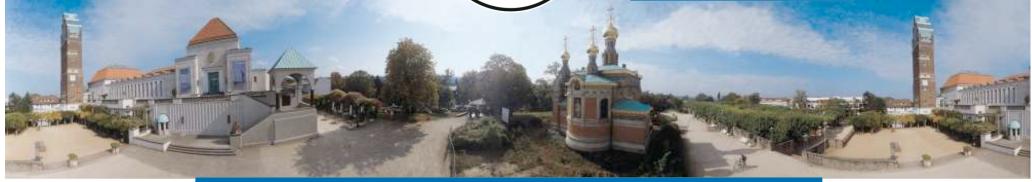
# The strong interaction at neutron-rich extremes

#### Achim Schwenk









**Helmholtz Alliance** 

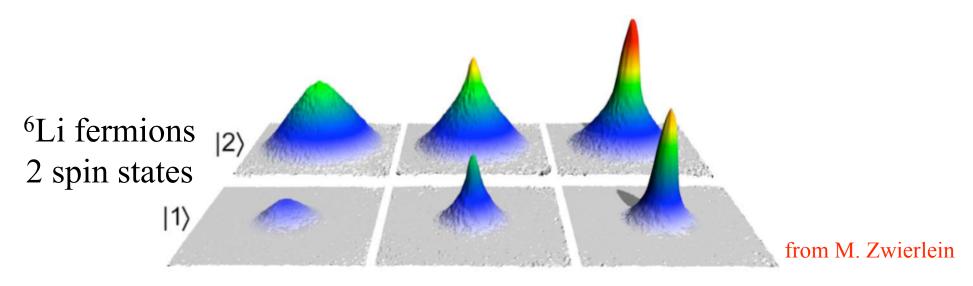
Extremes of Density and Temperature: Cosmic Matter in the Laboratory

## **ExtreMe Matter Institute EMMI**

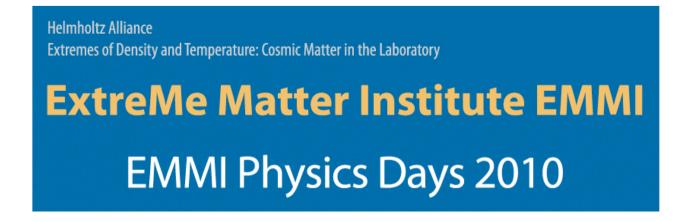
**EMMI Physics Days 2010** 

GSI, Darmstadt, Germany November 4 - 5, 2010

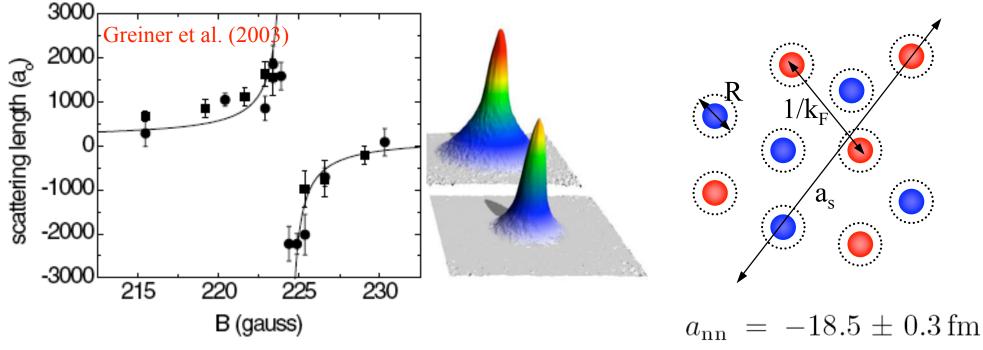
# The strong interaction at extreme low densities: Universal properties of neutrons and cold atoms



neutrons with same density and temperature have the same properties!



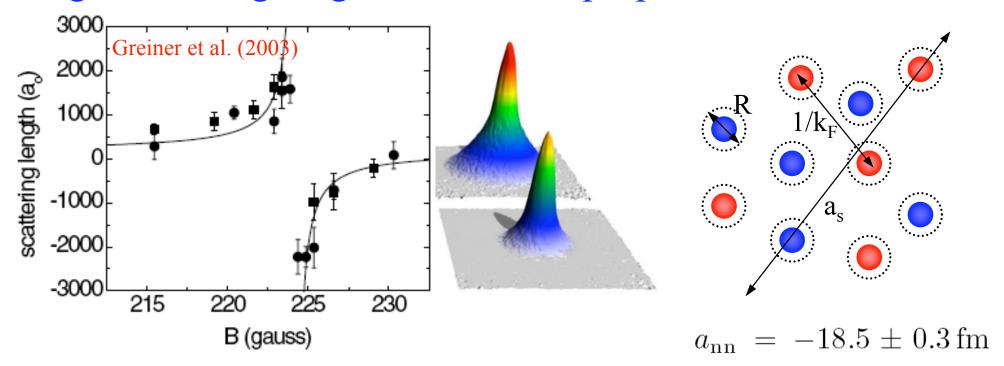
# Large scattering lengths: Universal properties at low densities



strong interactions via Feshbach resonances

 $a_{\rm nn} = -18.5 \pm 0.3 \, {\rm fm}$ large for neutrons

## Large scattering lengths: Universal properties at low densities



strong interactions via Feshbach resonances

large for neutrons

dilute Fermi system with large scattering length has universal properties

$$0 \leftarrow 1/a_s \ll k_{\rm F} \ll 1/r_e\,, 1/R\,, \ldots \rightarrow \infty$$
 strongly-interacting dilute

only Fermi momentum or density sets scale

physics is independent of interaction/system details: from dilute neutron matter to resonant <sup>6</sup>Li or <sup>40</sup>K atoms in traps

## Large scattering lengths: Universal thermodynamics

energy per particle 
$$\frac{E}{N} = \xi \left(\frac{E}{N}\right)_{\text{free}} = \xi \, \frac{3k_{\text{F}}^2}{10m}$$

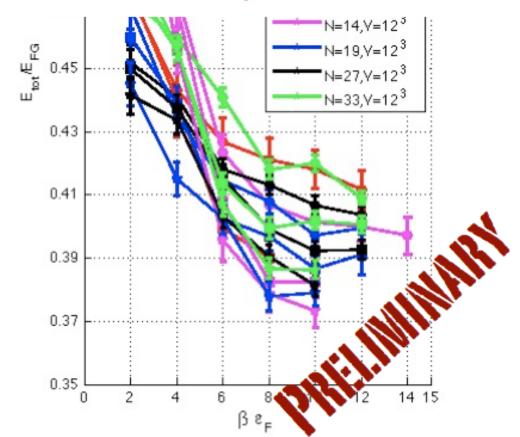
#### with universal Bertsch parameter ξ

Quantum Monte Carlo:  $\xi$ =0.40(1)

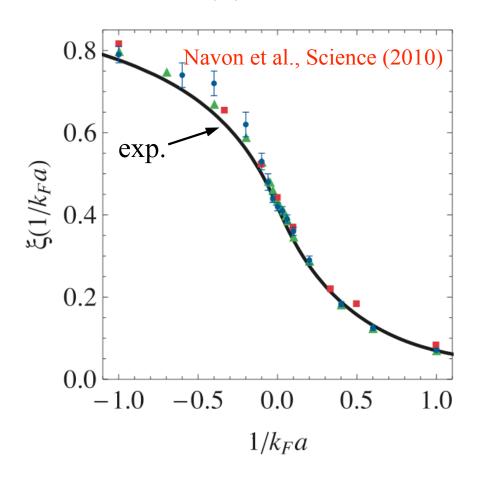
Gezerlis, Carlson (2009),...

lattice: Auxiliary-Field Hybrid MC

Drut, Gezerlis, Laehde @ 2010 UNEDF meeting



most precise results with  $^6$ Li  $\xi$ =0.39(2) and 0.41(2) cloud size and E(S) zuo, Thomas (2009)



# Large scattering lengths: Resonance superfluidity

phase transition to superfluid with universal critical temperature  $T_c \approx 0.15$ -0.2  $T_F$  Thomas et al., Science (2005),... Nascimbene et al., Nature (2010)

and universal pairing gap  $\Delta \approx 0.45~\epsilon_F$ 

Schirotzek et al. (2009), Carlson, Reddy (2008)

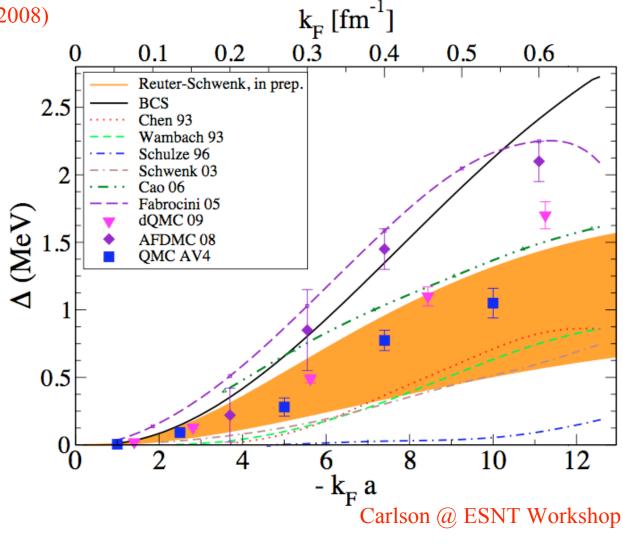
Superfluid pairing gap for low-density neutron matter

Progress from MC and other calculations

Differences to cold atoms:

effective range  $k_F r_e \sim 2$ 

(weak shell effects for resonant interactions)



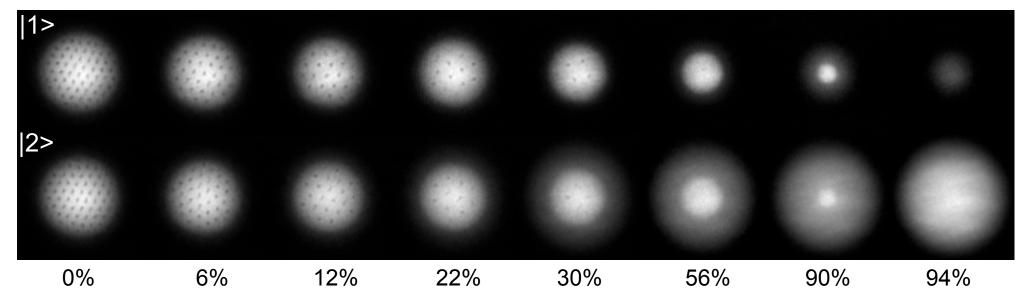
# Large scattering lengths: Asymmetric systems

polaron binding =  $\eta \epsilon_F$  with  $\eta \approx -0.6$ 

determines critical polarization of superfluidity (superfluid energy  $\xi$  vs. normal energy  $\eta$ )

Chevy (2006), Bulgac, Forbes (2007),...





Zwierlein et al., Science (2006)

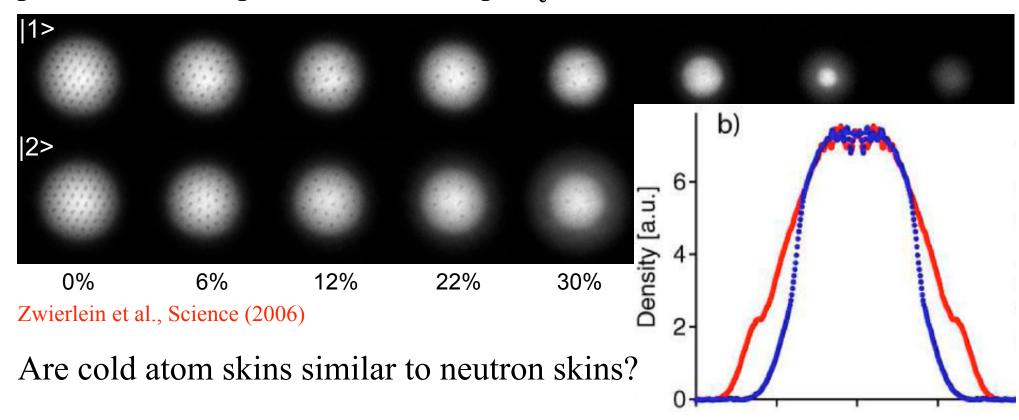
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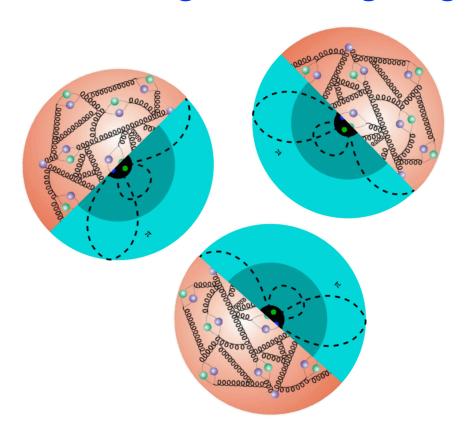
Chevy (2006), Bulgac, Forbes (2007),...

predicts critical polarization in traps  $P_c \sim 70\%$ 



Shin et al., (2006)

# The nuclear forces frontier — beyond universal large scattering length physics



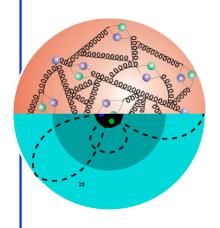
Helmholtz Alliance
Extremes of Density and Temperature: Cosmic Matter in the Laboratory

**ExtreMe Matter Institute EMMI** 

**EMMI Physics Days 2010** 

## $\Lambda$ / Resolution dependence of nuclear forces

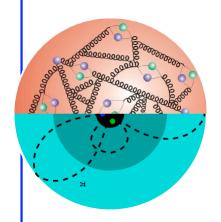
with high-energy probes: quarks+gluons



 $\Lambda_{pionless}$  momenta  $Q << m_{\pi}$ : pionless effective field theory large scattering length physics and corrections

## $\Lambda$ / Resolution dependence of nuclear forces

with high-energy probes: quarks+gluons



Effective theory for NN, 3N, many-N interactions and electroweak operators: resolution scale/ $\Lambda$ -dependent

$$H(\Lambda) = T + V_{NN}(\Lambda) + V_{3N}(\Lambda) + V_{4N}(\Lambda) + \dots$$

$$\begin{split} & \Lambda_{chiral} \\ & momenta \ Q \sim \lambda^{-1} \sim m_\pi \text{=} 140 \ \text{MeV: chiral effective field theory} \\ & neutrons \ \text{and protons interacting via pion exchanges} \\ & \text{and shorter-range contact interactions} \\ & \text{typical momenta in nuclei} \sim m_\pi \end{split}$$

$$\Lambda_{\text{pionless}}$$
  $Q \ll m_{\pi}$ 

Separation of scales: low momenta NN 3N 4N 
$$\frac{Q^0}{\Lambda^0}$$
  $\frac{Q^0}{\Lambda^0}$   $\frac{Q^0}{\Lambda^0}$ 

Separation of scales: low momenta  $\frac{1}{\lambda} = Q \ll \Lambda_b$  breakdown scale ~500 MeV

limited resolution at low energies, can expand in powers  $(Q/\Lambda_b)^n$ 

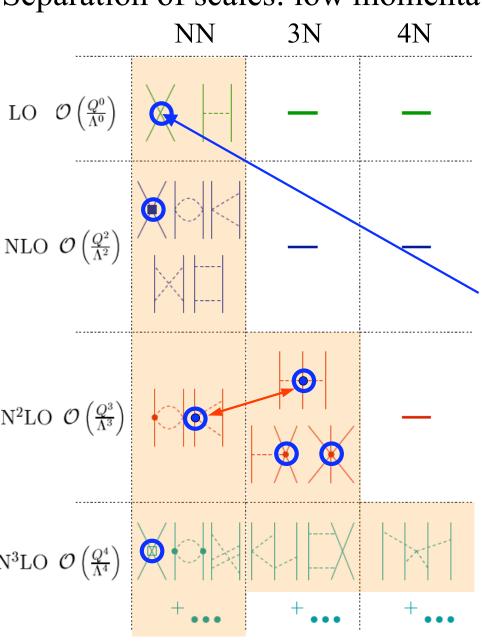
expansion parameter  $\sim 1/3$ 

Separation of scales: low momenta  $\frac{1}{\lambda} = Q \ll \Lambda_b$  breakdown scale ~500 MeV

 $\overline{\lambda} = Q \ll \Lambda_b$  breakdown scale ~500 MeV include long-range pion physics details at short distance not resolved capture in few short-range couplings, fit to experiment once,  $\Lambda$ -dependent

Separatio	n ot scat NN	es: 10W 1 3N	nomenta 4N
LO $\mathcal{O}\left(\frac{Q^0}{\Lambda^0}\right)$	<b>⋈</b>		
NLO $\mathcal{O}\left(rac{Q^2}{\Lambda^2} ight)$			
${ m N^2LO}~{\cal O}\left(rac{Q^3}{\Lambda^3} ight)$			
N <sup>3</sup> LO $\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$	<b>⋈</b>	+	+

Separation of scales: low momenta  $\frac{1}{\lambda} = Q \ll \Lambda_b$  breakdown scale ~500 MeV



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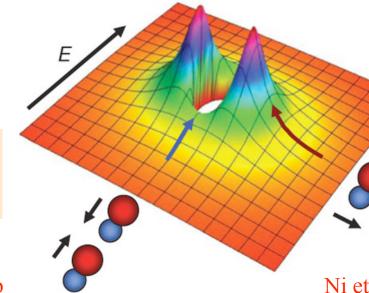
large scattering length physics

Separation of scales: low momenta  $\frac{1}{\lambda} = Q \ll \Lambda_b$  breakdown scale ~500 MeV

include long-range pion physics details at short distance not resolved capture in few short-range couplings, fit to experiment once,  $\Lambda$ -dependent

NLO  $\mathcal{O}\left(\frac{Q^2}{\Lambda^2}\right)$ 

pion tensor/dipole interactions, compare to cold polar molecules



Weinberg, van Kolck, Kaplan, Savage, Wise, Epelb

Ni et al., Nature (2010)

Separation of scales: low momenta  $\frac{1}{\gamma} = Q \ll \Lambda_b$  breakdown scale ~500 MeV

include long-range pion physics details at short distance not resolved capture in few short-range couplings, fit to experiment once,  $\Lambda$ -dependent systematic: can work to desired accuracy and obtain error estimates can connect to lattice QCD several open problems regarding renormalization and power counting

Separation of scales: low momenta  $\frac{1}{1} = Q \ll \Lambda_b$  breakdown scale ~500 MeV NN accurate reproduction of low-energy NN scattering at N<sup>3</sup>LO LO  $\mathcal{O}\left(\frac{Q^0}{\Lambda^0}\right)$ Phase Shift [deg] NLO  $\mathcal{O}\left(\frac{Q^2}{\Lambda^2}\right)$ Phase Shift [deg] -20 -30 Phase Shift [deg]  $^{1}D_{2}$  $N^2LO \mathcal{O}\left(\frac{Q^3}{\Lambda^3}\right)$ Phase Shift [deg]

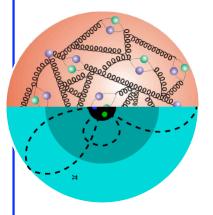
Lab. Energy [MeV]

Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Kaiser, Meissner, Meiss

## Nuclear forces and the Renormalization Group (RG)

RG evolution to lower resolution/cutoffs

$$H(\Lambda) = T + V_{NN}(\Lambda) + V_{3N}(\Lambda) + V_{4N}(\Lambda) + \dots$$



 $\Lambda_{
m chiral}$ 



## Nuclear forces and the Renormalization Group (RG)

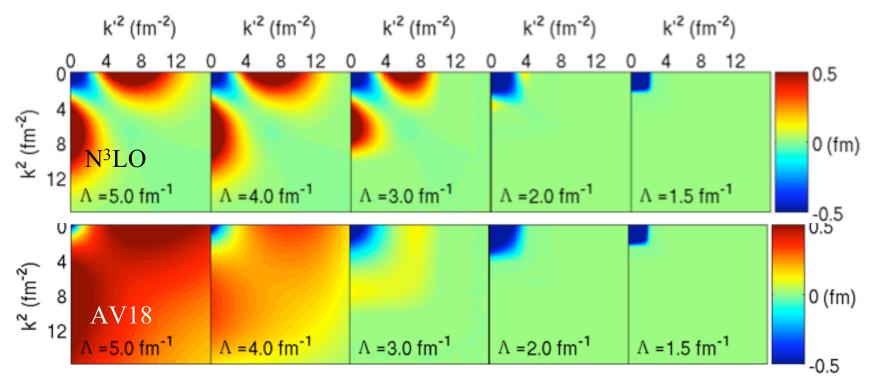
RG evolution to lower resolution/cutoffs

$$H(\Lambda) = T + V_{\rm NN}(\Lambda) + V_{\rm 3N}(\Lambda) + V_{\rm 4N}(\Lambda) + \dots$$

for NN interactions (preserves NN observables)

$$\frac{d}{d\Lambda} V_{\text{low } k}(k', k) = \frac{2}{\pi} \frac{V_{\text{low } k}(k', \Lambda) T_{\text{low } k}(\Lambda, k; \Lambda^2)}{1 - (k/\Lambda)^2}$$

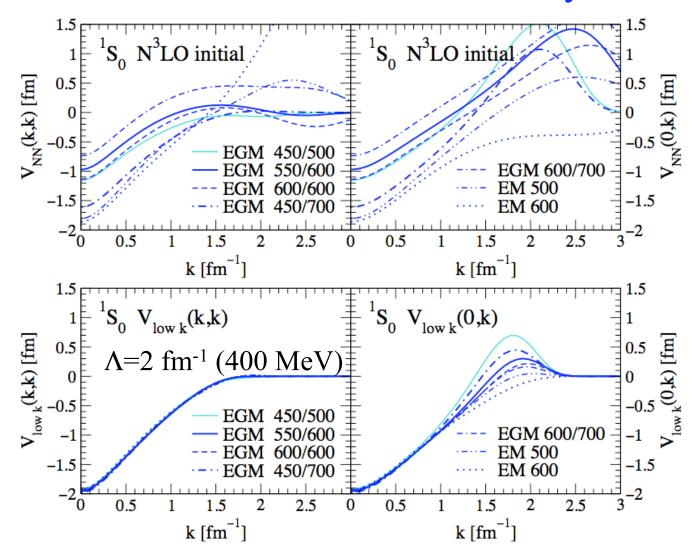
Bogner, Kuo, AS, Furnstahl,...



low-momentum interactions  $V_{low k}(\Lambda)$ 

RG decouples low-momentum physics from high momenta

## Low-momentum universality



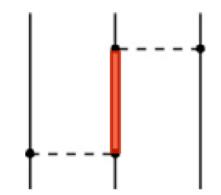
 $\approx$  universality from different chiral N<sup>3</sup>LO potentials

A similar NN universality can also be found in many-body systems, especially for properties involving neutrons.

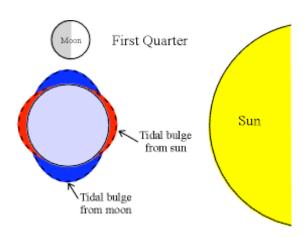
### Why are there three-nucleon (3N) forces?

Nucleons are finite-mass composite particles, can be excited to resonances

dominant contribution from  $\Delta(1232 \text{ MeV})$ 

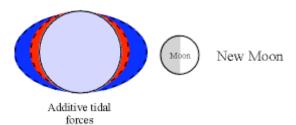


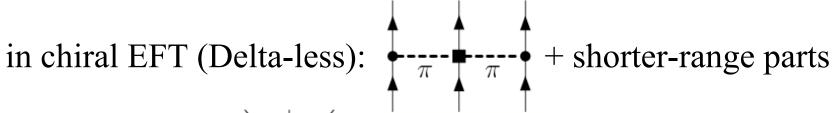
+ many shorter-range parts



Sun

tidal effects lead to 3-body forces in earth-sun-moon system

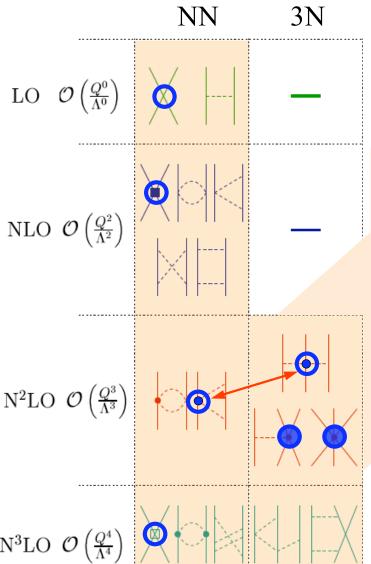




in pionless EFT:

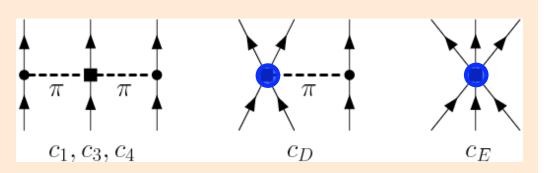


Separation of scales: low momenta  $\frac{1}{\lambda} = Q \ll \Lambda_b$  breakdown scale ~500 MeV



consistent **NN-3N** interactions

3N,4N: only 2 new couplings to N<sup>3</sup>LO



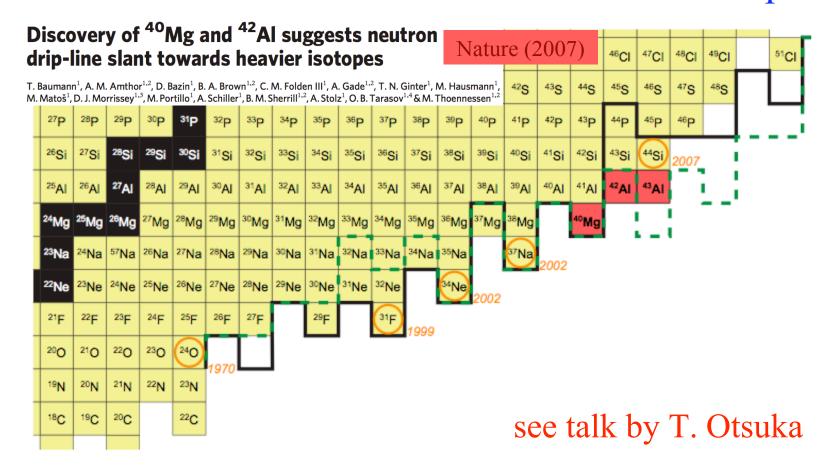
 $c_i$  from  $\pi N$  and NN Meissner et al. (2007)

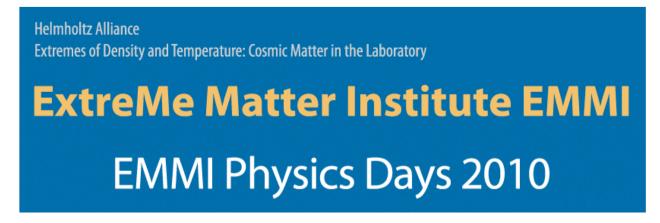
$$c_1 = -0.9^{+0.2}_{-0.5}$$
,  $c_3 = -4.7^{+1.2}_{-1.0}$ ,  $c_4 = 3.5^{+0.5}_{-0.2}$ 

single-
$$\Delta$$
:  $c_1=0$ ,  $c_3=-c_4/2=-3$  GeV<sup>-1</sup>

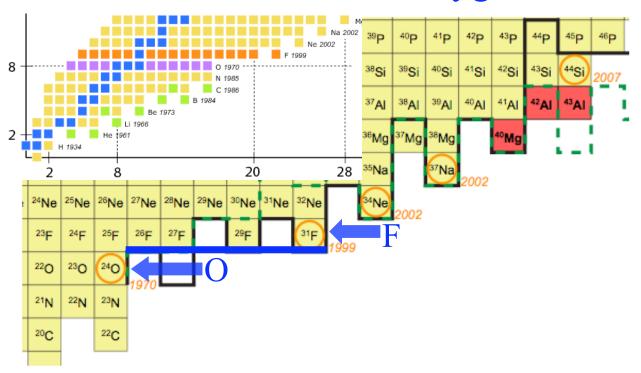
c<sub>D</sub>, c<sub>E</sub> fit to <sup>3</sup>H binding energy and <sup>4</sup>He radius (or <sup>3</sup>H beta decay half-life)

## Towards the limits of existence - the neutron drip-line

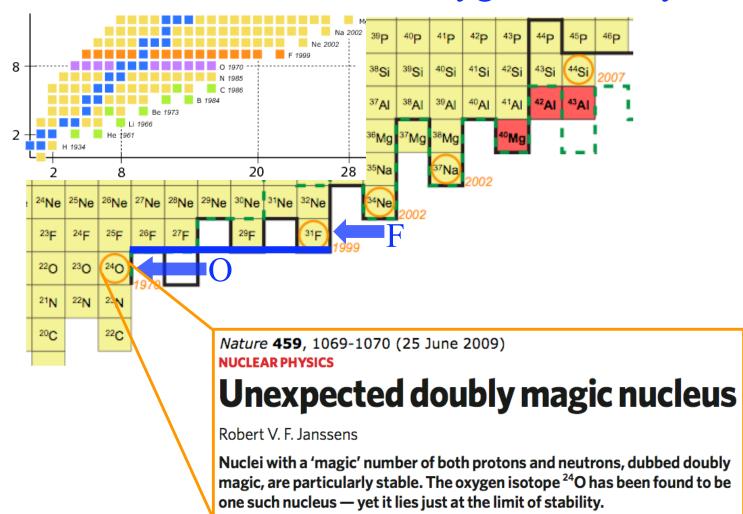




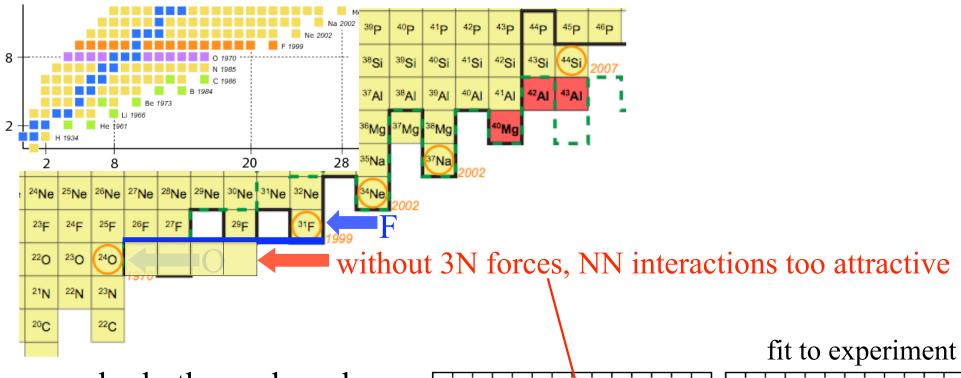
# The oxygen anomaly



## The oxygen anomaly

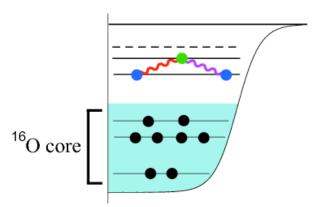


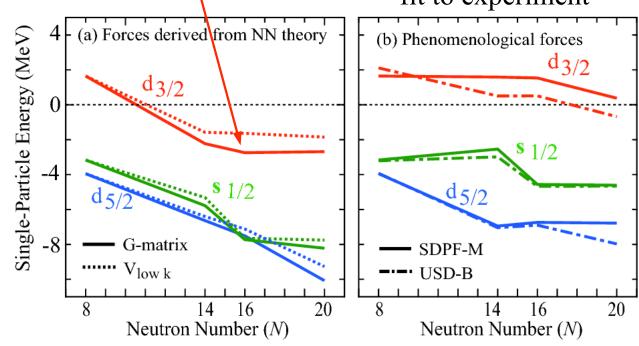
## The oxygen anomaly - not reproduced without 3N forces



many-body theory based on two-nucleon forces:

drip-line incorrect at <sup>28</sup>O





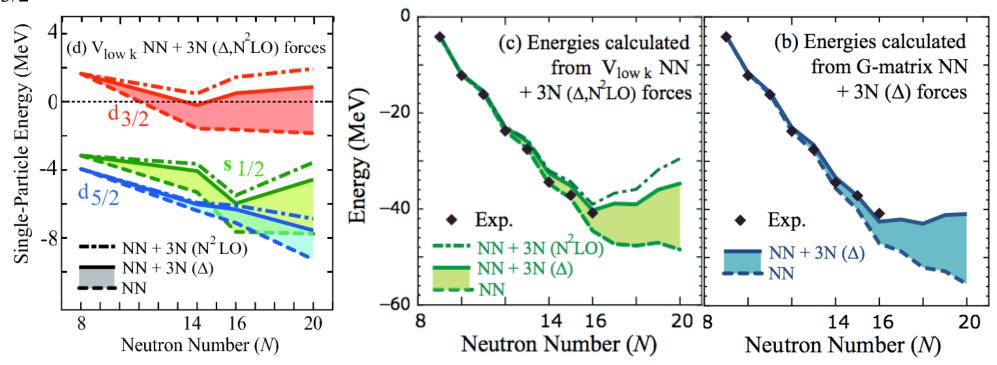
## The oxygen anomaly - impact of 3N forces

include "normal-ordered" 2-body part of 3N forces (enhanced by core A)

leads to repulsive interactions between valence neutrons can understand partly based on Pauli principle for  $3N(\Delta)$ 

 $3N(\Delta)$   $^{16}O core$ 

d<sub>3/2</sub> orbital remains unbound from <sup>16</sup>O to <sup>28</sup>O



first microscopic explanation of the oxygen anomaly

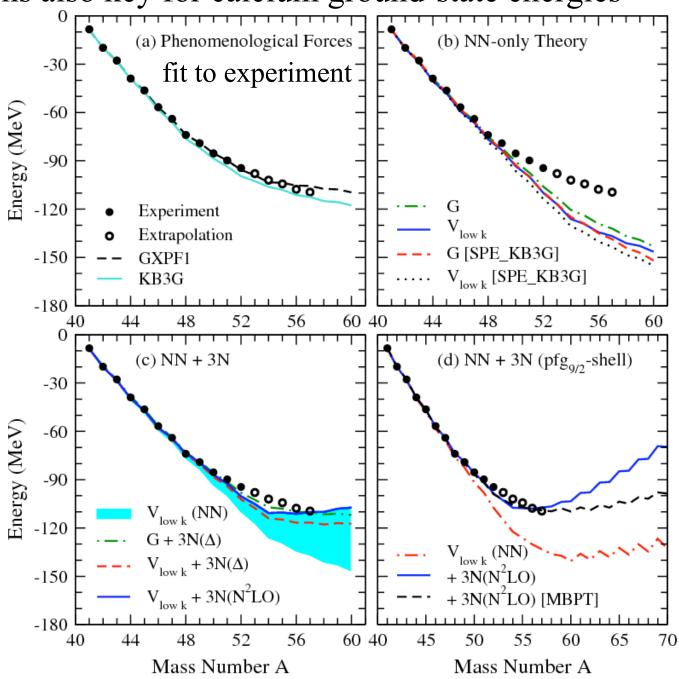
Otsuka, Suzuki, Holt, AS, Akaishi (2010)

## Evolution to neutron-rich calcium isotopes

repulsive 3N contributions also key for calcium ground-state energies

Holt, Otsuka, AS, Suzuki (2010)

mass measured to <sup>52</sup>Ca shown to exist to <sup>58</sup>Ca



## Evolution to neutron-rich calcium isotopes

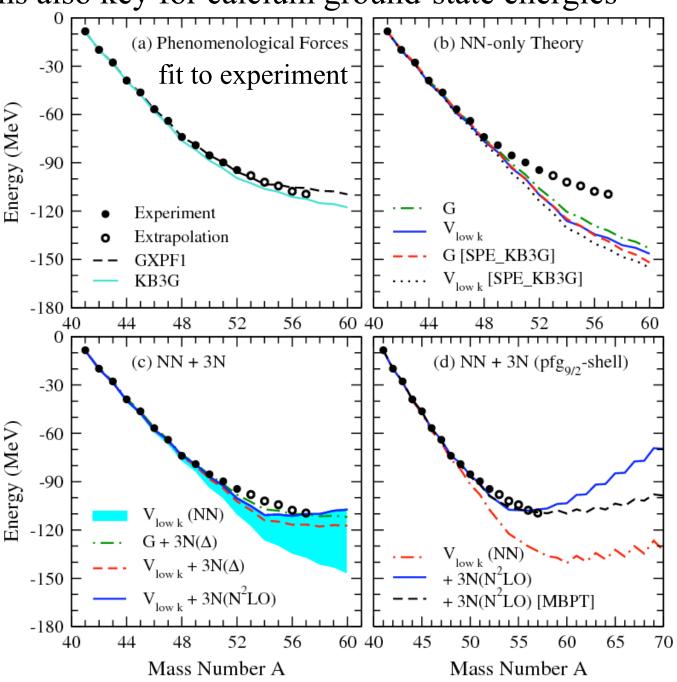
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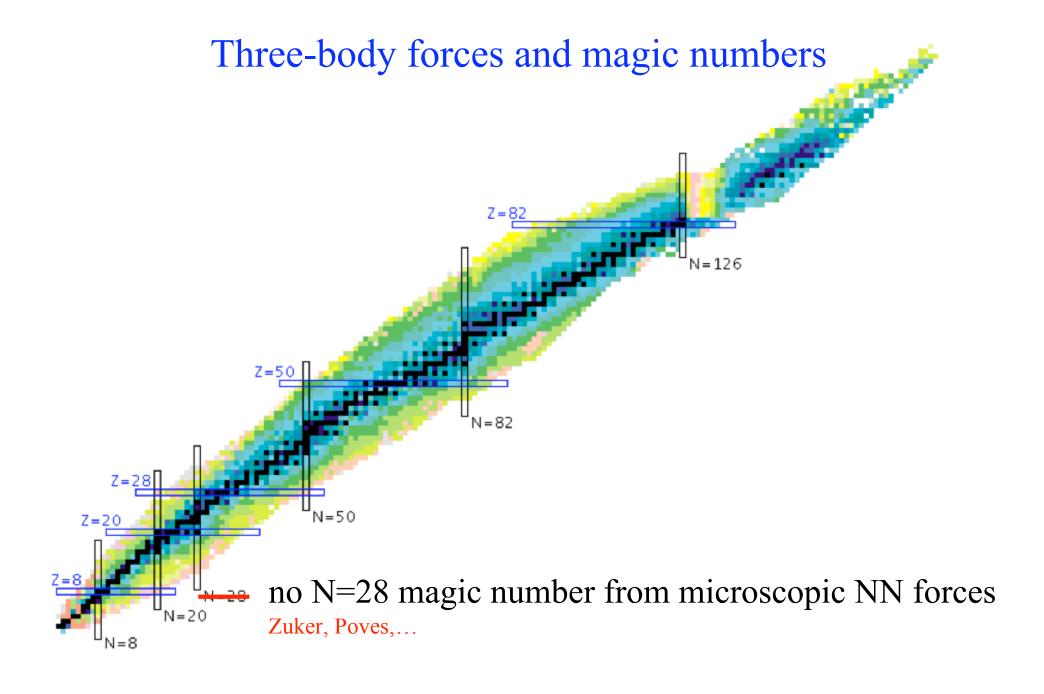
Holt, Otsuka, AS, Suzuki (2010)

mass measured to <sup>52</sup>Ca shown to exist to <sup>58</sup>Ca

predict drip-line around <sup>60</sup>Ca

continuum contributions need to be included





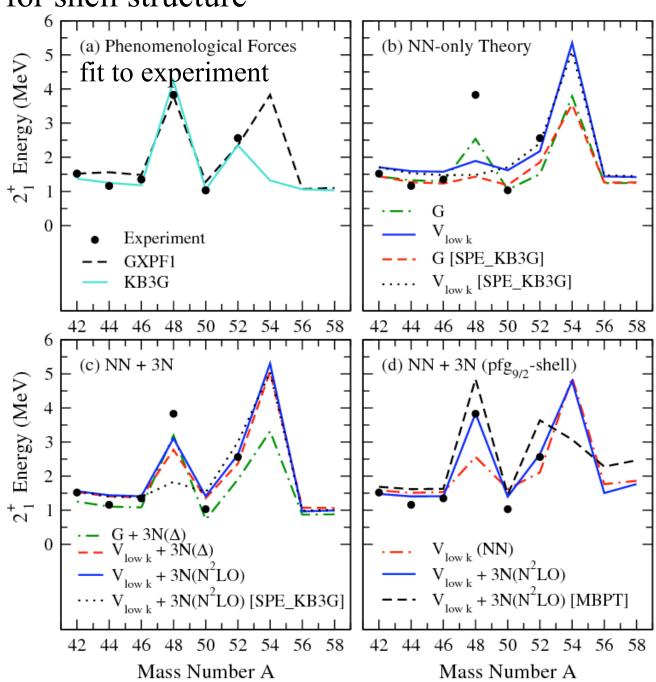
## Three-body forces and magic numbers

3N mechanism important for shell structure

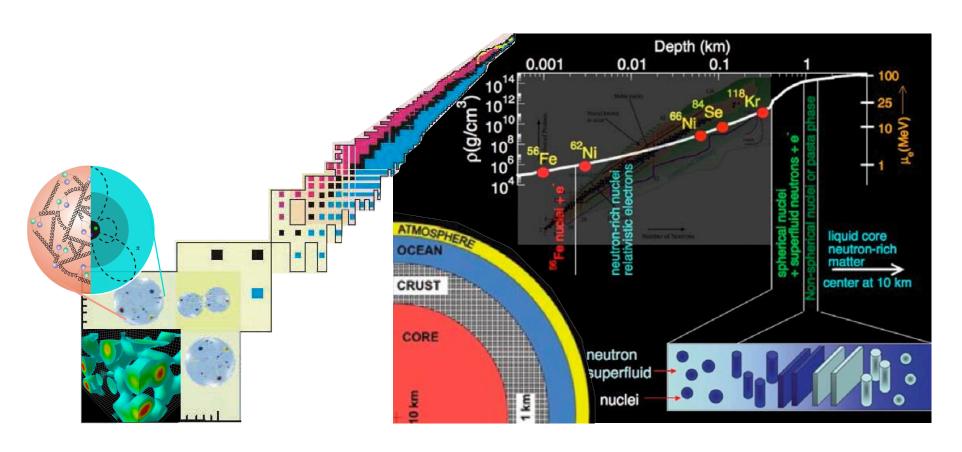
Holt, Otsuka, AS, Suzuki (2010)

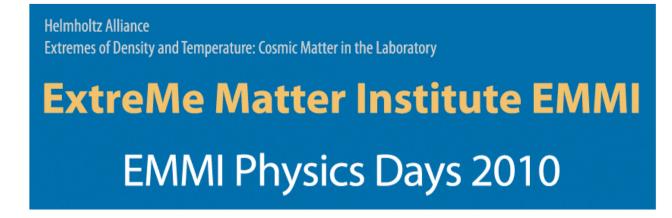
N=28 shell closure due to 3N forces and single-particle effects (<sup>41</sup>Ca)

N=34: predict high 2+ excitation energy in <sup>54</sup>Ca at 3-5 MeV



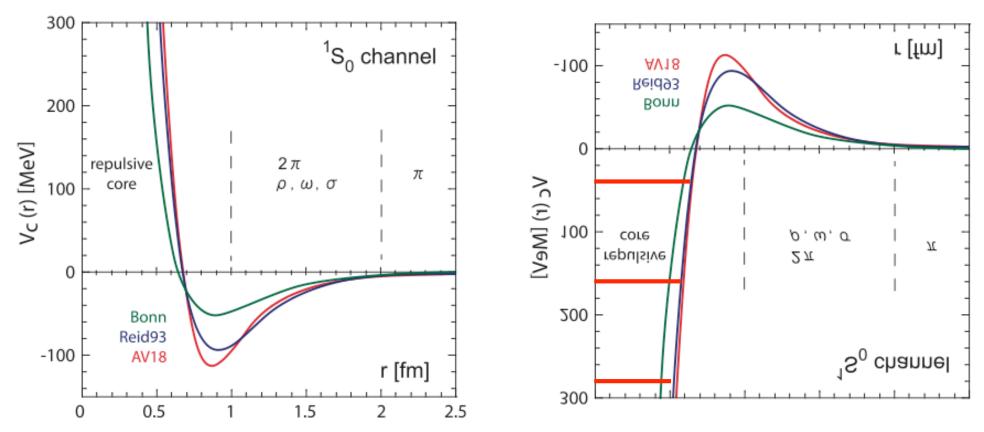
#### Extreme neutron-rich matter in stars





# Convergence with low-momentum interactions

large cutoffs lead to flipped-potential bound states, even for small - $\lambda V$  requires nonperturbative expansion, leads to slow convergence for nuclei

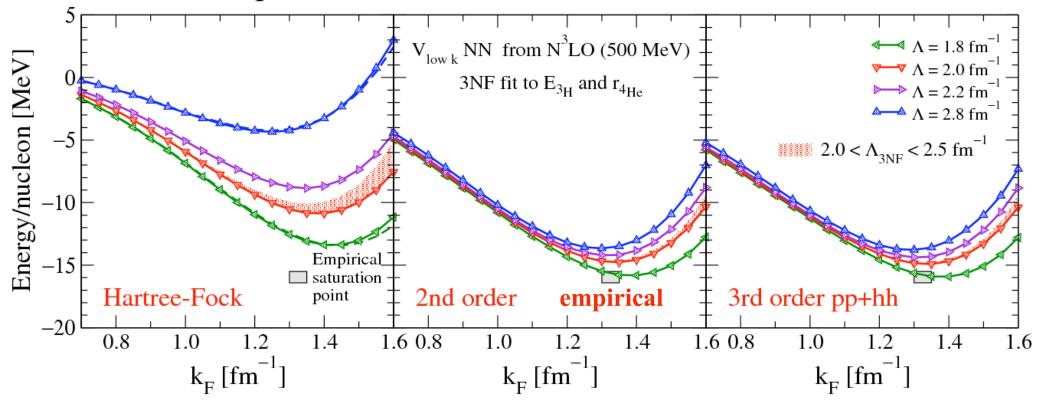


Weinberg eigenvalue analysis: two-body scattering becomes perturbative after RG evolution, except in channels with bound states

RG leads to improved convergence for nuclei and nuclear matter

# Advances in nuclear matter theory

Is nuclear matter perturbative with chiral EFT and RG evolution?



Hebeler, Bogner, Furnstahl, Nogga, AS, in prep. and (2009)

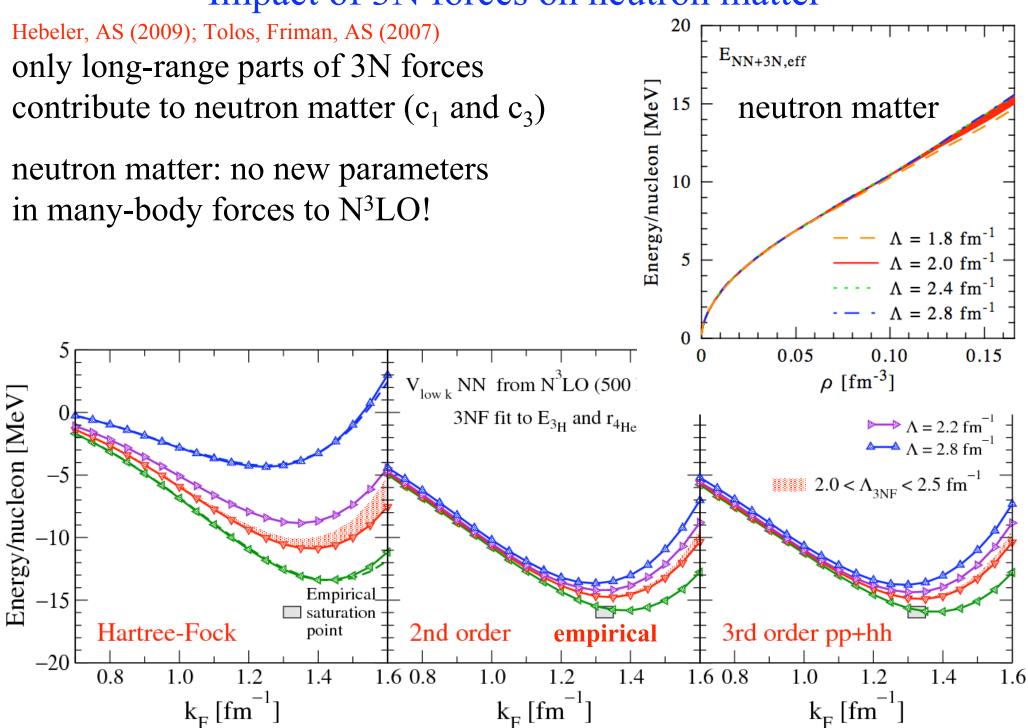
exciting: empirical saturation with theoretical uncertainties improved 3N treatment see also Holt, Kaiser, Weise (2010)

input to develop a universal energy density functional for all nuclei

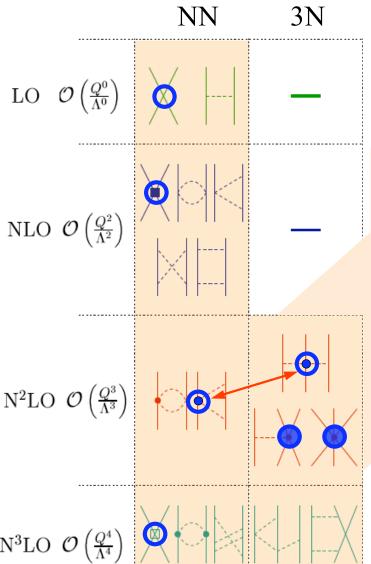
**UNEDF SciDAC Collaboration** 

Universal Nuclear Energy Density Functional

## Impact of 3N forces on neutron matter

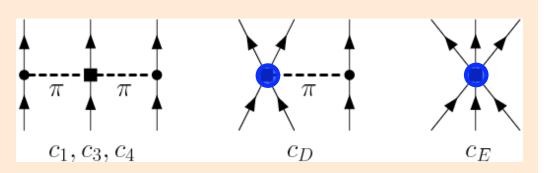


Separation of scales: low momenta  $\frac{1}{\lambda} = Q \ll \Lambda_b$  breakdown scale ~500 MeV



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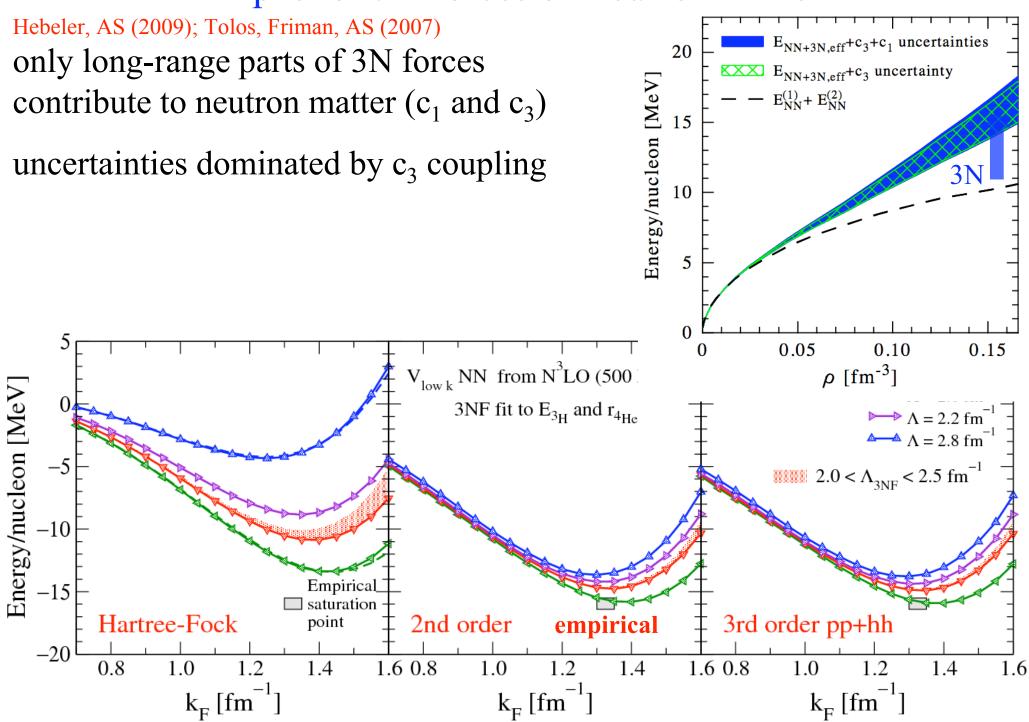
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,  $c_3 = -4.7^{+1.2}_{-1.0}$ ,  $c_4 = 3.5^{+0.5}_{-0.2}$ 

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c<sub>D</sub>, c<sub>E</sub> fit to <sup>3</sup>H binding energy and <sup>4</sup>He radius (or <sup>3</sup>H beta decay half-life)

## Impact of 3N forces on neutron matter



## Impact of 3N forces on neutron matter

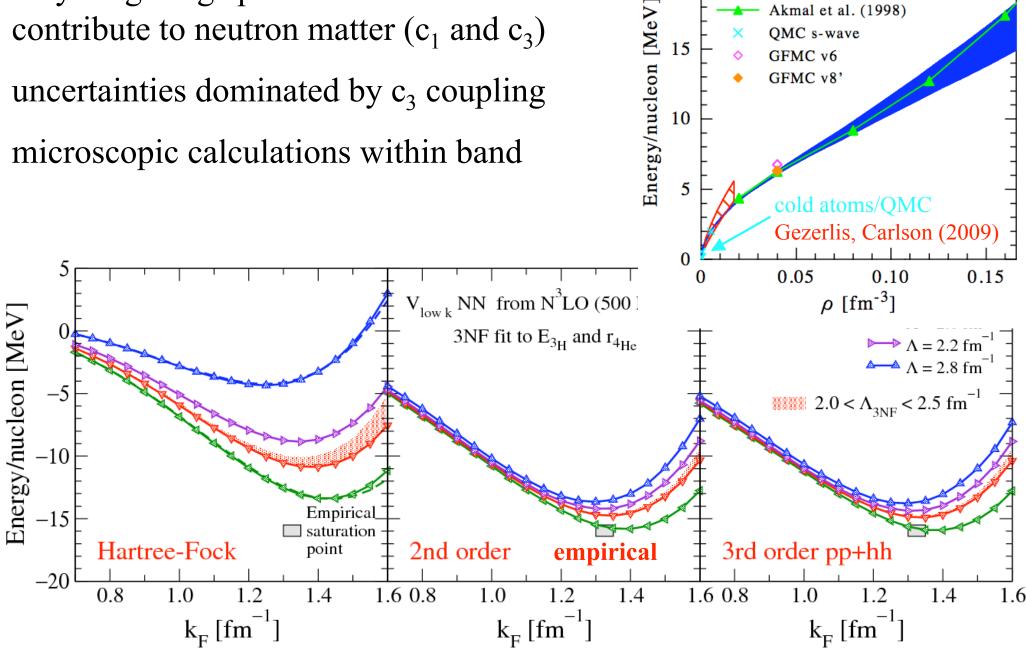
20

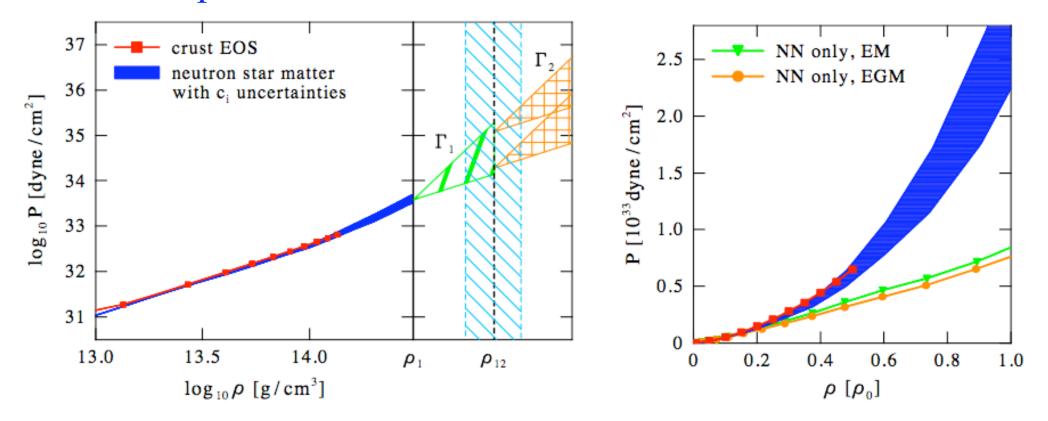
Schwenk+Pethick (2005)

Akmal et al. (1998)

QMC s-wave GFMC v6 GFMC v8'

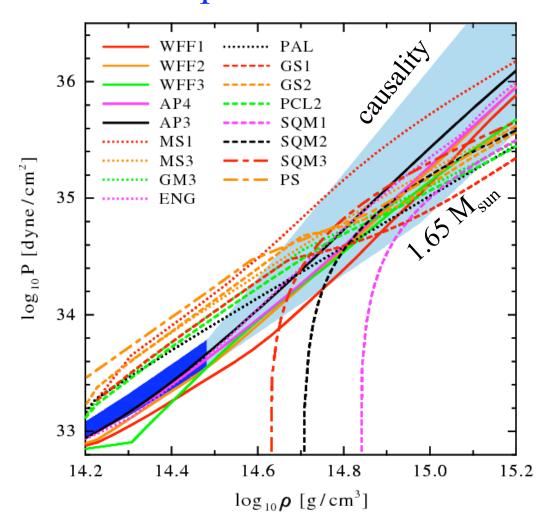
Hebeler, AS (2009); Tolos, Friman, AS (2007) only long-range parts of 3N forces contribute to neutron matter ( $c_1$  and  $c_3$ ) uncertainties dominated by c<sub>3</sub> coupling microscopic calculations within band



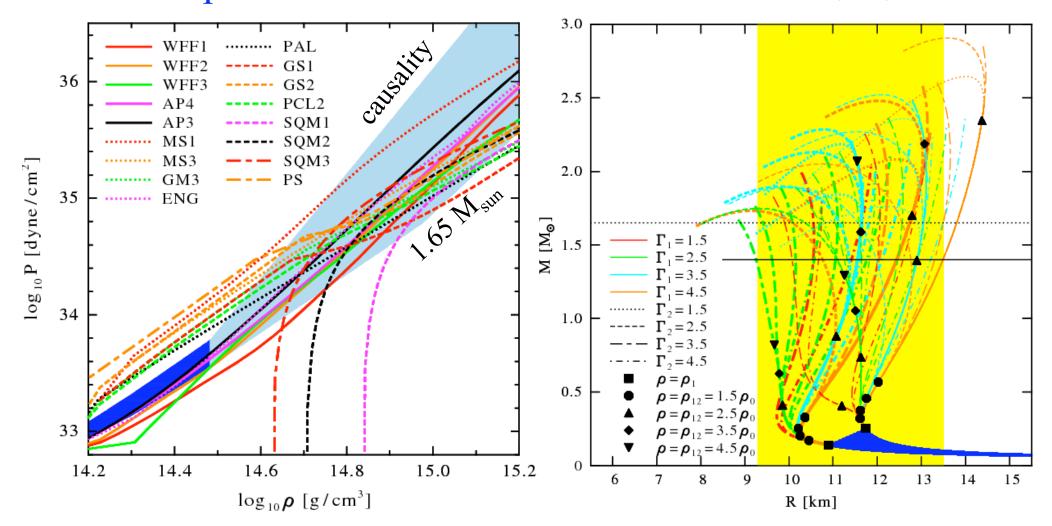


pressure below nuclear densities agrees with standard crust EOS only after 3N forces are included

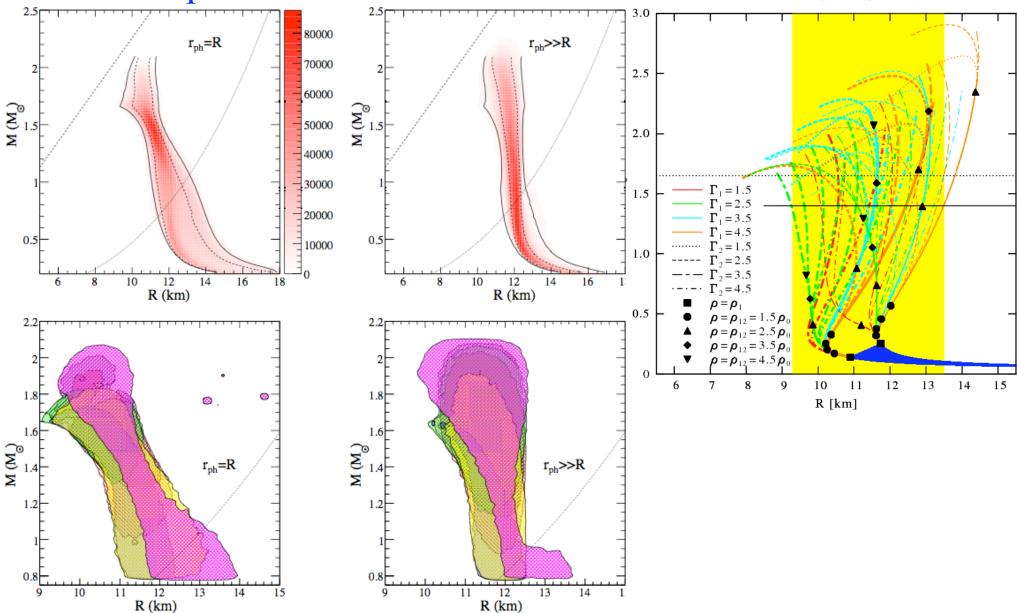
extend uncertainty band to higher densities using piecewise polytropes constrain polytropes by causality and require to support 1.65  $M_{sun}$  star



low-density pressure sets scale, our results reduce spread at nuclear densities in current neutron star modeling from factor 6 to ±25%



low-density pressure sets scale, our results reduce spread at nuclear densities in current neutron star modeling from factor 6 to  $\pm 25\%$  constrains neutron star radius to 9.7-13.9 km for M=1.4  $M_{sun}$ 



constrains neutron star radius to 9.7-13.9 km for M=1.4  $M_{sun}$  consistent with modeling of X-ray burst sources Steiner, Lattimer, Brown (2010)

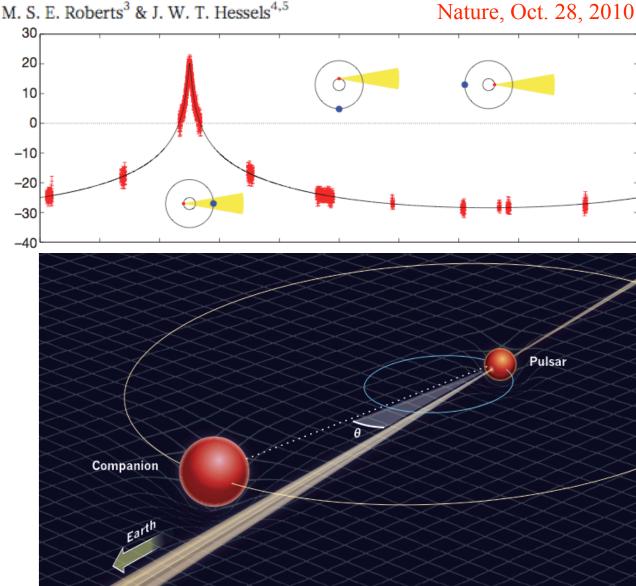
# A two-solar-mass neutron star measured using Shapiro delay

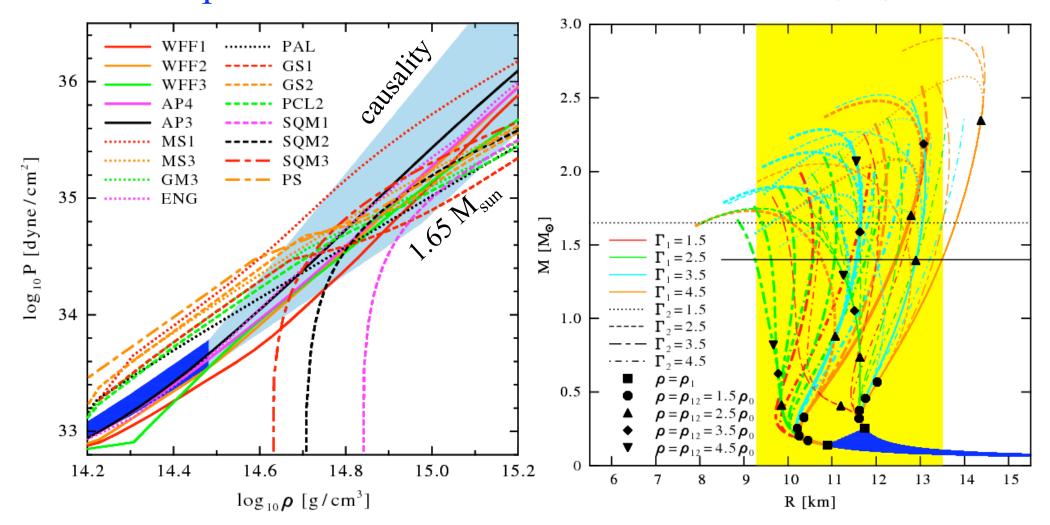
P. B. Demorest<sup>1</sup>, T. Pennucci<sup>2</sup>, S. M. Ransom<sup>1</sup>, M. S. E. Roberts<sup>3</sup> & J. W. T. Hessels<sup>4,5</sup>

direct measurement of neutron star mass from increase in travel time near companion

J1614-2230 most edge-on binary pulsar known (89.17°) + massive white dwarf companion (0.5 M<sub>sun</sub>)

heaviest neutron star with 1.97±0.04 M<sub>sun</sub>





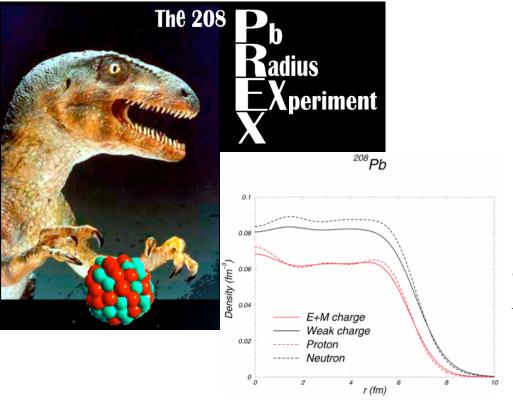
low-density pressure sets scale, our results reduce spread at nuclear densities in current neutron star modeling from factor 6 to  $\pm 25\%$  constrains neutron star radius to 9.7-13.9 km for M=1.4  $M_{sun}$  using 2.0  $M_{sun}$  constraint: 10.9-13.9 km (12% uncertainty!)

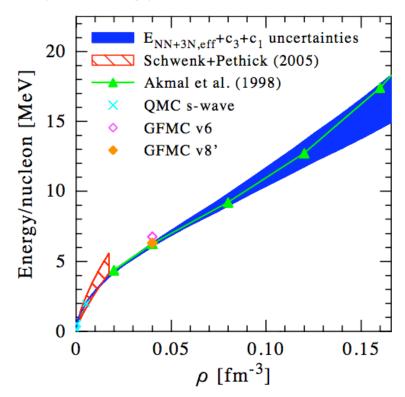
## Symmetry energy and neutron skin Hebeler, Lattimer, Pethick, AS (2010)

neutron matter band predicts range for symmetry energy 30.1-34.3 MeV

$c_1  [\mathrm{GeV}^{-1}]$	$c_3  [\mathrm{GeV}^{-1}]$	$\overline{S}_2 [{ m MeV}]$
-0.7	-2.2	30.1
-1.4	-4.8	34.4
NN-only EM		26.5
NN-onl	25.6	

and neutron skin of <sup>208</sup>Pb to 0.17±0.03 fm





compare to ±0.05 fm uncertainty goal of PREX @ JLAB



#### Thanks to collaborators!

J. Menendez



S.K. Bogner



R.J. Furnstahl,

K. Hebeler



A. Nogga



OAK RIDGE NATIONAL LABORATORY



T. Otsuka



T. Suzuki



Y. Akaishi



C.J. Pethick



J.M. Lattimer



D. Gazit

## Summary: From universal properties to neutron-rich extremes

Exciting era with advances on many fronts: development of effective field theory and the renormalization group

enables a unified description from nuclei to matter in astrophysics

universal properties and strong-interaction physics (corrections?)

3N forces are a frontier for neutron-rich nuclei/matter:

key to explain why <sup>24</sup>O is the heaviest oxygen isotope

N=28 magic number in calcium

dominant uncertainty of neutron (star) matter below nuclear densities, constraints on neutron star radii

intersections with atomic, condensed-matter, particle and astrophysics