



# Relativistic, quantum electrodynamic and nuclear effects in highly charged ions

Zoltán Harman

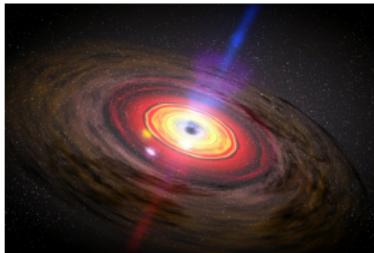
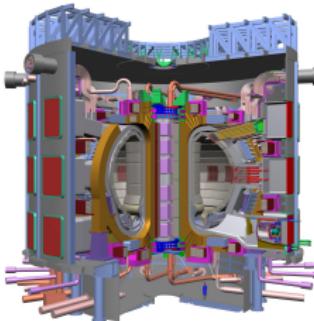
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ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für  
Schwerionenforschung GmbH, Germany

EMMI Physics Days  
November 4-5, 2010

# Motivation – Why to study HCI?



- Applications: magnetically confined **fusion** plasmas (tokamaks), **astrophysical** and technical plasmas
- sensitive tests of relativity, relativistic electron correlation
- tests of strong-field QED, nuclear effects (e.g. isotope shifts)

# Outline

- 1 Dielectronic recombination with H-like  $U^{91+}$
- 2 Higher-order resonant recombination processes
- 3 Two-center dielectronic recombination
- 4 X-ray resonant photoionization
- 5 Nuclear effects and isotope shifts
- 6 Summary

# People

- **Theory, Max-Planck-Institut für Kernphysik:**

Z. H., O. Postavaru, C. H. Keitel  
C. Müller, A. B. Voitkiv

- **Theory, St. Petersburg State University:**

I. I. Tupitsyn

- **Theory, Uni Heidelberg:**

A. N. Artemyev

- **ESR Experiment, GSI Darmstadt & Uni Giessen:**

C. Brandau, D. Bernhardt, C. Kozhuharov, A. Müller,  
S. Schippers, E. W. Schmidt, A. Gumberidze, F. Bosch,  
H.-J. Kluge, Th. Stöhlker, *et al.*

- **EBIT Experiment, Max-Planck-Institut für Kernphysik:**

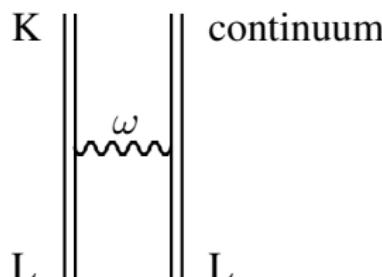
M. C. Simon, C. Beilmann, M. Schwarz, S. W. Epp, V. Mäckel,  
K. Kubicek, P. H. Mokler, H. Tawara, J. R. Crespo  
López-Urrutia, J. Ullrich, *et al.*

Hydrogenlike ions: highly stripped, only one electron left  
most simple system for studying electron interaction in a dynamic  
recombination process

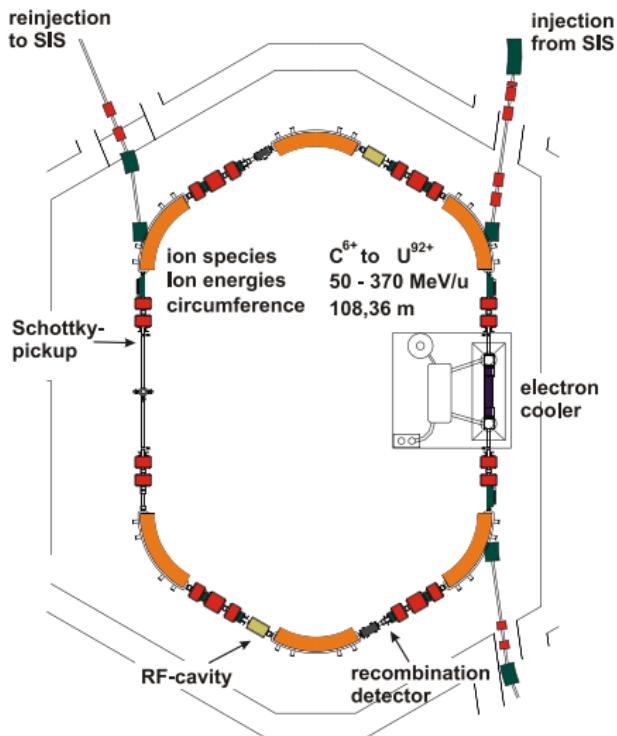
The electron-electron interaction operator is given as the sum of the  
Coulomb and Breit terms:

- static Coulomb interaction
- magnetic (Gaunt) term ("Lorentz force")
- scalar and transverse interaction retardation terms
- higher-order frequency-dependent retardation term

From: field theory, exchange of a virtual photon between the electrons

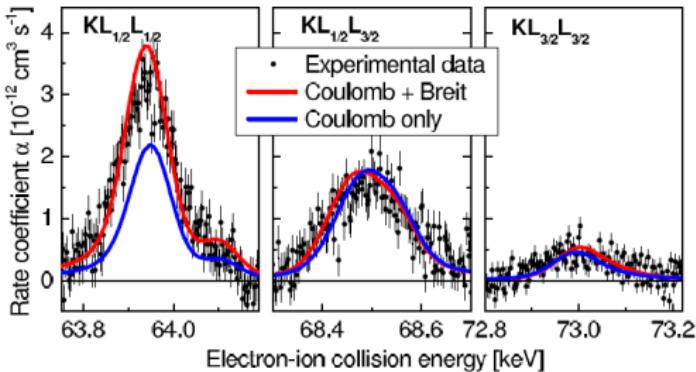
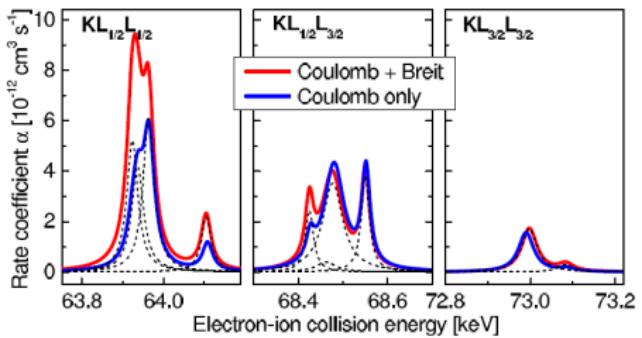


# DR with H-like U<sup>91+</sup> at the GSI ESR



- stochastically cooled ion beam
- electron cooler used as a target
- detection of recombined ions

(More on ESR: Alexandre Gumberidze,  
talk tomorrow 11:30)



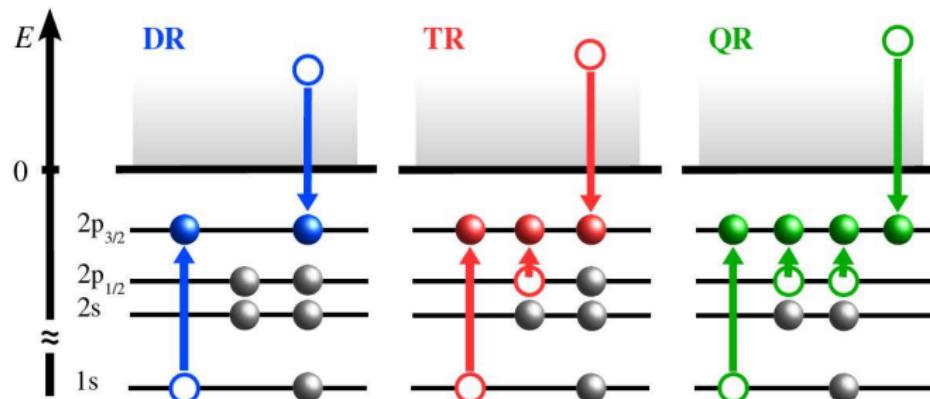
Agreement only if the Breit interaction is taken into account in the **dynamic** recombination process

D. Bernhardt, C. Brandau, Z. H., C. Kozuharov, A. Müller, *et al.*, submitted (2010)

# Higher-order resonant recombination processes

Dielectronic recombination (DR):

- Radiationless resonant capture of a continuum electron (inverse Auger effect)
- Radiative decay of the autoionizing state



3 electrons involved → trielectronic recombination (TR)

4 electrons involved → quadruelectronic recombination (QR)

## Total cross section

$T$ -matrix distorted wave formalism  $\rightarrow$  total cross section of DR (/TR/QR/. . .):

$$\sigma_{idf}(\varepsilon) = \frac{2\pi^2}{p^2} \frac{A_r^{df}}{\Gamma_d} L_d(\varepsilon) \frac{2J_d + 1}{2(2J_i + 1)} A_a^{di} \quad \text{with}$$

$$L_d = \frac{\Gamma_d/(2\pi)}{(E_i + \varepsilon - E_d)^2 + \Gamma_d^2/4}$$

The rates for autoionization and radiative decay are

$$A_a^{di} = \frac{2\pi}{2J_d + 1} \sum_{M_i m_s M_d} \int d\Omega_p |\langle \Phi_i J_i M_i, \vec{p} m_s | V_C + V_B | \Phi_d J_d M_d \rangle|^2 \rho_i$$

$$A_r^{df} = \frac{2\pi}{2J_d + 1} \sum_{M_d \lambda M_f} \int d\Omega_k \left| \left\langle \Phi_d J_d M_d | H_{er} | \Phi_f J_f M_f, \vec{k} \lambda \right\rangle \right|^2 \rho_f$$

# How to treat the electronic many-body problem?

Non-relativistically

Hartree: *independent particles*

$$\Psi_{1,2} = \phi_1^S \phi_2^S$$



Hartree-Fock: *Pauli principle*

$$\Psi_{1,2} = \frac{1}{\sqrt{2}} (\phi_1^S \phi_2^S - \phi_2^S \phi_1^S)$$



Multiconfiguration Hartree-Fock:

$$\Psi_{1,2} = \frac{1}{\sqrt{2}} \sum_{ij} c_{ij} (\phi_i^S \phi_j^S - \phi_j^S \phi_i^S)$$

Relativistically

Dirac-Hartree:

$$\Psi_{1,2} = \phi_1^D \phi_2^D$$



Dirac-(Hartree-)Fock:

$$\Psi(1, 2) = \frac{1}{\sqrt{2}} (\phi_1^D \phi_2^D - \phi_2^D \phi_1^D)$$



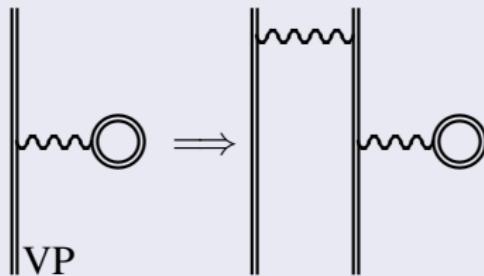
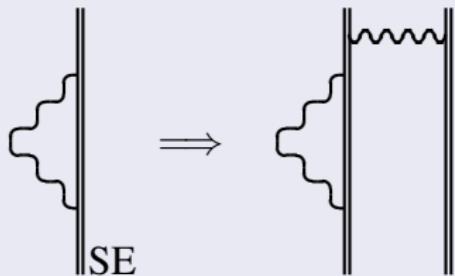
**Multiconfiguration Dirac-Fock:**

$$\Psi_{1,2} = \frac{1}{\sqrt{2}} \sum_{ij} c_{ij} (\phi_i^D \phi_j^D - \phi_j^D \phi_i^D)$$

# QED corrections: self-energy and vacuum polarization

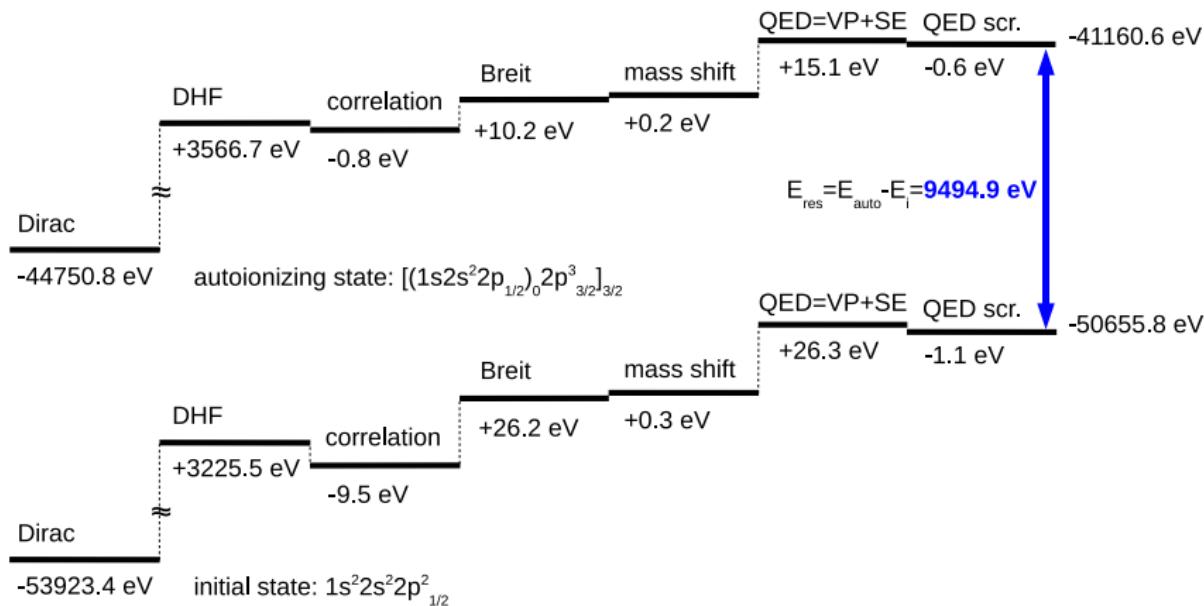
QED corrections:  $\approx 10\text{-}30 \text{ eV}$  in Kr

In many-body systems: interacting electrons, screening effects



Estimation of SE/VP screening by rescaling hydrogenic values

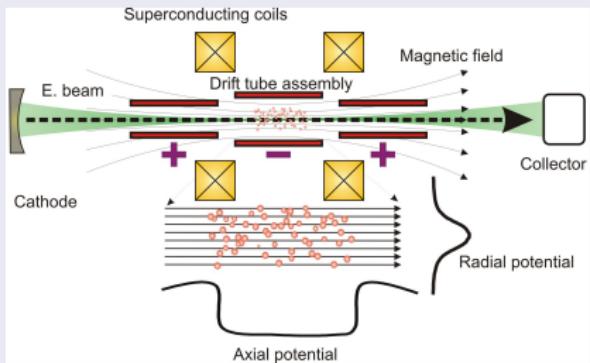
The bookkeeping part of theoretical physics:  
adding all the contributions to get one number



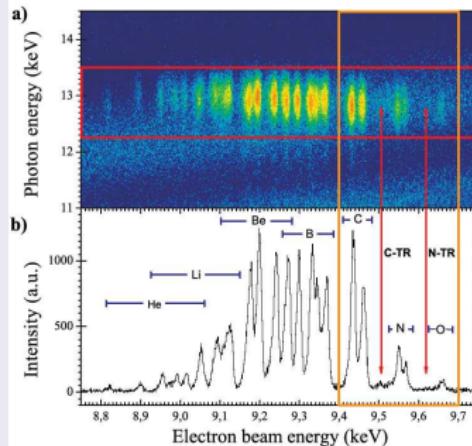
Line positions and cross sections are calculated

# Experiments at the Heidelberg EBIT

## Electron beam ion trap:

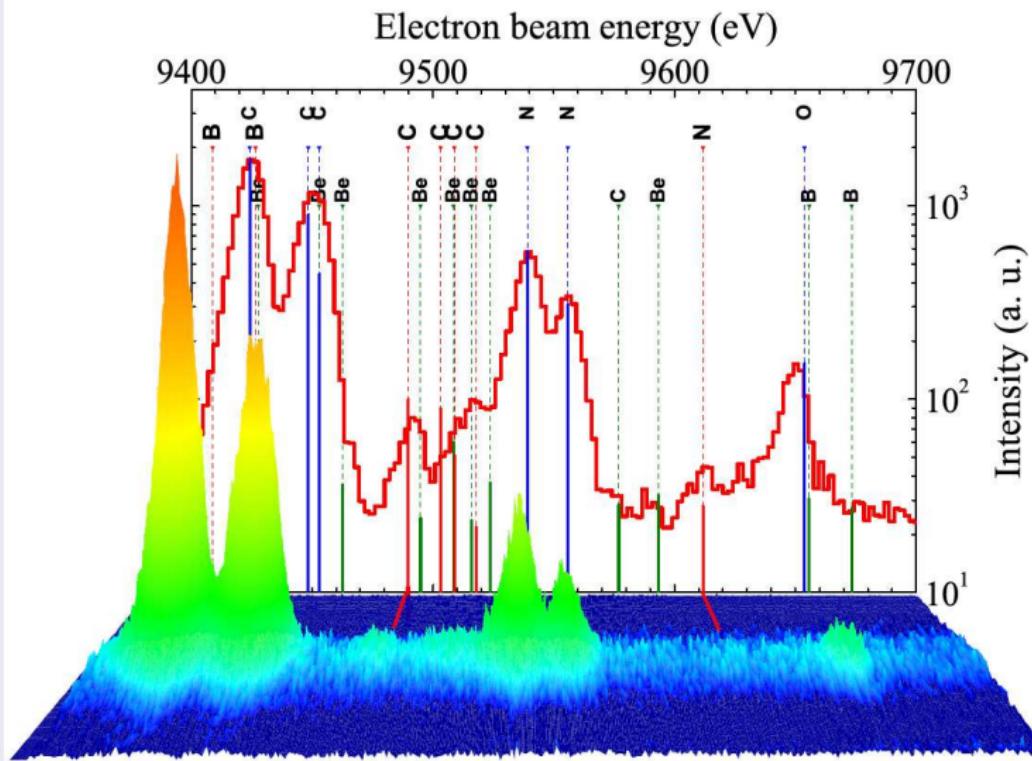


## Resonances in Kr ions

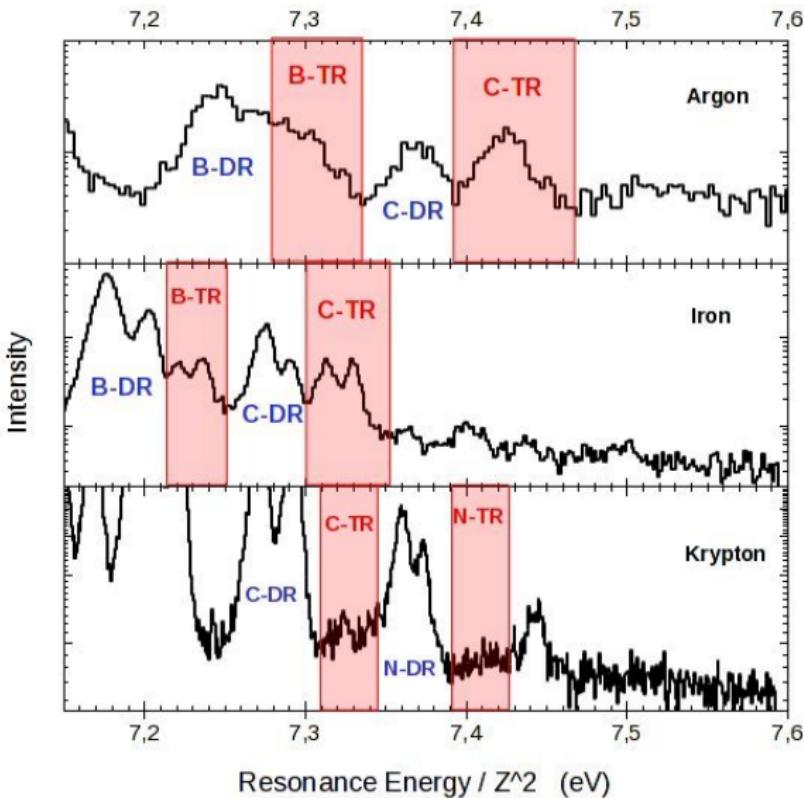


- C. Beilmann, O. Postavaru, L. H. Arntzen *et al.*, PRA **80**, 050702 (R) (2009)
- First observation of TR at the TSR of the MPIK: M. Schnell *et al.*, PRL, **91**, 043001 (2003)

## DR, TR, and QR resonances in Kr ions

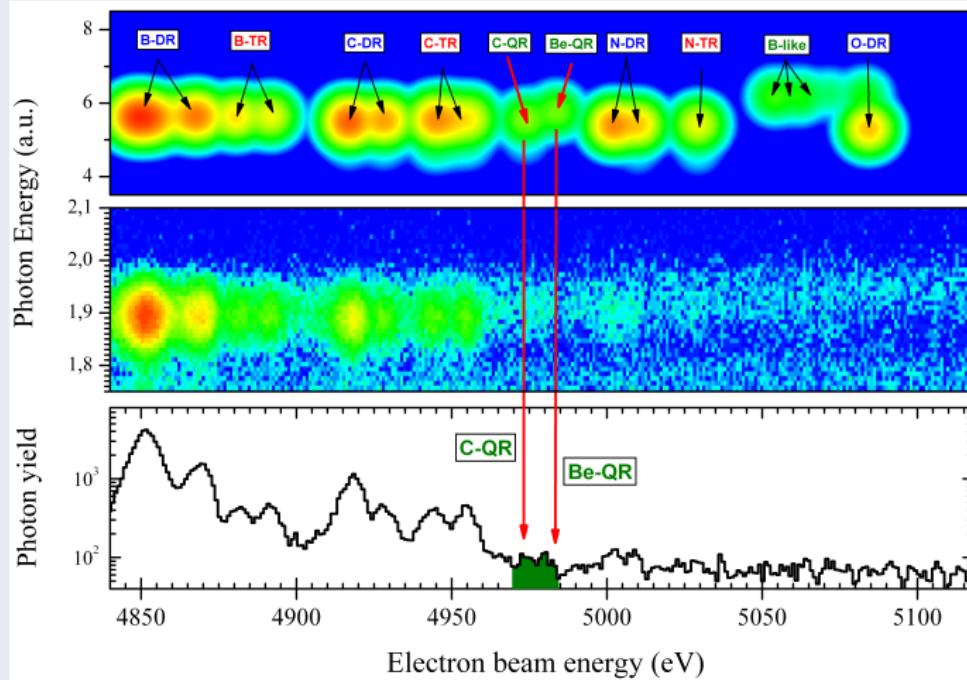


# Experimental results for lighter elements



# Quadruelectronic resonances in Fe ions

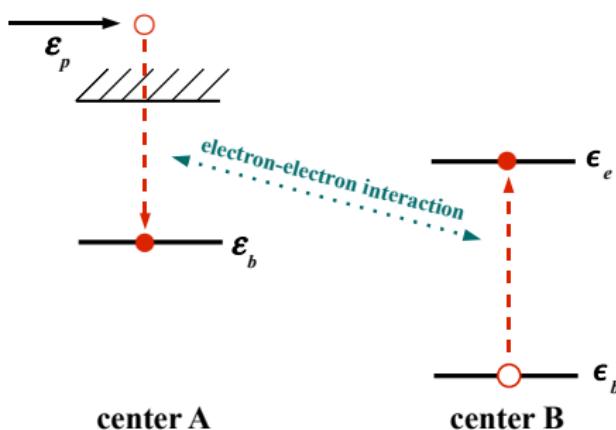
Fluorescence signal vs. electron energy **and** photon energy:



# Environment-assisted recombination

Two-center dielectronic recombination (2CDR) process:

- electron recombination into center A; simultaneous resonant excitation in center B (time-reversed inter-atomic Auger process)
- radiative decay of the excited state in center B



“Center”=atom/ion/molecule

Energy transfer between the centers: exchange of a virtual photon with  $\hbar\omega$

Competing process: radiative recombination (RR) with center A; 2CDR can exceed RR

Not observed yet (!)

## Ratio of the 2CDR at resonance and RR (approx.):

$$\frac{\sigma_{\text{2CDR}}}{\sigma_{\text{RR}}^{(A)}} \sim \left(\frac{c}{R\omega}\right)^6$$

- small distances  $R$  preferable  $\leftrightarrow$  large densities
- low energies preferable:  $\approx 1$  eV, outer shells, low charge states

### Examples and potential relevance:

$e^- + H^+ + He^+$  ( $1s-2p$ )

$e^- + H^+ + He$

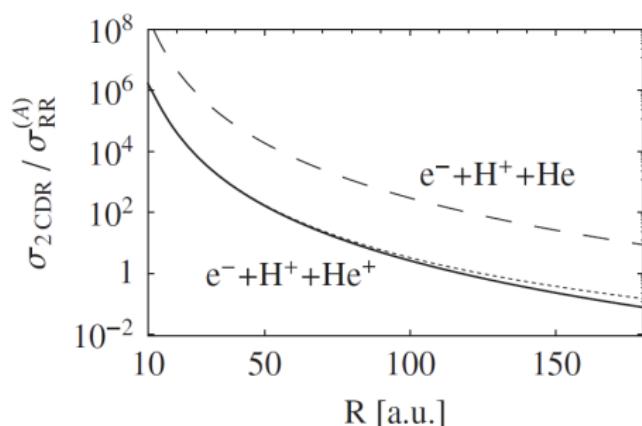
- astrophysics - Sun?

- thermonuclear fusion plasmas

- warm dense matter in

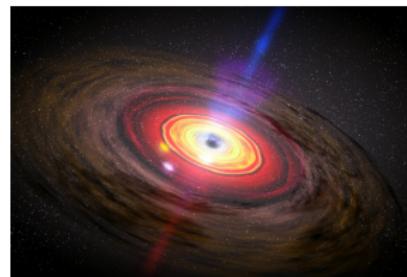
laser-solid interaction

$R=100 a_0 \approx 5$  nm



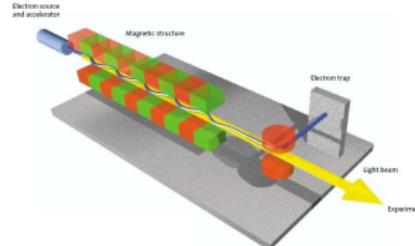
# Photoionization of HCl in the x-ray regime

- Application: **astrophysical plasmas**
- half of the baryonic matter in the universe is in the form of HCl's comprised in the Warm-Hot Intergalactic Medium (WHIM); PI may be the dominant ionization process
- analysis of ionic absorption spectra (C, N, O, Ne, Ar, Fe ions) → temperature, density, mass and velocity

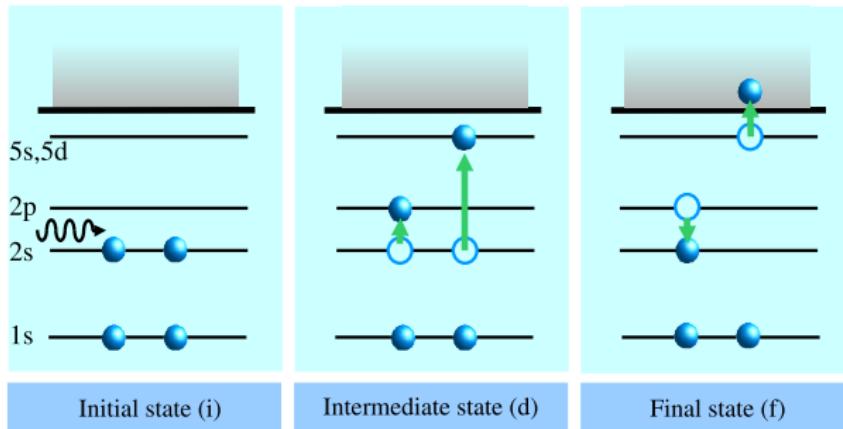


- PI data are important constituents of **astrophysical models**: reliable theories and measurements are called for
- interpretation of data from grating spectrometers on board of the Chandra and XMM Newton observatories

- **Photoionization of HCl:** novel light sources with high energies and brilliance (synchrotron radiation, XFEL lasers) → extending to higher ionic charge states
- high detection efficiency of the photoions: accurate measurements possible
- established methods: merged beam experiments, dual laser-produced plasmas, Penning traps
- new method: PI of HCl in an **electron beam ion trap (EBIT)**



Example: Be-like ions, excitation from the ground state  $1s^2 2s^2$



Excitation: two-electron one-photon transition

Photoionization via  $2p5s$ ,  $2p5d$ ,  $2p6s$ ,  $2p6d\dots$  states

Total cross section for a given photoionization channel

$i$  (initial)  $\rightarrow d$  (autoionizing)  $\rightarrow f$  (final):

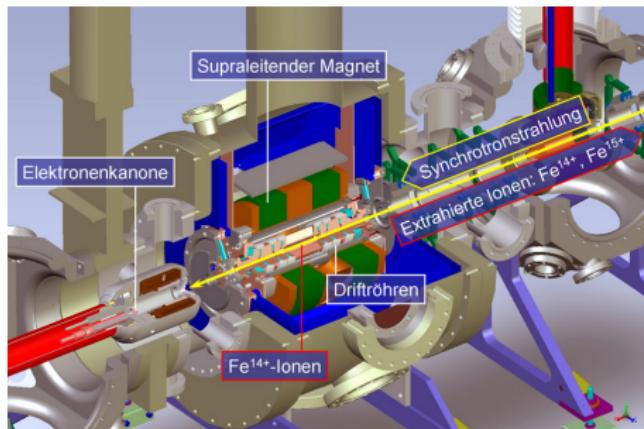
$$\sigma_{i \rightarrow d \rightarrow f}^{\text{PI}}(\hbar\omega) = \frac{2\pi^2 c^2 \hbar^4}{(\hbar\omega)^2} \frac{A_{id}^{re} A_{df}^a}{\Gamma_d} L_d(\hbar\omega)$$

Lorentzian line shape function:

$$L_d(\hbar\omega) = \frac{\Gamma_d/(2\pi)}{(\hbar\omega + E_i - E_d)^2 + \Gamma_d^2/4}$$

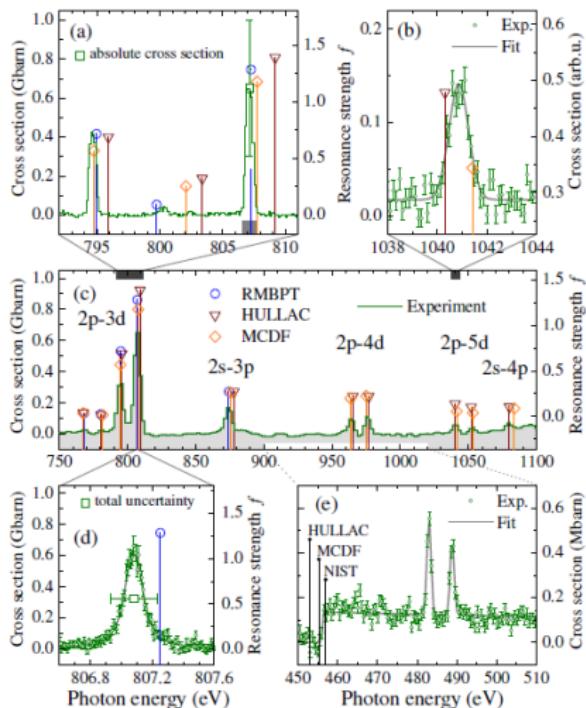
$A_d^r, A_d^a$ : radiative and Auger decay rates;  
 $A_{id}^{re}$ : radiative excitation rate

Experimental setup using synchrotron radiation and an EBIT equipped with ion extraction



HCI cloud area density:  $\approx 10^{10} \text{ cm}^{-2}$

# Experimental spectra for Mg-like Fe<sup>14+</sup> ( $E/\Delta E \approx 6500$ ); theory



- MCDF: this work  
10-20 eV dev.
- for the active galaxy  
NGC3783: velocity from  
Doppler shifts
- more brilliant x-ray beams  
from synchrotrons (e.g. Petra  
III) and x-ray free electron  
lasers (e.g. XFEL, LCLS) in  
the near future: multi-keV  
regime, multiphoton  
excitations

M. C. Simon, J. R. Crespo López-Urrutia, C. Beilmann, M. Schwarz, Z. H., *et al.*,  
PRL **105**, 183001 (2010)

M. C. Simon, M. Schwarz, S. W. Epp, C. Beilmann, B. L. Schmitt, et al., JPB. **43**,  
065003 (2010)

# Nuclear effects and isotope shifts

Nuclear property	Effect on electronic structure
charge $+Ze$	→ binding energy
size or radius $r_{\text{RMS}} = \sqrt{\langle r^2 \rangle}$	→ field shift (FS)
mass $M < \infty$ , nuclear recoil	→ normal and specific mass shift (MS)
polarizability	→ nuclear polarization (NP) shift
spin and magnetic moment	→ magnetic hyperfine splitting (HFS)
quadrupole moment	→ quadrupole hyperfine splitting
weak interaction	→ parity non-conservation (PNC)

These are usually small corrections to level and transition energies

Best studied by comparing two different isotopes, i.e. by measuring **isotope shifts**

E.g. two-photon laser spectroscopy:  $2s \rightarrow 3s$  transitions with short-lived Li isotopes →  $\delta r_{\text{RMS}}$  for  $^{8,9}\text{Li}$ ,  $^{9,11}\text{Li}$

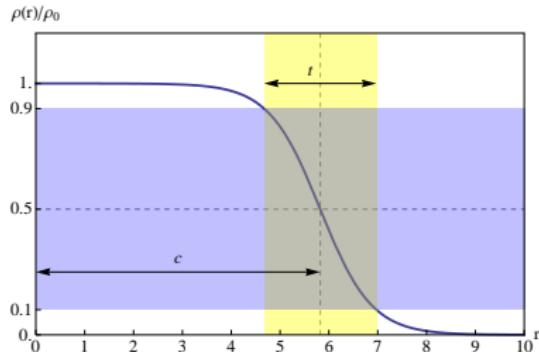
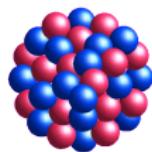
G. Ewald, W. Nörterhäuser, A. Dax, *et al.*, PRL (2004);

R. Sánchez, W. Nörterhäuser, G. Ewald, *et al.*, PRL (2006)

# Nuclear charge distribution

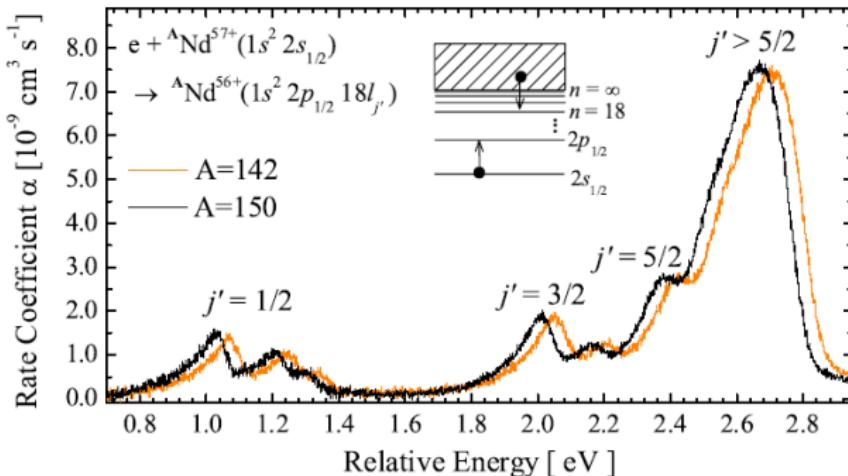
Fermi two-parameter distribution for describing the charge distribution of the considered nuclei:

$$\rho_{\text{nuc}}(r) = \frac{\rho_0}{1 + e^{(r-c)/a}}, \quad a = t 4 \ln 3$$



- Numerical integration of the many-electron Dirac equations with  $V_{\text{nuc}}$
- Nuclear size dependence of SE and VP terms: analytically calculated formulae
- Mass shifts taken into account

# DR isotope shift measurement at the ESR with Nd<sup>57+</sup>



High accuracy for low-lying resonances

$$\delta E_{\text{res}} = 40.2(3)(6) \text{ meV} \rightarrow \delta \langle r^2 \rangle = -1.36(1)(3) \text{ fm}^2$$

A new method to gain information on the nuclear size

C. Brandau, C. Kozhuharov, Z. H., A. Müller, S. Schippers *et al.*, PRL 100, 073201 (2008)

Similar measurements are planned with nuclei of non-zero spin:  
hyperfine splitting, nuclear moments

# Summary

- Dielectronic recombination with  $U^{91+}$ : **relativistic electron interaction**
- Theoretical and experimental K-LL DR and KL-LLL TR spectra for Be-, B-, C-, and N-like Kr ( $Z=36$ ), Fe ( $Z=26$ ), Ar ( $Z=18$ )  
Trielectronic and quadruelectronic recombination: **nonlinear electron-electron interaction**
- **Two-center** dielectronic recombination: new process in warm-dense matter and in chemical environments
- **Photoionization of HCl**: extending to higher charge states ( $14+$ ) and photon energies ( $\approx 1$  keV)  
identification and analysis of structures in astrophysical spectra
- **Isotope shifts** in HCl spectra: determination of nuclear properties
- Work in progress: QED and nuclear effects on the bound electron ***g* factor**

Grazie per l'attenzione!

Köszönöm a figyelmet!

Merci à tous pour votre attention!

Muchas gracias por su atención!

Mulțumesc pentru atenție!

Obrigado pela atenção!

Spasibo za vnimanie!

Thank you for your attention!

Vielen Dank für Ihre Aufmerksamkeit!