X(3872) in the hybrid model of charmonium and hadronic molecule

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in collaboration with

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Experimental and theoretical status of and perspectives for XYZ states, Online, 12-15 April 2021

Hadron structure: Constituent quark model

- Hadron = Quark composite system
- Ordinary Hadrons: Baryon (qqq) and Meson $(q\bar{q})$



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Exotic Hadrons ($\neq qqq$, $q\bar{q}$): Multiquark? Multihadron?



Constituent quark picture and beyond



N. Brambilla, et al. Eur. Phys. J.C 71 (2011)1534, S. Godfrey and N. Isgur, PRD32 (1985)189

Y. Yamaguchi (JAEA)

Constituent quark picture and beyond

• e.g. $c\bar{c}$ mesons (Charmonium) and Unexpected X, Y, Z



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Exotics ≠ cc̄ have been observed in the Experiments (BaBar, Belle, BESIII, LHCb,...) ⇒ Q. Structure? Physics?

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Recent reports of Exotic hadrons!

▷ X(6900) (cccc?)

LHCb, Science Bulletin 65 (2020) 1983



 $\triangleright Z_{cs} (c \overline{c} s \overline{u}?)$

BESIII PRL126, 102001 (2021)



▷ X_{0,1}(2900) (*c̄sud*?)

LHCb, PRL125, 242001 (2020), PRD102, 112003 (2020)



 $\triangleright P_c (uudc\bar{c}?),$

LHCb PRL115(2015)072001, PRL122(2019)222001



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LHCb PRL115(2015)072001, PRL122(2019)222001, 2012.10380



Observation of X(3872)!

► X(3872) (or $\chi_{c1}(3872)$) was reported by Belle in $B \to K\pi^+\pi^- J/\psi$ decay S.-K. Choi *et al.*, PRL91(2003)262001





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Mass, width and $I^G(J^{PC})$ from PDG PA. Zyla et al. (Particle Data Group), PTEP 2020, 083C01 (2020) Mass: 3871.69 ± 0.17 MeV $\mathbf{I}^G(\mathbf{J}^{\mathbf{PC}}) = 0^+(1^{++})$ Width: < 1.2 MeV

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Found in many types of experiments (summarized in Y.Y *et al.*, J.Phys.G 47 (2020)053001) **B-decay:** Belle, *BABAR*, LHCb,..., e^+e^- BESIII, ... *pp*: LHCb, CMS, ATLAS, ..., $p\bar{p}$ CDF (II) D0,... **PbPb:** CMS 2102.13048 [hep-ex]

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Exotic nature of $X(3872) \neq c\bar{c}$

> Isospin violation in the X(3872) decay [Belle, BABAR,...]

 $\frac{Br(X \to \pi^+ \pi^- \pi^0 J/\psi)_{I=0}}{Br(X \to \pi^+ \pi^- J/\psi)_{I=1}} = \mathbf{1.0} \pm \mathbf{0.4} \pm \mathbf{0.3}$

 \Rightarrow Isospin mixed state $!? \rightarrow \text{Exotic structure (e.g. } c\bar{c}q\bar{q})$ is expected

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- X(3872) near the $D\bar{D}^*$ thresholds

 $D^+D^{*-}_{--}$ 3879.84 MeV $D^0\bar{D}^{*0}_{--}$ 3871.69 MeV 3871.68 MeV

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• X(3872) near the $D\bar{D}^*$ thresholds



 $X(3872) = D\bar{D}^*$ molecule?

⇒ Hadronic molecule is expected near thresholds ⇒ X(3872) as $D\overline{D}^*$ molecule?

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Possible existence of a $c\bar{c}$ core component in X(3872)

Production in the pp collisions



• $(d\sigma/dp)_{X(3872)}$ is similar to that of ψ' rather than that of nuclei (a hadron composite)

 \Rightarrow Existence of a $c\bar{c}$ core component (~ 5%)?

Y. Yamaguchi (JAEA)

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 \Leftrightarrow



$$\left(B
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$$M_X = 3871.69 \pm 0.17$$
 MeV
 $\Gamma_X < 1.2$ MeV

[Belle]

$$M_X = 3872.9^{+0.6}_{-0.4} - 0.5$$
 MeV
 $\Gamma_X = 3.9^{+2.8}_{-1.4} + 0.2$ MeV
[BABAR]
 $M_X = 3875.1^{+0.7}_{-0.5} \pm 0.5$ MeV
 $\Gamma_X = 3.0^{+1.9}_{-1.4} \pm 0.9$ MeV

Inconsistency between $X \to J/\psi \pi \pi$ and $X \to D\bar{D}^*$ in the *B*-decay process?



$$\left(B
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$$\begin{array}{l} \text{[Belle]} \\ M_X = 3871.69 \pm 0.17 \text{ MeV} \\ \Gamma_X < 1.2 \text{ MeV} \end{array} \Leftrightarrow \begin{array}{l} \text{[Belle]} \\ M_X = 3872.9^{+0.6}_{-0.4} + 0.4 \text{ MeV} \\ \Gamma_X = 3.9^{+2.8}_{-1.4} + 0.2 \text{ MeV} \\ \text{[$BABAR]} \\ M_X = 3875.1^{+0.7}_{-0.5} \pm 0.5 \text{ MeV} \\ \Gamma_X = 3.0^{+1.9}_{-1.4} \pm 0.9 \text{ MeV} \\ \end{array} \\ \Rightarrow \quad (M_X, \Gamma_X)_{X \to J/\psi\pi\pi} < (M_X, \Gamma_X)_{X \to D\bar{D}^*} \end{array}$$

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Our model: a hybrid state of $c\bar{c}$ and hadronic molecules



• Core $c\bar{c}$ as $\chi_{c1}(2P)$ with $m_{c\bar{c}} = 3950$ MeV (quark model pred.)

 One pion exchange potential (OPEP) between D^(*) D
^(*) (Central + Tensor)

► Couplings of $c\bar{c} - D\bar{D}^*$, $D\bar{D}^* - J/\psi V \Rightarrow$ determined empirically

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Potentials

$$\begin{split} D^{(*)}\bar{D}^{(*)} - D^{(*)}\bar{D}^{(*)} : & V_{OPEP} \\ c\bar{c} - D\bar{D}^* : & u(\vec{q}\,) = \frac{g_{c\bar{c}}}{\sqrt{\Lambda_q}}L_{\Lambda_q}(\vec{q}\,) & \text{with Form factor } L_{\Lambda_q}(\vec{q}\,) = \frac{\Lambda_q^2}{\Lambda_q^2 + \vec{q}^2} \\ D\bar{D}^* - J/\psi\,V : & v(\vec{q},\vec{q}\,') = \frac{v_0}{\Lambda_q^2}L_{\Lambda_q}(\vec{q}\,)L_{\Lambda_q}(\vec{q}\,') \end{split}$$

 $^*g_{car{c}}$, v_0 determined empirically, $car{c}-J/\psi\,V$ is not considered because of the OZI rule

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Hamiltonian of
$$c\bar{c}$$
 + two mesons (*mm*): $H = \begin{pmatrix} m_{c\bar{c}} & U^{(mc)} \\ U^{(mc)} & H^{(mm)} \end{pmatrix}$

• Hamiltonian of two mesons:

$$H^{(mm)} = H_0 + V^{D\bar{D}^* - J/\psi V} + V_{\text{OPEP}}$$

• Matrices of the $c \bar{c} - (mm)$ potential and $D \bar{D} - J/\psi V$ potentials

$$U^{(mc)} = \begin{pmatrix} -u \\ u \\ 0 \\ 0 \end{pmatrix} \int_{J/\psi\omega}^{p\bar{}_{D^{*0}}} V^{D\bar{}_{D^{*}}-J/\psi V} = \begin{pmatrix} 0 & 0 & -v & -v \\ 0 & 0 & v & -v \\ -v & v & 0 & 0 \\ -v & -v & 0 & 0 \end{pmatrix}$$

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One pion exchange potential for $D^{(*)}\bar{D}^{(*)}$



- Effective $\pi D^{(*)} D^{(*)}$ Lagrangian $\mathcal{L}_{\pi H H} = i \frac{g_{\pi}}{2f_{\pi}} \operatorname{tr} \left[H_a \gamma_{\mu} \gamma_5 \vec{\tau} \cdot i \partial^{\mu} \vec{\pi} \bar{H}_a \right]$ $H_a = \frac{1 + \dot{p}}{2} \left[-P^*_{a\mu} \gamma^{\mu} + P^a \gamma_5 \right], \bar{H}_a = \gamma_0 H^{\dagger}_a \gamma_0$
- Nonzero energy transfer $\omega^2 \sim (m_{D^*} m_D)^2$ \rightarrow Effective π mass: $\mu^2 = m_{\pi}^2 - \omega^2 < 0$

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Effective
$$\pi D^{(*)} D^{(*)}$$
 Lagrangian
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• OPEP with the imaginary mass $m = i\mu$:

$$V_{OPEP}(r) = \left(\frac{g_{\pi}}{2f_{\pi}}\right)^2 \frac{1}{3} \begin{bmatrix} \vec{S}_1 \cdot \vec{S}_2 C(r;\mu) + S_{12}(\hat{r}) T(r;\mu) \end{bmatrix} \vec{\tau}_1 \cdot \vec{\tau}_2$$

Central Tensor

* the real part is employed, while the imaginary one is ignored ** the contact term is not considered

Form factor with Cutoff Λ (determined by the hadron size)

$$F(\vec{q}^{\,2}) = \frac{\Lambda_{\pi}^2 - m_{\pi}^2}{\Lambda_{\pi}^2 + \vec{q}^{\,2}}, \quad \Lambda_{\pi} \sim 1130 \text{ MeV}$$

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Parameter sets

Model	$g_{c\bar{c}}$	$\Lambda_q \; [\text{GeV}]$	g_{π}	Λ_{π} [GeV]	$rms(D^0 \overline{D}^{*0})$ [fm]	$rms(D^{\pm}\bar{D}^{*\mp})$ [fm]	BE [MeV]
only $c\bar{c} - D\bar{D}^*$	0.05110	0.5	_	—	8.39	1.56	0.16 (input)
only OPEP	—		0.55	1.791	8.25	1.44	0.16 (input)
Full	0.04445	0.5	0.55	1.13	8.36	1.59	0.16 (input)

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Mixing ratios: Comparison of models

Model	$c\bar{c}$	$D^0 \bar{D}^{*0}({}^3S)$	$D^0\bar{D}^{*0}(^3D)$	$D^{*0}\bar{D}^{*0}({}^5D)$	$D^{\pm}\bar{D}^{*\mp}(^{3}S)$	$D^{\pm}\bar{D}^{*\mp}(^{3}D)$	$D^{*\pm}\bar{D}^{*\mp}({}^5D)$
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• Major component: $D^0 ar{D}^{*0} ({}^3S) \sim 86.9\%$

- $c\bar{c}$ component $\sim 5.9\% \rightarrow$ consistent with result of the production ratio in the pp collision
- Only OPEP \Rightarrow Large cutoff $\Lambda_{\pi} \sim 1.8$ GeV is needed.

 $\Rightarrow c\bar{c} - D\bar{D}^* \text{ coupling helps to generate a bound state}$ Y. Yamaguchi (JAEA) XYZ state 2021 (12 Apr.)

Decay Spectrum



Decay spectrum

Studying peak structures of X(3872) in the B-decay process



M.Takizawa, S.Takeuchi PTEP2013,093D01, S.Takeuchi, K. Shimizu, M.Takizawa, PTEP2014,123D01

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M.Takizawa, S.Takeuchi PTEP2013,093D01, S.Takeuchi, K. Shimizu, M.Takizawa, PTEP2014,123D01 \Rightarrow computing the transfer from $c\bar{c}$ to two-meson final states (D, \bar{D}^*) , $(J/\psi, \rho)$,... $c\bar{c} - D\bar{D}^*$ coupling is allowed, while no direct $c\bar{c} - J/\psi V$ by the OZI rule $\leftrightarrow D\bar{D}^* - J/\psi V$ coupling is allowed

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- ► $c\bar{c} D\bar{D}^*$ coupling is allowed, while no direct $c\bar{c} J/\psi V$ by the OZI rule $\leftrightarrow D\bar{D}^* - J/\psi V$ coupling is allowed
- Complex scaling method (CSM) is employed T. Myo, et al., PPNP 79 (2014) 1-56 χ^{θ}_{α} : eigenfunction with the eigenvalue E_{α} obtained by CSM

Y. Yamaguchi (JAEA)

• Contribution from each final states, $D\bar{D}^*$, $J/\psi\rho$, $J/\psi\omega$

dW/dE S. Takeuchi, K. Shimizu, M. Takizawa PTEP2014, 123D01



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▶ $D^0 \overline{D}^{*0}$ peak is wide,

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D⁰D̄^{*0} peak is wide, while J/ψρ, J/ψω peaks are very narrow!
 Contribution from J/ψρ is smaller than that from J/ψω

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Narrow width of $J/\psi\rho$, $J/\psi\omega$ peaks

Suppression of the $c\bar{c} \rightarrow J/\psi V$ transition ($V = \rho, \omega$)

- 1. No direct $c\bar{c} \rightarrow J/\psi V \operatorname{coupling} \leftrightarrow \operatorname{OZI}$ rule
- 2. Rearrangement is necessary



 \Rightarrow suppression of the $(J/\psi\rho)_{I=1}$ decay

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X(3872) pole above the threshold

Previous works: PTEP**2013**,093D01 PTEP**2014**,123D01

BE= 0.16 MeV below $D^0 \overline{D}^{*0}$ threshold was assumed [PDG2012]

[PDG2020] $M_{X(3872)}$ above the $D^0 \bar{D}^{*0}$ threshold $D^+ D^{*-}_{--} 3879.84 \text{ MeV}$ $D^0 \bar{D}^{*0}_{--} 3871.69 \text{ MeV}$

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 $D^0 \overline{D} \stackrel{*^0}{=} - \frac{3871.69 \text{ MeV}}{3871.68 \text{ MeV}}$ \rightarrow Q. Can we obtain a pole above $D^0 \overline{D}^{*0}$ in the hybrid model?

Y. Yamaguchi (JAEA)

X(3872) pole in the LHCb analysis

▶ Poles by the Flatté approach (Only the $D^0 \overline{D}^{*0}$ channel is considered)

Im E [MeV] sheet IV sheet IV [m E [MeV] LHCb LHCb -0.2 1σ -0.4 2σ -0.6 = (0.06 - 0.13 i) MeV 2: -0.8 sheet II sheet III -1 -0.4 -0.20.2 -2 0 Re E [MeV] Re E [MeV]

R. Aaij et al., (LHCb collaboration) PRD102(2020)092005

Two poles are found

1. $E_{II} = 0.06 - 0.13i$ MeV on the physical sheet \rightarrow above $D^0 \bar{D}^{*0}$ threshold

2. $E_{III} = -3.58 - 1.22i$ MeV on the unphysical sheet

Y. Yamaguchi (JAEA)

Summary



- > X(3872) as a hybrid state of $c\bar{c}$ and hadronic molecules
- The X(3872) spectrum is reproduced reasonably.
 ⇒ About 6% cc̄ component ⇔ agreement with the production measurement
- ▶ The difference between $X \to J/\psi \pi \pi$ and $X \to D\bar{D}^*$ in the *B* decay spectrum.

 $\Rightarrow D\bar{D}^*$ peak is wide, while the peaks of $J/\psi\rho$, $J/\psi\omega$ peak are narrow.

 \blacktriangleright The pole can be above the $\bar{D}^0\bar{D}^{*0}$ threshold thanks to the finite width of ρ and ω

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Thank you for your kind attention.

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