

Multiscale approach to the physics of ion beam cancer therapy

Andrey Solov'yov



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Acknowledgements

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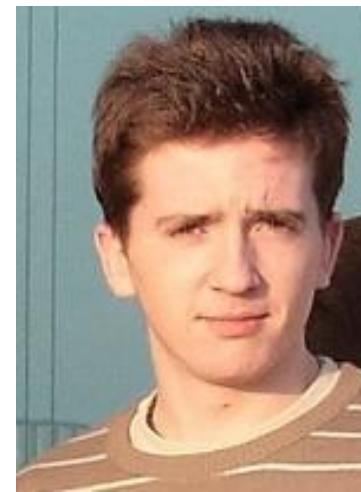
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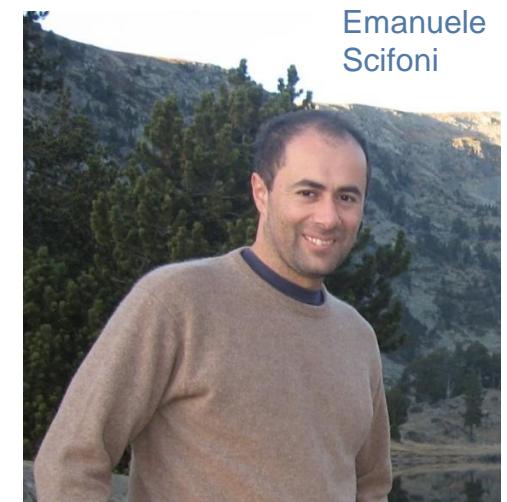
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Outline

- Introduction to the problem of radiation biodamage by ions and ion-beam cancer therapy
- Multiscale approach: current status and perspectives
 - Propagation of ions
 - Primary ionization in the medium
 - Propagation of secondary electrons, holes and free radicals
 - Mechanisms of DNA degradation
 - Radiobiological scale effects
- Conclusions and outlook

Ion beam cancer therapy principles in brief

- Fundamental properties of propagation of charged particles: - dE/dx (LET) $\sim Z_{\text{eff}}^2/E$
- Bragg peak in Linear Energy Transfer (LET) closer to the end of the ion's track
- Beams easily manipulated, better focusing



1992-2007

Technical development

Kraft et al. Rep.Mod.Phys. (2007)

^{12}C beam



Nov.2009 -
now

Clinical treatment

www.hit-centrum.de

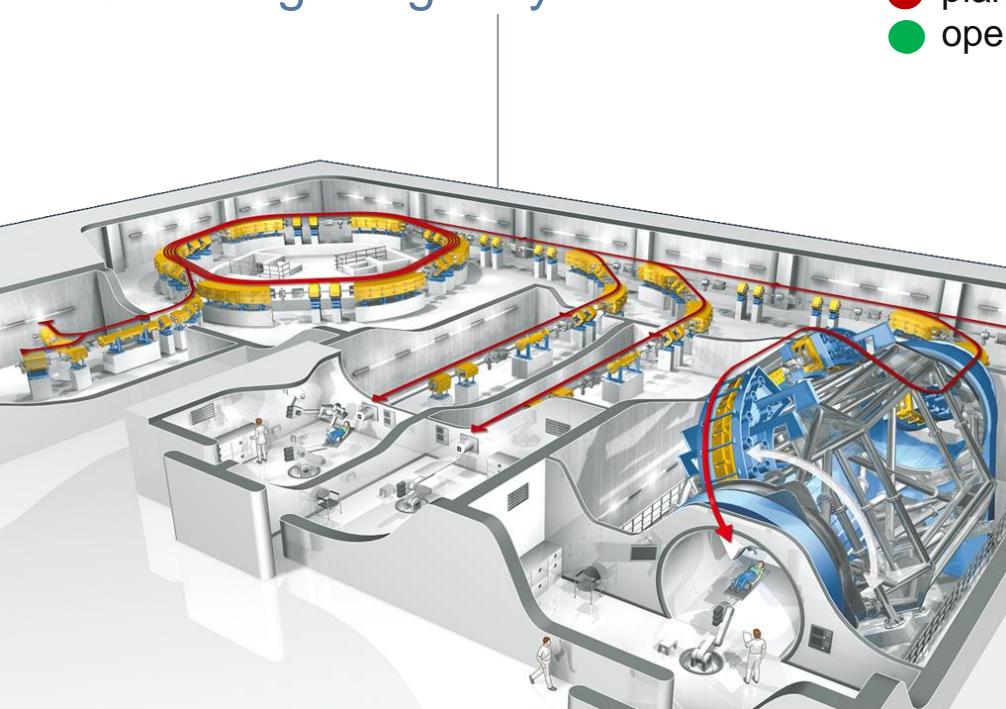


COST Action Nano-IBCT 2010- 2014
Nanoscale insights

E. Surdutovich, E. Scifoni, A.V. Solov'yov,
Mutation Research- Reviews, v.704, p.206-212 (2010)

Present and future of IBCT in Europe

- Heidelberg operational from November 2009, using carbon ions
- 1300 patients to be treated per year ...
- Full clinical integration
- Raster scanning world-wide first scanning ion gantry



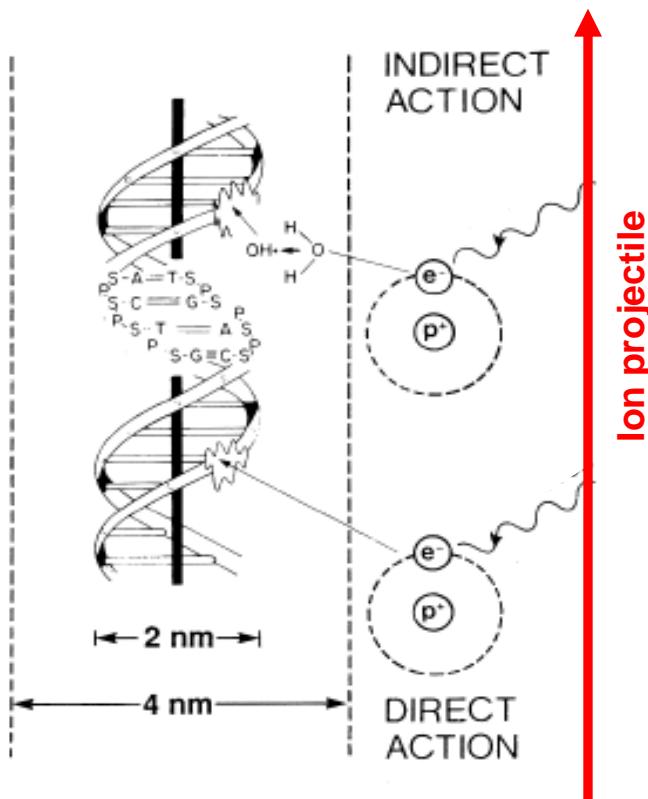
● planned
 ● operational



www.gsi.de

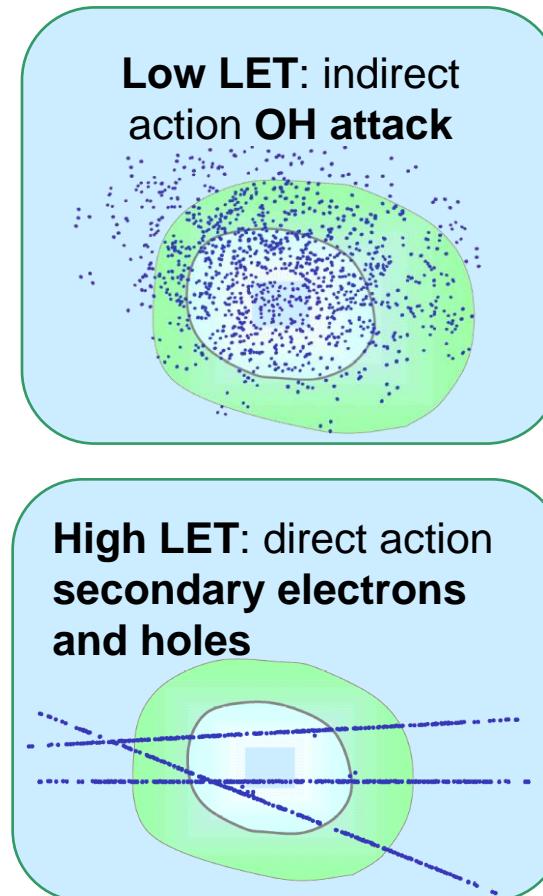
- In construction: Pavia, Kiel, Marburg, Lyon, Vienna, Catania, Maastricht...

Lack of understanding of molecular and nanoscale mechanisms of biodamage

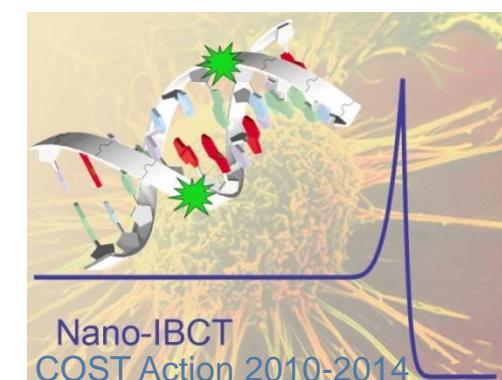
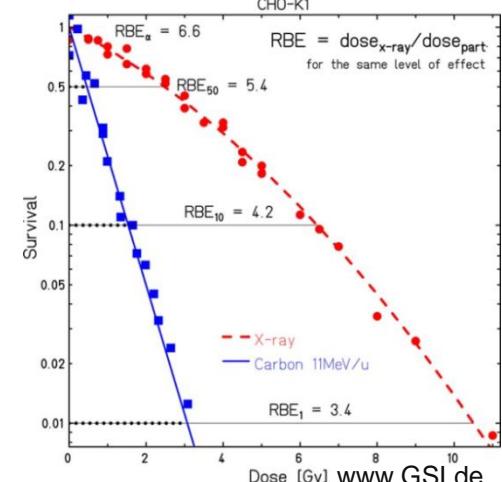


Single Strand Break - SSB
Double Strand Break - DSB

DSB's – the most pernicious DNA damage...
What processes on a microscopic level cause DSBs'?



Relative Biological Effectiveness (RBE): enhanced cell killing at the same dose



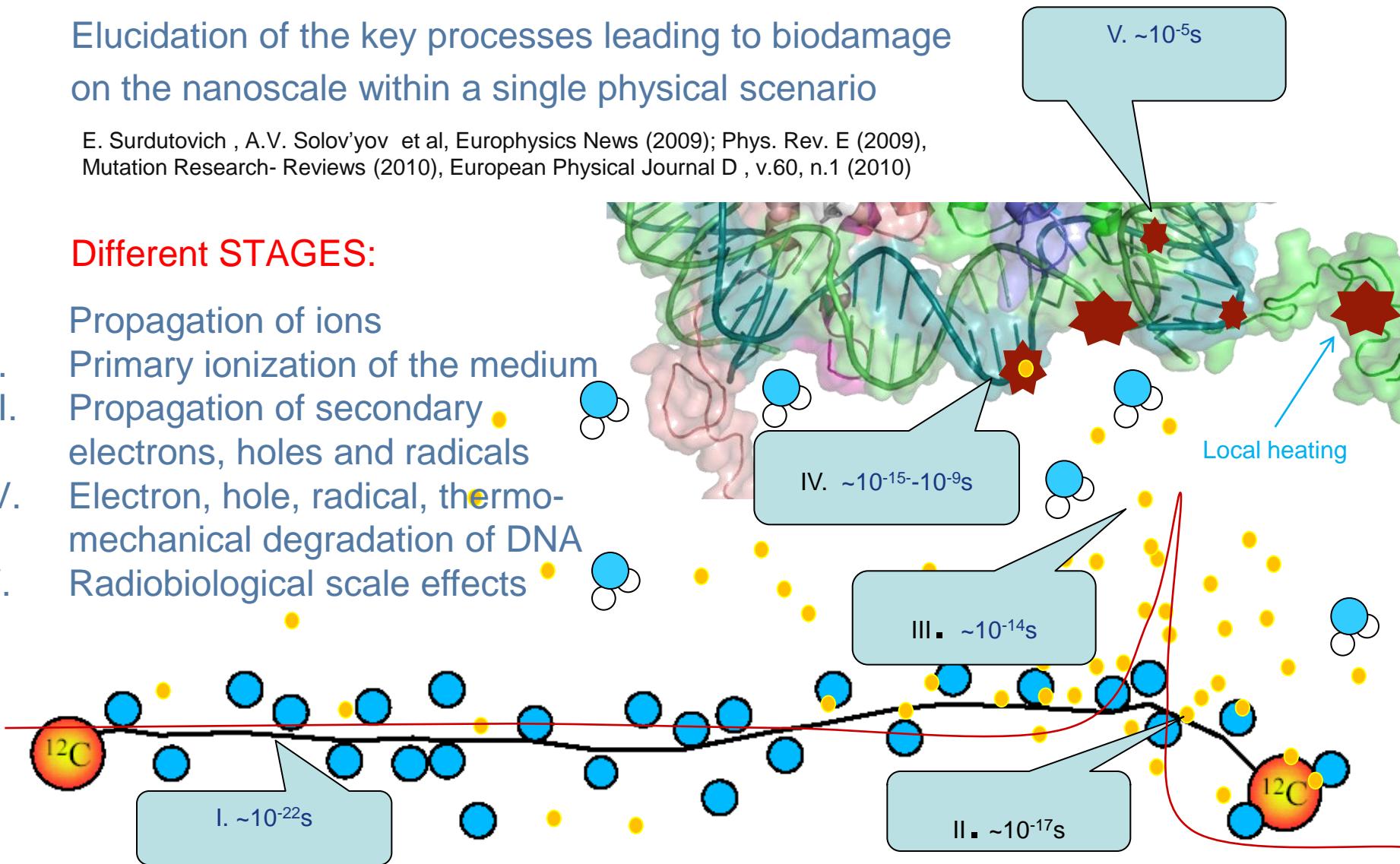
Multiscale approach to the physics of ion beam cancer therapy

Elucidation of the key processes leading to biodamage on the nanoscale within a single physical scenario

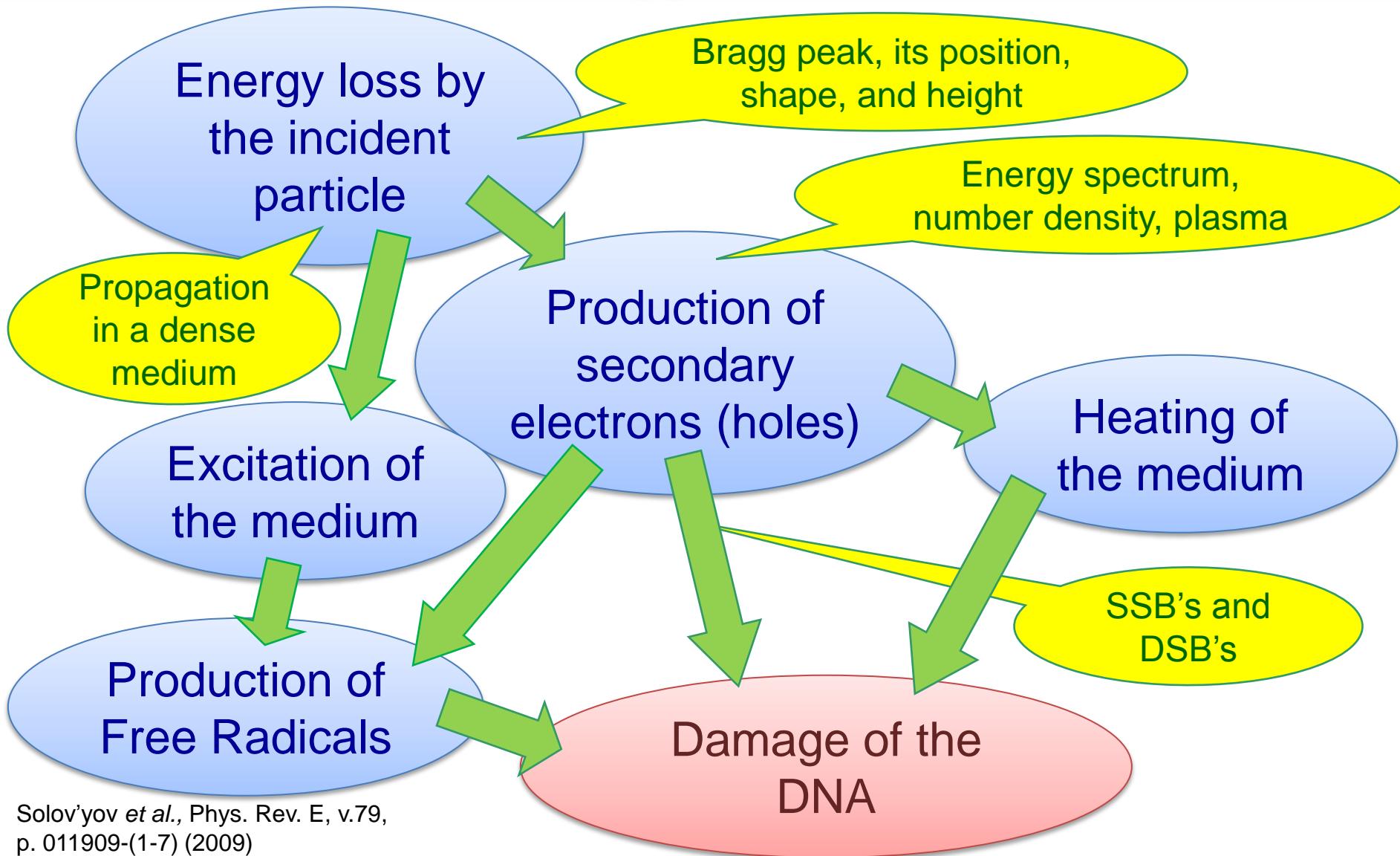
E. Surdutovich , A.V. Solov'yov et al, Europhysics News (2009); Phys. Rev. E (2009), Mutation Research- Reviews (2010), European Physical Journal D , v.60, n.1 (2010)

Different STAGES:

- I. Propagation of ions
- II. Primary ionization of the medium
- III. Propagation of secondary electrons, holes and radicals
- IV. Electron, hole, radical, thermo-mechanical degradation of DNA
- V. Radiobiological scale effects



Multi-scale approach to DNA damage in ion-beam cancer therapy



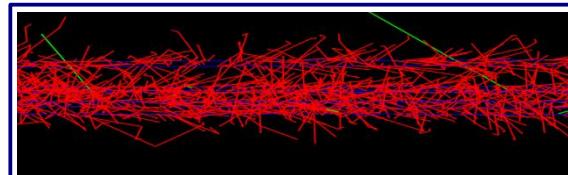
Solov'yov et al., Phys. Rev. E, v.79,
p. 011909-(1-7) (2009)
Europhysicsnews, v.40, n.2, p.21-24 (2009)

Outline

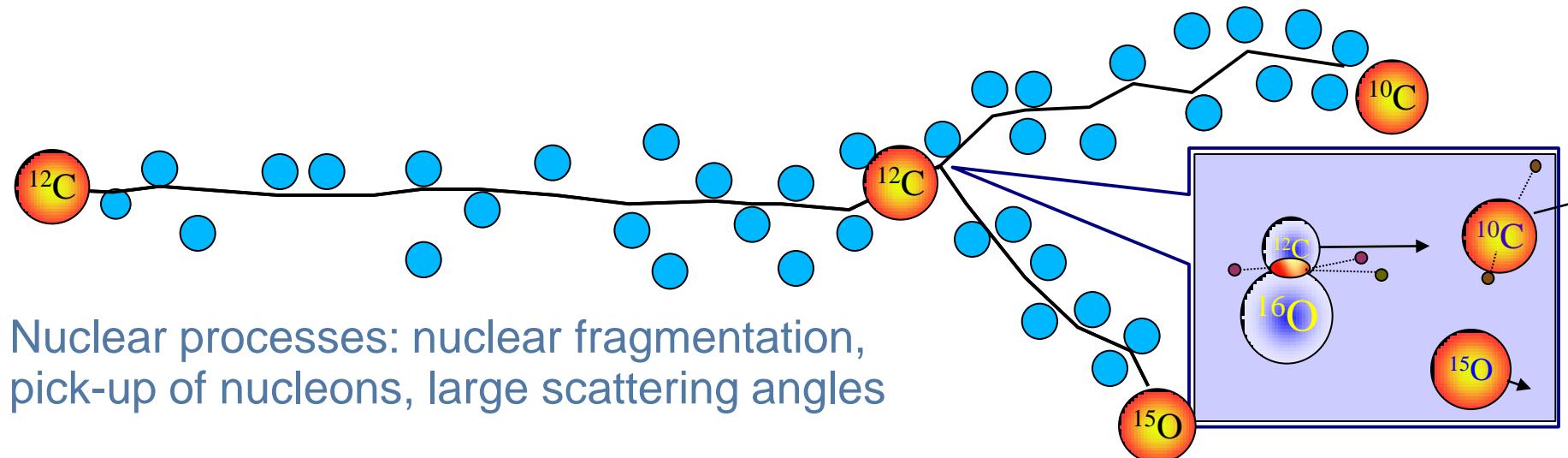
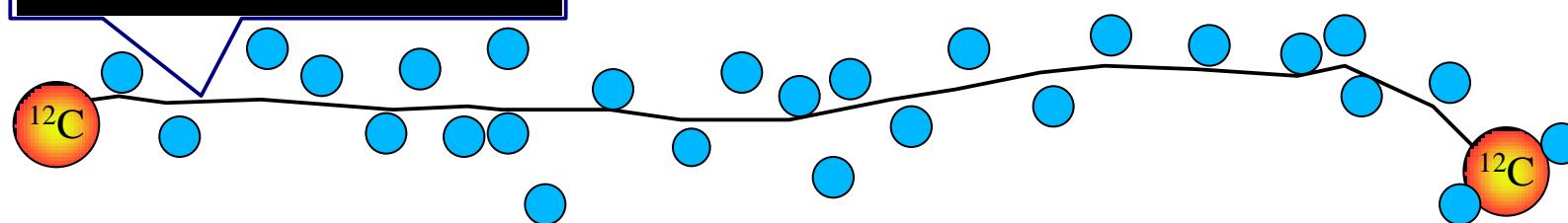
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Transport of ions in target medium

$$-\frac{dE}{dx} \sim 1/\beta^2 \sim 1/E$$



Atomic processes: elastic scattering, scattering with excitation, ionization, multiple Coulomb scattering, small scattering angles

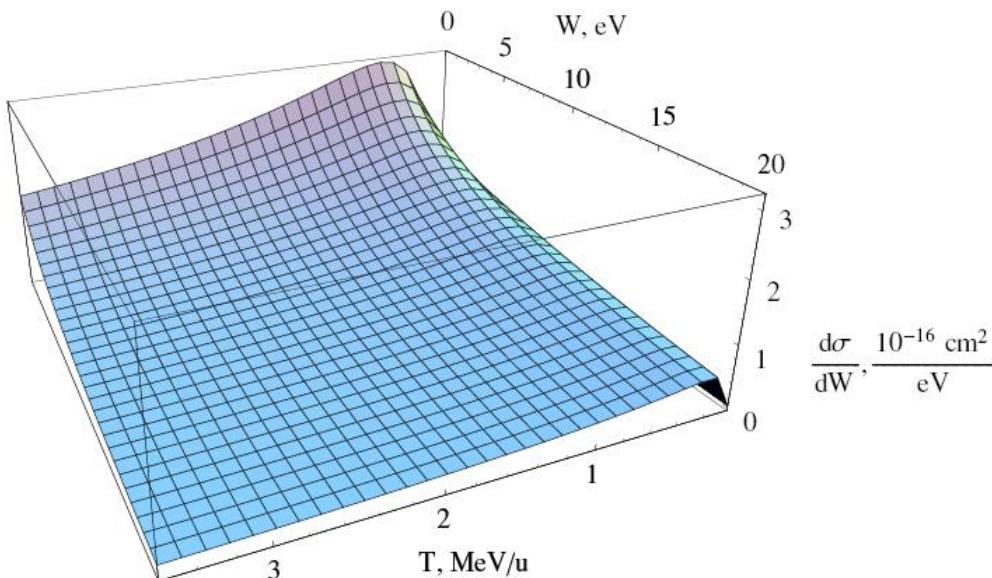


Nuclear processes: nuclear fragmentation, pick-up of nucleons, large scattering angles

Ion stopping in the medium

Ionization cross section (ICS) is the key quantity in modeling of ion energy loss and abundance of secondary electrons

ICS for water by A.V. Solov'yov et al NIMB 266,1623(2008);
E. Surdutovich et al EPJD, v.51, p.63-71 (2009)



Rudd's semi-empirical model

$$\frac{d\sigma}{dW} = \sum_i \frac{N_i}{I_i^2} z^2 f\left(\frac{W}{I_i}, \sqrt{\frac{m}{M}} \frac{T}{I_i}\right)$$

M.E. Rudd et al. Rev. of Mod. Phys., 64, 1992

Beyond Rudd's semi-empirical model

- Relativistic effects (γ up to 1.5) – relativization of the Rudd model

$$v^2 \rightarrow c^2 \beta^2; \ln(1+v^2) \rightarrow \ln\left(\frac{1+v^2}{1-\beta^2}\right) - \beta^2$$

Correct asymptote with relativistic Bethe-Bloch formula for energy loss;
Surdutovich, Scifoni, Solov'yov et al. EPJD , v.51, p.63-71 (2009)

Critical for the position of the Bragg peak

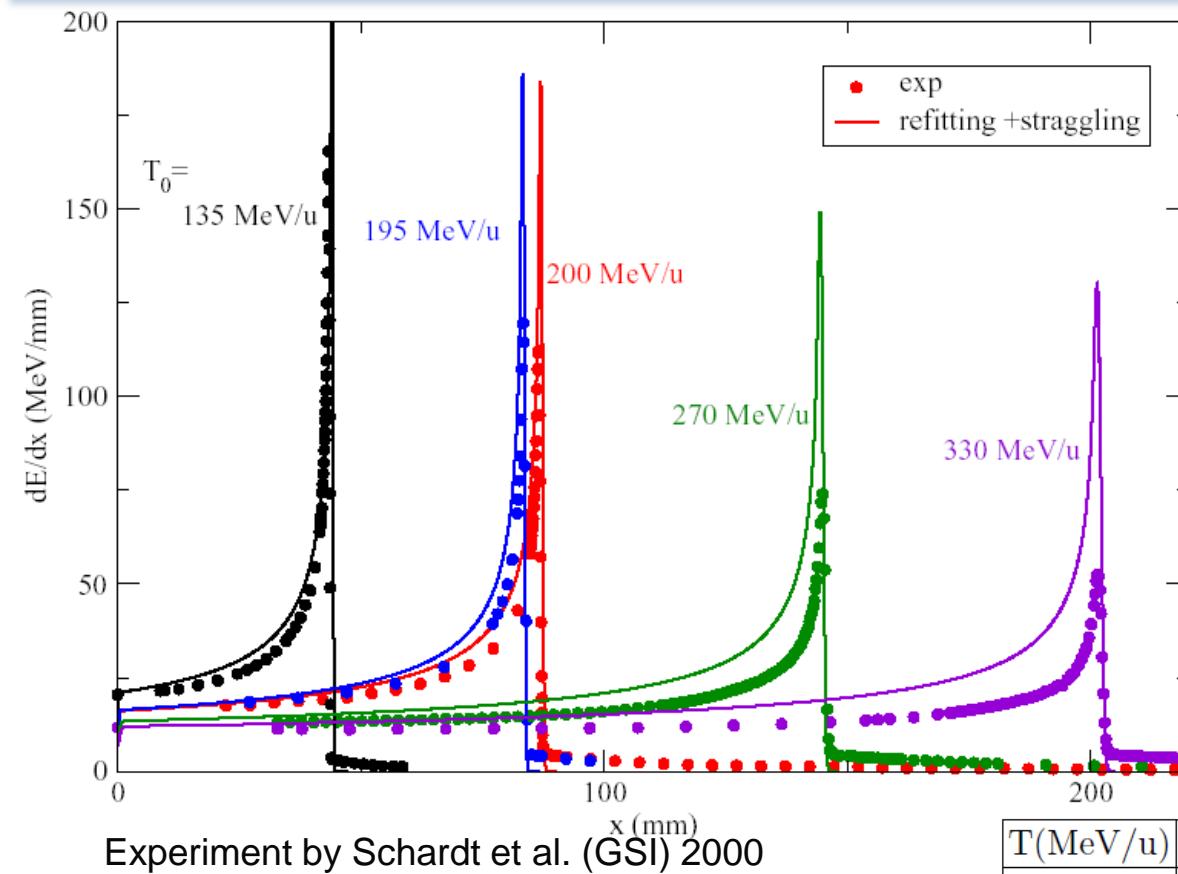
- Charge transfer

$$z_{eff} = z(1 - \exp(-125\beta z^{-2/3}))$$

Barkas H.W., 1963

Critical for the correct height of the Bragg peak

Tuning on the Bragg Peak for carbon ions



Experiment by Schardt et al. (GSI) 2000

Deviation <0.5mm,
i.e. within therapeutic
tolerance

E.Scifoni, E.Surdutovich, A.V. Solov'yov PRE,
v.81, p.021903-(1-7) (2010)

Good agreement
with experimental
positions of Bragg
peak for a large
range of initial
energies T_0

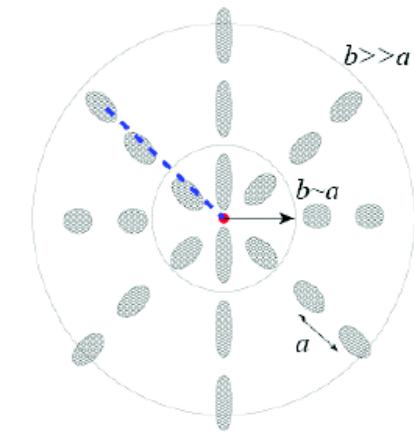
$$\sigma_{str}(x_0) = 0.012x_0^{0.951}/\sqrt{A}$$

$x_0 = \text{max penetration depth}$
 $A = \text{ion mass}$

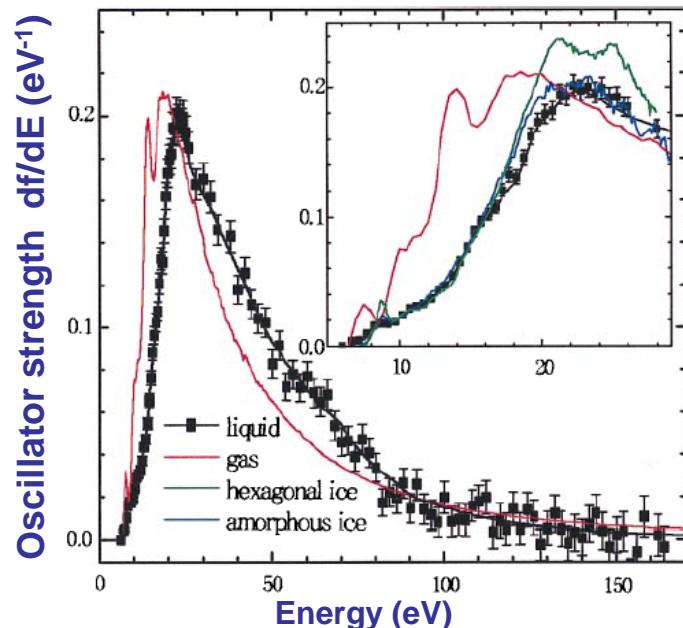
T (MeV/u)	Bragg peak position (mm)			
	us (single ion)	us (straggling)	Schardt	MCHIT
400	275.75	275.01	274.72	274.53
330	201.81	201.30	201.42	200.52
270	144.78	144.42	144.82	143.78
200	87.19	86.94	86.46	86.90
195	83.49	83.26	83.39	
135	44.20	44.05	43.34	43.73

Liquid water: ionization in the medium

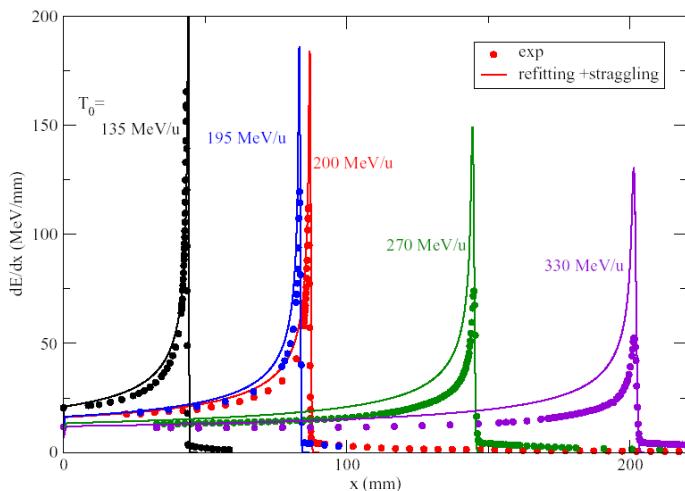
- Polarization and screening of the medium become important for **distant** collisions, with a net effect of suppression of these ionizations
- Optical experiments may provide a key quantity: the **generalized oscillator strength** $df/dE(E, q)$ at $q \rightarrow 0$, the co-called optical limit
- Description of the collision process depends on the energy range of the projectile:
 - » Fast ions $T > 0.2$ MeV/u, Born approx. (BA) is valid
 - » Slow ions $T < 0.2$ MeV/u, BA not fully valid



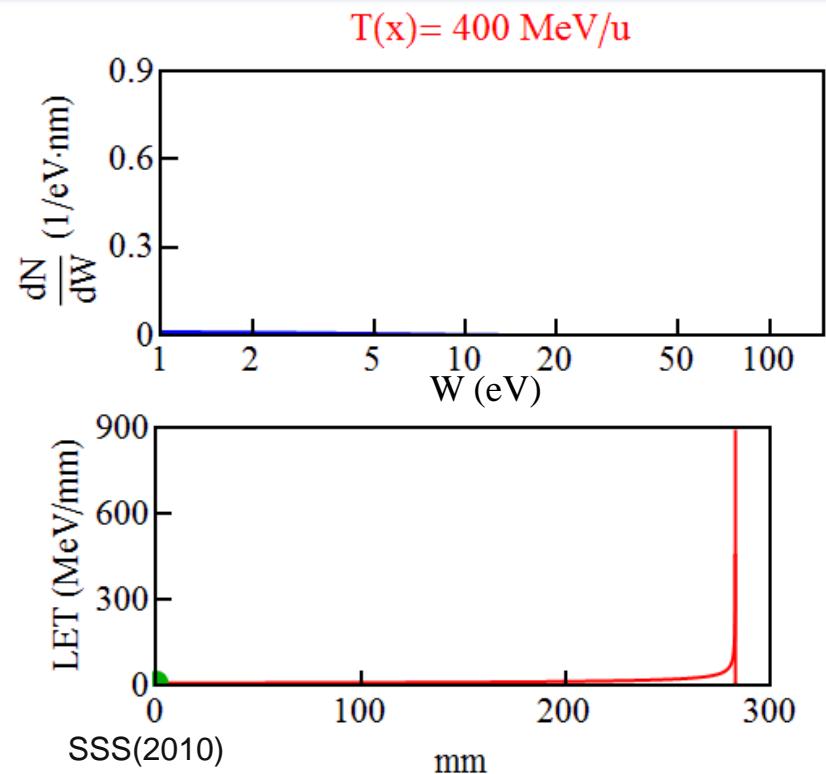
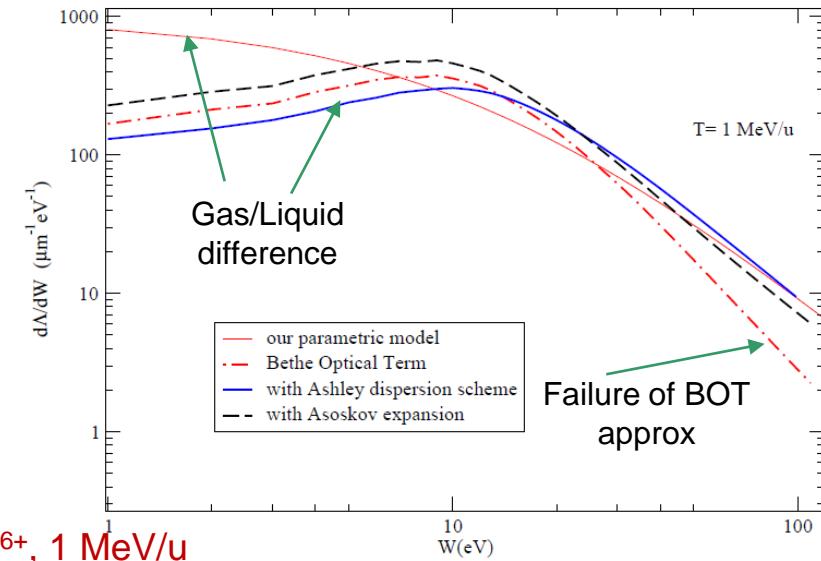
X ray water spectrum
Hayashi,Watanabe 2000



Secondary electrons production in liquid H₂O



Scifoni, Surdutovich, Solov'yov *PRE*, v.81, p.021903-(1-7) (2010);
 A.V. Solov'yov et al., *NIMB* 266, 1623(2008)



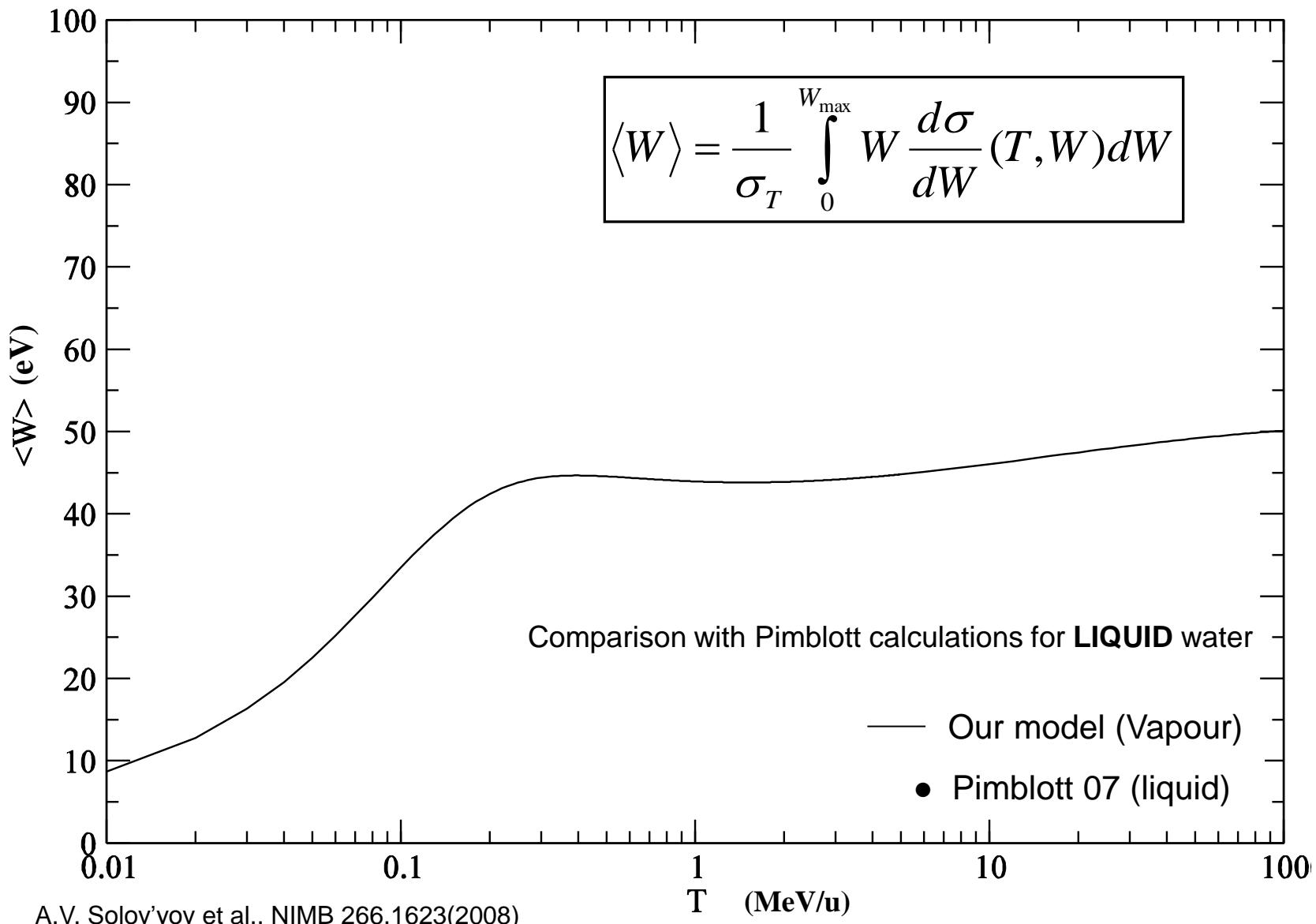
Parametric model

$$\frac{dN}{dW} = nz_{eff}^2(T) \sum_i f\left(\frac{W}{I_i}, \frac{T}{I_i}\right)$$

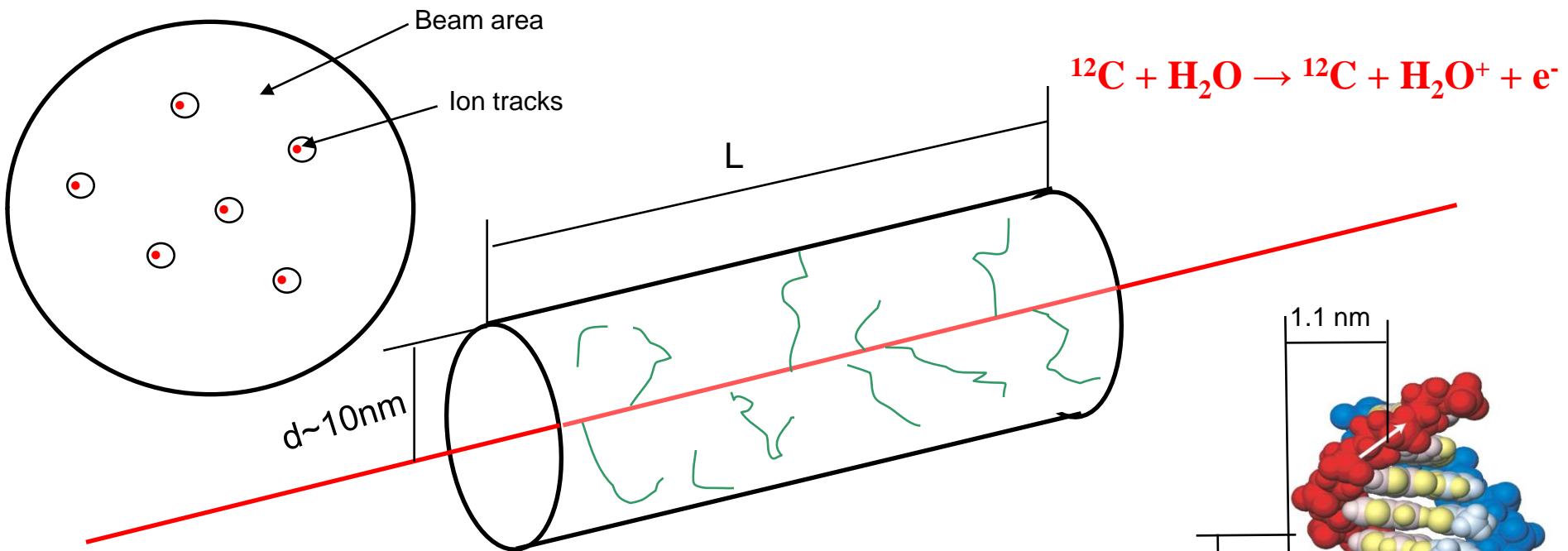
Dielectric function approach

$$\frac{dN}{dW} = \frac{z_{eff}^2(T)}{\pi a_0 T} \int_{k_{min}}^{k_{max}} \eta_2(E, k) \frac{dk}{k},$$

Average energy of secondary electrons



Electron plasma along a heavy ion track



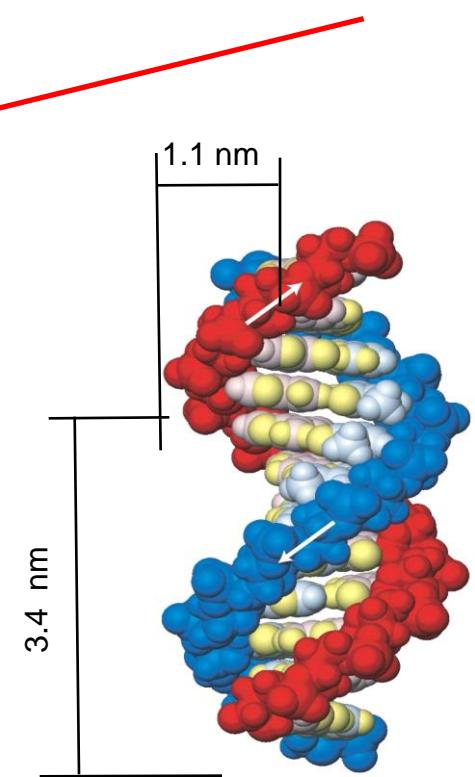
$$\sigma_T(0.1 \text{ MeV}) \approx 0.4 \text{ nm}^2$$

$$n_{\text{water}} = 33.4 \text{ nm}^{-3}$$

$$\text{Number of } e^- / \text{nm} = n_{\text{water}} \sigma_T \approx 13 / \text{nm}$$

$$e^- \text{ plasma density} \approx 13 / \pi d^2 = 0.04 \text{ nm}^{-3} = 4 \times 10^{19} \text{ cm}^{-3}$$

$$\text{Number of } e^- / \text{DNA conv.} = V_{\text{DNA conv}} n_{e^-} = 0.5$$



$$V_{\text{DNA conv}} = 13 \text{ nm}^3$$

E. Surdutovich, O. I. Obolensky, E. Scifoni, I. Pshenichnov, I. Mishustin, A. V. Solov'yov and W. Greiner, EPJD , v.51, p.63-71 (2009)

Targeting DNA with secondary electrons

$$N_{SSB} = \Gamma_{SSB}(W) \sum_k \int d\zeta d\vec{A} \cdot D\nabla \left[P(k, |\vec{r}|) \varepsilon(k, W) \frac{dN(\zeta, W)}{d\zeta} \right]$$

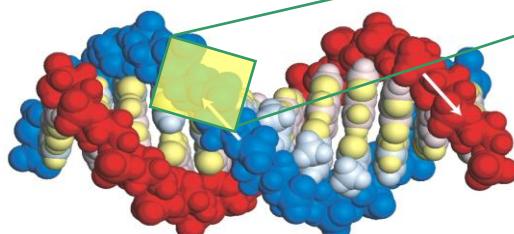
Probability of SSB by e^- with energy W

Element of the DNA's surface

Probability of random walk @ r

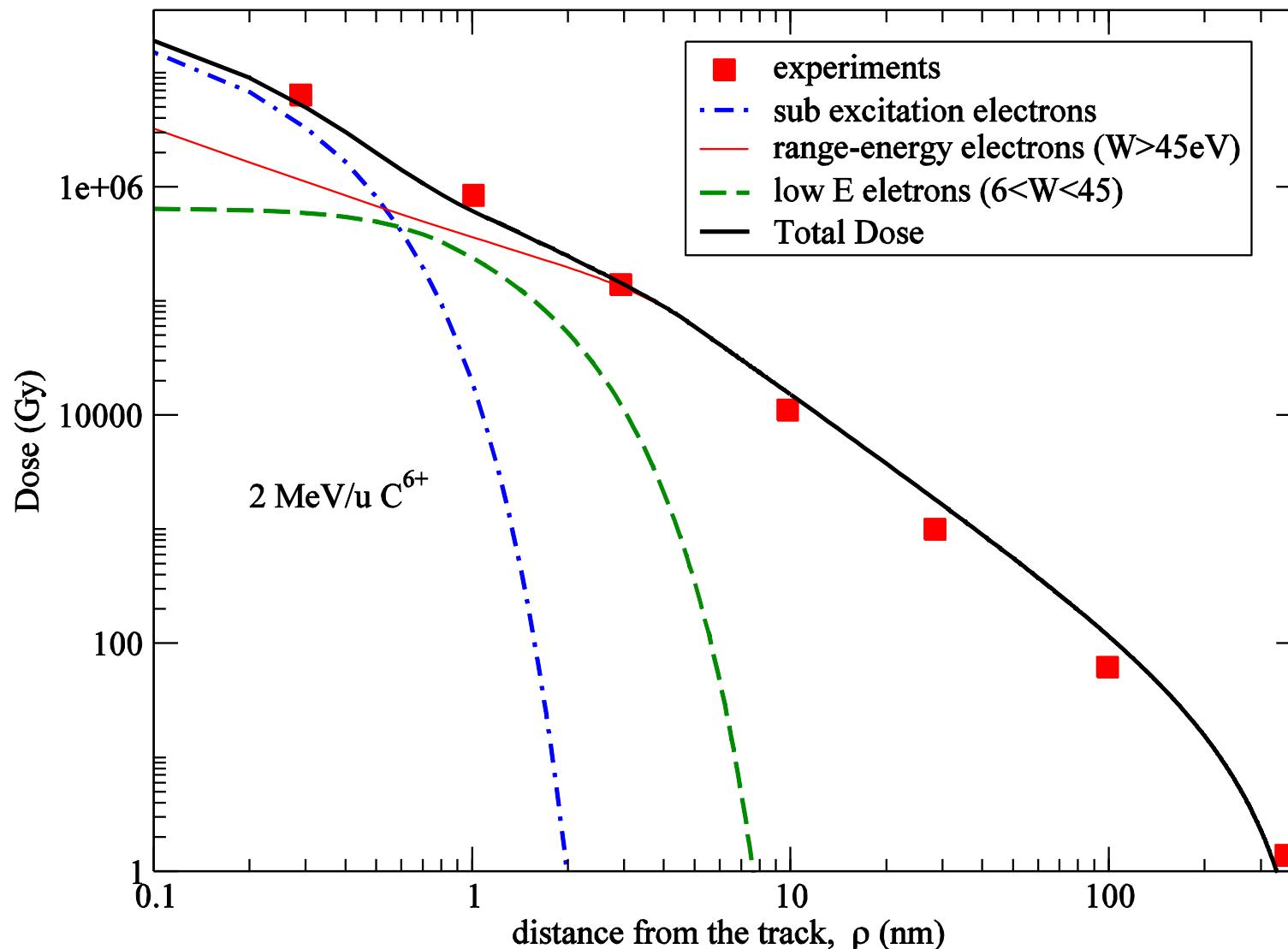
Number of secondary electrons with kinetic energy W produced at a given segment of projectile's trajectory

Attenuation due to inelastic collisions



Solov'yov et al., Phys. Rev. E, v.79, p. 011909-(1-7) (2009)

Radial Dose

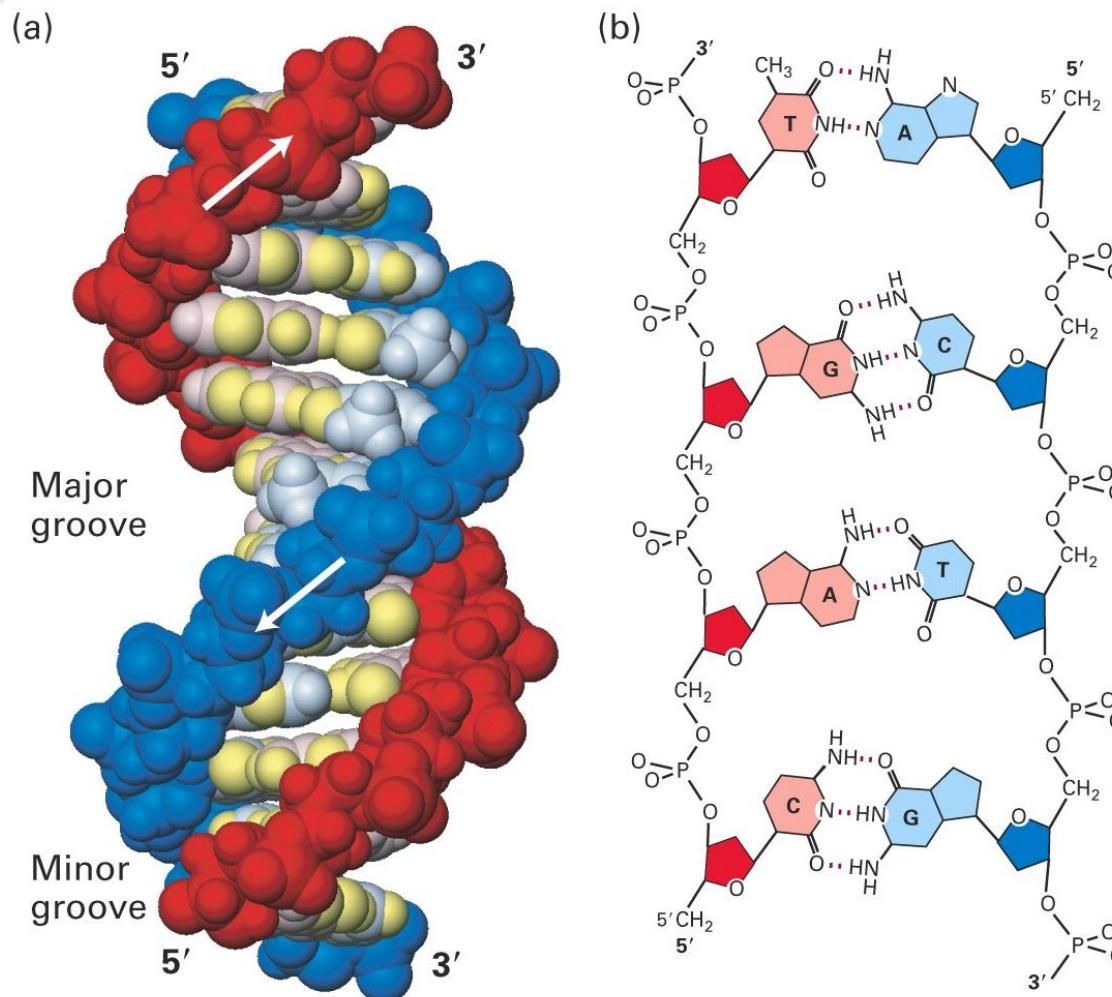


E. Scifoni , A.V. Solov'yov, E. Surdutovich, EPJD, v.60, p.115–119 (2010)

Outline

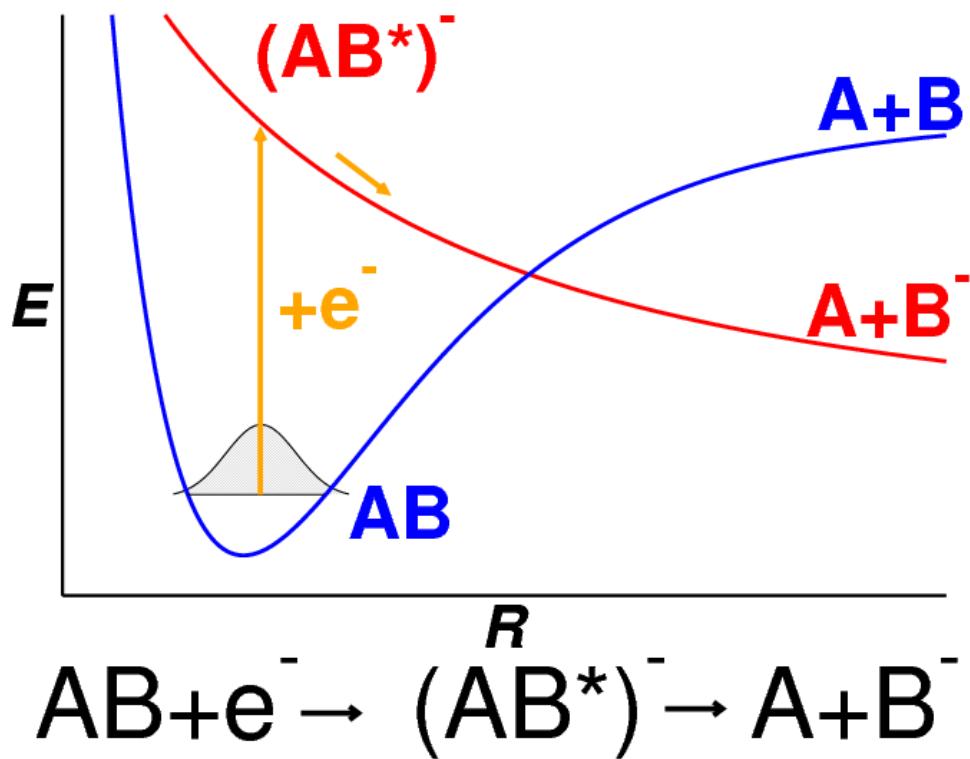
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 - ✓ Electron attack on DNA
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DNA structure



DNA's backbone is made of sugars and phosphate groups joined by ester bonds. Attached to each sugar is one of four types of molecules called bases. Each type of base on one strand forms a bond with just one type of base on the other strand. This is called complementary base pairing. A bonds only to T, and C bonds only to G.

Low energy dissociative electron attachment



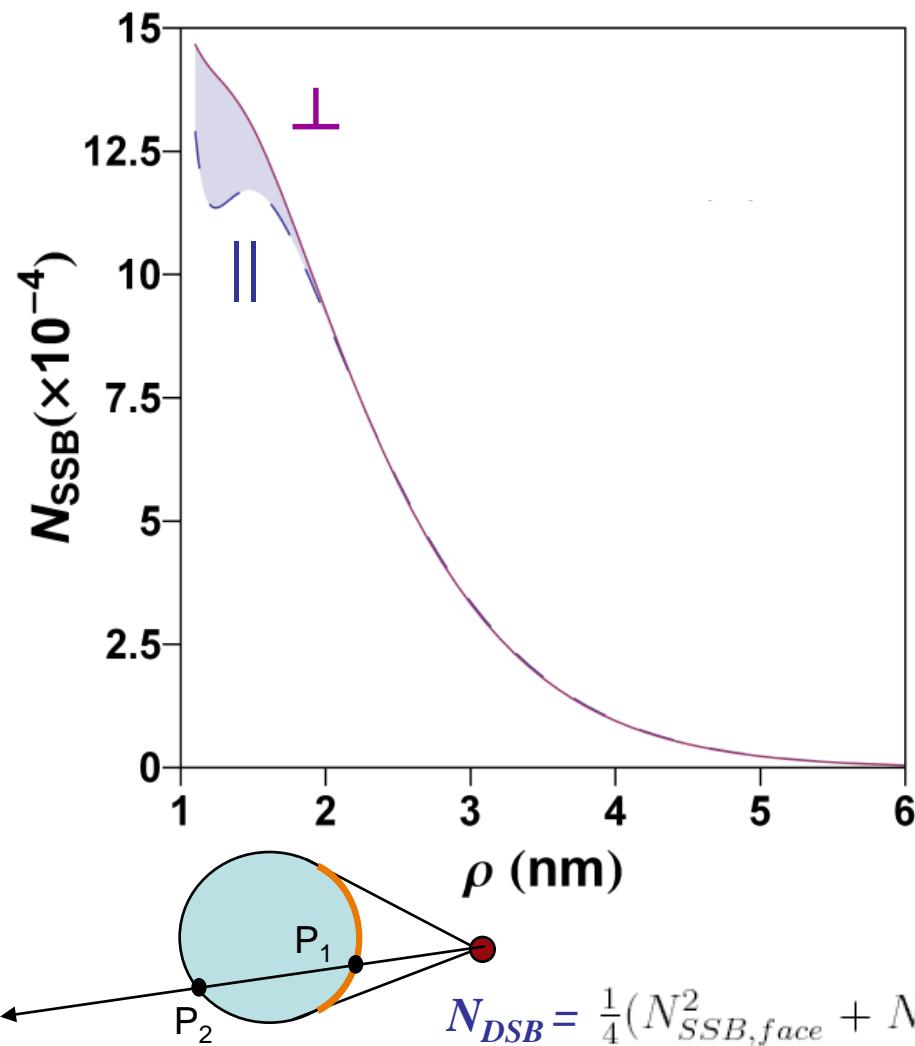
Many experiments (**Sanche**) and quantum calculations on the selected DNA components (**Gianturco, Greene**) confirm the presence of these resonances at low energy.

The resonant mechanism has **comparable** effect to higher energy mechanism

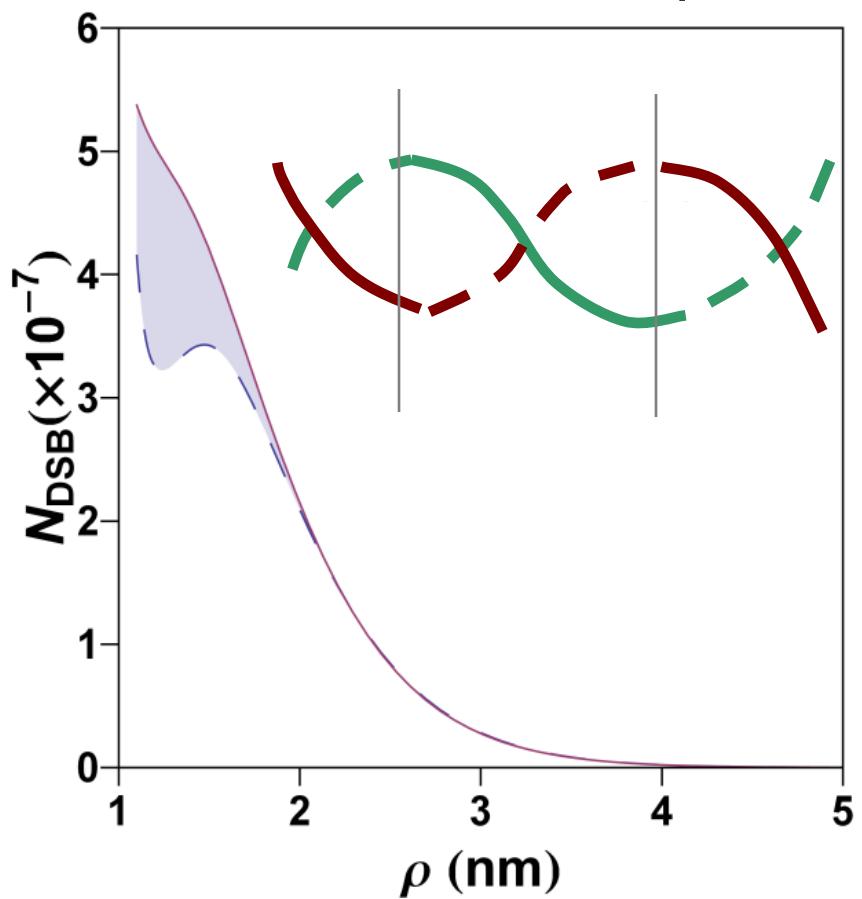
M.A. Huels, B. Boudaïffa, P. Cloutier, D. Hunting, L. Sanche, J. Am. Chem. Soc. 125, 4467 (2003)

DNA double strand breaks

Overall SSB

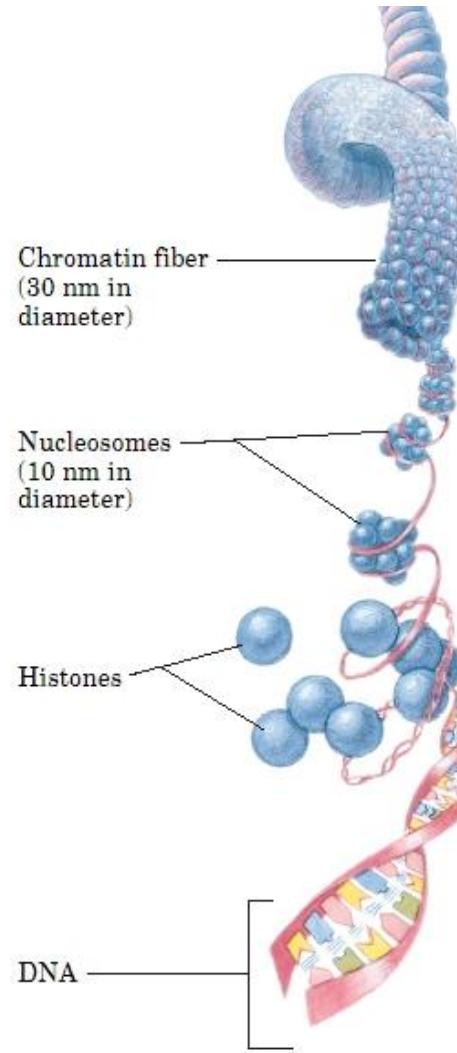
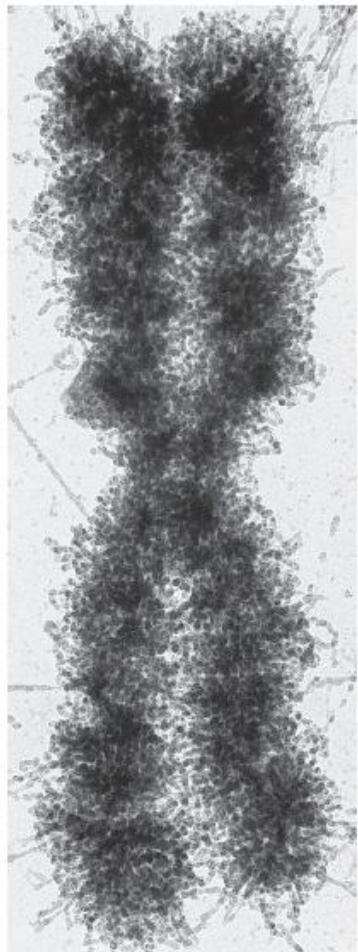


Overall DSB's due to separate e⁻



Solov'yov et al., Phys. Rev. E, v.79, p. 011909-(1-7) (2009)

Generalization for macromolecular targets: packing in histons, chromatid and chromosome



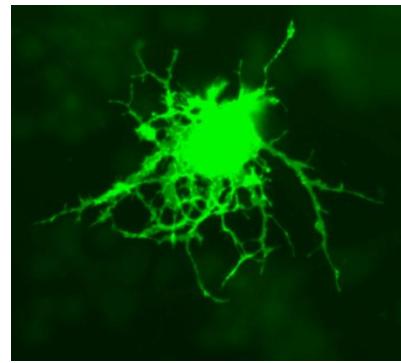
Averaging over orientations of a decamer all along the chain by introducing a distribution function for decamer orientations

Preliminary estimates on cell damage

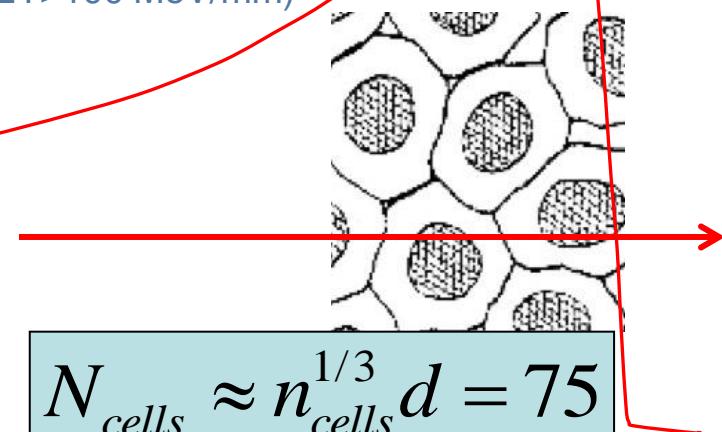
An example:

Glial cells comprise
90% of brain cells

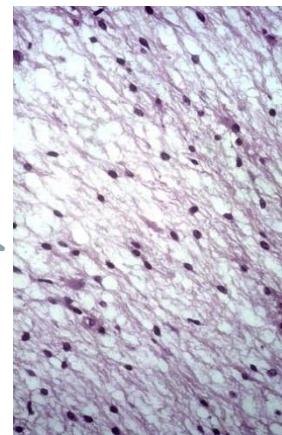
How many cells
may be affected by
a single ion near
Bragg peak?



The Bragg peak
width of a single
C-ion, $d=1\text{mm}$
($\text{LET} > 100 \text{ MeV/mm}$)



With how many
nuclei will this ion
interact ?



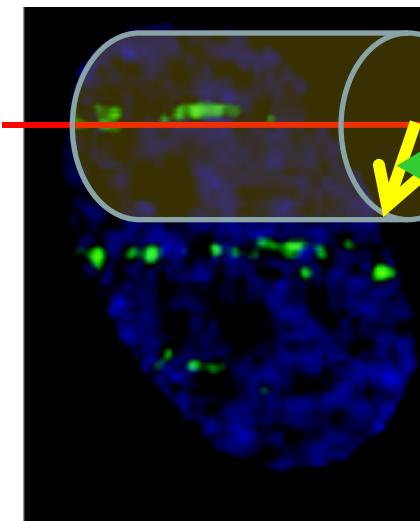
$$N_{cells} \approx n_{cells}^{1/3} d = 75$$

$$N_{nucl} \approx \frac{\sigma_n}{\sigma_c} N_{cell} = 12$$

Solov'yov *et al.*, Phys. Rev. E, v.79,
p. 011909-(1-7) (2009)

Number of DSB's per nucleus

1. Probability of a DSB vs. distance between the track and a convolution
2. Account for larger ionization cross sections inside the nucleus
3. Distribution of convolutions: uniform inside the nucleus for cells in interphase



Number
of DSB's

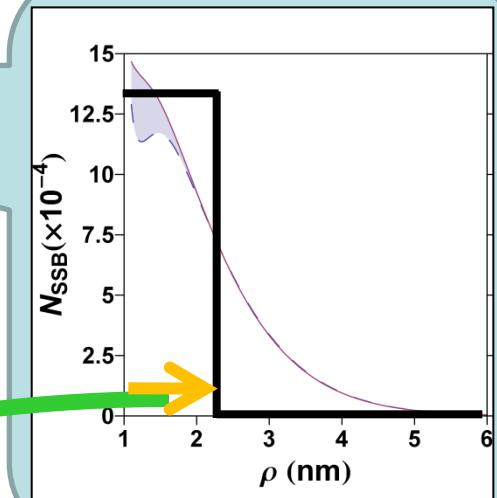
=

$$\times P_{DSB} \times n_{conv} \approx 3.5 \mu\text{m}^{-1}$$

$$\approx 18 \text{ per nucleus}$$

No experiments for glial cells; for **human fibroblasts**, GSI reports **2.6 DSBs/ μm**

Jakob, Splinter and Taucher-Scholz., Rad. Res. 171, 405 (2009).



$$n_{conv} \approx 4.3 \times 10^6 \mu\text{m}^{-3}$$

Solov'yov *et al.*, Phys. Rev. E, (2008)

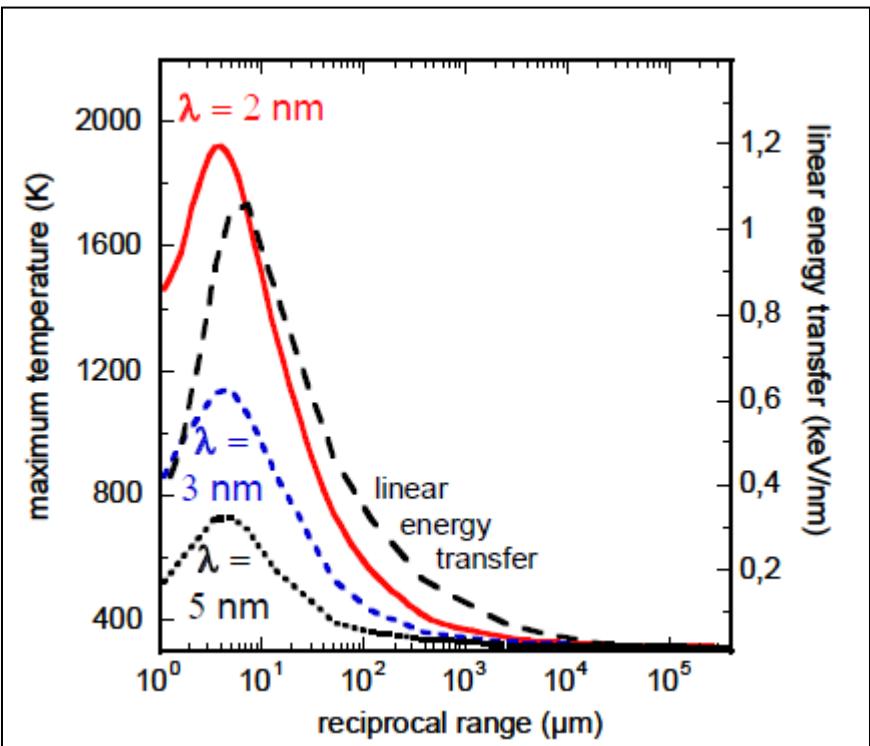
Heat transfer from the ion track: thermal spike model

Coupled thermal conductivity equations for electron and lattice components

$$C_e \frac{\partial T_e}{\partial t} = \nabla \cdot (K_e \nabla T_e) - g(T_e - T) + B(r, t)$$

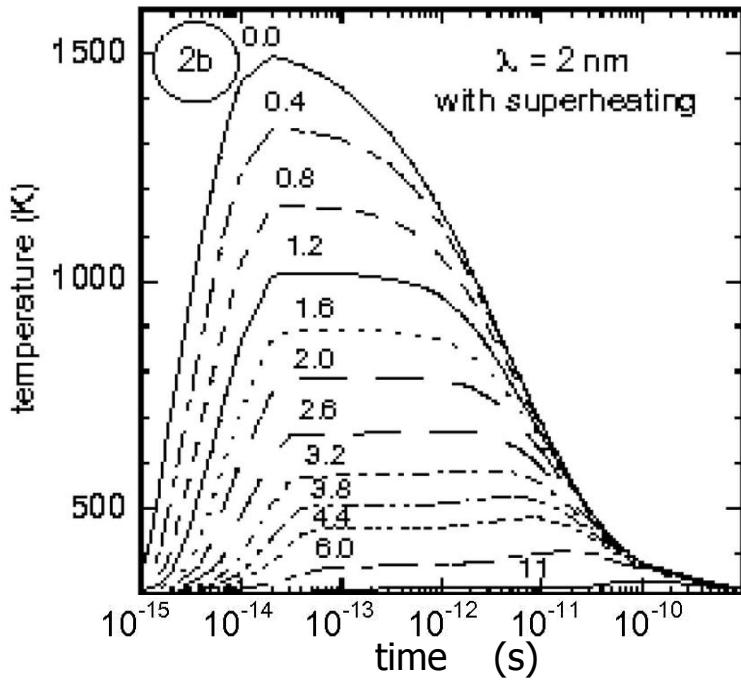
$$\rho C \frac{\partial T}{\partial t} = \nabla \cdot (K \nabla T) + g(T_e - T)$$

- $B(r, t)$ is the energy density supplied by the incident ion to the electronic system by ballistic collisions at a distance r from the track at time t .
- g – electron-phonon coupling
- λ –electron-phonon mean free path



M. Toulemonde, E. Surdutovich and A.V. Solov'yov ,
Phys. Rev.E, v.80, p. 031913 (2009).

Thermal and pressure spikes: mechanical mechanism of DNA damage?



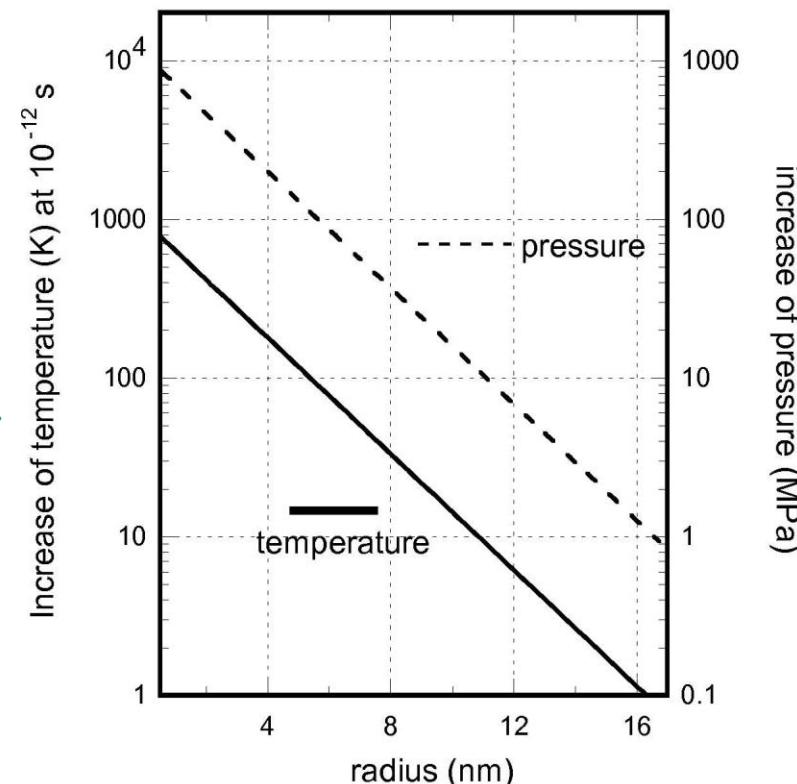
Temperature and pressure vs.
distance from the ion track

Estimated forces acting on sections
of DNA are over 100 pN.

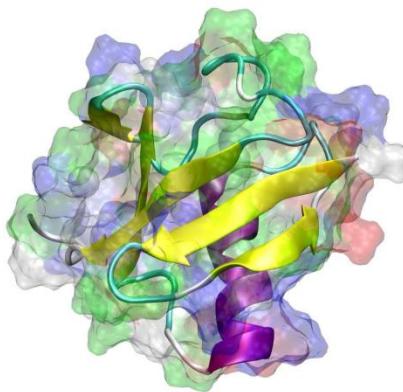
Can this cause strand breaks?

M. Toulemonde, E. Surdutovich and A.V. Solov'yov ,
Phys. Rev.E, v.80, p. 031913 (2009).

Temperature dependence on time in
the vicinity of Bragg peak for C ions at
different distances from the ion track

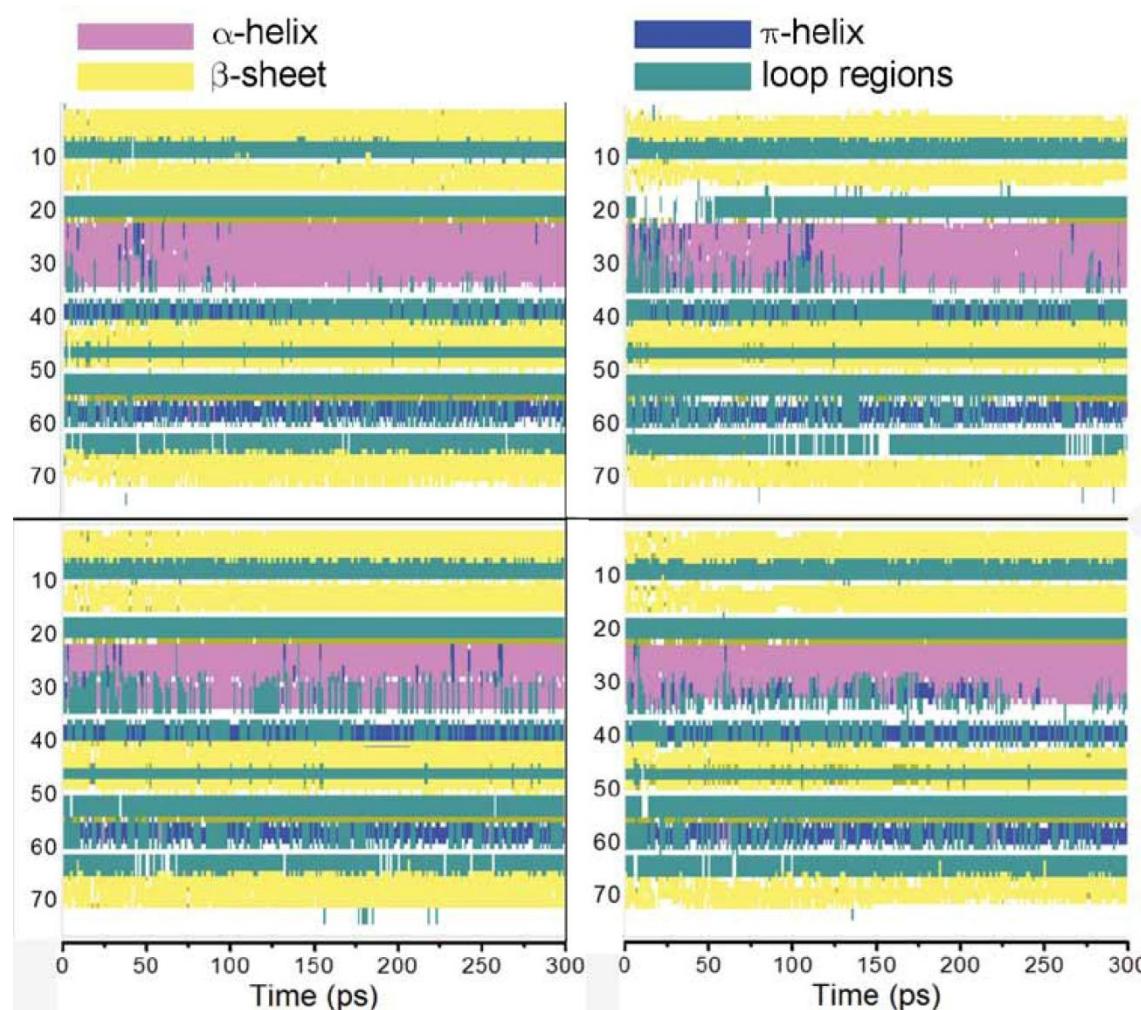


Effects of thermal spikes on biomolecules



Ubiquitin

- Explore molecular dynamics at the conditions of a thermal spike
- Explore effects of the corresponding temperature and pressure dependences on biomolecules and other molecules
- Consider the consequences of residual temperature increase

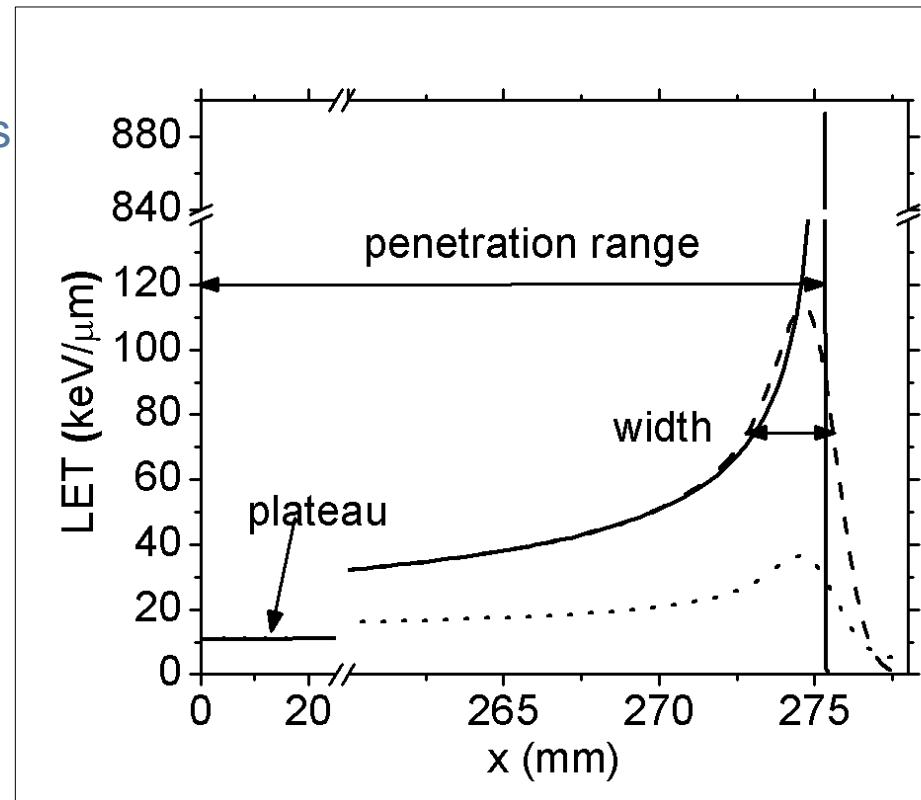


E. Surdutovich, A.V. Yakubovich and A.V. Solov'yov , Topical Issue on Biodamage, EPJD , v.60, n.1 (2010)

Energy deposition: geometry and pressure

Goal: model the environment following the thermal spike

- LET on the length of the Bragg peak =0.9keV/nm; determines the energy loss by an ion/nm
- Characteristic sizes:
 - Characteristic time (1000 electron collisions) is 0.1ps.
 - C-ion's energy =0.3 MeV/u
 - Per 0.1ps, it passes 0.8 μm
 - Secondary electrons wander within 1nm radius
- Strength of the explosion:
 - Pressure inside the cylinder is higher than 10^3 MPa
 - Pressure outside the cylinder is 0.1 MPa



$$\frac{P_{in}}{P_{out}} \gg \frac{\gamma + 1}{\gamma - 1} = 10, \gamma = \frac{C_p}{C_v} = 1.22$$

Holds up to t=1ns;

Hydrodynamics: strong explosion

L. I. Sedov (1946) and
J. von Neumann (1947)

$$\begin{aligned}\frac{\partial \rho}{\partial t} + \frac{\partial \rho v}{\partial r} + \frac{\rho v}{r} &= 0 \\ \frac{\partial v}{\partial t} + v \frac{\partial v}{\partial r} &= -\frac{1}{\rho} \frac{\partial P}{\partial r} \\ \left(\frac{\partial}{\partial t} + v \frac{\partial}{\partial r} \right) \ln \frac{P}{\rho^\gamma} &= 0\end{aligned}$$

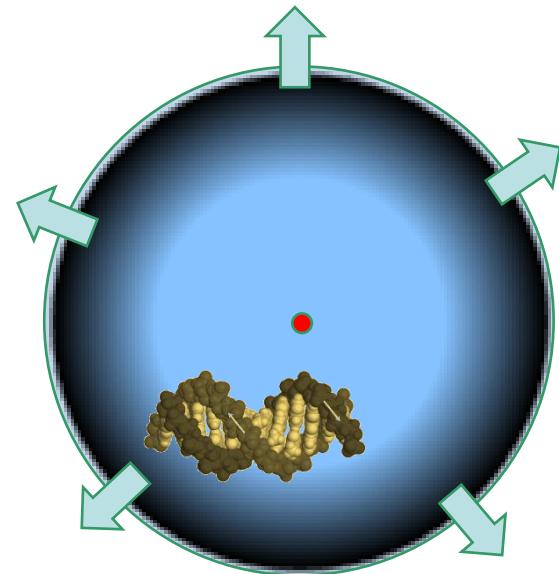
Hydrodynamic equations

$$\xi = \frac{r}{\beta \sqrt{t}} \left[\frac{\rho_1}{\varepsilon} \right]^{1/4}$$

Self-similar flow

E. Surdufovich and A.V. Solov'yov, PRE 82, 051915 (2010)

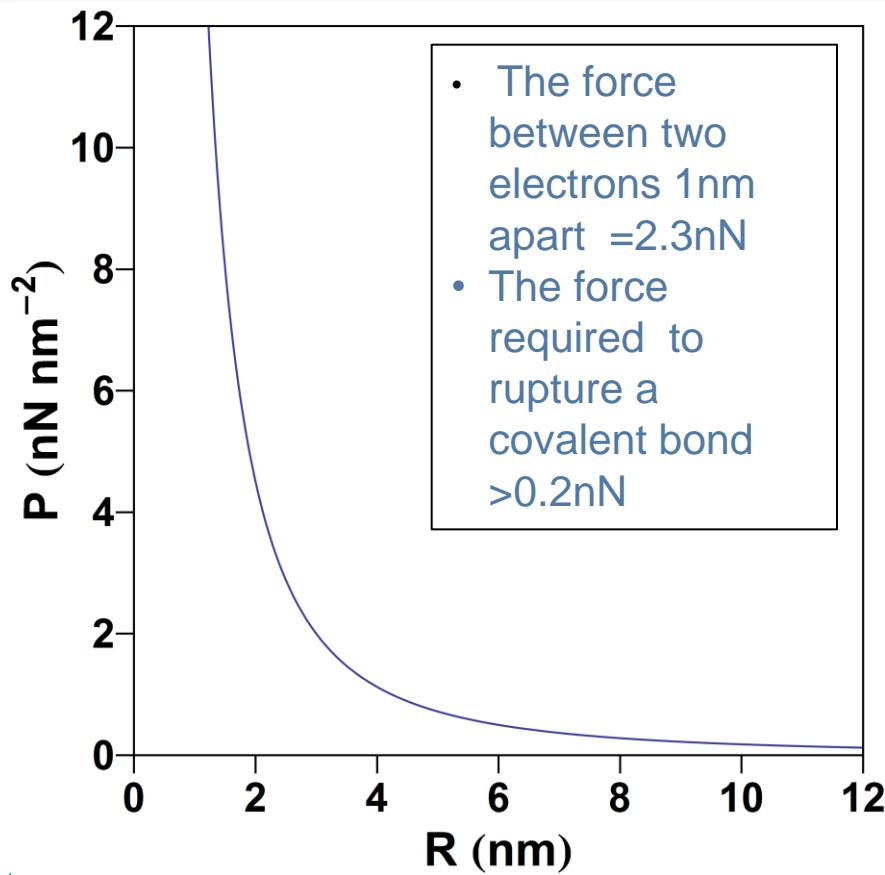
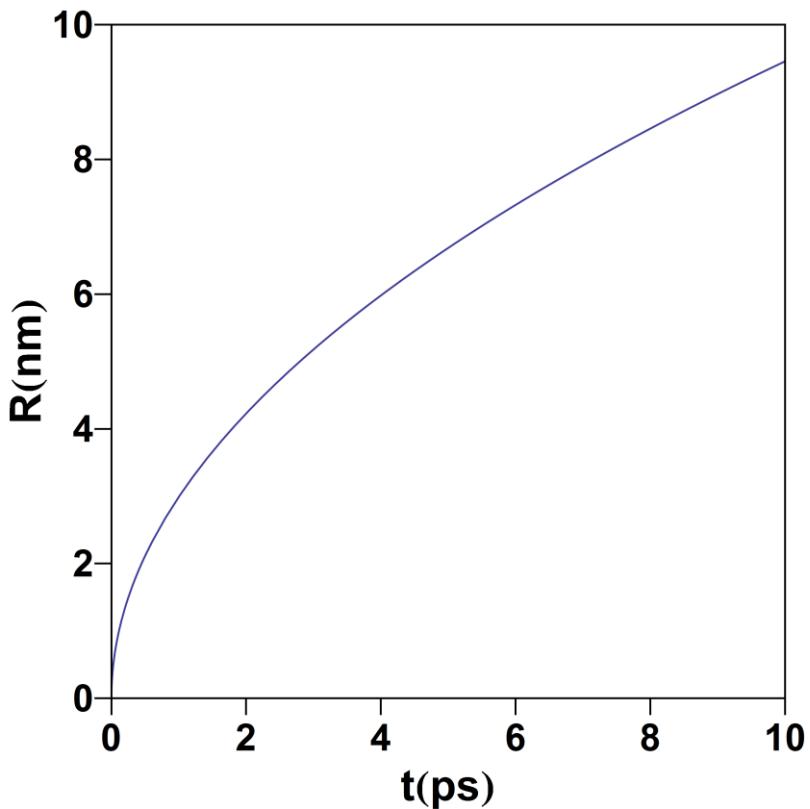
Cylindrical shockwave



Explore the conditions on
the wave front and in the
wake

On the front of shockwave

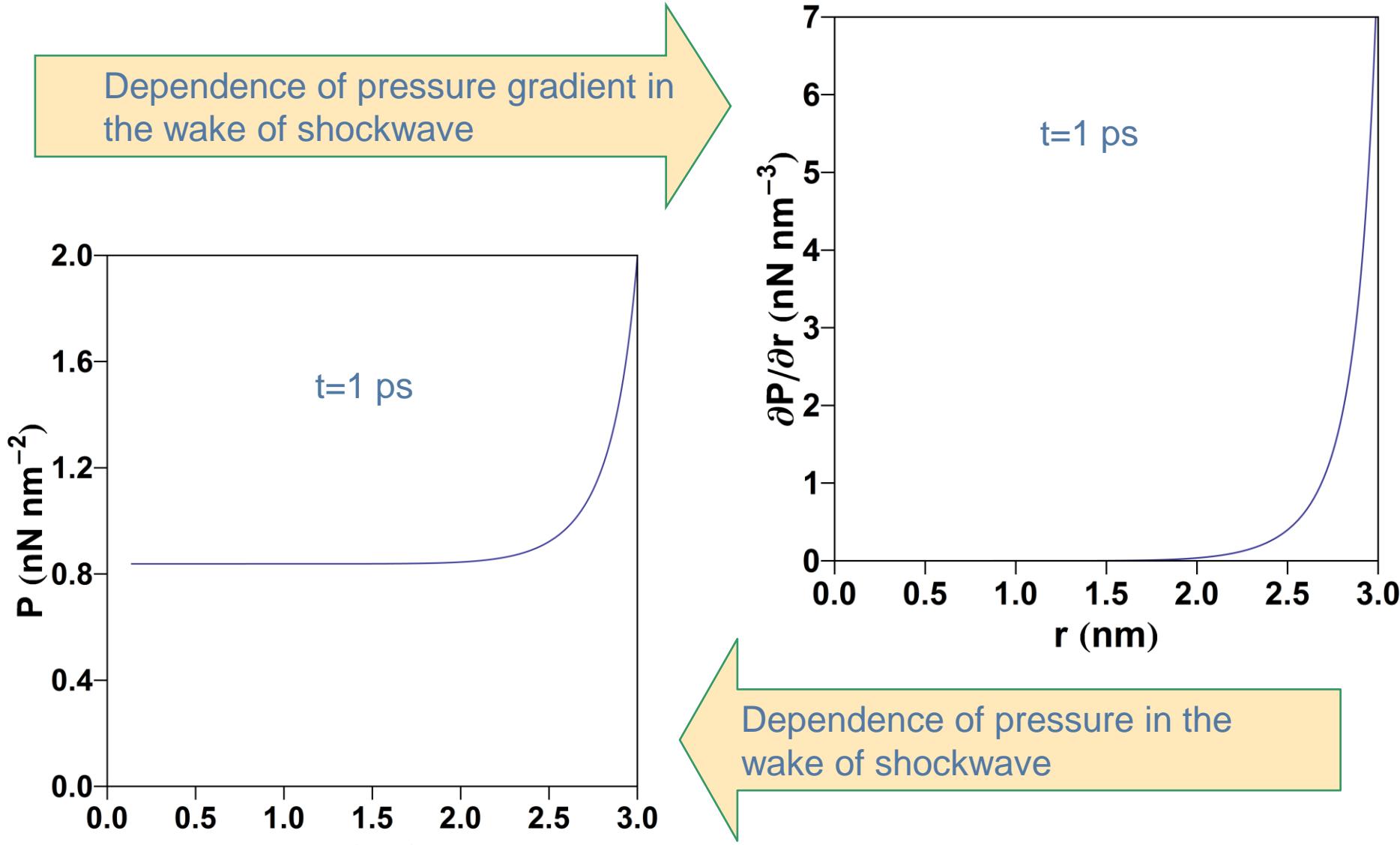
Dependence of pressure on the shockwave front on the front radius



Propagation of shockwave's front on time

E. Surdutovich and A.V. Solov'yov, PRE 82, 051915 (2010)

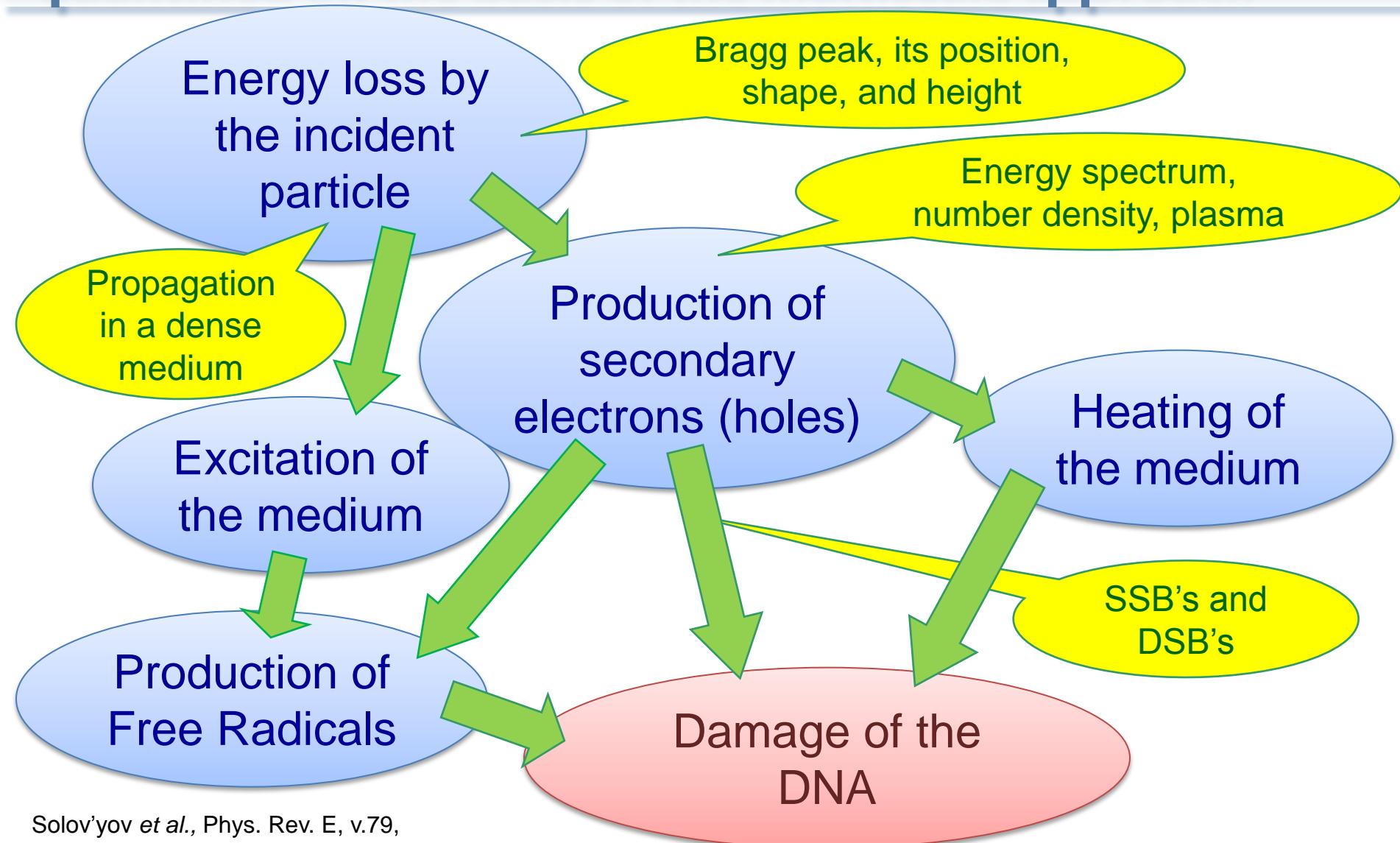
In the wake of shockwave



Dependence of pressure in the wake of shockwave

E. Surdutovich and A.V. Solov'yov, PRE 82, 051915 (2010)

Conclusion I: radiation biodamage can be quantified on the basis of the multiscale approach



Solov'yov et al., Phys. Rev. E, v.79,
p. 011909-(1-7) (2009)
Europhysicsnews, v.40, n.2, p.21-24 (2009)

Conclusion II: quantification of biodamage on the cellular scale require further research efforts

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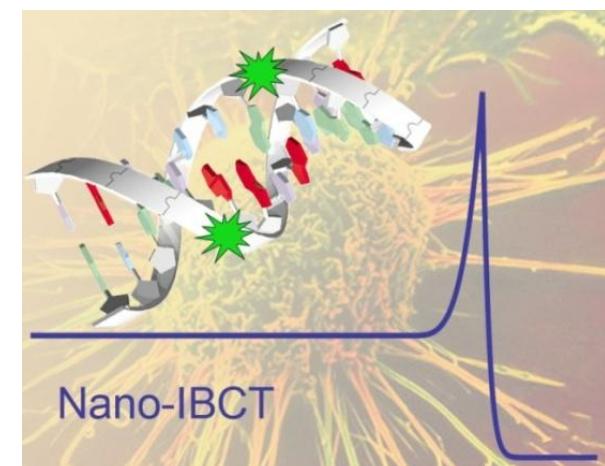
Atomic, Molecular,
Optical and Plasma
Physics

Molecular level assessments of
radiation biodamage

Topical issue edited by Isabella Baccarelli, Franco A. Gianturco,
Emanuele Scifoni, Andrey V. Solov'yov and Eugene Surdutovich



COST Action Nano-IBCT
'Nanoscale insights into ion beam cancer therapy'



Nano-IBCT

Memorandum of Understanding (MoU)
http://www.cost.esf.org/domains_actions/mpns/Actions/nano-ibct/%28glossary%29/off

MoU has been signed by 17 EU countries:
 Austria, Belgium, Czech Republik, Denmark, Finland, France, Germany, Hungary, Italy, Ireland, Netherlands, Portugal, Serbia, Slovakia, Spain, Sweden, Greatbritain;
other 4 EU countries are expected to sign

Australia, USA and Japan will join the COST Action

Thank you for your attention!