# Measurement of Hyperon-nucleon scatterings at J-PARC 

Koji Miwa (Tohoku univ.) for the E40 collaboration

## CONTENTS

## Introduction

" Brief review of hyperon-proton scattering experiment

- Importance of hyperon-nucleon (YN) interaction study
$\Sigma$ p scattering experiment at J-PARC
- Analysis status of scattering channels
- Prospects of spin observables

Future prospect on $\Lambda$ p scattering

## NN AND YN SCATTERING DATA



Rich data of pp, np scattering.
$\rightarrow$ Fundamental information to construct realistic model of nuclear force.
Y. Fujiwara, Y. Suzuki, C. Nakamoto, Prog. Part. Nucl. Phys. 58 (2007) 439


For YN case, data quality and quantity are insufficiént.

## DIFFICULTY OF YP SCATTERING (HYPERON BEAM)

Proton beam ( $10^{12 \sim 13}$ particle / pulse)

Secondary beam ( $\pi, \mathrm{K}$ ) (107 particle / pulse)
G. Alexander, et al. Phys. Rev. 173 (1968) 1452


In past experiment, imaging method was used to identify Yp scattering

## STRATEGY OF HYPERNUCLEAR PHYSICS

Nuclear physics
Known nuclear force


Unknown nuclear structure

Hypernuclear physics
Unknown YN interaction

Hyperon proton scattering
Lattice QCD


PRC 64044302 (2001)


Theoretical framework extended to $\mathrm{SU}_{\mathrm{F}}(3)$ symmetry

Expect from hypernuclear structure


EPJA33 (2007) 243 PTEP (2015) 081 D01
$\gamma$-ray spectroscopy


## NEUTRON STAR AND YN INTERACTION

Two-body YN scattering is essential to understand the internal structure of neutron star.

- Interaction at short range
- Basic information to derive 3 body force from hypernuclear structure

Hypernuclear physics based on Realistic YN interaction


[^0]
## YN INTERACTION IN S=-1 SECTOR

|  | ${ }^{1} \mathrm{E}$ or ${ }^{3} \mathrm{O}$ | ${ }^{3} \mathrm{E}$ or ${ }^{1} \mathrm{O}$ |
| :---: | :---: | :---: |
| $\mathrm{NN}(1=0)$ | ------ | (27) |
| NN (l=1) | (10*) | ----- |
| $\Lambda N(1=1 / 2)$ | $1 / \sqrt{10[(8 s)}+3(27)]$ | $1 / \sqrt{2}\left[-\left(8_{a}\right)+\left(10^{*}\right)\right]$ |
| $\Sigma \mathrm{N}(1=1 / 2)$ | $1 / \sqrt{10[3(8)}$ ) - (27)] | $1 / \sqrt{ } 2\left[\left(8{ }_{c}\right)+\left(10^{*}\right)\right]$ |
| $\Sigma \mathrm{N}(1=3 / 2)$ | (27) | (10) |

## Pauli effect in quark level

Large repulsive core is expected in quark based model (Lattice QCD)


$\Sigma^{+} p$ potential by Lattice QCD


## YN INTERACTION IN S=-1 SECTOR

|  | ${ }^{1} \mathrm{E}$ or ${ }^{3} \mathrm{O} \quad{ }^{3} \mathrm{E}$ or ${ }^{1} \mathrm{O}$ |
| :---: | :---: |
| N ( $\mathrm{l}=0)$ | Forbiddden by isospin symmetry (27) |
| $\mathrm{NN}(1=1)$ | (10*) $\leftrightarrows$ |
| $\Lambda \mathrm{N}(\mathrm{l}=1 / 2)$ | $1 / \sqrt{10[(8 s)}+3(27)] \quad 1 / \sqrt{ } 2\left[-\left(8_{a}\right)+\left(10^{*}\right)\right]$ |
| $\Sigma \mathrm{N}(1=1 / 2)$ | $1 / \sqrt{ } 10\left[3\left(8_{s}\right)-(27)\right] \quad 1 / \sqrt{2}\left[\left(8_{a}\right)+\left(10^{*}\right)\right]$ |
| $\Sigma \mathrm{N}(1=3 / 2)$ | (27) (10) |

New type of LS force appears in YN sector
Anti-symmetric LS ${ }^{(-)}$force : $\mathrm{V}_{\text {ALS }} L \cdot\left(s_{1}-s_{2}\right)$


Spin singlet $\stackrel{\text { LS }}{ }(-)$ Spin triplet

- Quark based model
- Large ALS originated from coupling between (8s) and (8a) by one gluon exchange
- Boson exchange model
- No large ALS contribution

Difference theoretical prediction

- Polarization, Analyzing power
- Cross section


## J-PARC E40:

## Measurement of $d \sigma / d \Omega$ of $\Sigma p$ scatterings

## Physics motivations

" Verification of repulsive force due to quark Pauli effect in the $\Sigma^{+} p$ channel

- Systematic study of the $\Sigma \mathrm{N}$ interaction by separating isospin channel


Large $\mathrm{d} \sigma / \mathrm{d} \Omega$ is predicted in $\Sigma^{+} \mathrm{p}$ channel
Measurement of $d \sigma / d \Omega$
Aim to detect 10,000 events

- $\Sigma^{+} p$ elastic scattering
- $\quad \Sigma^{-} p$ elastic scattering

H. Nemura et al.

Few-Body Syst (2013) 54:1223-1 226

- $\quad \Sigma^{-} p \rightarrow \Lambda n$ inelastic scattering


## Collaborators

Tohoku Univ. : T. Aramaki, N. Chiga, N. Fujioka, M. Fujita, R. Honda, M. Ikeda, Y. Ishikawa, H. Kanauchi, S. Kajikawa, T. Kitaoka, T. Koike, K. Matsuda, Y. Matsumoto, K. Miwa, S. Ozawa, T. Rogers, T. Sakao, T. Shiozaki, H. Tamura, J. Yoshida H. Umetsu, S. Wada

JAEA : S. Hasegawa, S. Hayakawa, K. Hosomi, Y. Ichikawa, K. Imai, H. Sako, S. Sato, K. Tanida , T.O. Yamamoto, KEK : Y. Akazawa, M. leiri, S. Ishimoto, I. Nakamura, S. Suzuki, H. Takahashi, T. Takahashi, M. Tanaka, M. Ukai RIKEN : H. Ekawa

Chiba Univ. : H. Kawai, M. Tabata

Kyoto Univ. : S. Ashikaga, T. Gogami, T. Harada, M. Ichikawa,
T. Nanamura, M. Naruki, K. Suzuki

Osaka Univ. : K. Kobayashi, S. Hoshino, Y. Nakada, R. Nagatomi, M. Nakagawa, A. Sakaguchi

RCNP : H. Kanda, K. Shirotori, T.N. Takahashi
Okayama Univ. : K. Yoshimura
Korea Univ. : J.K. Ahn, S.H. Kim, W.S. Jung, S.W. Choi, B.M. Kang OMEGA Ecole Polytechnique-CNRS/IN2P3 : S. Callier, C.d.L. Taille,
L. Raux

Joint Institute for Nuclear Research : P. Evtoukhovitch,


## E40 EXPERIMENTAL SETUP



DATA ANALYSIS

## $\Sigma$ beam identification

beam in $\mathrm{LH}_{2}$ target is tagged by the magnetic
CATCH


## $\Sigma$ beam identification

$\Sigma$ beam in $\mathrm{LH}_{2}$ target is tagged by the magnetic spectrometers
CATCH




## Recoil proton identification



## Proton event in $\Sigma^{-}$production



## Identification of np scattering from $\Sigma^{-}$decay



## Identification of $\Sigma^{-} \mathrm{p} \rightarrow \Lambda \mathrm{n}$ conversion



## Proton event in $\Sigma^{-}$production

$\Sigma$ p scattering


Background (scattering with decay products)


Conversion process


## Proton event in $\Sigma^{-}$production



Conversion process


## Proton event in $\Sigma^{-}$production

$\Sigma \mathrm{p}$ scattering
Background (scattering with decay products)


Conversion process

np scattering is good tool to check our systematics to measure $\mathrm{d} \sigma / \mathrm{d} \Omega$

We derive $d \sigma / d \Omega$ for np scattering

## $\Sigma^{+} \mathrm{p}$ analysis

$\Sigma^{+} p$ scattering

$\Sigma^{+}$p scattering event $\sim 4500$
$\Delta \mathrm{E}(\mathrm{\Sigma} \mathrm{P}$ scattering, 2proton event)


Analysis by T. Nanamura (Kyoto Univ.)
pp scattering from $\Sigma^{+}$decay
 $\Delta \mathbf{P}\left(\Sigma^{+} \rightarrow \mathbf{p} \pi^{0}\right.$ decay, pp scattering)

## Summary of E40 scattering channel

pp scattering (CATCH detector calibration) with proton beam
" p momentum 0.5, 0.55, 0.6, 0.65, 0.75, $0.85 \mathrm{GeV} / \mathrm{c}$
np scattering from $\Sigma^{-} \rightarrow \mathrm{n} \pi^{-}$decay

- $0.3<p_{\mathrm{n}}(\mathrm{GeV} / \mathrm{c})<0.6$
$\Sigma^{-} p$ elastic scattering
- $0.45<\mathrm{p}_{\text {г- }}(\mathrm{GeV} / \mathrm{c})<0.8$
: ~4,500 events
$\Sigma^{-} p \rightarrow \Lambda n$ scattering
- $0.45<\mathrm{p}_{\text {E- }}(\mathrm{GeV} / \mathrm{c})<0.8$
: ~2,500 events
$\Sigma^{+} p$ elastic scattering
- $0.4<p_{\Sigma+}(\mathrm{GeV} / \mathrm{c})<0.8$
- $\sim 4,500$ events ( $\Sigma^{+} \rightarrow p \pi^{0}$ decay mode)
- $\sim 3,500$ events ( $\Sigma^{+} \rightarrow n \pi^{+}$decay mode)


## Efficiency for recoil proton in pp scattering

pp scattering
@ 0.45, 0.5, 0.55, 0.6, 0.65, 0.75, 0.85 GeV/c


Identify pp scattering event from missing mass of $\mathrm{pp} \rightarrow \mathrm{pX}$ reaction


The other proton's angle and momentum can be predicted from missing momentum

Check whether measured angle and energy are consistent or not.

## CATCH Efficiency study

1



Efficiency


Efficiency
CATCH efficiency was measured as a function of

- Vertex position
" Angle
- Proton energy


We made an efficiency table of CATCH based on the pp scattering data and Geant4 simulation.

## pp differential cross section with proton beams


pp scattering data were used for

- CATCH detector calibration, efficiency study
- Consistency check for $\mathrm{d} \sigma / \mathrm{d} \Omega$ measurement

There are some inconsistencies at edge region of detector acceptance. But the obtained $d \sigma / d \Omega$ 's are reasonable.


## np scattering from $\Sigma^{-}$decay

Good practice to derive $\Sigma p$ scattering $d \sigma / d \Omega$


Beam does not pass though the whole size of target

- $\Sigma^{-}$decay
- Various production point and direction


We have to estimate total track length in the $\mathrm{LH}_{2}$ target from these information

- production vertex
- $\Sigma^{-}$or neutron momentum

$$
\frac{d \sigma}{d \Omega}=\frac{\sum_{i_{v t z}} \frac{N_{\text {scat }}\left(i_{v t z}, \cos \theta\right)}{\text { eff }\left(i_{v t z}, \cos \theta\right)}}{\text { density }}
$$

## Neutron beam from $\Sigma^{-}$decay

## np scattering

$K^{+}$

Neutron momentum is obtained from $\mathrm{P}_{\Sigma-}$ and $\mathrm{P}_{\pi-}$


Track length in LH2 target was obtained from decay position and neutron direction



## np scattering identification from $\Sigma^{-}$decay



## $\mathrm{d} \sigma / \mathrm{d} \Omega$ of np scattering

Differential cross section of np scattering


## $\Sigma$ p channels

Data are w/o acceptance correction


Analysis is on going to derive $d \sigma / d \Omega$ with acceptance and efficiency tables
$\Sigma$-p elastic scattering $(0.55<\mathrm{p}(\mathrm{GeV} / \mathrm{c})<0.65)$

$\Sigma+\mathrm{p}$ elastic scattering $(0.5<\mathrm{p}(\mathrm{GeV} / \mathrm{c})<0.6)$


## Analysis of $\Sigma^{-} p \rightarrow \Lambda n$ channel



Based on the vertex position and momentum estimated from spectrometer analysis, $\Sigma^{-}$flight length was estimated from a Monte Carlo simulation.


## Analysis of $\Sigma^{-} p \rightarrow \Lambda n$ channel (BG estimation)

The $\Sigma^{-} p \rightarrow \Lambda n$ channel is separated from other background events.
We are trying to understand the background contribution as much as possible.
We fit the four kinematical spectra simultaneously with sum of each contribution.


## Analysis of $\Sigma^{-} p \rightarrow \Lambda n$ channel (BG estimation)


$\Delta p(\mathrm{GeV} / \mathrm{c})$

$\Delta p(\mathrm{GeV} / \mathrm{c})$

$\Delta p(\mathrm{GeV} / \mathrm{c})$

$\Delta p(\mathrm{GeV} / \mathrm{c})$

There are still discrepancy in background estimation. But uncertainty of the background contamination in the An peak is not large.

We can maybe release $d \sigma / d \Omega$ spectra soon for $\Sigma^{-} p \rightarrow \Lambda n$ channel.

## $\Sigma^{-} P$ ELASTIC SCATTERING

In order to identify S-p scattering, much tighter cuts are necessary.

## Scattering vertex



## Decay order


vtz(decay)-vtz(scat)


$$
0.45<p(\mathrm{GeV} / \mathrm{c})<0.55,-0.1<\cos \theta<0.1
$$

Red shaded : Simulation
Open hist : data selected for $-20<\Delta \mathrm{E}_{\Sigma_{p}}<10$ region to enhance true $\Sigma^{-} p$ scattering event

In order to estimate the efficiency, the simulation program was updated based on the measured vertex resolution (spectrometer), angular resolution (CFT) and energy resolution (BGO).

## $\Sigma^{-P}$ ELASTIC SCATTERING

## Subtraction of background reaction



In order to get better $S / N$ ratio for identifying $\Sigma^{-} p$ elastic scattering event, these background events were rejected.



## $\Sigma^{-P}$ ELASTIC SCATTERING

$$
\Sigma^{-} p(0.55<p(\mathrm{GeV} / \mathrm{c})<0.65)
$$



## POLARIZATION OF HYPERON

## $\Sigma^{+}$polarization at production

- Hyperon can be produced with some polarization for the production plane
- Spin direction of hyperon can be identified from angular dependence of decay product



This method can be applied to $Y p$ scattering to derive polarization in $Y_{p}^{4}$ scattering

## Spin observables in E40

Progress of Theoretical Physics, Vol. 100, No. 5, November 1998
Scattering Observables of the $N N$ and $Y N$ Interactions in the $S U_{6}$ Quark Model

Tadashi Fujita, Yoshikazu Fujiwara, Choki Nakamoto*
and Yasuyuki Suzuki**

## Polarization in $\Sigma^{+} p \rightarrow \Sigma^{+} p$ <br> - U/D asymmetry of proton from $\Sigma^{+}$decay

## Polarization in $\Sigma^{-} p \rightarrow \Lambda n$

" U/D asymmetry of proton from $\Lambda$ decay


Theoretical calculation of spin observable in Yp scattering is not updated.
We want to collaborate (or ask) theorist for calculation of spin observable in Yp scattering.
$\Lambda$ p scattering with the $\pi^{-} p \rightarrow K^{0} \Lambda$ reaction

## Possibility of $\Lambda$ p scattering

$\Lambda$ identification by $\pi^{-} p \rightarrow K^{0} \Lambda$ reaction is one of milestones

- Study of neutron-rich $\Lambda$ hypernuclei
- Study of $\Lambda p$ scattering

New K ${ }^{0}$ identify method $\pi^{+}$: magnetic spectrometer $\pi^{-}$: CATCH

$\square$ We can keep large acceptance for $\mathrm{K}^{0}$


## Possibility of $\Lambda$ p scattering

$\Lambda$ identification by $\pi^{-} p \rightarrow K^{0} \Lambda$ reaction is one of milestones

- Study of neutron-rich $\Lambda$ hypernuclei
- Study of $\Lambda p$ scattering

New K ${ }^{0}$ identify method $\pi^{+}$: magnetic spectrometer $\pi^{-}$: CATCH


We can keep large acceptance for $\mathrm{K}^{0}$
Missing mass ( $\pi^{-} p \rightarrow K^{0} X$ reaction)


Accumulated $\Lambda$ events $\sim 200 \mathrm{k} \rightarrow$ a several $10 \Lambda$ p scattering might be collected

# FUTURE PROJECT ：$\Lambda$ P SCATTERING EXPERIMENT 

Planning to perform at（new）K 1.1 beamline $\Lambda$ production via $\pi^{-} p \rightarrow \mathrm{~K}^{0} \Lambda$ at $\mathrm{p}_{\pi-}=1.05 \mathrm{GeV} / \mathrm{c}$ －High polarization of $\Lambda$ beam


R．D．Baker et al．，Nucl．Phys．B141（1978） 29



ヘのスピンが半成平面に対して十分偏極している

## Conventional representation of elastic scattering

Scattering amplitude in $\frac{\vec{i}}{2}+\frac{\overrightarrow{1}}{2} \rightarrow \frac{\vec{i}}{2}+\frac{\overrightarrow{1}}{2}$ scattering : $\rightarrow 4 \times 4$ matrix
From S. Ishikawa et al.
PRC 69, 034001 (2004)
$\rightarrow 6$ components from the restriction of parity conservation and time-reversal invariance
spin-independent spin-spin symmetric LS $(\Delta S=0) \quad$ anti-symmetric $L S(\Delta S=1)$

T matrix

$$
\boldsymbol{M}=V_{\mathrm{c}}+V_{\sigma}\left(\boldsymbol{s}_{a} \cdot \boldsymbol{s}_{b}\right)+V_{\mathrm{SLS}}\left(\boldsymbol{s}_{a}+\boldsymbol{s}_{b}\right) \cdot \boldsymbol{L}+V_{\mathrm{ALS}}\left(\boldsymbol{s}_{a}-\boldsymbol{s}_{b}\right) \cdot \boldsymbol{L}+V_{\mathrm{T}}\left(\left[\boldsymbol{s}_{a} \otimes \boldsymbol{s}_{b}\right]^{(2)} \cdot \boldsymbol{Y}_{2}(\hat{\boldsymbol{r}})\right),
$$

Tensor

## Scalar amplitude



## Vector amplitude



## Tensor amplitude

The tensor amplitudes $T_{j}(j=1,2,3)$ are calculated as

$$
\begin{gathered}
T_{j}=\frac{1}{2}\left\langle\boldsymbol{k}_{\mathrm{f}}\right| V_{\mathrm{T}} Y_{2, j-1}\left|\boldsymbol{k}_{\mathrm{i}}\right\rangle, \\
T_{\alpha}=\frac{1}{\sqrt{6}} T_{1}+T_{3},
\end{gathered}
$$

We want to derive these scattering amplitudes separately.

$$
T_{\beta}=\frac{1}{\sqrt{6}} T_{1}-T_{3,47}
$$

## Scattering observables

From S. Ishikawa et al.
PRC 69, 034001 (2004)

Differential cross section $\quad \sigma(\theta)=\frac{1}{4} \operatorname{Tr}\left(\mathcal{M} \mathcal{M}^{\dagger}\right)=\left|U_{\alpha}\right|^{2}+\frac{3}{16}\left|U_{\beta}\right|^{2}+\frac{1}{2}\left(\left|S_{S L S}\right|^{2}+\left|S_{\mathrm{ALS}}\right|^{2}\right)+\frac{1}{4}\left|T_{1}\right|^{2}+\frac{1}{2}\left(\left|T_{2}\right|^{2}+\left|T_{3}\right|^{2}\right)$

Analyzing power (Polarization)

Depolarization

$$
A_{y}(a)=-\frac{4 \sqrt{2}}{N_{R}} \operatorname{Im}\left\{U_{\alpha}^{*}\left(S_{\alpha}+S_{\beta}\right)+\frac{1}{4} U_{\beta}^{*}\left(-S_{\alpha}+S_{\beta}\right) \quad-\frac{1}{2} T_{\alpha}^{*}\left(-S_{\alpha}+S_{\beta}\right)\right\}
$$

$$
D_{y}^{y}(a)=\frac{4}{N_{R}} \operatorname{Re}\left\{\frac{1}{2 \sqrt{3}}\left(U_{0}+\frac{1}{\sqrt{3}} U_{1}\right)^{*} U_{1}+\frac{1}{2}\left(U_{0}-\frac{1}{\sqrt{3}} U_{1}\right)^{*} \times\left(\frac{1}{\sqrt{6}} T_{1}+T_{3}\right)-S_{1}^{*} S_{2}+\frac{1}{2}\left|S_{3}\right|^{2}-\frac{1}{\sqrt{6}} T_{1}^{*}\left(\frac{1}{\sqrt{6}} T_{1}-T_{3}\right)-\frac{1}{2}\left|T_{2}\right|^{2}\right\}
$$

## Analyzing power (Polarization)

Scattering with polarized $\Lambda$
Measure polarization of scattered $\Lambda$ with unpolarized $\Lambda$


A lot of measurements of scattering observables enable us to investigate each matrix components

## SUMMARY OF EXPERIMENTAL CONDITION

Beam time : $\sim 1$ month ( $w / 20 \mathrm{M} /$ spill $\pi^{-}$beam intensity)

## $\Lambda$ beam

- Momentum range $0.4 \sim 0.8 \mathrm{GeV} / \mathrm{c}$
" $\Lambda$ beam yield : $\sim 17 \mathrm{M}(\sim 0.55 \%$ of produced $\Lambda$ can be tagged by SKS)
" ~100\% polarization


## $\Lambda$ p scattering

- Total cross section of 30 mb was assumed
- 10,000 $\Lambda$ p scattering events can be identified



Tagged momentum distribution


Analyzing power ( $\Lambda$ )


## UP/DOWN ASYMMETRY FOR DEPOLARIZATION MEASUREMENT

Polarization case $\left(P_{\text {beam }}=1, P=0, D_{y}^{y}=1\right)$


## SUMMARY

## Hyperon-proton scattering experiment has become possible.

- Study of two-body YN interaction from scattering experiment is very important to understand twobody interaction without uncertainty from many-body system.


## $\Sigma \mathrm{p}$ scattering experiment at J-PARC

- $\Sigma \mathrm{N}$ interaction is a key to understand the origin of repulsive core in nuclear force.
- $\Sigma^{-} p, \Sigma^{-} p \rightarrow \Lambda n, \Sigma^{+} p$ channels are clearly identified with much better statistics.
" Derived $d \sigma / d \Omega$ of $p p$ and np scattering are consistent with partial wave analysis.
- Analysis is on going to derive differential cross section.
- Polarization in $\Sigma^{+} p$ scattering might be obtained.

Prospect on $\Lambda p$ scattering

- $\Lambda$ production by $\pi^{-} p \rightarrow K^{0} \Lambda$ reaction was established.
- By using polarized $\Lambda$ beam, we are planning a new experiment to measure $d \sigma / d \Omega$ and spin obserbables such as $P_{\Lambda}$ and $D_{y y}$.

We want to keep close discussion with theorist to investigate YN interaction.

BACKUP

## Lp scattering : Event category

1. All particles can be detected Kinematical check by both $\Delta E_{\text {recoil pr }}$ and $\Delta p_{\Lambda}$

2. two protons can be detected

Kinematical check by $\Delta E_{\text {recoil }}$ p,

$\pi^{-}$can be identified by MissingMass $\left(\Lambda_{\text {scat }} \rightarrow p X\right)$
$\Lambda p$ scattering event might be identified even in the bad $S / N$ ratio in missing mass spectrum


Start the $\Lambda$ p scattering analysis from case 1
3. 1 proton and $1 \pi$-can be obtained Kinematical check by $\Delta p_{\Lambda}$ or $\Delta E_{\text {recoil pr }}$



## $\Lambda p$ scattering identification





## Future project : $\Lambda$ p scattering experiment

Planning to perform at (new) K1.1 beamline $\Lambda$ production via $\pi^{-} p \rightarrow \mathrm{~K}^{0} \Lambda$ at $\mathrm{p}_{\pi-}=1.05 \mathrm{GeV} / \mathrm{c}$
R.D. Baker et al. , Nucl. Phys. B141 (1978) 29




High polarization of $\Lambda$ beam

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$$

We want to derive these scattering amplitudes separately.

$$
T_{\beta}=\frac{1}{\sqrt{6}} T_{1}-T_{3,56}
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## Scattering observables

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Differential cross section $\quad \sigma(\theta)=\frac{1}{4} \operatorname{Tr}\left(\mathcal{M} \mathcal{M}^{\dagger}\right)=\left|U_{\alpha}\right|^{2}+\frac{3}{16}\left|U_{\beta}\right|^{2}+\frac{1}{2}\left(\left|S_{S L S}\right|^{2}+\left|S_{\mathrm{ALS}}\right|^{2}\right)+\frac{1}{4}\left|T_{1}\right|^{2}+\frac{1}{2}\left(\left|T_{2}\right|^{2}+\left|T_{3}\right|^{2}\right)$

Analyzing power (Polarization)

Depolarization

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A_{y}(a)=-\frac{4 \sqrt{2}}{N_{R}} \operatorname{Im}\left\{U_{\alpha}^{*}\left(S_{\alpha}+S_{\beta}\right)+\frac{1}{4} U_{\beta}^{*}\left(-S_{\alpha}+S_{\beta}\right) \quad-\frac{1}{2} T_{\alpha}^{*}\left(-S_{\alpha}+S_{\beta}\right)\right\}
$$

$$
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Scattering with polarized $\Lambda$
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## Scattering observables

Differential cross section

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\begin{aligned}
& \sigma(\theta)=\frac{1}{4} \operatorname{Tr}\left(\mathcal{M} \mathcal{M}^{\dagger}\right)=\left|U_{\alpha}\right|^{2}+\frac{3}{16}\left|U_{\beta}\right|^{2}+\frac{1}{2}\left(\left|S_{\mathrm{SLS}}\right|^{2}+\left|S_{\mathrm{ALS}}\right|^{2}\right)+\frac{1}{4}\left|T_{1}\right|^{2}+\frac{1}{2}\left(\left|T_{2}\right|^{2}+\left|T_{3}\right|^{2}\right) \\
& A_{y}(a)=-\frac{4 \sqrt{2}}{N_{R}} \operatorname{Im}\left\{U_{\alpha}^{*}\left(S_{\alpha}+S_{\beta}\right)+\frac{1}{4} U_{\beta}^{*}\left(-S_{\alpha}+S_{\beta}\right) \quad-\frac{1}{2} T_{\alpha}^{*}\left(-S_{\alpha}+S_{\beta}\right)\right\}
\end{aligned}
$$

Analyzing power (Polarization)

Depolarization

$$
D_{y}^{y}(a)=\frac{4}{N_{R}} \operatorname{Re}\left\{\frac{1}{2 \sqrt{3}}\left(U_{0}+\frac{1}{\sqrt{3}} U_{1}\right)^{*} U_{1}+\frac{1}{2}\left(U_{0}-\frac{1}{\sqrt{3}} U_{1}\right)^{*} \times\left(\frac{1}{\sqrt{6}} T_{1}+T_{3}\right)-S_{1}^{*} S_{2}+\frac{1}{2}\left|S_{3}\right|^{2}-\frac{1}{\sqrt{6}} T_{1}^{*}\left(\frac{1}{\sqrt{6}} T_{1}-T_{3}\right)-\frac{1}{2}\left|T_{2}\right|^{2}\right\}
$$


T. Fujita et al., PTP 100 (1998) 931

J. Haidenbauer et al., Phys. Rev. C72, (2ర̊05), 044005

## Summary of experimental condition

Beam time : $\sim 1$ month ( $w / 20 \mathrm{M} /$ spill $\pi^{-}$beam intensity)
$\Lambda$ beam

- Momentum range $0.4 \sim 0.8 \mathrm{GeV} / \mathrm{c}$
- $\Lambda$ beam yield : $\sim 17 M(\sim 0.55 \%$ of produced $\Lambda$ can be tagged by SKS)
- $\sim 100 \%$ polarization
$\Lambda$ p scattering
- Total cross section of 30 mb was assumed
- 10,000 $\Lambda$ p scattering events can be identified



## Definition of quantization axis

$$
\begin{aligned}
& \mathbf{Y}=\mathbf{P}_{\pi-} \times \mathbf{P}_{\text {K0 }} \\
& \mathbf{X}=\mathbf{Y} \times \mathbf{P}_{\text {Abeam }} \\
& \mathbf{Z}=\mathbf{P}_{\text {Abeam }}
\end{aligned}
$$


$\Lambda$ beam polarization

$$
\frac{1}{N_{0}} \frac{d N}{d \cos \theta_{p}}=\frac{1}{2}\left(1+\alpha P_{\text {beam }} \cos \theta_{p}\right)
$$



## Definition of quantization axis



Beam $\Lambda$ polarization

$$
P_{\text {beam }}=P(\phi=0) \times \cos \phi
$$

$$
\phi=0,180^{\circ} \text { Maximum polarization }
$$

$$
\phi=90,270^{\circ} \text { Unpolarized beam }
$$

Scattered $\Lambda$ polarization

$$
\begin{gathered}
P_{\text {scat }}=\frac{2}{\alpha} \frac{N_{U}-N_{D}}{N_{U}+N_{D}} \\
\alpha=0.750 \pm 0.009 \pm 0.004
\end{gathered}
$$

(BESIII M. Ablikim et al. arXiv:1808.08917 )

$$
\begin{gathered}
N_{U}: 0<\theta_{p}(d e g)<90 \\
N_{D}: 90<\theta_{p}(\text { deg })<180
\end{gathered}
$$

## Up/Down asymmetry for depolarization

No beam polarization case

Spin polarization in the final state
$P_{\text {scat }}=\frac{2}{\alpha} \frac{N_{U}-N_{D}}{N_{U}+N_{D}}=\frac{P+D_{y}^{y} P_{\text {beam }}}{1+P P_{\text {beam }}}$
$\mathrm{P}_{\text {beam }}$ : Polarization of beam

$$
P_{\text {beam }}=P(\phi=0) \times \cos \phi
$$

$P$ : Induced polarization by the unpolarized beam
$D_{y}{ }^{y}$ : Depolarization


## Up/Down asymmetry for depolarization

## Polarized case

Spin polarization in the final state
$P_{s c a t}=\frac{2}{\alpha} \frac{N_{U}-N_{D}}{N_{U}+N_{D}}=\frac{P+D_{y}^{y} P_{\text {beam }}}{1+P P_{\text {beam }}}$
$P_{\text {beam }}$ : Polarization of beam

$$
P_{\text {beam }}=P(\phi=0) \times \cos \phi
$$

$P$ : Induced polarization by the unpolarized beam
$D_{y}{ }^{y}$ : Depolarization


## Up/Down asymmetry for depolarization measurement

Polarization case $\left(P_{\text {beam }}=1, P=0, D_{y}^{y}=1\right)$



## Up/Down asymmetry for depolarization measurement



## SUMMARY

## Hyperon-proton scattering experiment has become possible.

- Study of two-body YN interaction from scattering experiment is very important to understand twobody interaction without uncertainty from many-body system.


## $\Sigma \mathrm{p}$ scattering experiment at J-PARC

- $\Sigma \mathrm{N}$ interaction is a key to understand the origin of repulsive core in nuclear force.
- $\Sigma^{-} p, \Sigma^{-} p \rightarrow \Lambda n, \Sigma^{+} p$ channels are clearly identified with much better statistics.
- Derived $d \sigma / d \Omega$ of $p p$ and $n p$ scattering are consistent with partial wave analysis.
- Analysis is on going to derive differential cross section.
- Polarization in $\Sigma^{+} p$ scattering might be obtained.

Prospect on $\Lambda p$ scattering

- $\Lambda$ production by $\pi^{-} p \rightarrow K^{0} \Lambda$ reaction was established.
- By using polarized $\Lambda$ beam, we are planning a new experiment to measure $d \sigma / d \Omega$ and spin obserbables such as $P_{\Lambda}$ and $D_{y y}$.

We want to keep close discussion with theorist to investigate YN interaction.

## YN INTERACTION IN S=-1 SECTOR

|  | ${ }^{1} \mathrm{E}$ or ${ }^{3} \mathrm{O} \quad{ }^{3} \mathrm{E}$ or ${ }^{1} \mathrm{O}$ |
| :---: | :---: |
| N ( $\mathrm{l}=0)$ | Forbiddden by isospin symmetry (27) |
| $\mathrm{NN}(1=1)$ | (10*) $\leftrightarrows$ |
| $\Lambda \mathrm{N}(\mathrm{l}=1 / 2)$ | $1 / \sqrt{10[(8 s)}+3(27)] \quad 1 / \sqrt{ } 2\left[-\left(8_{a}\right)+\left(10^{*}\right)\right]$ |
| $\Sigma \mathrm{N}(1=1 / 2)$ | $1 / \sqrt{ } 10\left[3\left(8_{s}\right)-(27)\right] \quad 1 / \sqrt{2}\left[\left(8_{a}\right)+\left(10^{*}\right)\right]$ |
| $\Sigma \mathrm{N}(1=3 / 2)$ | (27) (10) |

New type of LS force appears in YN sector
Anti-symmetric LS ${ }^{(-)}$force : $\mathrm{V}_{\text {ALS }} L \cdot\left(s_{1}-s_{2}\right)$


Spin singlet $\stackrel{\text { LS }}{ }(-)$ Spin triplet

- Quark based model
- Large ALS originated from coupling between (8s) and (8a) by one gluon exchange
- Boson exchange model
- No large ALS contribution

Difference theoretical prediction

- Polarization, Analyzing power
- Cross section


## Baryon-Baryon Interaction

Investigation of $B B$ interaction

- Basic information to describe the system with Hyperon
- Hypernuclei
- High density matter inside neutron star
" Origin of short-range core
Short range interaction by
Quark Cluster model
- Pauli effect in quark level
- Color magnetic interaction

YN, YY interactions show rich aspect especially in short range region

Forbidden state


## $\Sigma$ hypernucleus and $\Sigma \mathrm{N}$ interaction


T. Nagae et al., Phys. Rev. Lett. 80 (1998) 1605

## $\Sigma$ hyperon in light and medium heavy nuclei

$$
\Sigma \text {-nucleus potential } \quad U_{0}^{\Sigma}+U_{\tau}^{\Sigma}\left(T_{R} t_{\Sigma}\right) / A_{\text {core }}
$$

spin-isospin averaged potential Lane's term : isospin dependent potential
Spin-isospin averaged potential can be studied from shape of quasi-free $\Sigma^{-}$production spectrum in medium heavy nuclei

P.K. Saho et al.;Phys.Rev.C70(2004)044613

T. Harada, R. Honda, Y. Hirabayashi,

Phys. Rev. C97 (2018) 024601


[^0]:    Ultra high-resolution $\Lambda$ hypernuclear spectroscopy

