## Heavy-Light Tetraquarks with Lattice QCD

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Experimental and theoretical status of and perspectives for XYZ states
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## Motivation (1)

## Experimental background

- Experimentally observed states $Z_{b}(10610)^{+}$and $Z_{b}(10650)^{+}$
- Mass suggests a bottomonium state $\bar{b} b$ but would be electrically neutral
$\Rightarrow$ Quantum numbers can be described with four-quark structure


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## Theoretical study

- We study similar but less challenging systems
- Quark content: $\bar{Q} \bar{Q}^{\prime} q q^{\prime}$, here: $\bar{b} \bar{b} u d, \bar{b} \bar{b} u s, \bar{b} \bar{c} u d$
- In the limit $m_{Q} \rightarrow \infty$ stable tetraquark was shown
[J. Carlson, L. Heller and J. A. Tjon, Phys. Rev. D 37,744 (1988)]
[A. V. Manohar and M. B. Wise, Nucl. Phys. B 399, 17 (1993)]
[E. J. Eichten and C. Quigg, Phys. Rev. Lett. 119, no. 20, 202002 (2017)]
[M. Karliner and J. L. Rosner, Phys. Rev. Lett. 119, no.20, 202001 (2017)]


## Motivation (2)

Born-Oppenheimer study of doubly-heavy tetraquarks:

- i.e. static heavy quarks ( $\bar{b}$-quarks)
- Prediction of a bound tetraquark in $\bar{b} \bar{b} u d$ sector with $I\left(J^{P}\right)=0\left(1^{+}\right)$and $M_{\bar{b} \bar{b} u d}-\left(M_{B}+M_{B^{*}}\right) \approx-90 \mathrm{MeV}$
[Z. S. Brown and K. Orginos, Phys. Rev. D 86, 114506 (2012)]
[P. Bicudo et al. [ETMC], Phys. Rev. D 87, no. 11, 114511 (2013)]
[P. Bicudo, K. Cichy, A. Peters, B. Wagenbach, M. Wagner, Phys. Rev. D 92, no. 1, 014507 (2015)]
[P. Bicudo, J. Scheunert and M. Wagner, Phys. Rev. D 95, no. 3, 034502 (2017)]


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- Evidence for a $\bar{b} \bar{b} u d$ resonance in the $I\left(J^{P}\right)=0\left(1^{-}\right)$channel with $M_{\bar{b} \bar{b} u d}-\left(M_{B}+M_{B}\right) \approx+20 \mathrm{MeV}, \Gamma \approx 100 \mathrm{MeV}$
[P. Bicudo, M. Cardoso, A. Peters, M.P. and M. Wagner, Phys. Rev. D 96, no. 5, 054510 (2017)]


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Searching for doubly-heavy tetraquark bound states in full lattice QCD using Non-Relativistic QCD:

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[A. Francis, R. J. Hudspith, R. Lewis and K. Maltman, Phys. Rev. Lett. 118, no. 14, 142001 (2017)]
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- For $\bar{b} \bar{c} u d$, the predictions are not as clear $\rightarrow$ Might be weakly bound or no binding
- In our study: We apply a more extended operator basis
$\rightarrow$ Enables a better treatment of threshold states.
$\rightarrow$ More on next slides

[^0]
## Interpolating Operators (1)

## Investigated systems and quantum numbers

- $\bar{b} \bar{b} u d$ with $I\left(J^{P}\right)=0\left(1^{+}\right)$
$\rightarrow$ Most promising as it is closest to $\bar{Q} \bar{Q} q q$ with $m_{Q} \rightarrow \infty$
- $\bar{b} \bar{b} u s$ with $I\left(J^{P}\right)=\frac{1}{2}\left(1^{+}\right)$
$\rightarrow$ Slightly less promising as $d$ replaced by $s$
- $\bar{b} \bar{c} u d$ with $I\left(J^{P}\right)=0\left(1^{+}\right)$and $I\left(J^{P}\right)=0\left(0^{+}\right)$
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## For all systems, we use two types of interpolating operators

- Local operators; basically used in all previous studies
- Nonlocal operators; unique compared to all other studies on heavy-light tetraquarks


## Interpolating Operators (2)

## - Local operators:

- Four quarks at the same space-time position
- Jointly projected to zero momentum
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- Describe mesonic scattering structure
- Expectation:
- Local operators: good overlap to ground state (stable four-quark)
- Nonlocal operators: sizable overlap to first excited state (2 meson state)
$\Rightarrow$ Isolate ground state from higher excitations, especially first excited state


## Interpolating Operators for $\bar{b} \bar{b} u d$ and $\bar{b} \bar{b} u s$

## Interpolating Operators for $\bar{b} \bar{b} u d$

$$
I\left(J^{P}\right)=0\left(1^{+}\right)
$$

| relevant thresholds | $B^{*} B, B^{*} B^{*}(\approx+45 \mathrm{Mev})$ |
| :---: | :---: |
| local operators | $B^{*} B, B^{*} B^{*}$, diquark-antidiquark |
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## Interpolating Operators for $\bar{b} \bar{b} u s$

$$
I\left(J^{P}\right)=\frac{1}{2}\left(1^{+}\right)
$$

relevant thresholds $B^{*} B_{s}, B B_{s}^{*}(\approx$ equal $), B^{*} B_{s}^{*}(\approx+45 \mathrm{Mev})$ local operators $\quad B^{*} B_{s}, B B_{s}^{*}, B^{*} B_{s}^{*}$, diquark-antidiquark
nonlocal operators $\quad B^{*} B_{s}, B B_{s}^{*}, B^{*} B_{s}^{*}$

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$$
I\left(J^{P}\right)=0\left(1^{+}\right)
$$

| relevant thresholds | $B^{*} D, B D^{*}(\approx+95 \mathrm{Mev}), B^{*} D^{*}(\approx+140 \mathrm{Mev})$ |
| :---: | :---: |
| local operators | $B^{*} D, B D^{*}$, diquark-antidiquark |
| nonlocal operators | $B D^{*}, B^{*} D$ |
|  | $I\left(J^{P}\right)=0\left(0^{+}\right)$ |
| relevant thresholds | $B D, B^{*} D^{*}(\approx+185 \mathrm{Mev})$ |
| local operators | $B D$, diquark-antidiquark |
| nonlocal operators | $B D$ |

## Energy Spectrum for the $\bar{Q} \bar{Q}^{\prime} q q^{\prime}$ system

- Due to point-to-all propagators, only non-symmetric correlation matrix available (no scattering operator at source)
- Apply multi-exponential matrix fitting: employable also for non-symmetric matrices

$$
C_{j k}(t) \approx \sum_{n=0}^{N-1} Z_{j}^{n} Z_{k}^{n} \mathrm{e}^{-E_{n} t}, \quad \begin{gathered}
E_{n}: n \text {-th energy eigenvalue } \\
Z_{j}^{n}=\langle\Omega| \mathcal{O}_{j}|n\rangle: \text { overlap factor }
\end{gathered}
$$



Schematic representation of Wick contractions for different correlation matrix elements

## Fit Results for $\bar{b} \bar{b} u d$



Results for the lowest two $\bar{b} \bar{b} u d$ energy levels relative to the $B B^{*}$ threshold. Black box: local operator included. Red box: scattering operator included.

- Found evidence for bound state with $E_{\text {binding }}=-128 \mathrm{MeV}$
- First excited state corresponds to threshold


## Preliminary Results for $\bar{b} \bar{b} u s$



- Found evidence for bound state with $E_{\text {binding }} \approx-80 \mathrm{MeV}$
- First excited state corresponds to threshold


## Preliminary Results for $\bar{b} \bar{c} u d$



- No evidence for bound states in $\bar{b} \bar{c} u d$ systems
- Lowest energy level corresponds to threshold


## Overlap Factors for $\bar{b} \bar{b} u d$

For fixed $j: Z_{j}^{n}$ indicates relative importance of energy eigenstates $|n\rangle$

$$
\mathcal{O}_{j}^{\dagger}|\Omega\rangle=\sum_{n=0}^{\infty}|n\rangle\langle n| \mathcal{O}_{j}^{\dagger}|\Omega\rangle=\sum_{n=0}^{\infty} Z_{j}^{n}|n\rangle
$$







The normalized overlap factors $\left|\tilde{Z}_{j}^{n}\right|^{2}=\frac{\left|Z_{j}^{n}\right|^{2}}{\max _{m}\left(\left|Z_{j}^{m}\right|^{2}\right)}$ as determined on ensemble C005.

## Scattering Analysis and Chiral Extrapolation for $\bar{b} \bar{b} u d$

## Scattering Analysis

- Relate finite volume energy spectrum $E_{n}$ to infinite volume scattering amplitude
- Use Lüscher's formula to determine phase shift and infinite volume binding energy
- Confirmation that ground state is stable tetraquark.


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## Chiral Extrapolation



Fit of the pion-mass dependence of $E_{\text {binding. }}$. The vertical dashed line indicates the physical pion mass.

$$
\begin{aligned}
& E_{\text {binding }}\left(m_{\pi, \text { phys }}\right)=(-128 \pm 24 \pm 10) \mathrm{MeV} \\
& m_{\text {tetraquark }}\left(m_{\pi, \text { phys }}\right)=(10476 \pm 24 \pm 10) \mathrm{MeV}
\end{aligned}
$$

## Comparison of Different Results for $\bar{b} \bar{b} u d$



Comparison of $\bar{b} \bar{b} u d$ tetraquark binding energies with $I\left(J^{P}\right)=0\left(1^{+}\right)$(black: this work; blue: lattice NRQCD; red: lattice QCD computations of static $\bar{b} \bar{b}$ potentials and solving the

Schrödinger equation; green: effective field theories and potential models).

## Summary

- Study bound states in doubly heavy tetraquarks
- Consider local and nonlocal interpolating operators
- Apply a finite volume Lüscher analysis


## Summary

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- Predict a bound state in the $\bar{b} \bar{b} u d$ channel with in $I\left(J^{P}\right)=0\left(1^{+}\right)$ with $E_{\text {binding }}=(-128 \pm 24 \pm 10) \mathrm{MeV}$
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## Outlook

- Finalize evaluation of $\bar{b} \bar{b} u s$ outcomes
- More detailed analysis of threshold states in $\bar{b} \bar{c} u d \rightarrow$ candidates for resonances?


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## Thank You for Your Attention!

## Lattice Setup

- Use gauge link configuration generated by RBC and UKQCD collaboration
[Y. Aoki et al. [RBC and UKQCD Collaborations], Phys. Rev. D 83, 074508 (2011)]
[T. Blum et al. [RBC and UKQCD Collaborations],Phys. Rev. D 93, no. 7, 074505 (2016)]
- $2+1$ flavours domain-wall fermions and Iwasaki gauge action
- Five different ensembles which differ in

$$
\begin{array}{ll}
\text { lattice spacing } & a \quad \approx 0.083 \mathrm{fm} \ldots 0.114 \mathrm{fm}, \\
\text { lattice size } & L \\
\text { pion mass } & m_{\pi} \approx 2.65 \mathrm{fm} \ldots 5.48 \mathrm{fm}, \\
\Rightarrow \text { explore dependence on } L, m_{\pi}
\end{array}
$$

- Smeared point-to-all propagators for the up and down quarks


## Scattering Analysis

- Relate finite volume energy spectrum $E_{n}$ to infinite volume scattering amplitude for 2 energy levels in $T_{1}^{+}$ irrep
- Use Lüscher's formula and scattering momenta $k_{n}^{2}$ to determine phase shift
- Apply effective-range-expansion (ERE)


Plot of the effective-range-expansion for C005.
Blue curve: $a k \cot (\delta(k))+|a k|$.
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- Search bound state pole of scattering amplitude below threshold at

$$
\cot \delta_{0}\left(k_{\mathrm{BS}}\right)=i, \quad \text { so: } \quad-\left|k_{\mathrm{BS}}\right|=\frac{1}{a_{0}}-\frac{1}{2} r_{0}\left|k_{\mathrm{BS}}\right|^{2}
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$$

- Results essentially identical to the finite-volume energy levels
- Confirmation that ground state is stable tetraquark.


## GEP-results



Effective masses $a E_{\text {eff }, n}$ for $n=0,1$ as a function of $t / a$ for a $3 \times 3$ correlation matrix $\underline{\underline{\underline{1}}}=$


[^0]:    [A. Francis, R. J. Hudspith, R. Lewis and K. Maltman, Phys. Rev. Lett. 118, no. 14, 142001 (2017)]
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