

# $\Lambda(1405)$ and the antikaon-nucleon potential




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# Contents



## $\Lambda(1405)$ in meson-baryon scattering

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);

T. Hyodo, M. Niiyama, arXiv: 2010.07592 [hep-ph], to appear in PPNP



## $\bar{K}N$ potentials and their applications

K. Miyahara, T. Hyodo, PRC 93, 015201 (2016);

K. Miyahara, T. Hyodo, W. Weise, PRC 98, 025201 (2018)

### - Kaonic nuclei

S. Ohnishi, W. Horiuchi, T. Hoshino, K. Miyahara, T. Hyodo, PRC95, 065202 (2017)

### - Kaonic deuterium

T. Hoshino, S. Ohnishi, W. Horiuchi, T. Hyodo, W. Weise, PRC96, 045204 (2017)

### - $K^-p$ correlation function

Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise. PRL124, 132501 (2020)

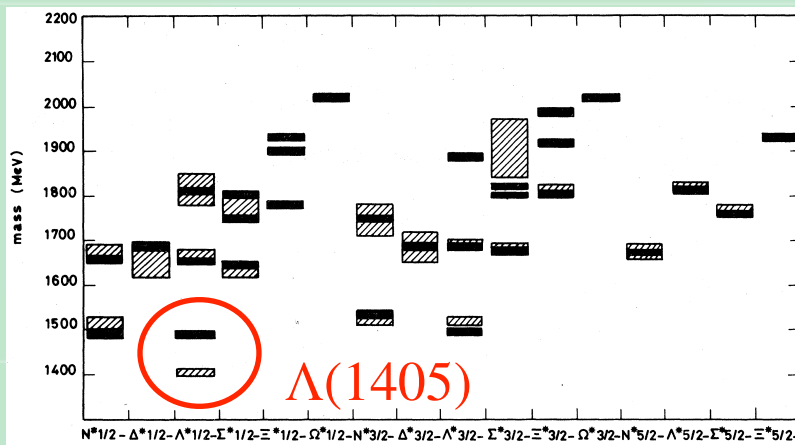
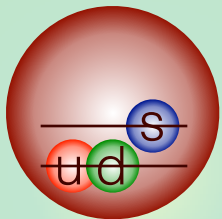


## Summary

# $\Lambda(1405)$ and $\bar{K}N$ scattering

$\Lambda(1405)$  does not fit in standard picture  $\rightarrow$  exotic candidate

N. Isgur and G. Karl, Phys. Rev. D18, 4187 (1978)

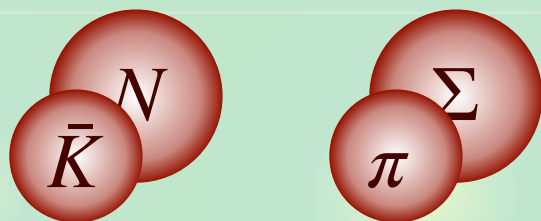


— : theory

▨ : experiment

## Resonance in coupled-channel scattering

- coupling to MB states



energy  $\uparrow$

—  $\bar{K}N$  threshold

▬  $\Lambda(1405)$

—  $\pi\Sigma$  threshold

Detailed analysis of  $\bar{K}N$ - $\pi\Sigma$  scattering is necessary.

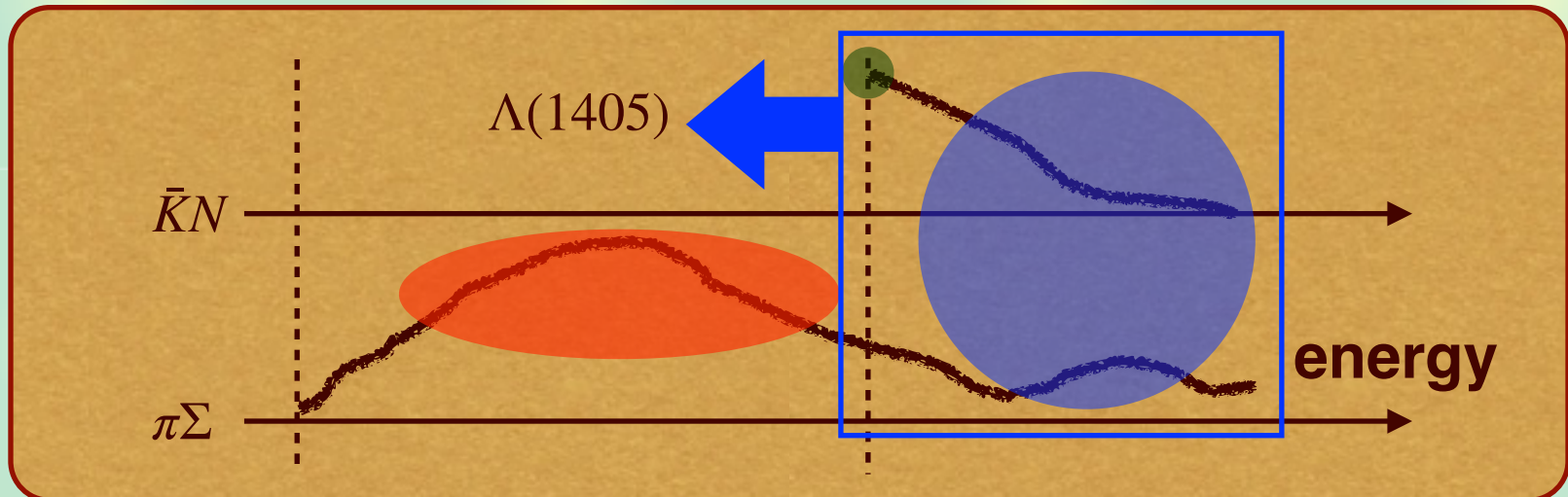
## Strategy for $\bar{K}N$ interaction

Above the  $\bar{K}N$  threshold : direct constraints

- $K^-p$  total cross sections (old data)
- $\bar{K}N$  threshold branching ratios (old data)
- $K^-p$  scattering length (new data : SIDDHARTA)

Below the  $\bar{K}N$  threshold: indirect constraints

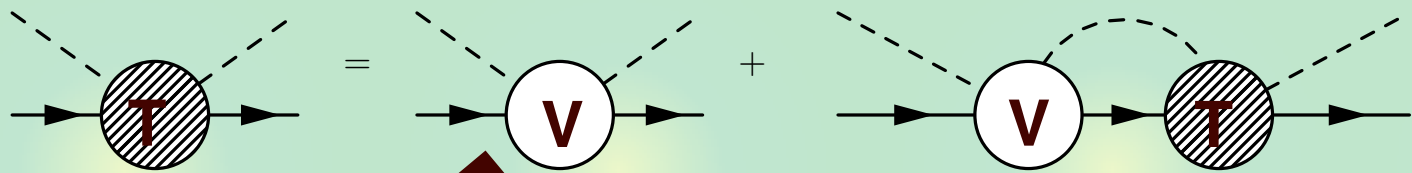
- $\pi\Sigma$  mass spectra (new data : LEPS, CLAS, HADES, ...)



# Construction of the realistic amplitude

Chiral SU(3) coupled-channels ( $\bar{K}N, \pi\Sigma, \pi\Lambda, \eta\Lambda, \eta\Sigma, K\Xi$ ) approach

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881 98 (2012)



Chiral perturbation theory

<p><b>1) TW term</b></p> <p><math>\mathcal{O}(p)</math></p> <p><b>6 cutoffs</b></p> <p><b>TW model</b></p>	<p><b>2) Born terms</b></p> <p><math>\mathcal{O}(p)</math></p> <p><b>TWB model</b></p>	<p><b>3) NLO terms</b></p> <p><math>\mathcal{O}(p^2)</math></p> <p><b>7 LECs</b></p> <p><b>NLO model</b></p>
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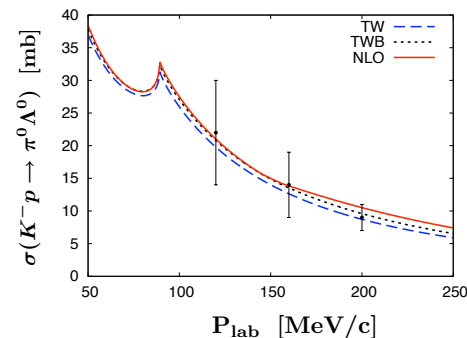
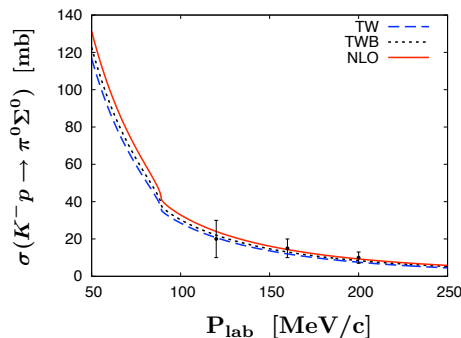
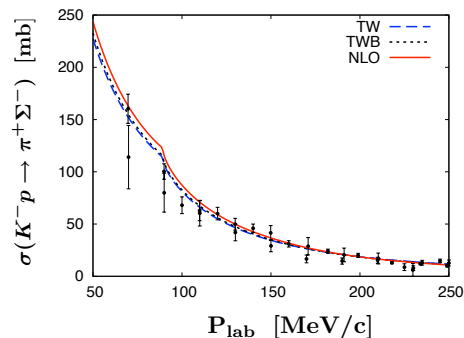
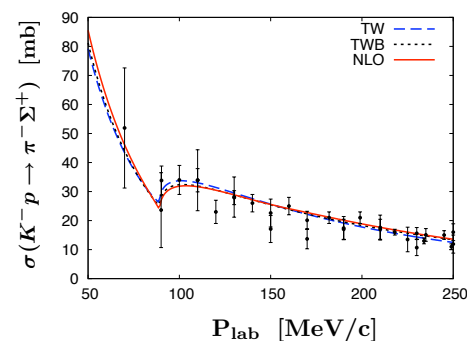
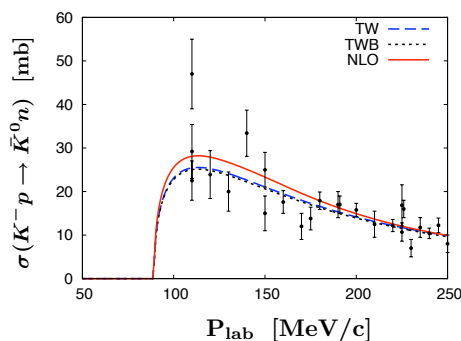
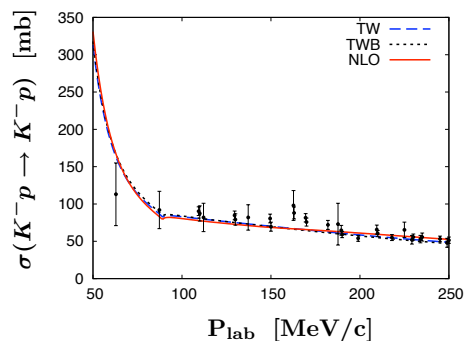
## Best-fit results

K at rest

	TW	TWB	NLO	Experiment
$\Delta E$ [eV]	373	377	306	$283 \pm 36 \pm 6$ [10]
$\Gamma$ [eV]	495	514	591	$541 \pm 89 \pm 22$ [10]
$\gamma$	2.36	2.36	2.37	$2.36 \pm 0.04$ [11]
$R_n$	0.20	0.19	0.19	$0.189 \pm 0.015$ [11]
$R_c$	0.66	0.66	0.66	$0.664 \pm 0.011$ [11]
$\chi^2/\text{d.o.f}$	1.12	1.15	0.96	

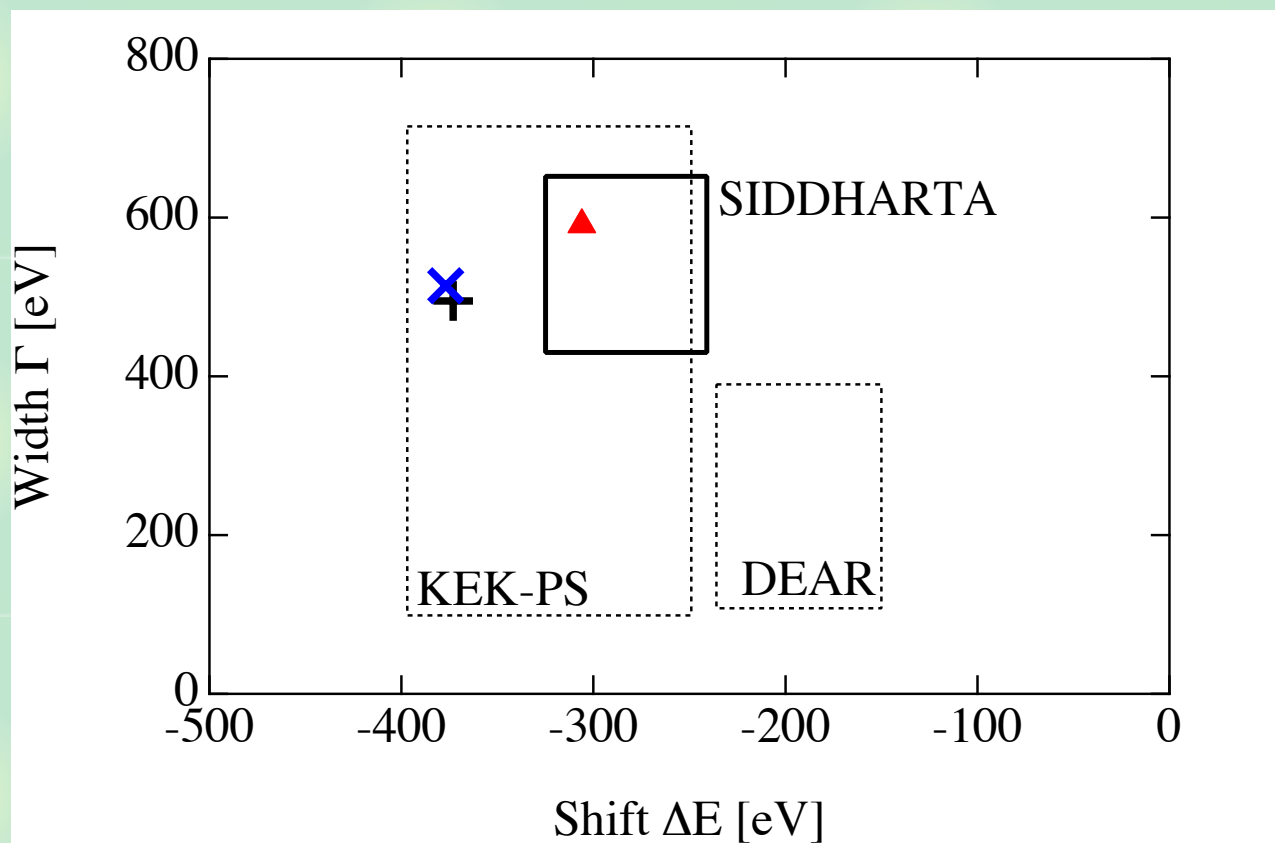
} SIDDHARTA

} Branching ratios

K<sup>-</sup>p cross sectionsAccurate description of all existing data ( $\chi^2/\text{d.o.f} \sim 1$ )

# Comparison with SIDDHARTA

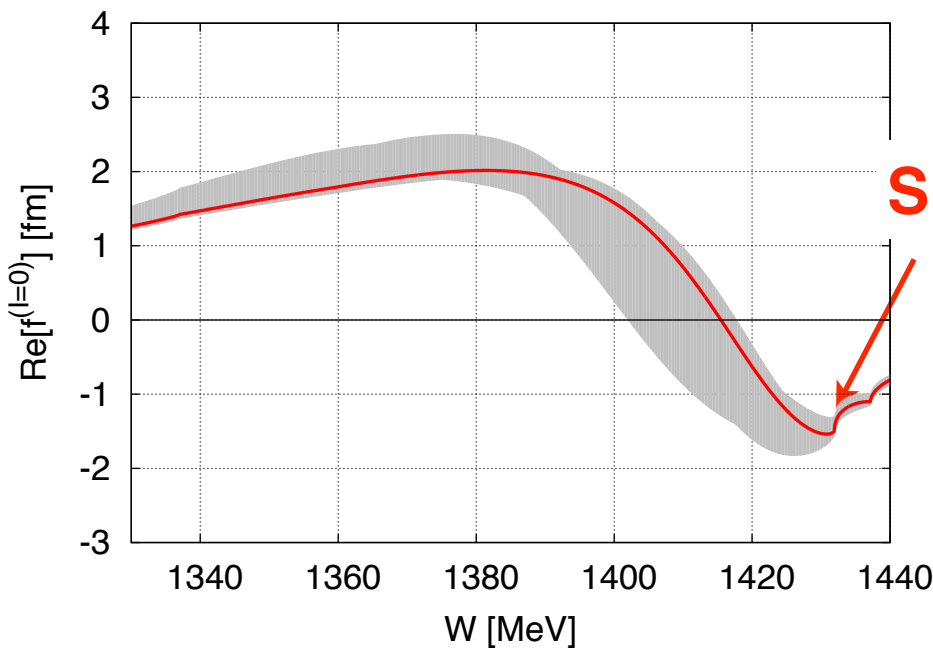
	<b>TW</b>	<b>TWB</b>	<b>NLO</b>
$\chi^2/\text{d.o.f.}$	<b>1.12</b>	<b>1.15</b>	<b>0.957</b>



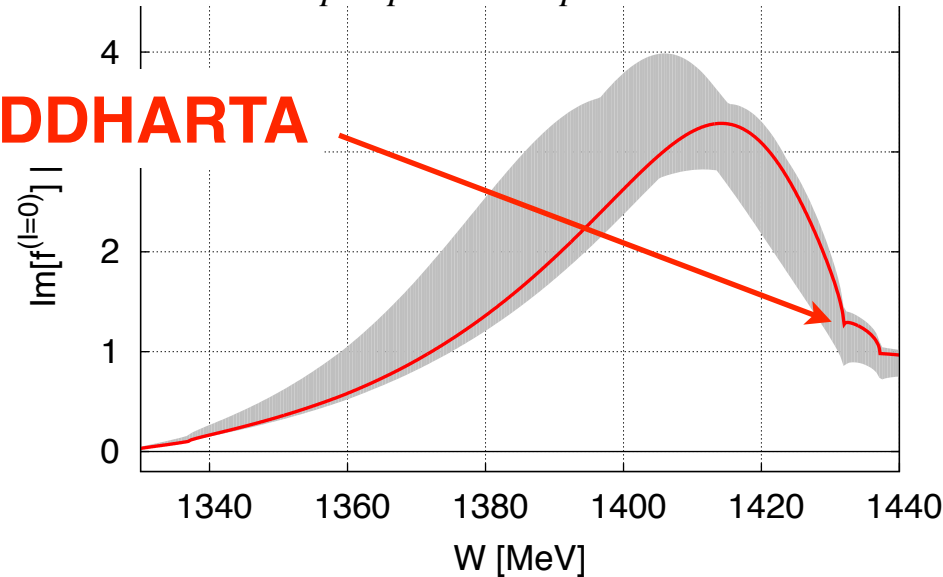
**TW** and **TWB** are reasonable, while best-fit requires **NLO**.

# Subthreshold extrapolation

## Uncertainty of $\bar{K}N \rightarrow \bar{K}N(I=0)$ amplitude below threshold



$$f^{(I=0)} = (f_{K^-pK^-p} + 2f_{K^-p\bar{K}^0n} + f_{\bar{K}^0n\bar{K}^0n})/2$$

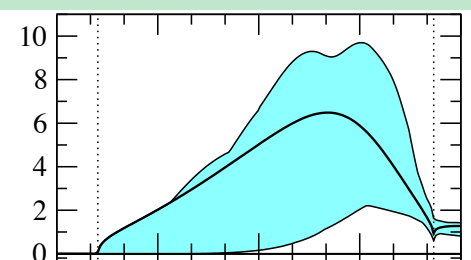
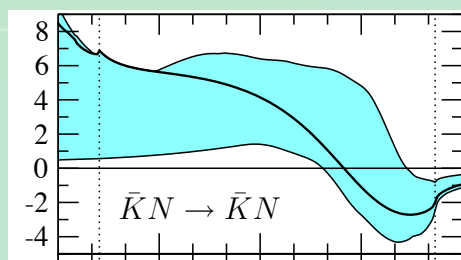


**SIDDHARTA**

Y. Kamiya, K. Miyahara, S. Ohnishi, Y. Ikeda, T. Hyodo, E. Oset, W. Weise, NPA 954, 41 (2016)

- c.f. without SIDDHARTA

R. Nissler, Doctoral Thesis (2007)



**SIDDHARTA is essential for subthreshold extrapolation.**



# Extrapolation to complex energy: two poles

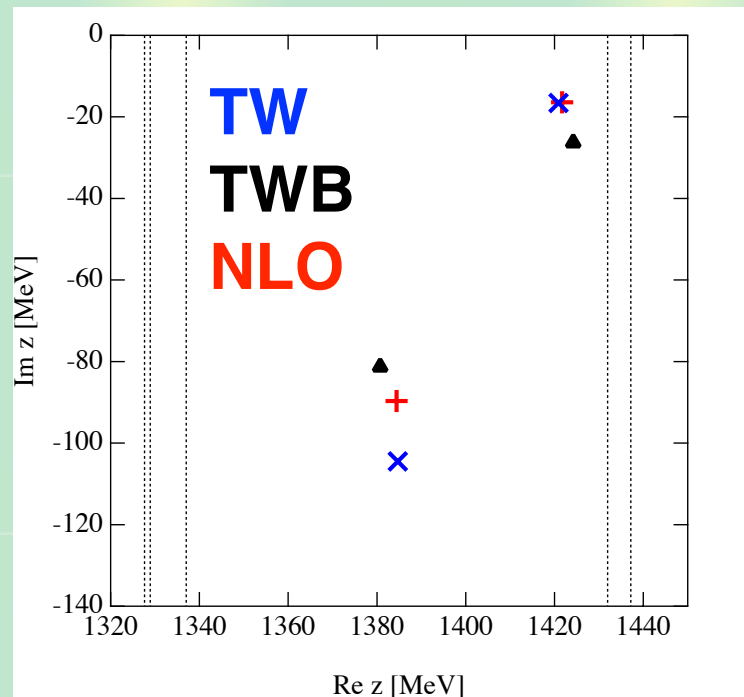
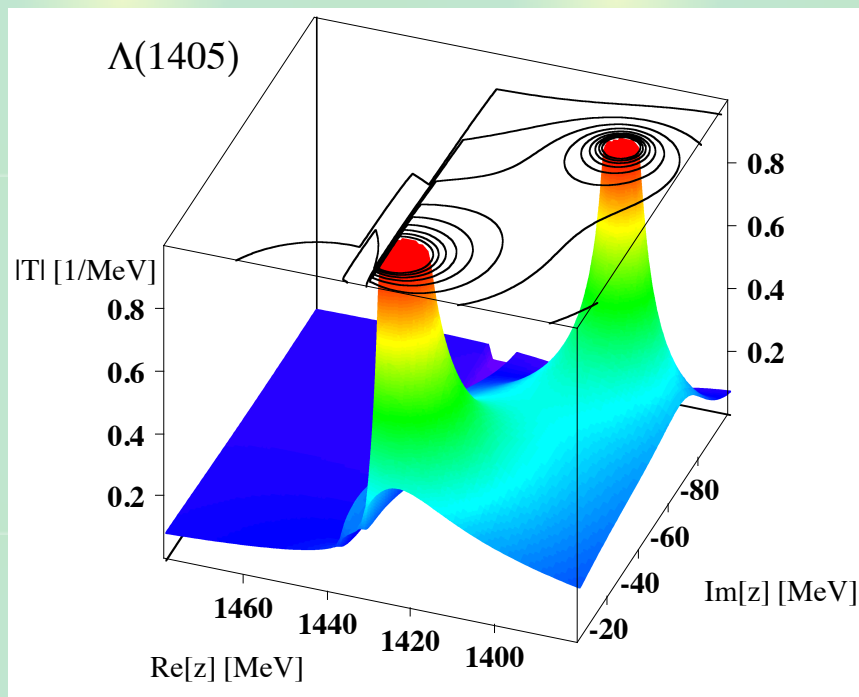
## Two poles : superposition of two eigenstates

J.A. Oller, U.G. Meißner, PLB 500, 263 (2001);

D. Jido, J.A. Oller, E. Oset, A. Ramos, U.G. Meißner, NPA 723, 205 (2003);

U.G. Meißner, Symmetry 12, 981 (2020); M. Mai, arXiv: 2010.00056 [nucl-th];

T. Hyodo, M. Niyama, arXiv: 2010.07592 [hep-ph], to appear in PPNP



T. Hyodo, D. Jido, PPNP 67, 55 (2012)

**NLO analysis confirms the two-pole structure.**

# PDG has changed

## 2020 update of PDG

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012); ▲

Z.H. Guo, J.A. Oller, PRC87, 035202 (2013); ✕

M. Mai, U.G. Meißner, EPJA51, 30 (2015) ■ ○

### - Particle Listing section:

Citation: P.A. Zyla *et al.* (Particle Data Group), Prog. Theor. Exp. Phys. **2020**, 083C01 (2020)

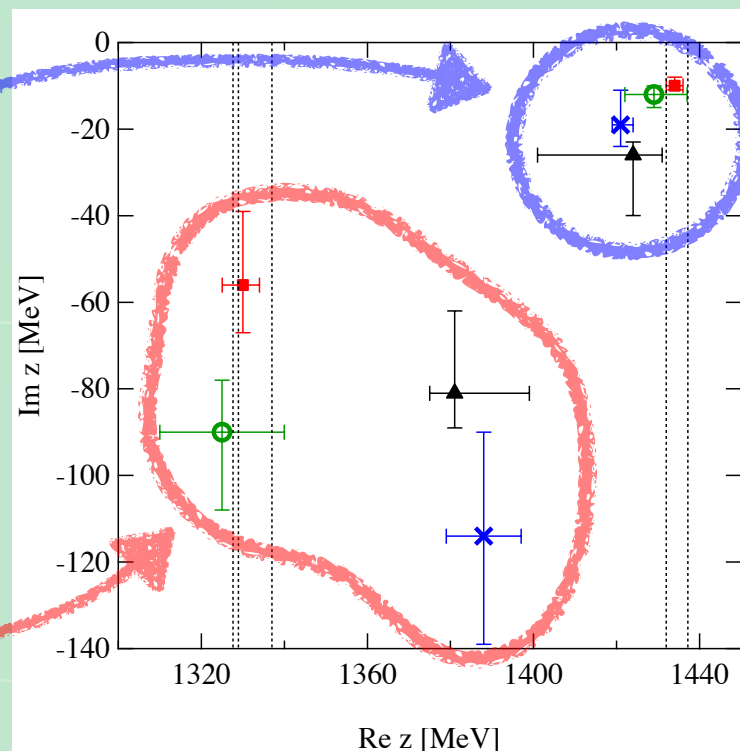
$\Lambda(1405) 1/2^-$

$I(J^P) = 0(\frac{1}{2}^-)$  Status: \*\*\*\*

Citation: P.A. Zyla *et al.* (Particle Data Group), Prog. Theor. Exp. Phys. **2020**, 083C01 (2020)

$\Lambda(1380) 1/2^-$

$J^P = \frac{1}{2}^-$  Status: \*\*  
**new!**



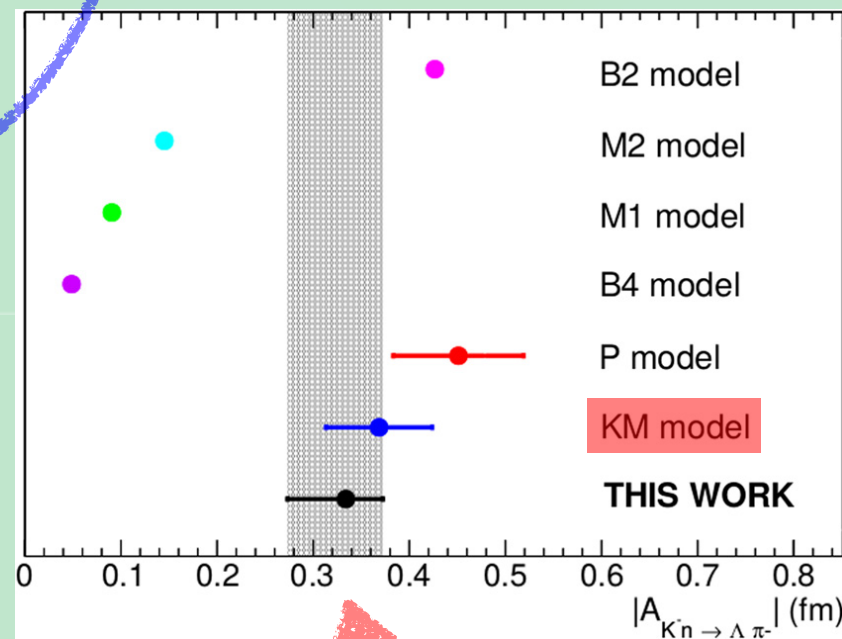
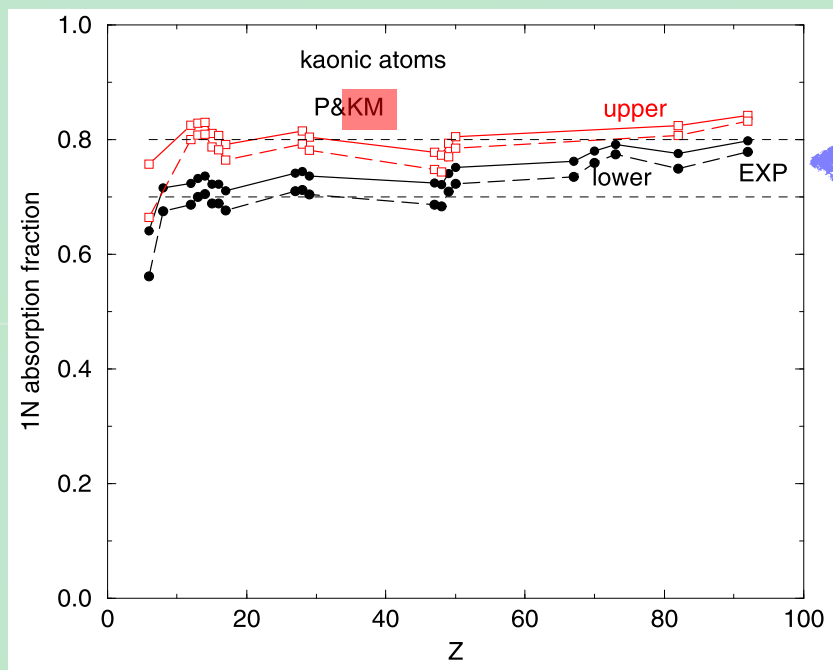
T. Hyodo, M. Niyama, arXiv: 2010.07592 [hep-ph], to appear in PPNP

- “ $\Lambda(1405)$ ” is no longer at 1405 MeV but  $\sim 1420$  MeV.
- Lower pole: two-star resonance  $\Lambda(1380)$

# Further check of amplitude

## Single-nucleon absorption on kaonic atoms

E. Friedman, A. Gal, NPA959, 66 (2017)



$|f_{K^- n \to \pi^- \Lambda}|$  from  $K^-$  absorption on  $^4\text{He}$  at DAΦNE

K. Piscicchia, *et al.*, PLB782, 339 (2018)

Our amplitude (**KM model**) is compatible with these analyses

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## $\bar{K}N$ potentials and their applications

K. Miyahara, T. Hyodo, PRC 93, 015201 (2016);  
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### - Kaonic nuclei

S. Ohnishi, W. Horiuchi, T. Hoshino, K. Miyahara, T. Hyodo, PRC95, 065202 (2017)

### - Kaonic deuterium

T. Hoshino, S. Ohnishi, W. Horiuchi, T. Hyodo, W. Weise, PRC96, 045204 (2017)

### - $K^-p$ correlation function

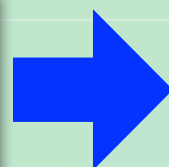
Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise. PRL124, 132501 (2020)

## Summary

# Construction of $\bar{K}N$ potentials

**Local  $\bar{K}N$  potential** is useful for various applications

meson-baryon amplitude  
(chiral SU(3) EFT)



Coupled-channel real  
 $\bar{K}N$   $\pi\Sigma$   $\pi\Lambda$  potential

K. Miyahara, T. Hyodo, W. Weise, PRC 98, 025201 (2018)



Single-channel complex  
 $\bar{K}N$  potential



$K^-p$  correlation function

K. Miyahara, T. Hyodo, PRC 93, 015201 (2016)



Kaonic nuclei

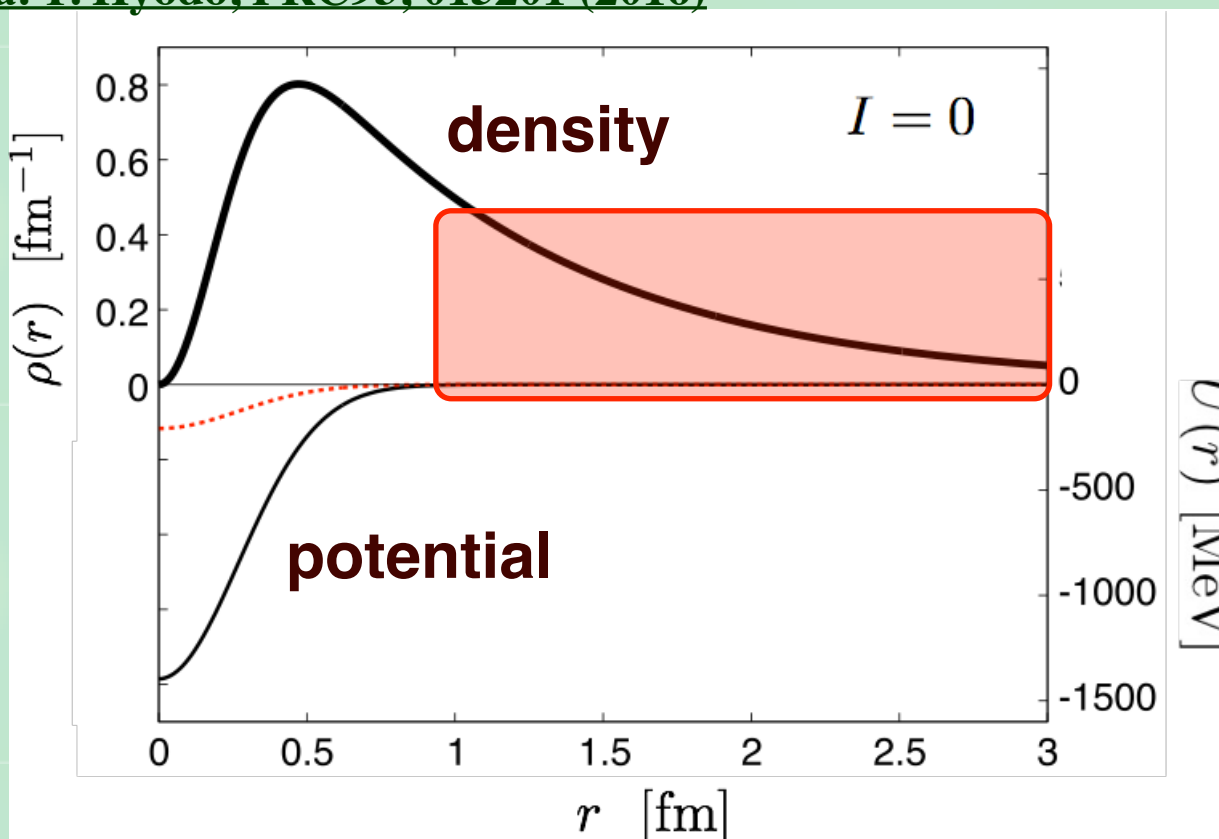


Kaonic deuterium

# Spatial structure of $\Lambda(1405)$

## $\bar{K}N$ wave function at $\Lambda(1405)$ pole

K. Miyahara, T. Hyodo, PRC93, 015201 (2016)



- substantial distribution at  $r > 1$  fm
- root mean squared radius  $\sqrt{\langle r^2 \rangle} = 1.44$  fm

The **size** of  $\Lambda(1405)$  is much **larger** than ordinary hadrons.



# Kaonic nuclei

## Rigorous few-body approach to $\bar{K}$ nuclear systems

S. Ohnishi, W. Horiuchi, T. Hoshino, K. Miyahara, T. Hyodo, PRC95, 065202 (2017).

- Stochastic variational method with correlated gaussians

$$\hat{V} = \hat{V}^{\bar{K}N}(\text{Kyoto } \bar{K}N) + \hat{V}^{NN}(\text{AV4}') \quad (\text{single channel})$$

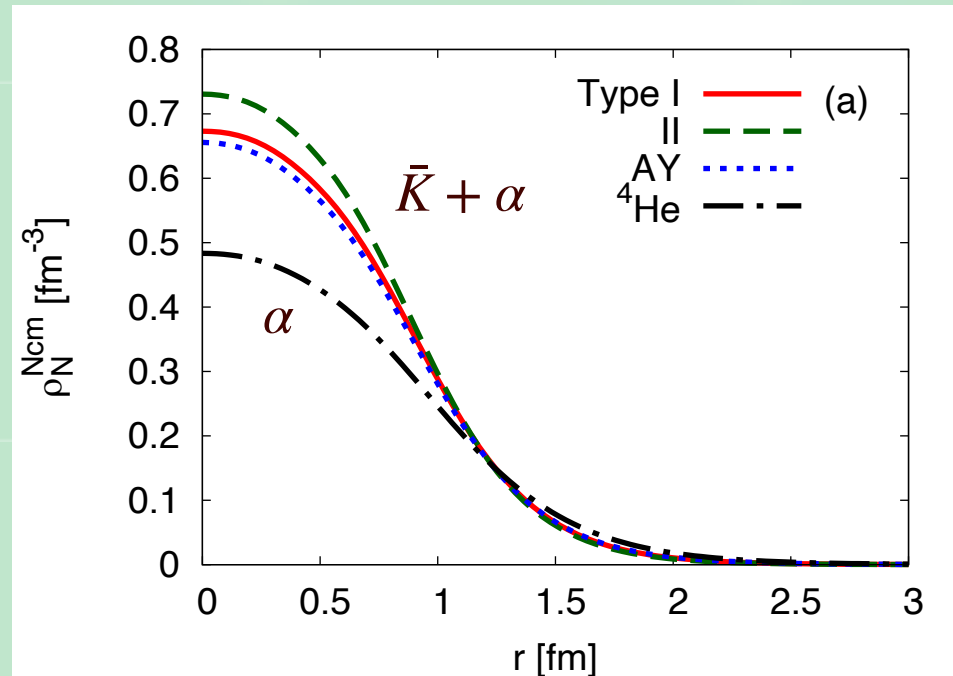
## Results for kaonic nuclei with $A = 2, 3, 4, 6$

	$\bar{K}NN$	$\bar{K}NNN$	$\bar{K}NNNN$	$\bar{K}NNNNN$
B [MeV]	25-28	45-50	68-76	70-81
$\Gamma_{\pi YN}$ [MeV]	31-59	26-70	28-74	24-76

- **quasi-bound** state below the lowest threshold
- decay width (**without multi- $N$  absorption**)  $\sim$  binding energy

# High density?

## Nucleon density distribution in four-nucleon system



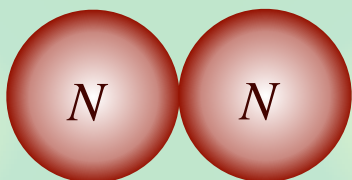
- central density increases (not substantially  $\leftarrow NN$  core)
- $B = 68-76$  MeV (Kyoto  $\bar{K}N$ )
- $B = 85-87$  MeV (AY)

Central density is **not always** proportional to  $B \leftarrow$  tail of w.f.

# Interplay between $NN$ and $\bar{K}N$ correlations 1

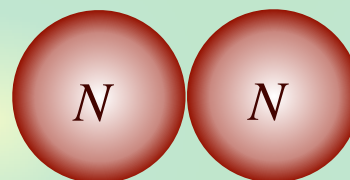
## Two-nucleon system

$${}^1S_0(I_{NN} = 1)$$

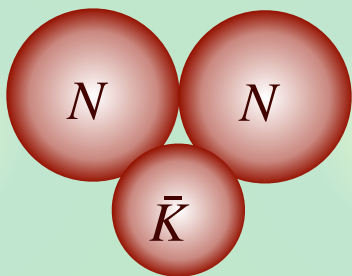


unbound

$${}^3S_1(I_{NN} = 0)$$

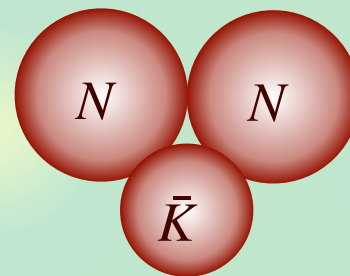


bound ( $d$ )



(quasi-)bound

$$\frac{\bar{K}N(I = 0)}{\bar{K}N(I = 1)} = 3$$



$\Lambda(1405)$

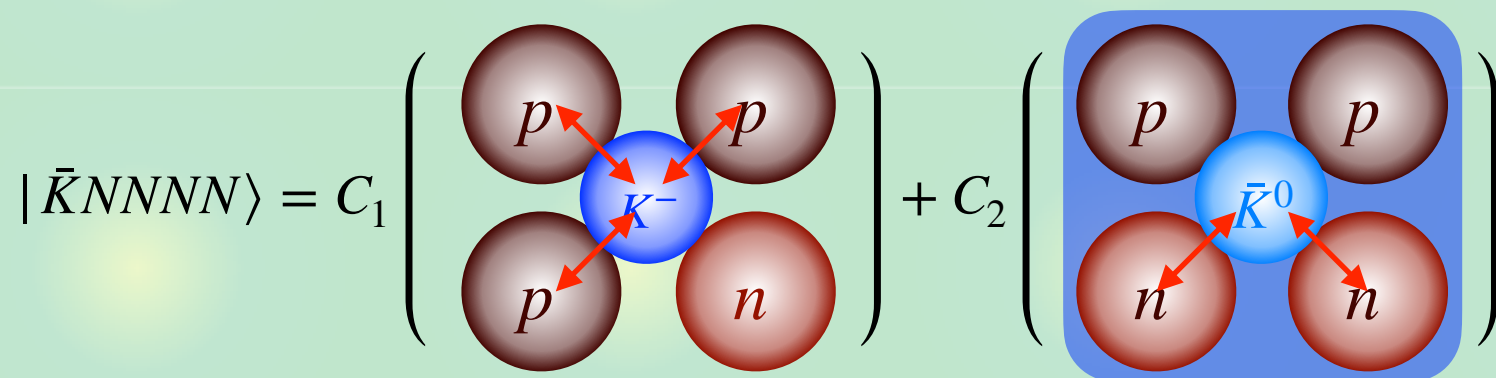
unbound

$$\frac{\bar{K}N(I = 0)}{\bar{K}N(I = 1)} = \frac{1}{3}$$

$NN$  correlation  $<$   $\bar{K}N$  correlation (also in  $A = 6$ )

# Interplay between $NN$ and $\bar{K}N$ correlations 2

Four-nucleon system with  $J^P = 0^-, I = 1/2, I_3 = +1/2$



## - $\bar{K}N$ correlation

$I = 0$  pair in  $K^-p$  (3 pairs) or  $\bar{K}^0n$  (2 pairs) :  $|C_1|^2 > |C_2|^2$

## - $NN$ correlation

$ppnn$  forms  $\alpha$  :  $|C_1|^2 < |C_2|^2$

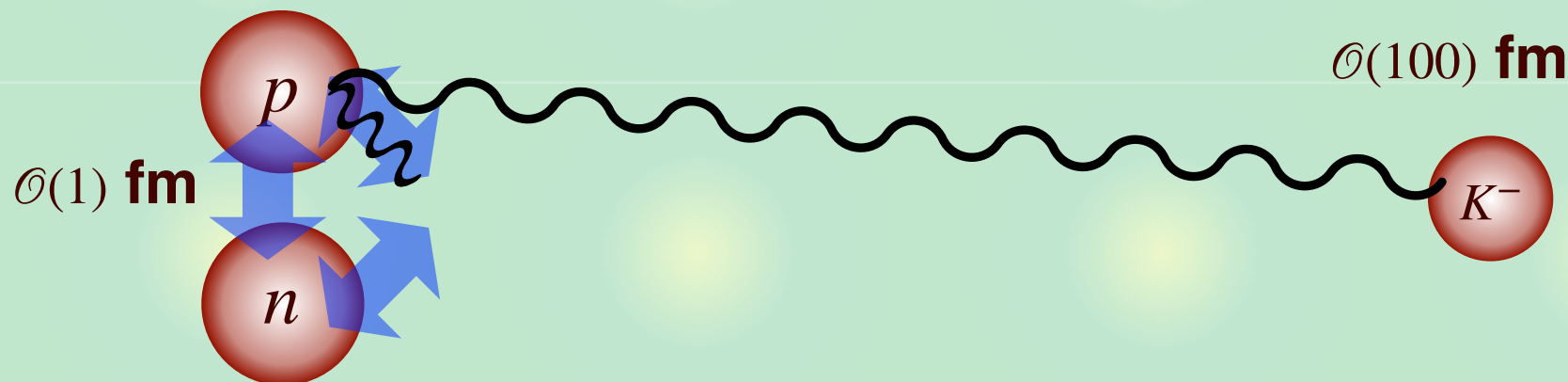
## - Numerical result

$$|C_1|^2 = 0.08, \quad |C_2|^2 = 0.92$$

$NN$  correlation  $>$   $\bar{K}N$  correlation

## Kaonic deuterium: background

$K^-pn$  system with **strong** + Coulomb interaction



- Experiments are planned at J-PARC E57, SIDDHARTA-2

Theoretical requirements:

- **Rigorous** three-body treatment of strong + Coulomb
- Inclusion of SIDDHARTA constraint (**realistic**  $\bar{K}N$ )
- c.f. advanced Faddeev calculations

P. Doleschall, J. Revai, N.V. Shevchenko, PLB 744, 105 (2015);

J. Revai, PRC 94, 054001 (2016)

# Check of kaonic hydrogen

Kaonic hydrogen ( $K^-p$ ) in the present setup?

- Deser-type formula is based on (systematic) expansion.
- $\bar{K}N$  potential is formulated with isospin symmetry.

Two-body calculation with physical masses

$$\begin{pmatrix} \hat{T} + \hat{V}^{\bar{K}N} + \hat{V}^{\text{EM}} & \hat{V}^{\bar{K}N} \\ \hat{V}^{\bar{K}N} & \hat{T} + \hat{V}^{\bar{K}N} + \Delta m \end{pmatrix} \begin{pmatrix} |K^-p\rangle \\ |\bar{K}^0n\rangle \end{pmatrix} = E \begin{pmatrix} |K^-p\rangle \\ |\bar{K}^0n\rangle \end{pmatrix}$$

**Result:**

- **consistent** with SIDDHARTA constraint
- Resummed Deser-type formula works reasonably for  $K^-p$ .

Mass	$E$ dependence	$\Delta E$ (eV)	$\Gamma$ (eV)
Physical	Self-consistent	283	607
Isospin	Self-consistent	163	574
Physical	$E_{\bar{K}N} = 0$	283	607
Expt. [31,32]		$283 \pm 36 \pm 6$	$541 \pm 89 \pm 22$

	$\Delta E$ (eV)	$\Gamma$ (eV)
Full Schrödinger equation	283	607
Improved Deser formula (18)	293	596
Resummed formula (19)	284	605



# Formulation

## Three-body calculation of $K^-d$ with physical masses

T. Hoshino, S. Ohnishi, W. Horiuchi, T. Hyodo, W. Weise, PRC96, 045204 (2017)

$$\begin{pmatrix} \hat{H}_{K^-pn} & \hat{V}_{12}^{\bar{K}N} + \hat{V}_{13}^{\bar{K}N} \\ \hat{V}_{12}^{\bar{K}N} + \hat{V}_{13}^{\bar{K}N} & \hat{H}_{\bar{K}^0nn} \end{pmatrix} \begin{pmatrix} |K^-pn\rangle \\ |\bar{K}^0nn\rangle \end{pmatrix} = E \begin{pmatrix} |K^-pn\rangle \\ |\bar{K}^0nn\rangle \end{pmatrix}$$

$$\hat{H}_{K^-pn} = \sum_{i=1}^3 \hat{T}_i - \hat{T}_{\text{cm}} + \hat{V}_{23}^{NN} + \sum_{i=2}^3 (\hat{V}_{1i}^{\bar{K}N} + \hat{V}_{1i}^{\text{EM}}) \text{Coulomb}$$

$$\hat{H}_{\bar{K}^0nn} = \sum_{i=1}^3 \hat{T}_i - \hat{T}_{\text{cm}} + \hat{V}_{23}^{NN} + \sum_{i=2}^3 \hat{V}_{1i}^{\bar{K}N} + \Delta m \text{ threshold difference}$$

- (single-channel) realistic  $\bar{K}N$  potential

K. Miyahara, T. Hyodo, PRC 93, 015201 (2016)

## Few-body technique

- a large number of correlated gaussian basis

Y. Suzuki, K. Varga, Lect. Notes Phys. M54, (1998)

# Kaonic deuterium: shift and width

## Results of the three-body calculation

T. Hoshino, S. Ohnishi, W. Horiuchi, T. Hyodo, W. Weise, PRC96, 045204 (2017)

- energy convergence
- ← a large number of basis

$N$	Re[ $E$ ] (MeV)
1677	-2.211689436
2194	-2.211722964
2377	-2.211732072
2511	-2.211735493
2621	-2.211737242
2721	-2.211737609
2806	-2.211737677
2879	-2.211737682

## Shift-width of the $1S$ state:

$$\Delta E - i\Gamma/2 = (670 - i508) \text{ eV}$$

keV      eV!

- No shift in  $2P$  state is shown by explicit calculation.
- Deser-type formula does **not** work accurately for  $K^-d$

c.f.) J. Revai, PRC 94, 054001 (2016)

	$\Delta E$ (eV)	$\Gamma$ (eV)
Full Schrödinger equation	670	1016
Improved Deser formula (18)	910	989
Resummed formula (19)	818	1188

## $I = 1$ dependence

### Study sensitivity to $I = 1$ interaction

- introduce parameter  $\beta$  to control the potential strength

$$\text{Re } \hat{V}^{\bar{K}N(I=1)} \rightarrow \beta \times \text{Re } \hat{V}^{\bar{K}N(I=1)}$$

### Vary $\beta$ within SIDDHARTA uncertainty of $K^-p$

- allowed region:  $-0.17 < \beta < 1.08$

(negative  $\beta$  may contradict with scattering data)

$\beta$	$K^-p$		$K^-d$	
	$\Delta E$	$\Gamma$	$\Delta E$	$\Gamma$
1.08	287	648	676	1020
1.00	283	607	670	1016
-0.17	310	430	506	980

- deviation of  $\Delta E$  of  $K^-d \sim 170$  eV

- Planned precision: 60 eV (30 eV) at J-PARC (SIDDHARTA-2)

Measurement of  $K^-d$  will provide **strong constraint** on  $I = 1$

# New data : $K^-p$ correlation function

## $K^-p$ total cross sections

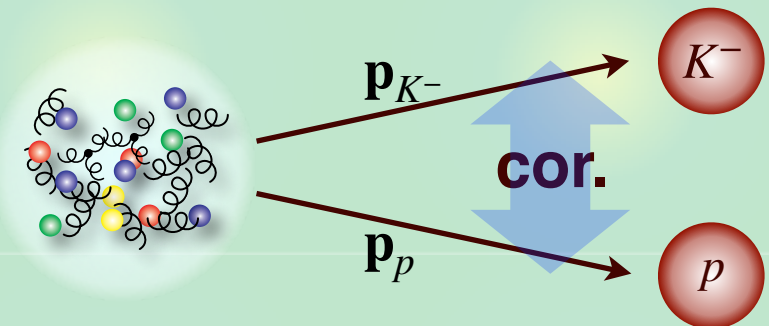
Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011)

- Old bubble chamber data

## $K^-p$ correlation function

S. Acharya *et al.* (ALICE), PRL 124, 092301 (2020)

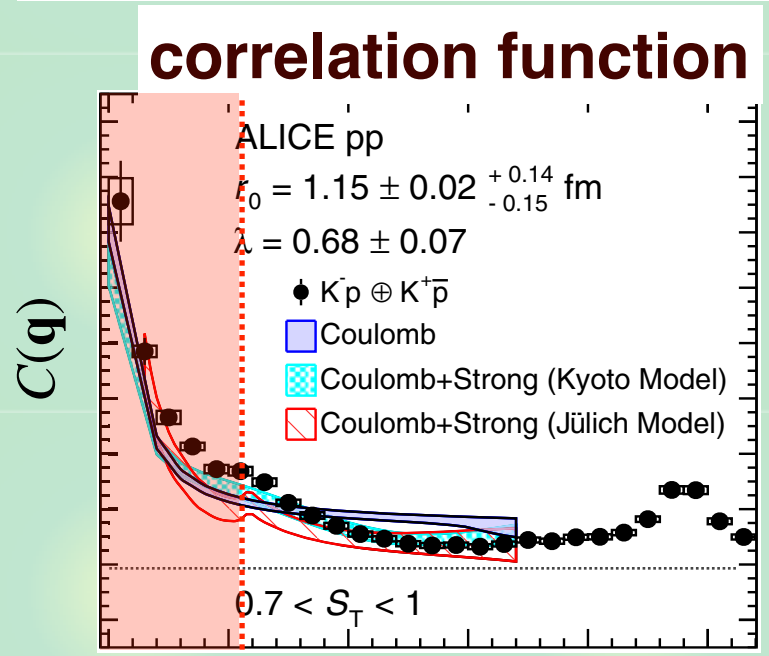
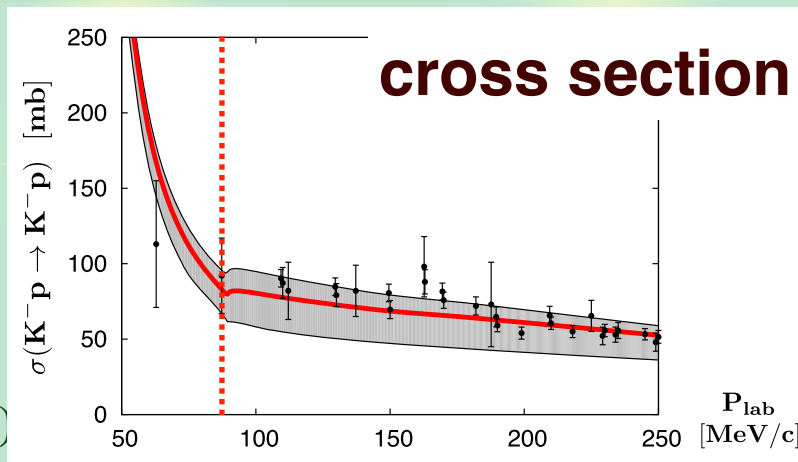
$$C(\mathbf{q}) = \frac{N_{K^-p}(\mathbf{p}_{K^-}, \mathbf{p}_p)}{N_{K^-}(\mathbf{p}_{K^-})N_p(\mathbf{p}_p)}$$



- Excellent **precision** ( $\bar{K}^0n$  cusp)

- Low-energy data **below**  $\bar{K}^0n$

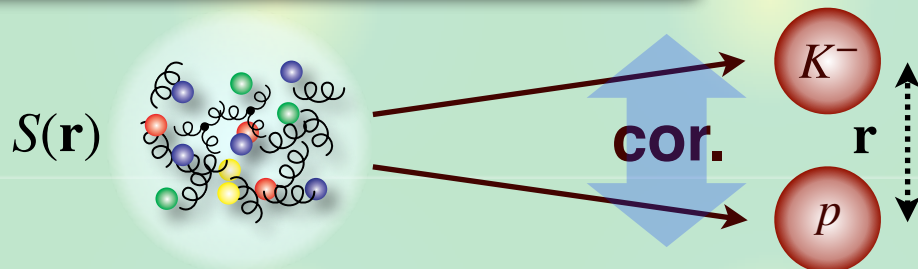
—> important constraint on  $\Lambda(1405)$  theories



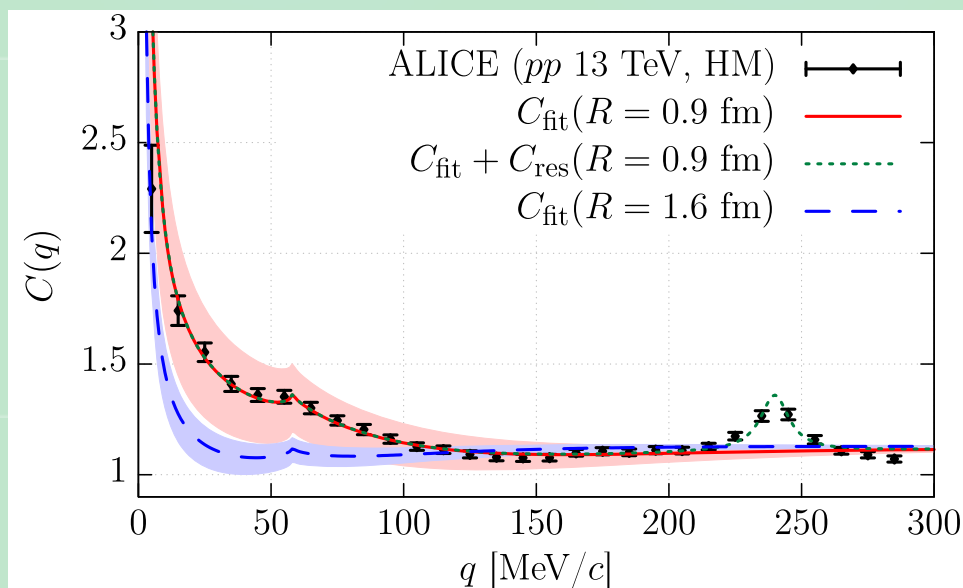
# Prediction from chiral SU(3) dynamics

Theoretical calculation of  $C(q)$

$$C(\mathbf{q}) \simeq \int d^3\mathbf{r} S(\mathbf{r}) |\Psi_{\mathbf{q}}^{(-)}(\mathbf{r})|^2$$



- wave function  $\Psi_{\mathbf{q}}^{(-)}(\mathbf{r})$  : coupled-channel  $\bar{K}N$ - $\pi\Sigma$ - $\pi\Lambda$  potential
- source function  $S(\mathbf{r})$  : estimated by  $K^+p$  data



Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise. PRL124, 132501 (2020)

**Correlation function is well reproduced.**

# Summary



**Pole structure of the  $\Lambda(1405)$  region is now well constrained by the experimental data.**

“ $\Lambda(1405)$ ”  $\rightarrow$   $\Lambda(1405)$  **and**  $\Lambda(1380)$

Y. Ikeda, T. Hyodo, W. Weise, PLB 706, 63 (2011); NPA 881, 98 (2012);

T. Hyodo, M. Niyama, arXiv: 2010.07592 [hep-ph], to appear in PPNP



**$\bar{K}N$  potentials are useful for various applications, Kaonic nuclei, Kaonic deuterium, and  $K^-p$  correlation function**

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Y. Kamiya, T. Hyodo, K. Morita, A. Ohnishi, W. Weise. PRL124, 132501 (2020)