



# $Z_{cs}^+$ exotics in $B^+ \rightarrow J/\psi \phi K^+$ decays at LHCb

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Experimental and theoretical status of and perspectives for XYZ states

# $Z_c$ and predictions of $Z_{cs}$

- ◆ Some  $Z_c$  states were observed from  $B$  decays or  $Y(4260)$
- ◆ Several papers have predicted the existence of  $Z_{cs}$  states in early time:

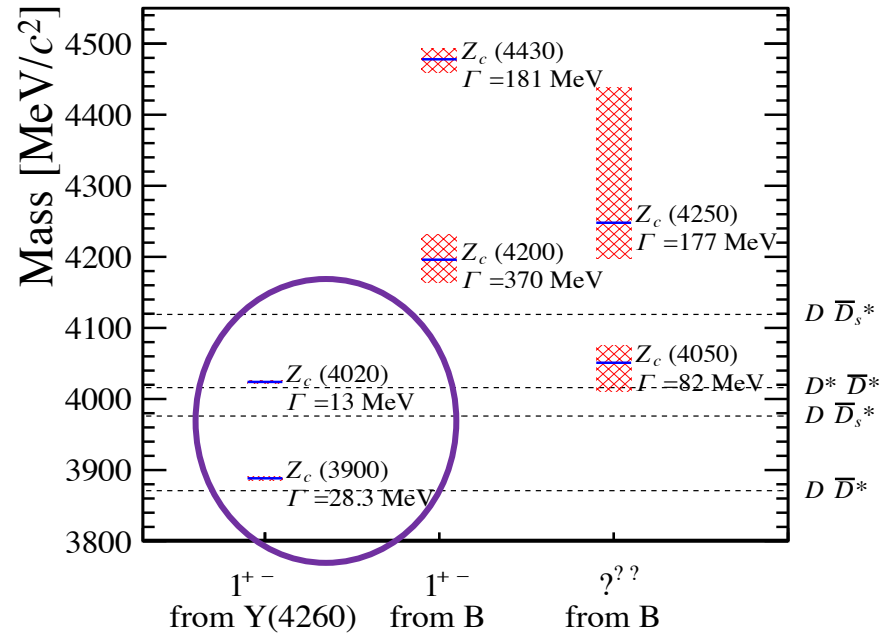
[J. Korean Phys. Soc. 55 (2009) 424]

[Phys. Rev. Lett. 110 (2013), no. 23 232001]

[Phys. Rev. D 88 (2013), no. 9 096014]

[Phys. Lett. B 798 (2019) 135022]

- ◆ The  $Z_{cs}$  states would be useful to distinguish different models:
  - Less exchange particles expected in the  $Z_{cs}$  molecule picture



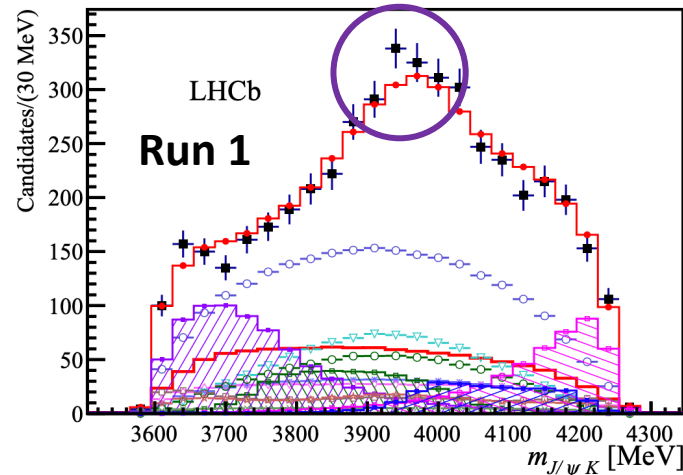
$Z_c$  states from  $Y(4260)$   
are narrow;  
Others from  $B$  decays are  
broad

Finally, as pointed out in ref. [13], the meson-meson molecular model cannot be used to describe heavy-light tetraquarks with non-null strangeness content. The reason is that one-pion-exchange cannot take place between strange and nonstrange heavy mesons, like  $B$  and  $B_s$ . Hidden-charm and bottom mesons with strangeness are also forbidden in the

[JHEP 04 (2020) 119]

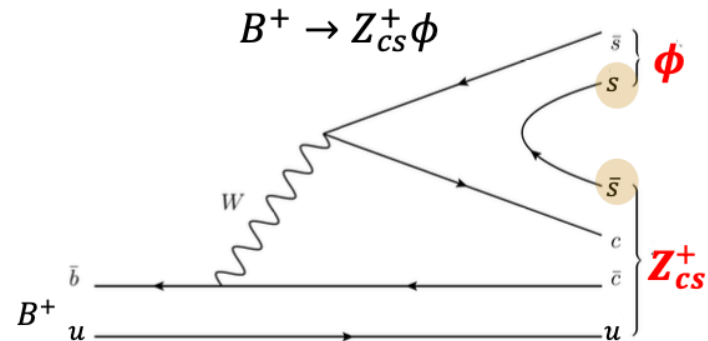
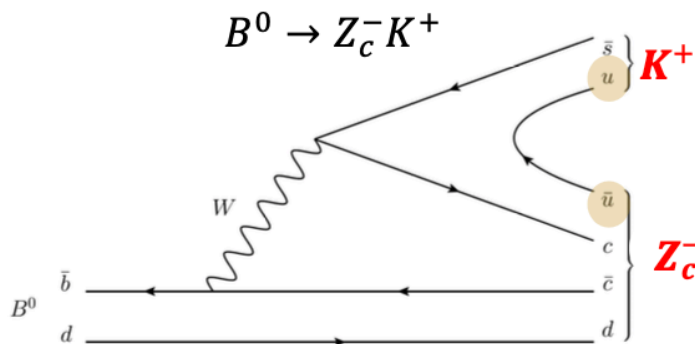
# Search for $Z_{cs}^+$ in $B^+ \rightarrow J/\psi \phi K^+$ decay

- First amplitude analysis for  $B^+ \rightarrow J/\psi \phi K^+$  was studied at LHCb using Run 1 sample.
- Hint of  $J/\psi K^+$  structure in Run 1 analysis, but not significant



[PRL. 118 (2017) 022003]

- $Z_{cs}^+$  in  $B^+ \rightarrow J/\psi \phi K^+$  decay has similar topology as  $Z_c^+$  in  $B^+ \rightarrow J/\psi K^- \pi^+$  decay, and  $P_c^+$  in  $\Lambda_b^0 \rightarrow J/\psi p K^-$



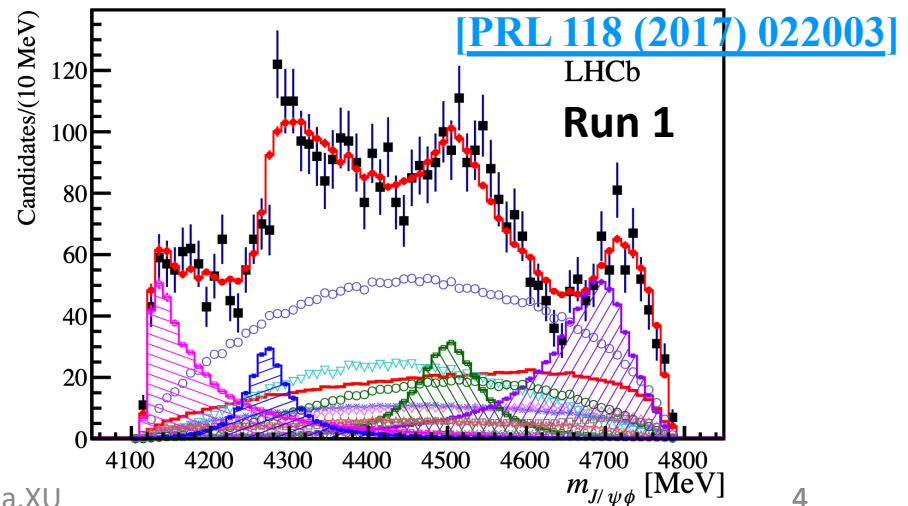
# $X$ states in $B^+ \rightarrow J/\psi \phi K^+$ decay

- The width of  $X(4140)$  is  $83 \pm 21^{+21}_{-14}$  MeV, larger than the value measured from other experiments.

[PRD 95 (2017) 012002]

Year	Experiment luminosity	$B \rightarrow J/\psi \phi K$ yield	Mass [MeV]	$X(4140)$ peak		
				Width [MeV]	Sign.	Fraction %
2008	CDF $2.7 \text{ fb}^{-1}$ [1]	$58 \pm 10$	$4143.0 \pm 2.9 \pm 1.2$	$11.7^{+8.3}_{-5.0} \pm 3.7$	$3.8\sigma$	
2009	<i>Belle</i> [22]	$325 \pm 21$	$4143.0$ fixed	$11.7$ fixed	$1.9\sigma$	
2011	CDF $6.0 \text{ fb}^{-1}$ [29]	$115 \pm 12$	$4143.4^{+2.9}_{-3.0} \pm 0.6$	$15.3^{+10.4}_{-6.1} \pm 2.5$	$5.0\sigma$	$14.9 \pm 3.9 \pm 2.4$
2011	LHCb $0.37 \text{ fb}^{-1}$ [21]	$346 \pm 20$	$4143.4$ fixed	$15.3$ fixed	$1.4\sigma$	$< 7$ @ 90%CL
2013	CMS $5.2 \text{ fb}^{-1}$ [25]	$2480 \pm 160$	$4148.0 \pm 2.4 \pm 6.3$	$28^{+15}_{-11} \pm 19$	$5.0\sigma$	$10 \pm 3$ (stat.)
2013	D0 $10.4 \text{ fb}^{-1}$ [26]	$215 \pm 37$	$4159.0 \pm 4.3 \pm 6.6$	$19.9 \pm 12.6^{+1.0}_{-8.0}$	$3.0\sigma$	$21 \pm 8 \pm 4$
2014	BaBar [24]	$189 \pm 14$	$4143.4$ fixed	$15.3$ fixed	$1.6\sigma$	$< 13.3$ @ 90%CL
2015	D0 $10.4 \text{ fb}^{-1}$ [27]	$p\bar{p} \rightarrow J/\psi \phi \dots$	$4152.5 \pm 1.7^{+6.2}_{-5.4}$	$16.3 \pm 5.6 \pm 11.4$	$4.7\sigma$ ( $5.7\sigma$ )	
Average			$4147.1 \pm 2.4$	$15.7 \pm 6.3$		

- Three other  $J/\psi \phi$  structures,  $X(4274)$ ,  $X(4500)$  and  $X(4700)$  were observed in Run 1 analysis.

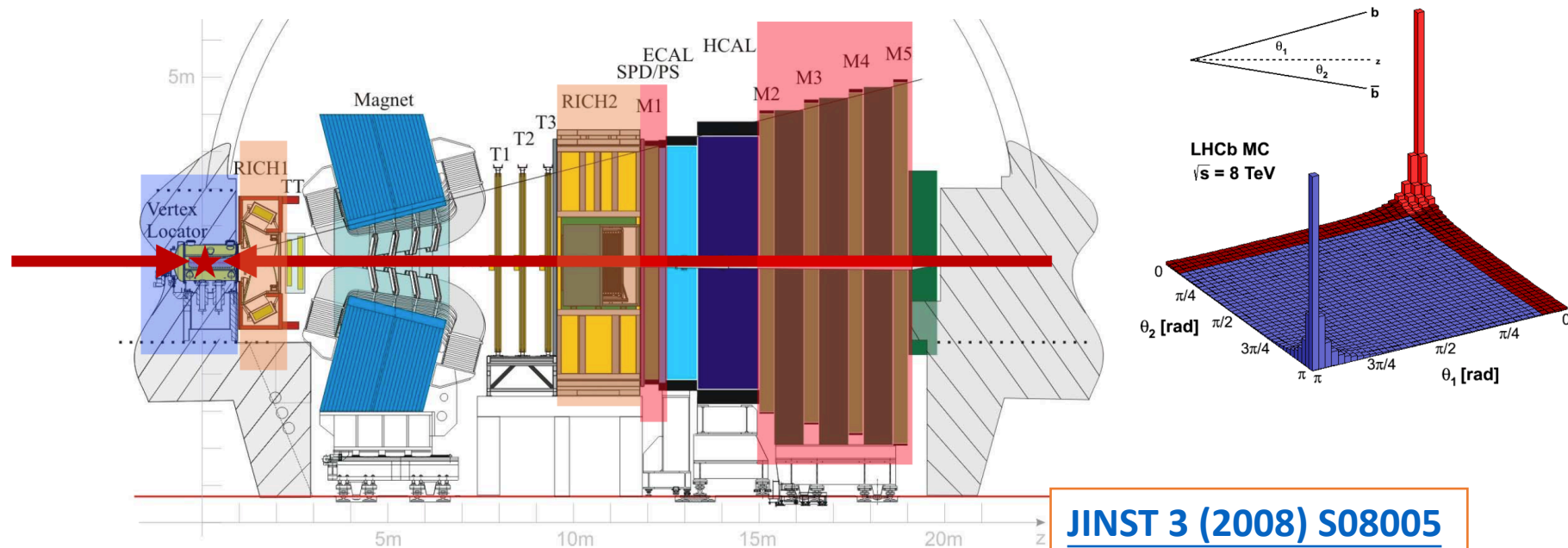




# LHCb detector

LHCb is a dedicated heavy flavor physics experiment at LHC

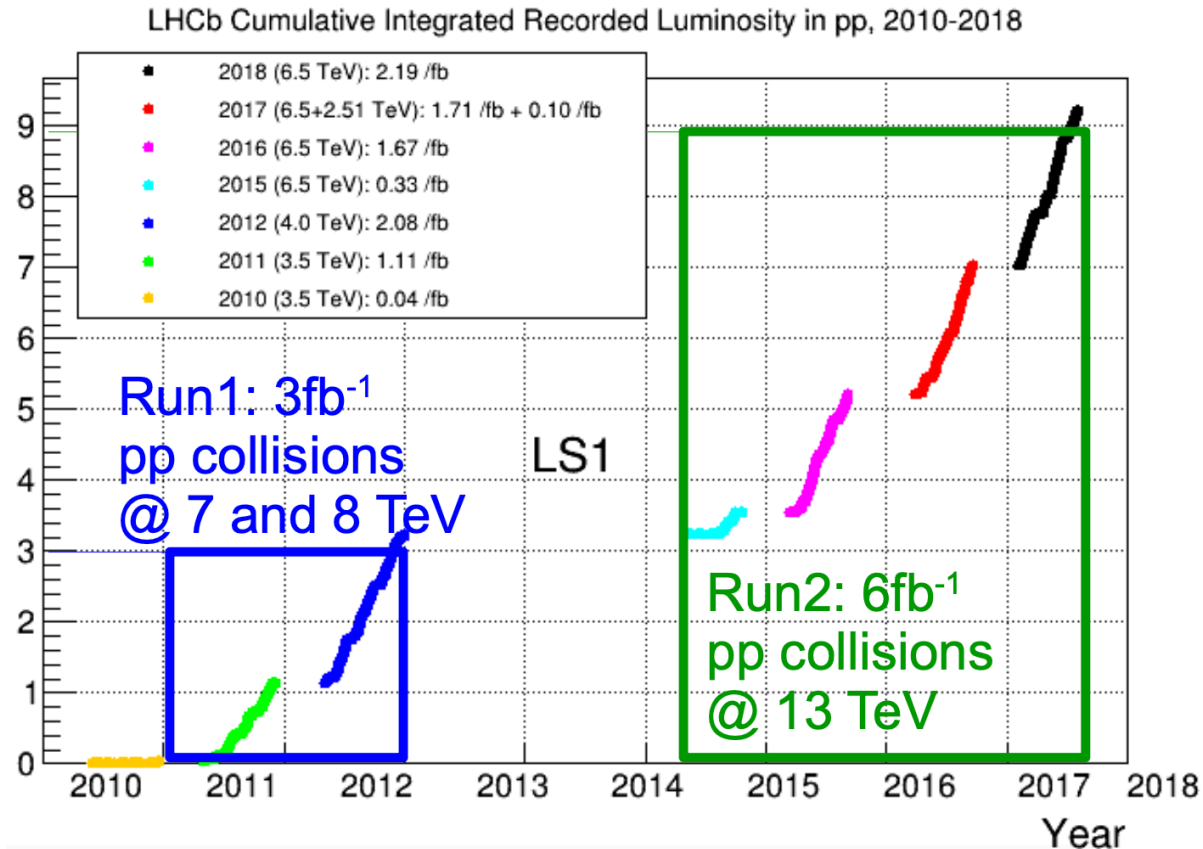
- $\sim 20,000/s$   $b\bar{b}$  generated at LHCb point in Run 2
- A single-arm forward region spectrometer covering  $2 < \eta < 5$



<b>Vertex:</b>	$\sigma_{IP} = 20 \mu\text{m}$
<b>Time:</b>	$\sigma_{\tau} = 45 \text{ fs}$ for $B_s^0 \rightarrow J/\psi\phi$ or $D_s^+\pi^-$
<b>Momentum:</b>	$\Delta p/p = 0.4 \sim 0.6\%$ (5 – 100 GeV/c)
<b>Mass :</b>	$\sigma_m = 8 \text{ MeV}/c^2$ for $B \rightarrow J/\psi X$ (constrained $m_{J/\psi}$ )
<b>Hadron ID:</b>	$\varepsilon(K \rightarrow K) \sim 95\%$ mis-ID $\varepsilon(\pi \rightarrow K) \sim 5\%$
<b>Muon ID:</b>	$\varepsilon(\mu \rightarrow \mu) \sim 97\%$ mis-ID $\varepsilon(\pi \rightarrow \mu) \sim 1 - 3\%$
<b>ECAL:</b>	$\Delta E/E = 1 \oplus 10\%/\sqrt{E \text{ (GeV)}}$

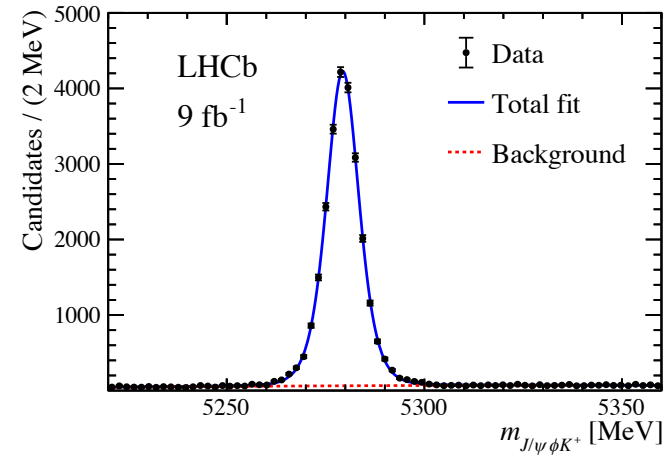
# LHCb data taking

Large data sample used to study  $B^+ \rightarrow J/\psi \phi K^+$  at LHCb

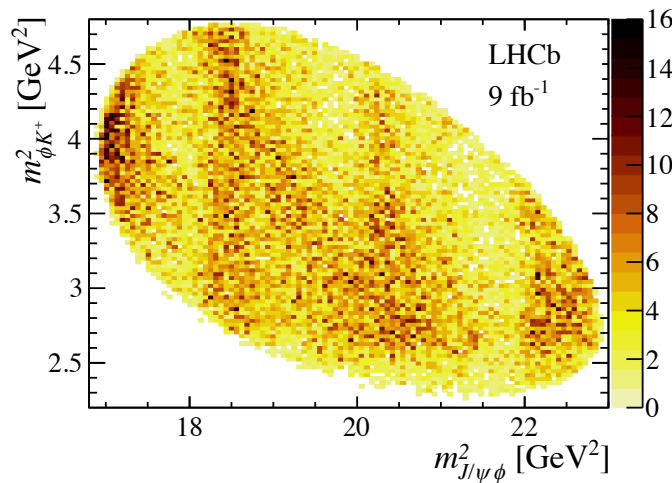


[Taken from LHCb public website]

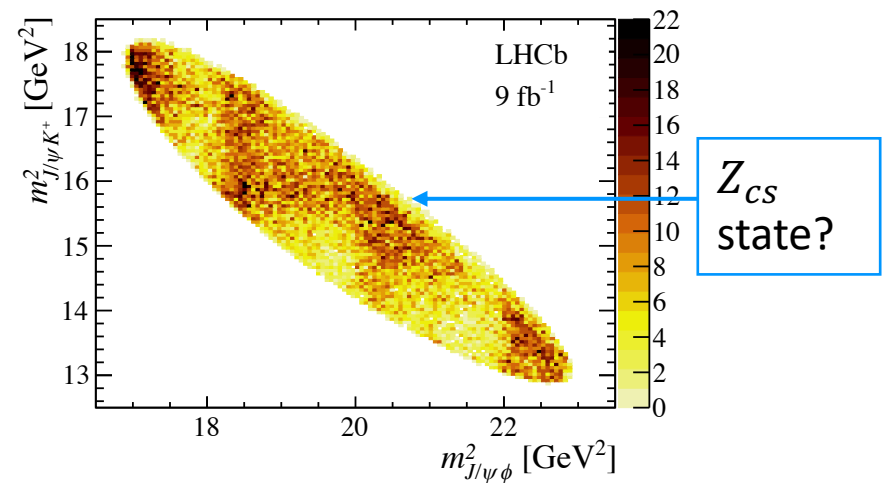
- Selections are optimized
  - $\sim 24\text{k}$  signal (6 times larger than the previous publication)
  - overall background fraction  $\sim 4\%$  (a factor of 6 smaller)



- Clear structures in Dalitz plot



Several  $X$  states along  $m_{J/\psi\phi}$ .



**Clearly visible: 4 structures in  $J/\psi\phi$  and an obvious structure in  $J/\psi K$**

# 6-D amplitude fit approach

- **Signal** and **background** components in PDF:

$$\begin{aligned}
 -\ln L(\vec{\omega}) &= -\sum_i \ln [(1 - \beta) \mathcal{P}_{\text{sig}}(m_{\phi K \ i}, \Omega_i | \vec{\omega}) + \beta \mathcal{P}_{\text{bkg}}(m_{\phi K \ i}, \Omega_i)] \\
 &= -\sum_i \ln \left[ (1 - \beta) \frac{|\mathcal{M}(m_{\phi K \ i}, \Omega_i | \vec{\omega})|^2 \Phi(m_{\phi K \ i}) \epsilon(m_{\phi K \ i}, \Omega_i)}{I(\vec{\omega})} + \beta \frac{\mathcal{P}_{\text{bkg}}^u(m_{\phi K \ i}, \Omega_i)}{I_{\text{bkg}}} \right] \\
 &= -\sum_i \ln \left[ |\mathcal{M}(m_{\phi K \ i}, \Omega_i | \vec{\omega})|^2 + \frac{\beta I(\vec{\omega})}{(1 - \beta) I_{\text{bkg}}} \frac{\mathcal{P}_{\text{bkg}}^u(m_{\phi K \ i}, \Omega_i)}{\Phi(m_{\phi K \ i}) \epsilon(m_{\phi K \ i}, \Omega_i)} \right] + N \ln I(\vec{\omega}) .
 \end{aligned}$$

- Each decay chain is described by 6 observables:

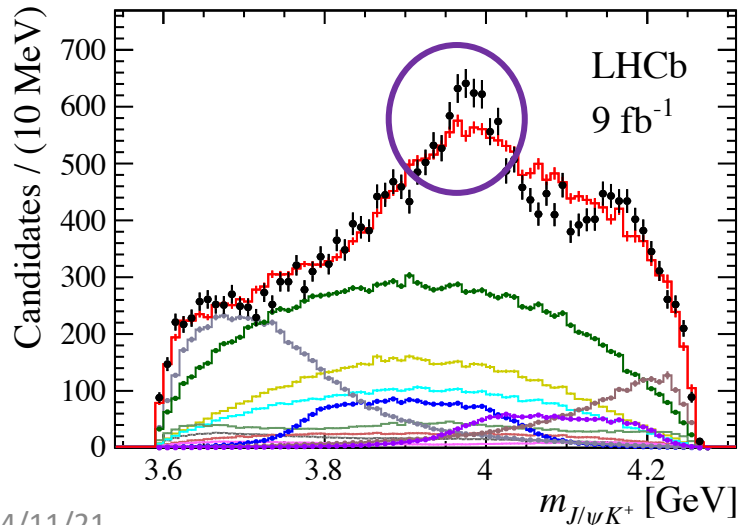
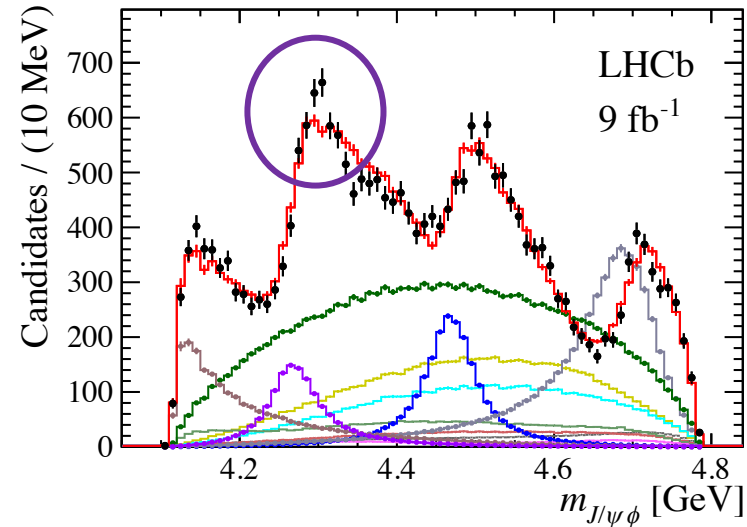
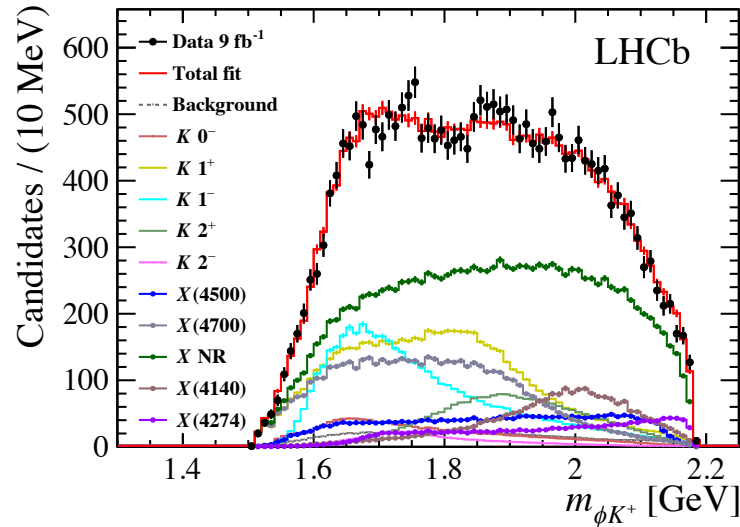
$$\Phi \equiv (m_{\phi K}, \theta_{K^*}, \theta_{J/\psi}, \hat{\theta}_{\phi}, \Delta\varphi_{K^*, J/\psi}, \Delta\varphi_{K^*, \phi})$$

Where  $\theta$  denotes the helicity angle, and  $\Delta\varphi$  is the angle between two decay chains.

- Resonance lineshape: Breit-Wigner (default), simplified K-matrix or Flatté functions for systematic uncertainties.

- The fit model of Run 1 analysis tested first:

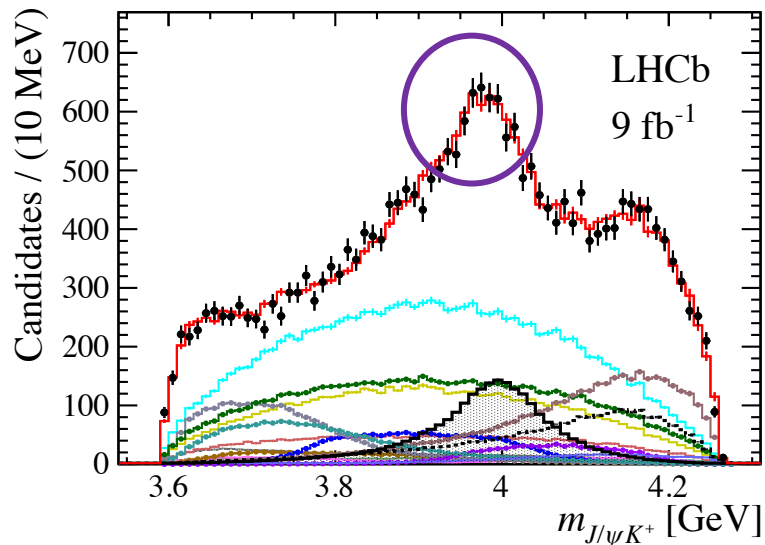
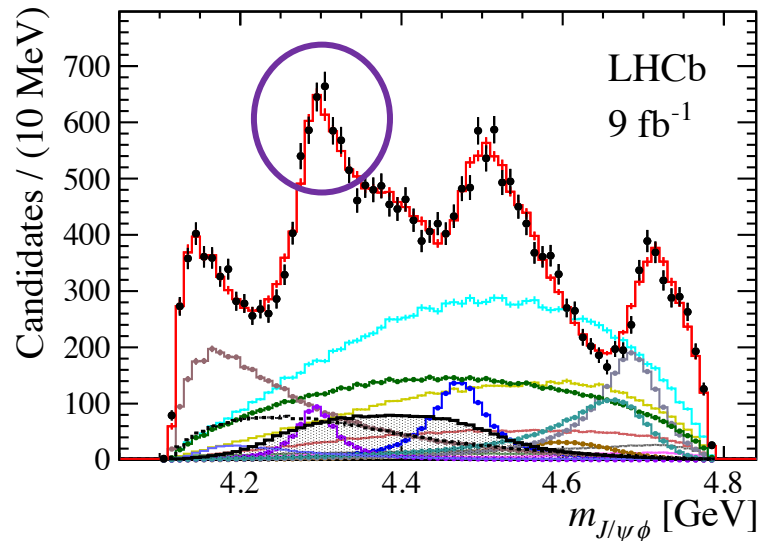
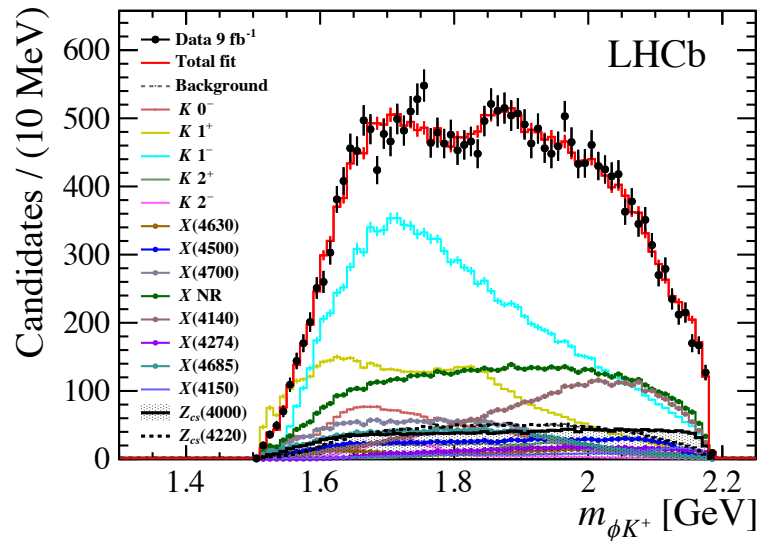
[PRL 118 (2017) 022003]



- Some deficiencies in describing the  $m_{J/\psi\phi}$  and  $m_{J/\psi K}$  are obvious.
- The fit model needs to be improved.

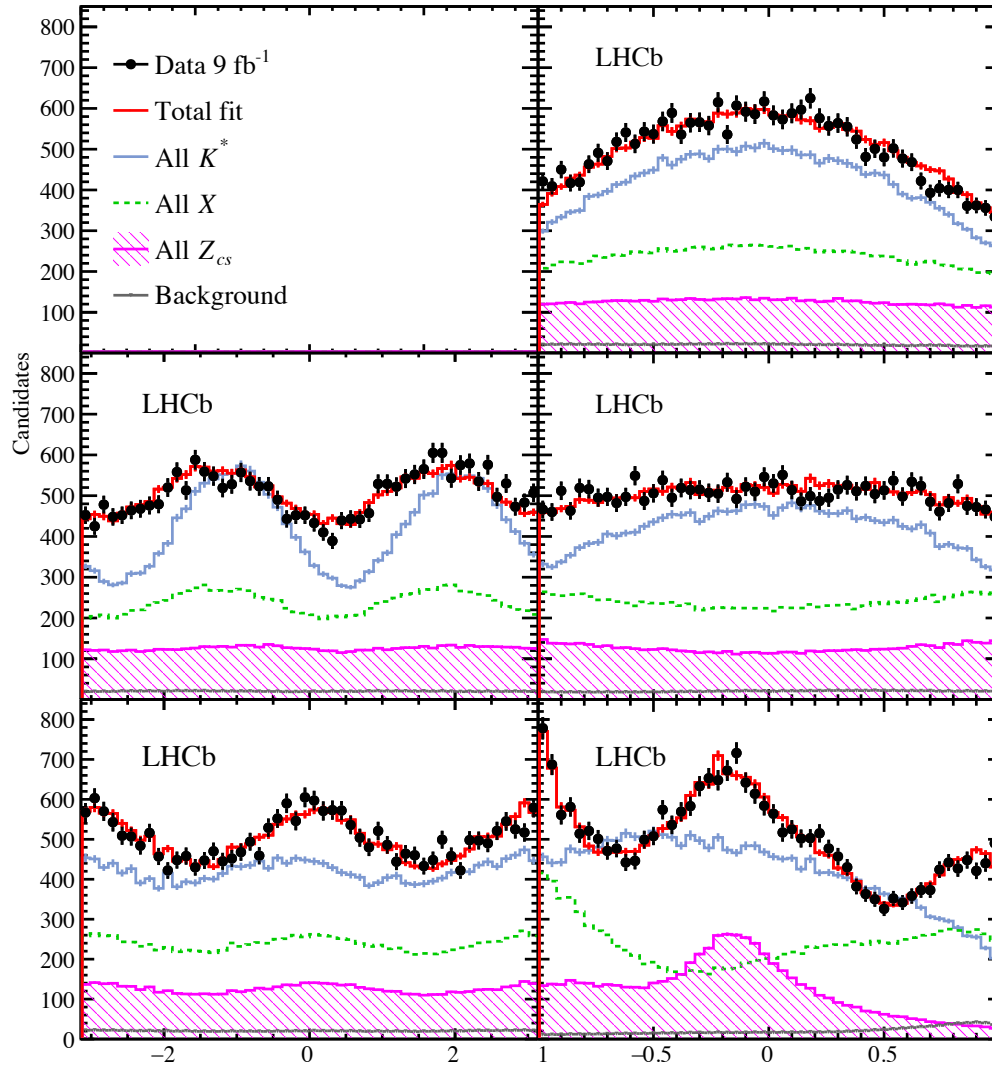
# Test new exotics

- Data cannot be described well by improving the  $K^*$  model
- New exotic states ( $X$  and  $Z_{cs}^+$ ) of different  $J^P$  were tested:
  - $1^+ Z_{cs}$  and  $1^+ X$ , giving the largest improvements, were first included.
  - Several states giving large fit improvements were also included in the default model: a second  $Z_{cs}$  and  $X$  states
- The default model includes  $9K^* + 7X + 1X(NR) + 2Z_{cs}$



- Data is well described.

- Angular distributions in the  $\phi K$  decay chain are also well described.





Contribution	Significance [ $\times\sigma$ ]	$M_0$ [MeV]	$\Gamma_0$ [MeV]	FF [%]	Fit fraction
$X(2^-)$					
$X(4150)$	4.8 (8.7)	$4146 \pm 18 \pm 33$	$135 \pm 28^{+59}_{-30}$	$2.0 \pm 0.5^{+0.8}_{-1.0}$	
$X(1^-)$					
$X(4630)$	5.5 (5.7)	$4626 \pm 16^{+18}_{-110}$	$174 \pm 27^{+134}_{-73}$	$2.6 \pm 0.5^{+2.9}_{-1.5}$	
All $X(0^+)$				$20 \pm 5^{+14}_{-7}$	
$X(4500)$	20 (20)	$4474 \pm 3 \pm 3$	$77 \pm 6^{+10}_{-8}$	$5.6 \pm 0.7^{+2.4}_{-0.6}$	
$X(4700)$	17 (18)	$4694 \pm 4^{+16}_{-3}$	$87 \pm 8^{+16}_{-6}$	$8.9 \pm 1.2^{+4.9}_{-1.4}$	
$NR_{J/\psi\phi}$	4.8 (5.7)			$28 \pm 8^{+19}_{-11}$	
All $X(1^+)$				$26 \pm 3^{+8}_{-10}$	
$X(4140)$	13 (16)	$4118 \pm 11^{+19}_{-36}$	$162 \pm 21^{+24}_{-49}$	$17 \pm 3^{+19}_{-6}$	
$X(4274)$	18 (18)	$4294 \pm 4^{+3}_{-6}$	$53 \pm 5 \pm 5$	$2.8 \pm 0.5^{+0.8}_{-0.4}$	
$X(4685)$	15 (15)	$4684 \pm 7^{+13}_{-16}$	$126 \pm 15^{+37}_{-41}$	$7.2 \pm 1.0^{+4.0}_{-2.0}$	
All $Z_{cs}(1^+)$				$25 \pm 5^{+11}_{-12}$	
$Z_{cs}(4000)$	15 (16)	$4003 \pm 6^{+4}_{-14}$	$131 \pm 15 \pm 26$	$9.4 \pm 2.1 \pm 3.4$	
$Z_{cs}(4220)$	5.9 (8.4)	$4216 \pm 24^{+43}_{-30}$	$233 \pm 52^{+97}_{-73}$	$10 \pm 4^{+10}_{-7}$	

- Two  $Z_{cs}^+ \rightarrow J/\psi K^+$  states were observed, both significance  $> 5\sigma$
- New  $X(4630)$  and  $X(4685)$  were observed, both significance  $> 5\sigma$
- Previous results using Run 1 sample were confirmed with large significance

# $J^P$ analysis

- “prefer”  $J^P$  gives best log-likelihood ( $\ln\mathcal{L}$ )
- The  $\ln\mathcal{L}$  difference between the “prefer” one and alternative hypotheses used to estimate significance:

$J^P$	$0^+$	$0^-$	$1^+$	$1^-$	$2^+$	$2^-$
<b>X(4630)</b>	$6.7\sigma$	$5.3\sigma$	$5.8\sigma$	prefer	$5.9\sigma$	<b><math>3.0\sigma</math></b>
X(4500)	prefer	$18\sigma$	$18\sigma$	$18\sigma$	$18\sigma$	$18\sigma$
X(4700)	prefer	$18\sigma$	$18\sigma$	$18\sigma$	$14\sigma$	$17\sigma$
X(4140)	$14\sigma$	$15\sigma$	prefer	$14\sigma$	$13\sigma$	$14\sigma$
X(4274)	$18\sigma$	$18\sigma$	prefer	$18\sigma$	$18\sigma$	$18\sigma$
<b>X(4685)</b>	$16\sigma$	$16\sigma$	prefer	$15\sigma$	$16\sigma$	$15\sigma$
<b>Z<sub>cs</sub>(4000)</b>	-	$17\sigma$	prefer	$17\sigma$	$15\sigma$	$16\sigma$
<b>Z<sub>cs</sub>(4220)</b>	-	$8.6\sigma$	prefer	<b><math>2.4\sigma</math></b>	$4.9\sigma$	$5.7\sigma$

- $J^P$  assignments to the previous 4X states are confirmed to be correct, with improved significance
- **Z<sub>cs</sub>(4000)** and X(4685)  $J^P$  are determined to be  **$1^+$**   $> 15\sigma$ .
- $J^P$  of the other two new X states not well determined (difference  $< 5\sigma$ ).
- Z<sub>cs</sub>(4220)  $1^+$  and  $1^-$  cannot be distinguished.

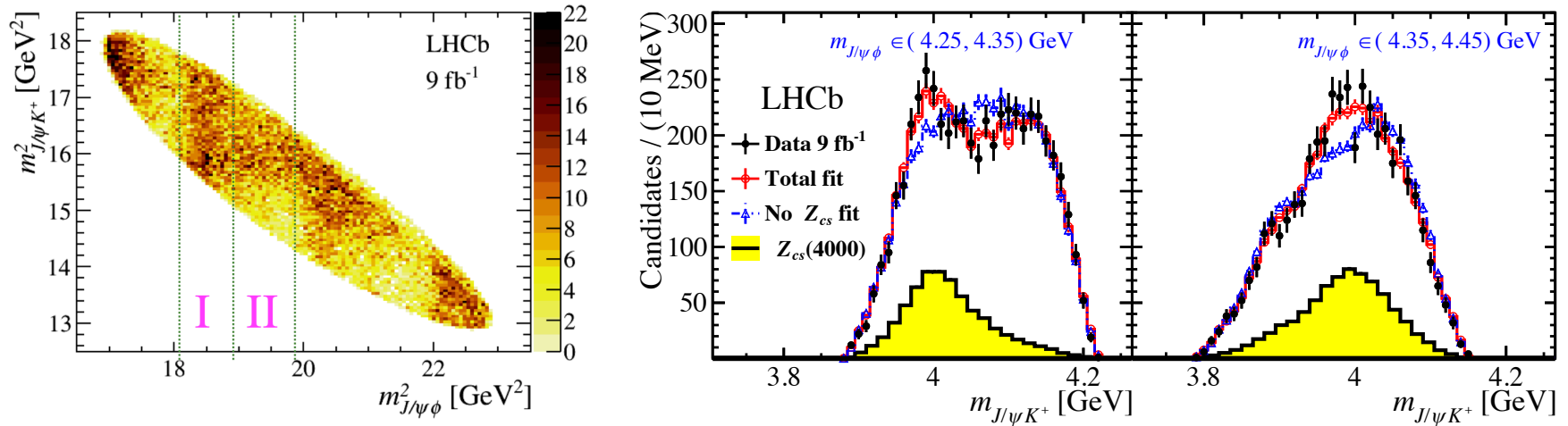
Following systematic sources are considered :

Source	$Z(4000)$			$X(4685)$		
	$M_0$	$\Gamma_0$	FF	$M_0$	$\Gamma_0$	FF
Fixed $M_0$ & $\Gamma_0$	-0.22	-3.60	-0.83	-0.14	2.72	0.25
$\chi^2_{\text{IP}}$ smearing	0.21	1.01	0.09	-0.53	1.11	0.12
Right sideband	0.01	0.58	0.11	-0.13	1.07	-0.13
Left sideband	-0.30	-1.16	-0.24	-0.09	-2.21	0.09
$\beta = 0.043$	-0.06	-0.00	0.01	0.01	-0.70	-0.09
$\beta = 0.037$	-0.02	0.26	0.02	-0.33	0.21	0.03
L0 Trigger	0.45	0.58	0.19	-0.58	1.12	0.11
PID efficiency	-1.06	-1.82	-0.69	-0.82	-4.42	-0.26
MC size	2.39	9.93	1.54	3.02	7.00	0.65
$\phi$ window	-4.71	-23.91	-2.75	8.60	-26.60	-1.17
Non $\phi$ subtraction	-2.87	-18.39	-1.79	12.40	-39.80	-1.80
Poly NR	-4.24	-16.36	-2.56	4.26	-22.07	-1.28
$X$ NR(1 <sup>+</sup> )	1.49	-21.25	-2.53	-15.72	35.54	3.84
$X$ NR(2 <sup>+</sup> )	2.16	3.09	1.26	1.88	-6.87	-0.03
BW $d=1.5$	-0.29	-5.27	-0.58	0.29	1.55	2.14
BW $d=4.5$	0.08	1.81	0.04	0.06	-3.53	-1.06
$L$	2.75	-3.19	-1.18	2.45	-24.33	-1.48
$X(4140)$ Flatté	0.52	-2.80	-0.45	-3.77	15.14	1.37
Extended model	-2.35	-6.66	-1.16	-3.61	-6.53	-0.94
Additional $X$	-0.68	2.07	0.30	0.74	-3.11	-0.18
1 <sup>-</sup> $Z$	-14.00	-21.09	-3.46	-9.41	-5.60	-1.52
$K^*$ BW	0.08	-0.66	-0.32	-0.06	-8.09	-0.82
K-Matrix	-3.75	-20.80	-2.85	4.10	-11.95	-0.06
$Z_{cs}(4000)$ Flatté	0.18		2.83	-0.85	2.79	0.18
Background model	0.10	-0.32	-0.12	-1.04	-1.72	-0.15
Total	(-14.26 , +3.85 )	(-26.26 , +26.26 )	(-3.43 , +3.41 )	(-16.05 , +12.82 )	(-40.85 , +36.72 )	(-1.96 , +3.92 )

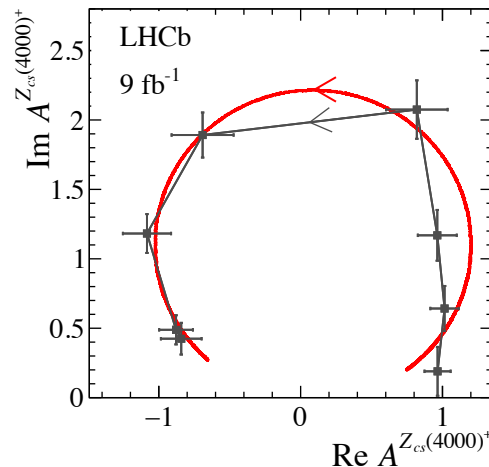
[\[arXiv:2103.01803\]](https://arxiv.org/abs/2103.01803)

- Take the maximum deviation among the **model variations**, then add other sources in quadrature.
- Different  $J^P$  assignments of  $Z_{cs}(4220)^+$  gives largest systematic uncertainty on  $Z_{cs}(4000)^+$ .

- The fit projection onto  $m(J/\psi K^+)$  in two slices of  $m(J/\psi \phi)$



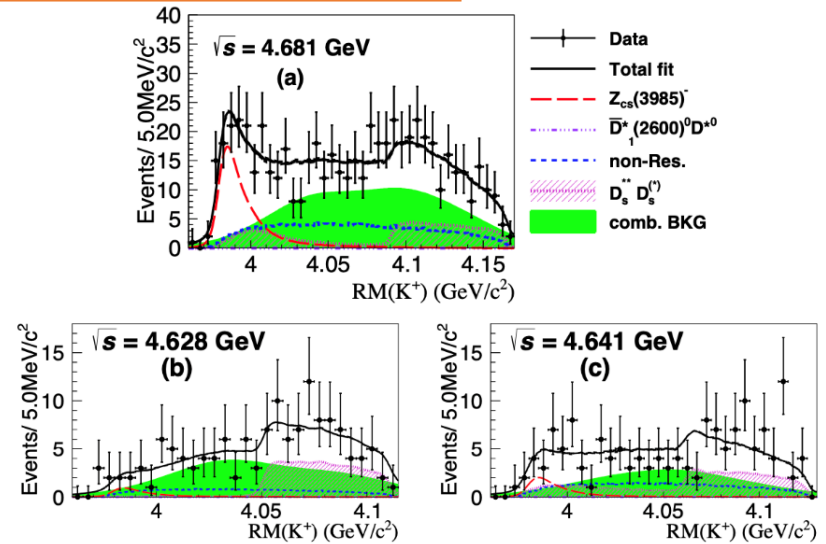
- Argand plot supports resonance character of  $Z_{cs}(4000)^+$ , obtained from lineshape-independent fitting.



# Comparison with BESIII

- BESIII experiment recently reported  $5.3\sigma$  observation of a very narrow  $Z_{cs}^-$  in  $D_s D^* + D D_s^*$  mass distributions
- Tests are applied:
  - Fix  $Z_{cs}(4000)^+$  to BESIII's result, log-likelihood is much worse.
  - Adding  $Z_{cs}(3985)^-$  on the default model almost doesn't improve the fit likelihood
- No evidence that  $Z_{cs}(4000)^+$  state is the same as the  $Z_{cs}(3985)^-$  seen by BESIII.

## BESIII results



$$m_{\text{pole}}(Z_{cs}(3985)^-) = (3982.5^{+1.8}_{-2.6} \pm 2.1) \text{ MeV}/c^2,$$

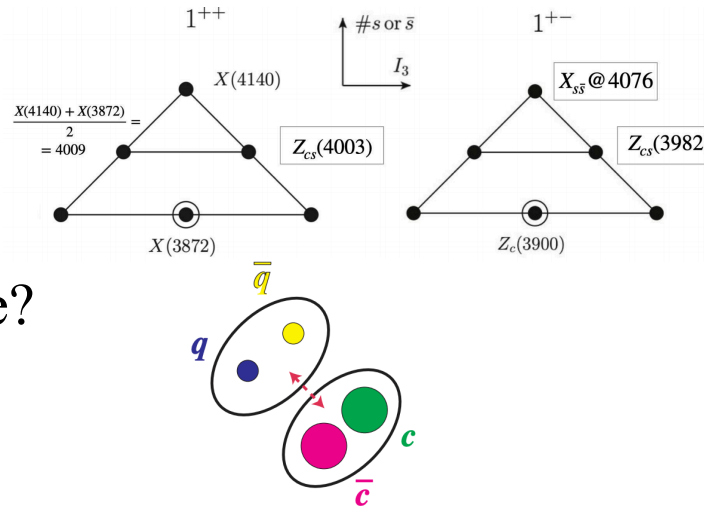
$$\Gamma_{\text{pole}}(Z_{cs}(3985)^-) = (12.8^{+5.3}_{-4.4} \pm 3.0) \text{ MeV}.$$

[\[arXiv:2011.07855\]](https://arxiv.org/abs/2011.07855)

# Theory interpretations to $Z_{cs}^+(4000)$

## ➤ Compact tetraquark model?

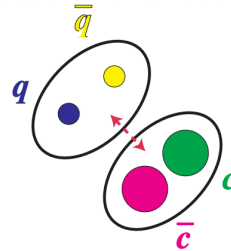
[\[arXiv:2103.08331\]](#)



## ➤ Hadronic molecule?

[\[arXiv:2011.08725\]](#)

[\[arXiv:2103.08586\]](#)

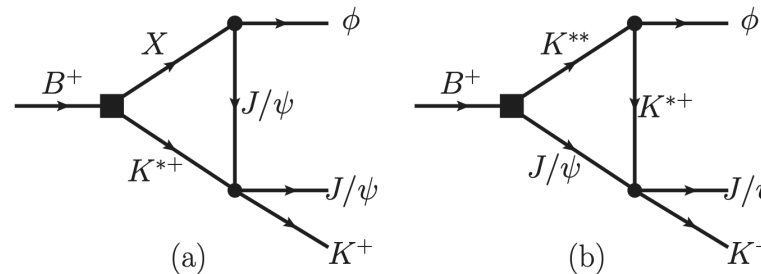


## ➤ $Z_{cs}(3985)$ and $Z_{cs}(4000)$ are same state in coupled-channels model?

[\[arXiv:2103.07871\]](#)

## ➤ Threshold cusps?

[\[arXiv:2103.05282\]](#)



➤ ... ..

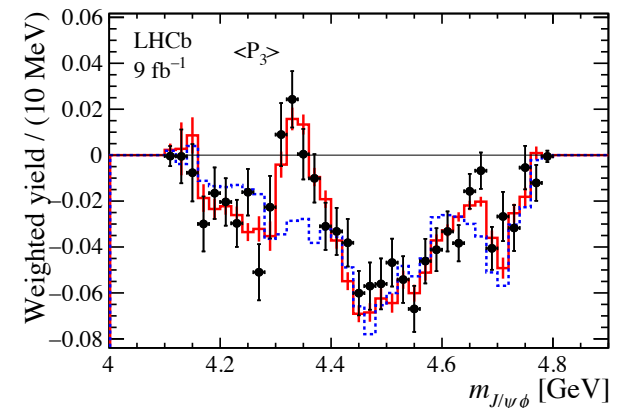
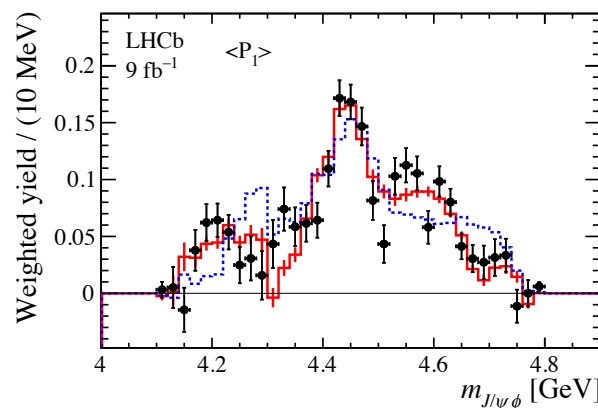
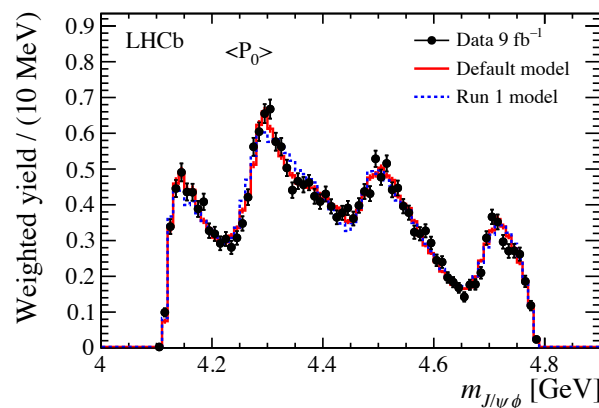
- The measured mass of  $X(4140)$  is  $4118 \pm 11_{-36}^{+19}$  MeV, with width  $162 \pm 21_{-49}^{+24}$  MeV, not very narrow; the mass is around the threshold of  $J/\psi\phi$ .

No evidence of a narrow threshold resonance at  $J/\psi\phi$  in our data

[Details in Sebastian's Talk](#)

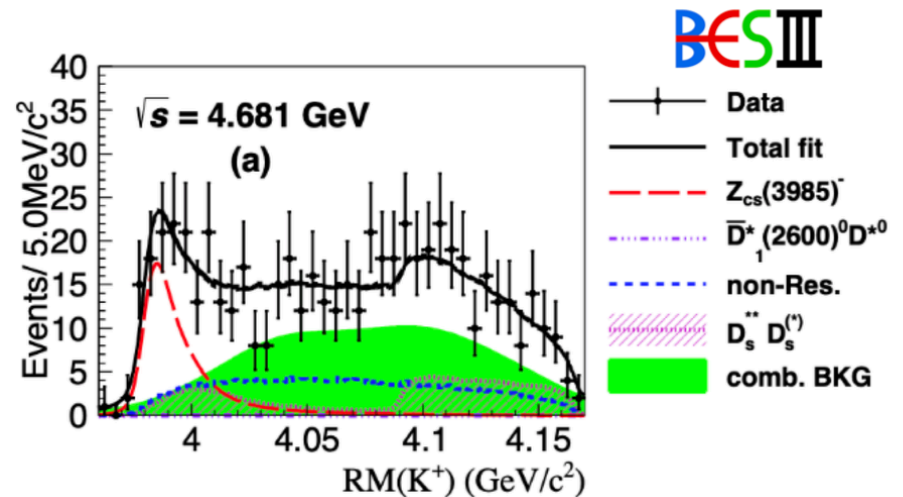
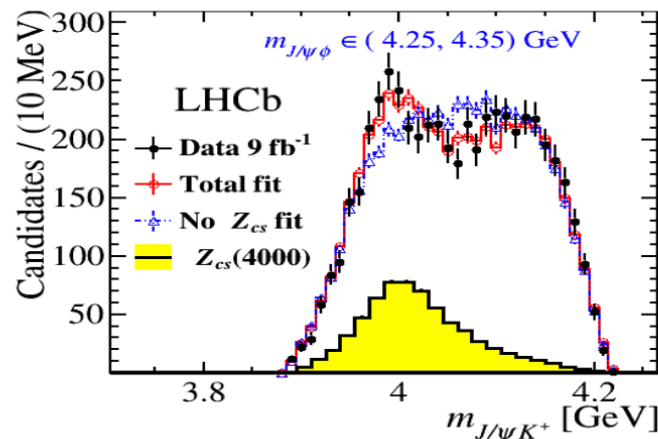
- Comparing the unnormalized Legendre moments of **Run 1 model** and **updated model**, new  $X(4630)$  and  $X(4685)$  are required.

$$\langle P_\ell^U \rangle = \sum_{i=1}^{N_{\text{events}}} \frac{1}{\epsilon_i} P_\ell(\cos \theta)_i$$



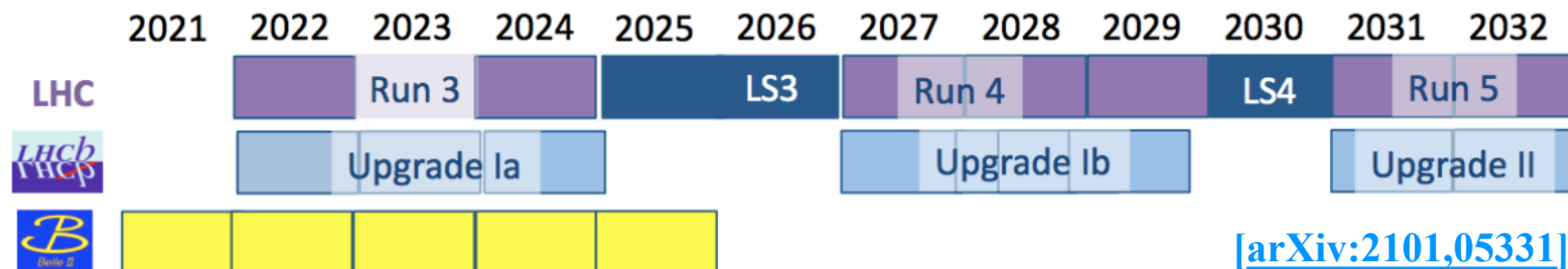
# Summary

- Four new  $J/\psi K^+$  and  $J/\psi \phi$  structures are observed  $B^+ \rightarrow J/\psi \phi K^+$ 
  1. 4  $X$  states observed in Run 1 data sample are confirmed, and  $J^P$  determined with higher significance.
  2.  $Z_{cs}(4000)^+ \rightarrow J/\psi K^+$  state is observed for first time, the significance is around  $15\sigma$ , and  $J^P = 1^+$  is also determined; another broad  $Z_{cs}(4220)^+$  is also observed
  3. A new  $1^+$   $X(4685)$  is  $> 15\sigma$ , and new  $X(4630) > 5\sigma$
- Understanding of  $Z_{cs}(4000)^+$  and  $Z_{cs}(3985)^-$  may shed light on molecular and compact tetraquarks





# Prospects



- LHCb is now boosting the data to a new level
    - Expect to **7x** more data (**14x** hadronic events) by 2029 than current, half of these by 2024
    - The  $J^P$  of  $Z_{cs}^+$  (4220) could be determined with larger data sample
    - The  $J^{PC}$  of  $X(4630)$  might be  **$1^{-+}$** , which is arousing interest
- [\[arXiv:2103.03127\]](https://arxiv.org/abs/2103.03127)
- If the  $Z_{cs}^+$  (4000) observed at LHCb and  $Z_{cs}^-$  (3985) observed at BESIII are same state?
  - The  $Z_{cs}^0$ , isospin partner of  $Z_{cs}^+$ , can be searched at LHCb

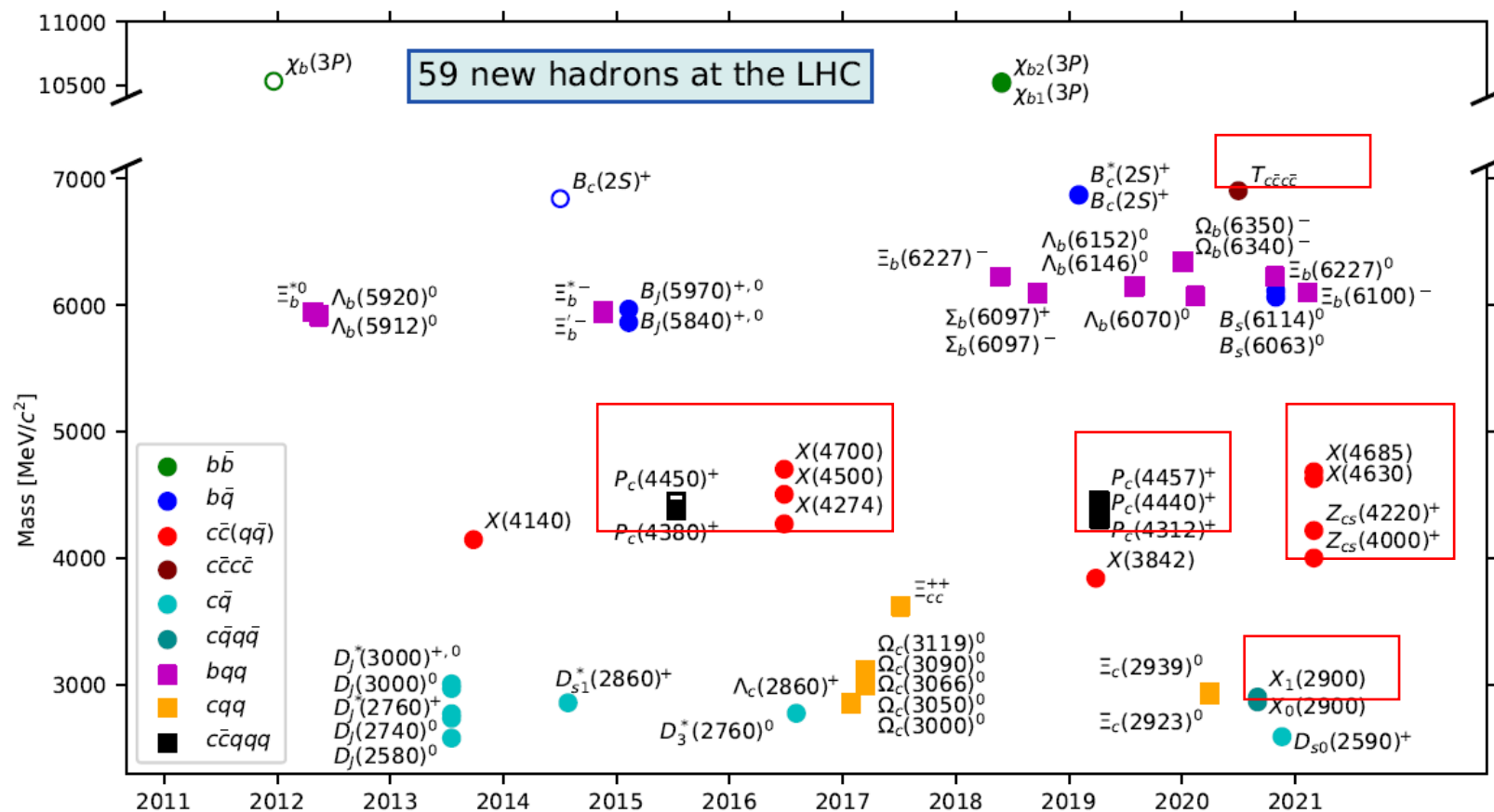
Thanks for listening

Backup

# Many exotic states observed at LHCb

- 59 new hadron states (conventional & exotic) observed at LHC, most of them discovered at LHCb

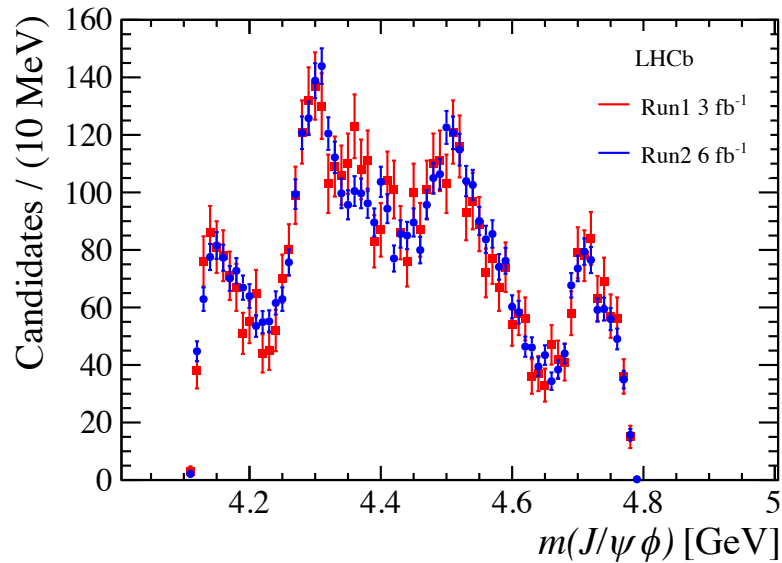
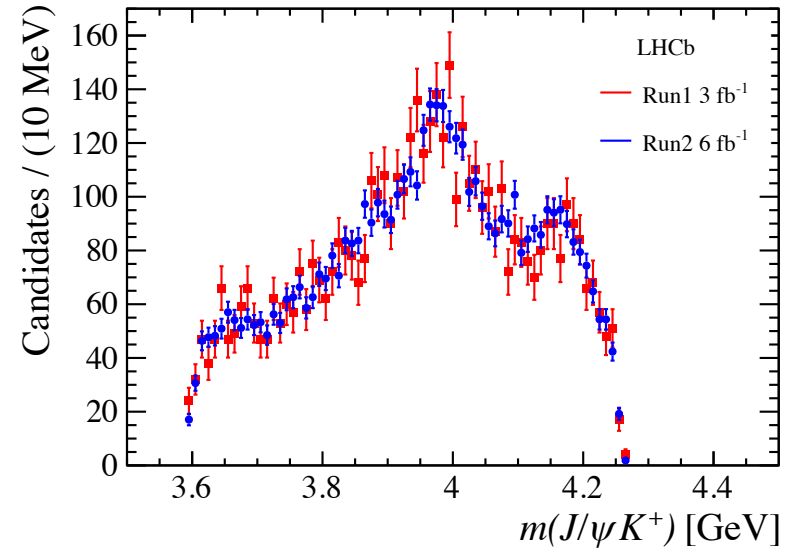
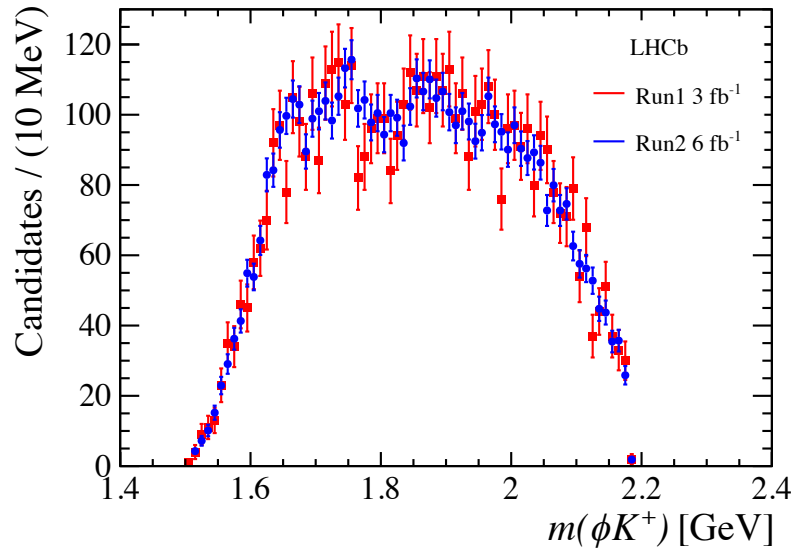
[Taken from CERNCOURIER]



**Diagram of discovery** The ATLAS, CMS and LHCb collaborations have discovered 59 new hadronic states so far – the most recent being the four tetraquarks reported in this article. Credit: CERN

# Run 1 and Run 2 comparison

[arXiv:2103.01803]



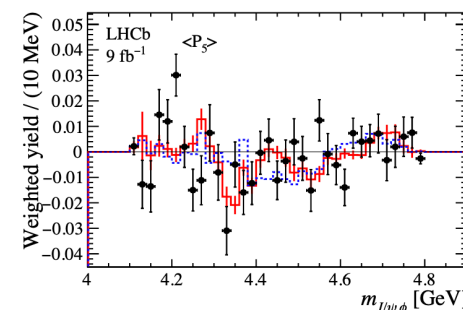
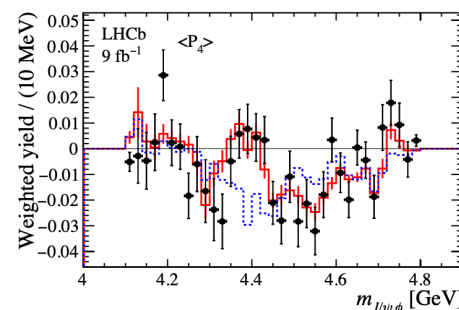
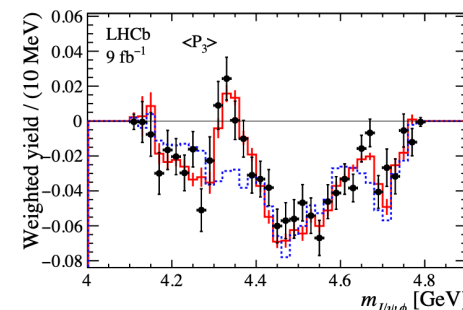
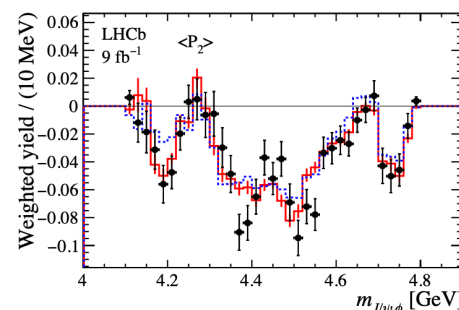
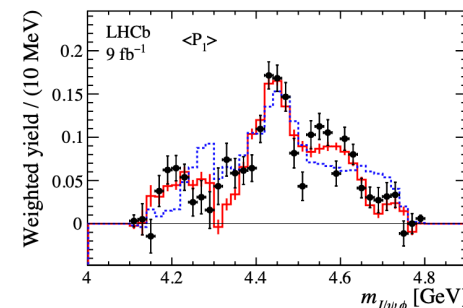
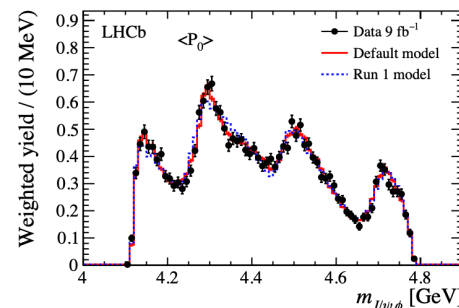
➤ Unnormalized Legendre moments:

$$\langle P_\ell^U \rangle = \sum_{i=1}^{N_{\text{events}}} \frac{1}{\epsilon_i} P_\ell(\cos \theta)$$

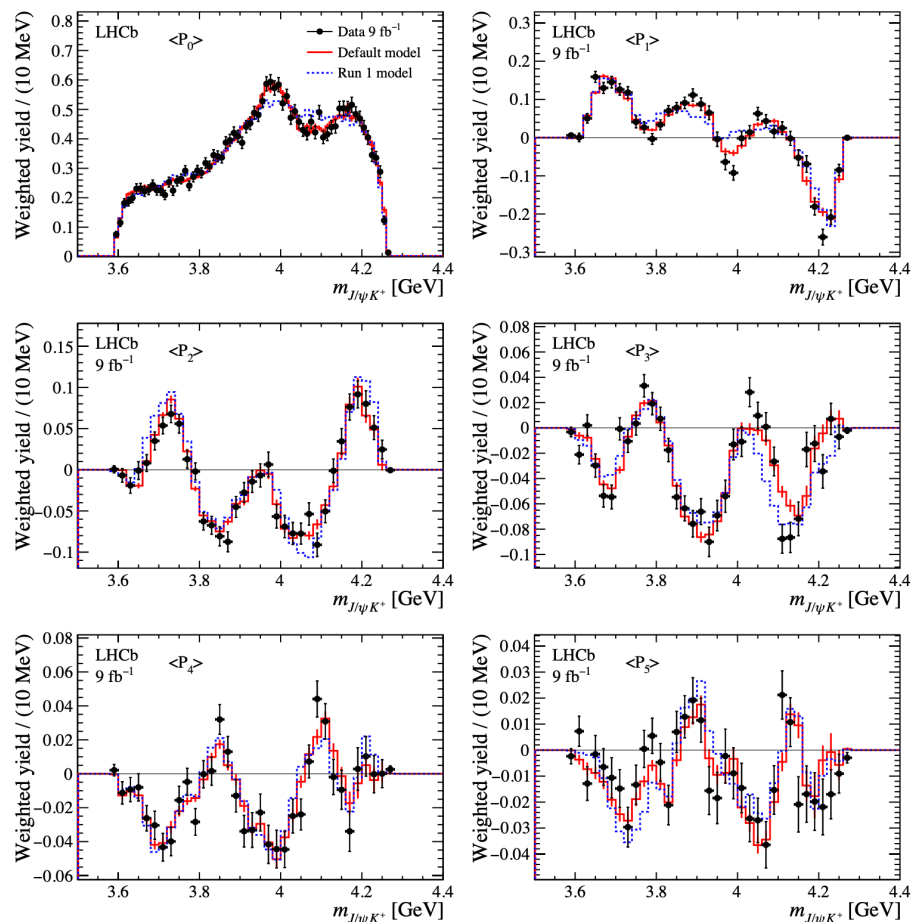
Legendre polynomial of order  $l$  and efficiency for each event  $i$ .

The moments distribution is obtained by a  $\frac{1}{\epsilon_i} P_l(\cos \theta)$  weight.

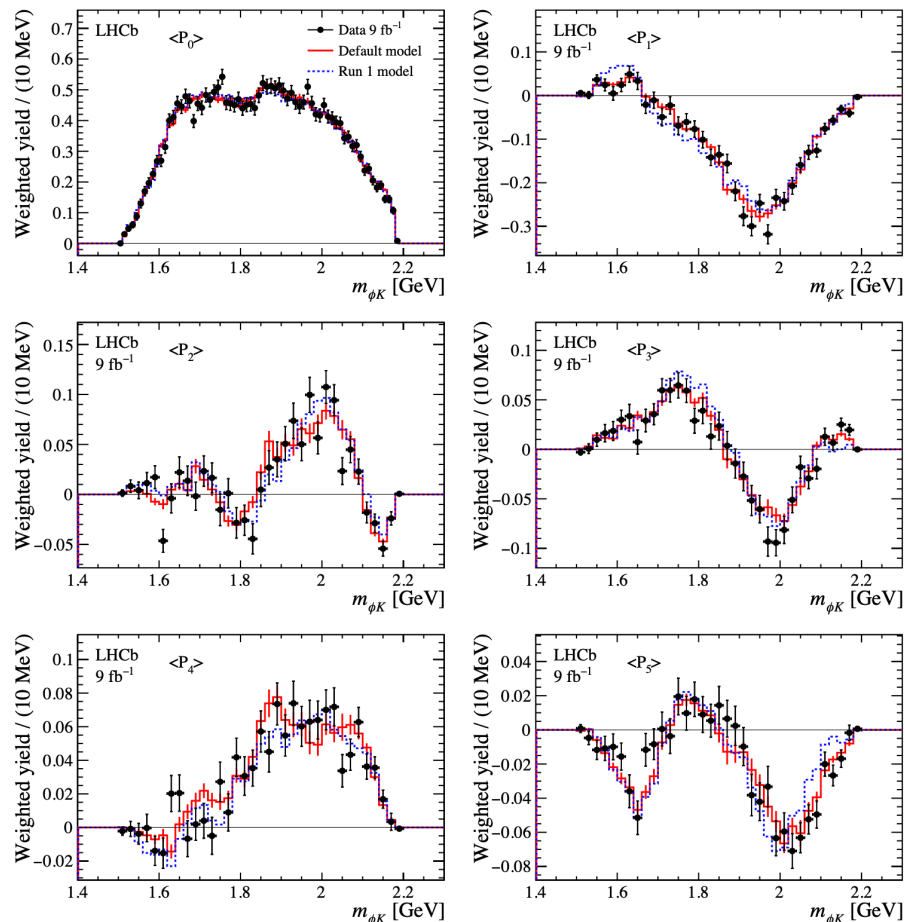
Angular moments of  
 $J/\psi\phi$  helicity angle



## Angular moments of $J/\psi K^+$ helicity angle



## Angular moments of $\phi K^+$ helicity angle



# Systematic uncertainty 1

- To evaluate uncertainties due to the **fixed masses and widths of known  $K^*$**  resonances: free the masses and widths but impose Gaussian constraints to the PDG values.
- $\chi^2_{IP}$  of  $B^+$  is not well modeled, smeared to match the data.
- To explore uncertainty in the background model, vary the  $B^+$  sideband window.
- The uncertainty in the background fraction  $\beta$ : change background shape to exponential function.
- Vary the Blatt-Weisskopf barrier factor  $d$  (hadron-size parameter).
- **Vary the smallest allowed orbital momentum** in the resonance description function, associate the L dependent term with each LS coupling.

$$R_{K_n^*}(m_{K\phi}) = \underbrace{B'_{L_B^{K_n^*}}(p, p_0, d) \left( \frac{p}{M_B} \right)^{L_B^{K_n^*}}}_{\text{Angular momentum barrier factor}} \underbrace{\text{BW}(m_{K\phi} | M_0^{K_n^*}, \Gamma_0^{K_n^*})}_{\text{Relative Breit-Wigner function}} \underbrace{B'_{L_{K_n^*}}(q, q_0, d) \left( \frac{q}{M_0^{K_n^*}} \right)^{L_{K_n^*}}}_{\text{Orbital momentum}}$$



# Systematic uncertainty 2

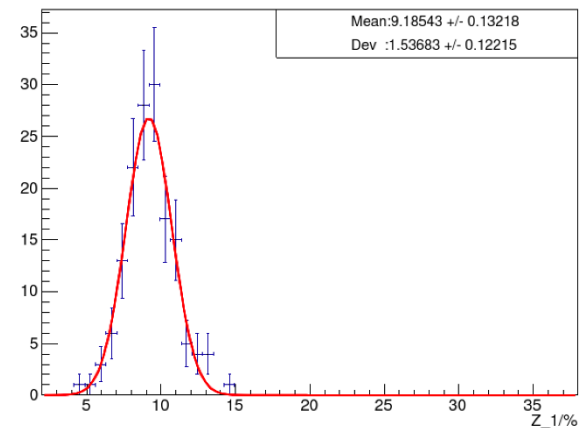
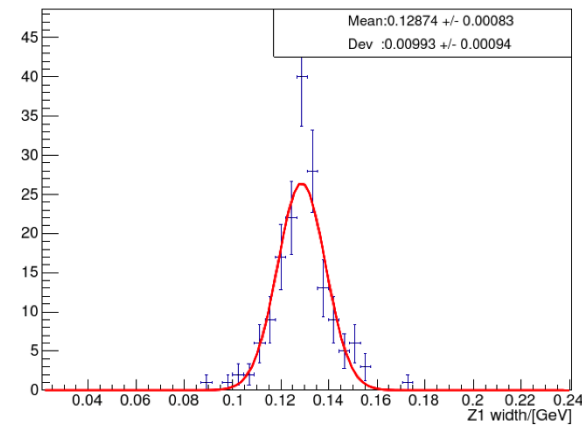
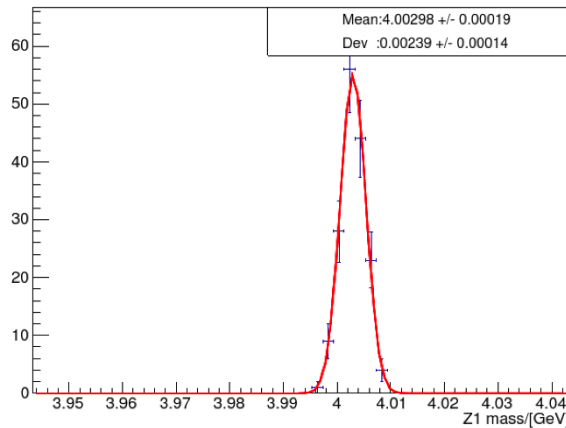
- Uncertainty due to the choice of NR component, change the constant parameterization to exponential function.
- $1^+$  or  $2^+$  NR  $X$  contributions are optionally introduced.
- The difference between nominal model and extended model.
- Flatté function to parameterize  $X(4140)$  or  $Z_{cs}(4000)$  to replace BW function.

$$\text{Flatte}_X(m|M_0, g_{J/\psi\phi}, g_{D_s^*D_s}) = \frac{1}{M_0^2 - m^2 - iM_0(g_{J/\psi\phi}\rho_{J/\psi\phi} + g_{D_s^*D_s}\rho_{D_s^*D_s})},$$

- Additional  $X$  states with different  $J^P$  in the extended model.
- Neglected no- $\phi$  contribution: 1) Change the  $\phi$  mass window from  $\pm 15\text{MeV}$  to  $\pm 7\text{MeV}$ , 2) sFit to subtract no- $\phi$  contribution is performed as alternative to cFit
- Modification of  $K^*$  width: as the partial width to  $\phi K$  is unknown, try a fit with mass dependence of the width driven by the lowest allowed decay channel, which is  $K\pi$  for the natural spin-parity and  $K\omega$  for others.

# Systematic uncertainty 3

- L0 trigger : change the requirement from L0 global decision to TOS on L0 Muon or Dimuon decision.
- PID correction : the nominal PID calibration is performed from PIDcorr package, recalibrate from PIDGen package. (small)
- MC size : simulation sample size is studied from bootstrap method, 200 bootstrapping samples are generated, and the deviations are taken as systematic uncertainties ( see plots below )



# Systematic uncertainty 4

- As an alternative to the 2D factorization of 6D background PDF, decompose the background density into multidimensional moments in the  $K^*$  decay chain variables (this uncertainty is small)
- K-Matrix model :
  1. Some  $K^*$  with the same  $J^P$  are overlapping, we use a simple K-Matrix formula to describe them as alternative

$$RK M_n(m|M_{0n}, \Gamma_{0n}) = \frac{\frac{1}{M_{0n}^2 - m^2}}{1 - i(\sum_j \frac{M_{0j}\Gamma_{0j}(m)}{M_{0j}^2 - m^2} + f_{sc} \cdot \rho(m))},$$

denominator sums over the same  $J^P$   $K^*$  resonances,  $f_{sc}$  accounts for possible non-resonance contribution. This fit didn't change the conclusion.

2. Alternative K-Matrix model with two coupling channels are tested, used to describe the  $2^1 P_1$  and  $2^3 P_1$   $K^*$  resonances

$$\mathcal{K}_{ba}(s) = \sum_R \frac{g_b^R g_a^R}{M_R^2 - s} + \sum_{i=0}^{N_{b.g.}} b_{ba}^{(i)} s^i$$

more floating parameters are included, the nominal model is stable.