

Momentum Dependent Mean-Fields for (Anti-)Hyperons

Th. Gaitanos, A. Chorozidou



ΤΜΗΜΑ ΦΥΣΙΚΗΣ

ΑΡΙΣΤΟΤΕΛΕΙΟ
ΠΑΝΕΠΙΣΤΗΜΙΟ
ΘΕΣΣΑΛΟΝΙΚΗΣ



Gaitanos & Kaskulov, NPA 940 (2015) 181, NPA 899 (2013) 133
Gaitanos & Chorozidou, NPA (2021) in press

Momentum Dependent Mean-Fields for (Anti-)Hyperons

- Introduction
- The Non-Linear Derivative (NLD) model
- Basic properties: nuclear EoS & p, \bar{p} -optical potentials
- γ properties: density & momentum dependent optical potentials
- $\bar{\gamma}$ properties: density & momentum dependent optical potentials

Momentum Dependent Mean-Fields for (Anti-)Hyperons



- » Introduction
- » The Non-Linear Derivative (NLD) model
- » Basic properties: nuclear EoS & p, \bar{p} -optical potentials
- » γ properties: density & momentum dependent optical potentials
- » $\bar{\gamma}$ properties: density & momentum dependent optical potentials

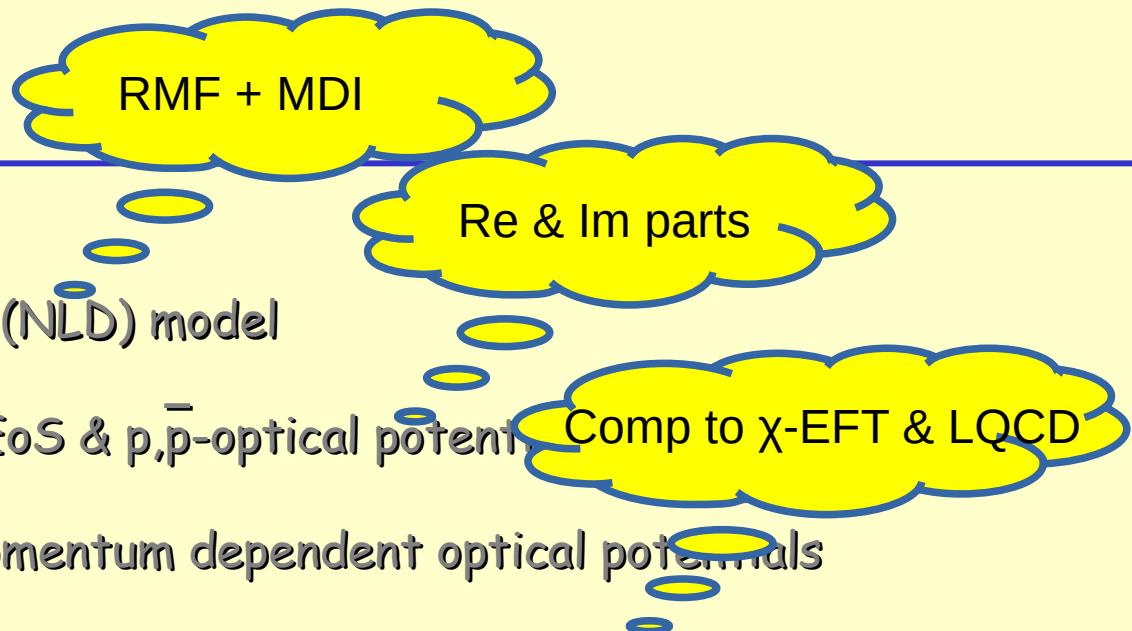
Momentum Dependent Mean-Fields for (Anti-)Hyperons

- » Introduction
- » The Non-Linear Derivative (NLD) model
- » Basic properties: nuclear EoS & p, \bar{p} -optical potentials
- » γ properties: density & momentum dependent optical potentials
- » $\bar{\gamma}$ properties: density & momentum dependent optical potentials



Momentum Dependent Mean-Fields for (Anti-)Hyperons

- » Introduction
- » The Non-Linear Derivative (NLD) model
- » Basic properties: nuclear EoS & p, \bar{p} -optical potent. Comp to χ -EFT & LQCD
- » γ properties: density & momentum dependent optical potentials
- » $\bar{\gamma}$ properties: density & momentum dependent optical potentials



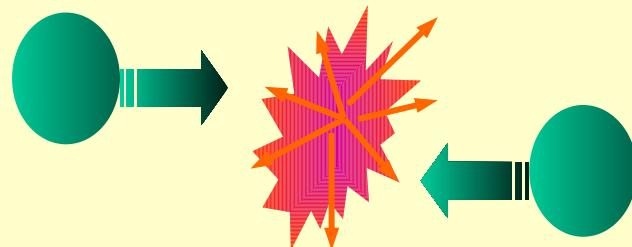
Introduction...

Important for astrophysics

explore EoS far beyond saturation (high ρ , high τ -asymm, $\Lambda/\Sigma/\Xi/\Omega$)

Heavy-ion collisions

(collective flow, meson production)



Densities of fireball for HIC@SIS:
 $\rho \sim (2\text{-}3)\rho_0$

Neutron stars (mass & radius)



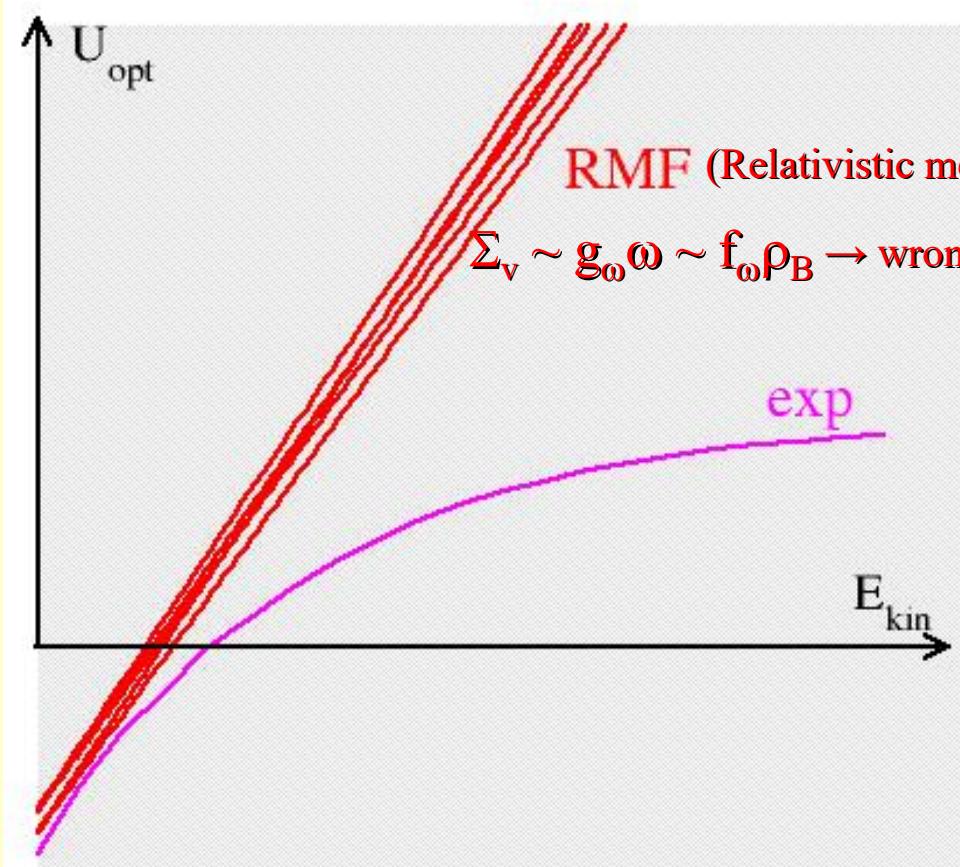
Densities in static NS: $\rho \sim (8\text{-}10)\rho_0$

- ⇒ In high-density matter (+kinematics) → particles with high-momenta p
- ⇒ Not only density dependence, but also **momentum dependence (MD)** essential
- ⇒
- ⇒ Not only nucleon-EoS, but also **hyperon-EoS** essential
- ⇒
- ⇒ Not only hyperon-density dependence, but also **hyperon-momentum dependence** essential

Introduction...

In-medium proton Schrödinger-equivalent $\text{Re}(U_{\text{opt}})$

$$U_{\text{opt}} = \frac{E}{m} \Sigma_v - \Sigma_s + \frac{1}{2m} (\Sigma_s^2 - \Sigma_v^2)$$



DBHF & Dirac-phenomenology for nucleons:

Proton-opt. pot. well known
saturating fields (particular vector) with rising p

Solutions so far:

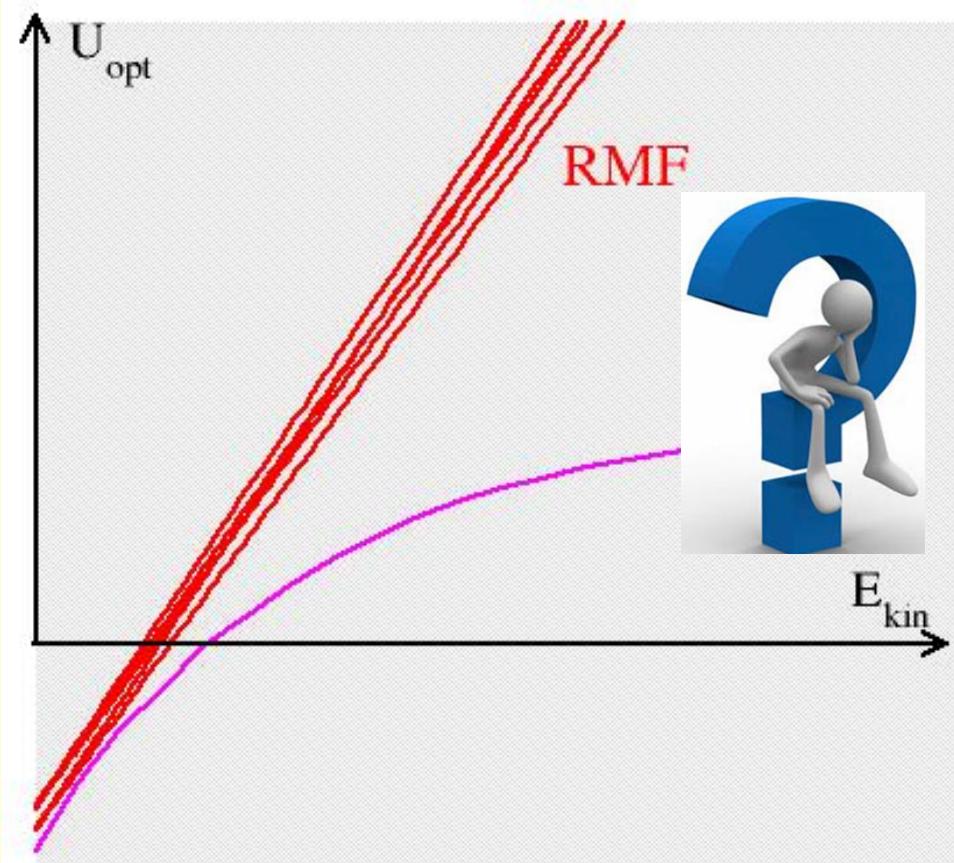
→ non-local (Hartree-Fock) contributions to RMF (Hartree) mean-field
Weber, Blättel, Cassing et al., Nucl. Phys. A539 (1992) 713

→ first-order derivative coupling terms into the interaction Lagrangian
S. Typel, Phys. Rev. C71, 064301 (2005)

Introduction...

In-medium **hyperon** Schrödinger-equivalent $\text{Re}(U_{\text{opt}})$

$$U_{\text{opt}} = \frac{E}{m} \Sigma_v - \Sigma_s + \frac{1}{2m} (\Sigma_s^2 - \Sigma_v^2)$$



Dirac-phenomenology for hyperons:
?
rare experimental scattering data so far

The Non-Linear Derivative (NLD) model...

NLD Lagrangian : as in conventional Relativistic Hadrodynamics (RHD)

$$\mathcal{L} = \frac{1}{2} \sum_B \left[\bar{\Psi}_B \gamma_\mu i \vec{\partial}^\mu \Psi_B - \bar{\Psi}_B i \overleftrightarrow{\partial}^\mu \gamma_\mu \Psi_B \right] - \sum_B m_B \bar{\Psi}_B \Psi_B + \sum_{m=\sigma,\omega,\rho} \mathcal{L}_{int}^m.$$

For the baryon octet:

$$\Psi_B = (\Psi_N, \Psi_\Lambda, \Psi_\Sigma, \Psi_\Xi)^T$$

The Non-Linear Derivative (NLD) model...

NLD Lagrangian : as in conventional Relativistic Hadrodynamics (RHD)

$$\mathcal{L} = \frac{1}{2} \sum_B \left[\bar{\Psi}_B \gamma_\mu i \vec{\partial}^\mu \Psi_B - \bar{\Psi}_B i \overleftrightarrow{\partial}^\mu \gamma_\mu \Psi_B \right] - \sum_B m_B \bar{\Psi}_B \Psi_B + \sum_{m=\sigma,\omega,\rho} \mathcal{L}_{int}^m.$$

For the baryon octet:

$$\Psi_B = (\Psi_N, \Psi_\Lambda, \Psi_\Sigma, \Psi_\Xi)^T$$

Interaction Lagrangian : as in conventional RHD

$$\mathcal{L}_{int}^m = \sum_B \frac{g_{mB}}{2} \left[\bar{\Psi}_B \Gamma_m \Psi_B \varphi_m + \varphi_m \bar{\Psi}_B \Gamma_m \Psi_B \right],$$

For ($\varphi_m = \sigma, \omega, \rho$)-baryon interaction with corresponding vertices $\Gamma_m = \mathbb{1}, \gamma^\mu, \dots$

The Non-Linear Derivative (NLD) model...

NLD Lagrangian : as in conventional Relativistic Hadrodynamics (RHD)

$$\mathcal{L} = \frac{1}{2} \sum_B \left[\bar{\Psi}_B \gamma_\mu i \vec{\partial}^\mu \Psi_B - \bar{\Psi}_B i \overleftrightarrow{\partial}^\mu \gamma_\mu \Psi_B \right] - \sum_B m_B \bar{\Psi}_B \Psi_B + \sum_{m=\sigma,\omega,\rho} \mathcal{L}_{int}^m.$$

For the baryon octet: $\Psi_B = (\Psi_N, \Psi_\Lambda, \Psi_\Sigma, \Psi_\Xi)^T$

Interaction Lagrangian : as in conventional RHD + non-linear derivative operators

$$\mathcal{L}_{int}^m = \sum_B \frac{g_{mB}}{2} \left[\bar{\Psi}_B \overleftarrow{\mathcal{D}}_B \Gamma_m \Psi_B \varphi_m + \varphi_m \bar{\Psi}_B \Gamma_m \overrightarrow{\mathcal{D}}_B \Psi_B \right],$$

For ($\varphi_m = \sigma, \omega, \rho$)-baryon interaction with corresponding vertices $\Gamma_m = \mathbb{1}, \gamma^\mu, \dots$

The Non-Linear Derivative (NLD) model...

NLD Lagrangian : as in conventional Relativistic Hadrodynamics (RHD)

$$\mathcal{L} = \frac{1}{2} \sum_B \left[\bar{\Psi}_B \gamma_\mu i \vec{\partial}^\mu \Psi_B - \bar{\Psi}_B i \overleftarrow{\partial}^\mu \gamma_\mu \Psi_B \right] - \sum_B m_B \bar{\Psi}_B \Psi_B + \sum_{m=\sigma,\omega,\rho} \mathcal{L}_{int}^m.$$

For the baryon octet: $\Psi_B = (\Psi_N, \Psi_\Lambda, \Psi_\Sigma, \Psi_\Xi)^T$

Interaction Lagrangian : as in conventional RHD + non-linear derivative operators

$$\mathcal{L}_{int}^m = \sum_B \frac{g_{mB}}{2} \left[\bar{\Psi}_B \overleftarrow{\mathcal{D}}_B \Gamma_m \Psi_B \varphi_m + \varphi_m \bar{\Psi}_B \Gamma_m \overrightarrow{\mathcal{D}}_B \Psi_B \right],$$

For ($\varphi_m = \sigma, \omega, \rho$)-baryon interaction with corresponding vertices $\Gamma_m = \mathbb{1}, \gamma^\mu, \dots$

Non-linear derivative operators : Taylor expansion of partial derivatives ξ

$$\overrightarrow{\mathcal{D}}_B := \mathcal{D} \left(\overrightarrow{\xi}_B \right), \quad \overleftarrow{\mathcal{D}}_B := \mathcal{D} \left(\overleftarrow{\xi}_B \right) \text{ with } \overrightarrow{\xi}_B = -\frac{v^\alpha i \vec{\partial}_\alpha}{\Lambda_B}, \quad \overleftarrow{\xi}_B = \frac{i \overleftarrow{\partial}_\alpha v^\alpha}{\Lambda_B}$$

v^α auxiliarly 4-vector choosen such to get p-dependence

The Non-Linear Derivative (NLD) model...

NLD Lagrangian : as in conventional Relativistic Hadrodynamics (RHD)

$$\mathcal{L} = \frac{1}{2} \sum_B \left[\bar{\Psi}_B \gamma_\mu i \vec{\partial}^\mu \Psi_B - \bar{\Psi}_B i \overleftrightarrow{\partial}^\mu \gamma_\mu \Psi_B \right] - \sum_B m_B \bar{\Psi}_B \Psi_B + \sum_{m=\sigma,\omega,\rho} \mathcal{L}_{int}^m.$$

For the baryon octet: $\Psi_B = (\Psi_N, \Psi_\Lambda, \Psi_\Sigma, \Psi_\Xi)^T$

Interaction Lagrangian : as in conventional RHD + non-linear derivative operators

$$\mathcal{L}_{int}^m = \sum_B \frac{g_{mB}}{2} \left[\bar{\Psi}_B \overleftarrow{\mathcal{D}}_B \Gamma_m \Psi_B \varphi_m + \varphi_m \bar{\Psi}_B \Gamma_m \overrightarrow{\mathcal{D}}_B \Psi_B \right],$$

For ($\varphi_m = \sigma, \omega, \rho$)-baryon interaction with corresponding vertices $\Gamma_m = \mathbb{1}, \gamma^\mu, \dots$

Non-linear derivative operators : Taylor expansion of partial derivatives ξ

$$\overrightarrow{\mathcal{D}}_B := \mathcal{D} \left(\overrightarrow{\xi}_B \right), \quad \overleftarrow{\mathcal{D}}_B := \mathcal{D} \left(\overleftarrow{\xi}_B \right) \text{ with } \overrightarrow{\xi}_B = -\frac{v^\alpha i \vec{\partial}_\alpha}{\Lambda_B}, \quad \overleftarrow{\xi}_B = \frac{i \overleftarrow{\partial}_\alpha v^\alpha}{\Lambda_B}$$

v^α auxiliarly 4-vector choosen such to get p-dependence

cut-off, will regulate the high-momentum tail of RMF fields

The Non-Linear Derivative (NLD) model...

NLD Lagrangian: contains higher field derivatives: $\mathcal{L}(\varphi_r, \partial_{\alpha_1}\varphi_r, \partial_{\alpha_1\alpha_2}\varphi_r, \dots, \partial_{\alpha_1\dots\alpha_n}\varphi_r)$

→ Generalized Euler-Lagrange equations:

$$\frac{\partial \mathcal{L}}{\partial \varphi_r} + \sum_{i=1}^n (-)^i \partial_{\alpha_1\dots\alpha_i} \frac{\partial \mathcal{L}}{\partial (\partial_{\alpha_1\dots\alpha_i} \varphi_r)} = 0$$

The Non-Linear Derivative (NLD) model...

NLD Lagrangian: contains higher field derivatives: $\mathcal{L}(\varphi_r, \partial_{\alpha_1}\varphi_r, \partial_{\alpha_1\alpha_2}\varphi_r, \dots, \partial_{\alpha_1\dots\alpha_n}\varphi_r)$

→ Generalized Euler-Lagrange equations:

$$\frac{\partial \mathcal{L}}{\partial \varphi_r} + \sum_{i=1}^n (-)^i \partial_{\alpha_1\dots\alpha_i} \frac{\partial \mathcal{L}}{\partial (\partial_{\alpha_1\dots\alpha_i} \varphi_r)} = 0$$

→ Generalized Noether-Theorem: conserved current

$$J^\mu = -i [\mathcal{K}_r^\mu \varphi_r + \mathcal{K}_r^{\mu\sigma_1} \partial_{\sigma_1} \varphi_r + \mathcal{K}_r^{\mu\sigma_1\sigma_2} \partial_{\sigma_1\sigma_2} \varphi_r + \dots + \mathcal{K}_r^{\mu\sigma_1\dots\sigma_n} \partial_{\sigma_1\dots\sigma_n} \varphi_r]$$

with the following tensors

$$\mathcal{K}_r^{\mu\sigma_1\dots\sigma_m} = \sum_{i=1}^n (-)^{i+1} \prod_{j=1}^{i-1} \partial_{\alpha_j} \frac{\partial \mathcal{L}}{\partial (\partial_{\mu\alpha_j\sigma_1\dots\sigma_m} \varphi_r)}.$$

The Non-Linear Derivative (NLD) model...

NLD Lagrangian: contains higher field derivatives: $\mathcal{L}(\varphi_r, \partial_{\alpha_1}\varphi_r, \partial_{\alpha_1\alpha_2}\varphi_r, \dots, \partial_{\alpha_1\dots\alpha_n}\varphi_r)$

→ Generalized Euler-Lagrange equations:

$$\frac{\partial \mathcal{L}}{\partial \varphi_r} + \sum_{i=1}^n (-)^i \partial_{\alpha_1\dots\alpha_i} \frac{\partial \mathcal{L}}{\partial (\partial_{\alpha_1\dots\alpha_i} \varphi_r)} = 0$$

→ Generalized Noether-Theorem: conserved current

$$J^\mu = -i [\mathcal{K}_r^\mu \varphi_r + \mathcal{K}_r^{\mu\sigma_1} \partial_{\sigma_1} \varphi_r + \mathcal{K}_r^{\mu\sigma_1\sigma_2} \partial_{\sigma_1\sigma_2} \varphi_r + \dots + \mathcal{K}_r^{\mu\sigma_1\dots\sigma_n} \partial_{\sigma_1\dots\sigma_n} \varphi_r]$$

with the following tensors

$$\mathcal{K}_r^{\mu\sigma_1\dots\sigma_m} = \sum_{i=1}^n (-)^{i+1} \prod_{j=1}^{i-1} \partial_{\alpha_j} \frac{\partial \mathcal{L}}{\partial (\partial_{\mu\alpha_j\sigma_1\dots\sigma_m} \varphi_r)}.$$

infinite series resp. to higher-order field derivatives, but...

The Non-Linear Derivative (NLD) model...

NLD Lagrangian: contains higher field derivatives: $\mathcal{L}(\varphi_r, \partial_{\alpha_1}\varphi_r, \partial_{\alpha_1\alpha_2}\varphi_r, \dots, \partial_{\alpha_1\dots\alpha_n}\varphi_r)$

→ Generalized Euler-Lagrange equations:

$$\frac{\partial \mathcal{L}}{\partial \varphi_r} + \sum_{i=1}^n (-)^i \partial_{\alpha_1\dots\alpha_i} \frac{\partial \mathcal{L}}{\partial (\partial_{\alpha_1\dots\alpha_i} \varphi_r)} = 0$$

→ Generalized Noether-Theorem: conserved current

$$J^\mu = -i [\mathcal{K}_r^\mu \varphi_r + \mathcal{K}_r^{\mu\sigma_1} \partial_{\sigma_1} \varphi_r + \mathcal{K}_r^{\mu\sigma_1\sigma_2} \partial_{\sigma_1\sigma_2} \varphi_r + \dots + \mathcal{K}_r^{\mu\sigma_1\dots\sigma_n} \partial_{\sigma_1\dots\sigma_n} \varphi_r]$$

with the following tensors

$$\mathcal{K}_r^{\mu\sigma_1\dots\sigma_m} = \sum_{i=1}^n (-)^{i+1} \prod_{j=1}^{i-1} \partial_{\alpha_j} \frac{\partial \mathcal{L}}{\partial (\partial_{\mu\alpha_j\sigma_1\dots\sigma_m} \varphi_r)}.$$

All infinite series can be resummed to compact expressions

The Non-Linear Derivative (NLD) model...

→ Dirac equation for nucleons $[\gamma_\mu(i\partial^\mu - \Sigma^\mu) - (m - \Sigma_s)] \Psi = 0$ with selfenergies

$$\Sigma^\mu = g_\omega \omega^\mu \vec{\mathcal{D}} + g_\rho \vec{\tau} \cdot \vec{\rho}^\mu \vec{\mathcal{D}} + \dots$$

$$\Sigma_s = g_\sigma \sigma \vec{\mathcal{D}} + \dots$$

(up to terms containing derivatives of the meson fields)

The Non-Linear Derivative (NLD) model...

→ Dirac equation for nucleons $[\gamma_\mu(i\partial^\mu - \Sigma^\mu) - (m - \Sigma_s)]\Psi = 0$ with selfenergies

$$\Sigma^\mu = g_\omega \omega^\mu \vec{\mathcal{D}} + g_\rho \vec{\tau} \cdot \vec{\rho}^\mu \vec{\mathcal{D}} + \dots$$

$$\Sigma_s = g_\sigma \sigma \vec{\mathcal{D}} + \dots$$

(up to terms containing derivatives of the meson fields)

→ Meson field equations:

$$\begin{aligned}\partial_\alpha \partial^\alpha \sigma + m_\sigma^2 \sigma + \frac{\partial U}{\partial \sigma} &= \frac{1}{2} g_\sigma [\bar{\Psi} \vec{\mathcal{D}} \Psi + \bar{\Psi} \vec{\mathcal{D}} \Psi] , \\ \partial_\mu F^{\mu\nu} + m_\omega^2 \omega^\nu &= \frac{1}{2} g_\omega [\bar{\Psi} \vec{\mathcal{D}} \gamma^\nu \Psi + \bar{\Psi} \gamma^\nu \vec{\mathcal{D}} \Psi] , \\ \partial_\mu \vec{G}^{\mu\nu} + m_\rho^2 \vec{\rho}^\nu &= \frac{1}{2} g_\rho [\bar{\Psi} \vec{\mathcal{D}} \gamma^\nu \vec{\tau} \Psi + \bar{\Psi} \vec{\tau} \gamma^\nu \vec{\mathcal{D}} \Psi]\end{aligned}$$

The Non-Linear Derivative (NLD) model...

→ Dirac equation for nucleons $[\gamma_\mu(i\partial^\mu - \Sigma^\mu) - (m - \Sigma_s)]\Psi = 0$ with selfenergies

$$\Sigma^\mu = g_\omega \omega^\mu \vec{\mathcal{D}} + g_\rho \vec{\tau} \cdot \vec{\rho}^\mu \vec{\mathcal{D}} + \dots$$

$$\Sigma_s = g_\sigma \sigma \vec{\mathcal{D}} + \dots$$

(up to terms containing derivatives of the meson fields)

→ Meson field equations:

$$\begin{aligned} \partial_\alpha \partial^\alpha \sigma + m_\sigma^2 \sigma + \frac{\partial U}{\partial \sigma} &= \frac{1}{2} g_\sigma [\bar{\Psi} \overset{\leftarrow}{\mathcal{D}} \Psi + \bar{\Psi} \overset{\rightarrow}{\mathcal{D}} \Psi] , \\ \partial_\mu F^{\mu\nu} + m_\omega^2 \omega^\nu &= \frac{1}{2} g_\omega [\bar{\Psi} \overset{\leftarrow}{\mathcal{D}} \gamma^\nu \Psi + \bar{\Psi} \gamma^\nu \overset{\rightarrow}{\mathcal{D}} \Psi] , \\ \partial_\mu \vec{G}^{\mu\nu} + m_\rho^2 \vec{\rho}^\nu &= \frac{1}{2} g_\rho [\bar{\Psi} \overset{\leftarrow}{\mathcal{D}} \gamma^\nu \vec{\tau} \Psi + \bar{\Psi} \vec{\tau} \gamma^\nu \overset{\rightarrow}{\mathcal{D}} \Psi] \end{aligned}$$

→ Energy-momentum tensor:

$$\begin{aligned} T^{\mu\nu} &= \frac{1}{2} \bar{\Psi} \gamma^\mu i \overset{\rightarrow}{\partial}^\nu \Psi - \frac{1}{2} \bar{\Psi} i \overset{\leftarrow}{\partial}^\nu \gamma^\mu \Psi \\ &\quad + \frac{1}{2} \sum_m g_m [\bar{\Psi} \Gamma_m \overset{\rightarrow}{\mathcal{D}}^\mu i \overset{\rightarrow}{\partial}^\nu \Psi + \bar{\Psi} i \overset{\leftarrow}{\partial}^\nu \overset{\leftarrow}{\mathcal{D}}^\mu \Gamma_m \Psi] \varphi_m - g^{\mu\nu} \mathcal{L} + \dots . \end{aligned}$$

The NLD model: RMF approach to INM...

→ Plane wave Ansatz for Ψ and $\bar{\Psi}$ with $\mathcal{D} = \mathcal{D}(p)$

$$\Sigma_{vi}^\mu = g_\omega \omega^\mu \mathcal{D} + g_\rho \tau_i \rho^\mu \mathcal{D}, \quad \Sigma_{si} = g_\sigma \sigma \mathcal{D}$$

$$m_\sigma^2 \sigma + \frac{\partial U}{\partial \sigma} = g_\sigma \sum_{i=p,n} \left\langle \bar{\Psi}_i \mathcal{D} \Psi_i \right\rangle = g_\sigma \rho_s$$
$$m_\omega^2 \omega = g_\omega \sum_{i=p,n} \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\omega \rho_0$$
$$m_\rho^2 \rho = g_\rho \sum_{i=p,n} \tau_i \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\rho \rho_I .$$

$$T^{\mu\nu} = \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int d^3 p \frac{\Pi_i^\mu p^\nu}{\Pi_i^0} - g^{\mu\nu} \langle \mathcal{L} \rangle$$

$$\Pi_i^\mu = p_i^{*\mu} + m_i^* \left(\partial_p^\mu \Sigma_{si} \right) - \left(\partial_p^\mu \Sigma_{vi}^\beta \right) p_{i\beta}^*$$

The NLD model: RMF approach to INM...

→ Plane wave Ansatz for Ψ and $\bar{\Psi}$ with $\mathcal{D} = \mathcal{D}(p)$

$$\Sigma_{vi}^\mu = g_\omega \omega^\mu \mathcal{D} + g_\rho \tau_i \rho^\mu \mathcal{D}, \quad \Sigma_{si} = g_\sigma \sigma \mathcal{D}$$

meson-field equations

$$m_\sigma^2 \sigma + \frac{\partial U}{\partial \sigma} = g_\sigma \sum_{i=p,n} \left\langle \bar{\Psi}_i \mathcal{D} \Psi_i \right\rangle = g_\sigma \rho_s$$

$$m_\omega^2 \omega = g_\omega \sum_{i=p,n} \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\omega \rho_0$$

$$m_\rho^2 \rho = g_\rho \sum_{i=p,n} \tau_i \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\rho \rho_I$$

Equation of State (EoS)

$$\varepsilon = \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int_{|\vec{p}| \leq p_{F_i}} d^3 p E(\vec{p}) - \langle \mathcal{L} \rangle$$

$$P = \frac{1}{3} \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int_{|\vec{p}| \leq p_{F_i}} d^3 p \frac{\vec{\Pi}_i \cdot \vec{p}}{\Pi_i^0} + \langle \mathcal{L} \rangle$$

Features of NLD...

→ Plane wave Ansatz for Ψ and $\bar{\Psi}$ with $\mathcal{D} = \mathcal{D}(p)$

$$\Sigma_{vi}^\mu = g_\omega \omega^\mu \mathcal{D} + g_\rho \tau_i \rho^\mu \mathcal{D}, \quad \Sigma_{si} = g_\sigma \sigma \mathcal{D}$$

meson-field equations

$$m_\sigma^2 \sigma + \frac{\partial U}{\partial \sigma} = g_\sigma \sum_{i=p,n} \left\langle \bar{\Psi}_i \mathcal{D} \Psi_i \right\rangle = g_\sigma \rho_s$$

$$m_\omega^2 \omega = g_\omega \sum_{i=p,n} \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\omega \rho_0$$

$$m_\rho^2 \rho = g_\rho \sum_{i=p,n} \tau_i \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\rho \rho_I$$

cut-off Λ regulates:

1) DD & MD of selfenergies

Equation of State (EoS)

$$\varepsilon = \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int_{|\vec{p}| \leq p_{F_i}} d^3 p E(\vec{p}) - \langle \mathcal{L} \rangle$$

$$P = \frac{1}{3} \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int_{|\vec{p}| \leq p_{F_i}} d^3 p \frac{\vec{\Pi}_i \cdot \vec{p}}{\Pi_i^0} + \langle \mathcal{L} \rangle$$

Features of NLD...

→ Plane wave Ansatz for Ψ and $\bar{\Psi}$ with $\mathcal{D} = \mathcal{D}(p)$

$$\Sigma_{vi}^\mu = g_\omega \omega^\mu \mathcal{D} + g_\rho \tau_i \rho^\mu \mathcal{D}, \quad \Sigma_{si} = g_\sigma \sigma \mathcal{D}$$

meson-field equations

$$m_\sigma^2 \sigma + \frac{\partial U}{\partial \sigma} = g_\sigma \sum_{i=p,n} \left\langle \bar{\Psi}_i \mathcal{D} \Psi_i \right\rangle = g_\sigma \rho_s$$

$$m_\omega^2 \omega = g_\omega \sum_{i=p,n} \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\omega \rho_0$$

$$m_\rho^2 \rho = g_\rho \sum_{i=p,n} \tau_i \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\rho \rho_I$$

Equation of State (EoS)

$$\varepsilon = \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int_{|\vec{p}| \leq p_{F_i}} d^3 p E(\vec{p}) - \langle \mathcal{L} \rangle$$

$$P = \frac{1}{3} \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int_{|\vec{p}| \leq p_{F_i}} d^3 p \frac{\vec{\Pi}_i \cdot \vec{p}}{\Pi_i^0} + \langle \mathcal{L} \rangle$$

cut-off Λ regulates:

1) DD & MD of selfenergies

2) DD of meson-field sources
(particularly for ω -field)

Features of NLD...

→ Plane wave Ansatz for Ψ and $\bar{\Psi}$ with $\mathcal{D} = \mathcal{D}(p)$

$$\Sigma_{vi}^\mu = g_\omega \omega^\mu \mathcal{D} + g_\rho \tau_i \rho^\mu \mathcal{D}, \quad \Sigma_{si} = g_\sigma \sigma \mathcal{D}$$

meson-field equations

$$m_\sigma^2 \sigma + \frac{\partial U}{\partial \sigma} = g_\sigma \sum_{i=p,n} \left\langle \bar{\Psi}_i \mathcal{D} \Psi_i \right\rangle = g_\sigma \rho_s$$

$$m_\omega^2 \omega = g_\omega \sum_{i=p,n} \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\omega \rho_0$$

$$m_\rho^2 \rho = g_\rho \sum_{i=p,n} \tau_i \left\langle \bar{\Psi}_i \gamma^0 \mathcal{D} \Psi_i \right\rangle = g_\rho \rho_I$$

Equation of State (EoS)

$$\varepsilon = \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int_{|\vec{p}| \leq p_{F_i}} d^3 p E(\vec{p}) - \langle \mathcal{L} \rangle$$

$$P = \frac{1}{3} \sum_{i=p,n} \frac{\kappa}{(2\pi)^3} \int_{|\vec{p}| \leq p_{F_i}} d^3 p \frac{\vec{\Pi}_i \cdot \vec{p}}{\Pi_i^0} + \langle \mathcal{L} \rangle$$

cut-off Λ regulates:

1) DD & MD of selfenergies

2) DD of meson-field sources
(particularly for ω -field)

3) fully thermodynamic consistent

Basic properties: nuclear EoS & opt. potentials...

Parameters

| | \vec{D} | cut-off | Λ_s [GeV] | Λ_v [GeV] | g_σ | g_ω | g_ρ | b [fm $^{-1}$] | c | m_σ [GeV] | m_ω [GeV] | m_ρ [GeV] |
|-----|---|---|------------------------|------------------------|------------|------------|----------|-----------------------|---------|-----------------------|-----------------------|---------------------|
| NLD | $\frac{1}{1 + \sum_{j=1}^4 (\zeta_j^\alpha i \vec{\partial}_\alpha)^2}$ | $\frac{\Lambda^2}{\Lambda^2 + \vec{p}^2}$ | 0.95 | 1.125 | 10.08 | 10.13 | 3.50 | 15.341 | -14.735 | 0.592 | 0.782 | 0.763 |

Comparison with other models

| Model | ρ_{sat} [fm $^{-3}$] | E_b [MeV/A] | K [MeV] | a_{sym} [MeV] | L [MeV] | K_{sym} [MeV] | K_{asy} [MeV] | |
|------------------|--------------------------------|--------------------|----------------|----------------------|----------------|----------------------|----------------------|----------------------------|
| NLD | 0.156 | -15.30 | 251 | 30 | 81 | -28 | -514 | |
| NL3* | 0.150 | -16.31 | 258 | 38.68 | 125.7 | 104.08 | -650.12 | → Lalazissis |
| DD | 0.149 | -16.02 | 240 | 31.60 | 56 | -95.30 | -431.30 | → Typel |
| D ³ C | 0.151 | -15.98 | 232.5 | 31.90 | 59.30 | -74.7 | -430.50 | |
| DBHF | 0.185 | -15.60 | 290 | 33.35 | 71.10 | -27.1 | -453.70 | → Li, Machleidt, Brockmann |
| | 0.181 | -16.15 | 230 | 34.20 | 71 | 87.36 | -340 | → Fuchs |
| empirical | 0.167 ± 0.019 | -16 ± 1 | 230 ± 10 | 31.1 ± 1.9 | 88 ± 25 | - | -550 ± 100 | THEIA Seminar, 10/02/2021 |

Basic properties: nuclear EoS & opt. potentials...

Parameters

| | \vec{D} | cut-off [GeV] | monopole form | g_ρ | b [fm $^{-1}$] | c | m_σ [GeV] | m_ω [GeV] | m_ρ [GeV] |
|-----|---|---|---------------|---|----------------------|-----|---------------------|---------------------|-------------------|
| NLD | $\frac{1}{1 + \sum_{j=1}^4 (\zeta_j^\alpha i \vec{\partial}_\alpha)^2}$ | $\frac{\Lambda^2}{\Lambda^2 + \vec{p}^2}$ | 0.95 1.125 | 10.08 10.13 3.50 15.341 -14.735 | | | 0.592 | 0.782 | 0.763 |

Comparison with other models

| Model | ρ_{sat} [fm $^{-3}$] | E_b [MeV/A] | K [MeV] | a_{sym} [MeV] | L [MeV] | K_{sym} [MeV] | K_{asy} [MeV] | |
|------------------|-------------------------------|------------------|--------------|--------------------|--------------|--------------------|--------------------|----------------------------|
| NLD | 0.156 | -15.30 | 251 | 30 | 81 | -28 | -514 | |
| NL3* | 0.150 | -16.31 | 258 | 38.68 | 125.7 | 104.08 | -650.12 | → Lalazissis |
| DD | 0.149 | -16.02 | 240 | 31.60 | 56 | -95.30 | -431.30 | → Typel |
| D ³ C | 0.151 | -15.98 | 232.5 | 31.90 | 59.30 | -74.7 | -430.50 | |
| DBHF | 0.185 | -15.60 | 290 | 33.35 | 71.10 | -27.1 | -453.70 | → Li, Machleidt, Brockmann |
| | 0.181 | -16.15 | 230 | 34.20 | 71 | 87.36 | -340 | → Fuchs |
| empirical | 0.167 ± 0.019 | -16 ± 1 | 230 ± 10 | 31.1 ± 1.9 | 88 ± 25 | - | -550 ± 100 | |

Basic properties: nuclear EoS & opt. potentials...

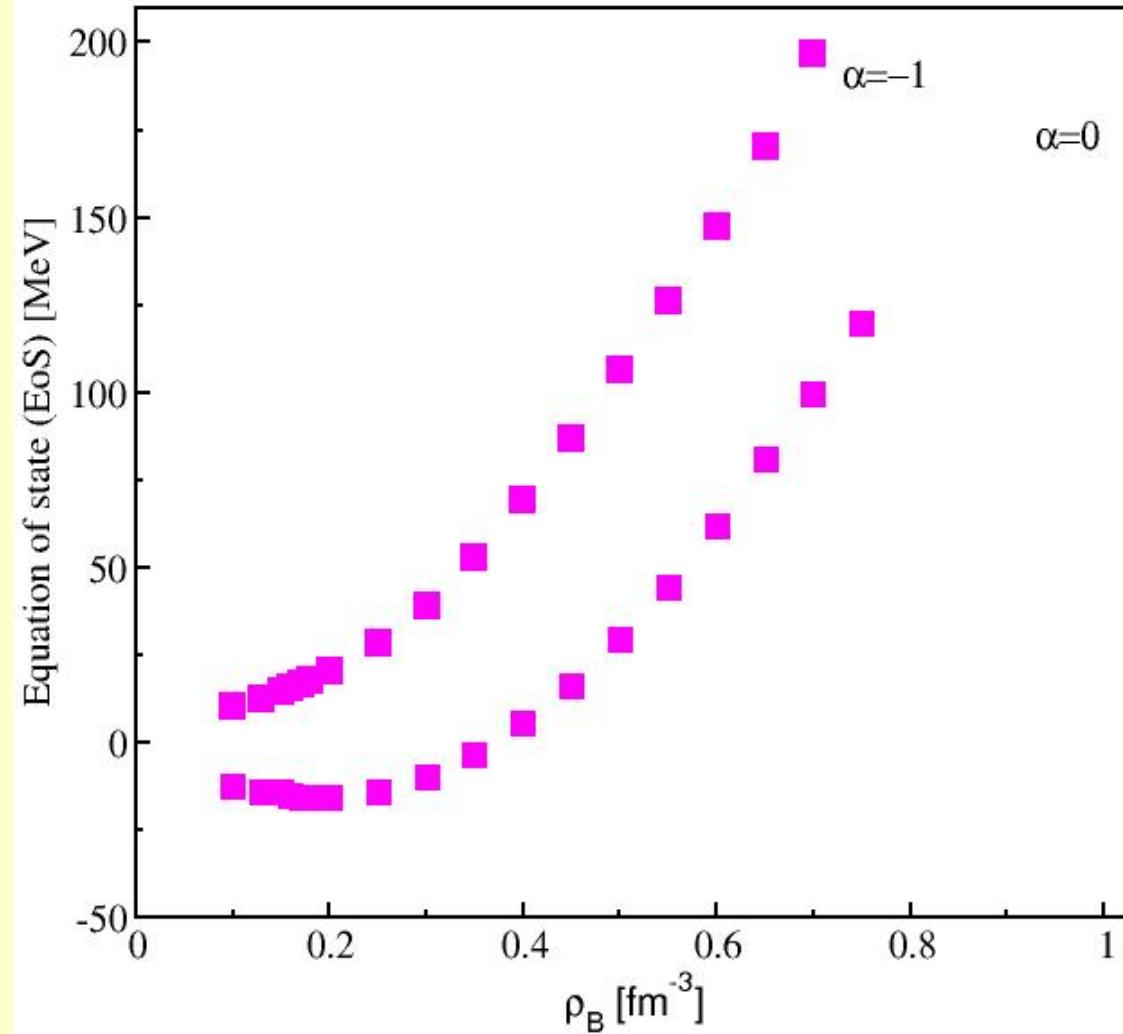
Parameters

| | \vec{D} | cut-off | Λ_s [GeV] | Λ_v [GeV] | g_σ | g_ω | g_ρ | b [fm $^{-1}$] | c | m_σ [GeV] | m_ω [GeV] | m_ρ [GeV] |
|-----|---|---|-----------------------|-----------------------|------------|------------|----------|-----------------------|---------|----------------------|----------------------|--------------------|
| NLD | $\frac{1}{1 + \sum_{j=1}^4 (\zeta_j^\alpha i \vec{\partial}_\alpha)^2}$ | $\frac{\Lambda^2}{\Lambda^2 + \vec{p}^2}$ | 0.95 | 1.125 | 10.08 | 10.13 | 3.50 | 15.341 | -14.735 | 0.592 | 0.782 | 0.763 |

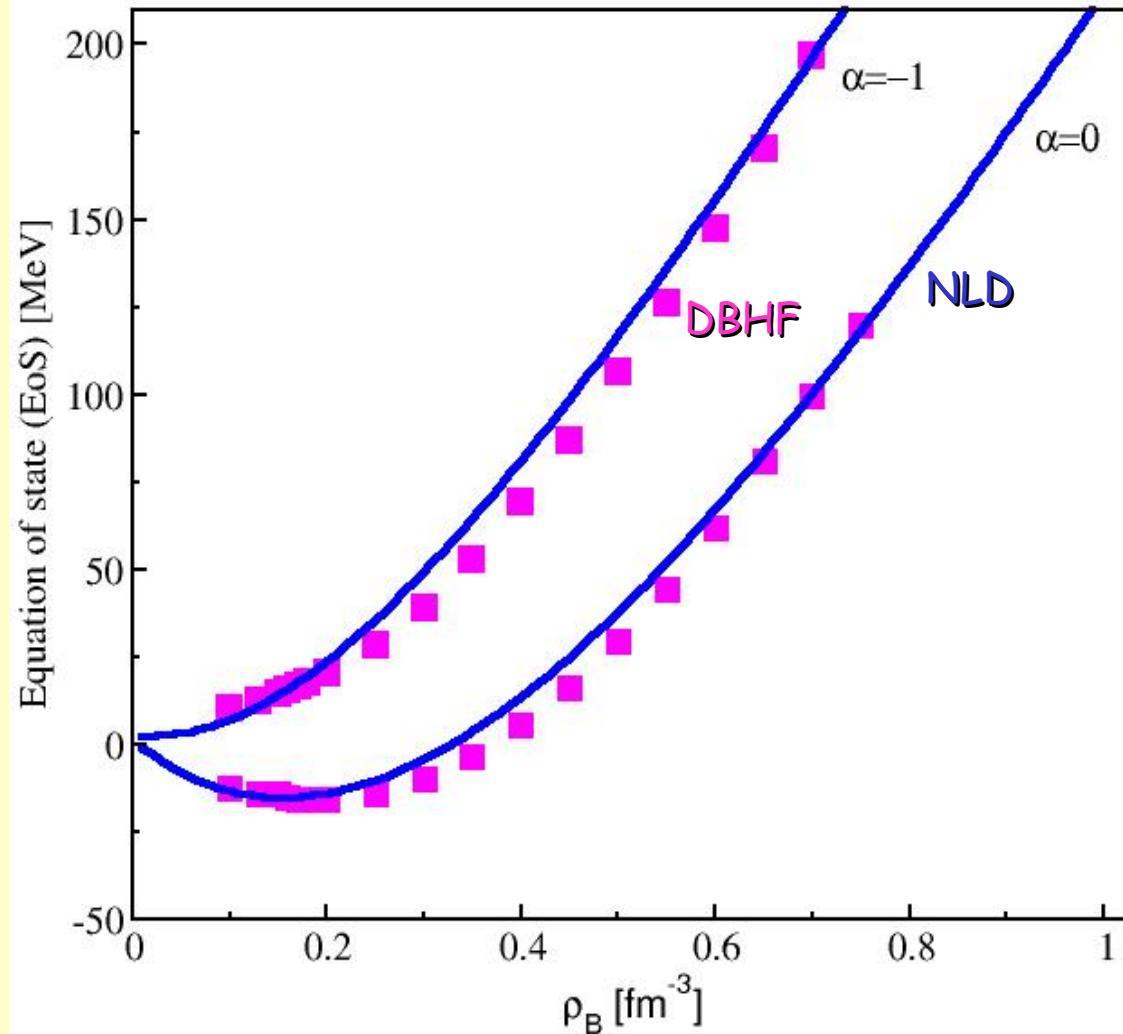
Comparison with other models

| Model | ρ_{sat} [fm $^{-3}$] | E_b [MeV/A] | K [MeV] | a [fm 3] | soft EoS at ρ_{sat} , but stiff at high ρ relevant for NS! | | | | |
|------------------|-------------------------------|-------------------|---------------|--------------------|---|--------|----------------|----------------------------|--|
| NLD | 0.156 | -15.30 | 251 | 30 | 81 | -28 | -514 | | |
| NL3* | 0.150 | -16.31 | 258 | 38.68 | 125.7 | 104.08 | -650.12 | → Lalazissis | |
| DD | 0.149 | -16.02 | 240 | 31.60 | 56 | -95.30 | -431.30 | → Typel | |
| D ³ C | 0.151 | -15.98 | 232.5 | 31.90 | 59.30 | -74.7 | -430.50 | | |
| DBHF | 0.185 | -15.60 | 290 | 33.35 | 71.10 | -27.1 | -453.70 | → Li, Machleidt, Brockmann | |
| | 0.181 | -16.15 | 230 | 34.20 | 71 | 87.36 | -340 | → Fuchs | |
| empirical | 0.167 ± 0.019 | -16 ± 1 | 230 ± 10 | 31.1 ± 1.9 | 88 ± 25 | - | -550 ± 100 | | |

Basic properties: nuclear EoS & opt. potentials...

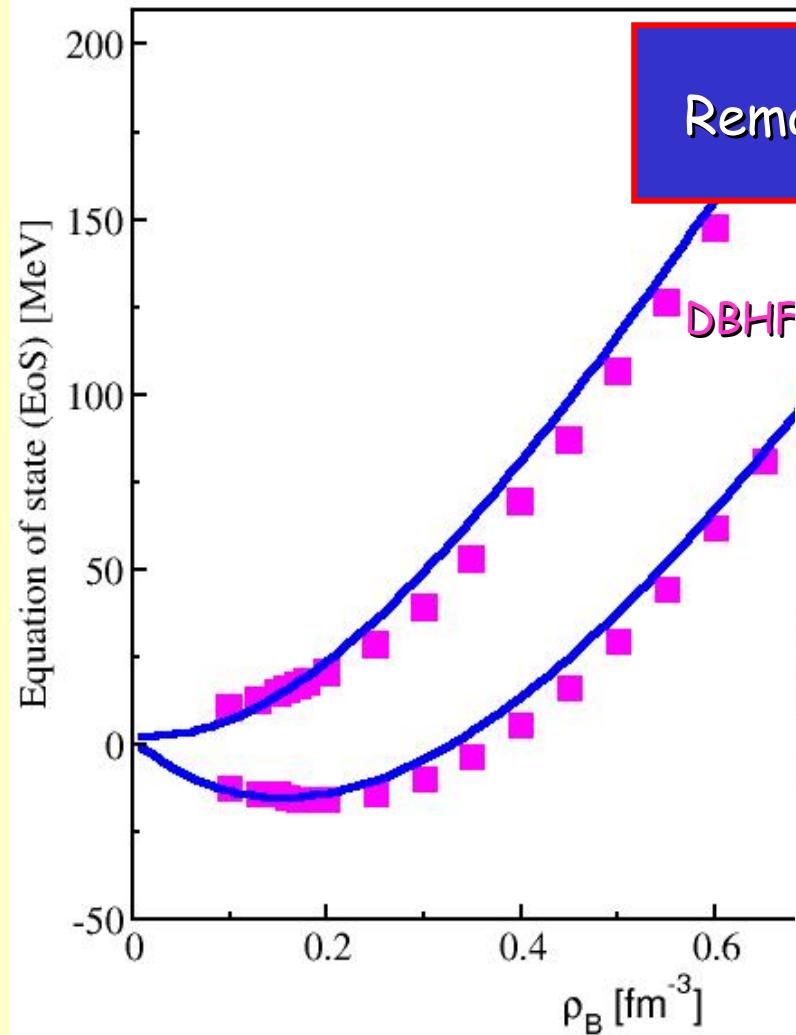


Basic properties: nuclear EoS & opt. potentials...

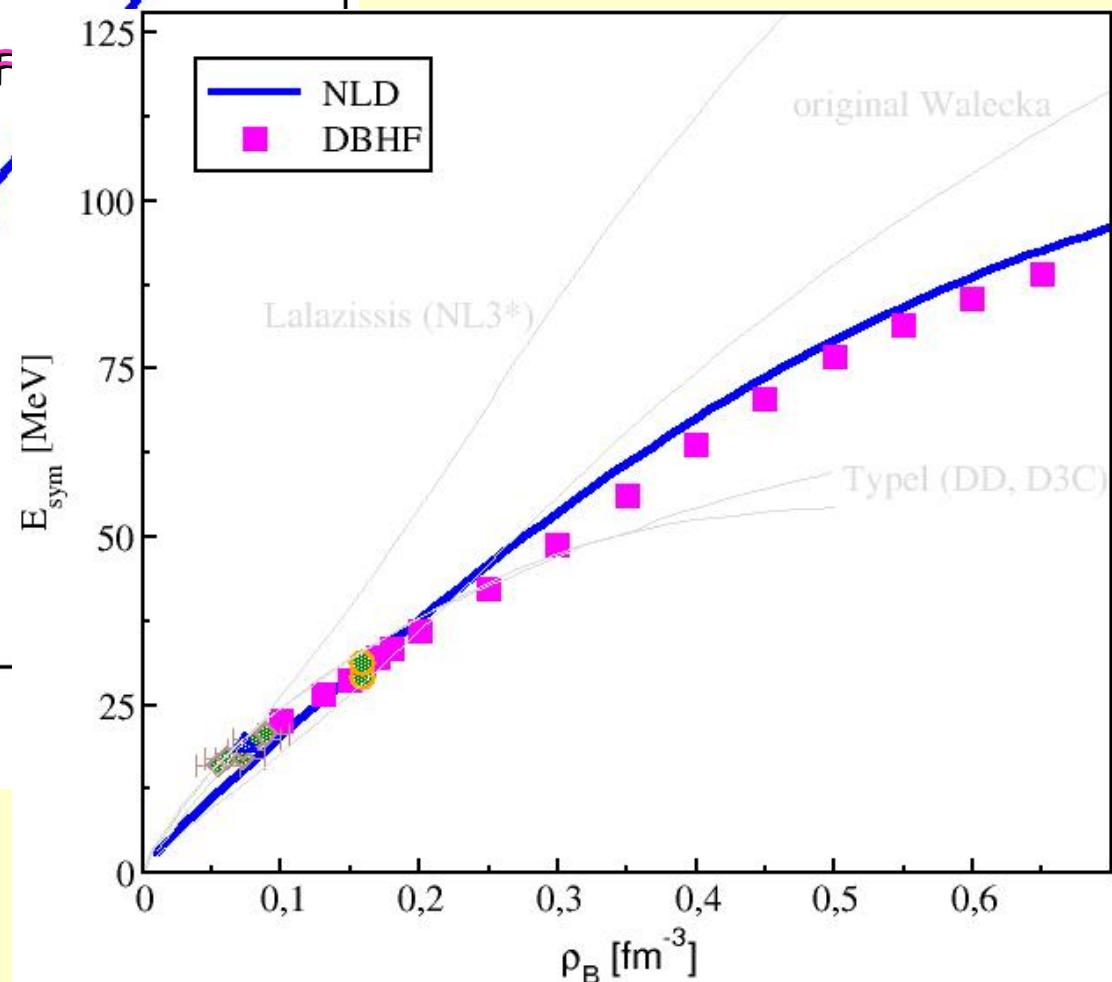


Remarkable comparison with microscopic DBHF !

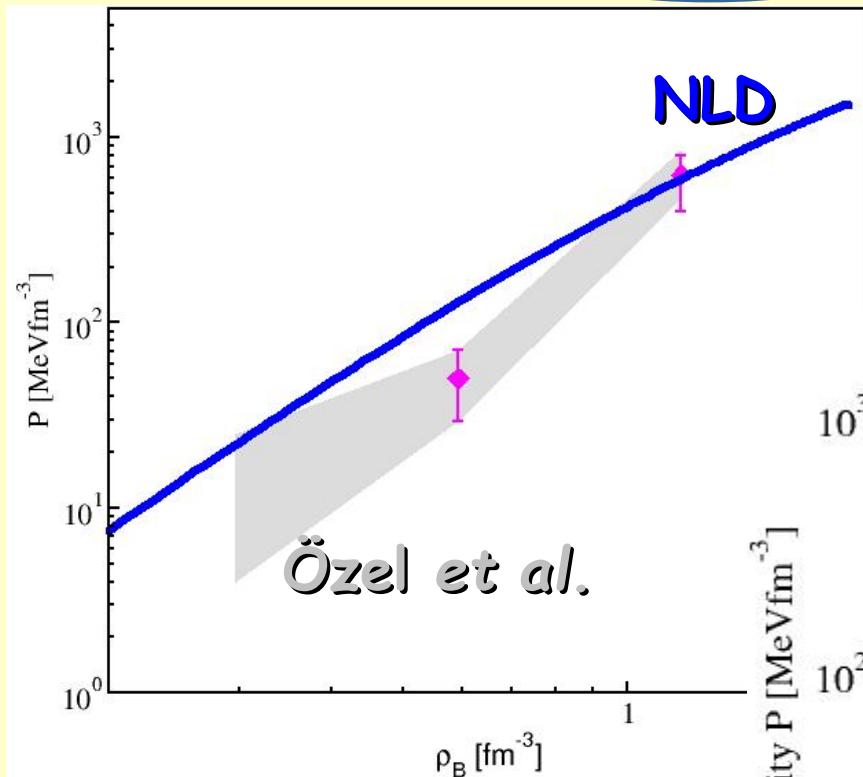
Basic properties: nuclear EoS & opt. potentials...



Remarkable comparison with microscopic DBHF !



Basic properties: nuclear EoS & opt. potentials...

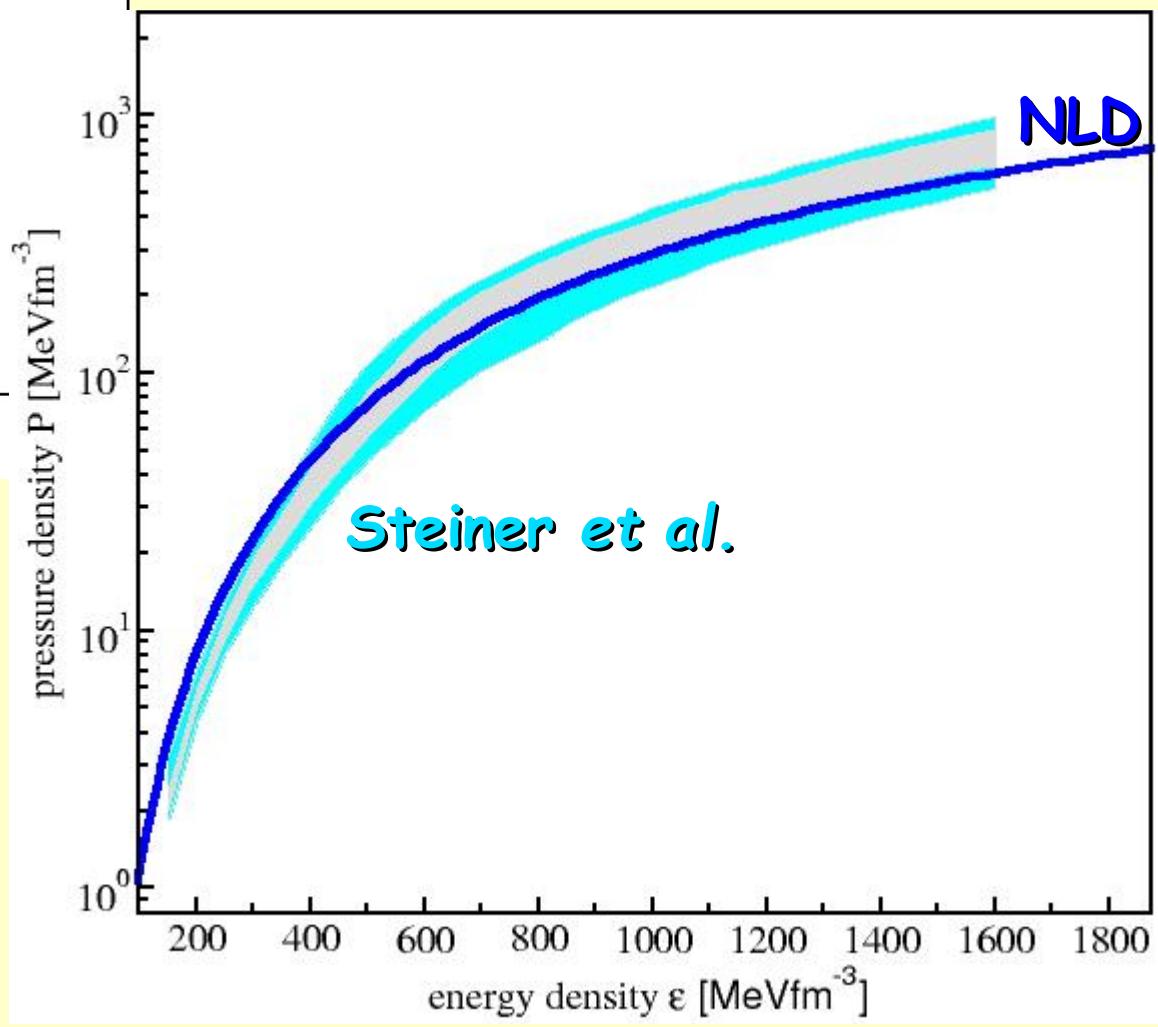


Consistent with analyses of F. Özel...

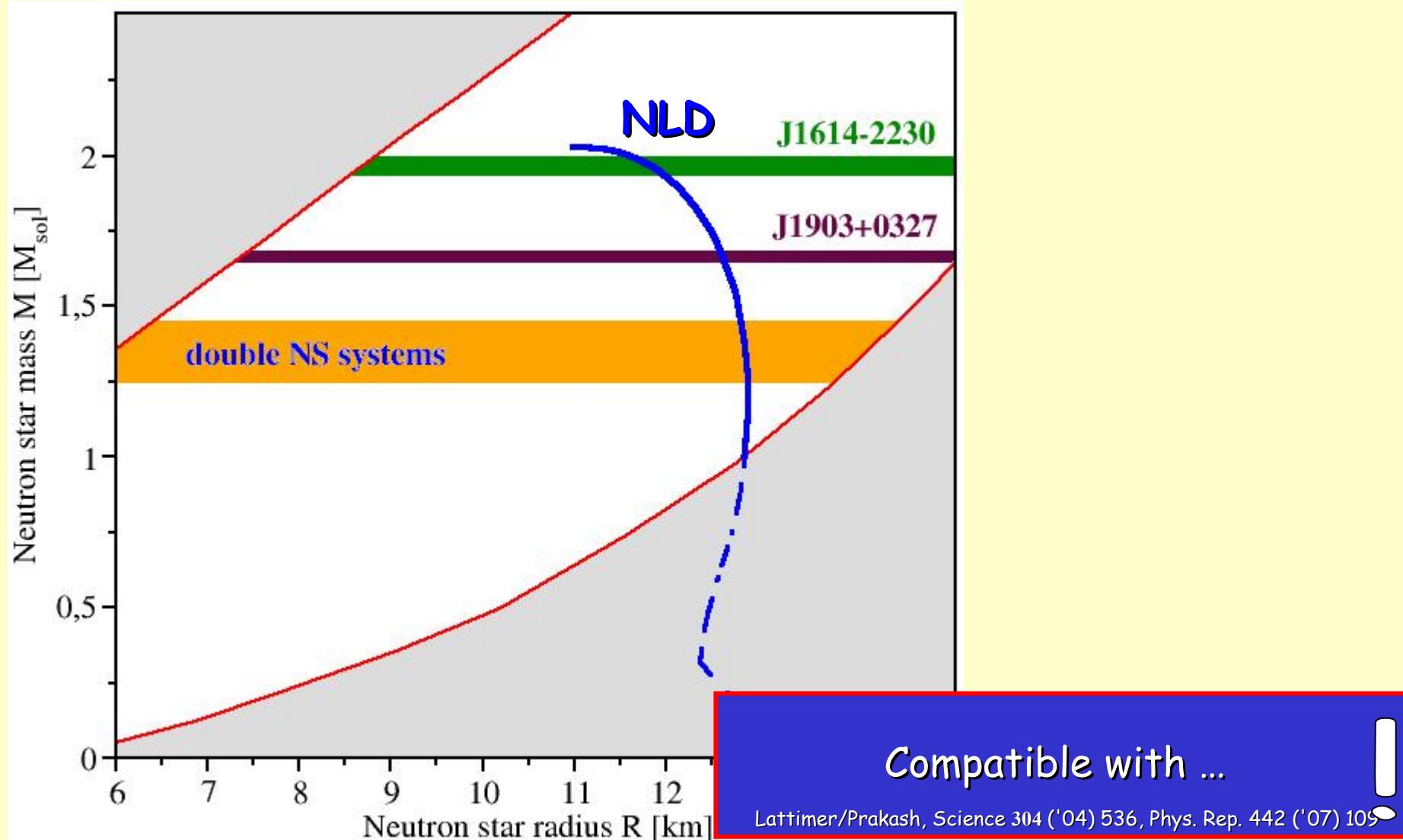
Phys. Rev. D82, 101301 (2010).

... and A.W. Steiner

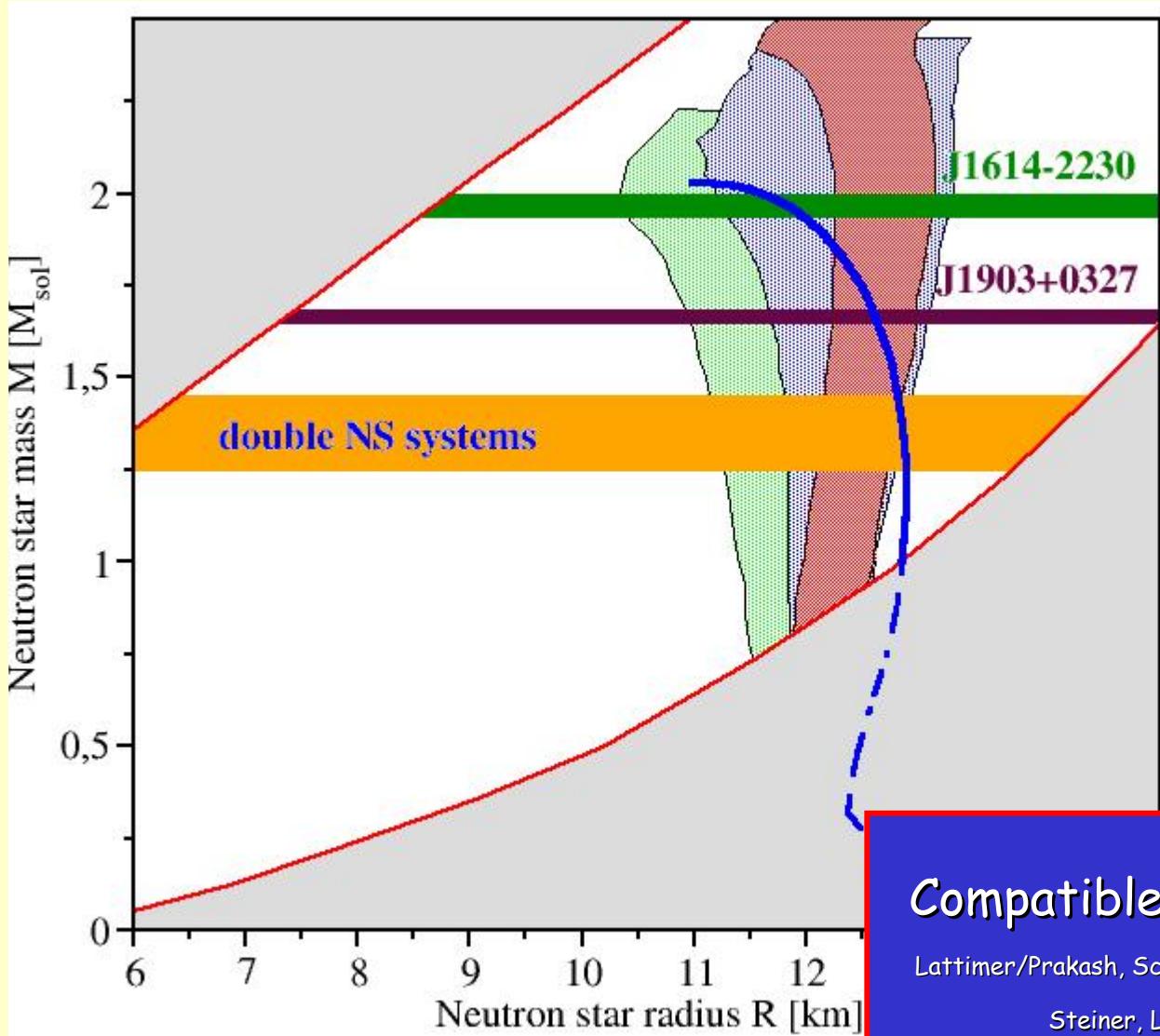
Astrophys. J. 722, 33 (2010).



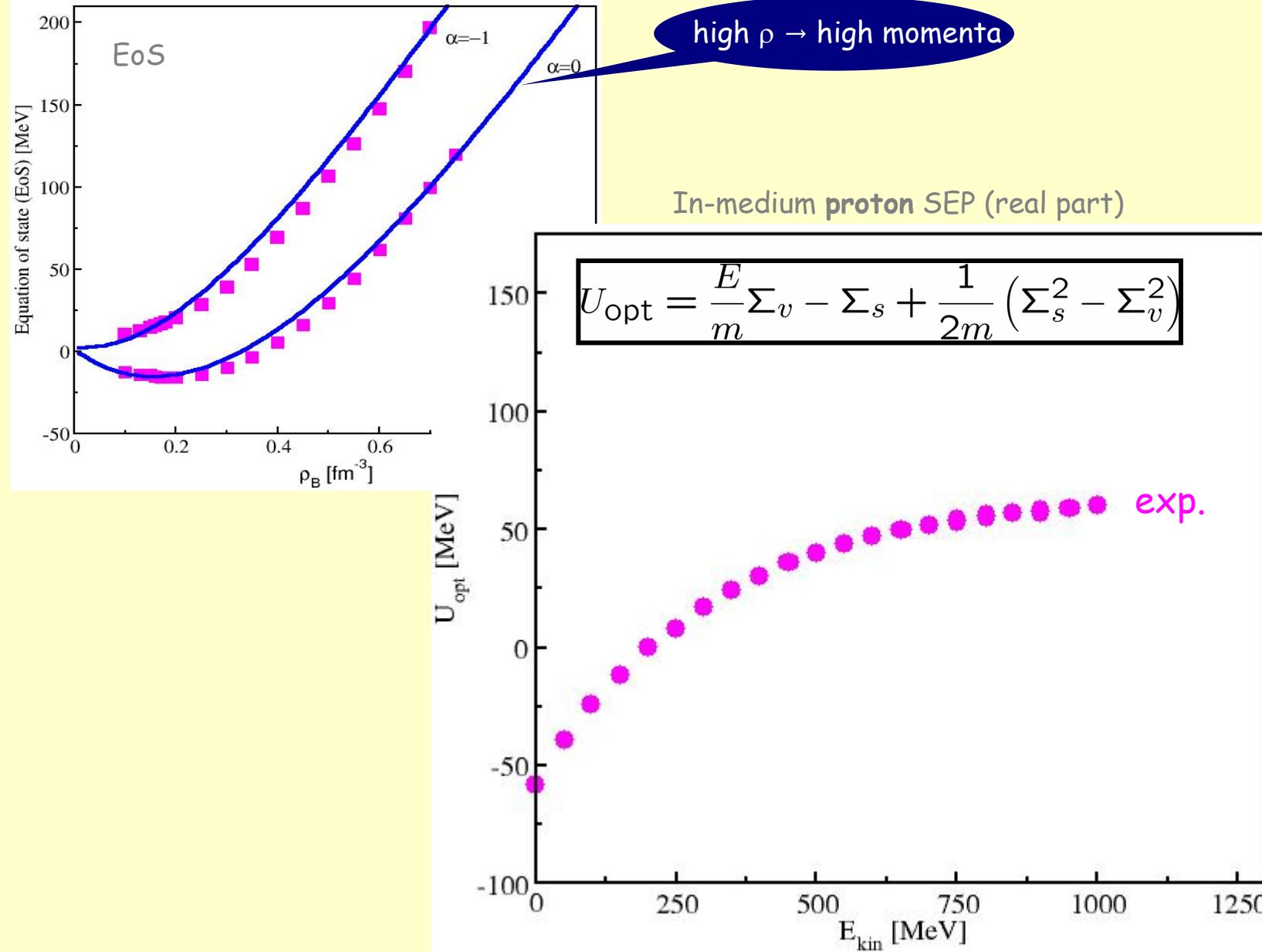
Basic properties: nuclear EoS & opt. potentials...



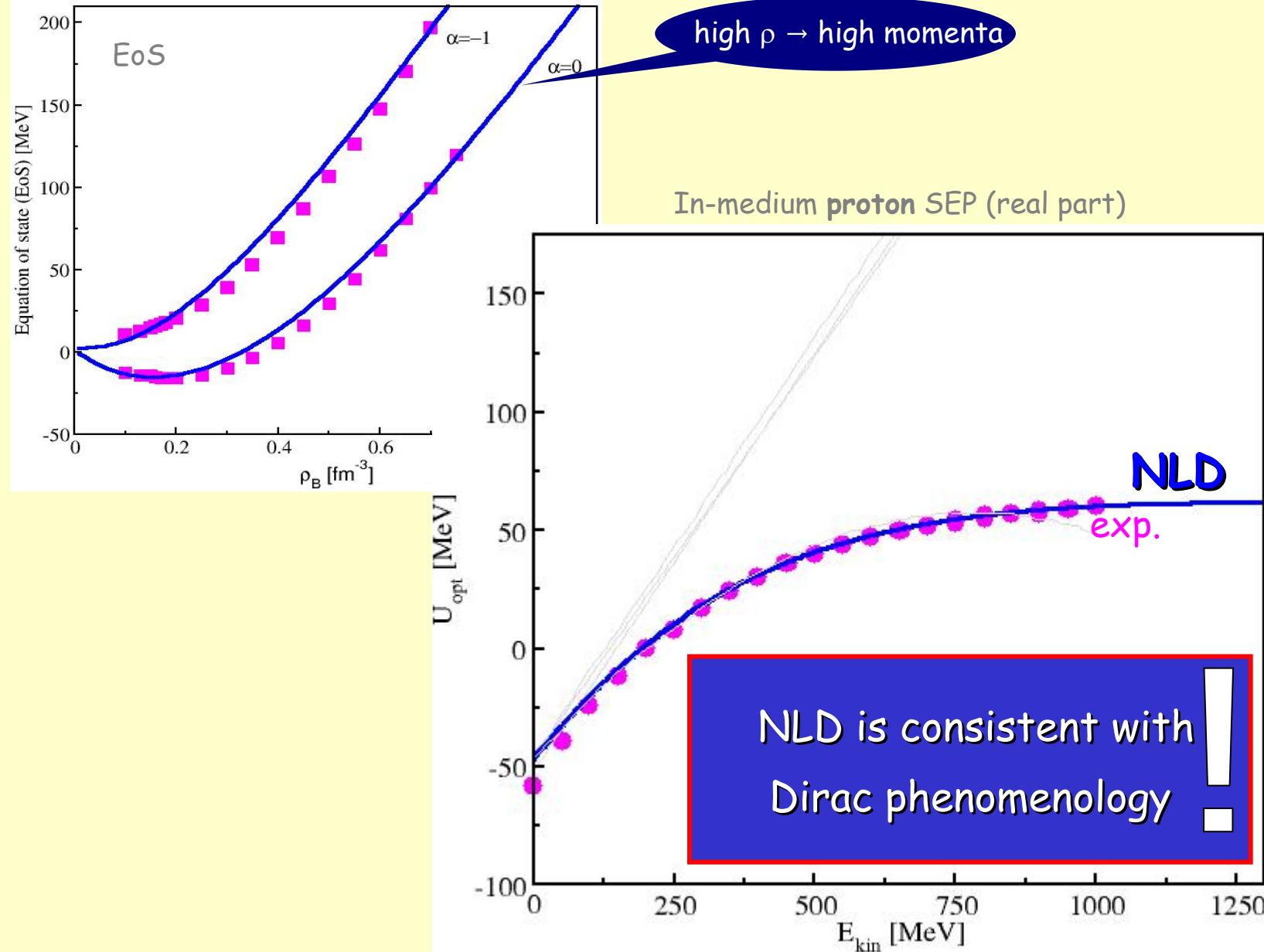
Basic properties: nuclear EoS & opt. potentials...



Basic properties: nuclear EoS & opt. potentials..

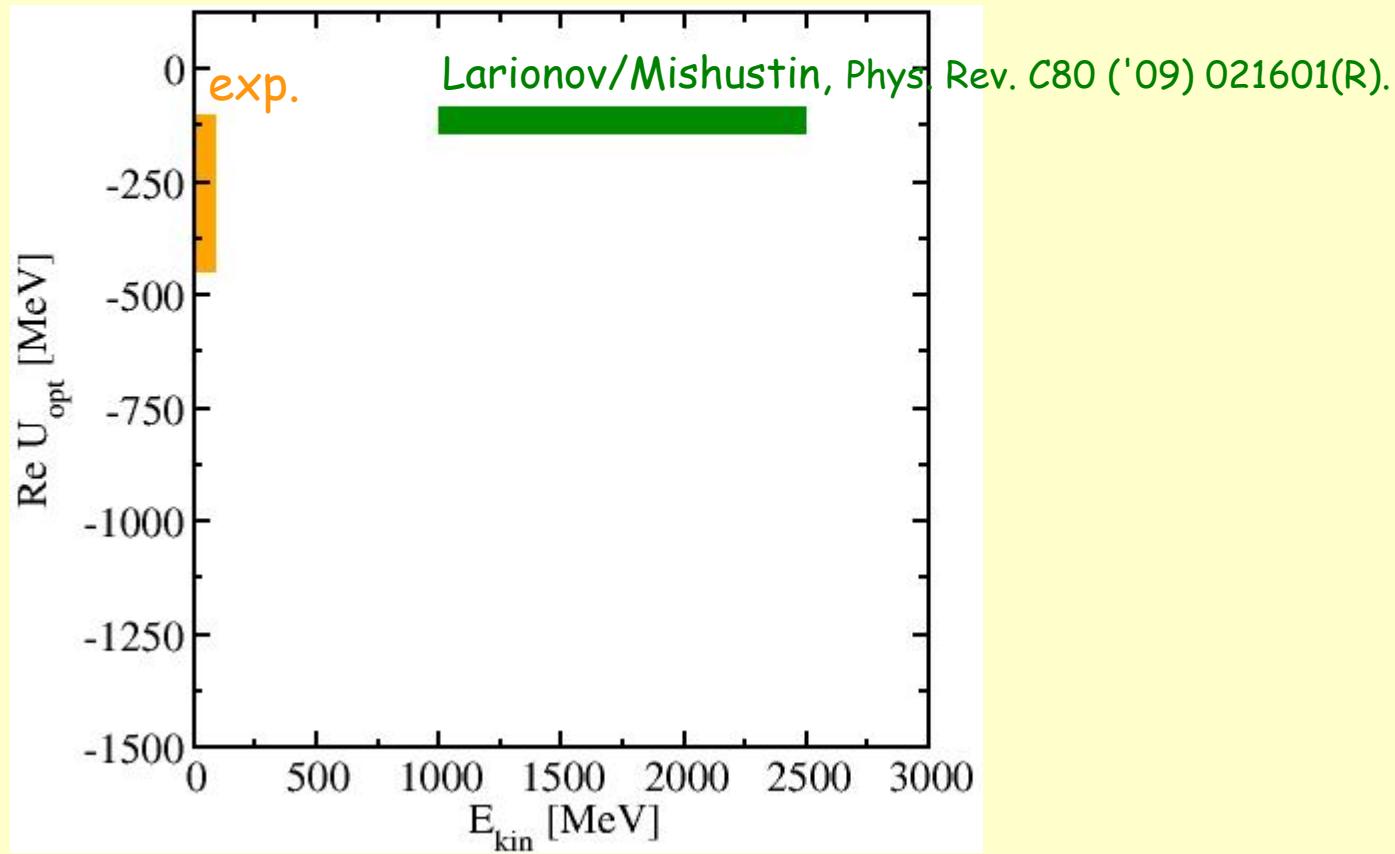


Basic properties: nuclear EoS & opt. potentials..



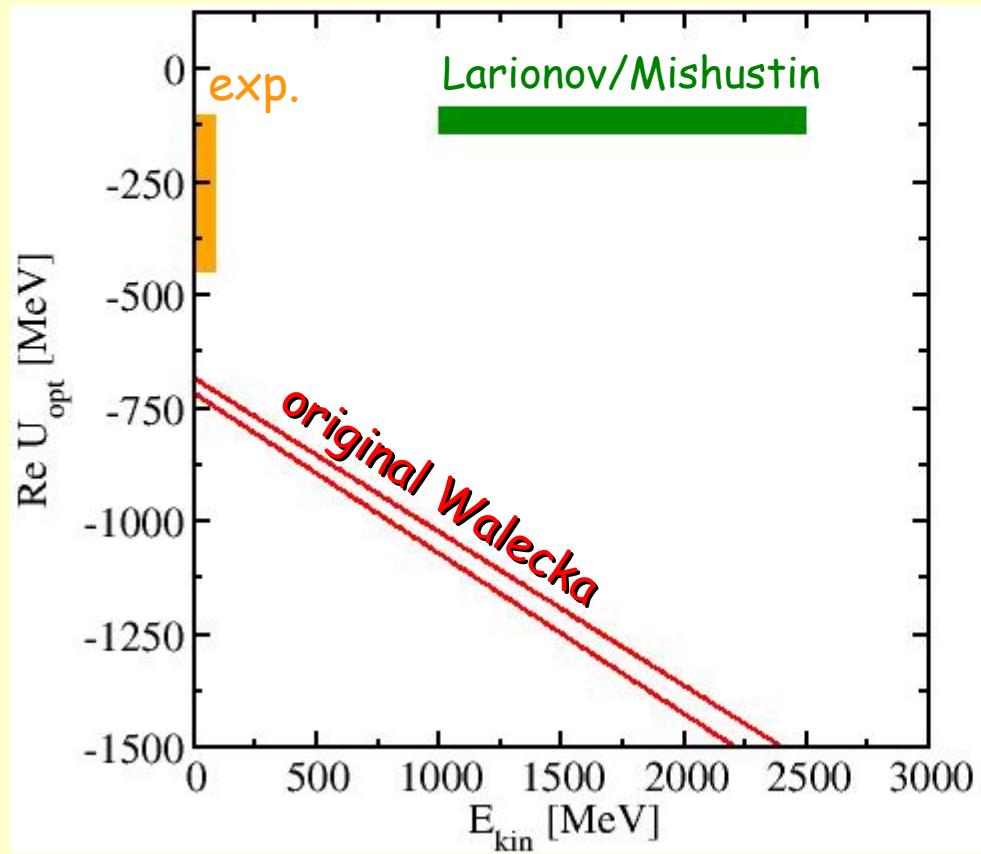
Basic properties: nuclear EoS & opt. potentials..

In-medium anti-proton SEP (real part)



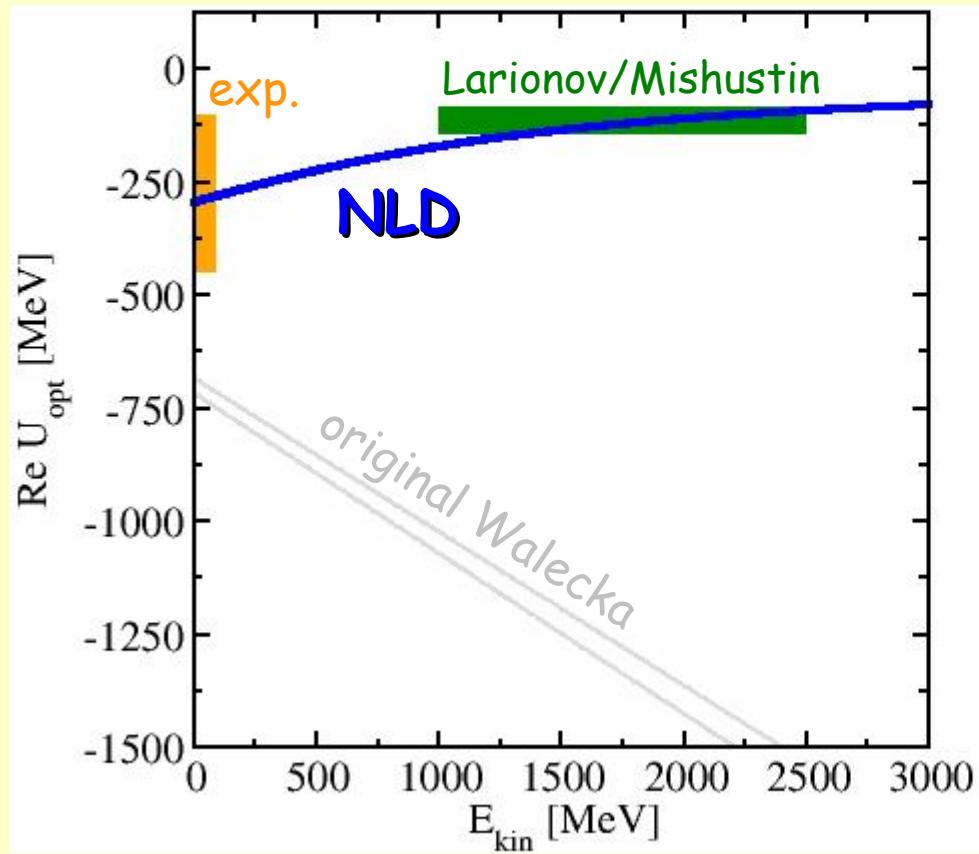
Basic properties: nuclear EoS & opt. potentials..

In-medium anti-proton SEP (real part)



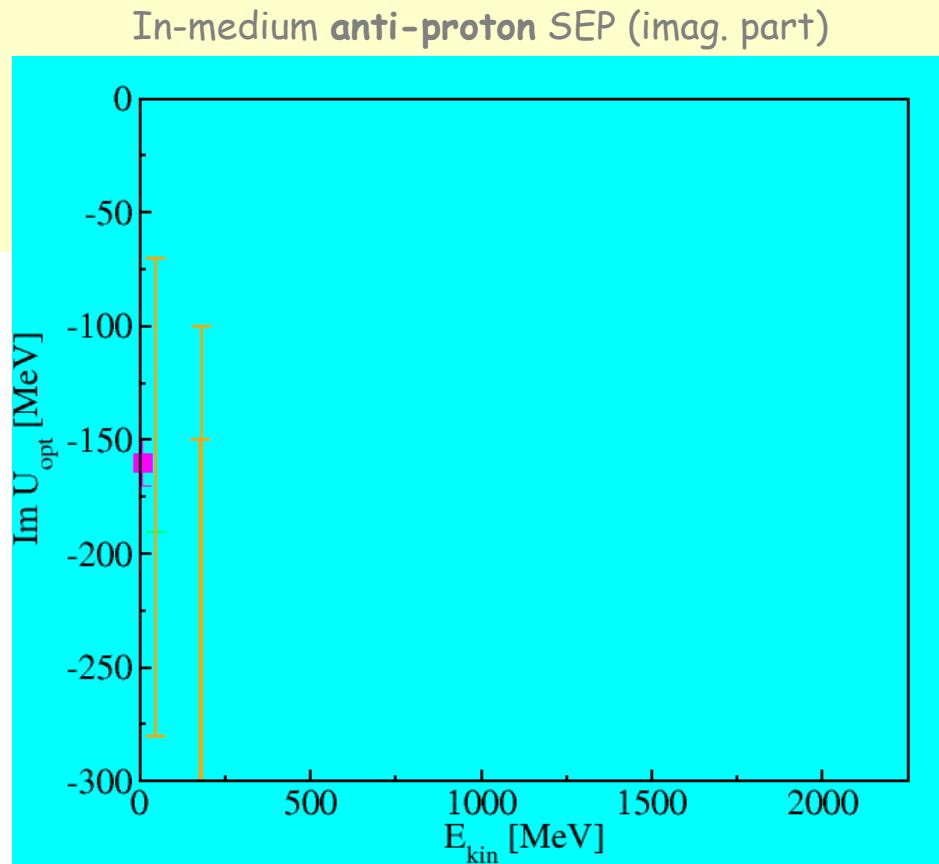
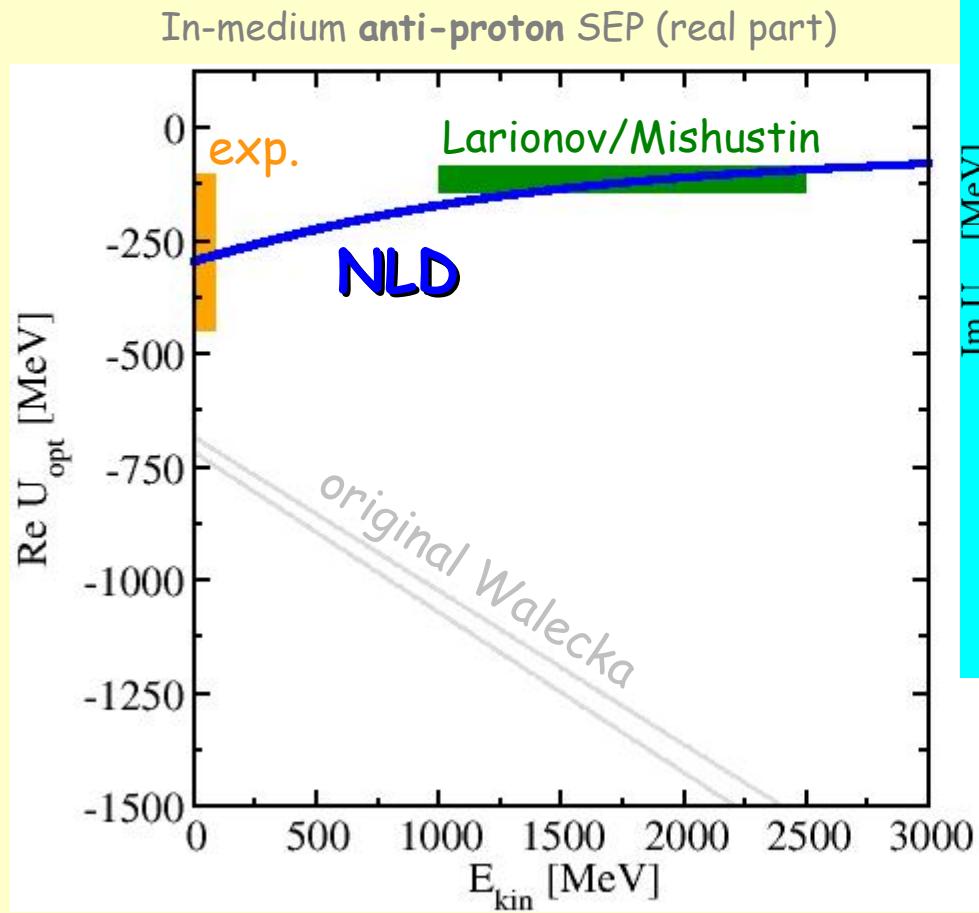
Basic properties: nuclear EoS & opt. potentials..

In-medium anti-proton SEP (real part)



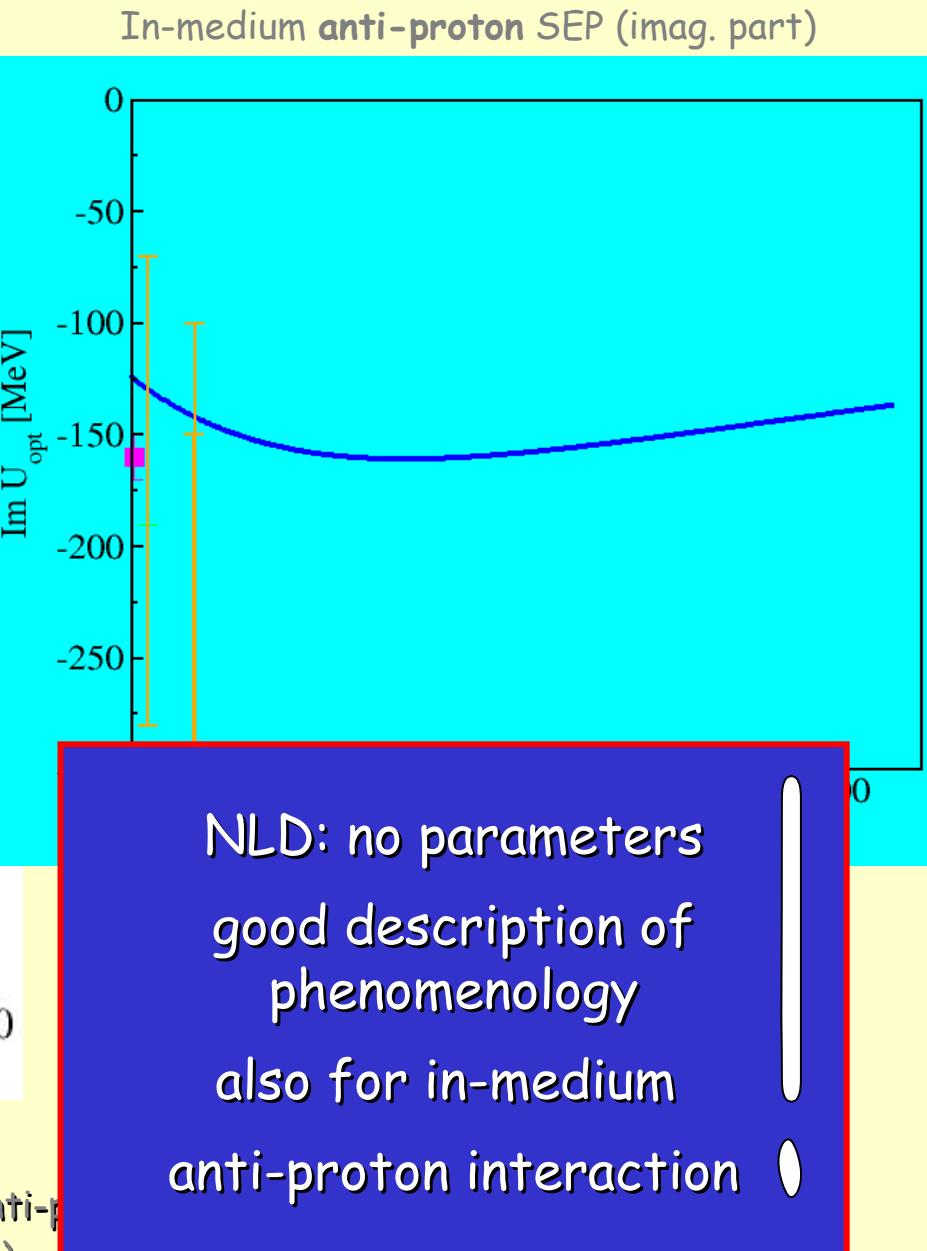
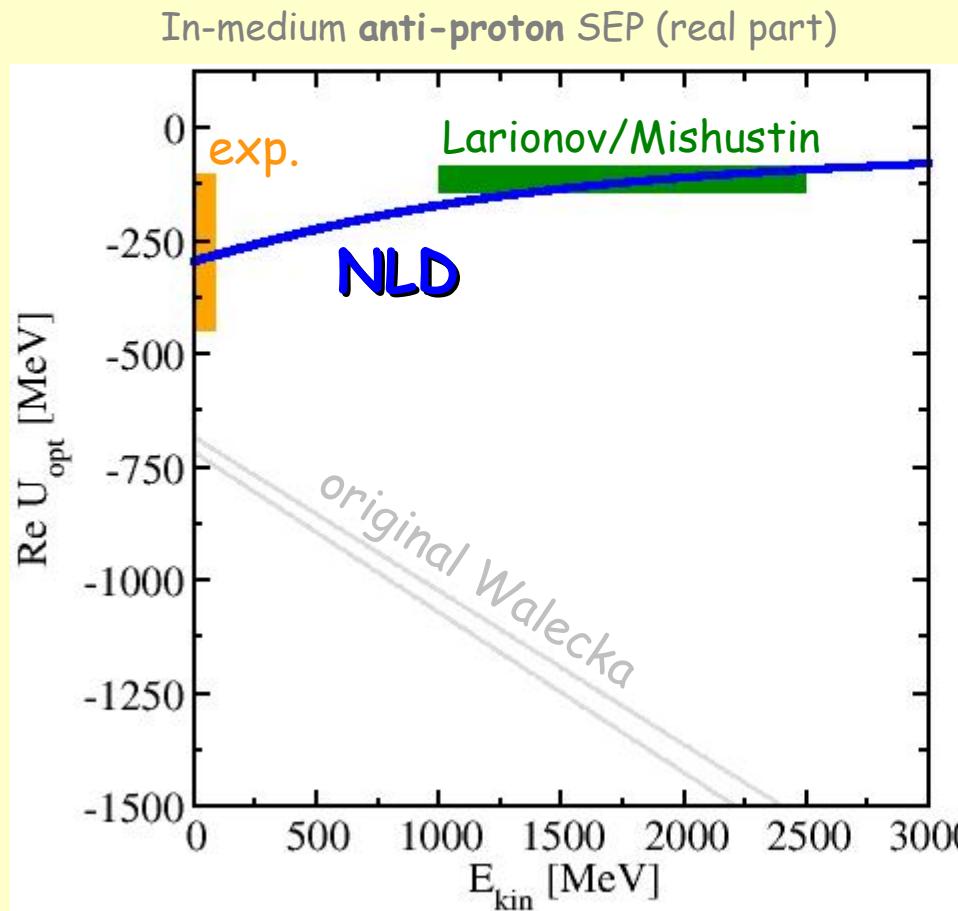
Also: NLD provides the imaginary part of SEP for anti-proton in-medium interactions using dispersion relation (without subtractions)

Basic properties: nuclear EoS & opt. potentials..



Also: NLD provides the imaginary part of SEP for anti-proton in-medium interactions using dispersion relation (without subtractions)

Basic properties: nuclear EoS & opt. potentials..

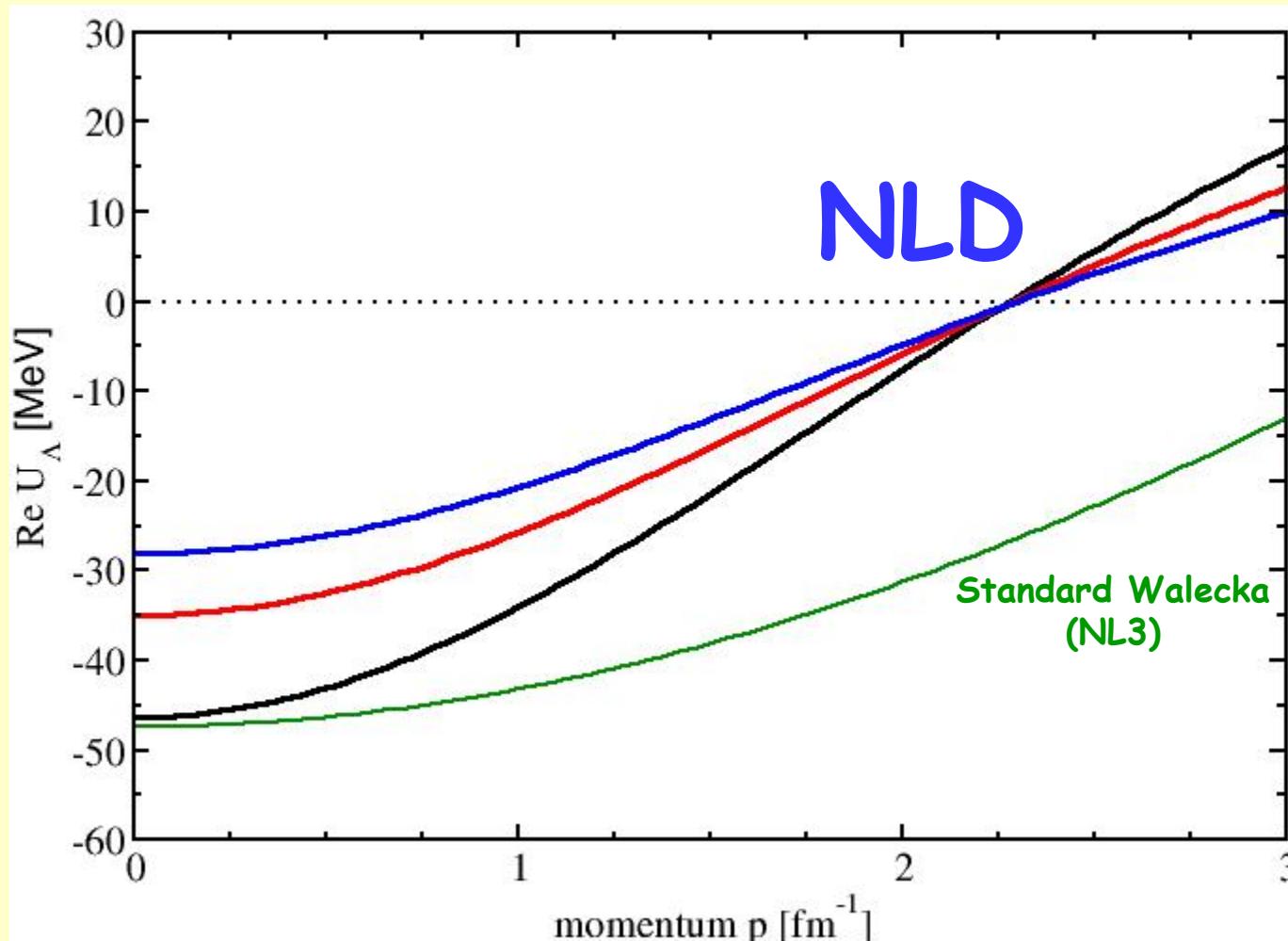


Also: NLD provides the imaginary part of SEP for anti- p using dispersion relation (without subtractions)

Hyperon properties: optical potentials...

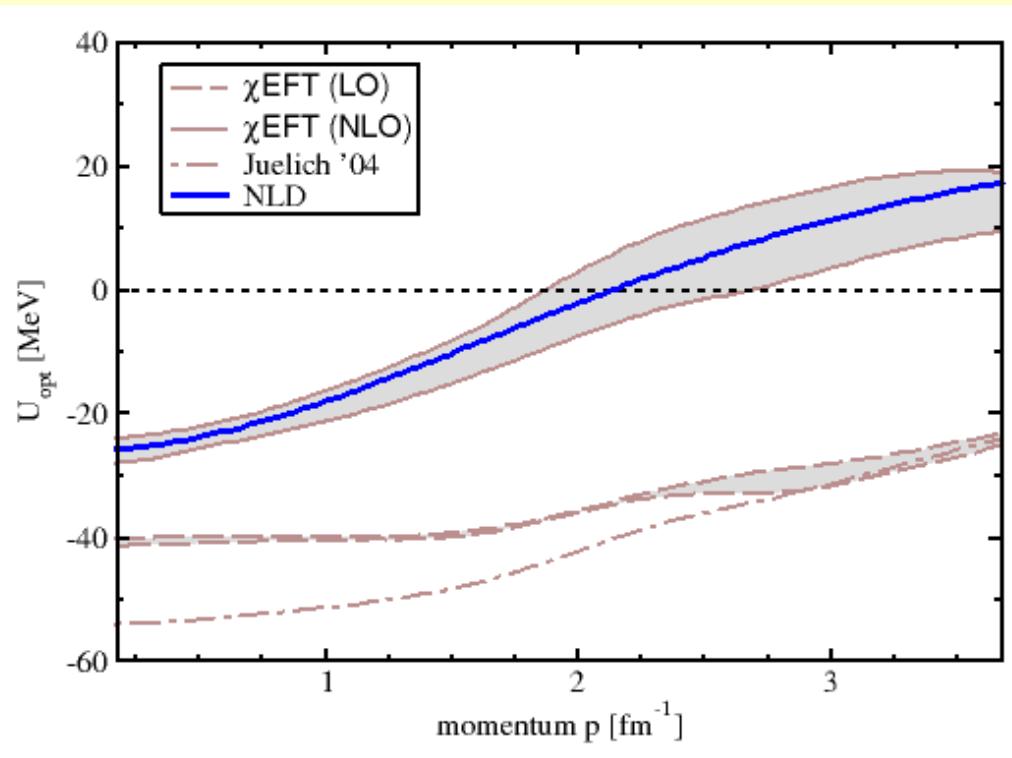
NLD + SU(6) for standard meson-nucleon couplings

Hyperon cut-off regulates MDI



Hyperon properties: Λ -optical potentials...

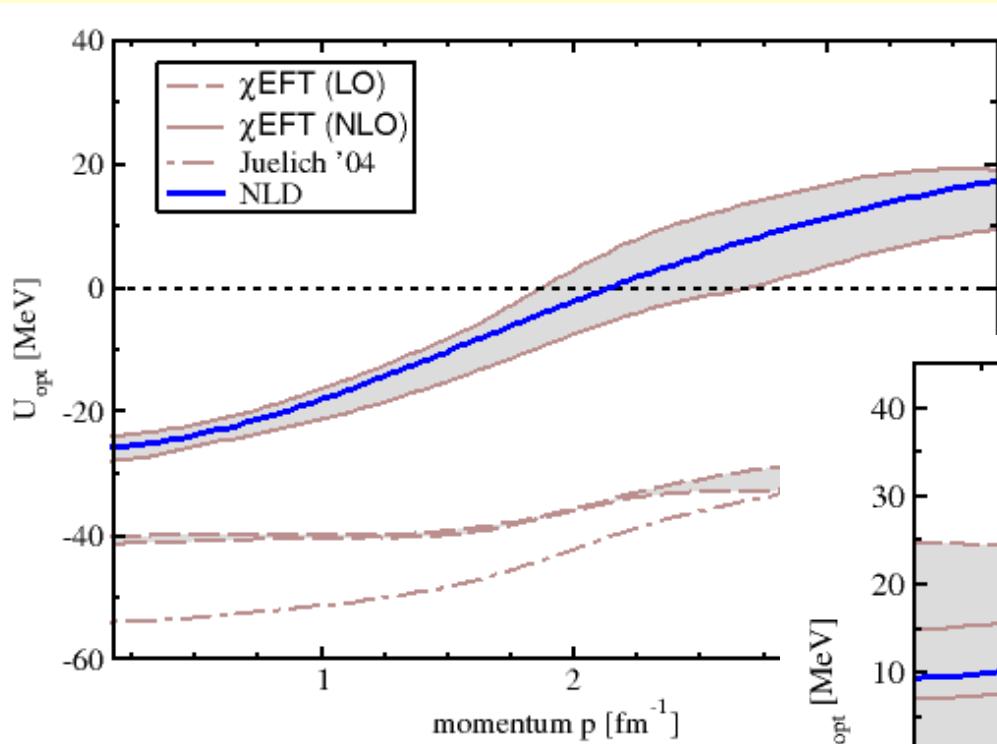
SU(6) for standard meson-nucleon couplings + NLD (monopole forms)
SNM, saturation density, adjust to χ EFT



Λ Uopt: cut-offs ~ 0.75 GeV

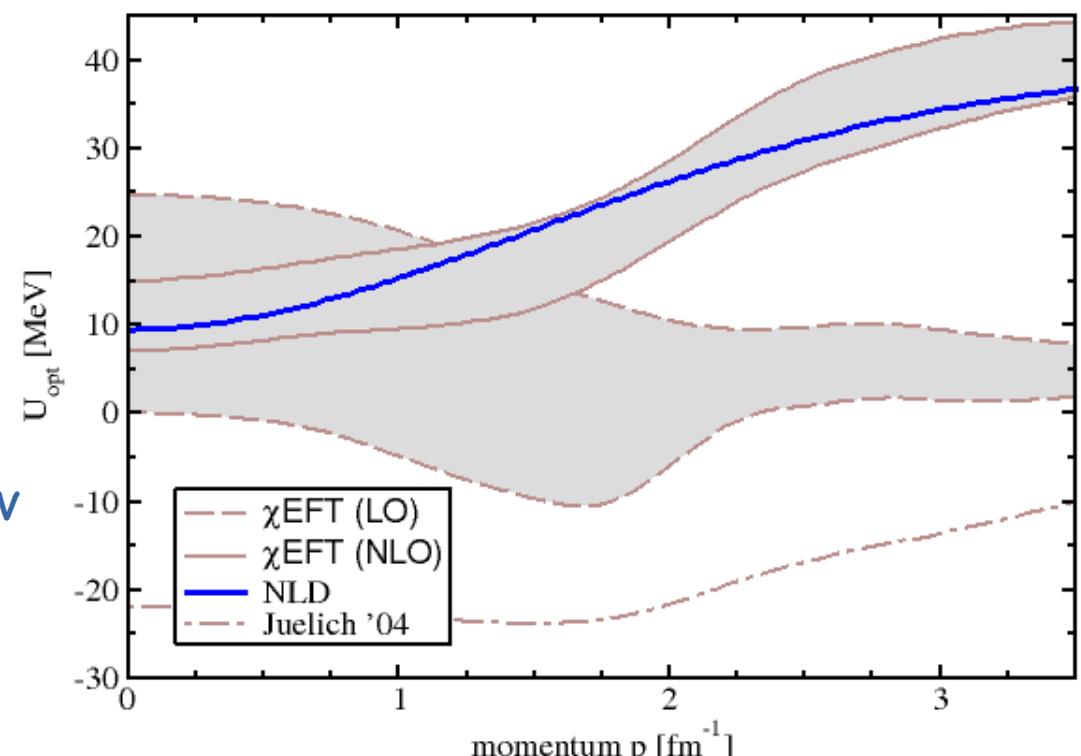
Hyperon properties: Λ, Σ -optical potentials...

SU(6) for standard meson-nucleon couplings + NLD (monopole forms)
SNM, saturation density, adjust to χ EFT



χ EFT: EPJA52 (2016) 15

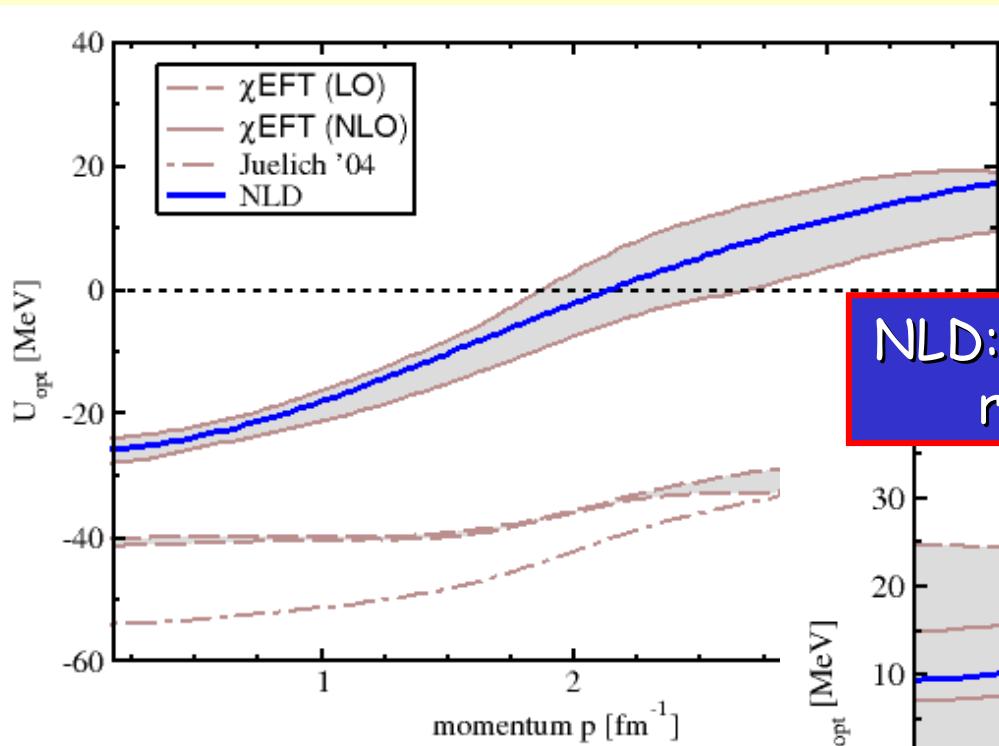
Λ Uopt: cut-offs ~ 0.78 GeV



Σ Uopt: cut-offs ~ 0.8 GeV

Hyperon properties: Λ, Σ -optical potentials...

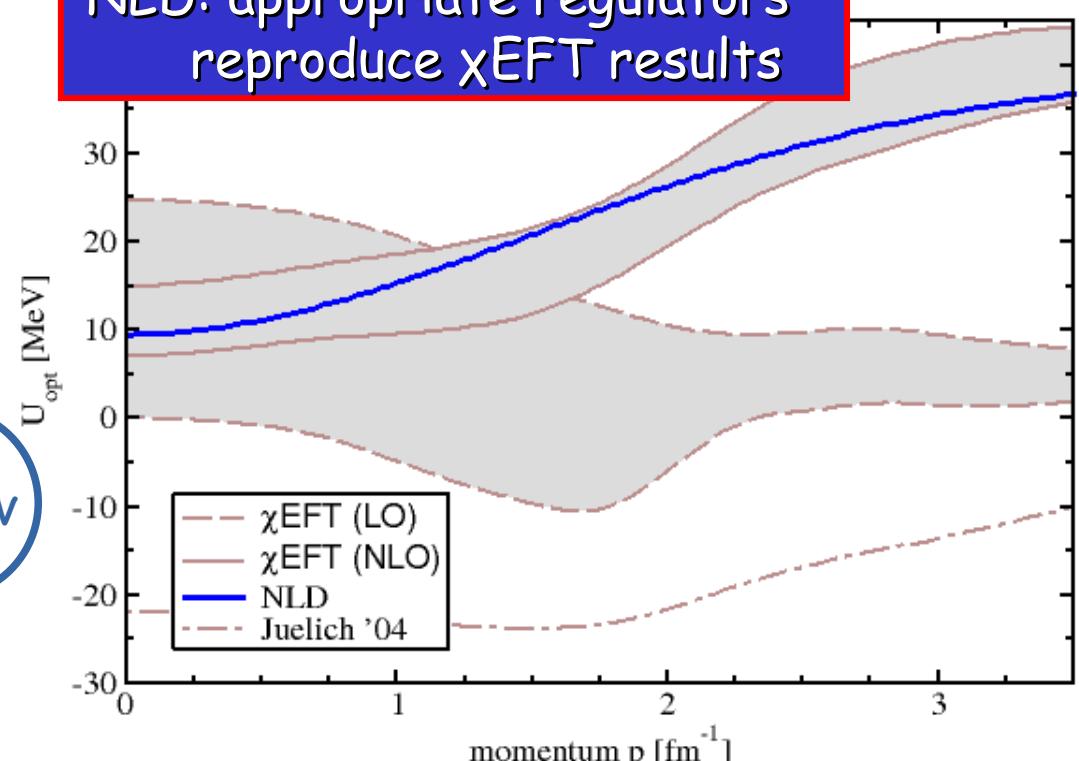
SU(6) for standard meson-nucleon couplings + NLD (monopole forms)
SNM, saturation density, adjust to χ EFT



χ EFT: EPJA52 (2016) 15

Λ Uopt: cut-offs ~ 0.78 GeV

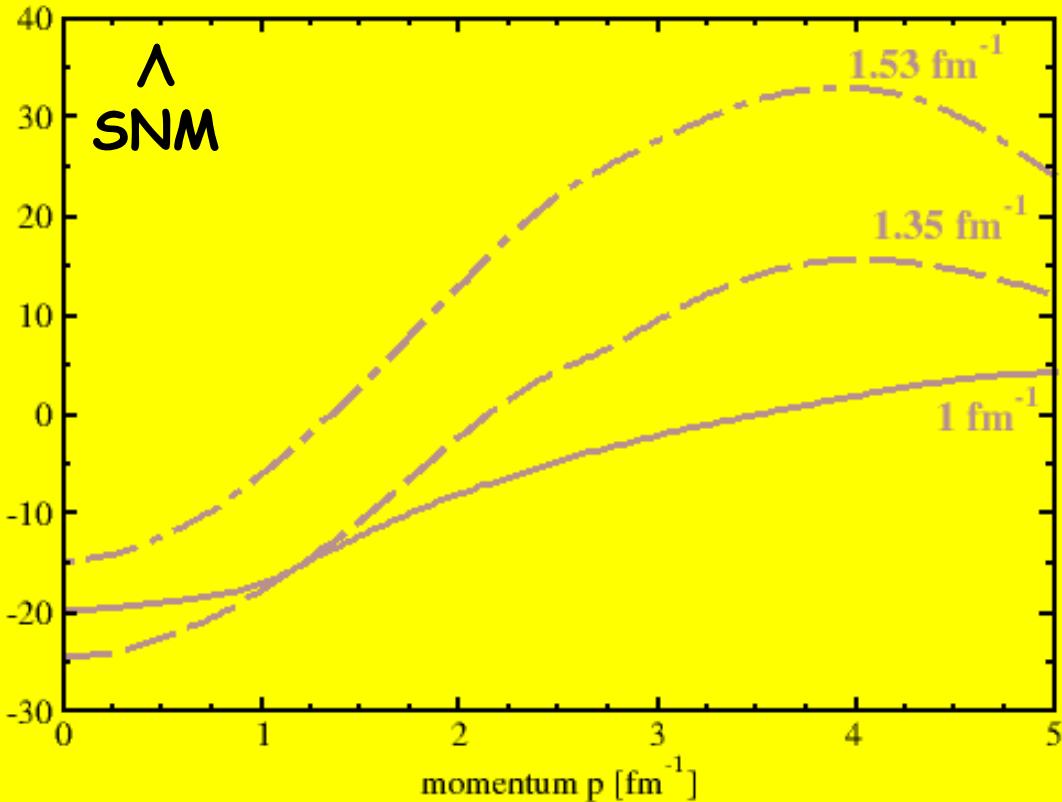
NLD: appropriate regulators
reproduce χ EFT results



Σ Uopt: cut-offs ~ 0.8 GeV

Hyperon properties: Λ -optical potentials...

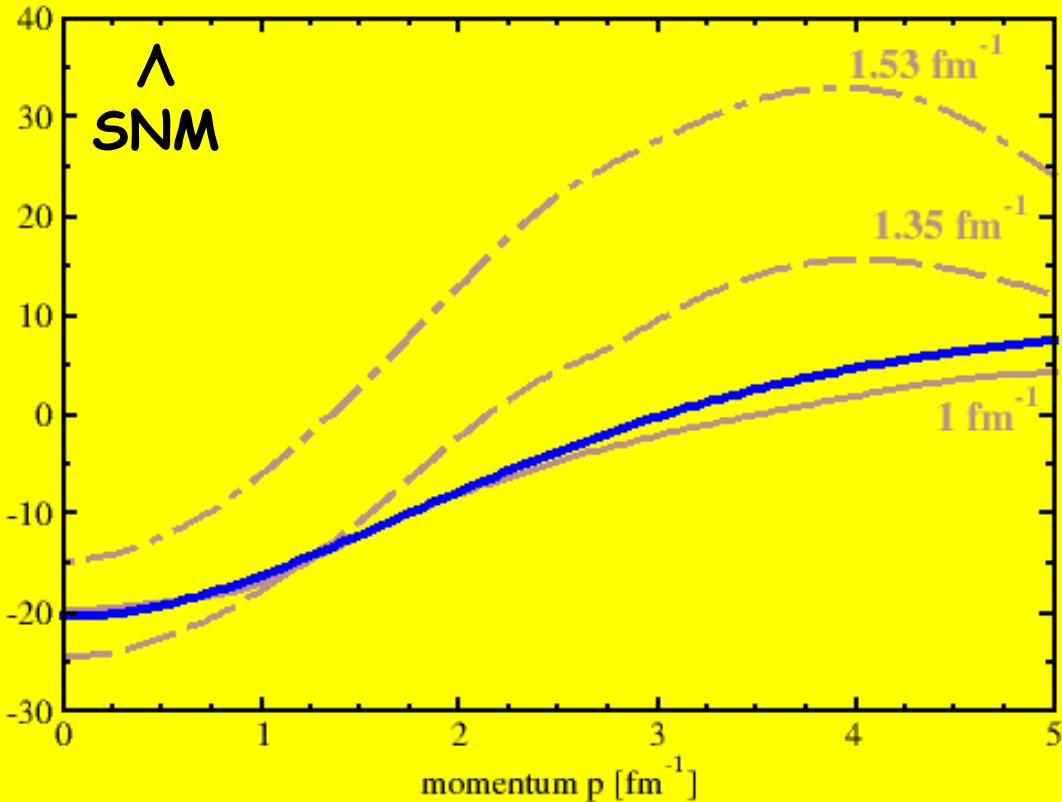
NLD predictions: density & momentum dependence



\times EFT: EPJA52 (2016) 15

Hyperon properties: Λ -optical potentials...

NLD predictions: density & momentum dependence

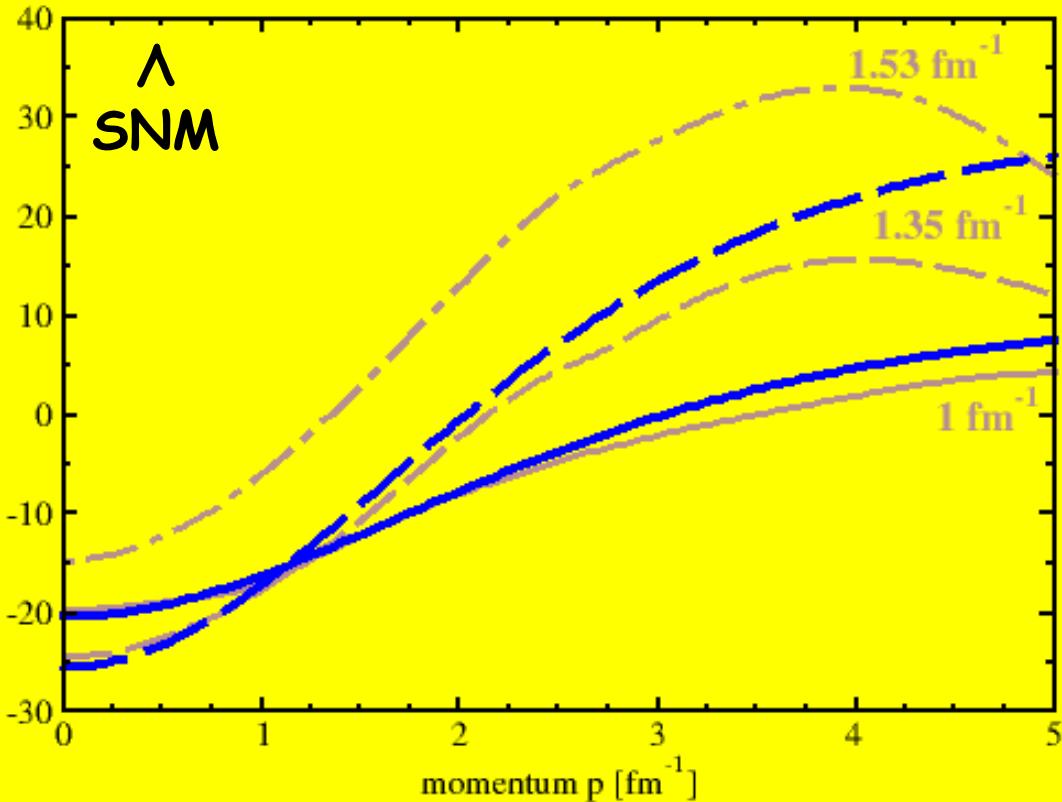


\times EFT: EPJA52 (2016) 15

NLD results

Hyperon properties: Λ -optical potentials...

NLD predictions: density & momentum dependence

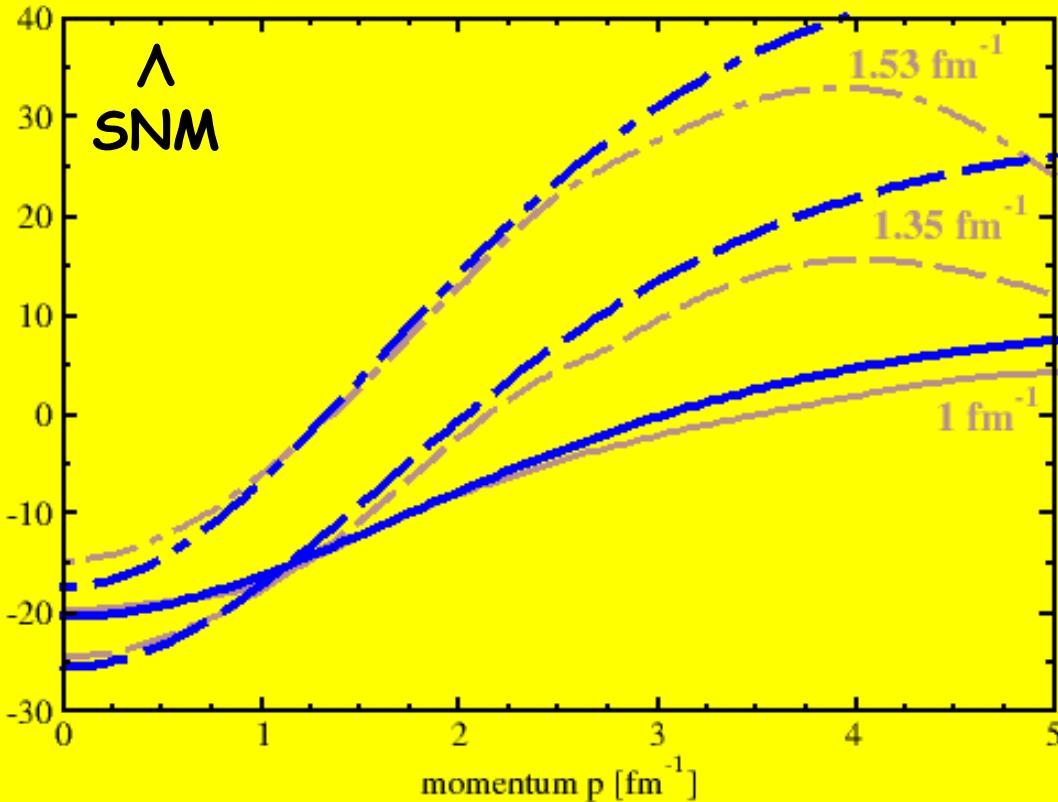


\times EFT: EPJA52 (2016) 15

NLD results

Hyperon properties: Λ -optical potentials...

NLD predictions: density & momentum dependence

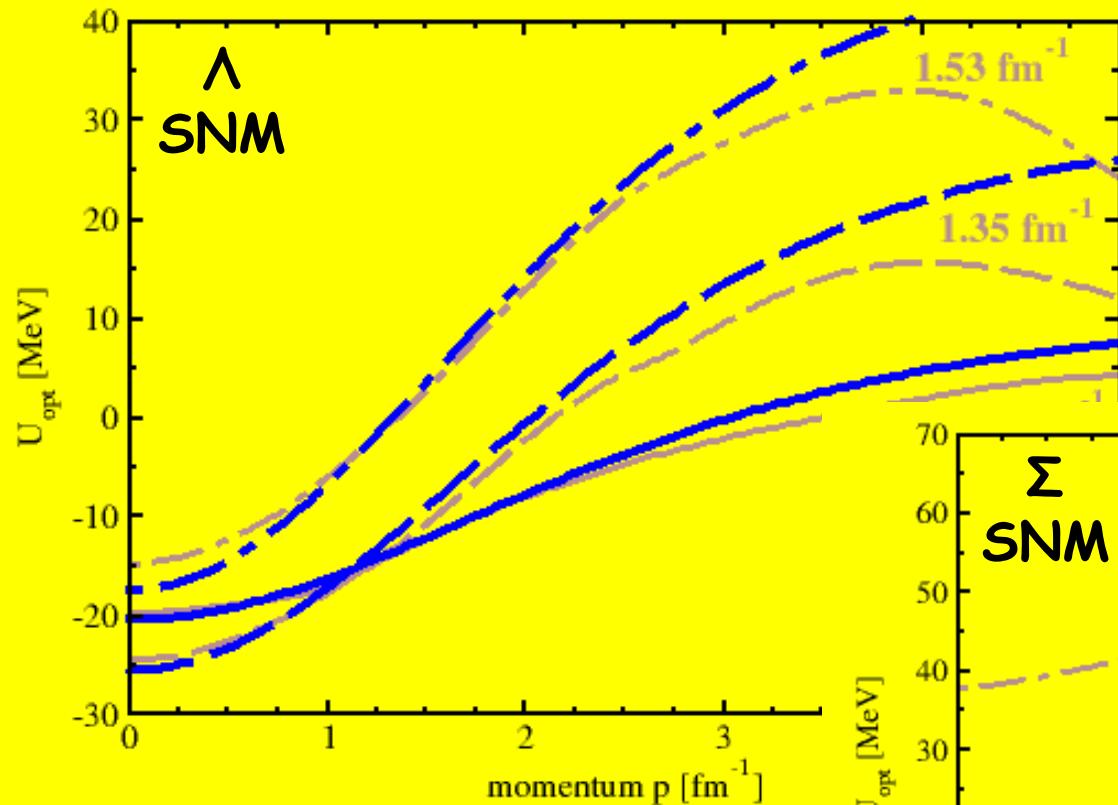


\times EFT: EPJA52 (2016) 15

NLD results

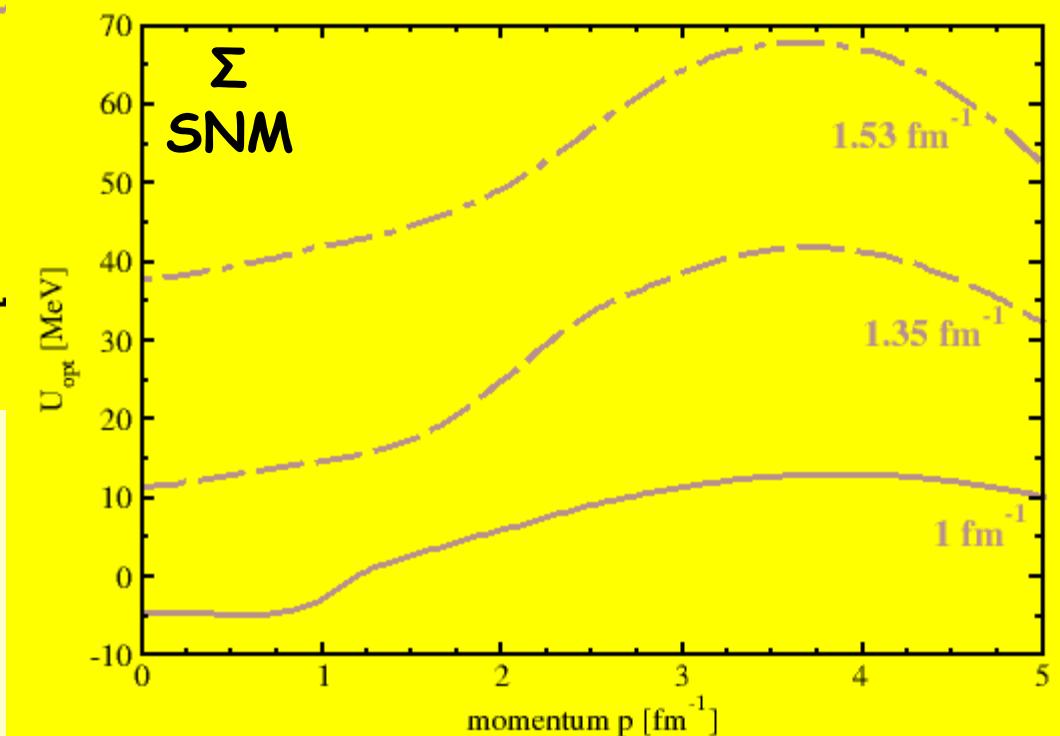
Hyperon properties: Λ, Σ -optical potentials...

NLD predictions: density & momentum dependence



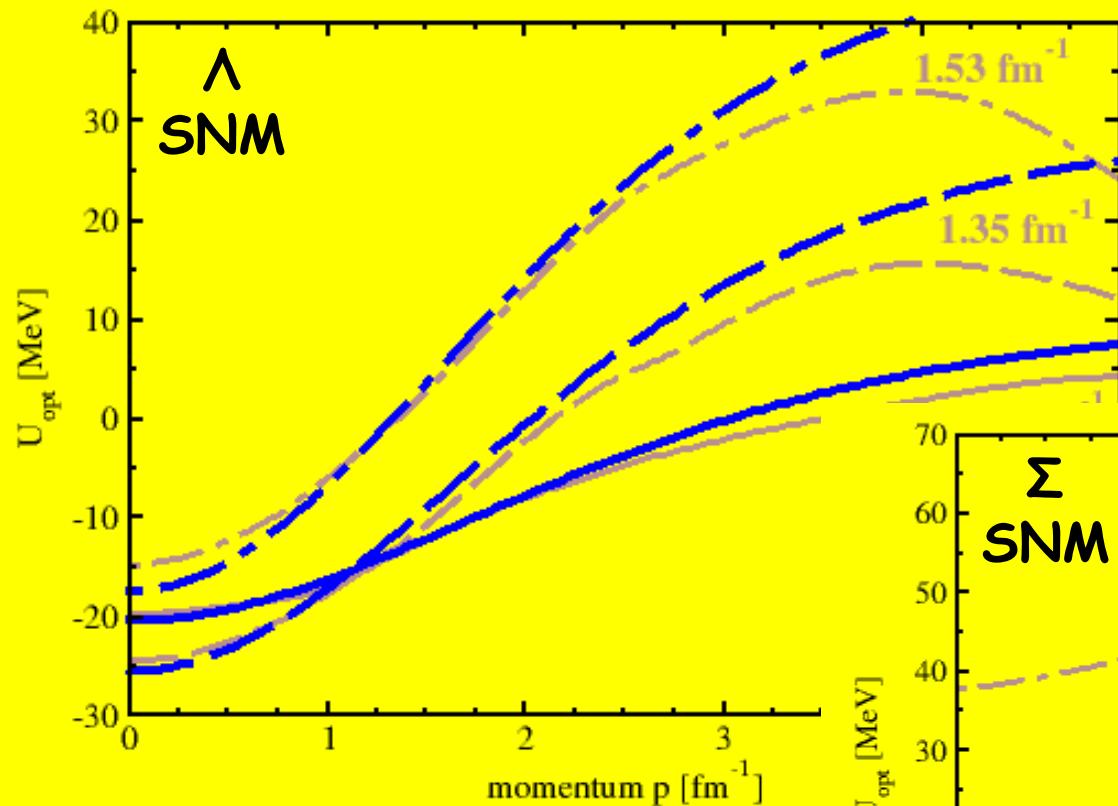
\times EFT: EPJA52 (2016) 15

NLD results



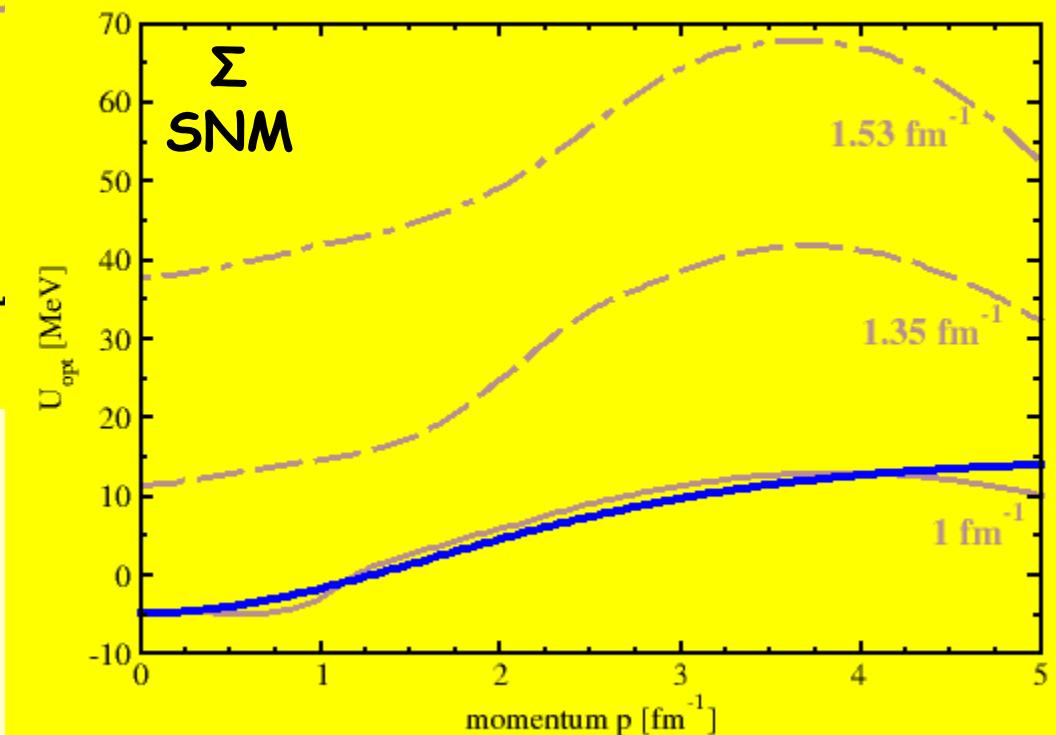
Hyperon properties: Λ, Σ -optical potentials...

NLD predictions: density & momentum dependence



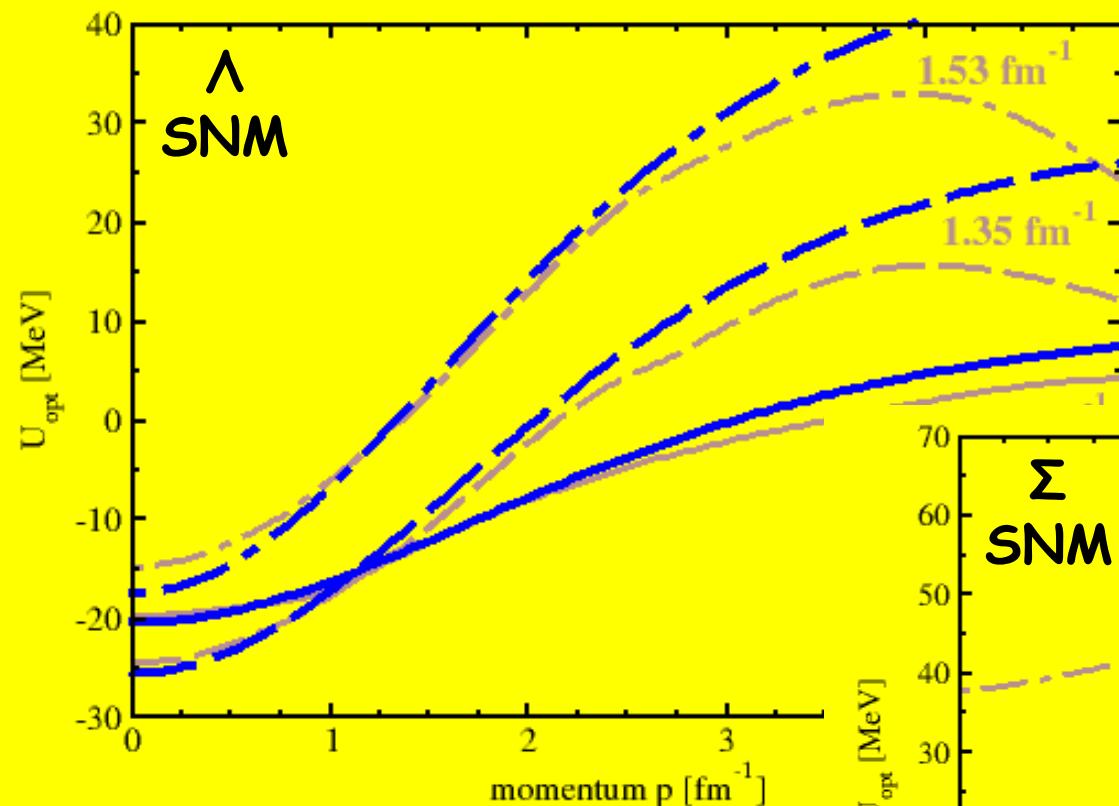
\times EFT: EPJA52 (2016) 15

NLD results



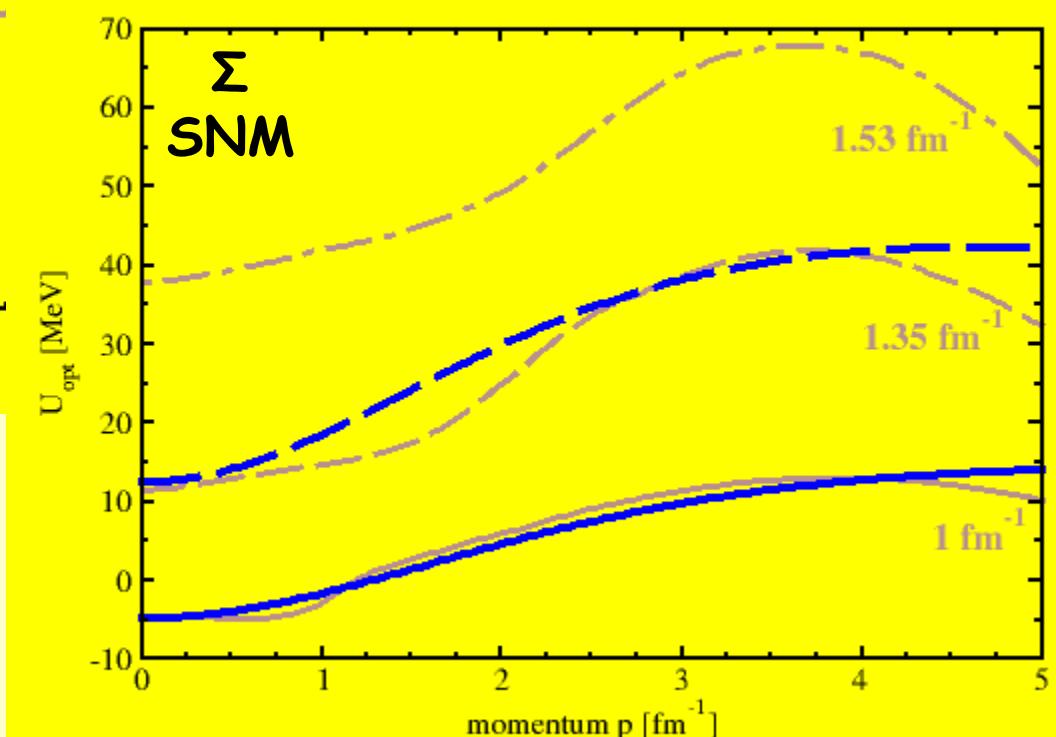
Hyperon properties: Λ, Σ -optical potentials...

NLD predictions: density & momentum dependence



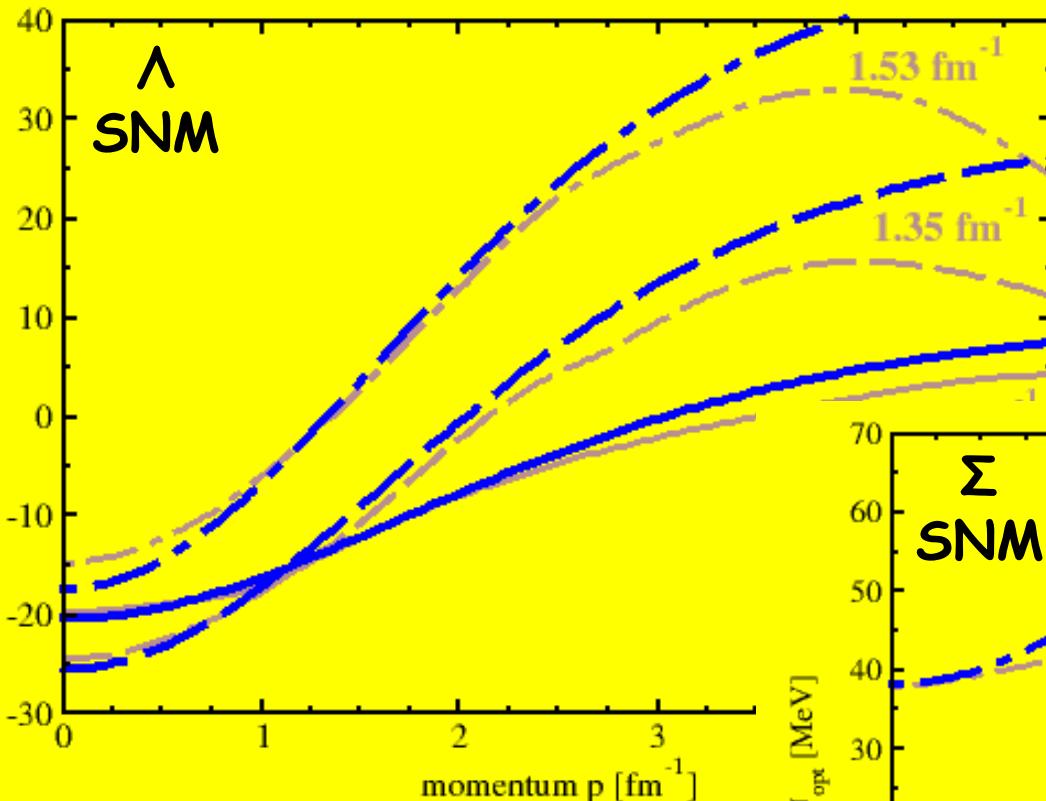
\times EFT: EPJA52 (2016) 15

NLD results



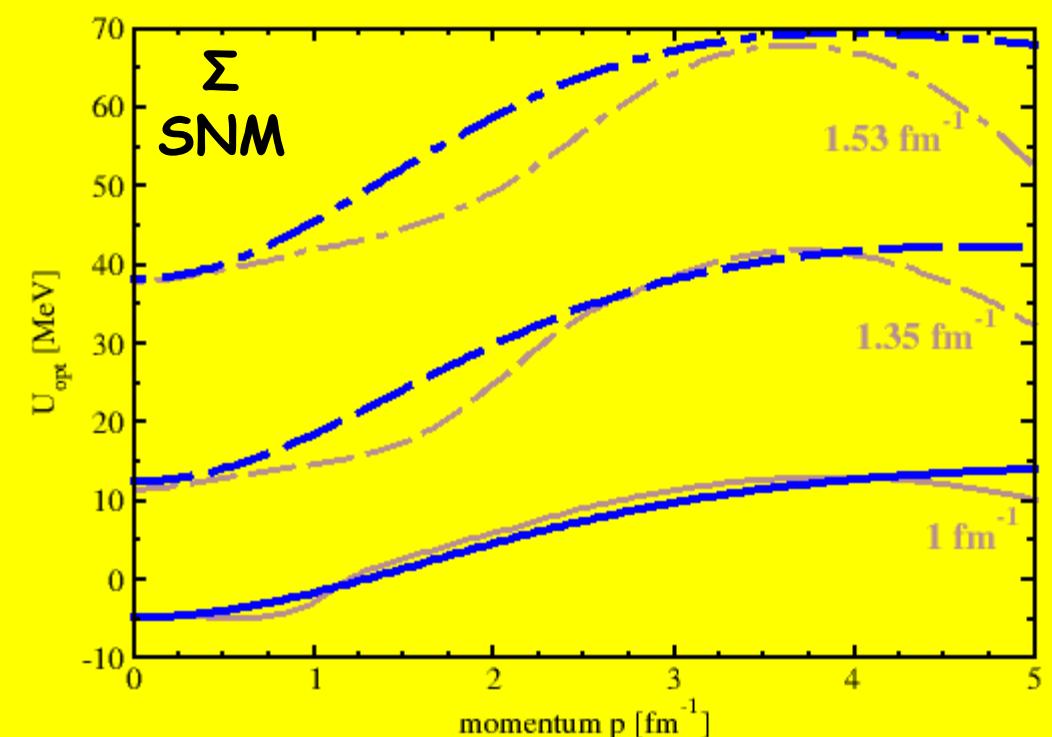
Hyperon properties: Λ, Σ -optical potentials...

NLD predictions: density & momentum dependence



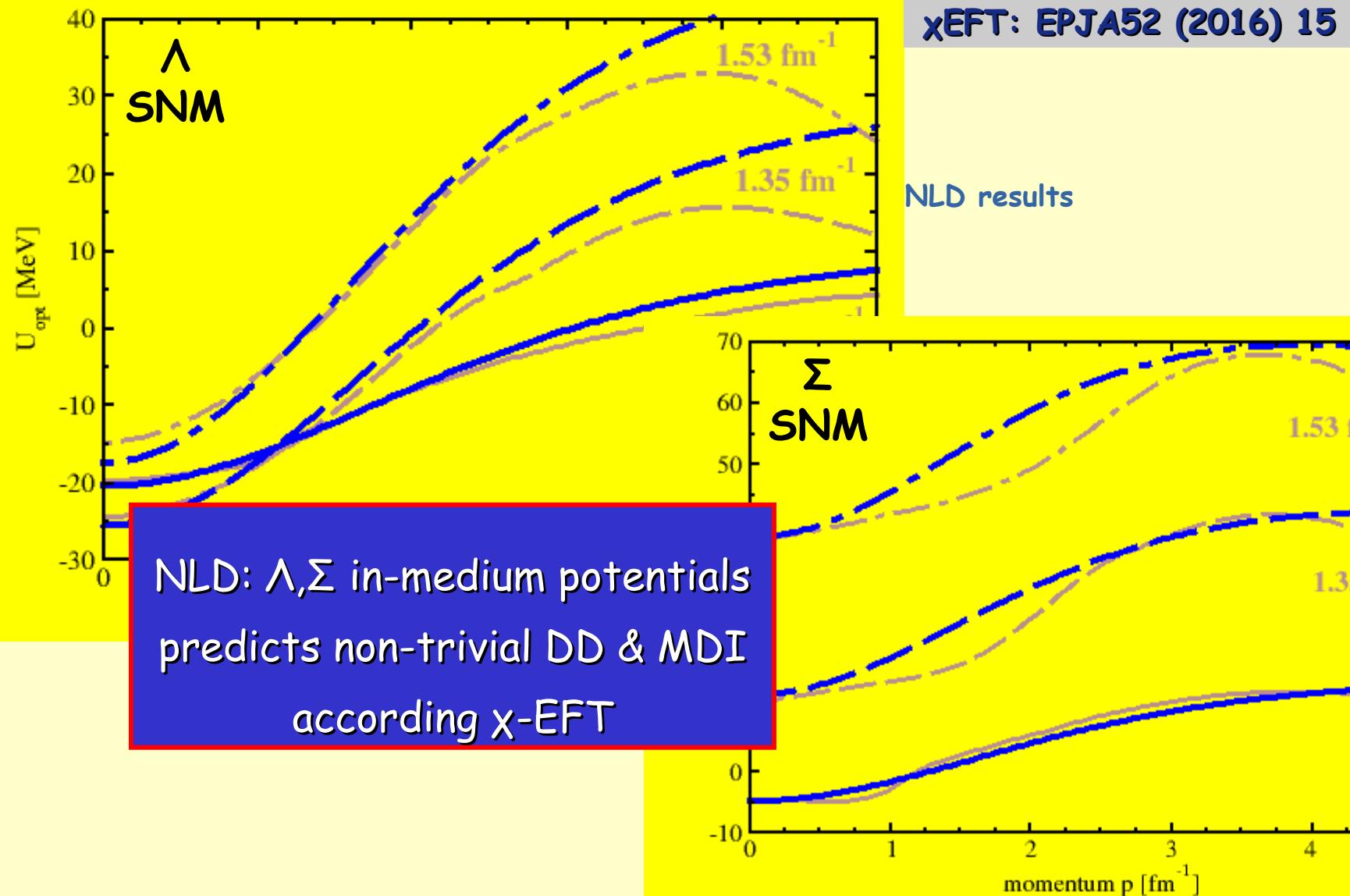
\times EFT: EPJA52 (2016) 15

NLD results



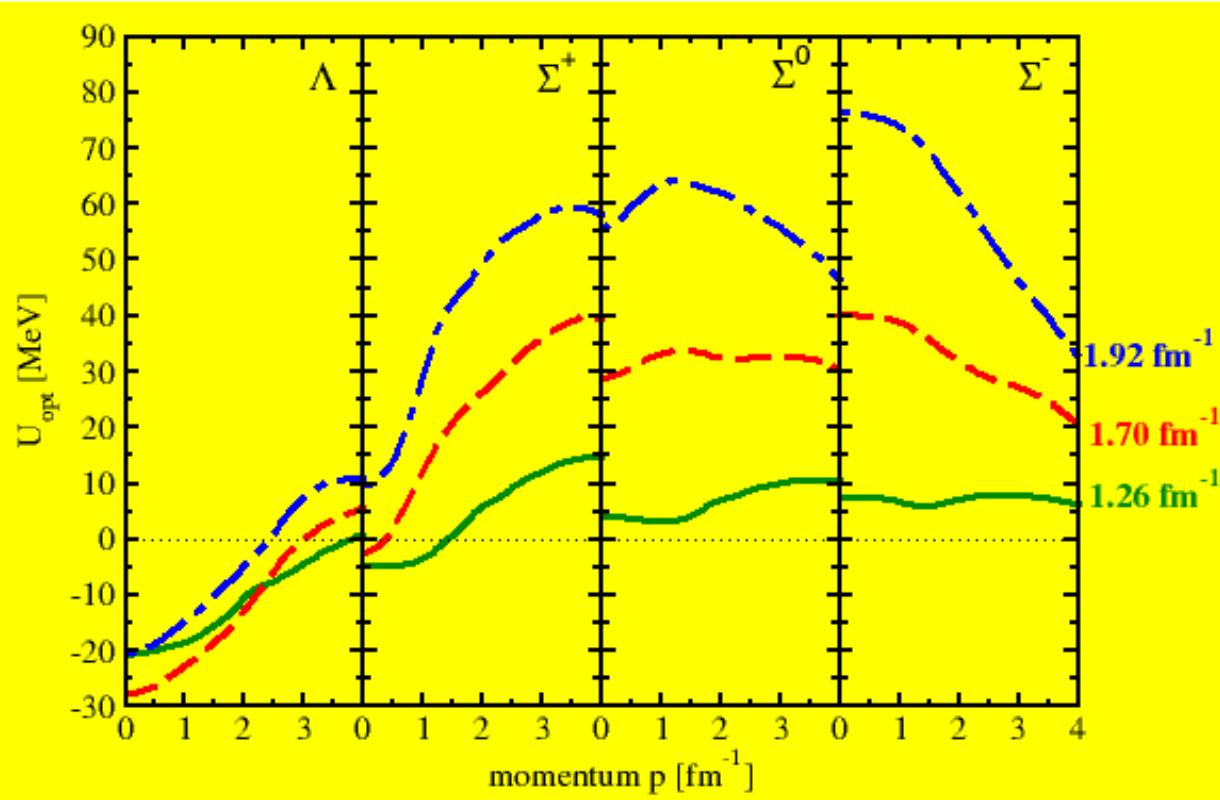
Hyperon properties: Λ, Σ -optical potentials...

NLD predictions: density & momentum dependence



Hyperon properties: Λ, Σ -optical potentials...

NLD predictions density & momentum dependence
pure neutron matter



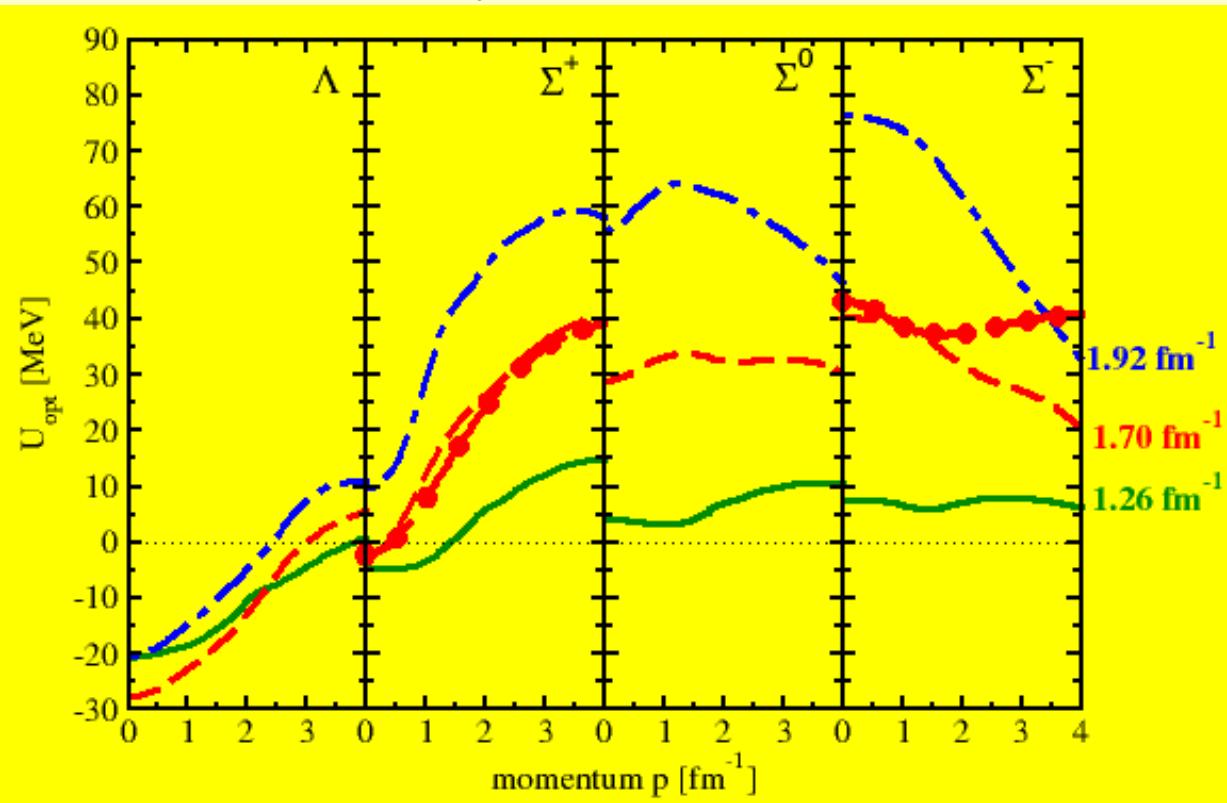
xEFT: EPJA52 (2016) 15

x-EFT: non-trivial MDI

Hyperon properties: Λ, Σ -optical potentials...

NLD fit:

density & momentum dependence
pure neutron matter

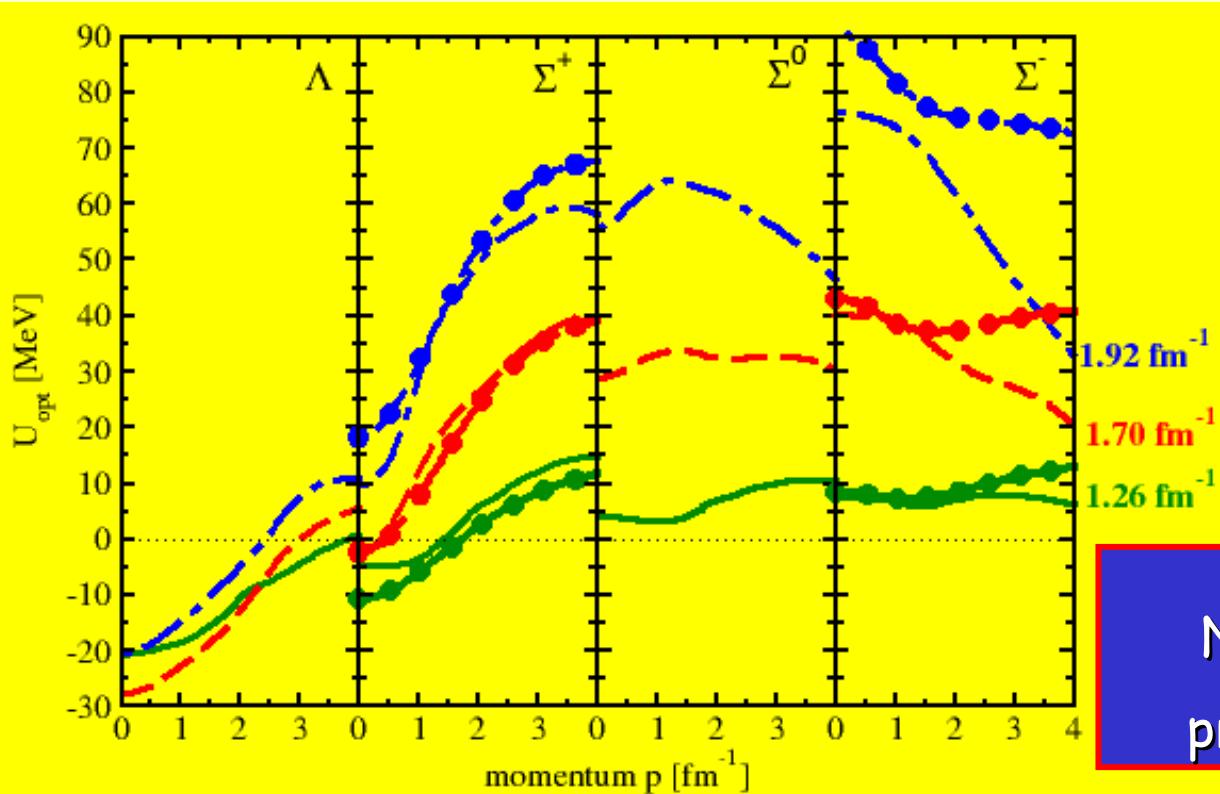


Curves: x-EFT
symbols: NLD

x-EFT: non-trivial MDI

Hyperon properties: Λ, Σ -optical potentials...

NLD predictions: density & momentum dependence
pure neutron matter



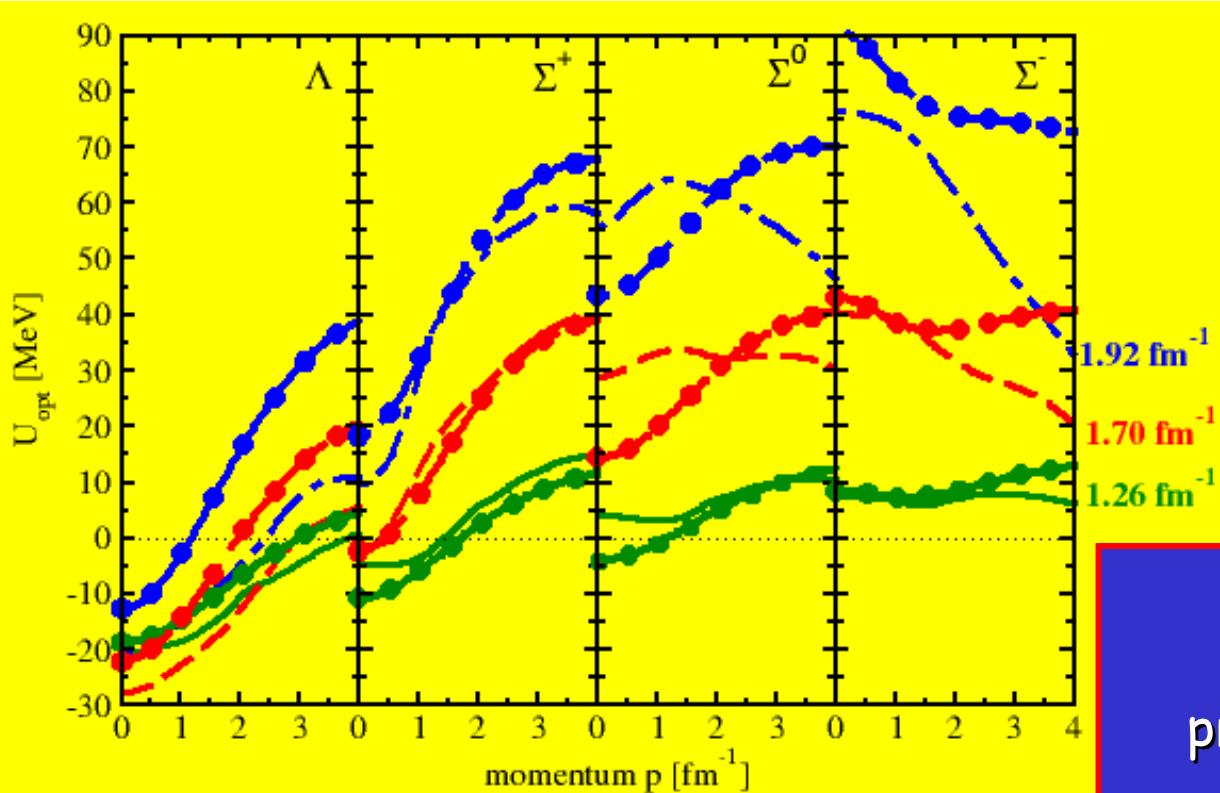
Curves: x-EFT
symbols: NLD

x-EFT: non-trivial MDI

NLD: charged hyperons
predicts non-trivial MDI

Hyperon properties: Λ, Σ -optical potentials...

NLD predictions: density & momentum dependence
pure neutron matter



Curves: x-EFT
symbols: NLD

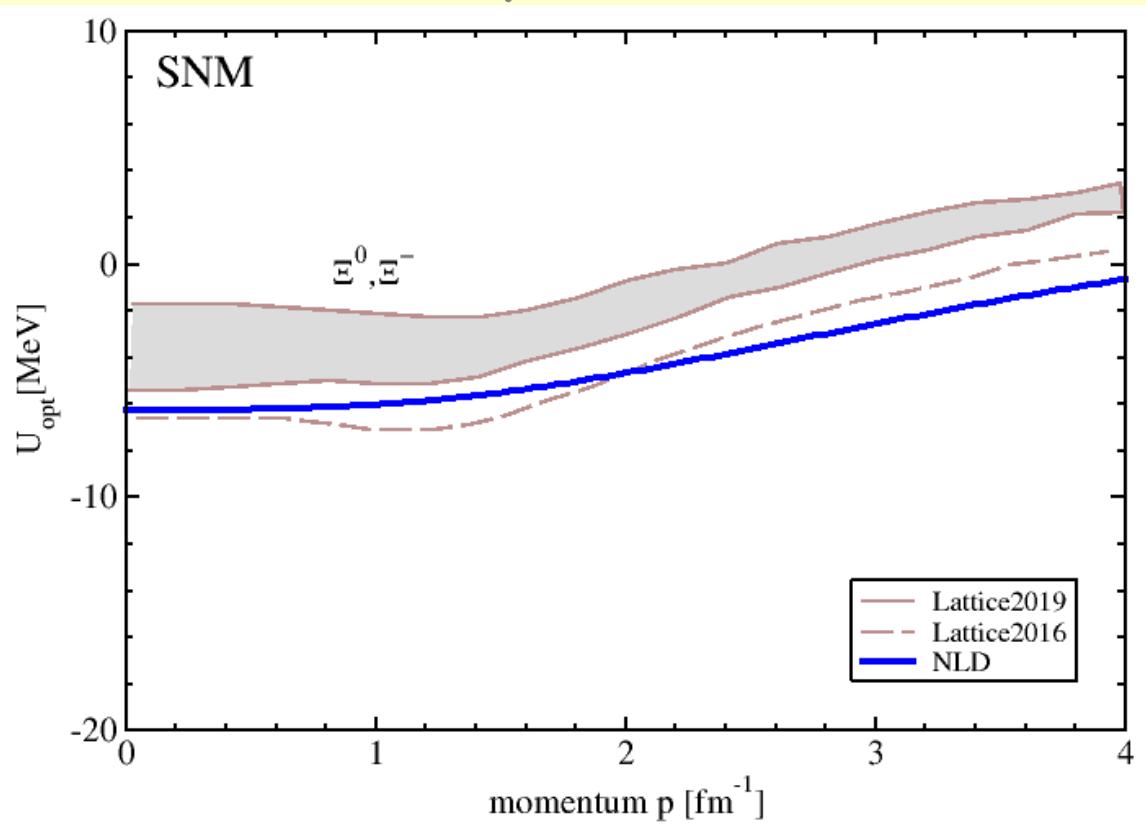
x-EFT: non-trivial MDI

NLD: all hyperons
predicts non-trivial MDI
adequate description

Hyperon properties: Ξ -optical potentials...

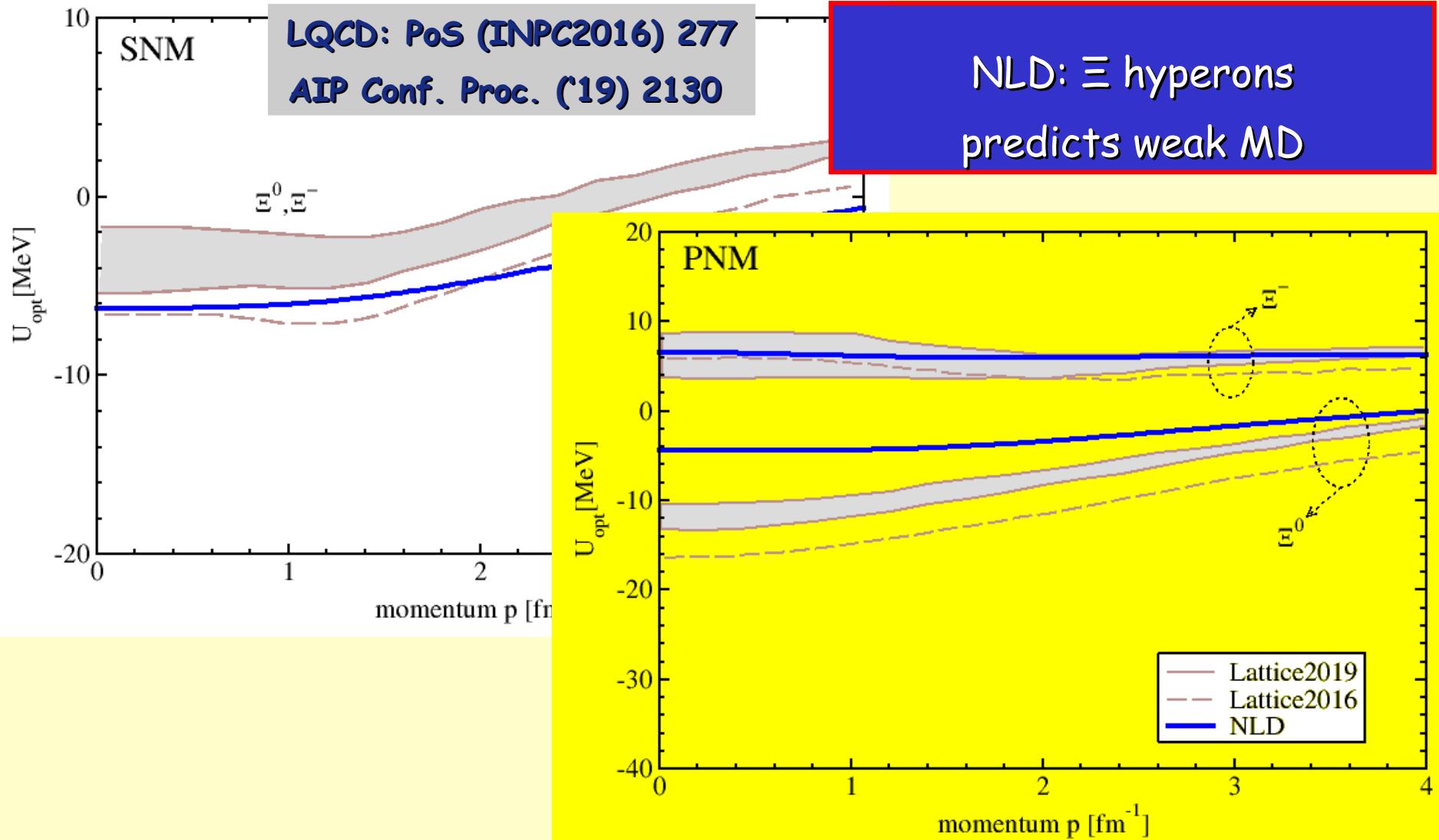
NLD fit:

density & momentum dependence
symmetric NM



Hyperon properties: Ξ -optical potentials...

NLD predictions: density & momentum dependence
SNM & pure neutron matter



Anti-Hyperon properties: optical potentials...

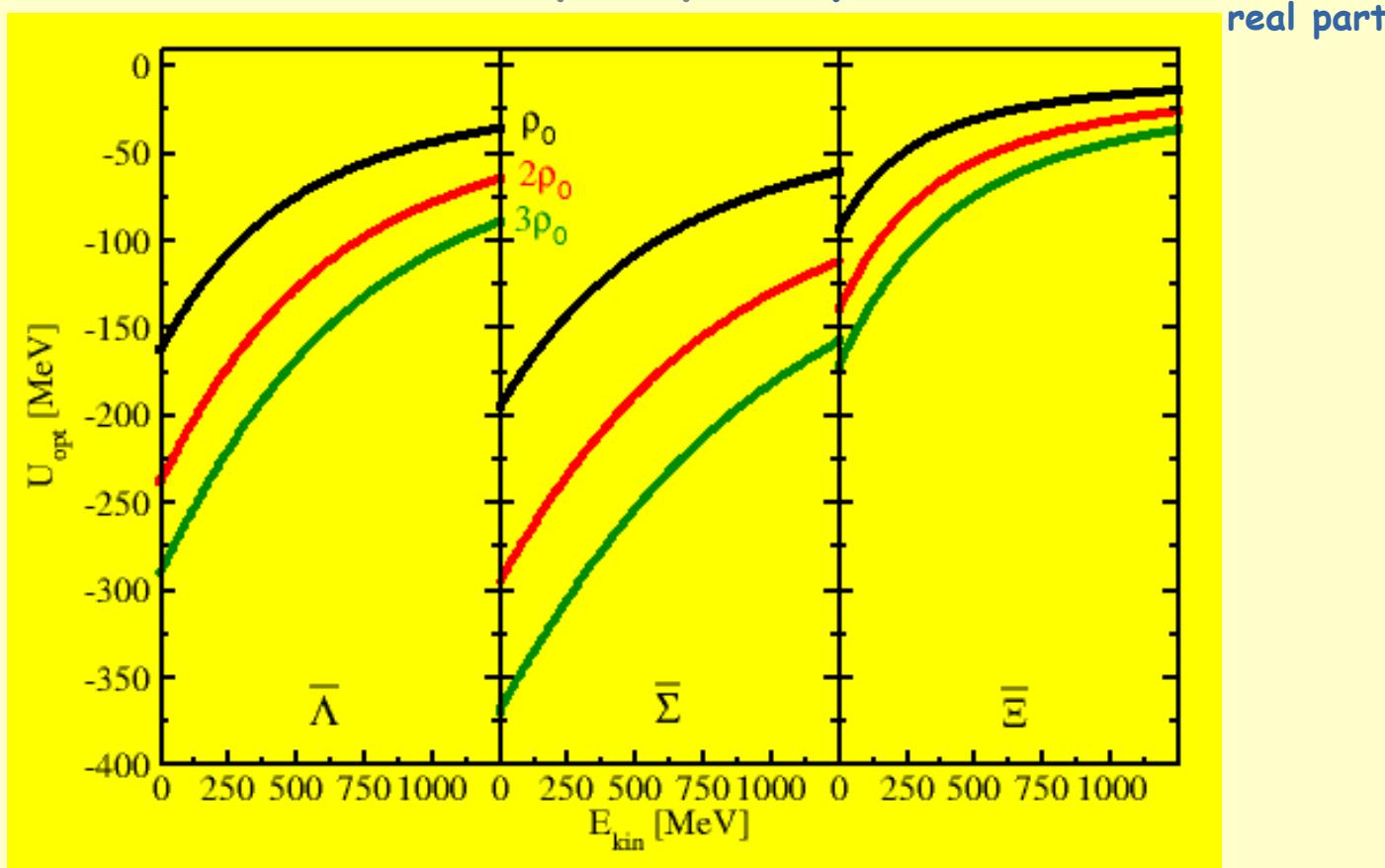
NLD predictions: density & momentum dependence
G-parity

Anti-Hyperon properties: optical potentials...

NLD predictions: density & momentum dependence
G-parity, no parameters

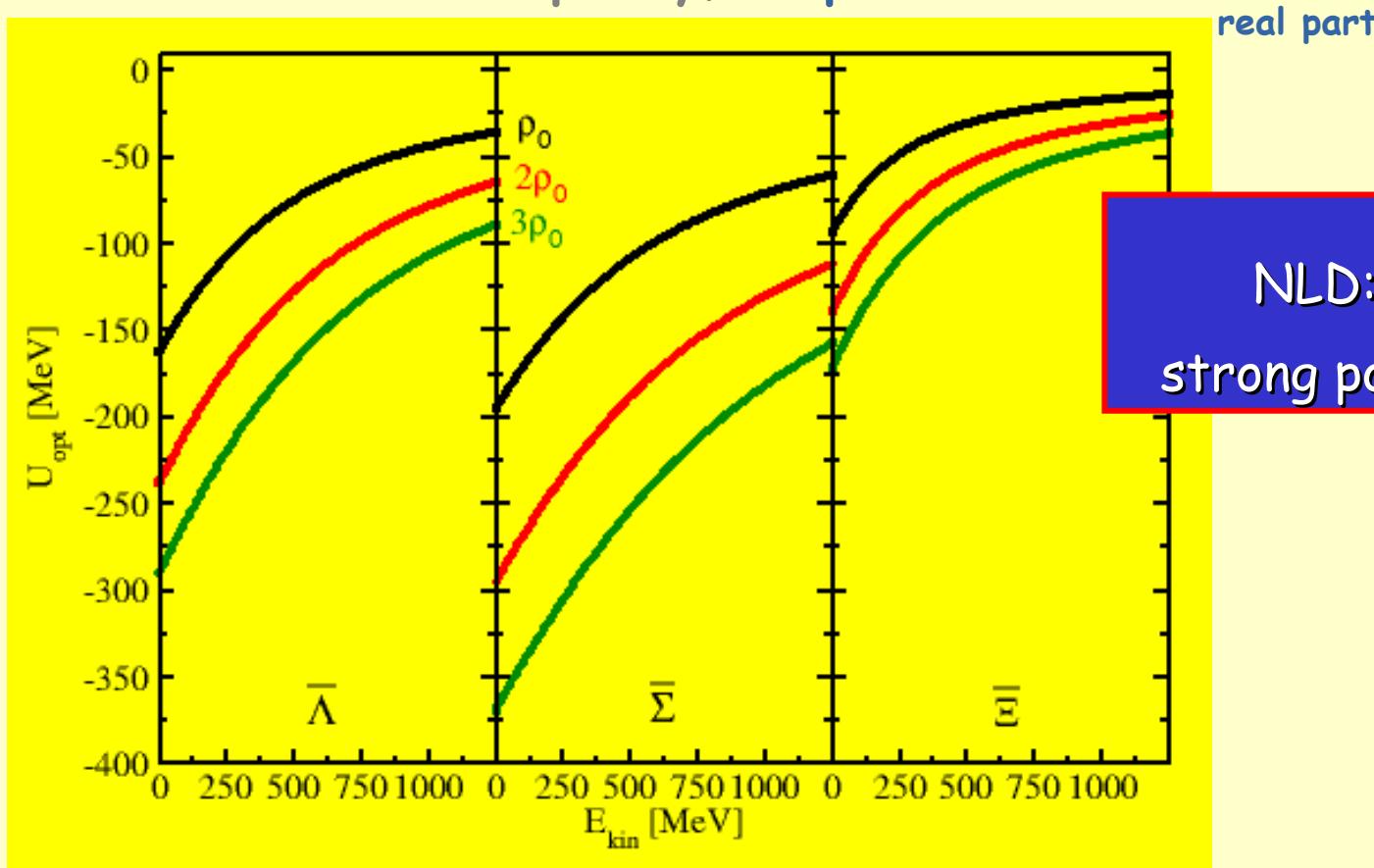
Anti-Hyperon properties: optical potentials...

NLD predictions: density & momentum dependence
G-parity, no parameters



Anti-Hyperon properties: optical potentials...

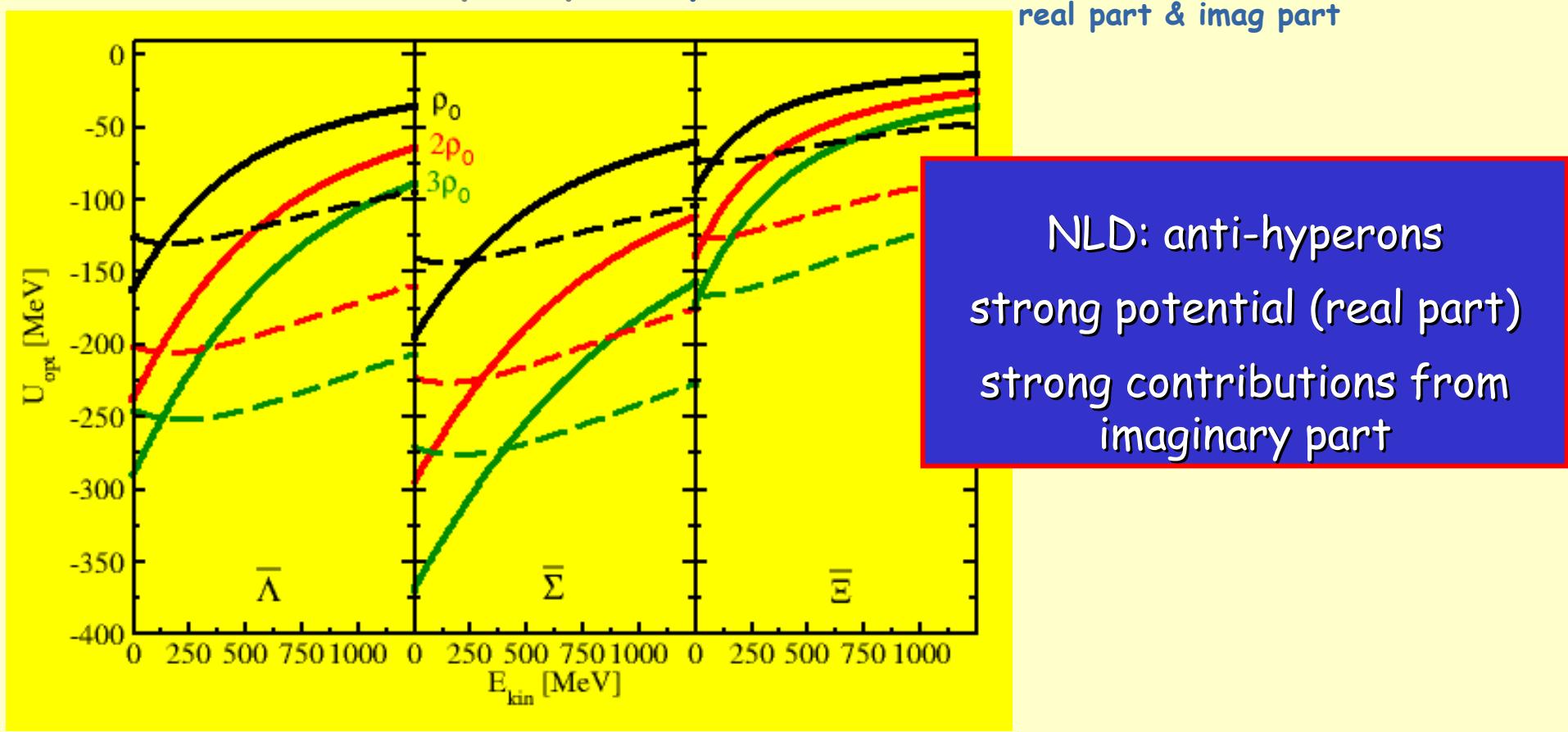
NLD predictions: density & momentum dependence
G-parity, no parameters



NLD: anti-hyperons
strong potential (real part)

Anti-Hyperon properties: optical potentials...

NLD predictions: density & momentum dependence
G-parity, no parameters



Final remarks & outlook...

→ NLD model

- keeping simplicity (RMF) to describe complexity (non-linear ρ & p dependences)
- realized by covariant introduction of regulators on a Lagrangian level
- in RMF: cut-off Λ regulates high ρ - & p -components of mean-fields
- cut-off Λ regulates also p -dependence of hyperon opt. pot.!

→ NLD Results

- EoS soft at low ρ ($K \sim 250$ MeV), but stiff at high ρ
compatible with all recent observations of EoS & NS
- Correct MD for in-medium proton (!) and (!) antiproton interactions
- compatible with recent results from x-EFT for hyperons in matter
- strong potentials for anti-Y & strong contributions from imag. parts

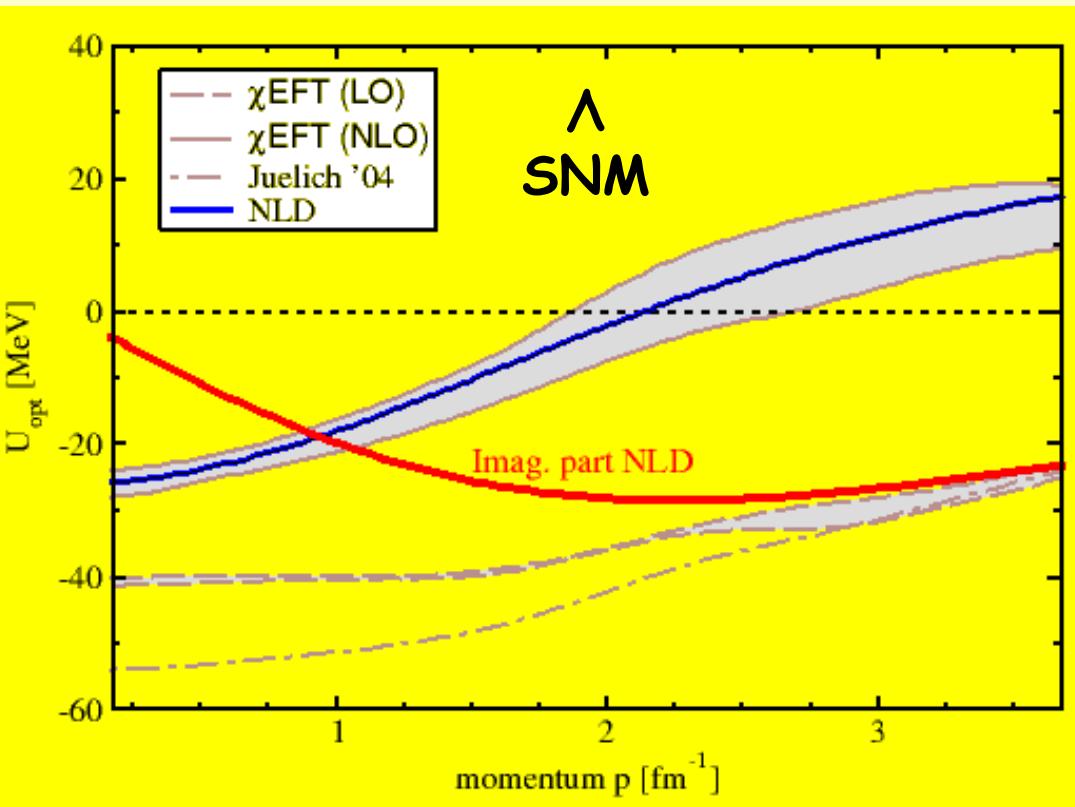
→ Under progress developments: include NLD mean-fields in...

- transport model for HADES ($\pi+A$ induced reactions)
- transport model for PANDA ($\bar{p}-A$ & $\Xi-A$ induced reactions)
- application to β -equilibrated matter for NS

Back up slides

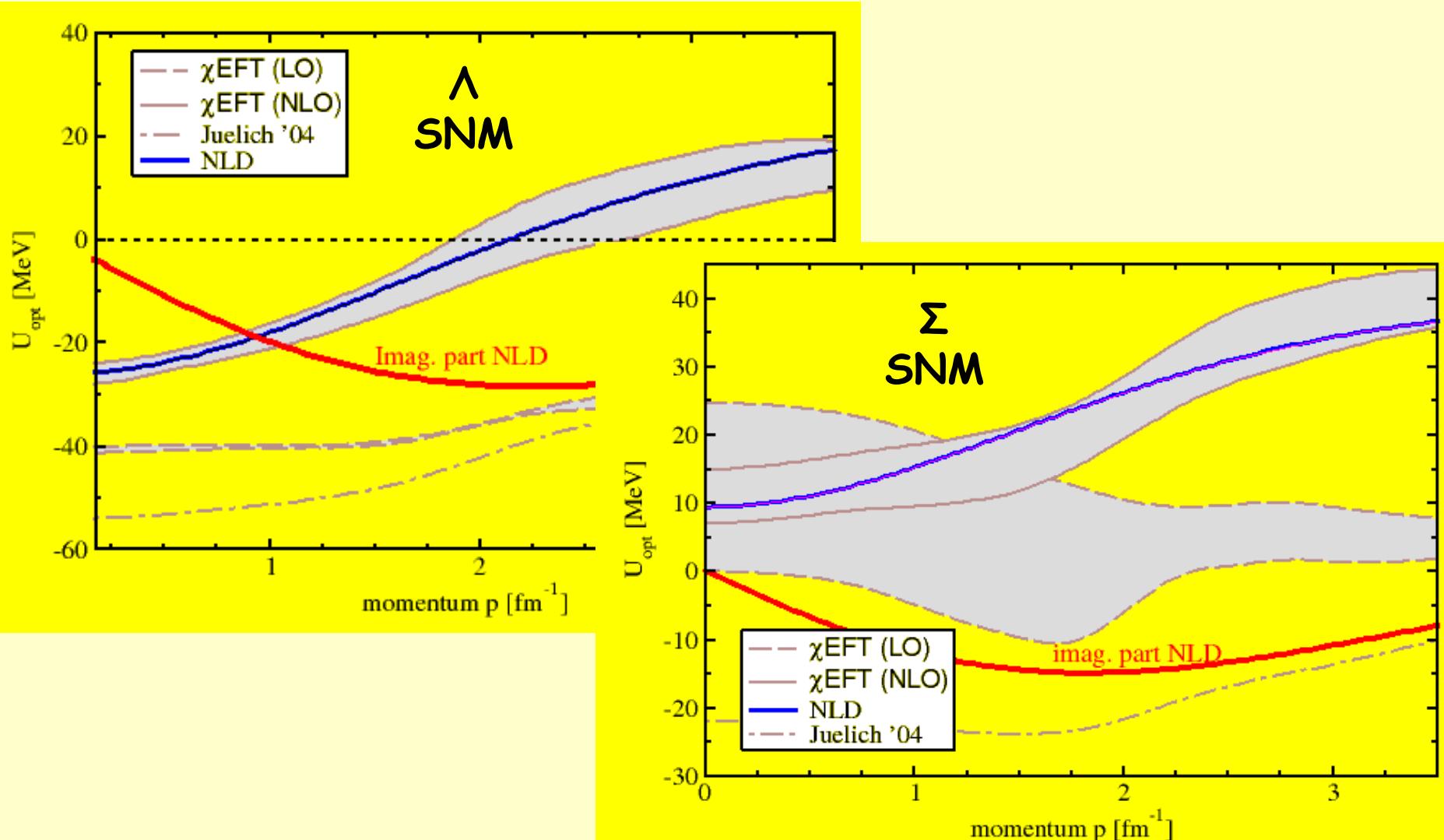
Anti-Hyperon properties: optical potentials...

NLD predictions: density & momentum dependence
G-parity, no parameters



Anti-Hyperon properties: optical potentials...

NLD predictions: density & momentum dependence
G-parity, no parameters



Explore in-medium Λ -pot: HADES experiment...

Key Information

- ▶ HADES Spectrometer at GSI
- ▶ Secondary π^- -beam 1.7 GeV/c

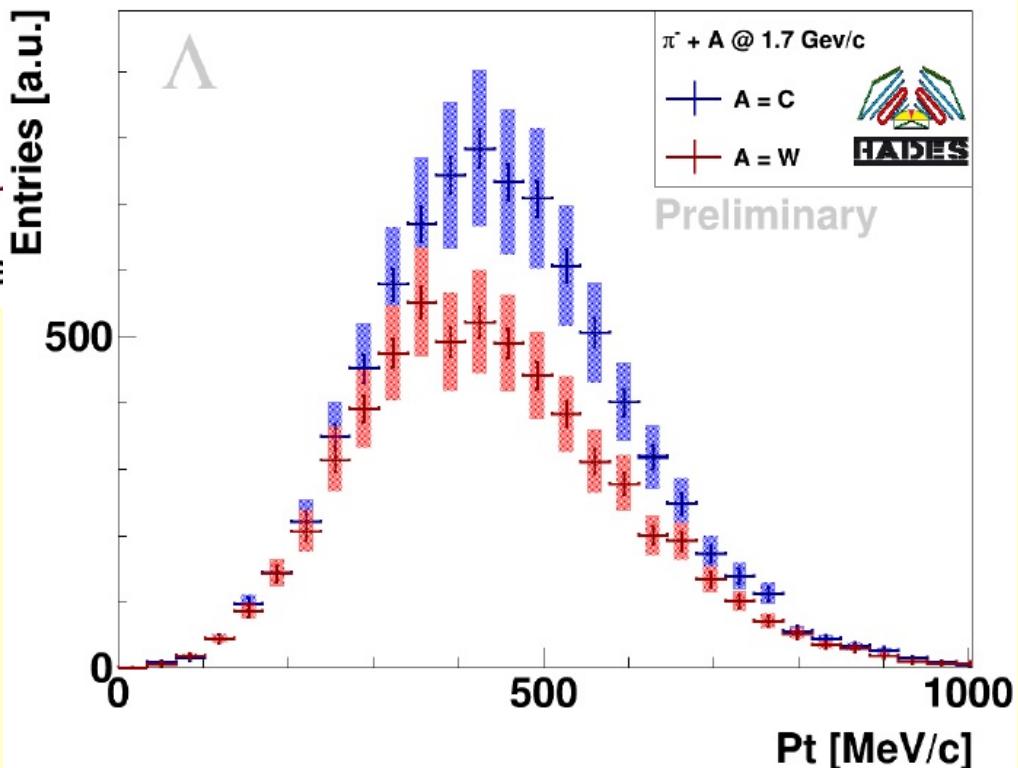
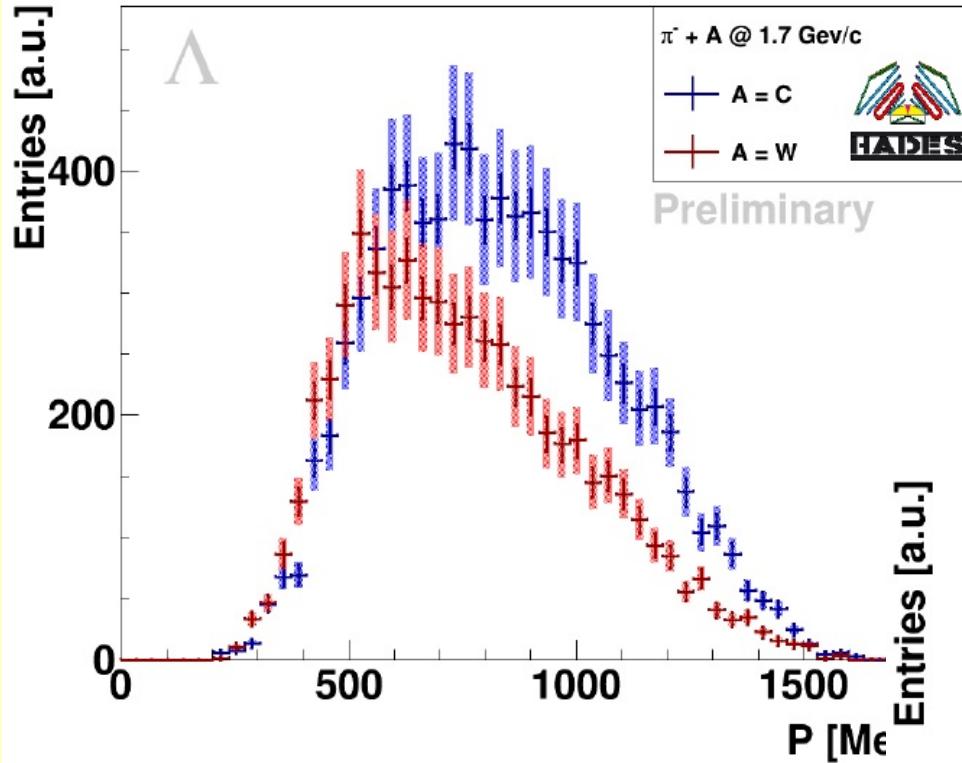
| ⌚ Target | W | C |
|--------------------------------|--------|--------|
| ➡ Segment Length [mm] | 2.4 | 7.2 |
| 📦 ρ [g/cm^3] | 19.3 | 1.85 |
| ● A | 183.84 | 12.011 |
| 📊 Statistics [$\times 10^8$] | 1.69 | 2.00 |

Idea:

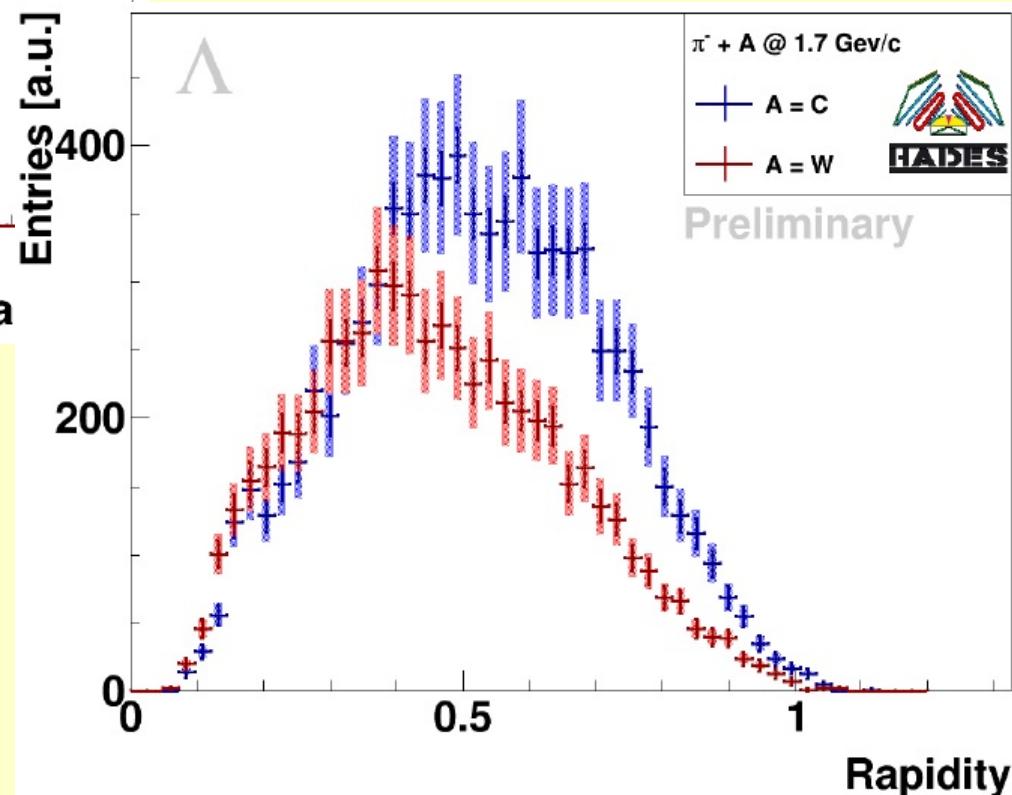
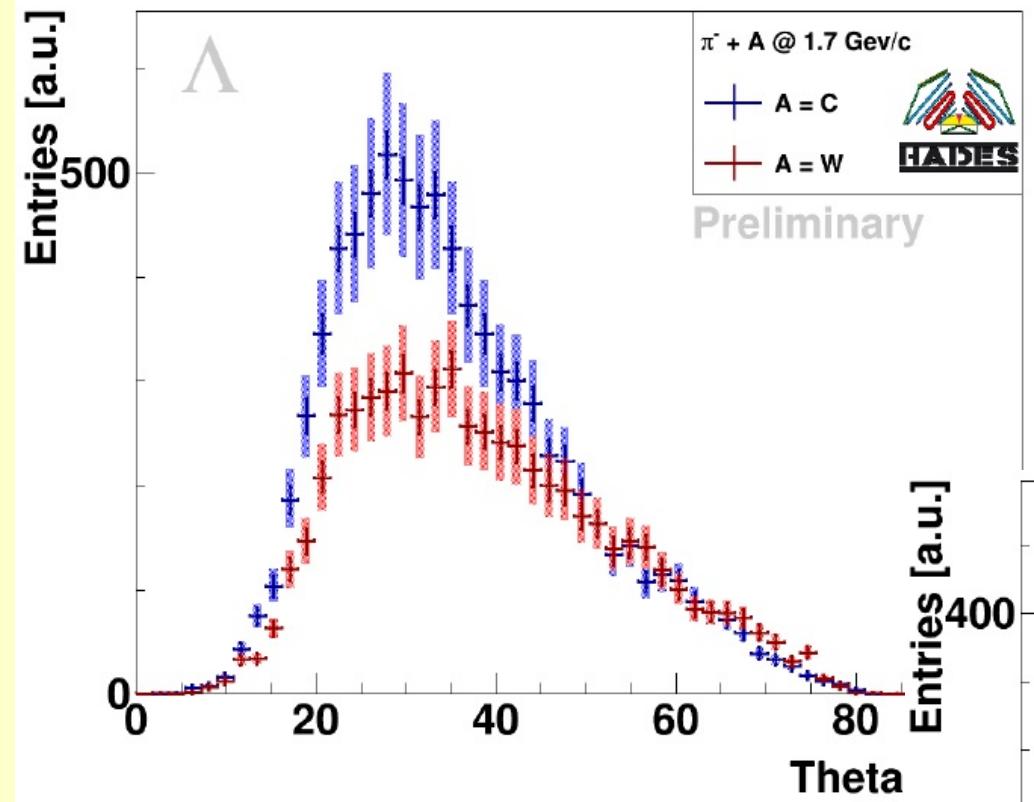
- ▶ Search for Charge Pattern 2+ 2- ($\Lambda \rightarrow p + \pi^-$, $K^0 \rightarrow \pi^+ + \pi^-$)
- ▶ Make best assignment of double π^- occurrence by minimizing:
 $\Delta M_\Lambda = M_{INV}(p + \pi^-) - M(\Lambda)_{PDG}$
 $\Delta M_{K0} = M_{INV}(\pi^+ + \pi^-) - M(K0)_{PDG}$
For all π^- Combination
- ▶ Cut on 2D ΔM_Λ vs. ΔM_{K0}

Icons from: <https://www.flaticon.com/>

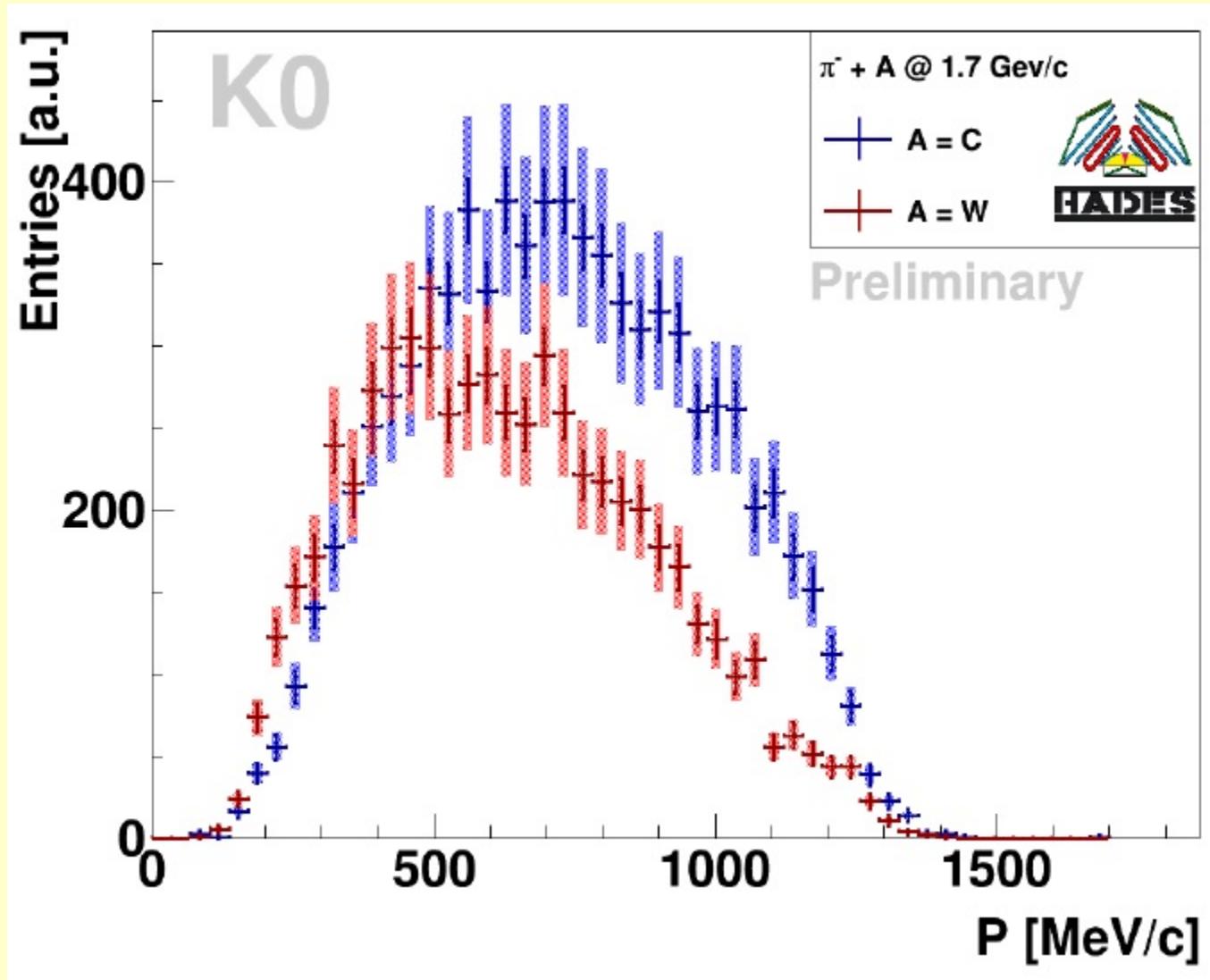
Explore in-medium Λ -pot: HADES new data...



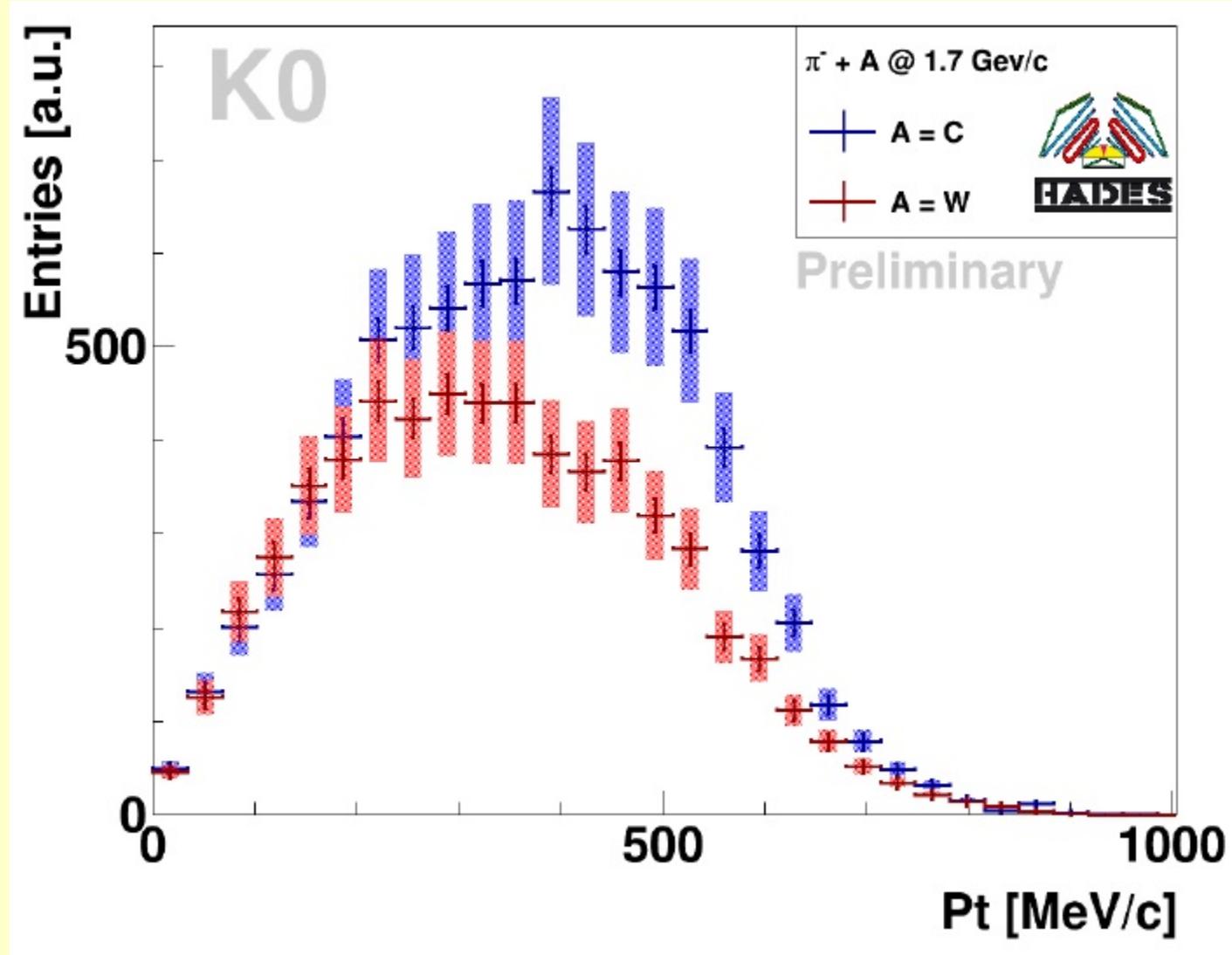
Explore in-medium Λ -pot: HADES new data...



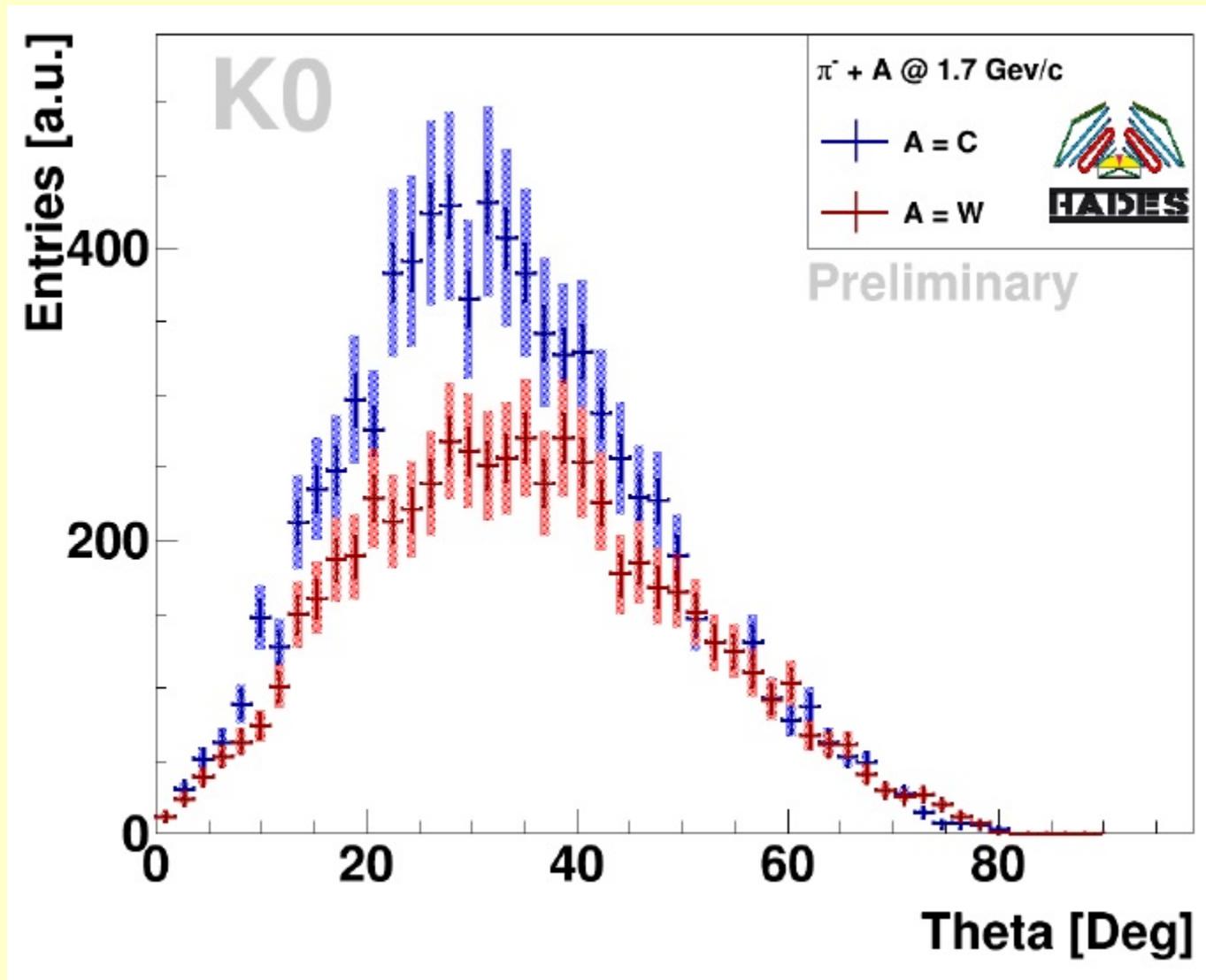
Explore in-medium K-pot: HADES new data...



Explore in-medium K-pot: HADES new data...



Explore in-medium K-pot: HADES new data...



Explore in-medium K-pot: HADES new data...

