics. Does the preplasma become transparent for  $\Delta t = 500$  ps?



Vielen Dank