## Electromagnetic Probes as Multi-Messengers

#### Sebastian Scheid Goethe University Frankfurt EMMI Symposium, 1st August 2022



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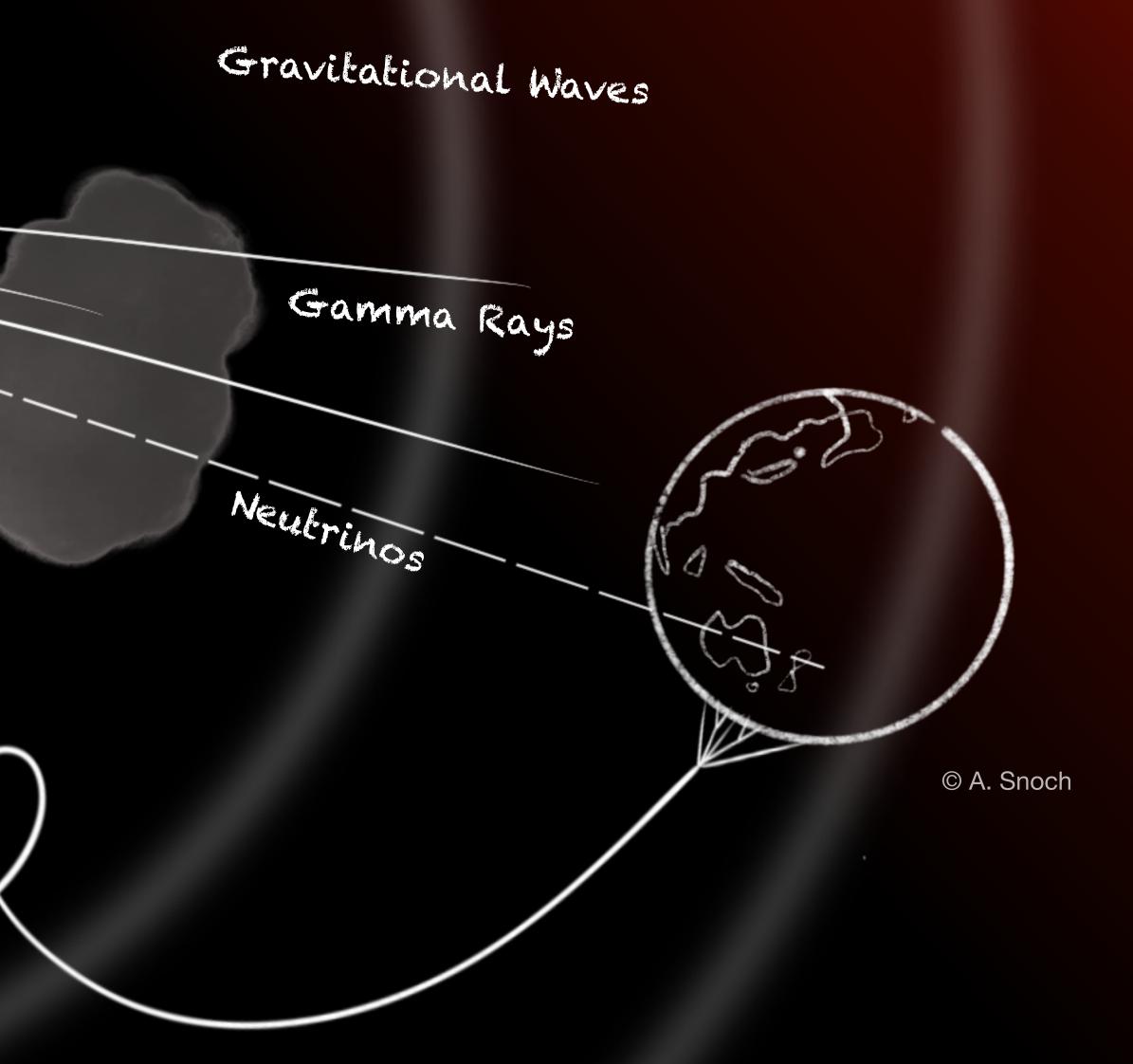


ECD **FSP**ALICE Erforschung von Universum und Materie



### Multi Messengers

Cosmic Rays



## Electromagnetic Probes as Multi-Messengers of Heavy Ion Collisions

#### Sebastian Scheid Goethe University Frankfurt EMMI Symposium, 1st August 2022





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FSP ALICE ECD Erforschung von Universum und Materie

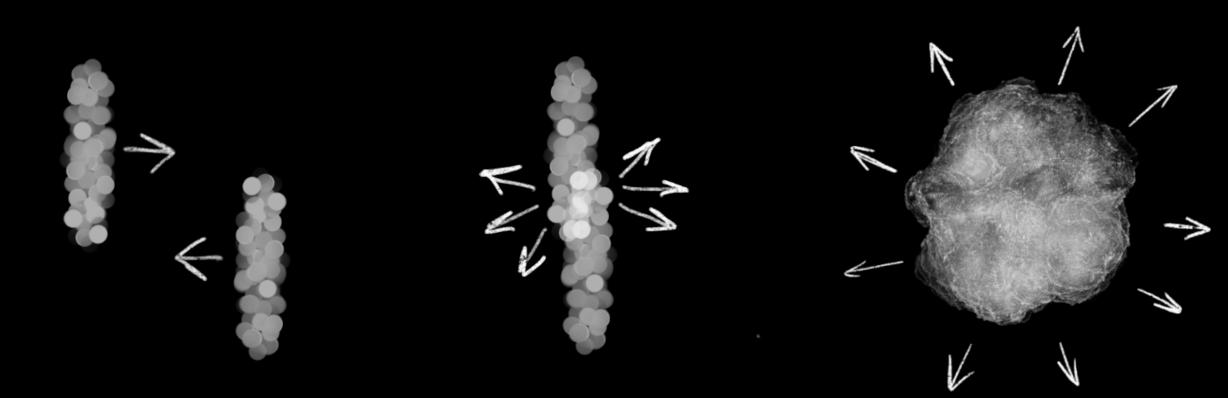


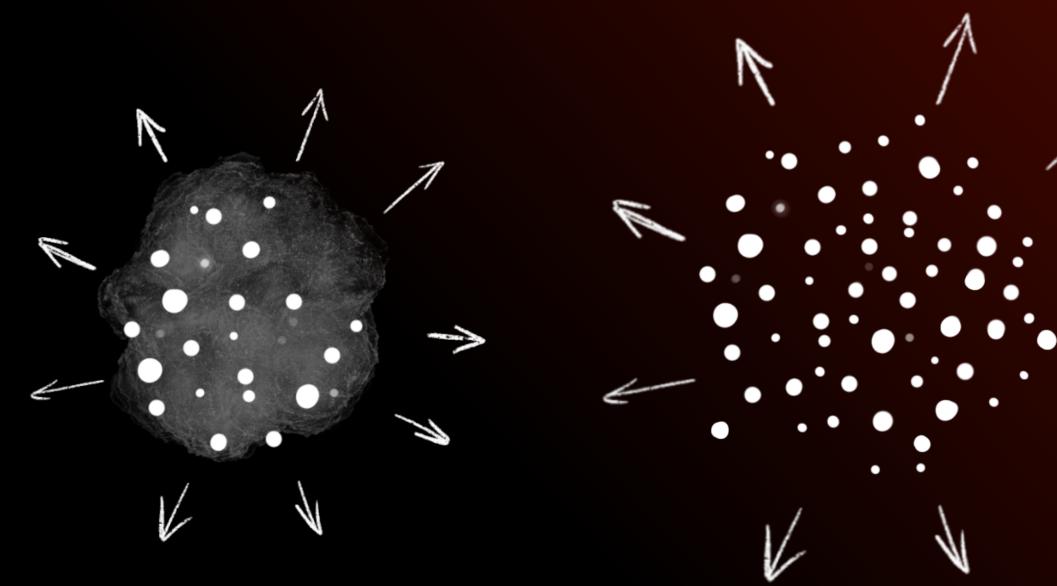
#### Act I Introduction

## Act II • A selection of (recent) results Act III • Perspectives for the future



## The little big bang



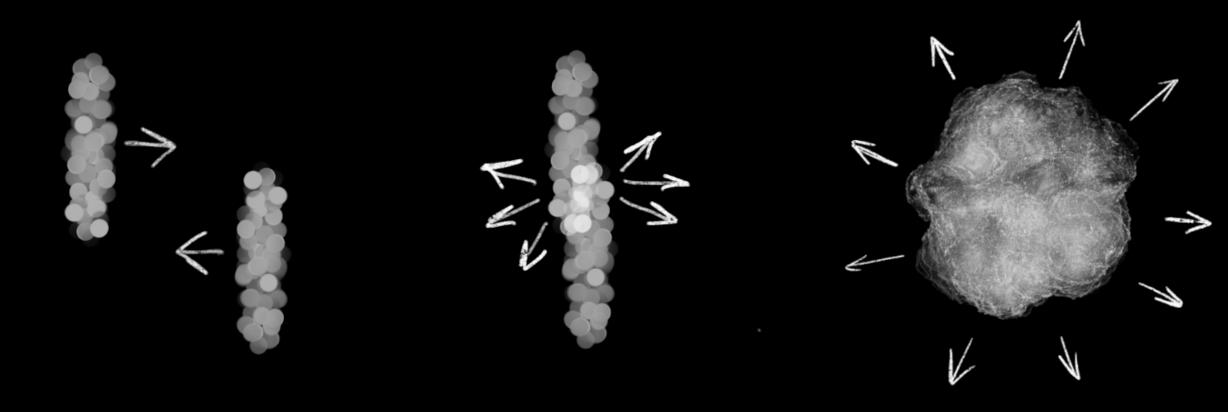


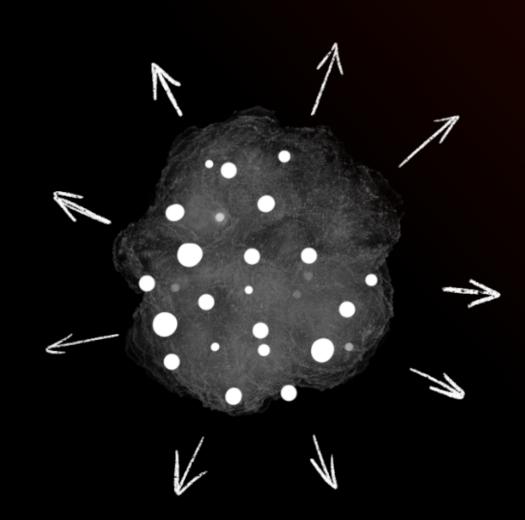


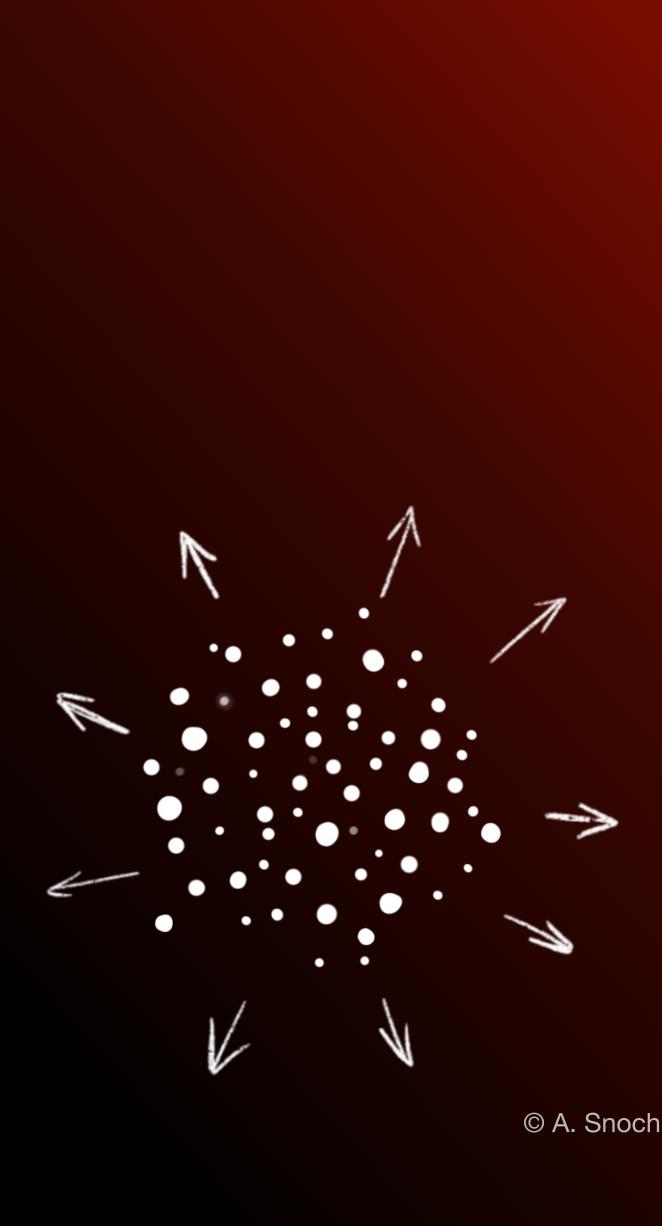


## Why EM probes

We want to measure the properties of strongly interacting matter
Why would we use probes that do not interact strongly?



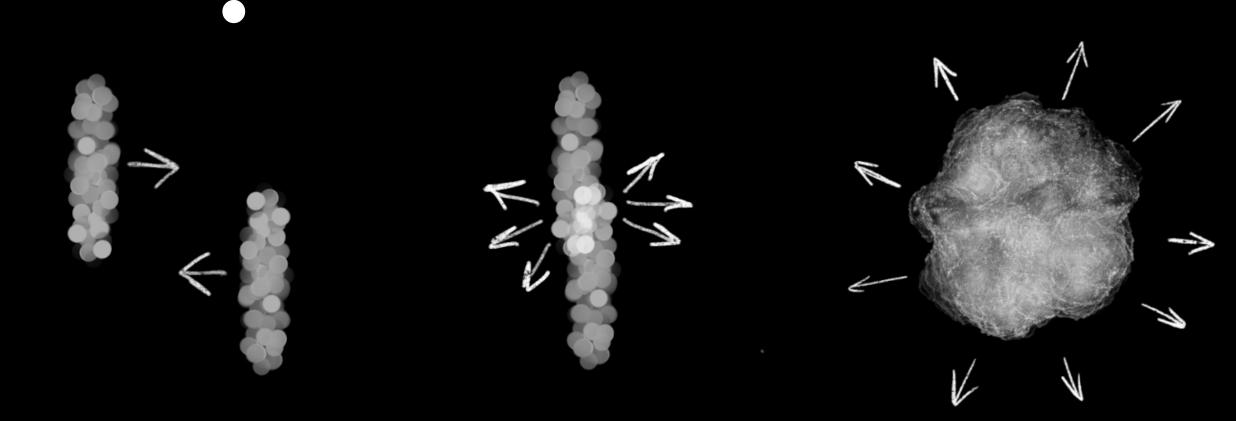






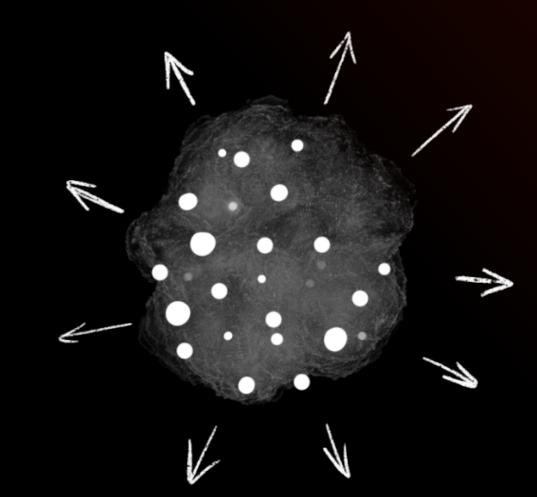
## Why EM probes

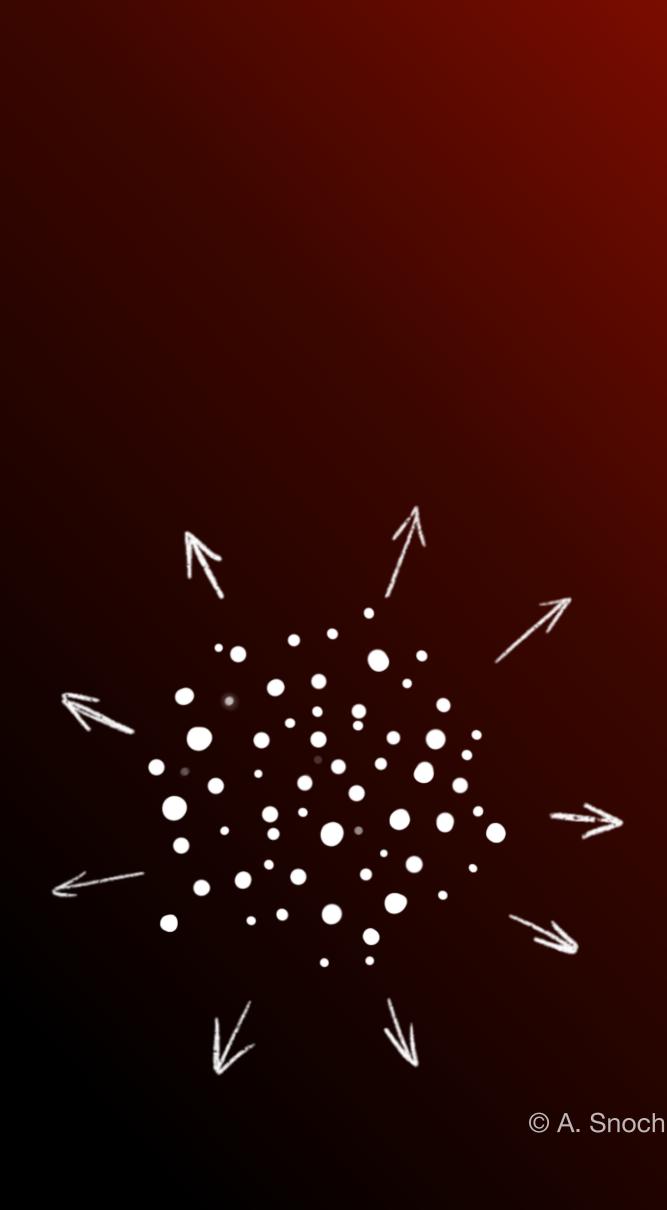
We want to measure the properties of strongly interacting matterWhy would we use probes that do not interact strongly?



#### Hadronic probes have limited sensitivity to early stages of the collision

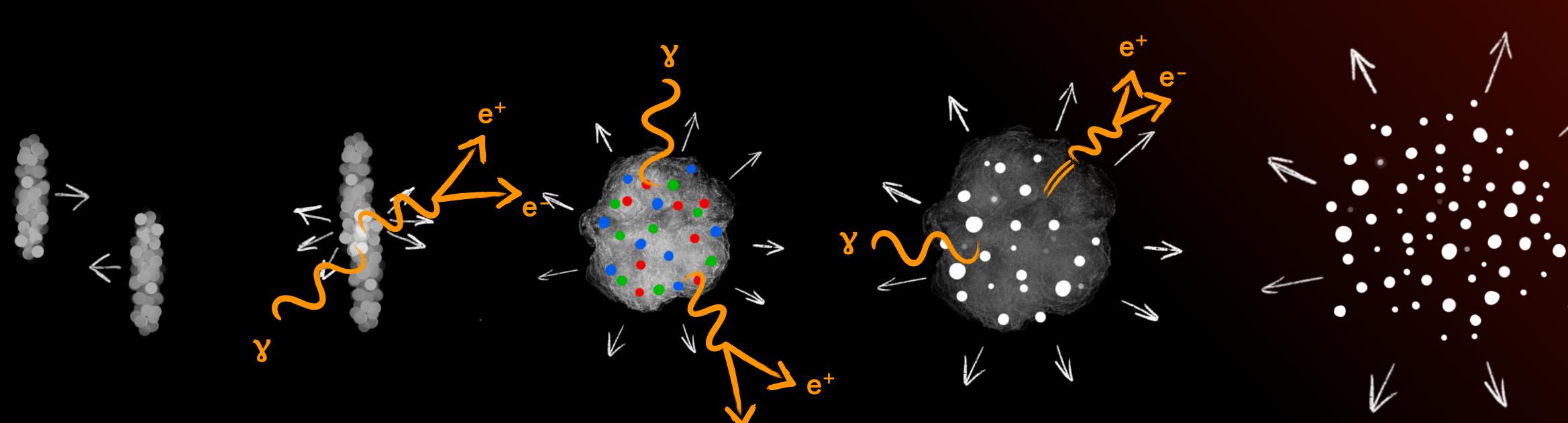
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## Why EM probes



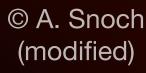
#### Pre-equilibrium

Quark-gluon plasma

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#### ma Hadron Gas





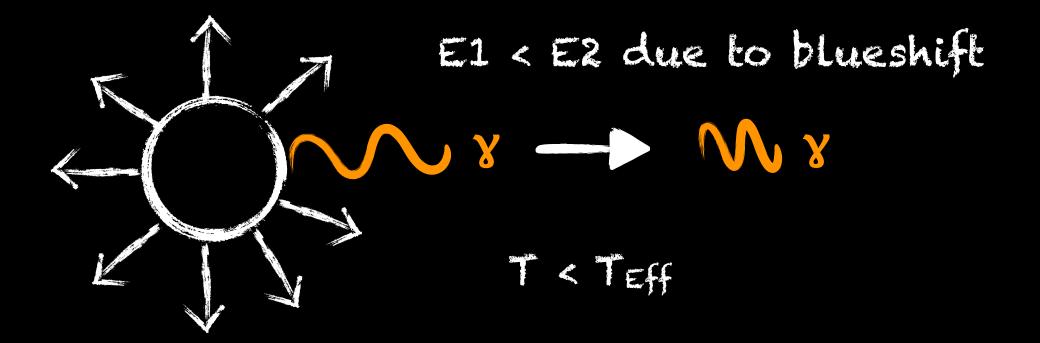


## The Sources - Real direct Photons

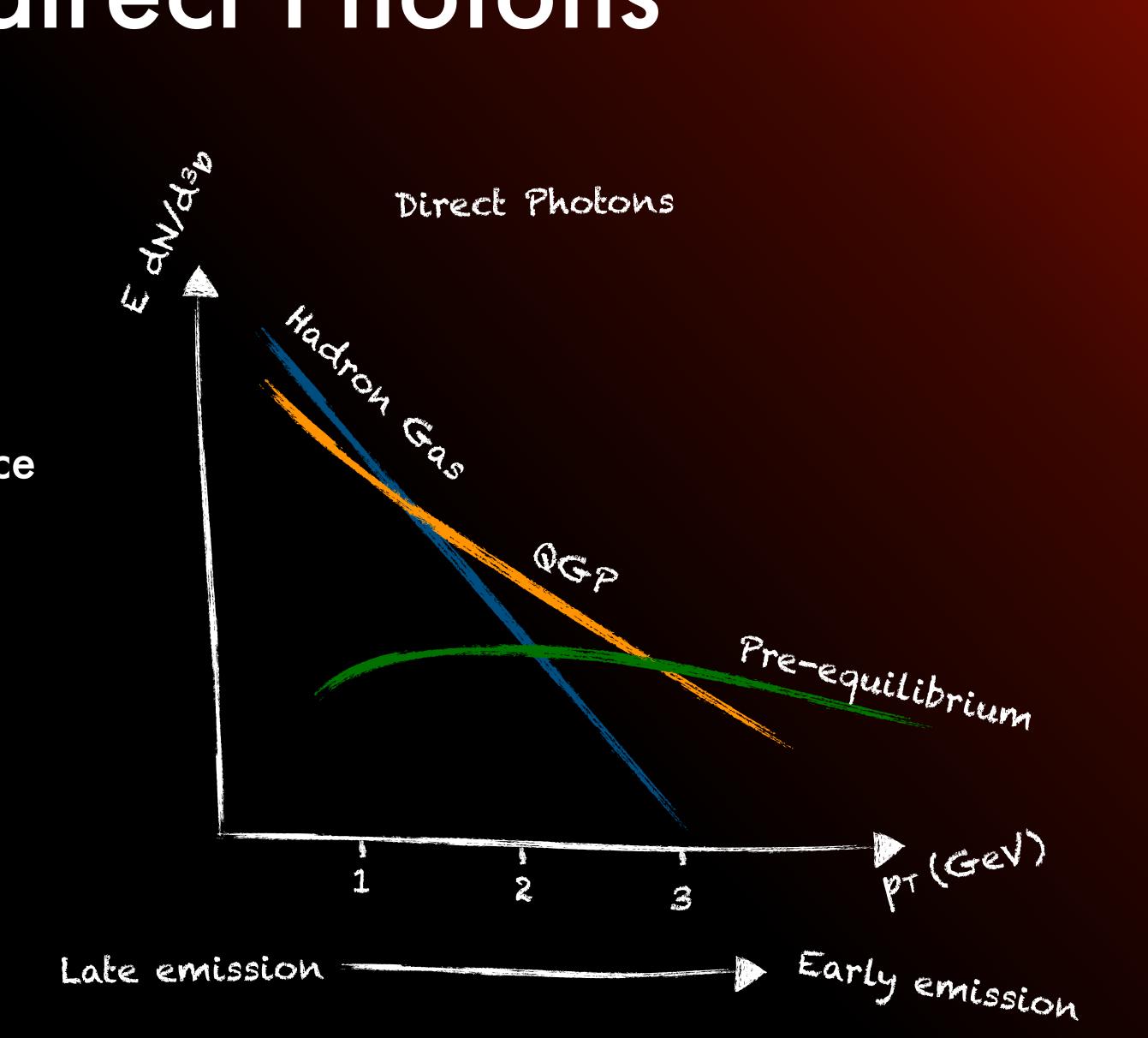
The different sources of EM radiation

• Thermal (QGP, Hadron gas)

Hard scattering and pre-equilibrium
 Inverse slope of thermal spectrum ∝ T of source
 Sources populate different p<sub>T</sub> ranges



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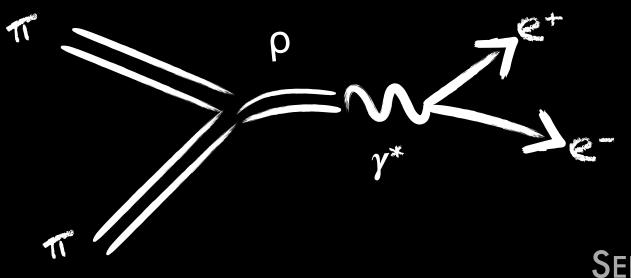
## The Sources - Dielectrons

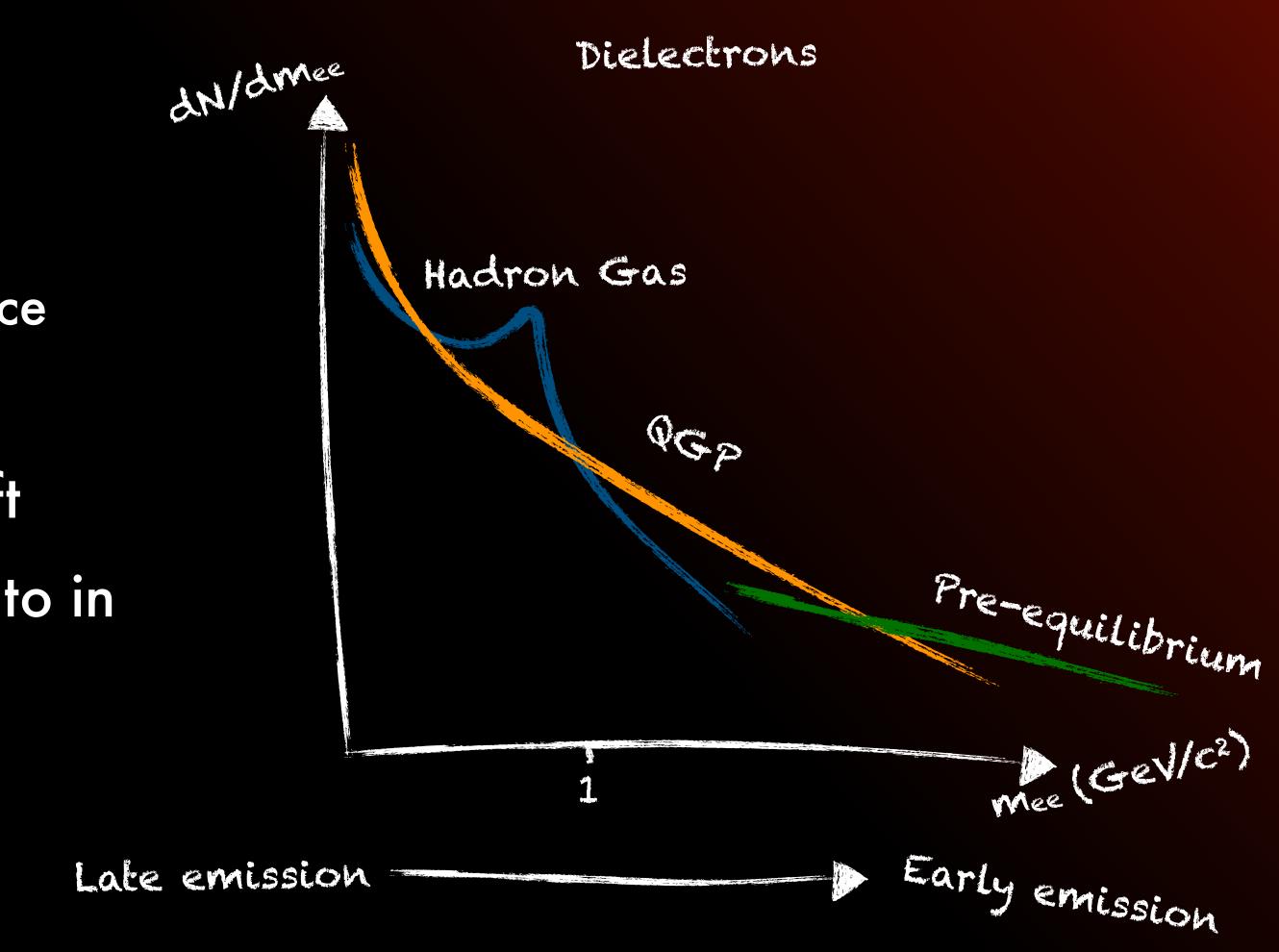
The different sources of EM radiation

Thermal (QGP, Hadron gas)

 Hard scattering and pre-equilibrium Inverse slope of thermal spectrum  $\propto$  T of source Sources populate different mass ranges

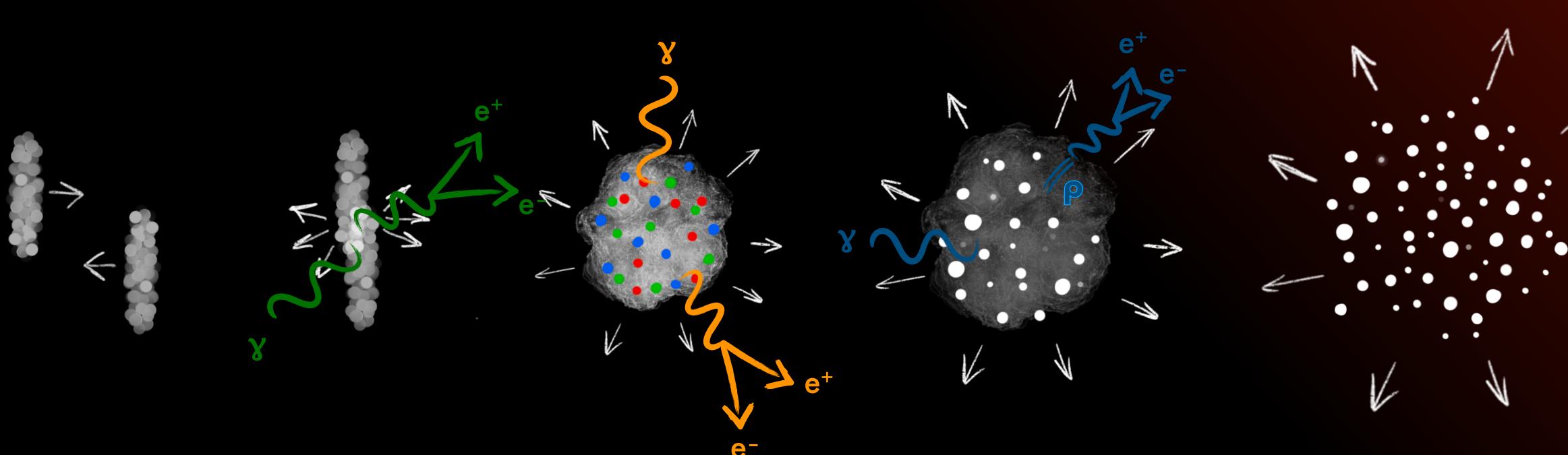
- Invariant mass unaffected by blueshift
- Structure in hadron gas → sensitivity to in medium spectral function of p







### The Sources



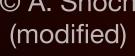
#### Pre-equilibrium

Quark-gluon plasma

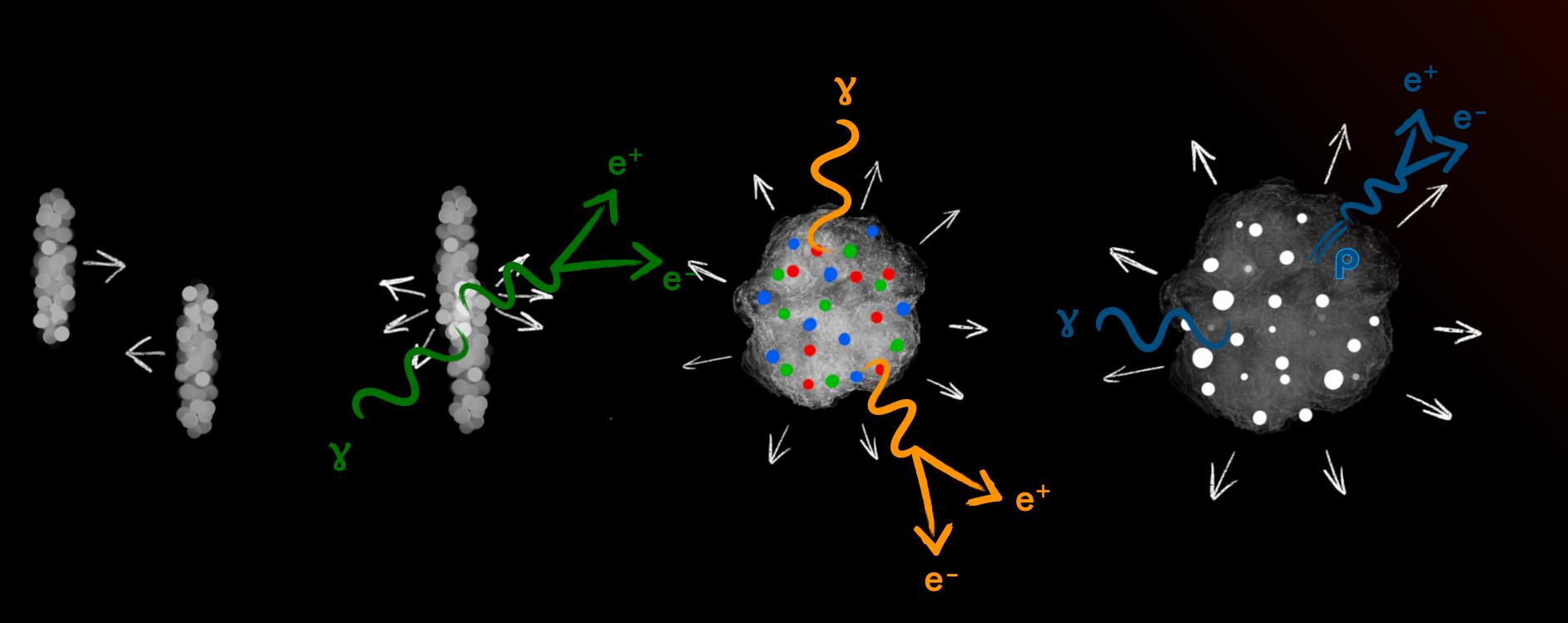
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#### ma Hadron Gas





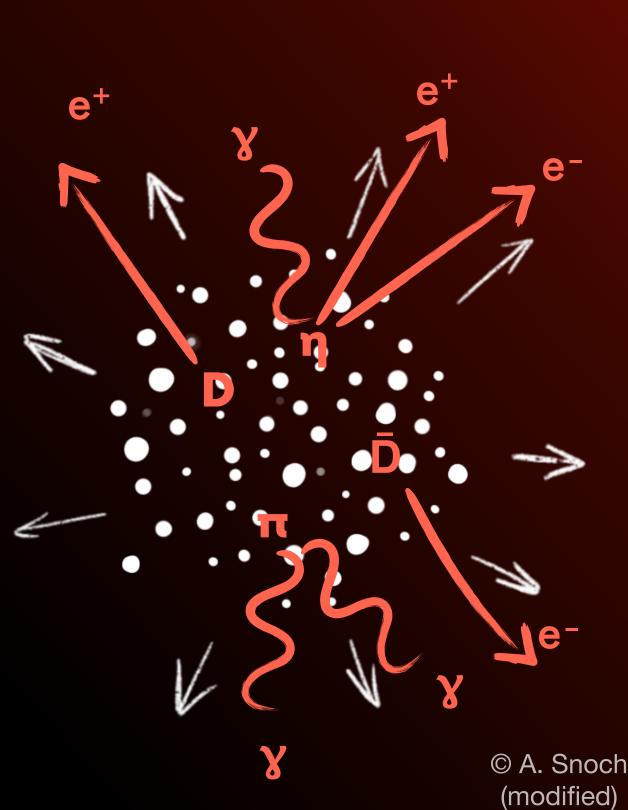
## The Sources and background



#### Pre-equilibrium

Quark-gluon plasma

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#### Hadron Gas

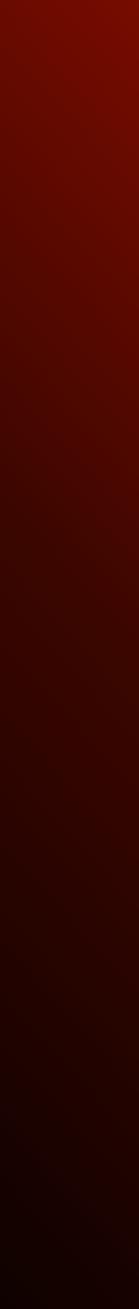
Late stage hadron decays





## The Backgrounds

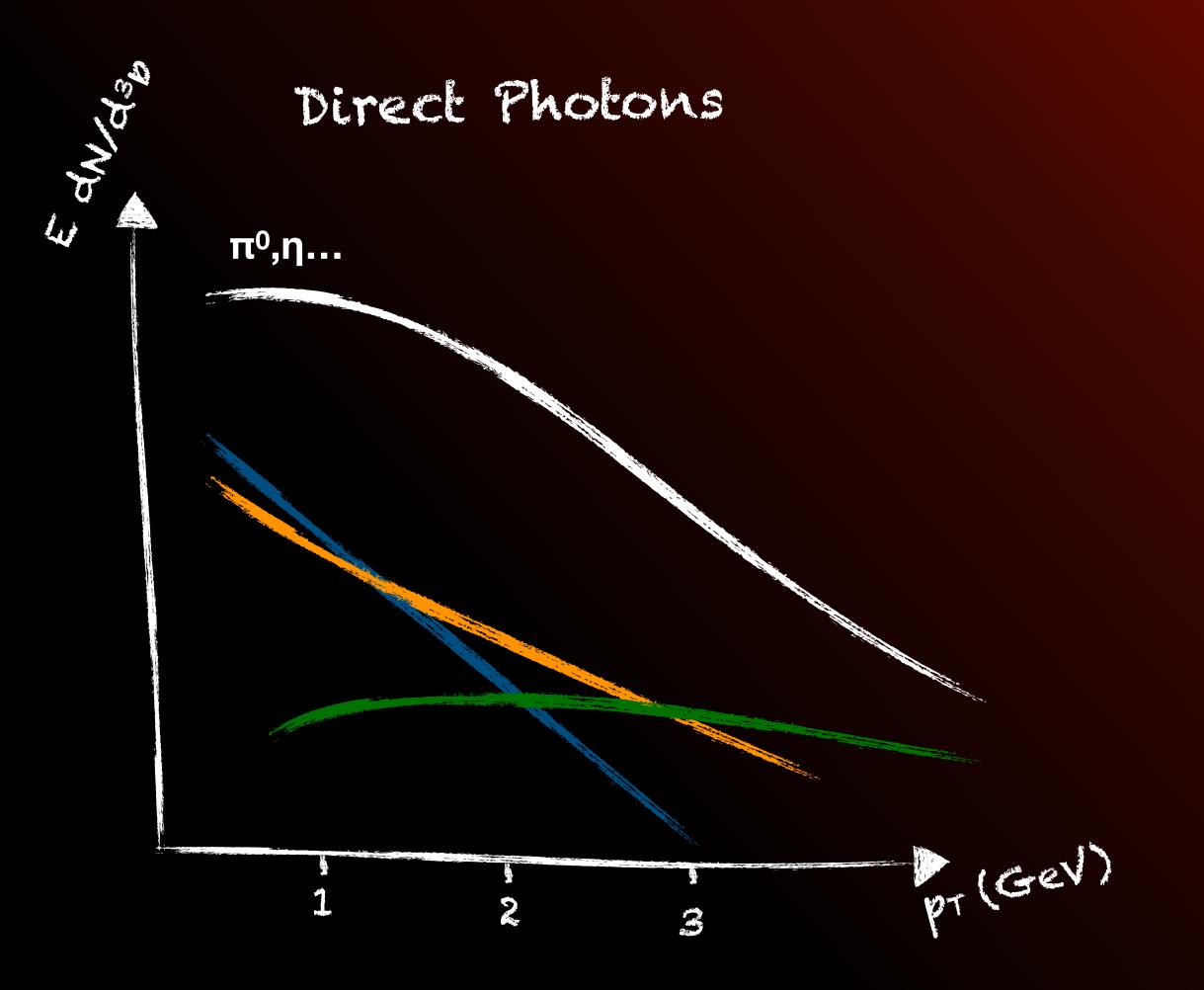
- Dileptons: Combinatorial background of "fake" pairs
- Irreducible physical backgrounds  $\pi^0 \rightarrow \gamma \gamma$  or  $\gamma e^+ e^-$ ,  $\eta \rightarrow \gamma \gamma$  or  $\gamma e e^+ e^$  $c\bar{c} \rightarrow D\bar{D} \rightarrow e^+ e^- X Y$  (semi leptonic decays)
- A precice understanding of the sources of these backgrounds is needed to model and subtract them
- In the case of open heavy flavour experimental techniques can be employed to reduce them (Use the finite decay length)





## The hadronic Cocktail

- Best possible estimate of the estimate of contributions from hadronic decays
- Based on independent measurements
- Parameterised spectra used to model contribution in dielectron and photon spectra



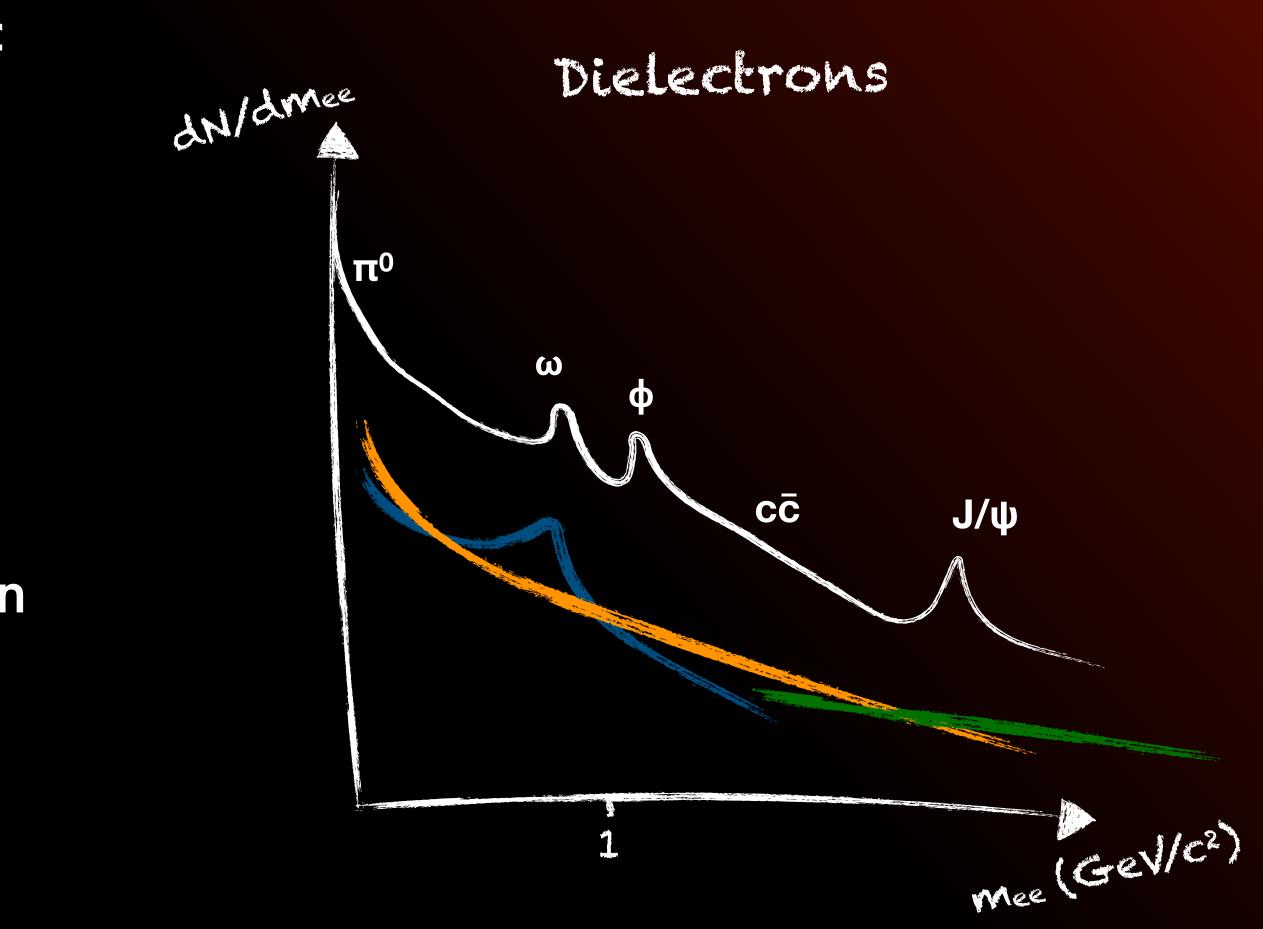
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## The hadronic Cocktail

- Best possible estimate of the estimate of contributions from hadronic decays
- Based on independent measurements
- Parameterised spectra used to model contribution in dielectron and photon spectra
- Heavy flavour challenging to describe in heavy ion collisions (cold nuclear matter + energy loss)







## Summary I

their information to the detectors undisturbed

Momentum and mass allow selection of different sources of EM radiation and production time in the collision

Dielectrons vs direct photons

- Additional dimension (invariant mass)
- No blueshift in T measurement
- Smaller statistical precision
- Additional backgrounds

### Electromagnetic probes are produced in all stages of heavy ion collisions and carry

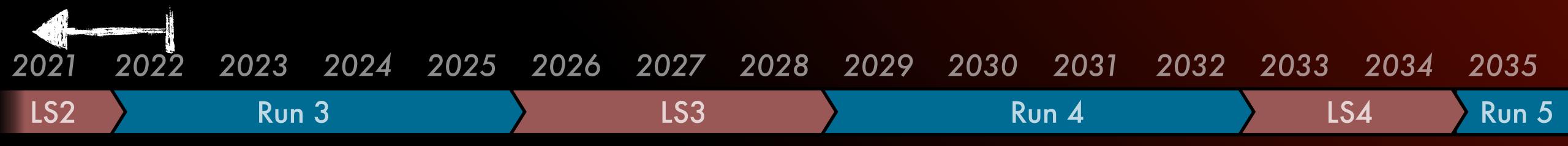


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### Where are we now?



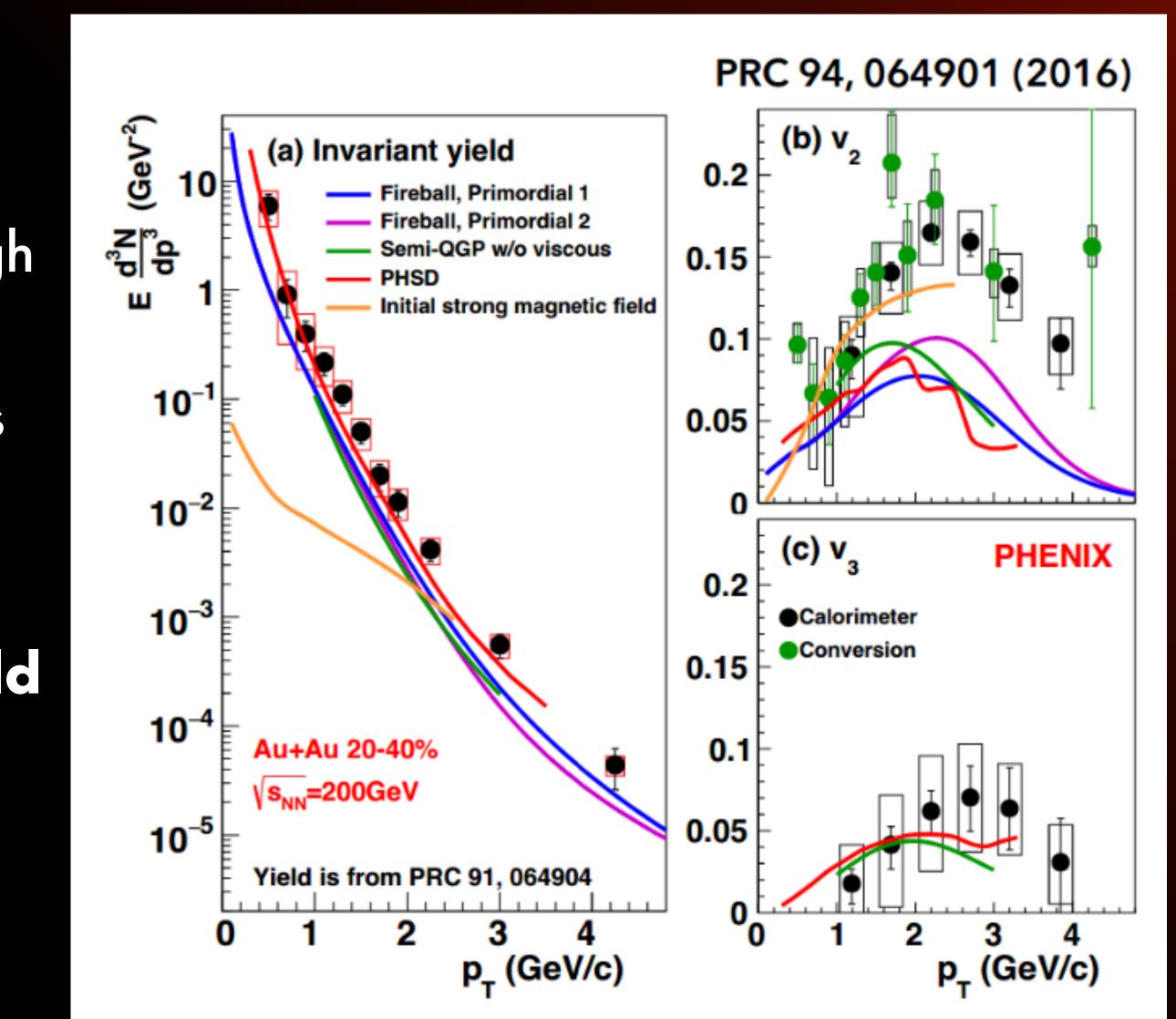


## The direct photon puzzle

Idea:

- Large yield → Early production due to high medium temperature (Partonic source)
- Large v<sub>2</sub> → Late emission since flow needs to build up (hadronic source)

Models are not able to describe yield and elliptic flow v<sub>2</sub> simultaneously

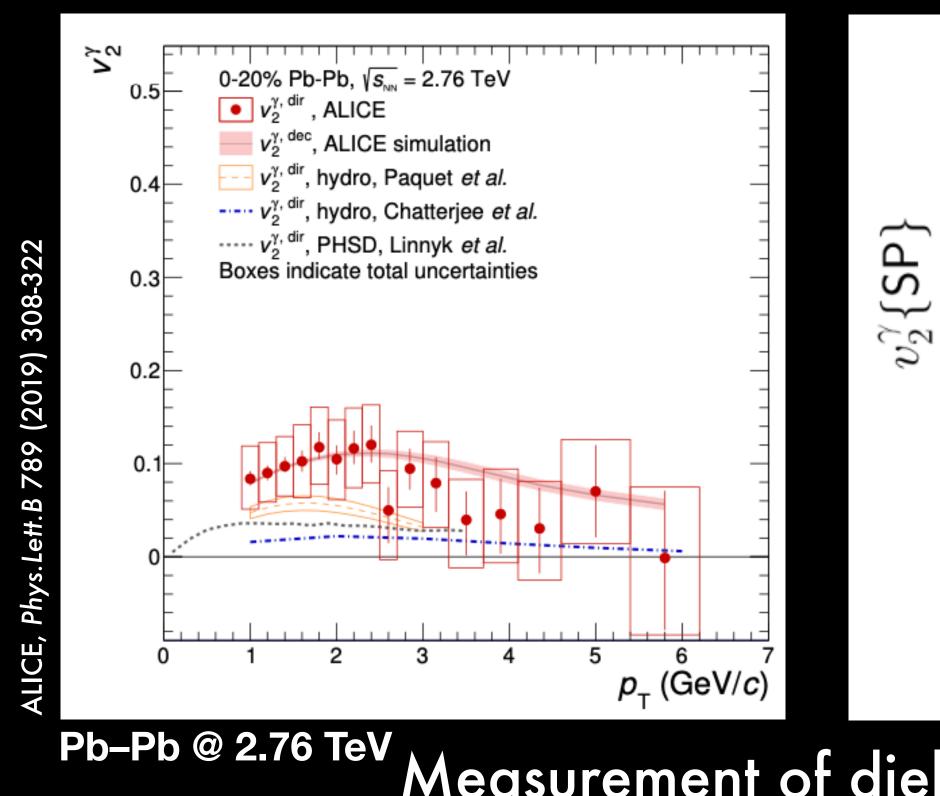




## The direct photon puzzle

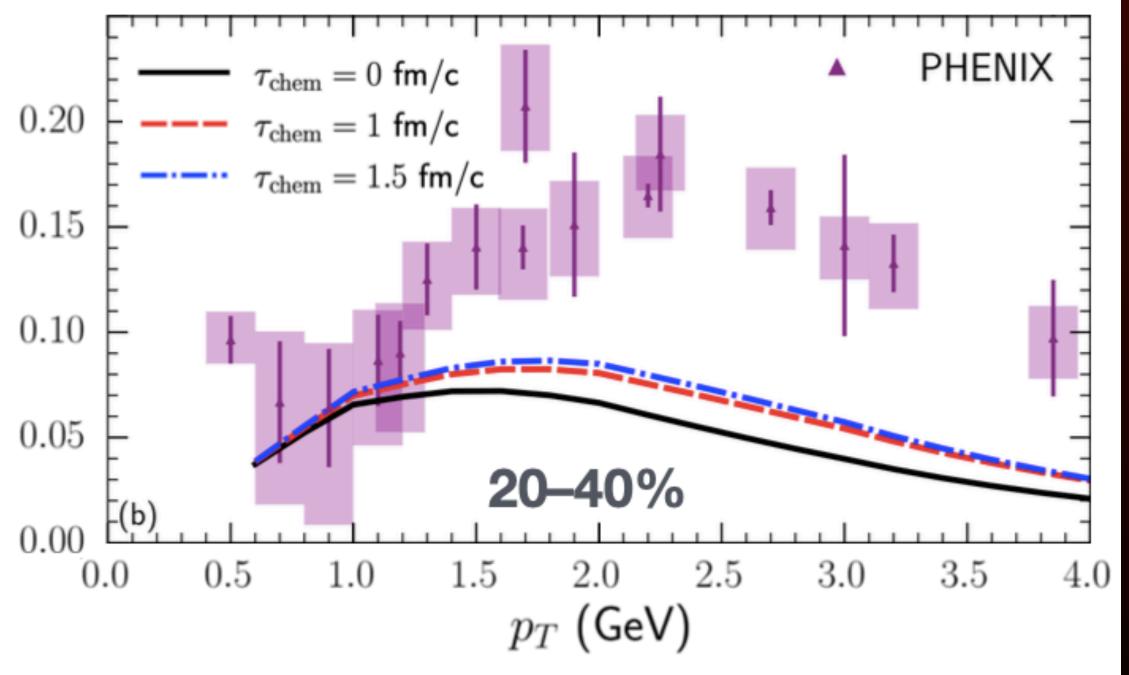
#### $v_{2,\gamma,dir}$ of direct photons at RHIC (PHENIX) and LHC (ALICE) $v_{2,\pi}$

#### Photon production dominated by late stage?

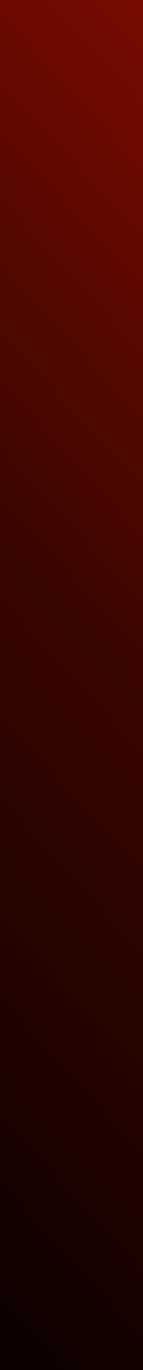


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Au-Au @ 200 GeV



#### Measurement of dielectron $v_2$ would be helpful

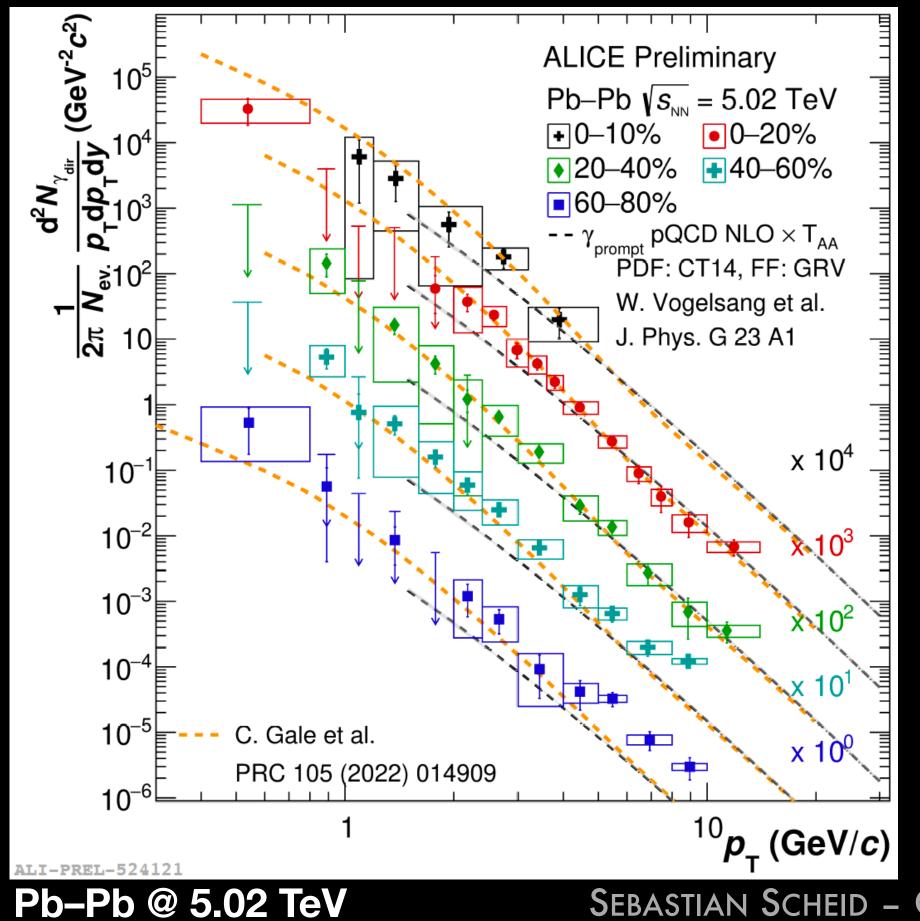


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### ALICE - Direct Photons

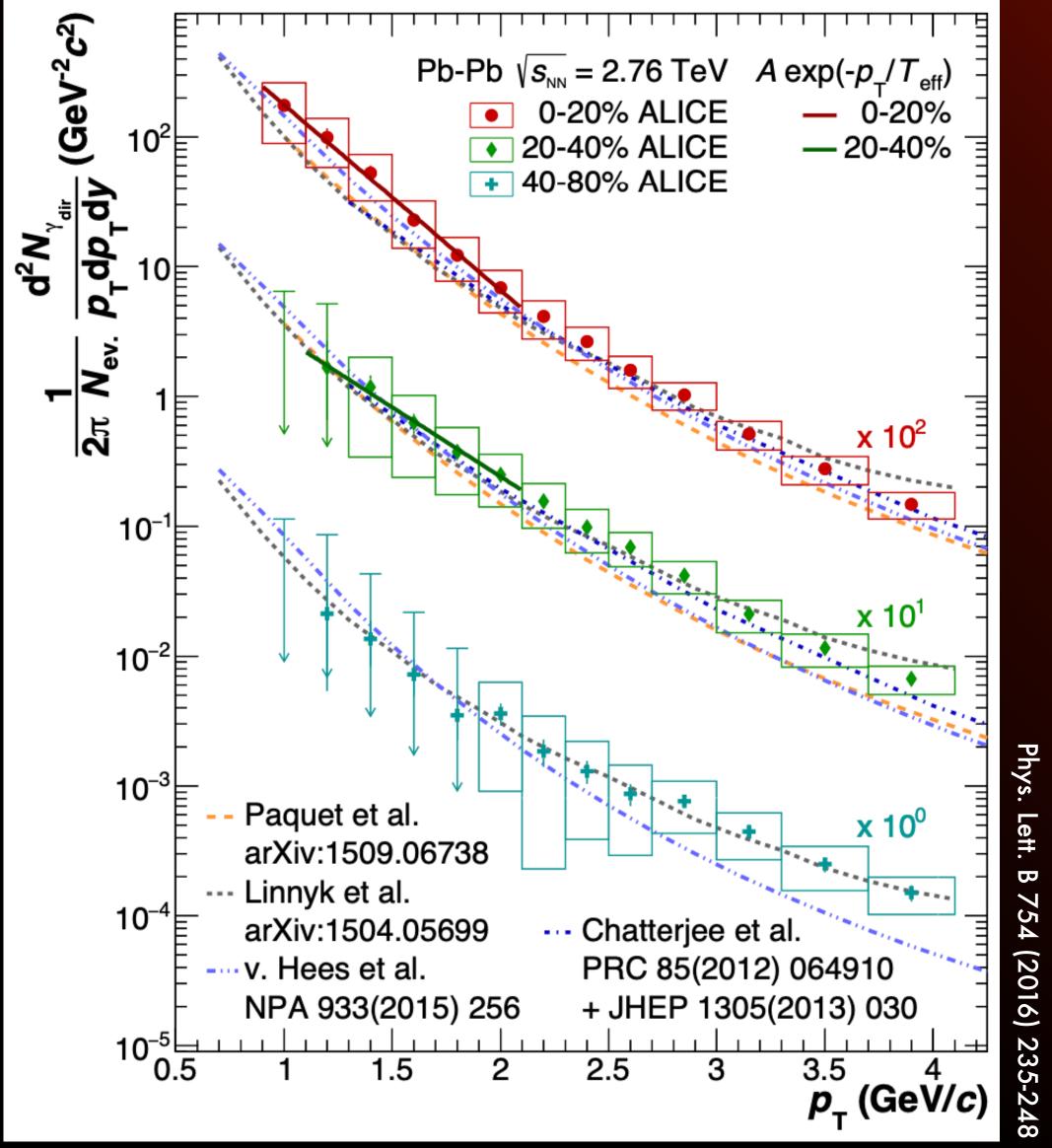
#### Data in agreement with models

#### T<sub>eff</sub> results from 2.76 TeV: $0-20\%: 297 \pm 12 \text{ (stat)} \pm 41 \text{ (syst)}$



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**Pb-Pb @ 2.76 TeV** 

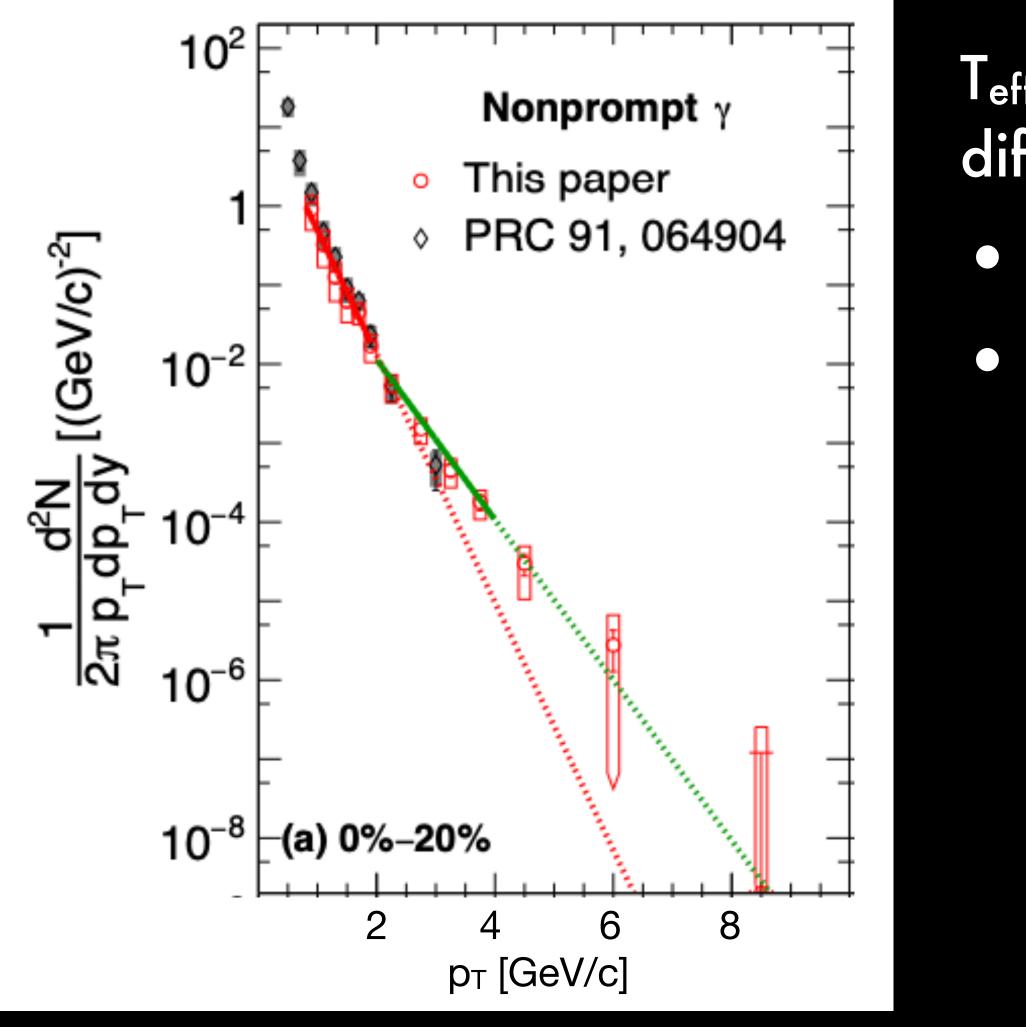


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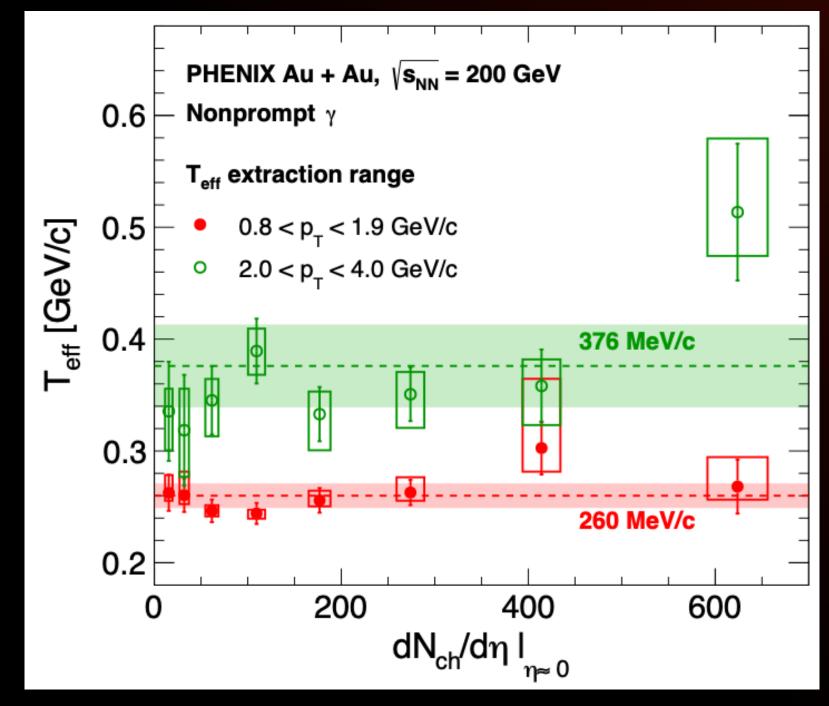
## PHENIX - Direct Photons

#### Au-Au @ 200 GeV



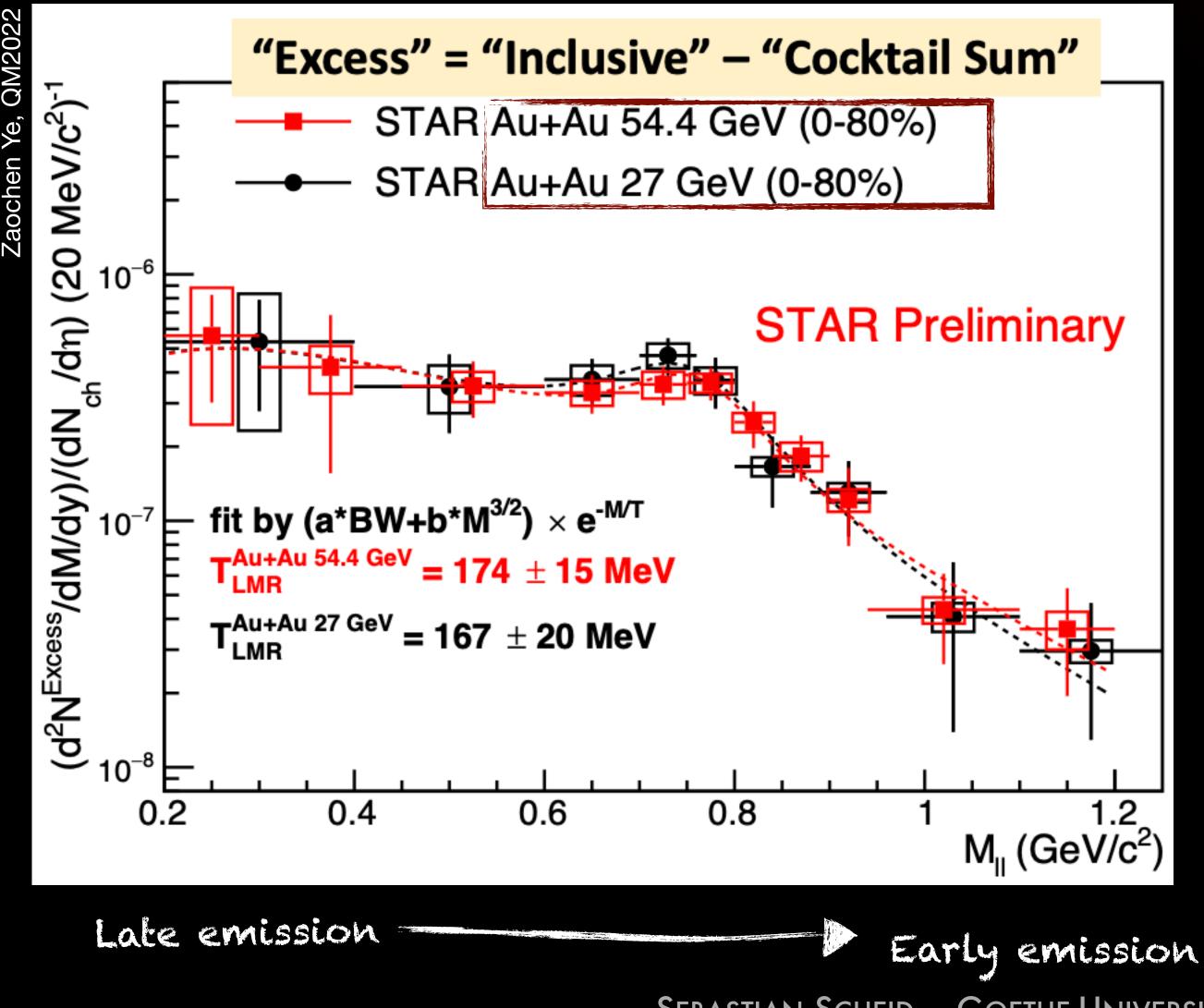
PHENIX, arXiv: 2203.17187

- $T_{eff}$  extracted from direct photon measurement in different  $p_T$  ranges
- Increases with pt
- Earlier emission? Different source?





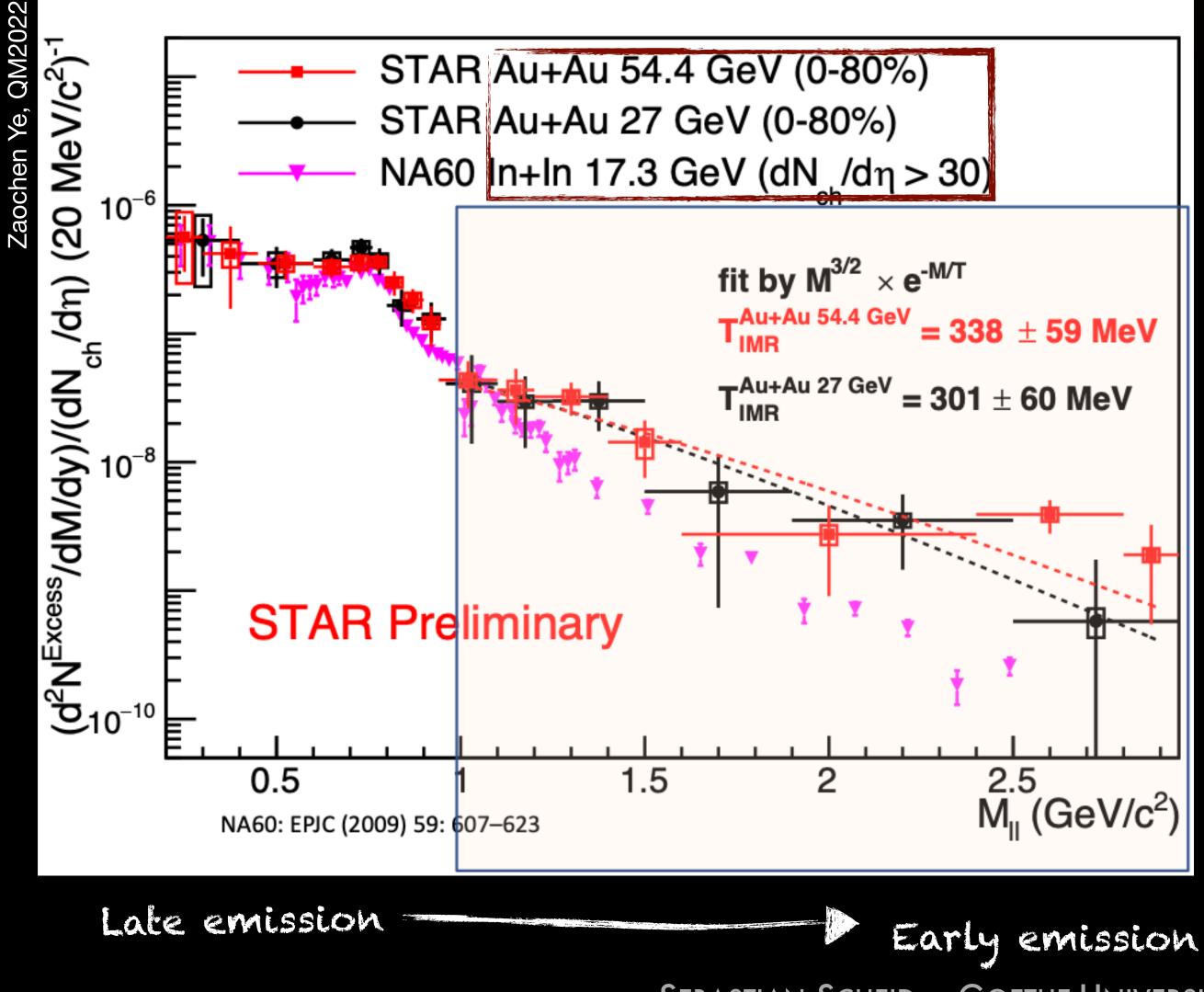
## STAR - Die ectrons



STAR sees clear enhancement wrt to hadronic cocktail Fit of p Lineshape (Breit-Wigner) + thermal Similar result with fit to NA60 In-In data at 17.3 GeV (165±4 MeV)

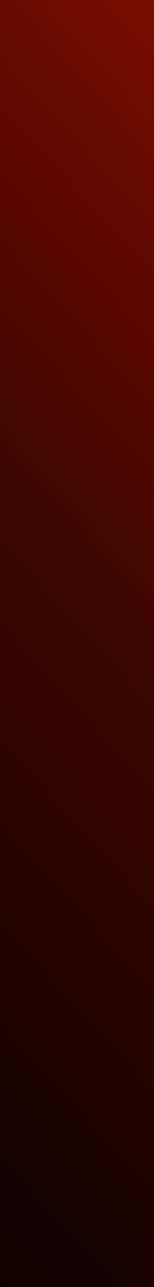


## STAR - Dielectrons II



#### T extracted at higher mass significantly higher NA60 = 205 ± 15 MeV

Early emission Sebastian Scheid – Goethe-Universität Frankfurt – EM Probes

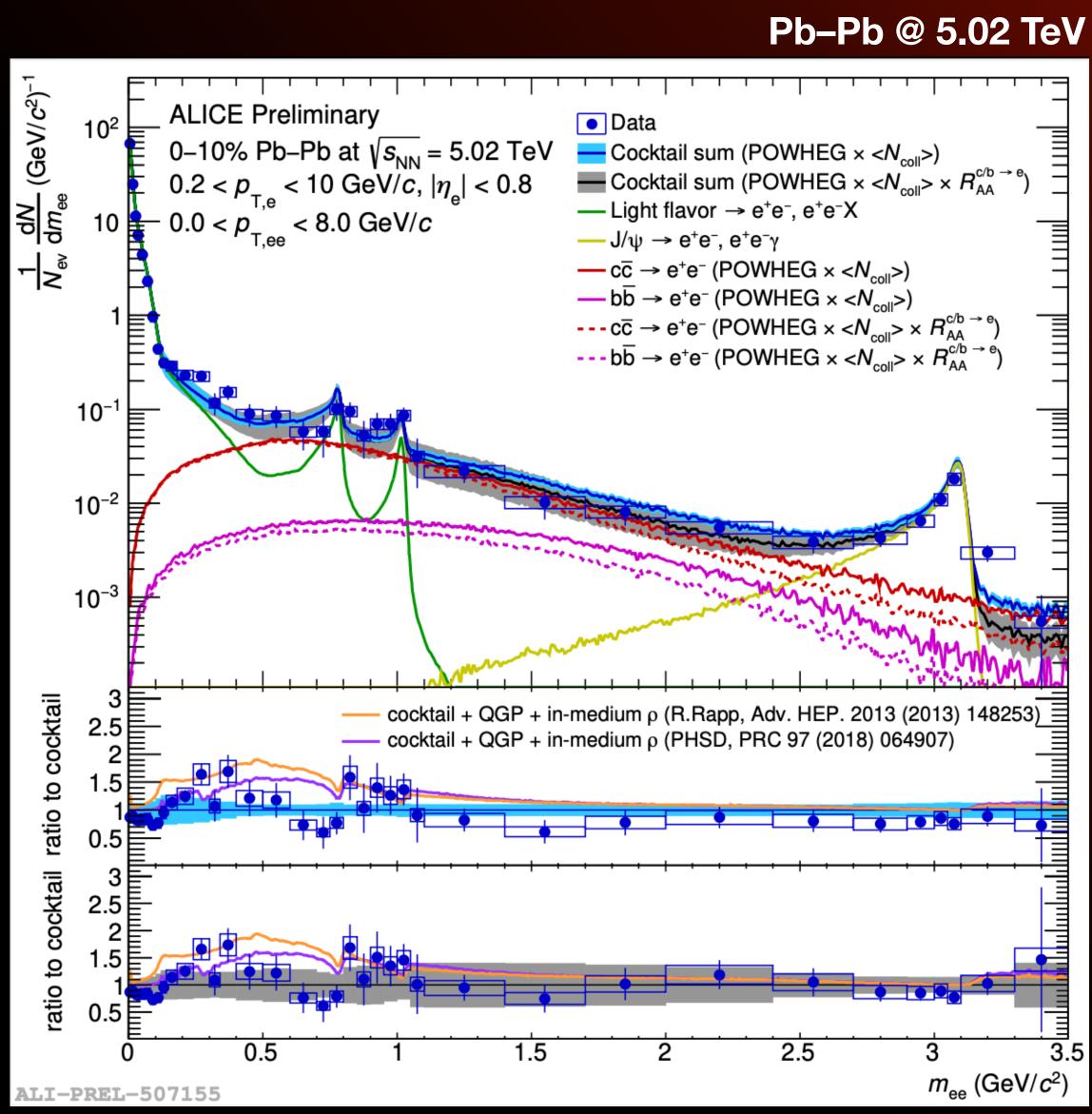




## ALICE - Dielectrons

- Comparison of dielectron spectrum with hadronic cocktail
- HF cross section much higher than at RHIC
- Two different approaches for HF
- N<sub>coll</sub> scaled from pp measurement
- Additional  $R_{AA}$  of  $c/b \rightarrow e$

Handeling HF contribution crucial



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## Distance of Closest Approach (DCA)

Heavy flavour hadrons have a finite decay length (150 - 500 µm)

$$DCA_{ee} = \sqrt{\frac{DCA_1^2 + DCA_2^2}{2}}$$

with  $DCA_{1/2}$  normalised to respective resolution

First steps in pp in Pb-Pb collisions by ALICE

- DCA DCA2

Resolution

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## Summary II

Possibility to extract temperature in different mass and momentum regions at RHIC energies

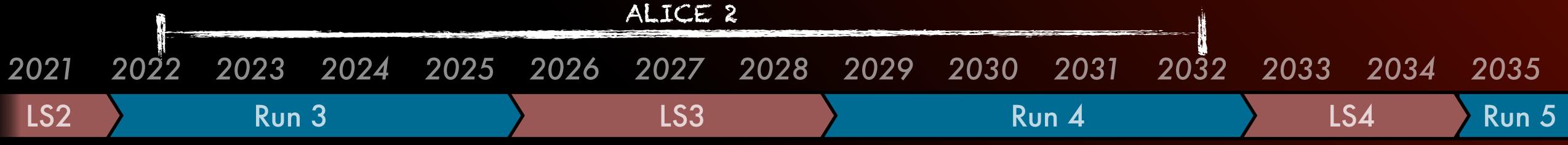
- Higher mass/p<sub>T</sub> yields in all cases higher temperature
- Can be attributed to earlier emission and possibly pre-equilibrium sources
- High temperature in IMR (300 MeV) by STAR could be due pre-equilibrium

Measurements at LHC challenging

- Small signal to background ratio
- Large heavy-flavour background (dielectrons)

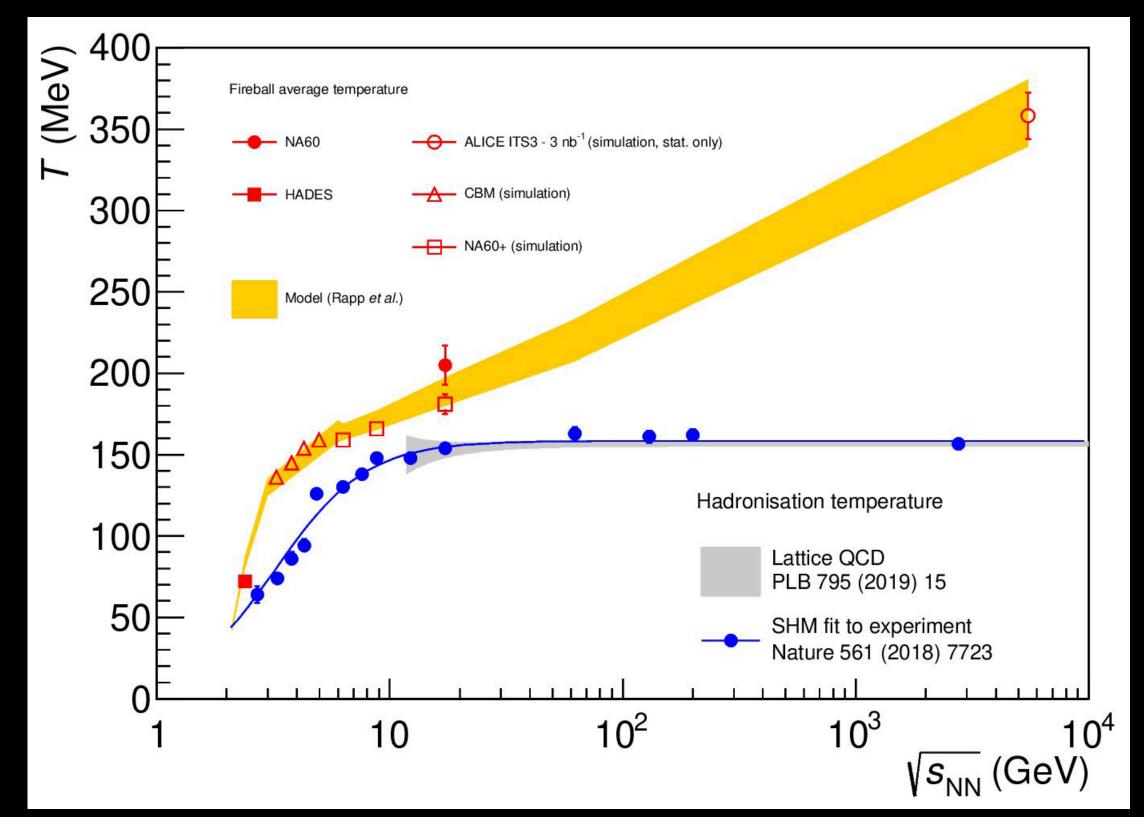


#### What is in the near future? Now - End of LHC Run 4 (2032)



## The situation end of Run 4

#### Inclusive temperature measurement of the QGP

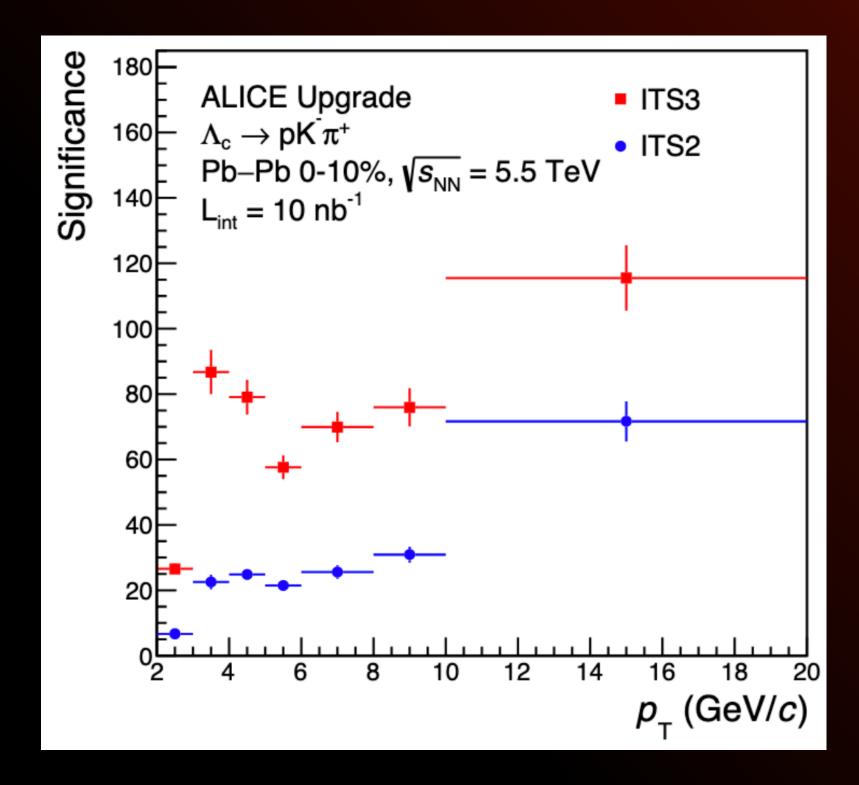


NA60, AIP Conf.Proc. 1322 (2010) 1, 1-10 ALICE, CERN-LHCC-2019-018 CBM, Nucl. Phys. A 982 (2019) 163 NA60+, SPSC-EOI-019

R. Rapp et al., Phys. Lett. B 753 (2016) 568 HADES, Nature Physics 15 (2019) 10, 1040-1045 T. Galatyuk et al., Eur. Phys. J. A52 (5) (2016) 131 Lattice QCD, Phys. Lett. B 795 (2019) 15 SHM, Nature 561 (2018) 7723, 321-330

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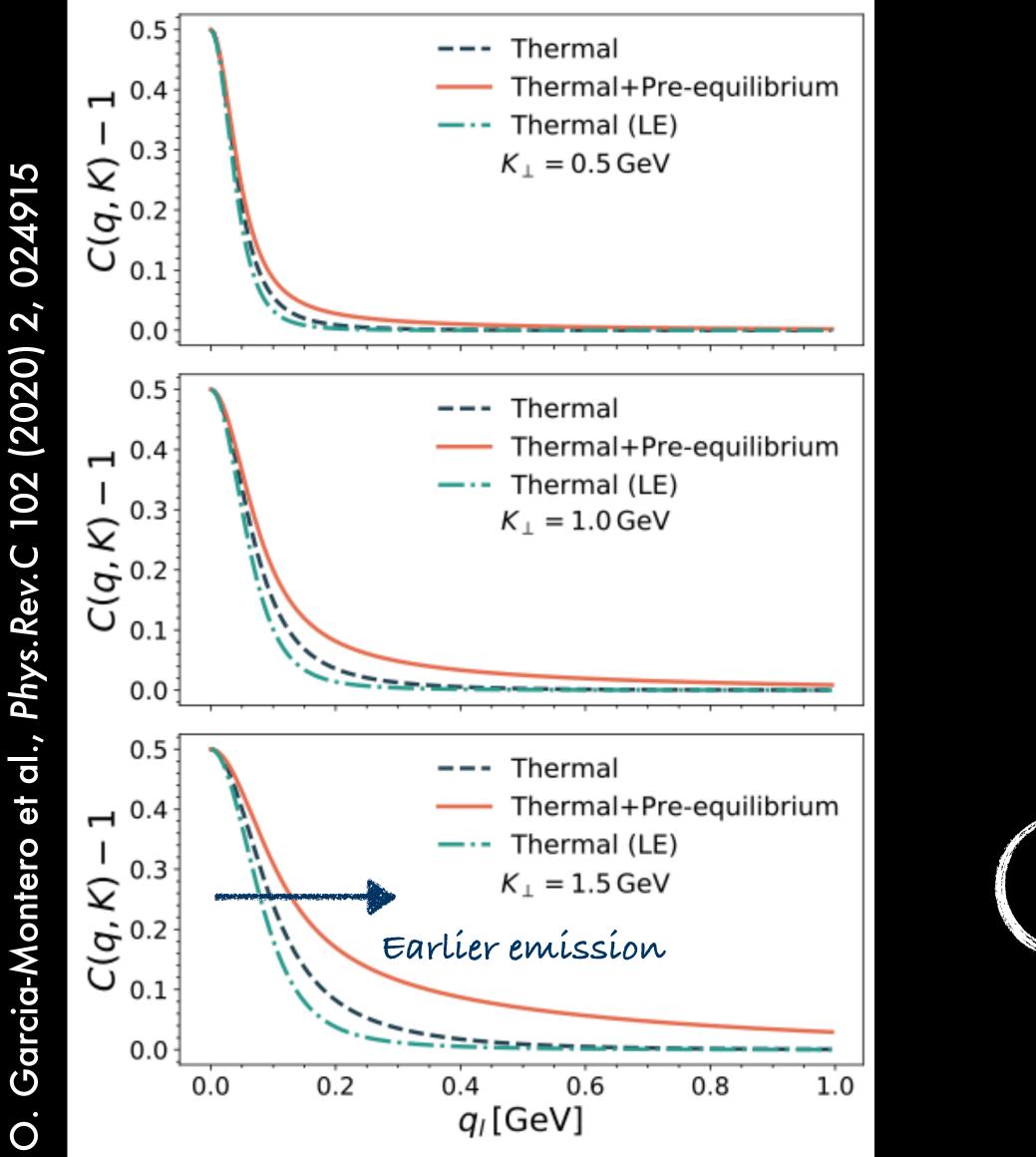


#### Precise measurements of heavy flavour hadron spectra





## The situation end of Run 4

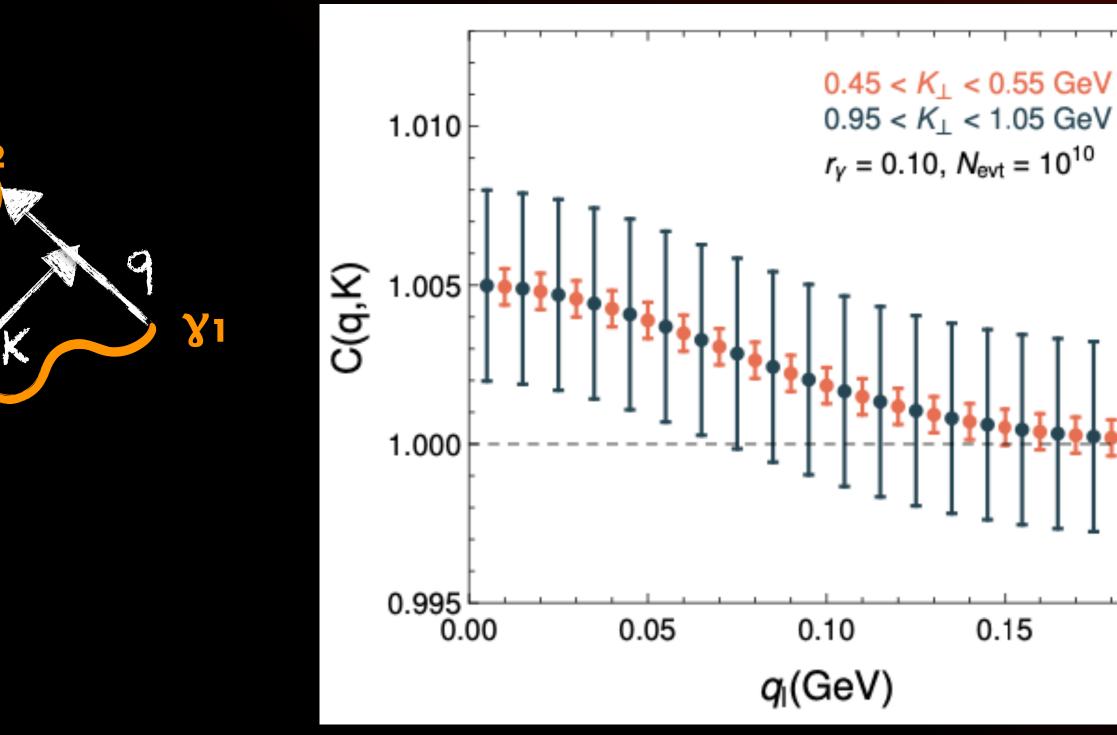


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**Direct photon HBT correlations** 

- Sensitivity to size of emitting source (production time)
- Not clear if precision sufficient to disentangle sources

0-20% Pb-Pb @ 2.76 TeV



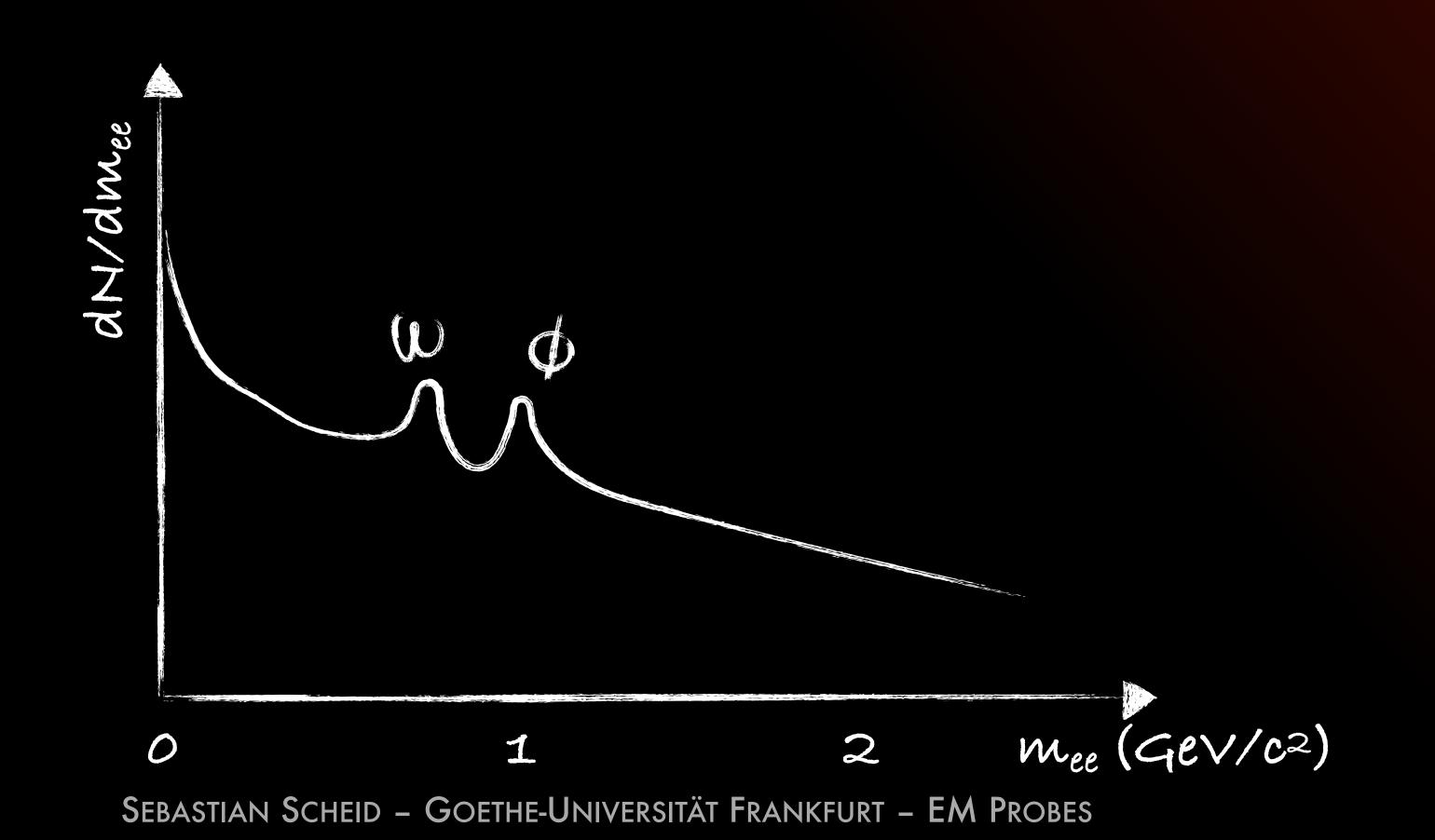






#### Where to keep looking? In the 2030s

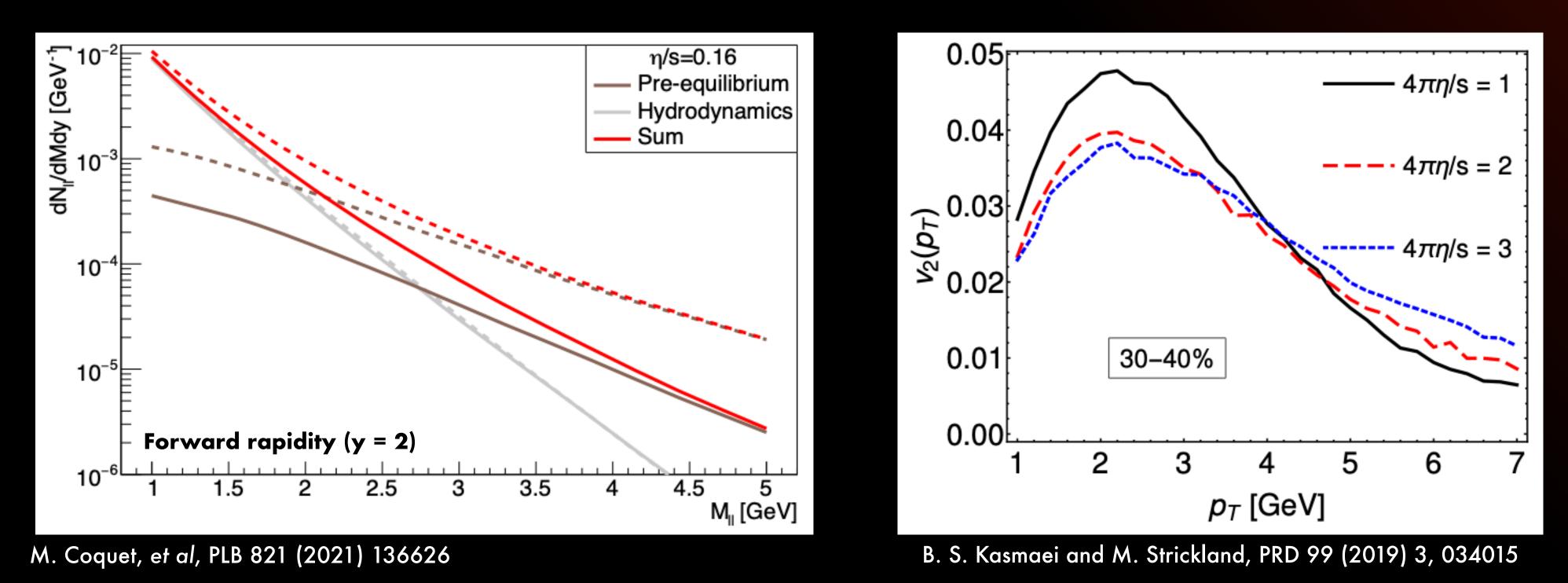






### The pre-equilibrium High mass and pt

Sensitivity to initial anisotropies and  $\eta/s$ , not accessible with hadronic probes  $\rightarrow$  Sensitivity not reached in RUN 3/4 (Now - 2032)



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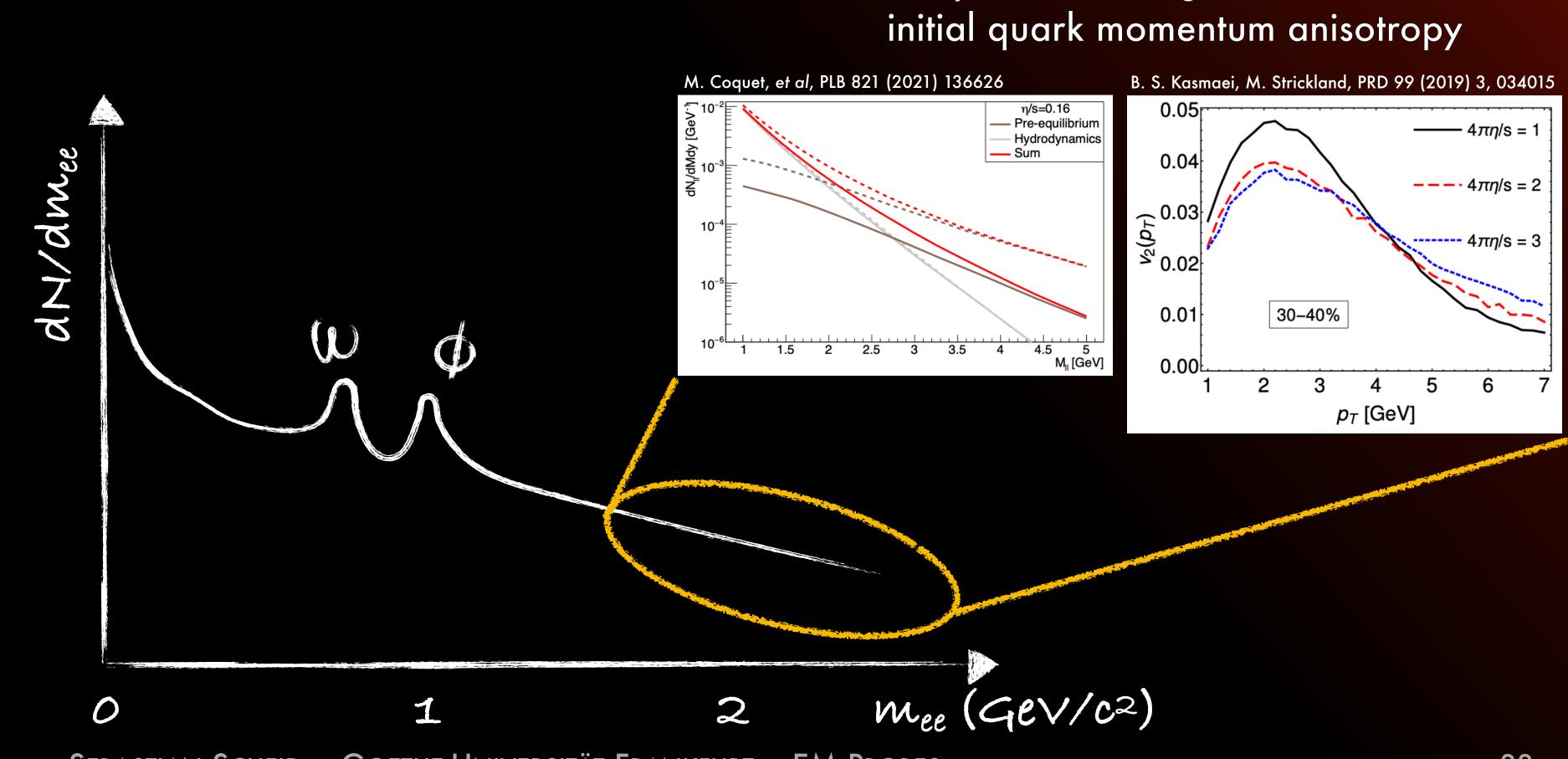
# Dileptons from earliest collision stage more abundant than the QGP contribution at higher masses







## Where to keep looking?

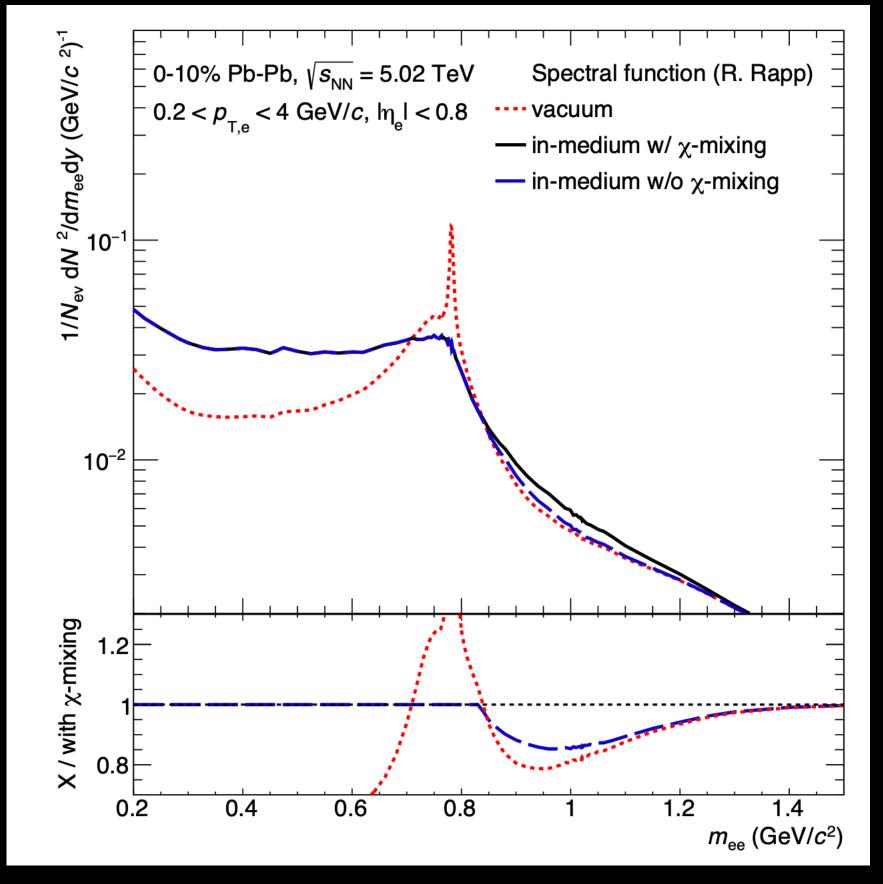


#### High mass and pt Early collision stages & initial quark momentum anisotrop



# The hadronic phase The in medium $\rho$





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Can we learn something from more precision in the p spectral function?

- Chiral mixing of  $\rho$ -a<sub>1</sub> changes the shape of dielectron spectrum in 0.9 <  $m_{ee}$  1.4 GeV/c<sup>2</sup>
- Measurement with precision higher >10% necessary

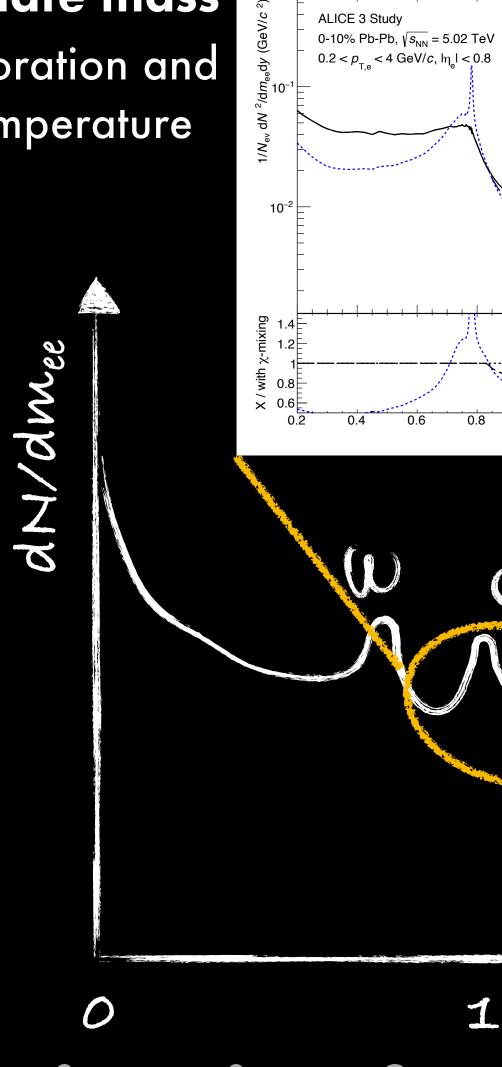




## Where to keep looking?

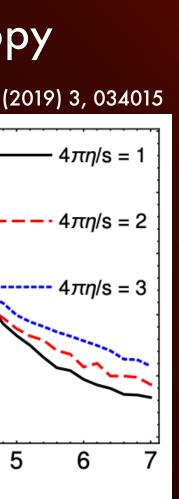
#### Intermediate mass

#### Chiral restoration and medium temperature



#### R. Rapp, private communication

#### ρ spectral functior in vacuum High mass and pt — in medium w/ $\chi$ -mixing — in medium w/o $\chi$ -mixing Early collision stages & initial quark momentum anisotropy M. Coquet, et al, PLB 821 (2021) 136626 B. S. Kasmaei, M. Strickland, PRD 99 (2019) 3, 034015 dN//dMdy [GeV-10 η/s=0.16 0.05r - Pre-equilibrium Hydrodynamics Sum 0.04 0.03 (1 *a*)<sup>2</sup> 0.02 1.4 10 $m_{\rm ee} \, ({\rm GeV}/c^2)$ 10-0.01 30-40% 4.5 3.5 M<sub>II</sub> [GeV] 0.00 2 3 *p*<sub>7</sub> [GeV] $m_{ee}$ (GeV/c<sup>2</sup>) 2





### Electric Conductivity Small mass and pt

#### Electric conductivity of the medium

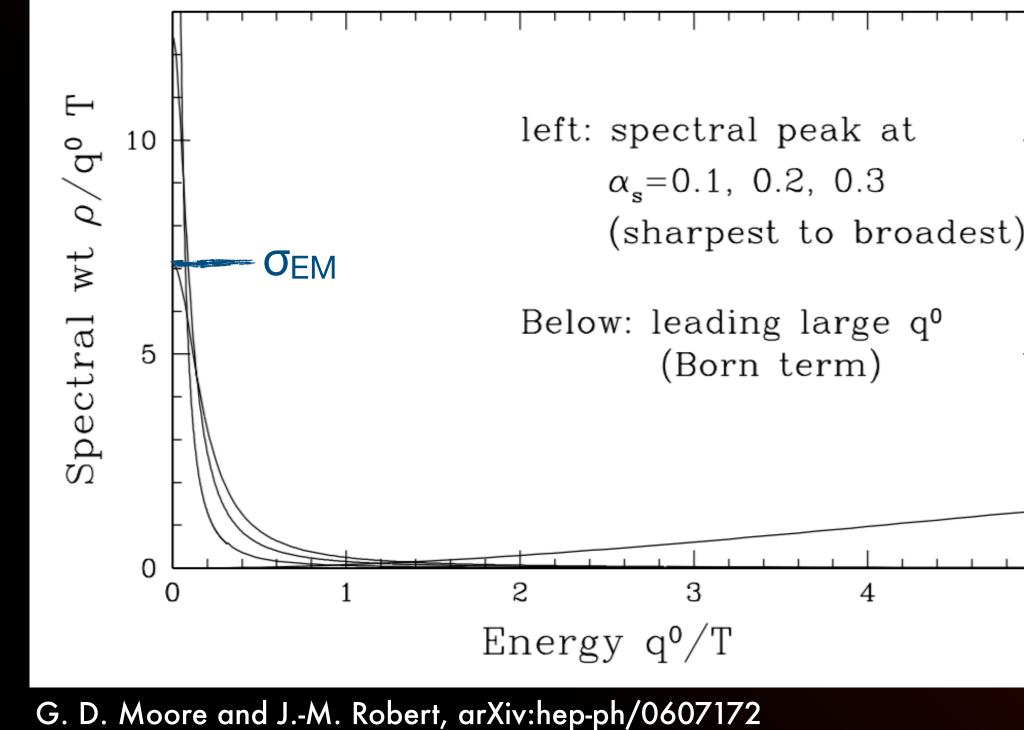
Fundamental property of the medium

- Wide range of theoretical values around
- Input for CME calculations

Connected to dielectron spectral function

$$\sigma_{\rm EM} = -e^2 \lim_{q_0 \to 0} \left( \frac{\partial}{\partial q_0} \operatorname{Im}\Pi(q_0, q = 0, T) \right)$$

Transport peak in the limit of very low mass and pt











## **Electric Conductivity** Small mass and pt

## Electric conductivity of the medium

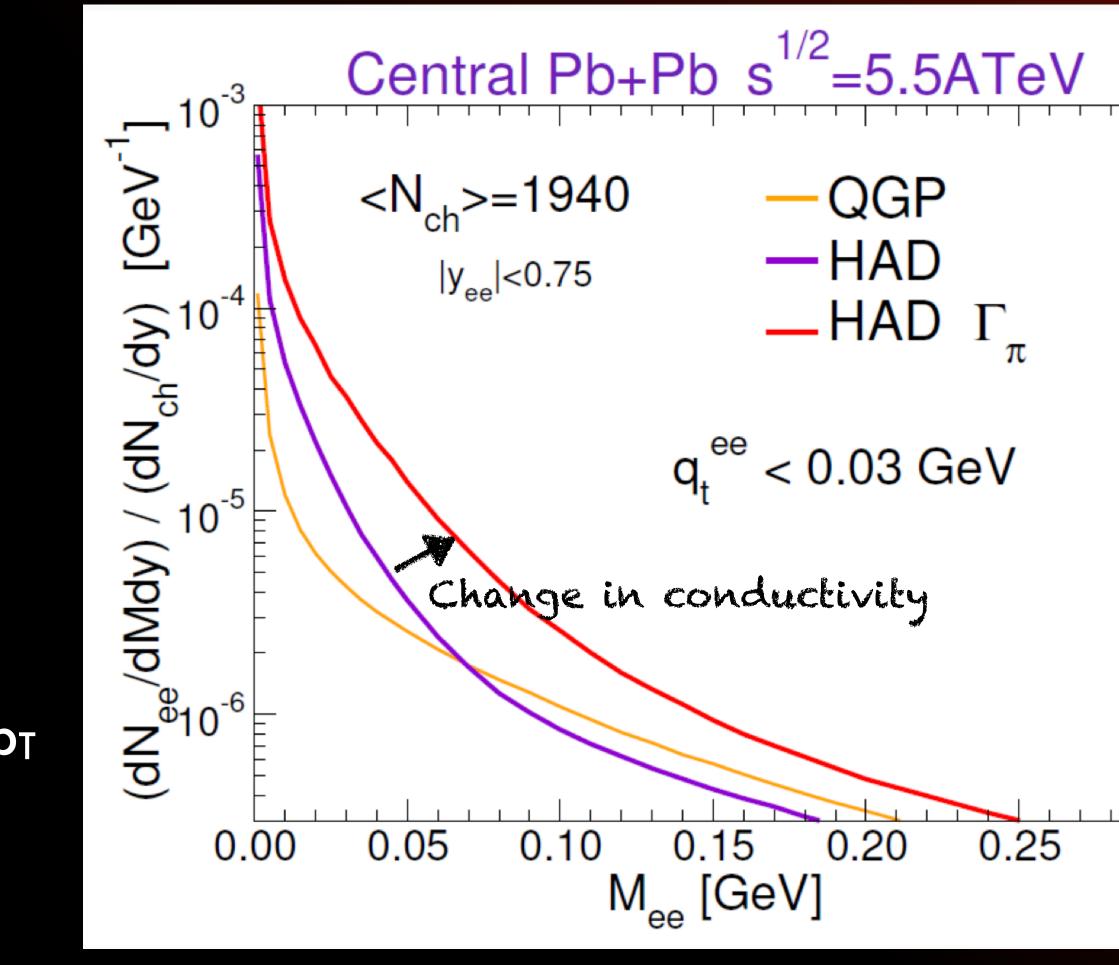
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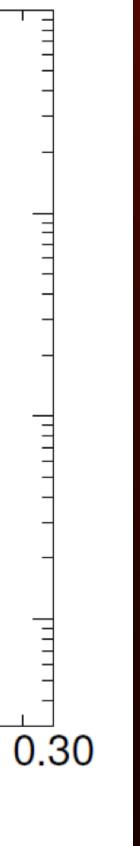
Connected to dielectron spectral function

$$\sigma_{\rm EM} = -e^2 \lim_{q_0 \to 0} \left( \frac{\partial}{\partial q_0} \operatorname{Im}\Pi(q_0, q = 0, T) \right)$$

Transport peak in the limit of very low mass and  $p_T$ Sensitivity of spectral function to conductivity not only at  $q_0 \rightarrow 0$  limit

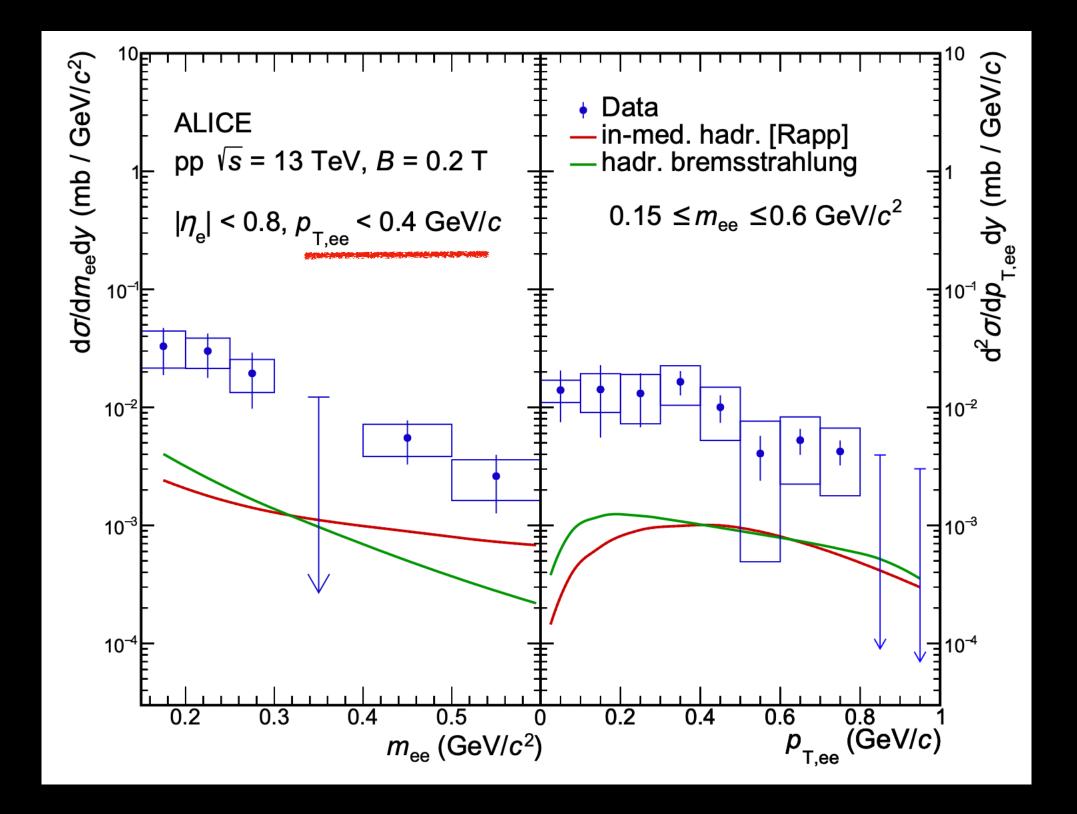


R. Rapp, EMMI RRTF, 13.9.2021





## Soft Dielectrons



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## pp collisions at $\sqrt{s} = 13$ TeV

- Lower B field than nominal  $\rightarrow$  smaller p<sub>T</sub> possible
- Excess of low p<sub>T</sub> dielectrons observed
- Not described by thermal radiation/hadronic bremsstrahlung

At very low pt: y produced mainly via inner Bremsstrahlung

Can be calculated via Low theorem if in- and outgoing particles can be measured





# Where to keep looking?

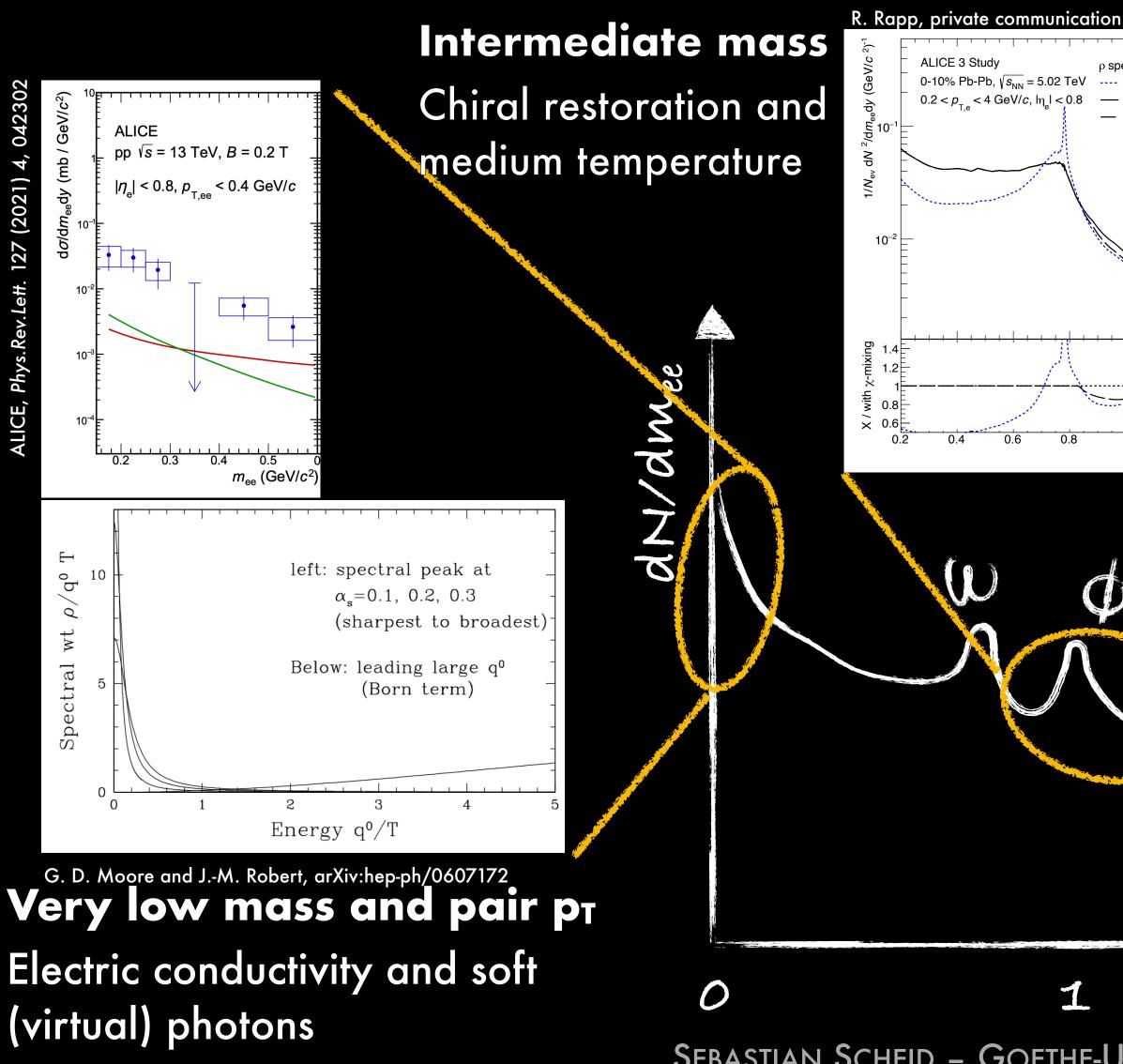
## The softest Photons

Experiment	Year	Collision energy	Photon p <sub>T</sub>	Photon/Brems Ratio	Detection method	Reference
π+p	1979	10.5 GeV	p⊤ < 30 MeV/c	<b>1.25 ± 0.25</b>	Bubble chamber	<u>Goshawk et al.,</u> PRL 43,1065 (1979
K+p WA27, CERN	1984	70 GeV	p⊤ < 60 MeV/c	4.0 ± 0.8	Bubble chamber (BEBC)	<u>Chiliapnikov et al.,</u> <u>PLB 141, 276 (1984</u>
π+p CERN, EHS, NA22	1991	250 GeV	р <sub>т</sub> < 40 MeV/c	6.4 ± 1.6	Bubble chamber (RCBC)	<u>Botterweck et al.,</u> Z. Phys. C 51, 541 (19
K+p CERN, EHS, NA22	1991	250 GeV	p⊤ < 40 MeV/c	6.9 ± 1.3	Bubble chamber (RCBC)	Botterweck et al., Z. Phys. C 51, 541 (19
π⁻p CERN, WA83, OMEGA	1993	280 GeV	p <sub>T</sub> < 10 MeV/c (0. 2 < E <sub>γ</sub> < 1 GeV)	7.9 ± 1.4	Calorimeter	<u>Banerjee et al.,</u> PLB 305, 182 (1993
p–Be	1993	450 GeV	р⊤ < 20 MeV/c	< 2	Pair conversion, calorimeter	<u>Antos et al.,</u> <u>Z. Phys. C 59, 547 (19</u>
p–Be, p–W	1996	18 GeV	р <sub>т</sub> < 50 MeV/c	< 2.65	Calorimeter	<u>Lissauer et al.,</u> <u>PRC 54, 1918 (1996</u>
π⁻p CERN, WA91, OMEGA	1997	280 GeV	p <sub>T</sub> < 20 MeV/c (0. 2 < E <sub>γ</sub> < 1 GeV)	7.8 ± 1.5	Pair conversion	<u>Belogianni et al.,</u> <u>PLB 408, 487 (1997</u>
π⁻p CERN, WA91, OMEGA	2002	280 GeV	p <sub>T</sub> < 20 MeV/c (0. 2 < E <sub>γ</sub> < 1 GeV)	5.3 ± 1.0	Pair conversion	<u>Belogianni et al.,</u> <u>PLB 548, 122 (2002</u>
pp CERN, WA102, OMEGA	2002	450 GeV	$p_T < 20 \text{ MeV/c}$ (0. 2 < E <sub><math>\gamma</math></sub> < 1 GeV)	4.2 ± 0.8	Pair conversion	<u>Belogianni et al.,</u> PLB 548, 129 (2002
e⁺e⁻→2 jets CERN, Delphi	2006	91 GeV (CM)	p <sub>T</sub> < 80 MeV/c (0. 2 < E <sub>γ</sub> < 1 GeV)	4.0 ± 0.3 ± 1.0	Pair conversion	<u>DELPHI,</u> EPJC 47, 273 (2006
e+e-→µ+µ- CERN, Delphi	2008	91 GeV (CM)	p⊤ < 80 MeV/c (0. 2 < E <sub>γ</sub> < 1 GeV)	~1	Pair conversion	<u>DELPHI,</u> EPJC 57, 499 (2008
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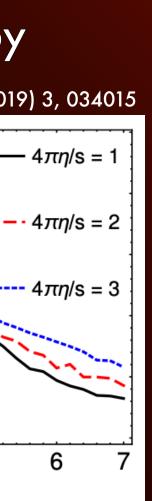
K. Reygers EMMI RRTF, Sep 14th 2021



# Where to keep looking?



#### in vacuum High mass and pt in medium w/ γ-mixing — in medium w/o χ-mixing Early collision stages & initial quark momentum anisotropy M. Coquet, et al, PLB 821 (2021) 136626 B. S. Kasmaei, M. Strickland, PRD 99 (2019) 3, 034015 dN<sub>l</sub>/dMdy [GeV<sup>-1</sup>] η/s=0.16 0.05r Pre-equilibrium Hydrodynamics Sum 0.04 0.03 (1 *a*)<sup>2</sup> 0.02 $m_{\rm ee}$ (GeV/ $c^2$ ) 10 10-0.01 30-40% 4.5 M<sub>II</sub> [GeV] 0.00 2 3 *р*<sub>7</sub> [GeV] $m_{ee}$ (GeV/c<sup>2</sup>) 2





# Thinking about an experiment

## What do we need?

- Sufficient statistics for multi dimensional analysis
- Measure photons and electrons with high efficiency and accuracy
- Rejection of heavy flavour contribution in dielectrons
- Low p<sub>T</sub> photon reconstruction

## How do we get that?

- Possibility to run at high rates
- Particle identification and tracking over a broad momentum range and large acceptance Photon measurements via conversion and calorimetry





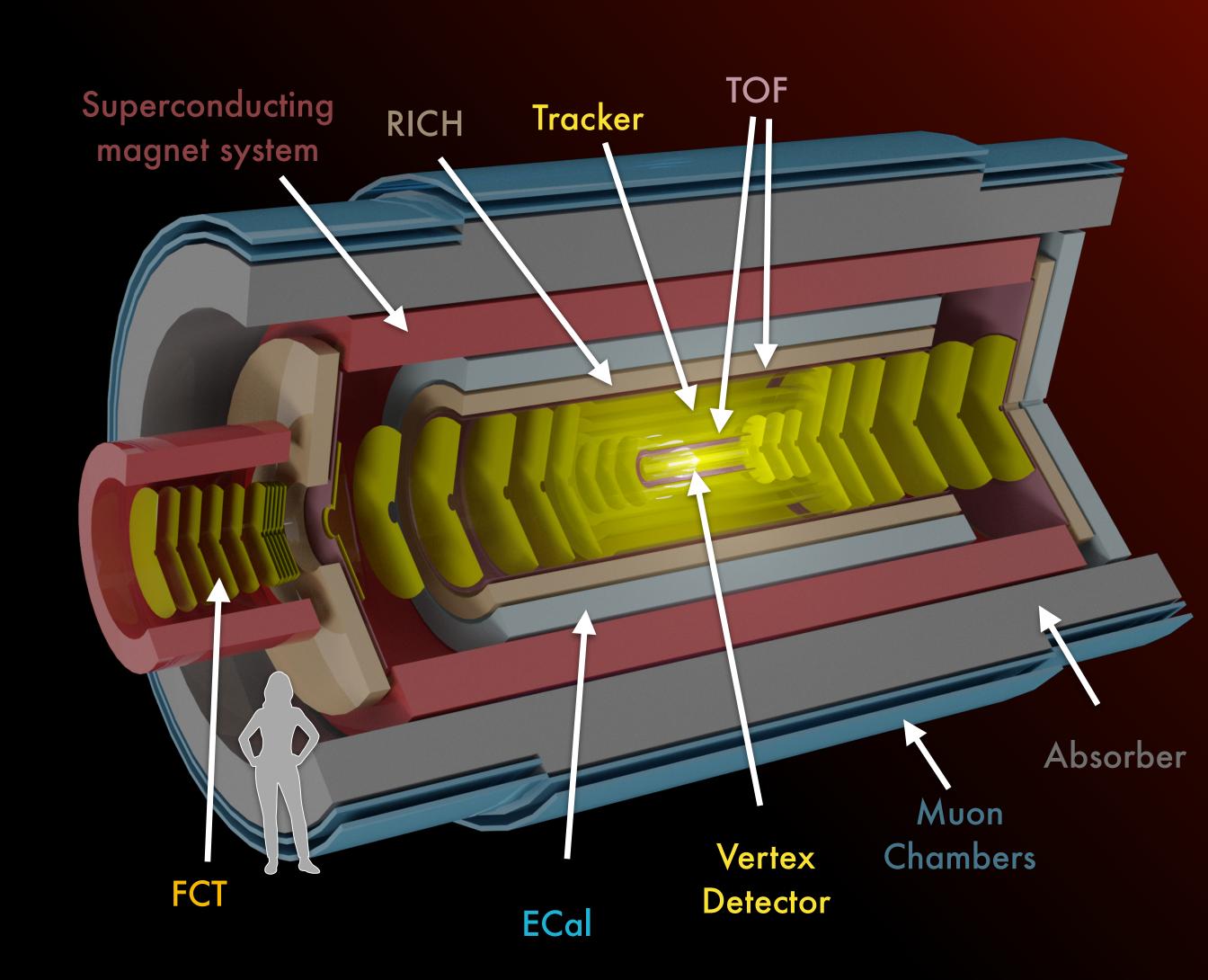
High precision tracking

**Electron Identification** 

**Electromagnetic Calorimeter** 

Forward Photon Converter

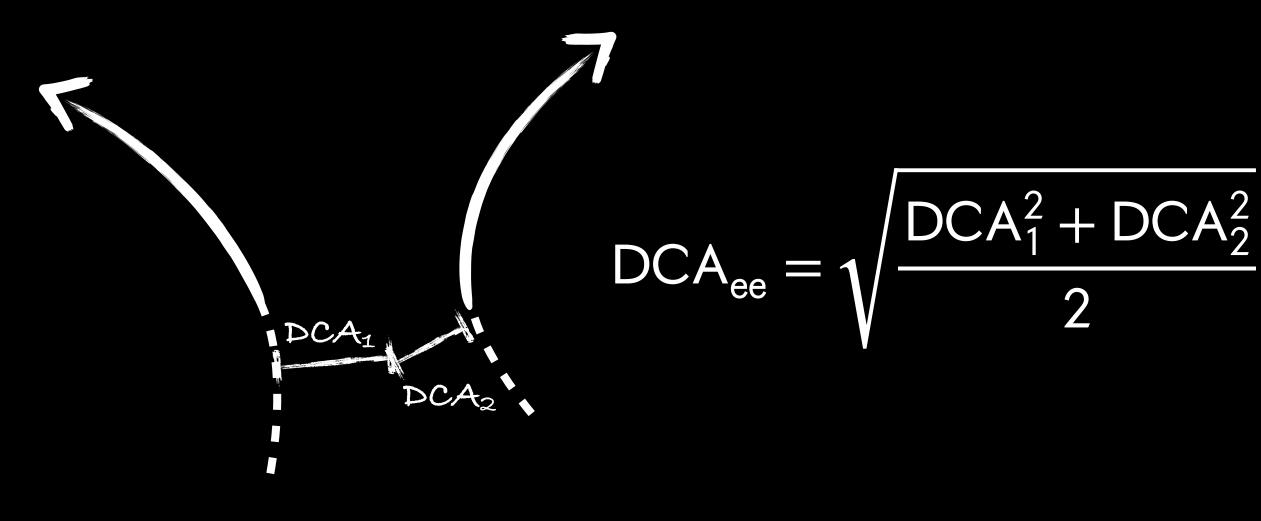
Talk by Jochen Klein





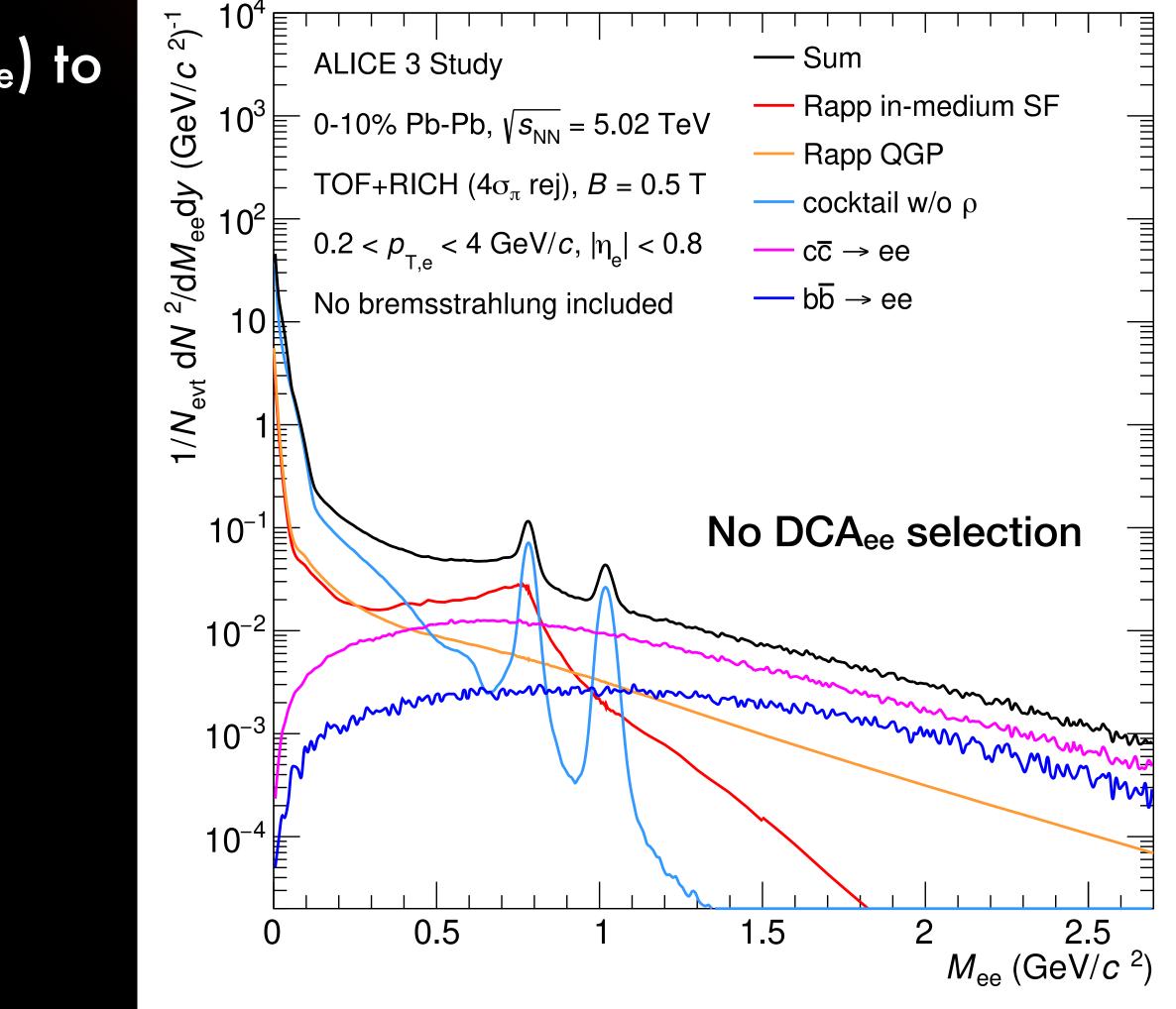
## Dielectrons

Use pair distance of closest approach (DCA<sub>ee</sub>) to reject heavy-flavour backgrounds



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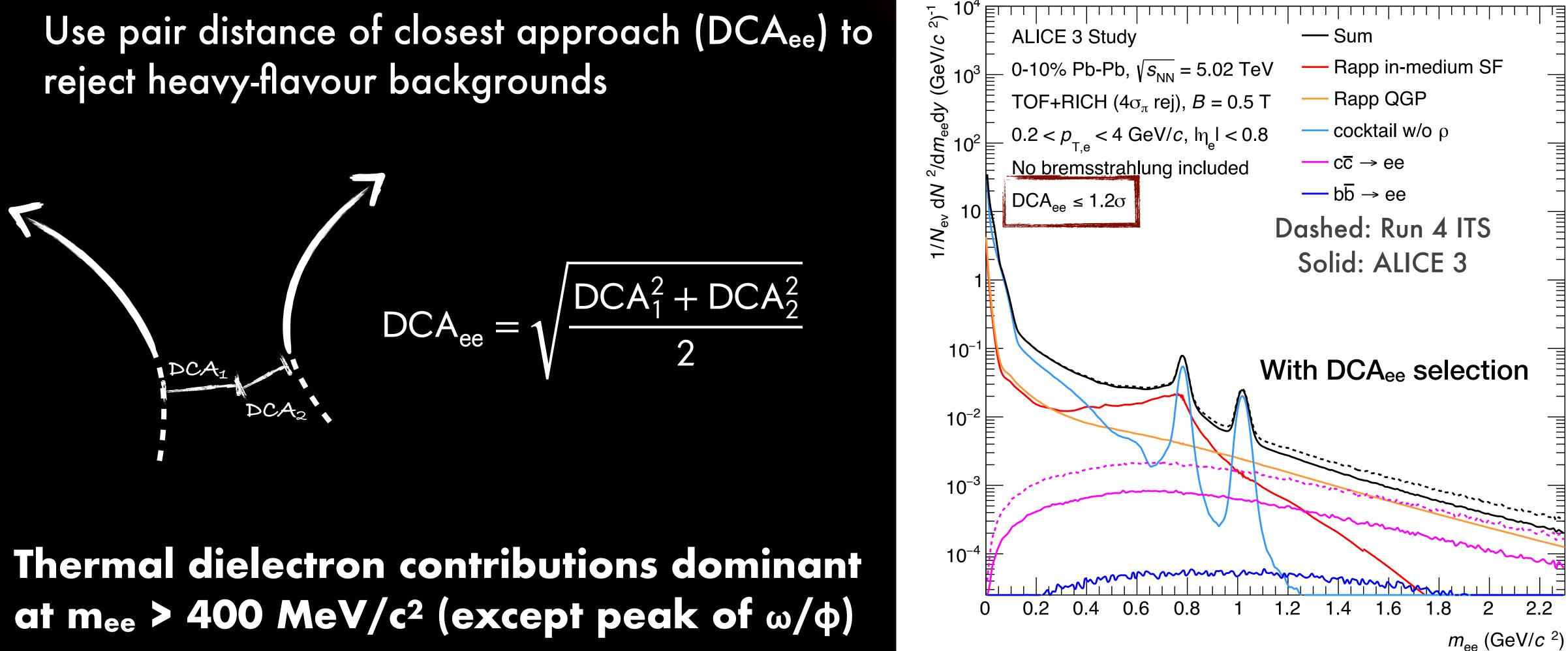
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ALI-SIMUL-498024

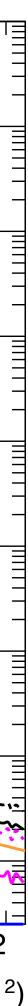


## Dielectrons



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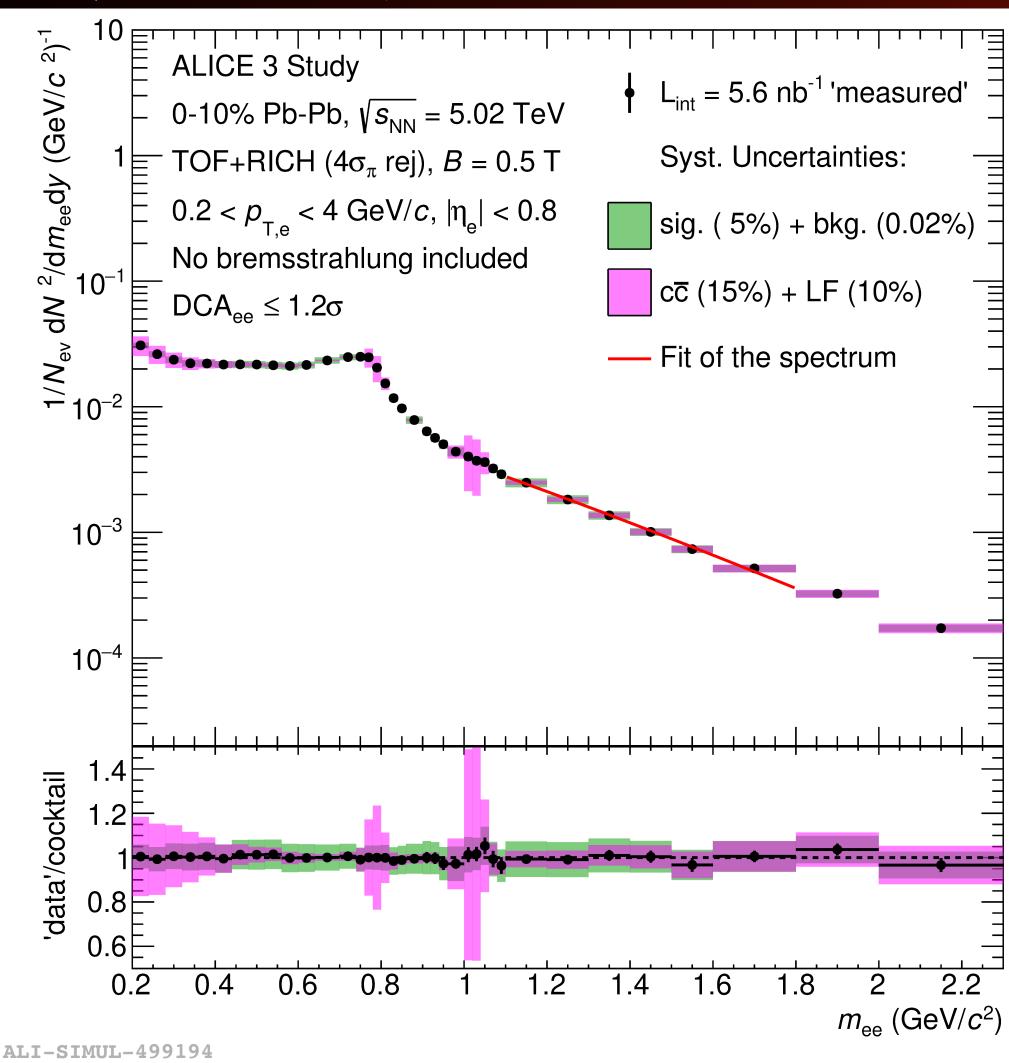
## Excess yield The Temperature fit

Thermal spectrum from QGP and hadron gas after subtraction of hadronic components

Temperature estimation via exponential fit  $\propto (m_{ee}T)^{\frac{3}{2}}exp(-m_{ee}/T)$ 

In region dominated by QGP radiation

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## Excess yield The Temperature fits

Differential analysis in p<sub>T</sub> gives access to early phase of the collision

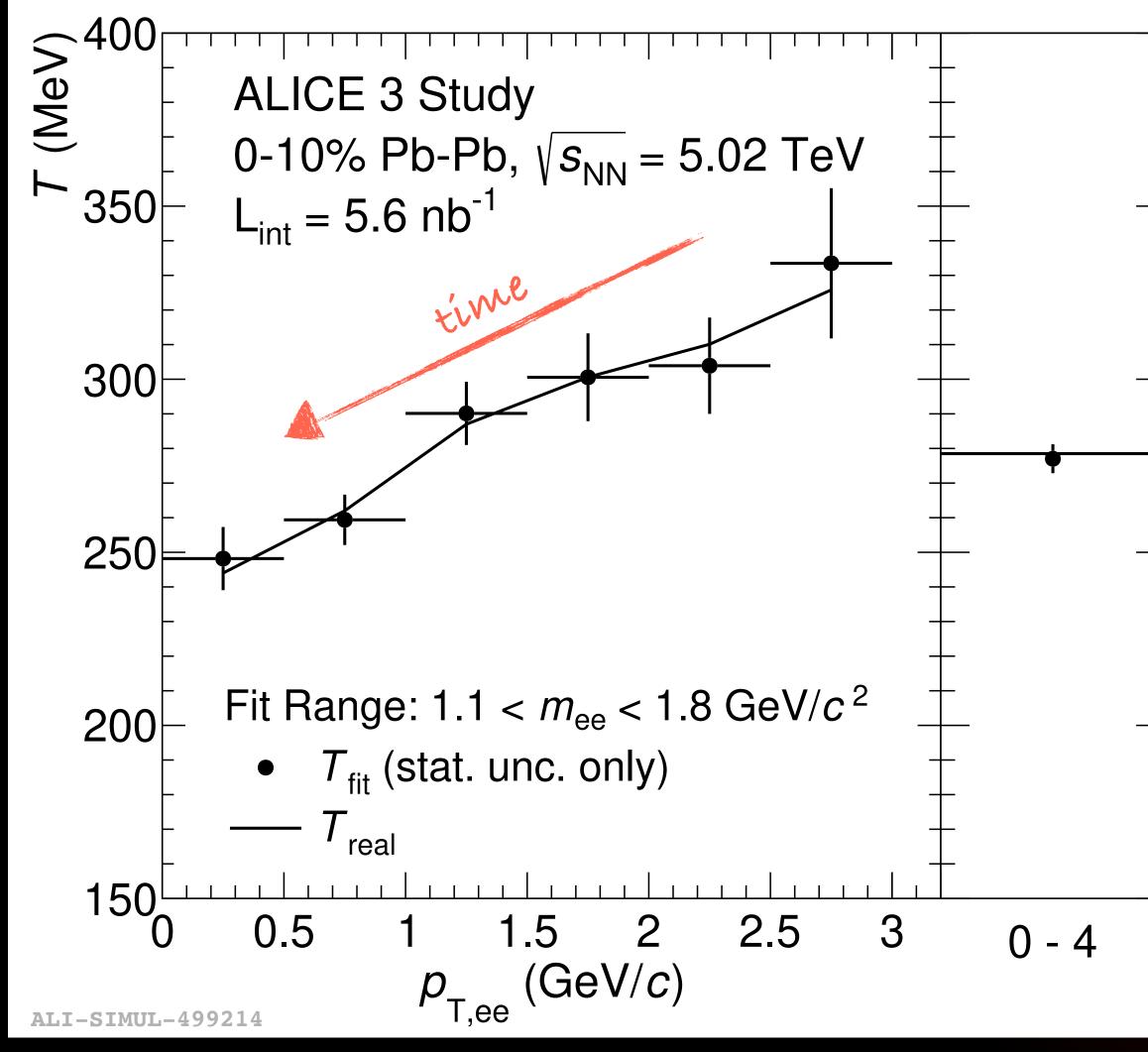
## Probe time evolution of the QGP

Complementary measurement to photon spectrum

Integrated: ±1.5% stat. ±2% syst. uncertainty (ITS 3:  $\pm 4\%$  stat.  $^{+4\%}_{-10\%}$  syst. uncertainty)

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### **CERN-LHCC-2022-009**



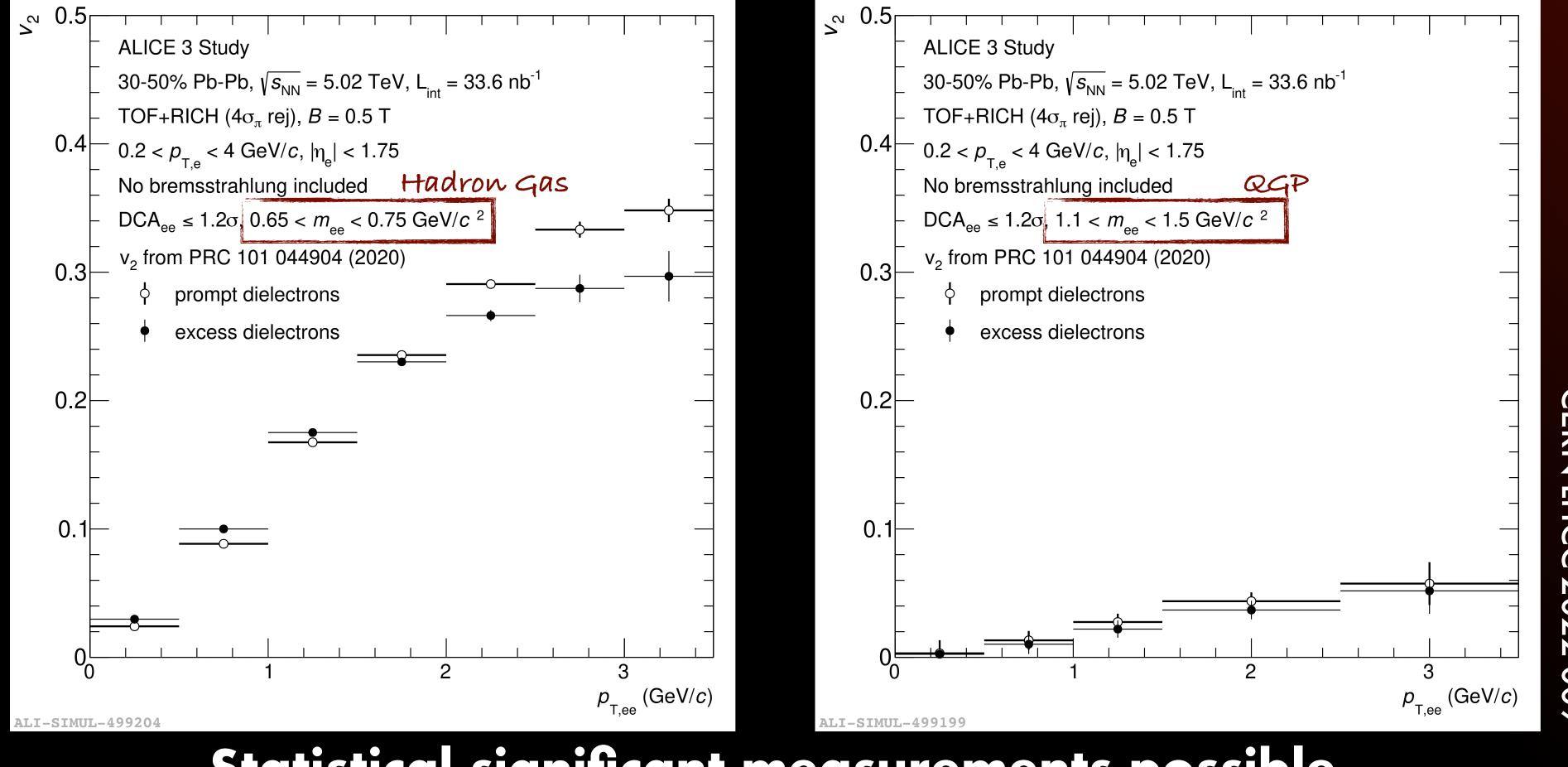






# Elliptic flow

Projection for dielectron v<sub>2</sub> in mass regions with dominant thermal contributions



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### Statistical significant measurements possible



## Electric conductivity A first look

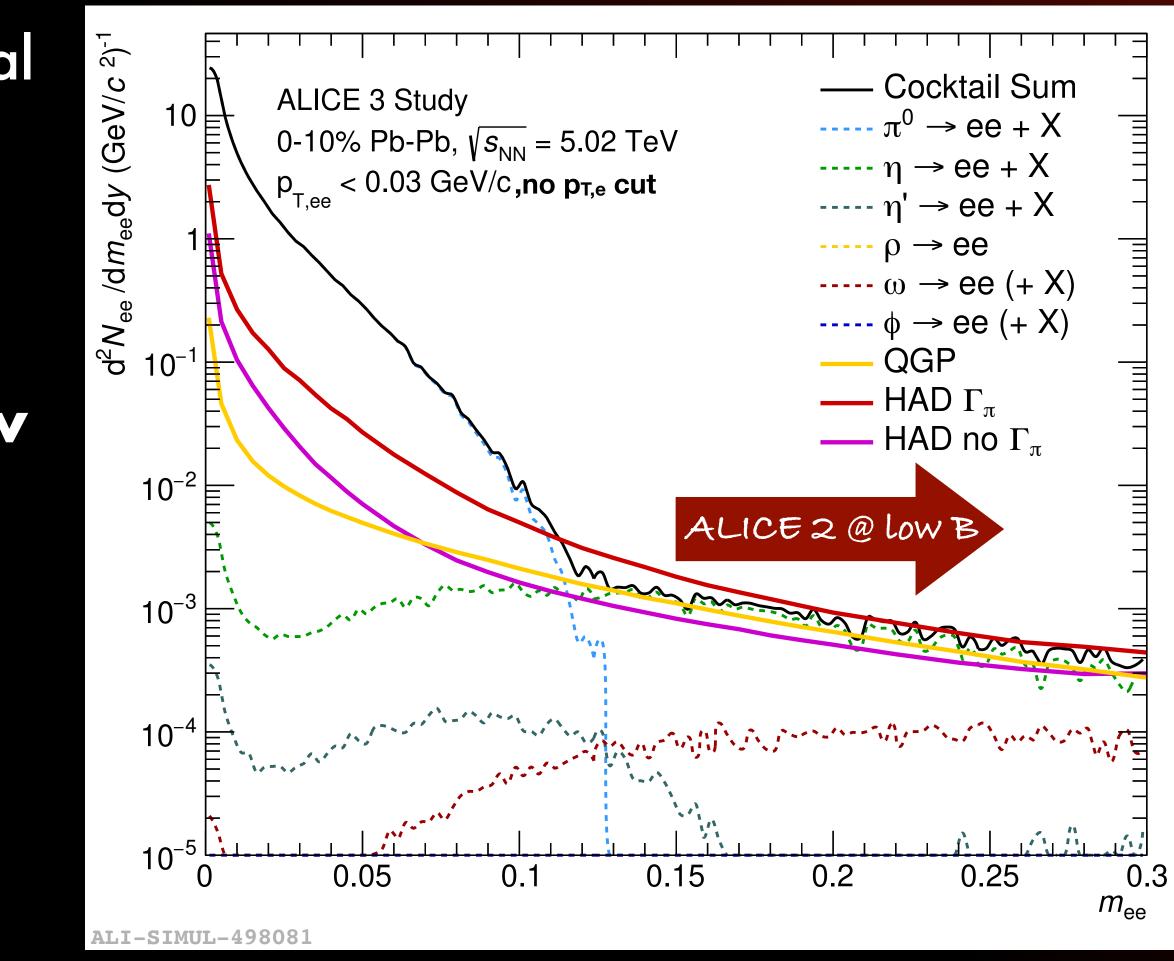
Very low pair pt selection enhances thermal contributions wrt hadronic cocktail

Width of thermal contribution sensitive to electric conductivity

Acceptance at very low p<sub>T,ee</sub> and low m<sub>ee</sub> needed

Crucial: precise measurement of  $\eta$  ( $\pi^0$ ) to smallest pt

What are other backgrounds?



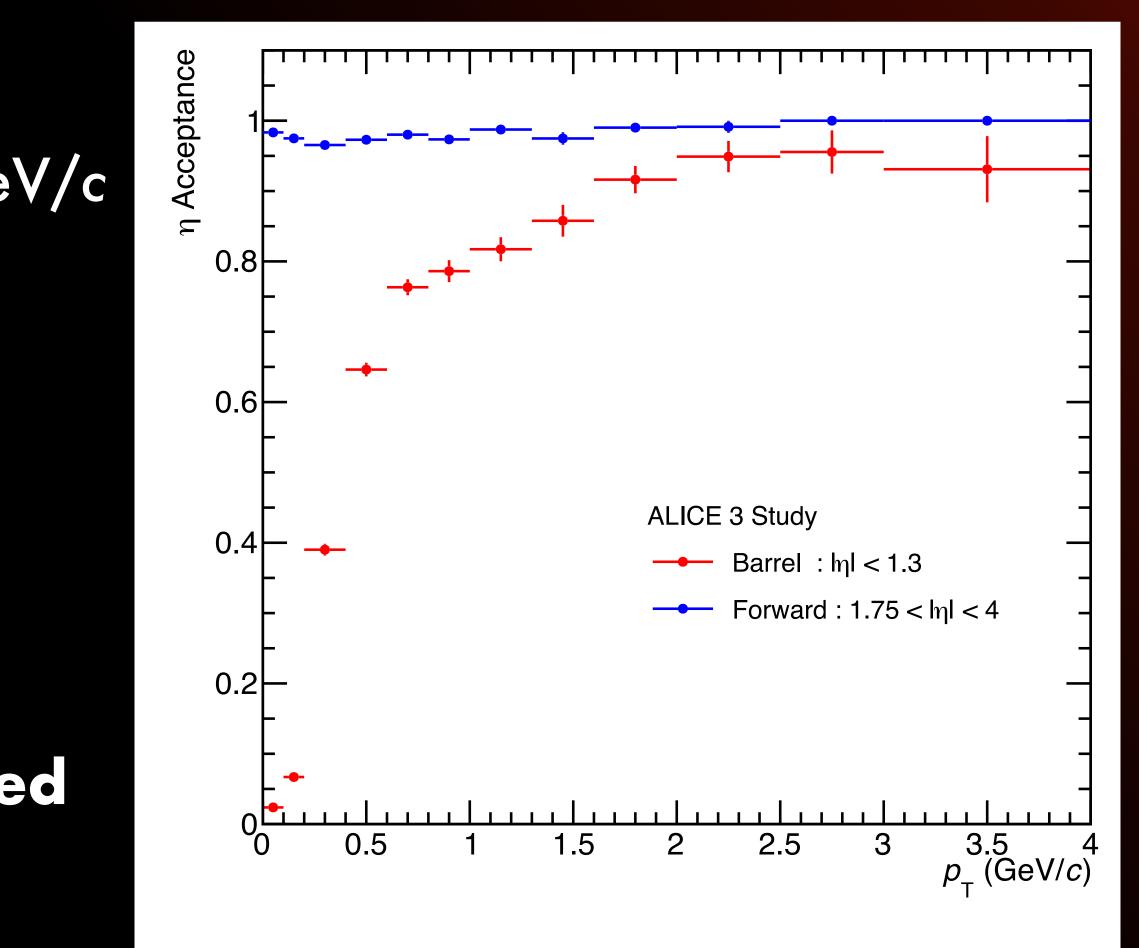


## meson acceptance

- Mid rapidity:
- Photons are required to have  $p_T > 0.1 \text{ GeV/c}$ Limited to low pt
- Forward:
- Photon p > 0.1 Gev/c
- No limitation at low pt

Drastic reduction of uncertainties of direct photon measurements expected



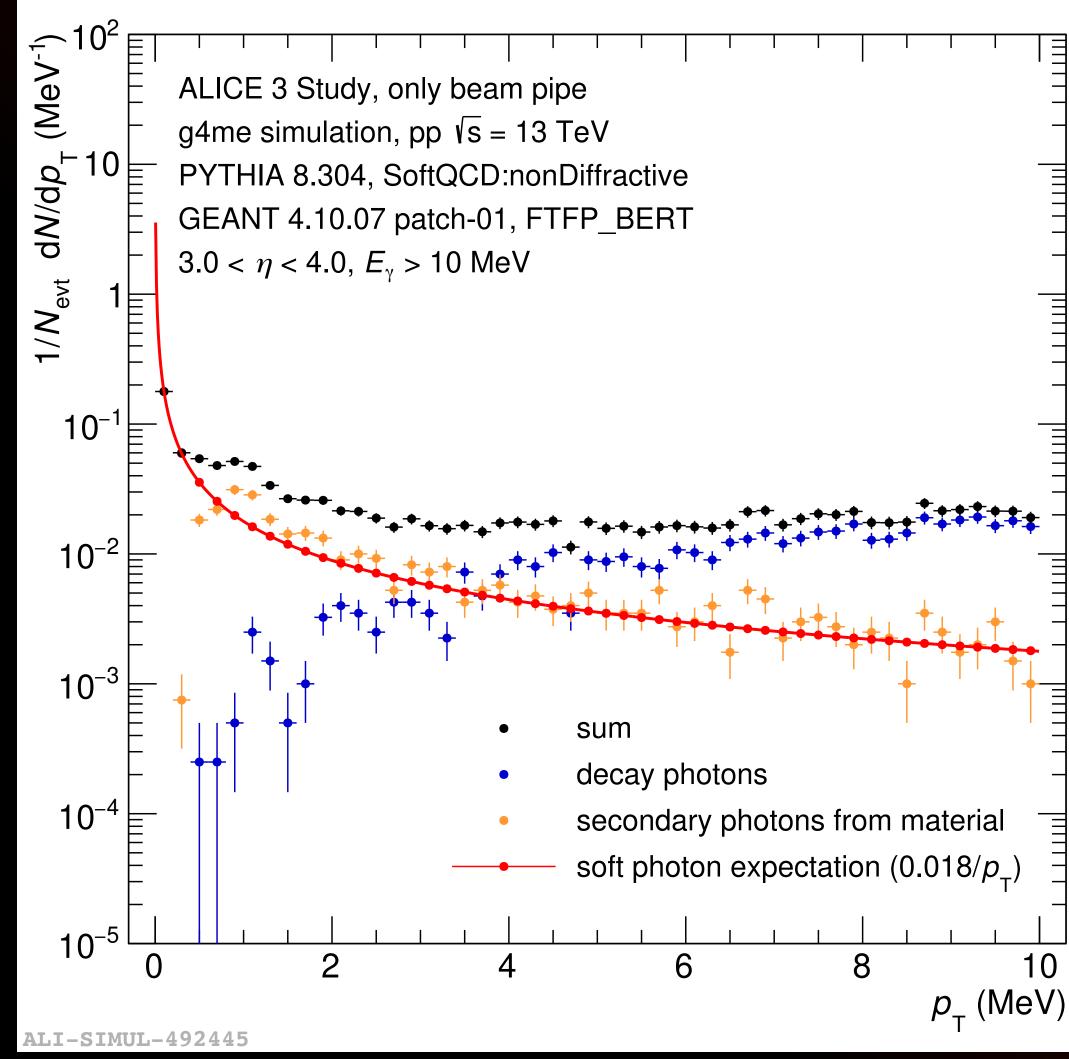


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## Soft photons Signal and Background

- At low pt bremsstrahlung dominant over decay photons
- Electron veto reduces contribution from material
- Additional modification of beam pipe possible to reduce material contribution







# Conclusion

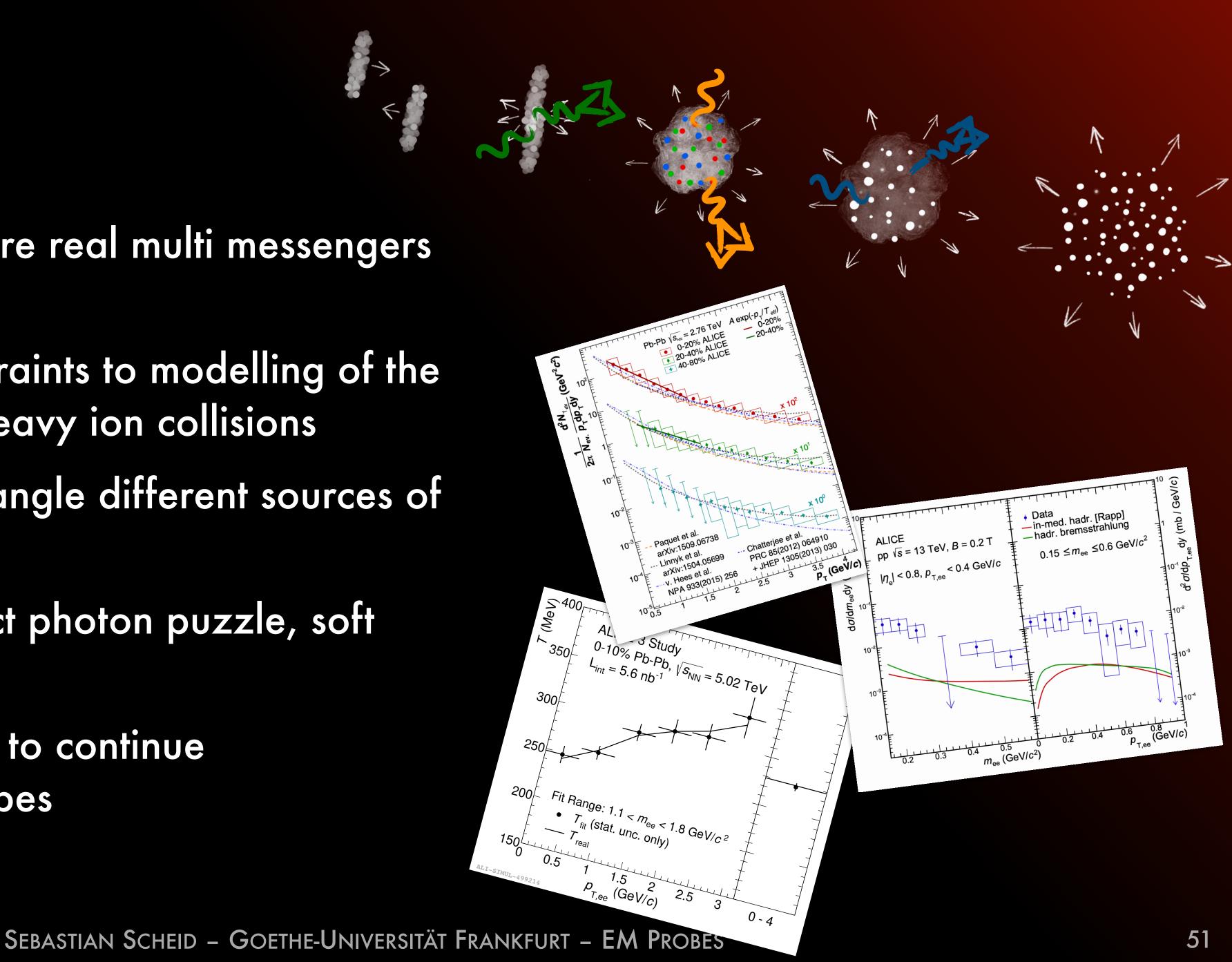
Electromagnetic probes are real multi messengers of heavy ion collisions

Measurements give constraints to modelling of the space time evolution of heavy ion collisions

 Necessary to disentangle different sources of EM probes

Still open questions (direct photon puzzle, soft photons)

Compelling physics cases to continue measurements of EM probes



Backup

## ALICE 3 A selection of topics

## Precision measurements of dileptons

- mechanisms of chiral symmetry restoration in the quark-gluon plasma
- evolution of the quark-gluon plasma

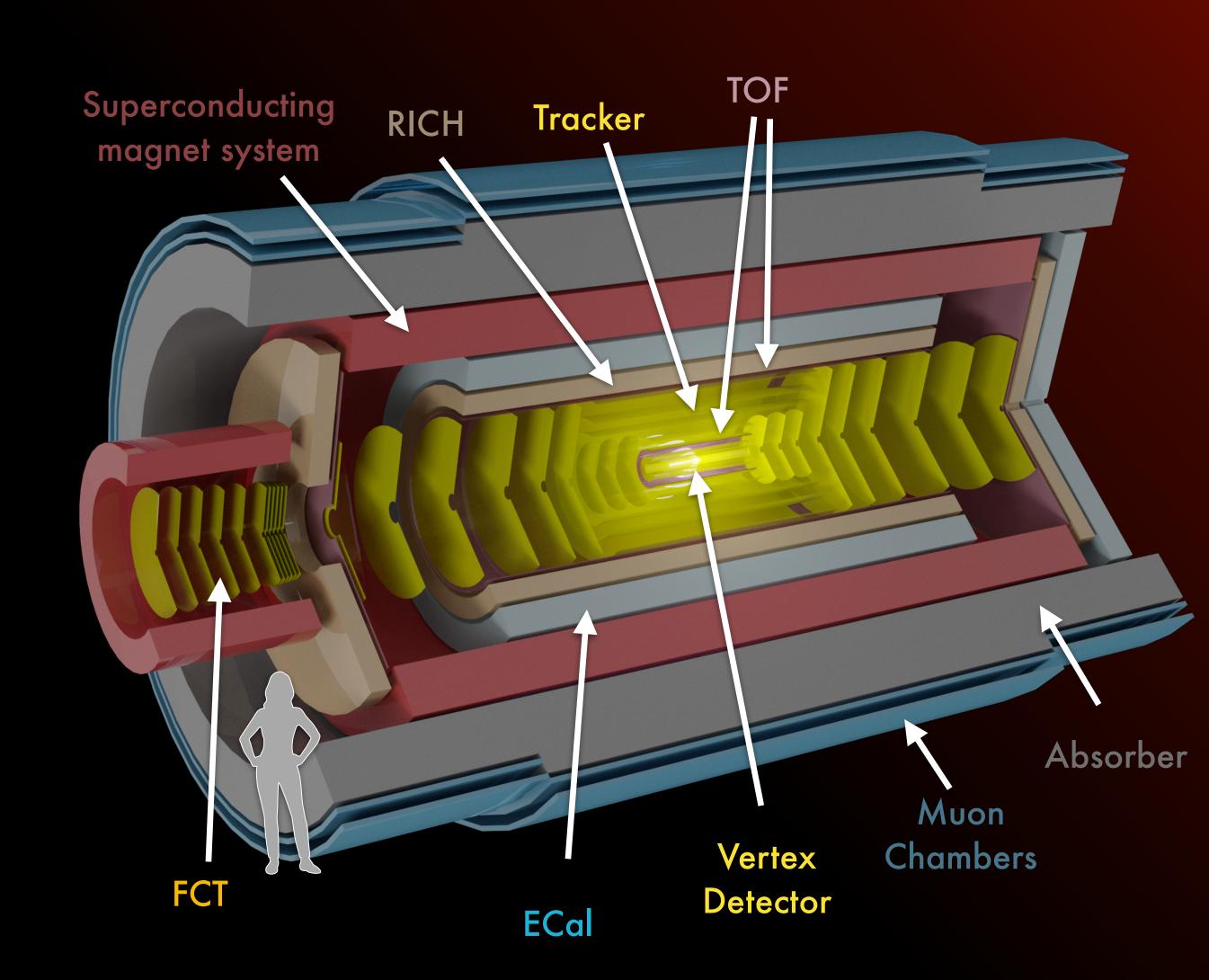
## Systematic measurements of (multi-)heavy-flavour hadrons

- transport properties in the quarkgluon plasma
- mechanisms of hadronisation from the quark-gluon plasma

Hadron correlations

**Super Soft Photon Conversions** 

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## Soft photons Motivation

At very low pt: y produced mainly via inner Bremsstrahlung

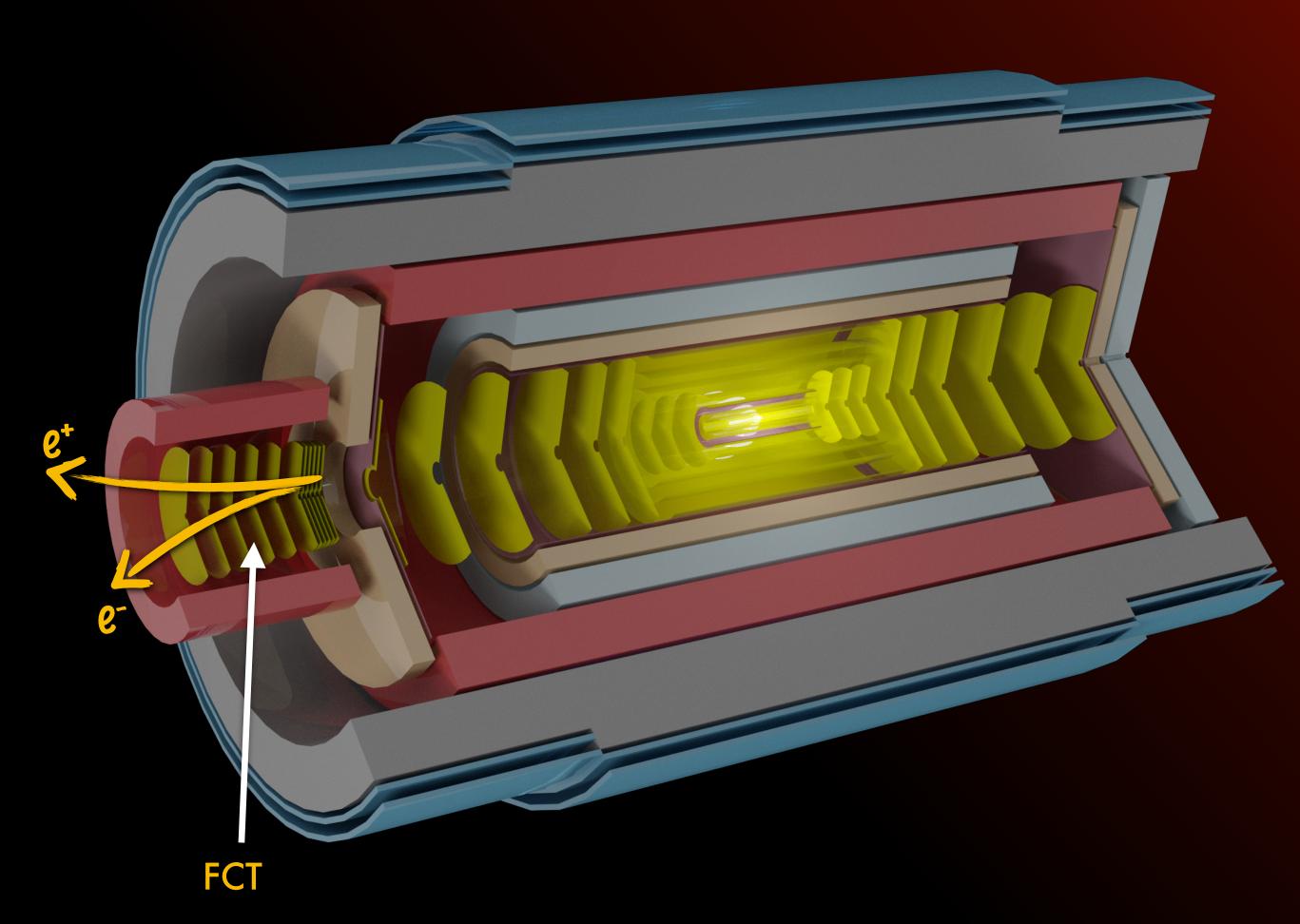
Can be calculated via Low theorem if inand out-going particles can be measured

Previous measurements:

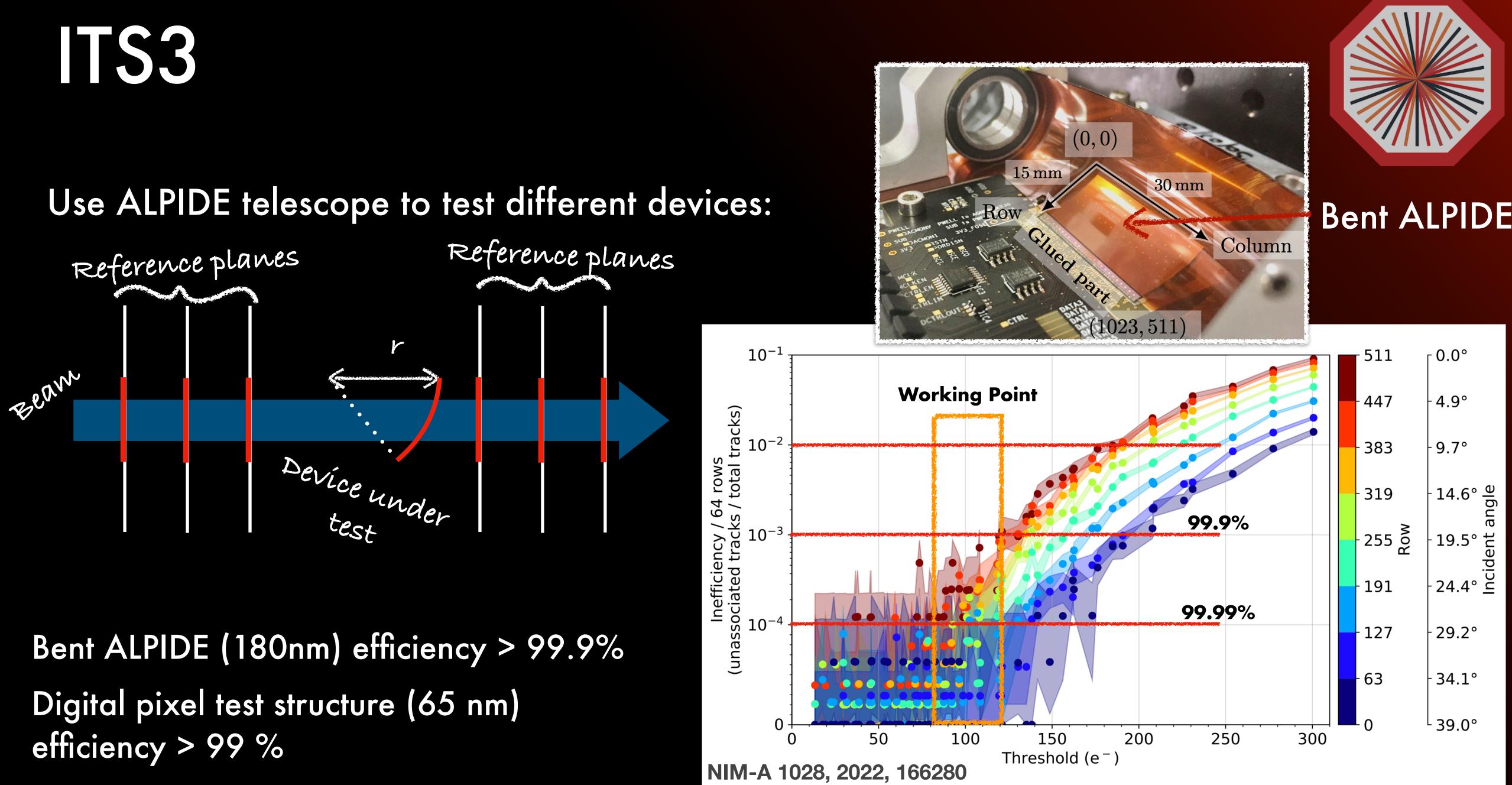
Factor 5 difference

 $\rightarrow$  no discrepancies seen in ee $\rightarrow$ µµY

## To measure lowest $p_T \rightarrow$ measure in forward direction

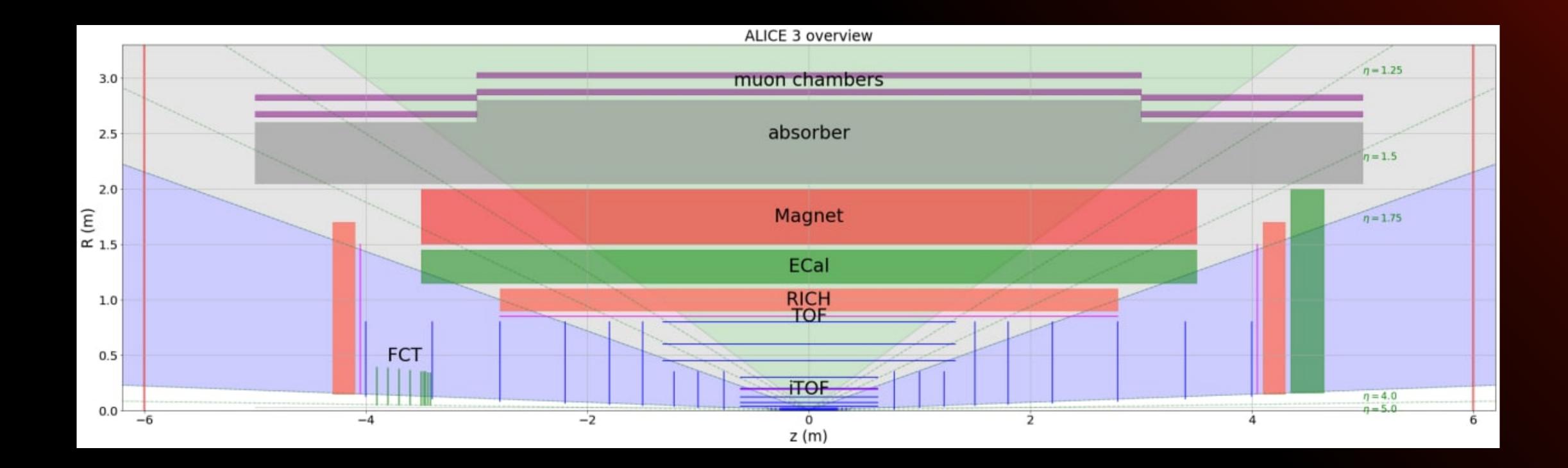








# A possible ALICE 3 setup



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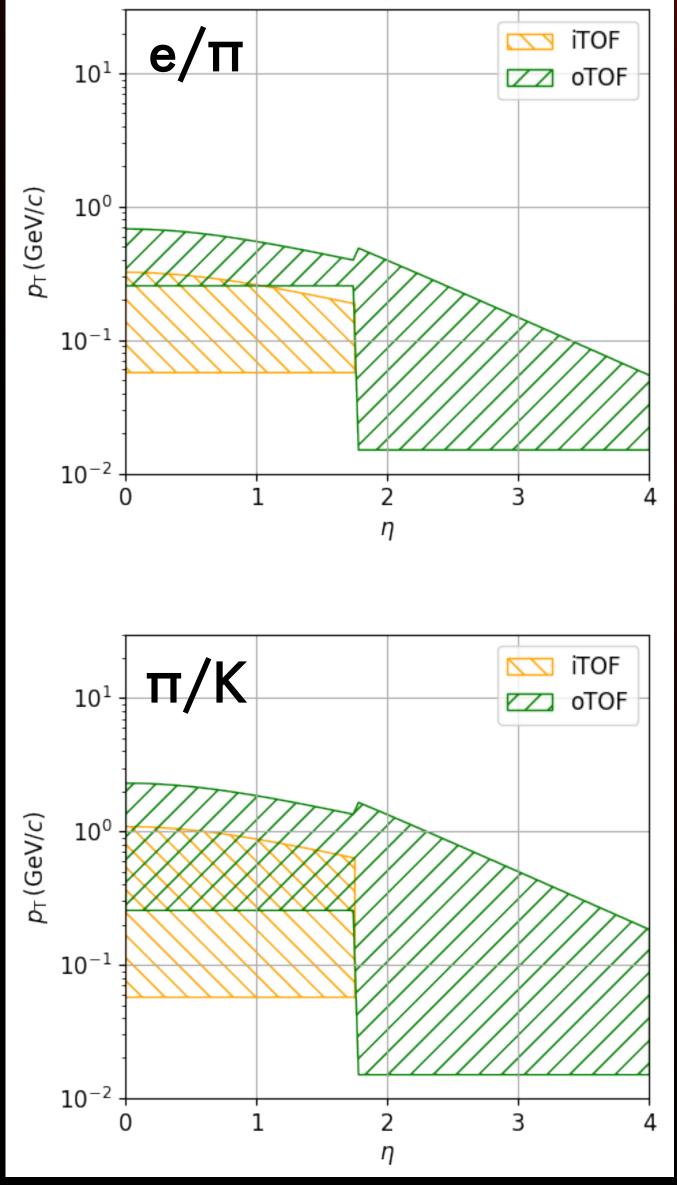


## ALICE 3 Particle identification TOF

Separation power  $\propto \sigma_{\rm tof}$ 

- distance and time resolution crucial
- larger radius results in lower p<sub>T</sub> bound
- 2 barrel + 1 forward TOF layers
  - outer TOF at  $R \approx 85$  cm
  - inner TOF at  $R \approx 19$  cm
  - forward TOF at  $z \approx 405$  cm
- Silicon timing sensors ( $\sigma_{TOF} \approx 20 \text{ ps}$ )

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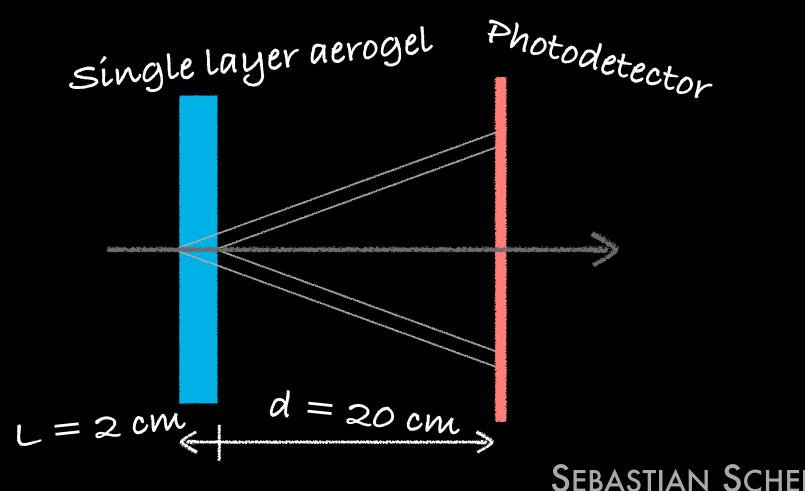


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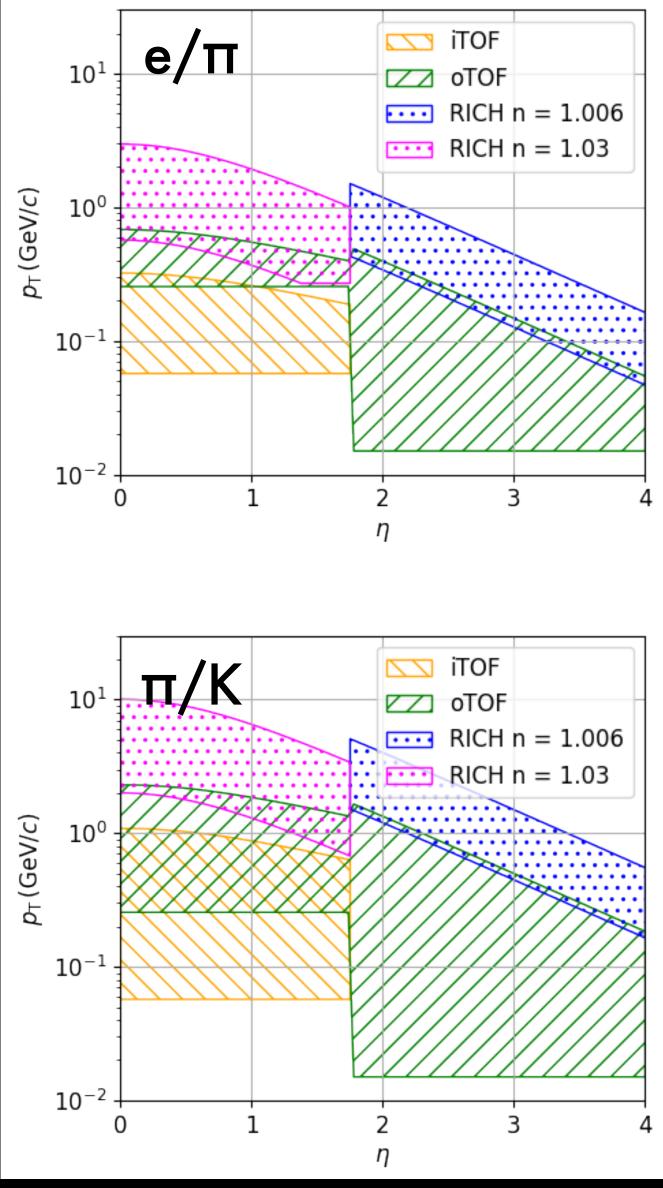
# ALICE 3

Particle identification

- Extend PID reach of outer TOF to higher pt with Cherenkov
  - aerogel radiator to ensure continuous coverage from TOF  $\rightarrow$  refractive index n = 1.03 (barrel)  $\rightarrow$  refractive index n = 1.006 (forward)



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# Excess yield

Thermal spectrum from QGP and hadron gas after subtraction of hadronic components

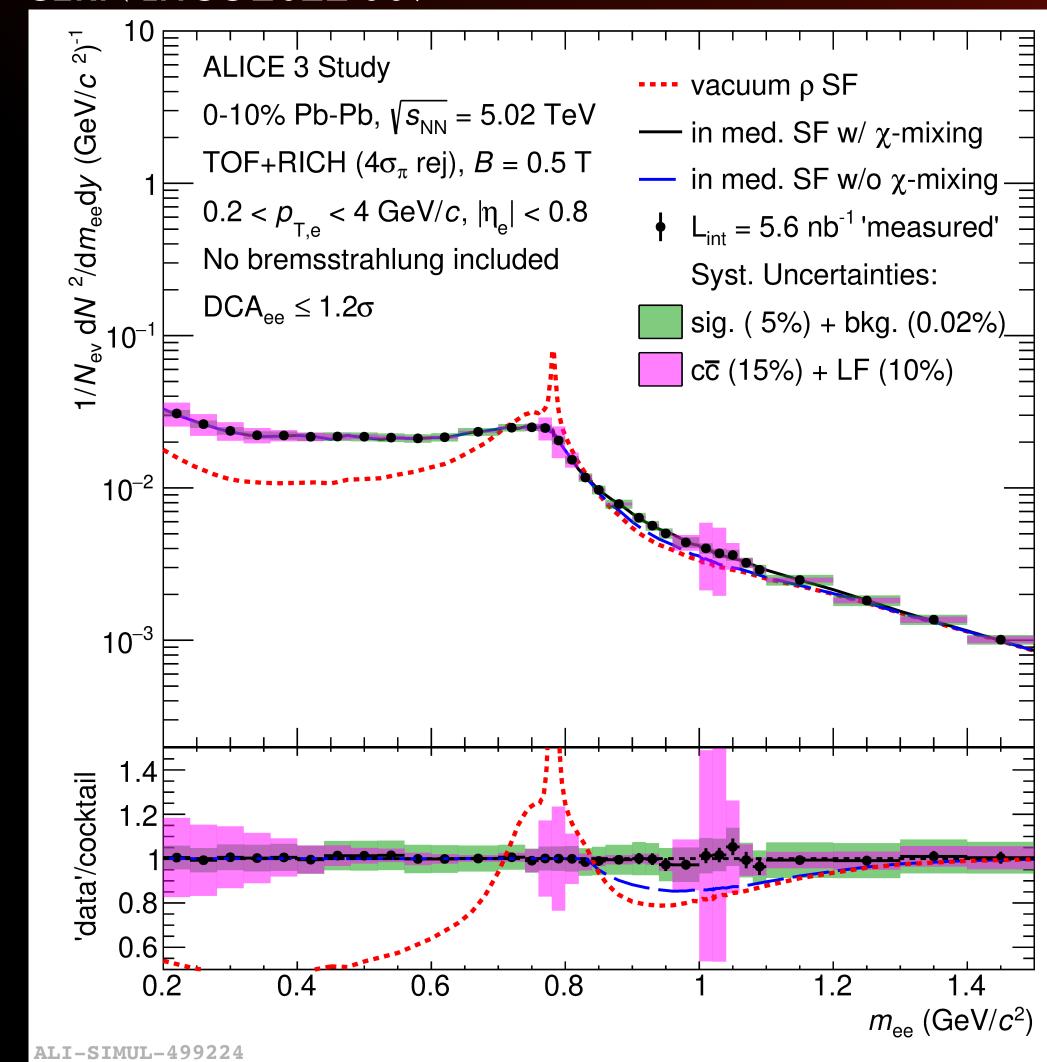
Compared to different  $\rho$  spectral functions:

- Including p-a1 mixing (reference)
- Without mixing of p and a1
- Vacuum p spectral function

Dominated by systematic uncertainties

Precision of data sufficient to measure 15% effect in  $\rho$ -a<sub>1</sub> mixing region

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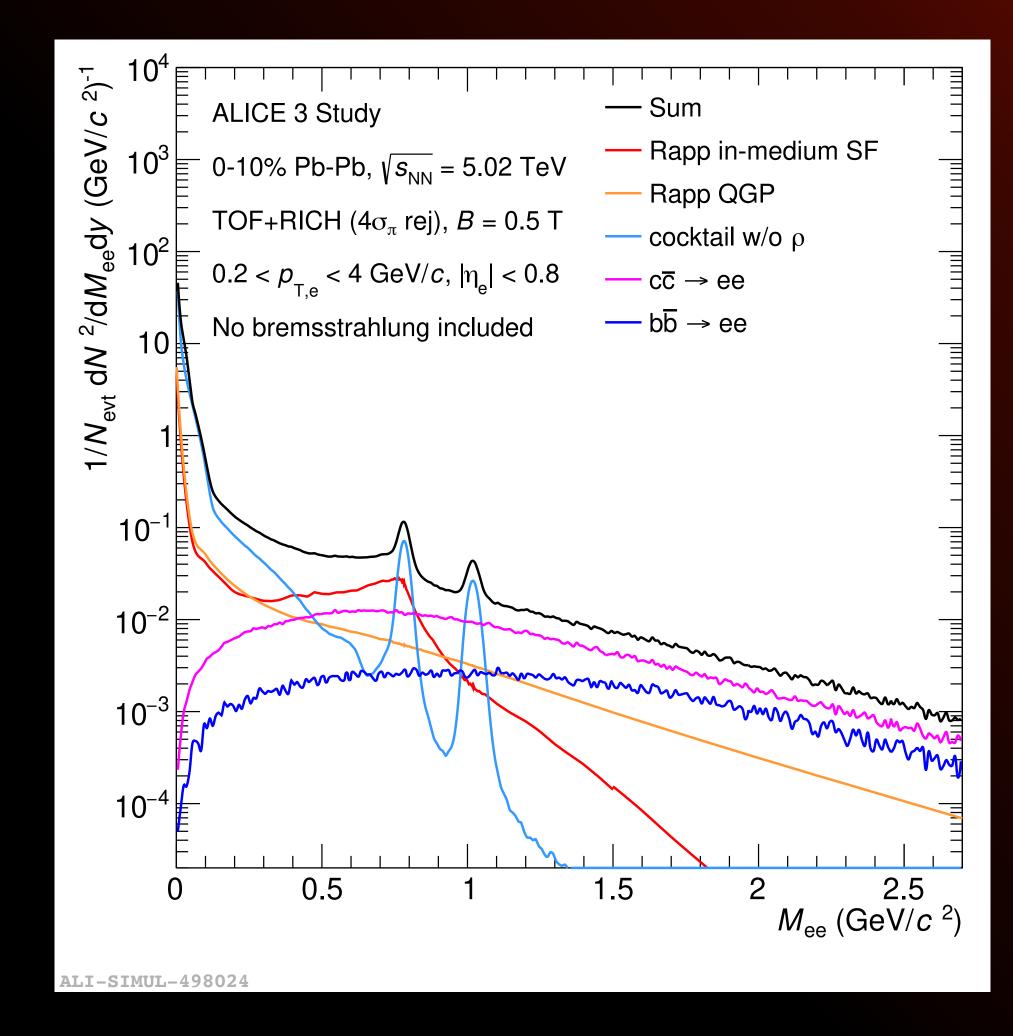
# Dielectron cocktail

Cocktail ingredients:

- Light flavour: Measurements of particle spectra
- Correlated HF pairs: Measurement in  $pp + N_{coll}$  scaling + CNM (EPSO9)
- Thermal radiation: Calculations by Ralf Rapp

Large charm contribution of wide mass range

 $\rightarrow$  dominant m<sub>ee</sub> > 0.9 GeV/c<sup>2</sup>





# Distance of Closest Approach

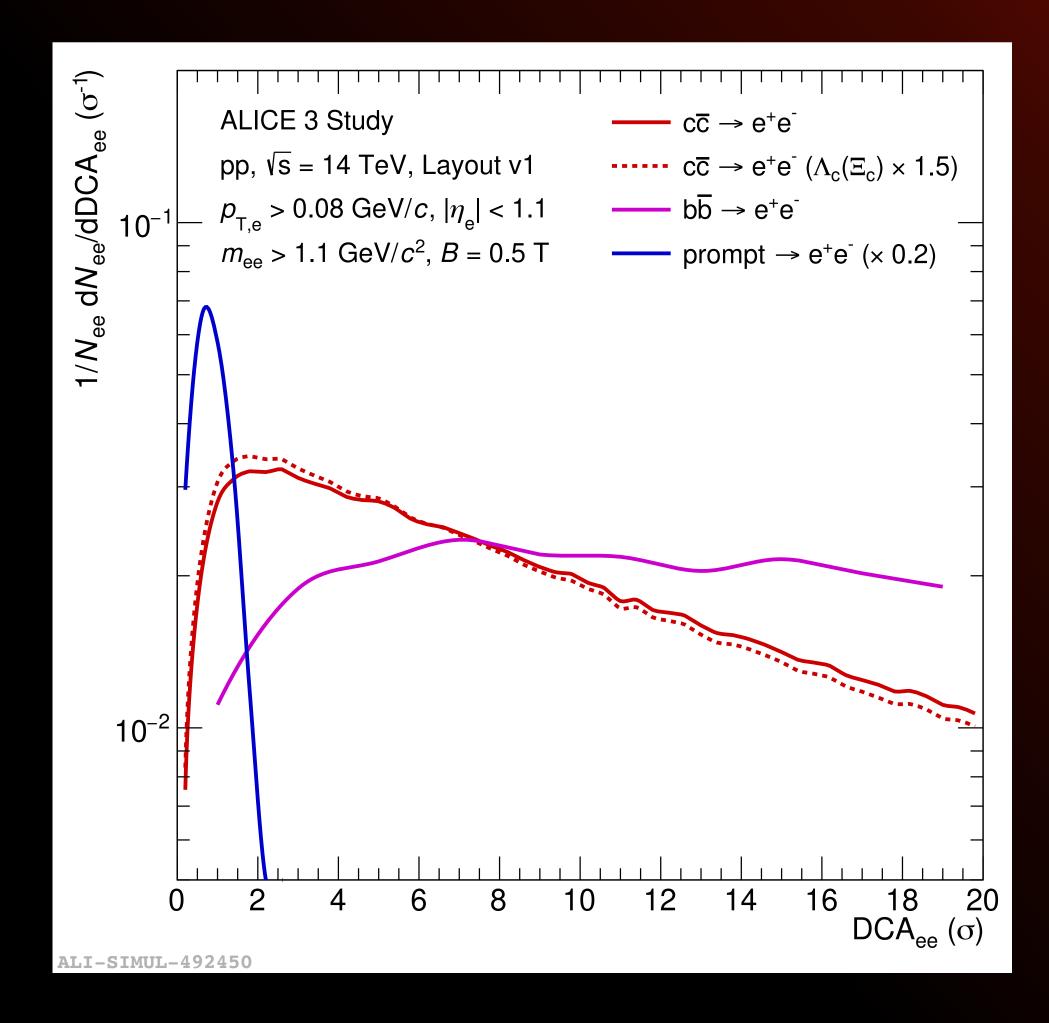
DCA<sub>ee</sub> templates for prompt, charm and beauty contributions

Ordering by decay length of mother particles as expected:

Prompt < Charm < Beauty</pre>

Changing hadron chemistry in the charm sector can change the shape of the template

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## Distance of Closest Approach Charm hadron contributions

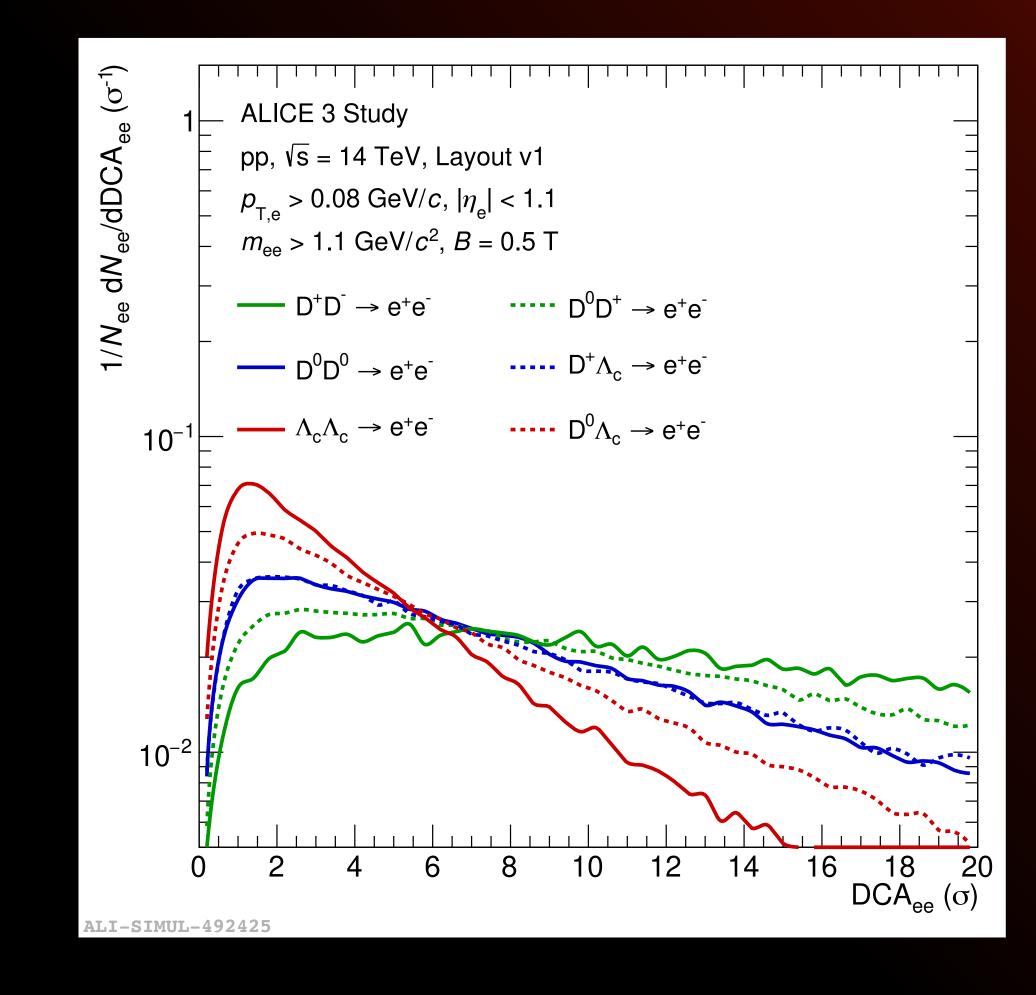
Sensitivity can be demonstrated on the level of different decay length of charm hadron species

- ст ( $\Lambda_c$ ) = 60 $\mu$ m
- ст ( $D^{0}$ ) = 150 µm
- $cT (D^{+}) = 300 \mu m$

Final template for charm from weighted sum of single contributions

Input are the measured branching ratios and fragmentation functions

Estimation of uncertainty underinvetigation



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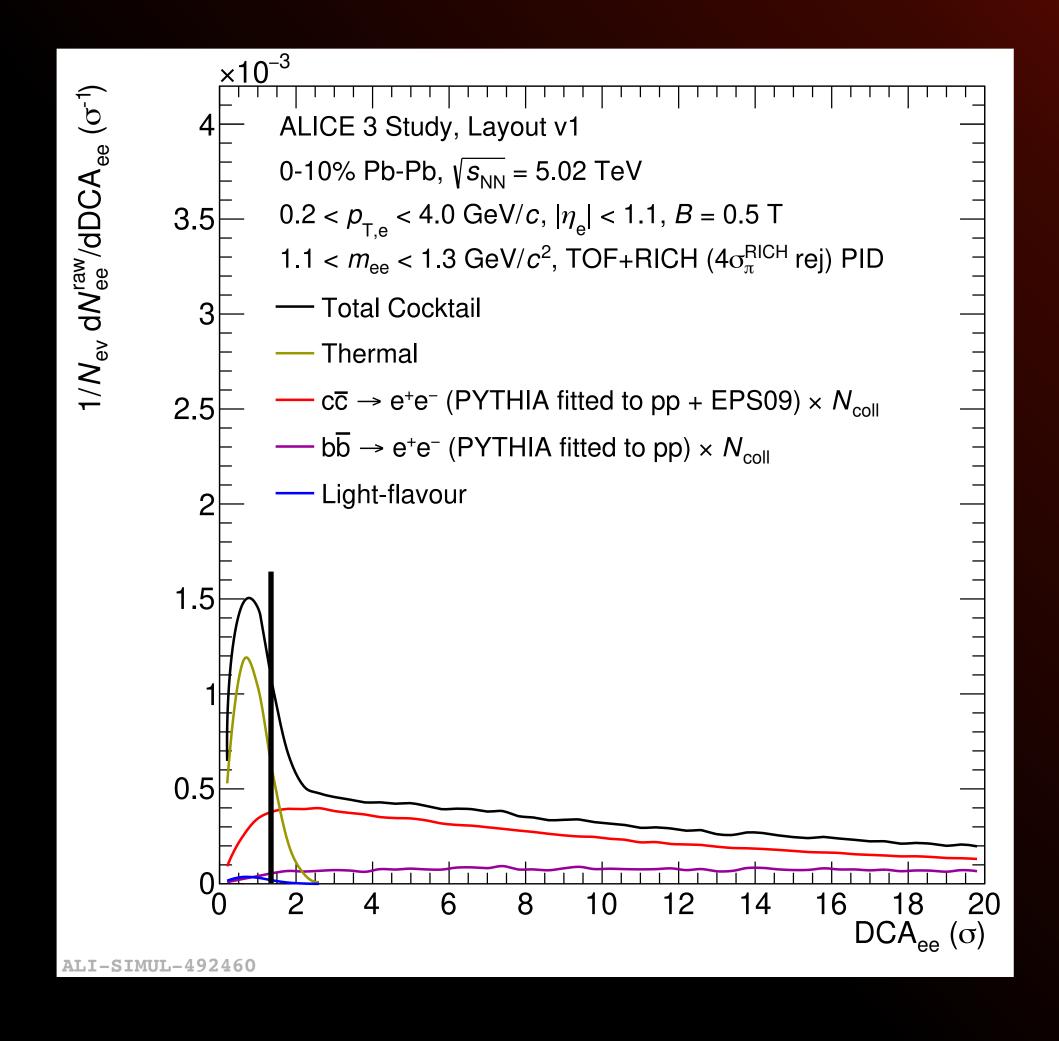


# Distance of Closest Approach

Expected raw DCA<sub>ee</sub> contribution

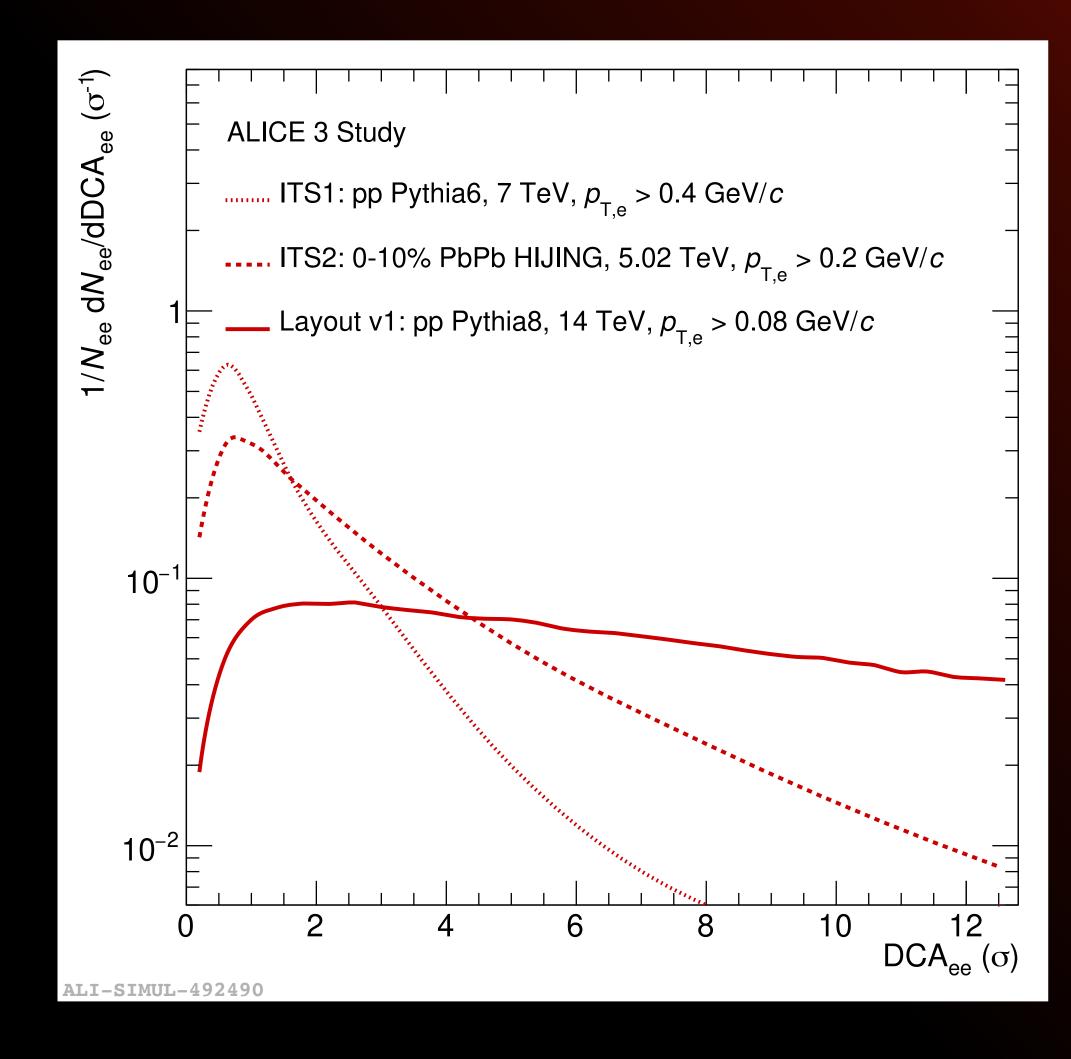
Good separation of prompt and HF contributions

Selection at  $1.2\sigma$  rejects: 27% prompt, 94% charm, and 98% of beauty ITS 3: 27% prompt, 83% charm





## Distance of Closest Approach Comparison with RUN1/3 performance





# Eliptic flow

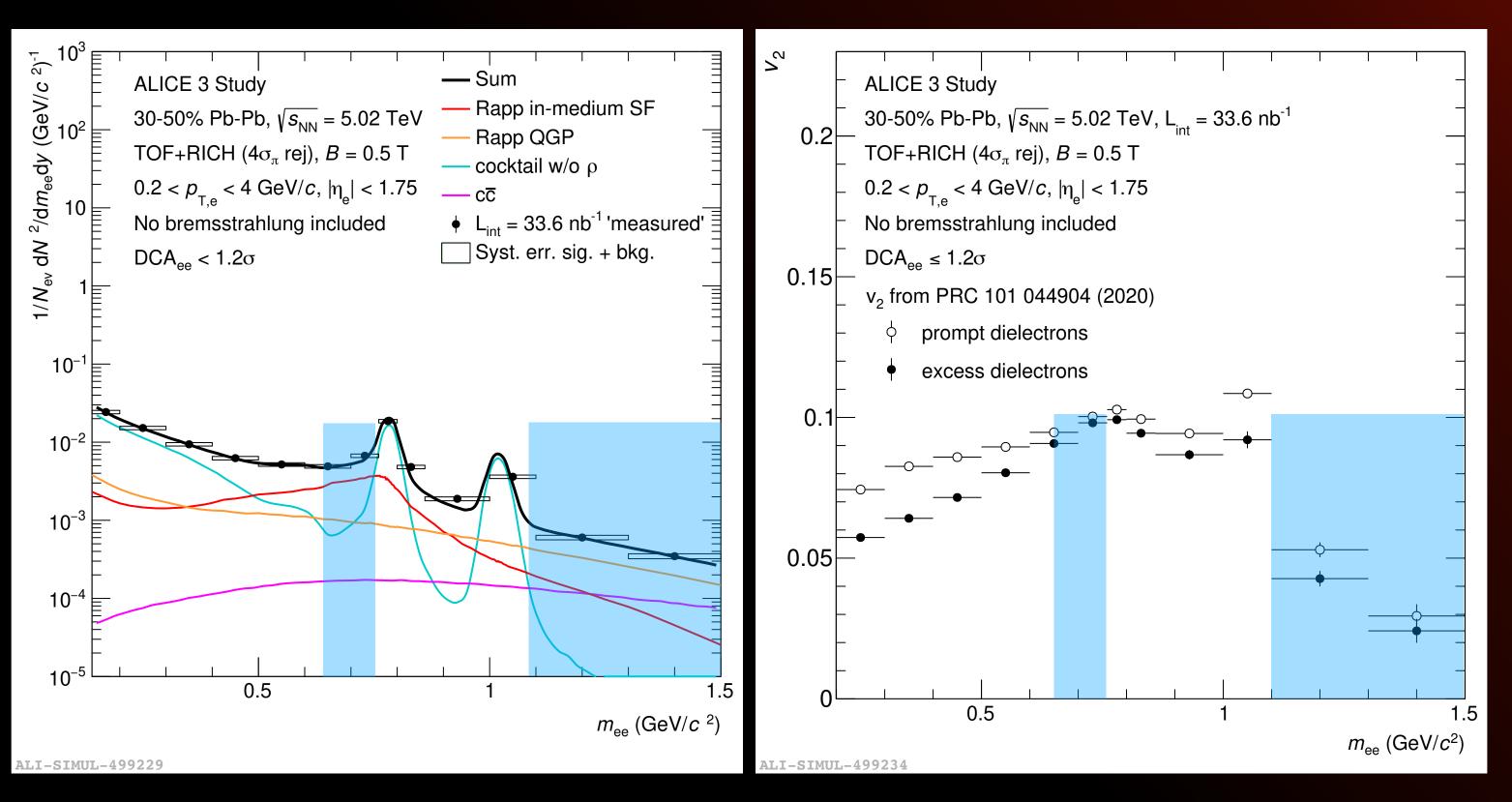
Yield and elliptic flow  $v_2$  in semi-central (30-50%) Pb-Pb collisions for 6 years

Important measurement together with photon  $v_2$ 

Elliptic flow (hadronic and thermal) based on calculations by Gojko Vojanovic et al

Differential measurement statistical possible

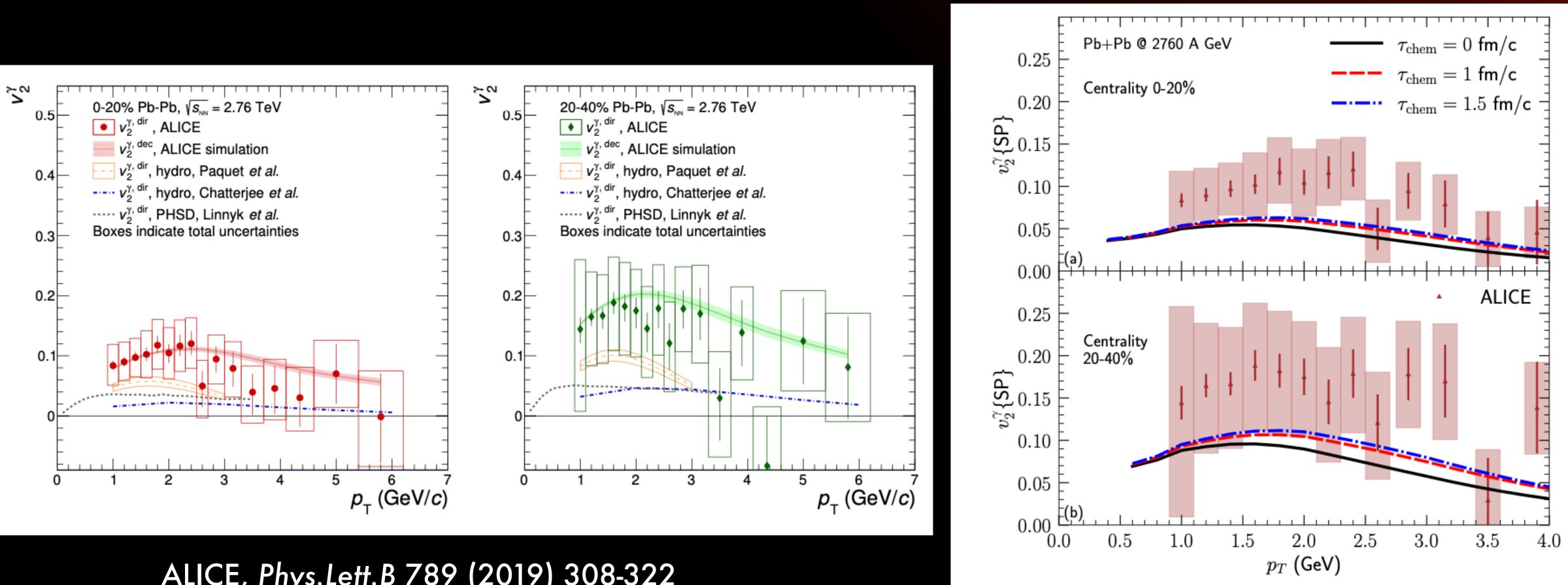
Look at regions with large contributions from hadron gas or QGP



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## Photon v2 in ALICE Subtitle



ALICE, Phys.Lett.B 789 (2019) 308-322



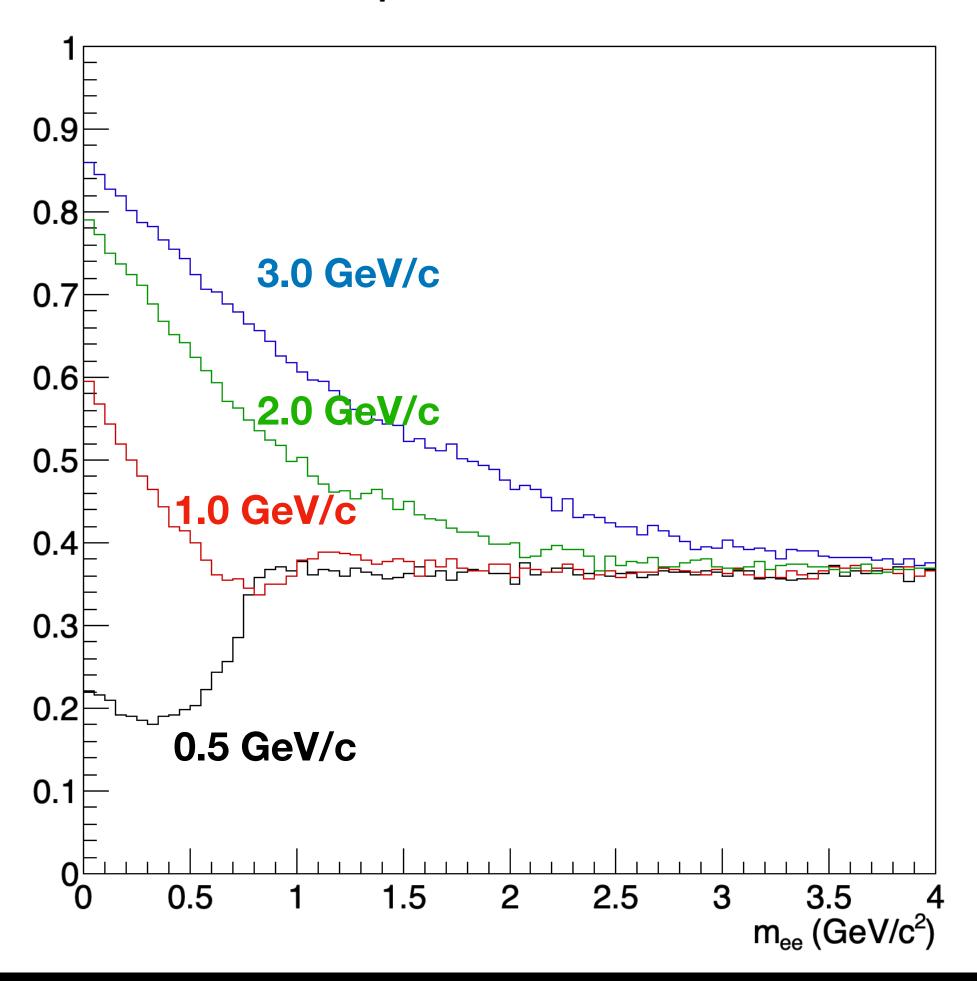
# Acceptance Correction

Generate pairs with given mass and pT (|y| < 0.8)Let decay into dielectron Reject based on kinematic cuts  $(p_T > 0.2 \text{ GeV/c}, |\eta| < 0.8)$ 

This would be applied by dividing the spectrum by the given values.

So the correction would make the spectra harder at larger pT, thus resulting in higher Temperatures.

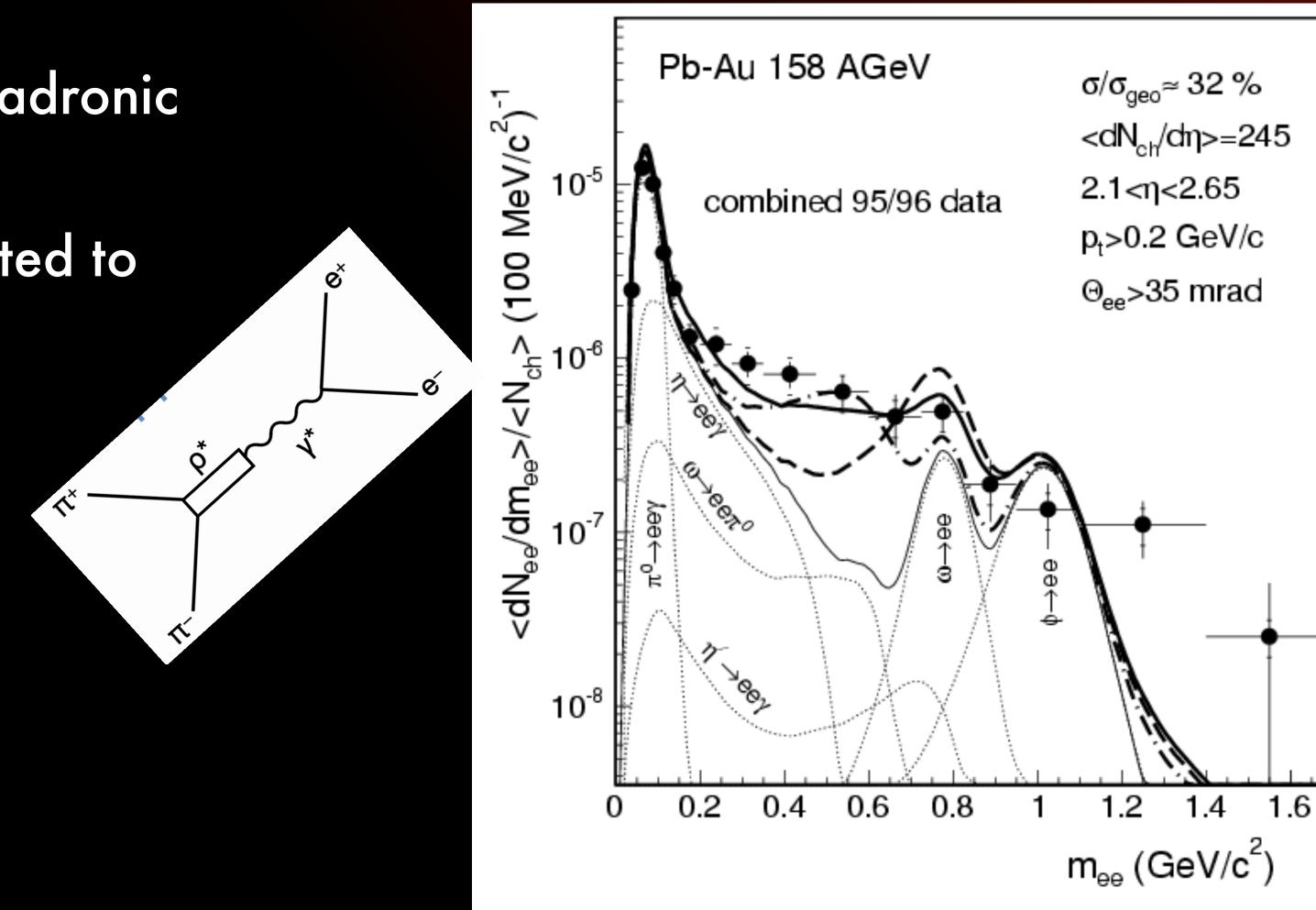
### **Acceptance Correction**





# The SPS - CERES

- Comparison of e+e- data with hadronic cocktail shows excess
- Enhancement at low mass attributed to  $\pi\pi$  annihilation?
- Spectral shape dominated by p
- Possible contributions:
- Vacuum p (just more than in pp, disfavoured)
- Medium modified p
- Dropping mass
- Melting of the spectral function





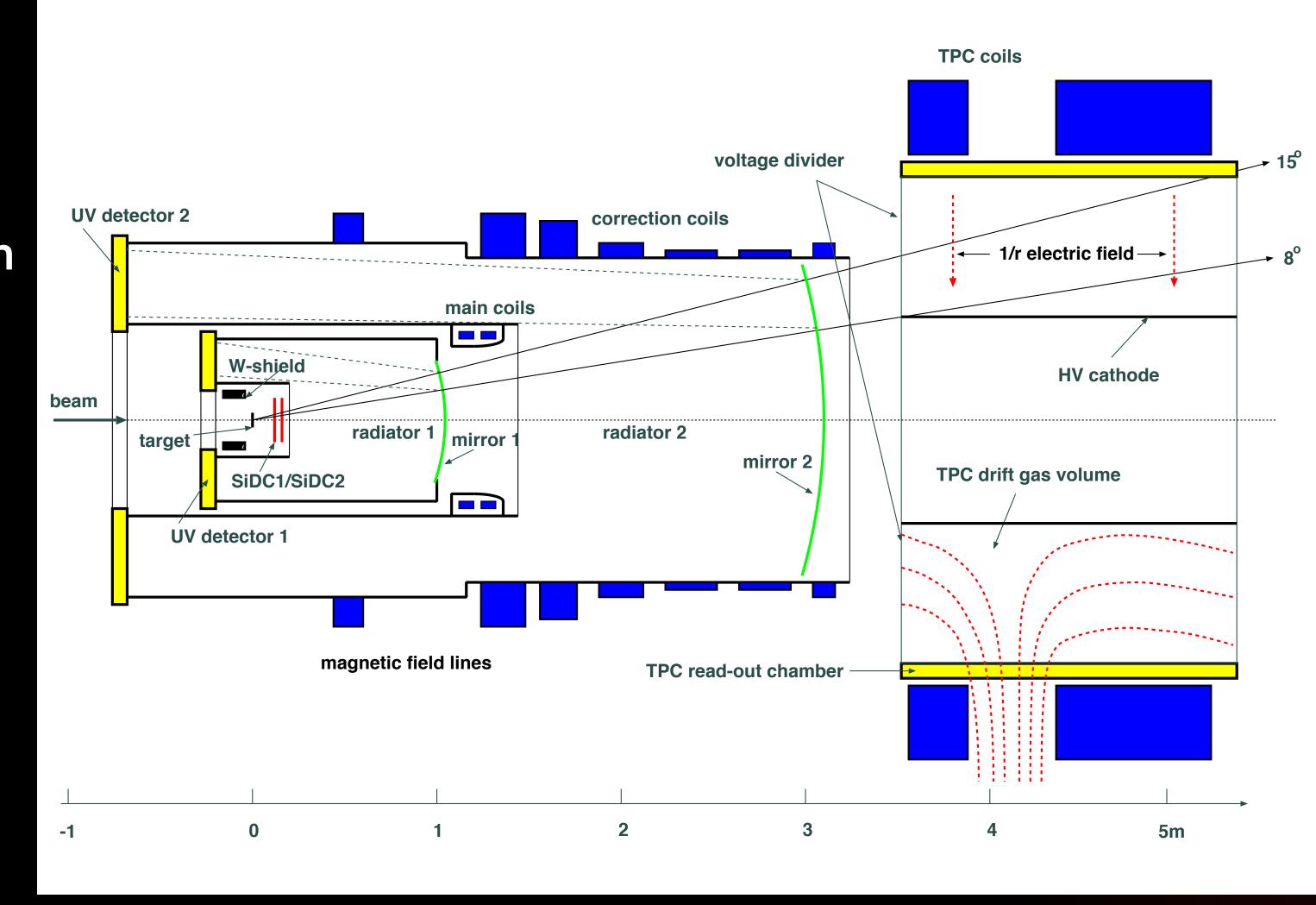


# The SPS - CERES 2.0

Upgrade of the CERES experiment with a TPC

Better in momentum/mass resolution

dE/dx gives possibility for better electron identification





# The SPS - CERES

Persistance of previous results

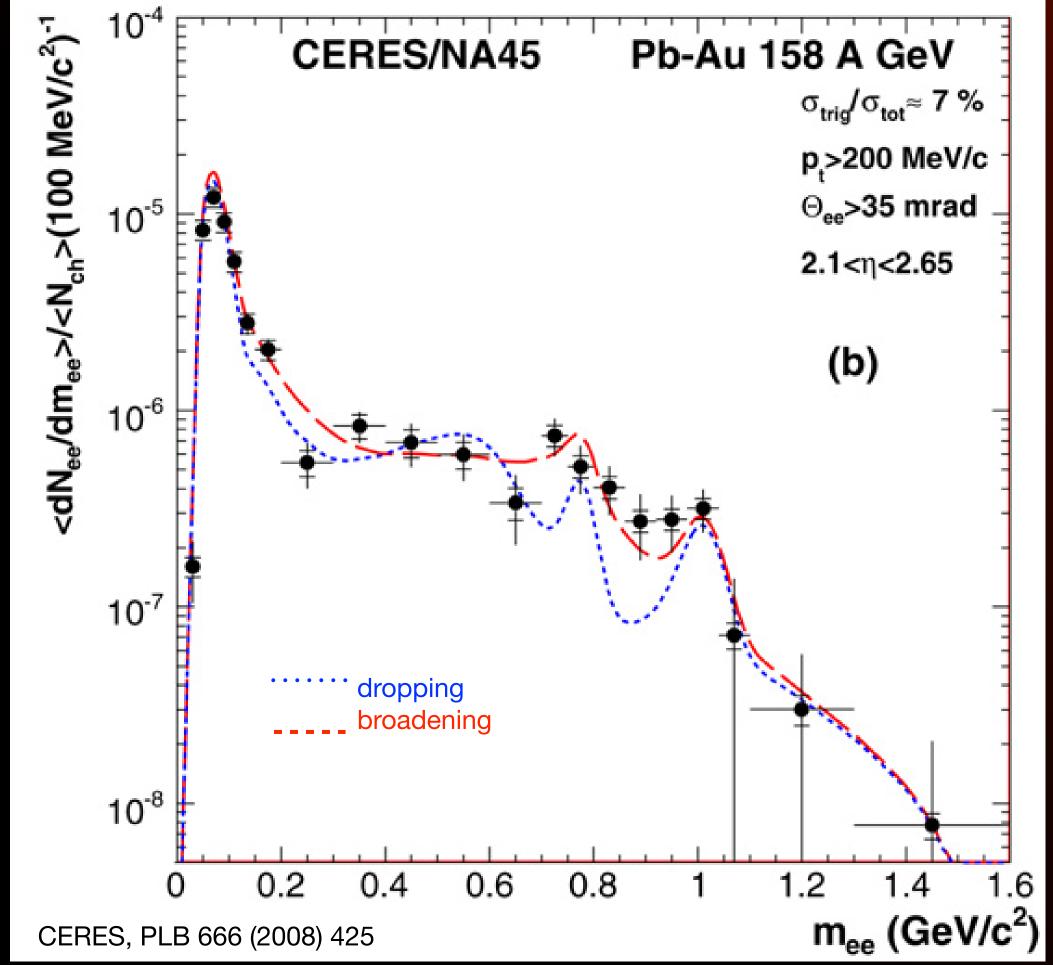
Vacuum  $\rho$  not enough to reproduce data

In medium modification:

• Broadening p spectral function (Rapp and Wambach)

• Dropping p mass (Brown et al.)

Additional mass resolution and statistical precision needed to conclude on modification



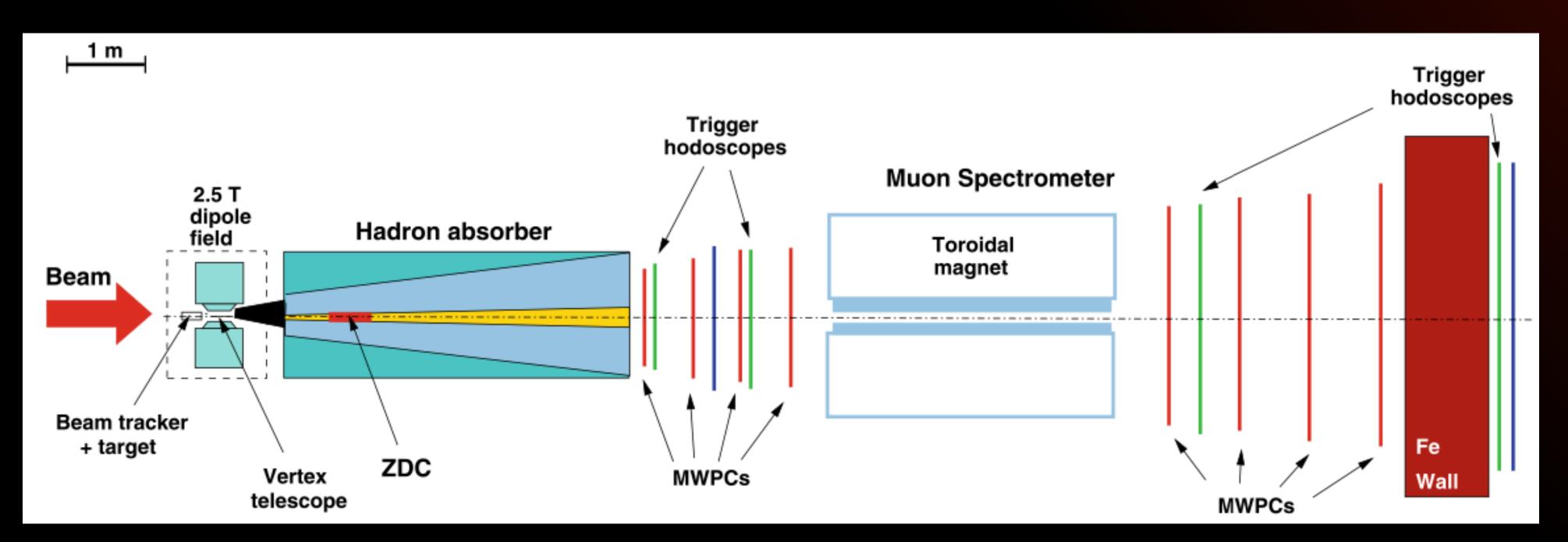
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# The SPS - NA60

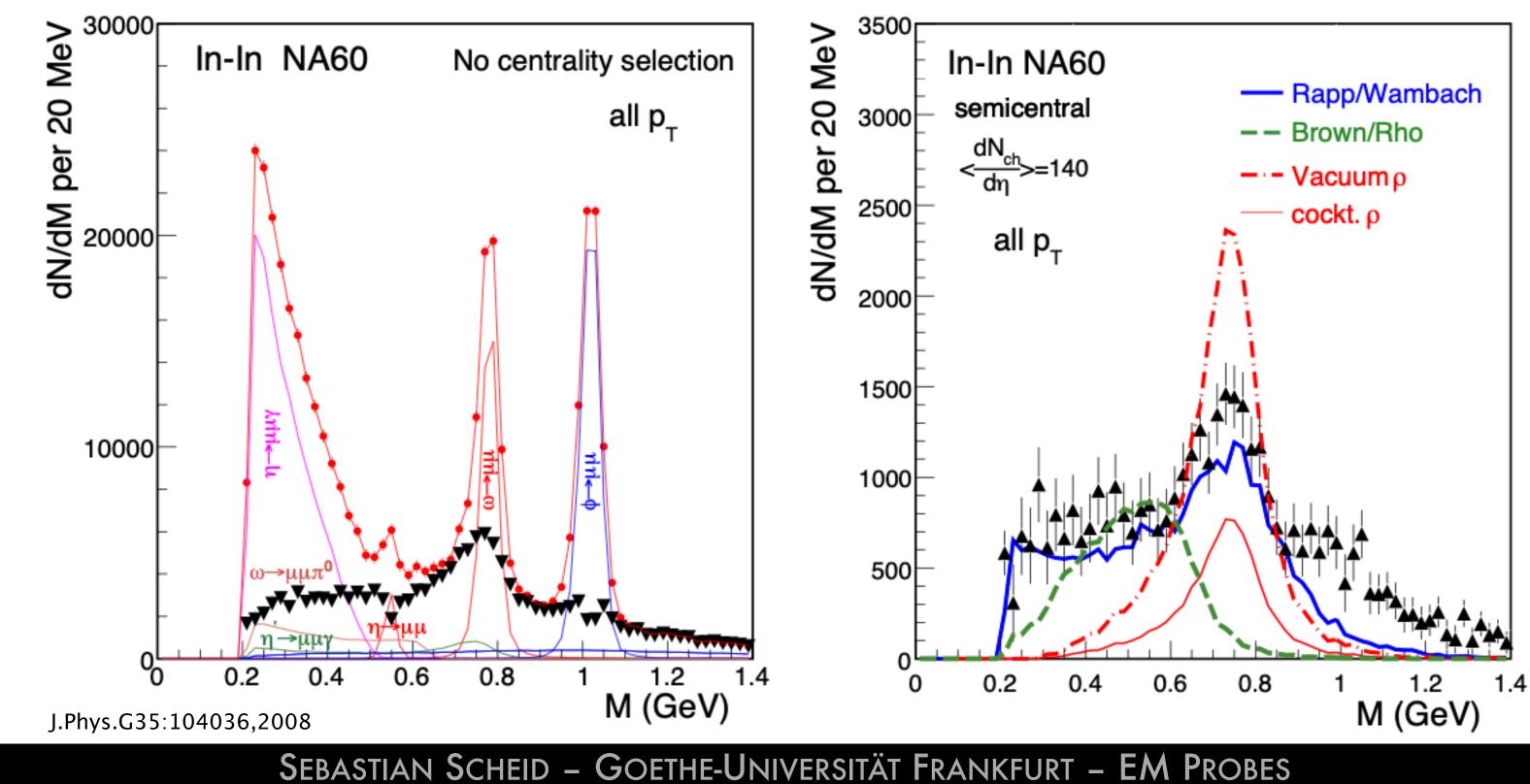
Muons instead of electrons  $\rightarrow$  Trigger possible Addition of vertex telescope to NA50 muon spectrometer gives possibility to match muons to tracks in the tracker → Higher mass resolution and information of muon origin





# SPS - NA60

Muons instead of electrons  $\rightarrow$  Trigger possible to tracks in the tracker  $\rightarrow$  Higher mass resolution and information of muon origin



## Addition of vertex telescope to NA50 muon spectrometer gives possibility to match muons



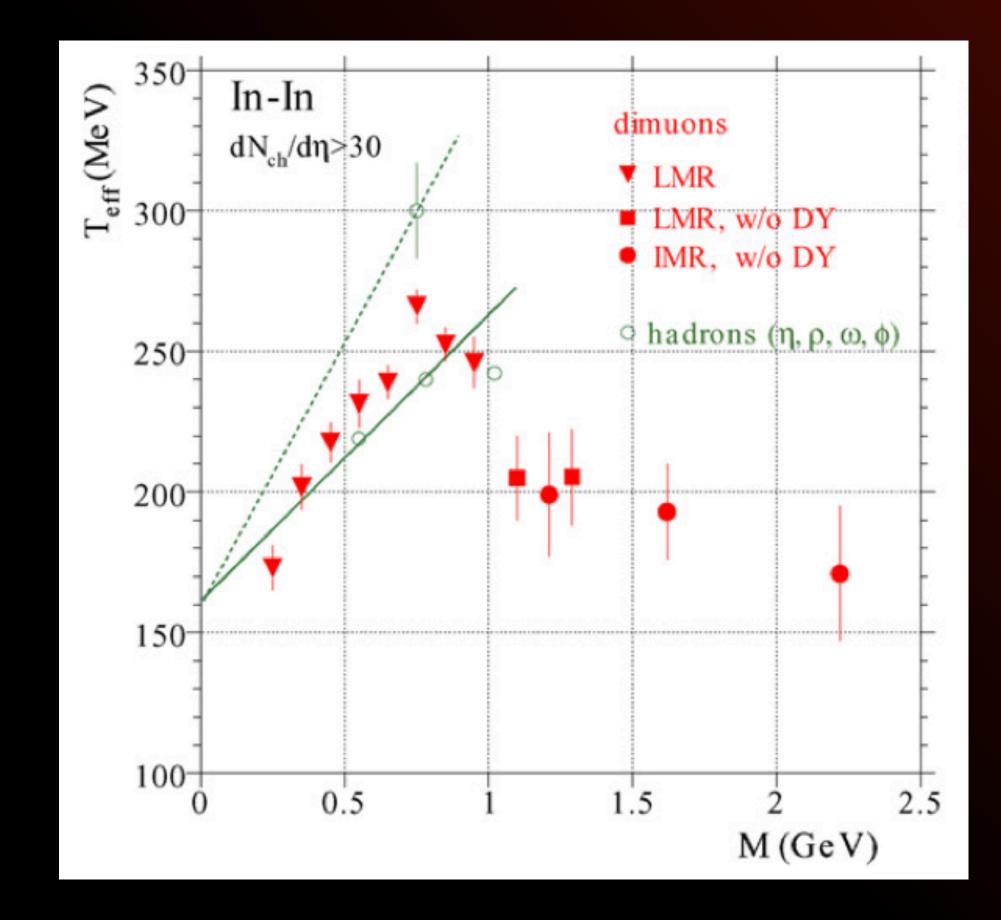


# SPS - NA60

- NA60 also measured the effective temprature as a function of mass Fit of mT inverse slope in mass intervals Low mass:
- T<sub>eff</sub> rises with mass, comparable with hadron → Hadronic Source

Intermediate mass:

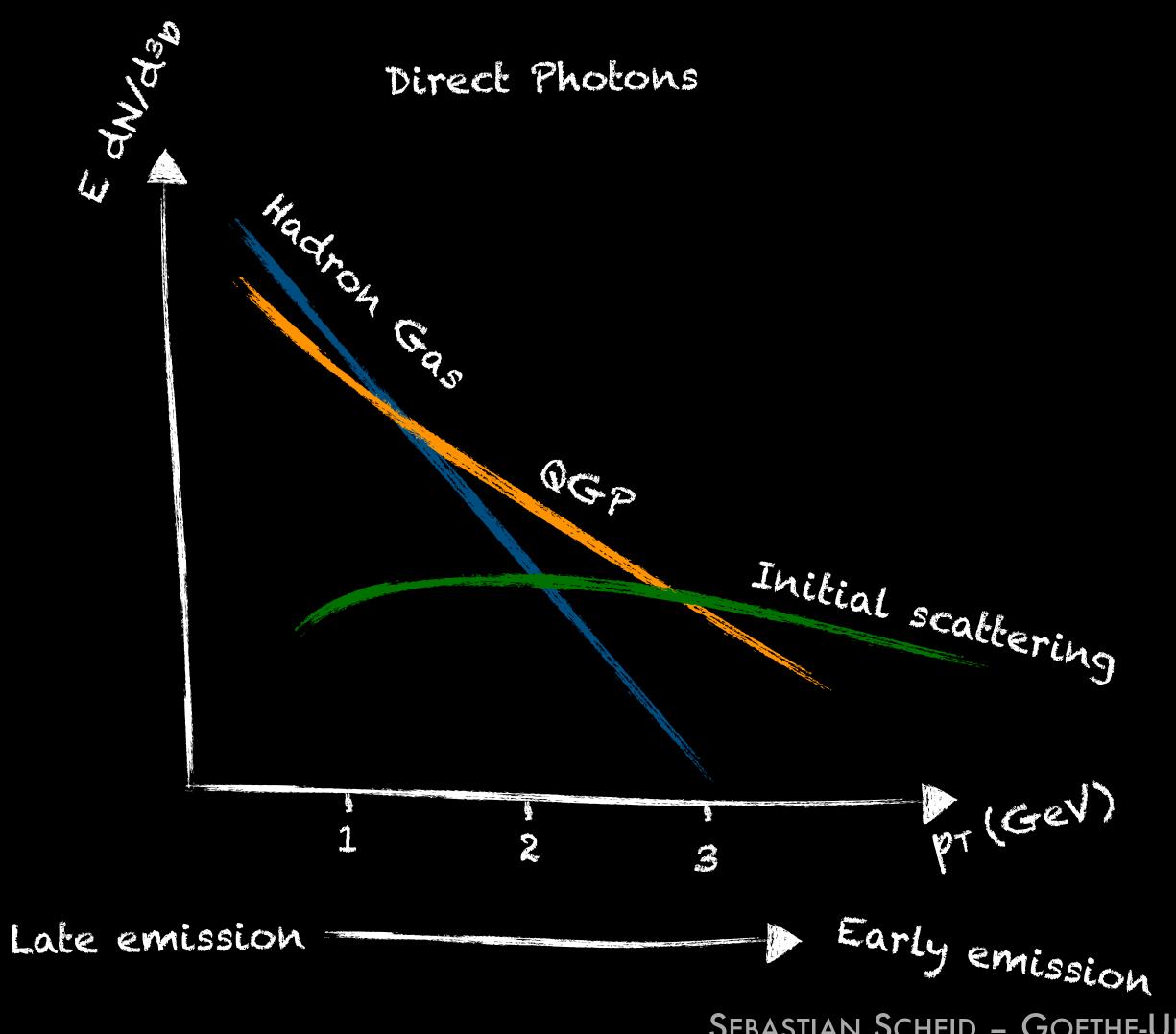
Drop and almost constant value of ~200  $MeV \rightarrow Partonic source$ 

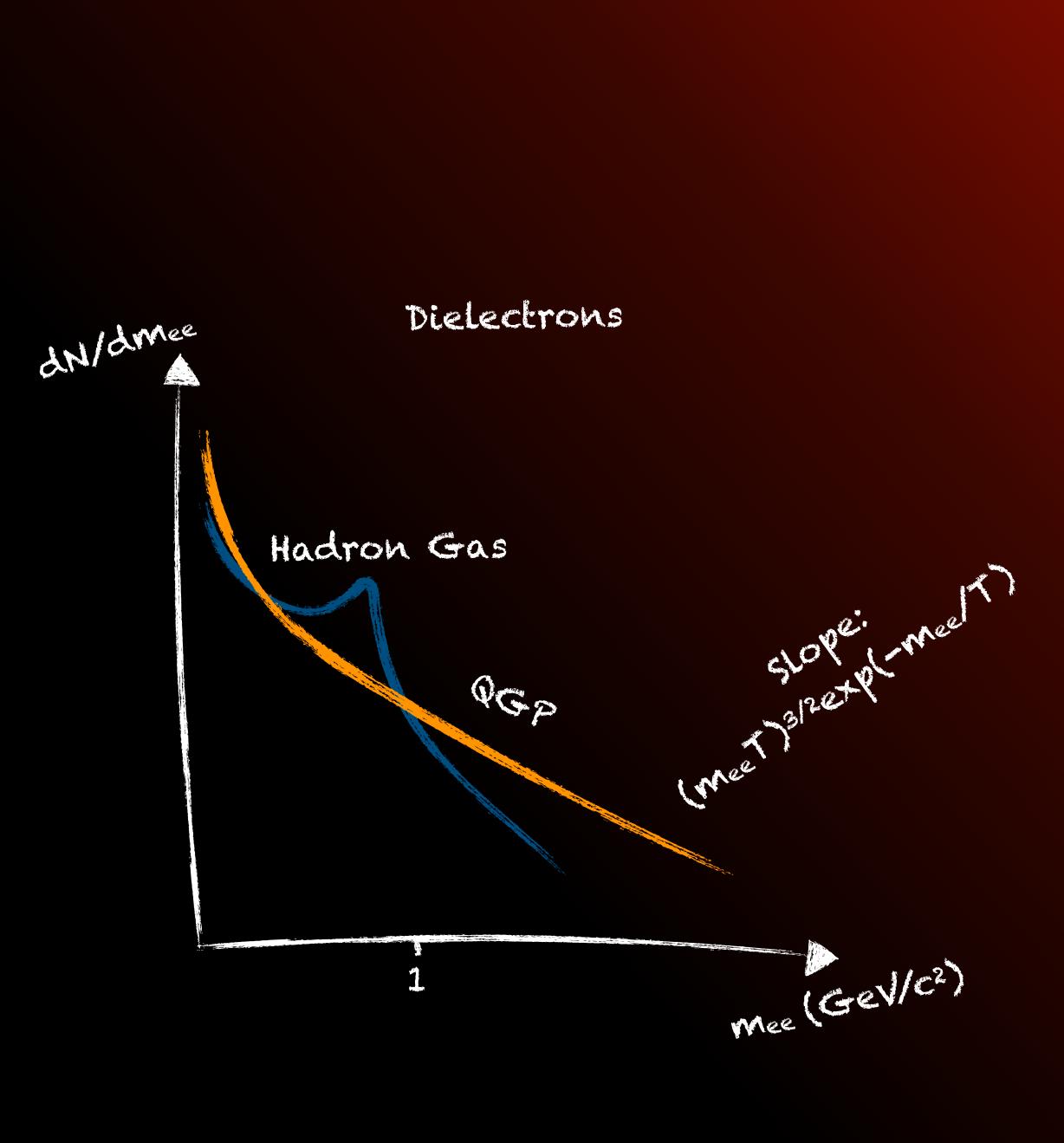


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## The Sources





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