

E. Thomé

Outline

Introduction to Hyperons

Existing Experimental Data

Simulations fo PANDA

Precession of the Polarisation Vector in the Magnetic Field

Conclusions

Antihyperon Hyperon Physics for the PANDA Experiment

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1 Introduction to Hyperons

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Simulations for PANDA

• $\bar{p}p \rightarrow \bar{\Lambda}\Lambda$

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- $\bar{p}p \rightarrow \bar{\Xi}^+ \Xi^-$
- $\bar{p}p \rightarrow \bar{\Omega}^+ \Omega^-$
- $\bar{\mathrm{pp}} \rightarrow \bar{\Lambda}_c^- \Lambda_c^+$
- CP Violation



Precession of the Polarisation Vector in the Magnetic Field

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 $\bar{p}p\to \bar{Y}Y$

What are the relevant degrees of freedom to describe a process? Quarks and gluons or hadronic degrees of freedom?



PANDA (1.5-15 GeV) is in the transition regime between the two descriptions.

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Hyperon Spin Variables

Hyperons decay with a parity violating weak decay to a spin 1/2 baryon and aspin 0 meson $\Rightarrow L = 0, 1$ (L = 1, 2 for Ω). Decay asymmetry parameters

$$\alpha = \frac{2\text{Re}(A_{\mathcal{S}}^*A_{\mathcal{P}})}{|A_{\mathcal{S}}|^2 + |A_{\mathcal{P}}|^2}, \quad \beta = \frac{2\text{Im}(A_{\mathcal{S}}^*A_{\mathcal{P}})}{|A_{\mathcal{S}}|^2 + |A_{\mathcal{P}}|^2}, \quad \gamma = \frac{|A_{\mathcal{S}}|^2 - |A_{\mathcal{P}}|^2}{|A_{\mathcal{S}}|^2 + |A_{\mathcal{P}}|^2}$$
(1)

Polarisation P is given from angular distribution

$$I(\cos\Theta) = \frac{1}{2}(1 + \alpha P \cos\Theta)$$
 (2)

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In $\bar{p}p \rightarrow \bar{Y}Y$ spin correlations are also be measured. For $\bar{\Lambda}\Lambda$ this relates to how the spins of the $\bar{s}s$ pair are correlated. (β and γ appears in the angular distribution of the decay of the daughter particle)



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Definition of the Hyperon Rest System



The hyperon rest systems are defined to make maximum use of C and P invariance

$$I_{\bar{p}p}(\Theta_{\bar{\Lambda}}, \hat{k}_{\bar{p}}, \hat{k}_{p}) = \frac{I_{0}}{64\pi^{3}} \begin{pmatrix} 1 \\ +P_{y}(\bar{\alpha}\hat{k}_{\bar{p}y} + \alpha\hat{k}_{py}) \\ +C_{xx}\bar{\alpha}\alpha\hat{k}_{\bar{p}x}\hat{k}_{px} \\ +C_{yy}\bar{\alpha}\alpha\hat{k}_{\bar{p}y}\hat{k}_{py} \\ +C_{zz}\bar{\alpha}\alpha\hat{k}_{\bar{p}z}\hat{k}_{pz} \\ +C_{xz}\bar{\alpha}\alpha\hat{k}_{\bar{p}z}\hat{k}_{px} \\ +C_{zx}\bar{\alpha}\alpha\hat{k}_{\bar{p}z}\hat{k}_{px} \end{pmatrix}$$
(3)



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Data from PS185 at LEAR

Largest amount of data, \sim 40k events,from PS185 experiment at LEAR.





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Spin Correlations



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Singlet Fraction

$$SF = \frac{1}{4}(1 + Cxx - Cyy + Czz) \tag{4}$$



 $\overline{\Lambda}\Lambda$ (and consequently the \overline{ss}) are always produced with parallel spins.



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CP Violation Parameters

$$D = \frac{\Gamma - \overline{\Gamma}}{\Gamma + \overline{\Gamma}}$$

$$A = \frac{\Gamma \alpha + \overline{\Gamma} \overline{\alpha}}{\Gamma \alpha - \overline{\Gamma} \overline{\alpha}} \approx \frac{\alpha + \overline{\alpha}}{\alpha - \overline{\alpha}}$$

$$B = \frac{\Gamma \beta + \overline{\Gamma} \overline{\beta}}{\Gamma \beta - \overline{\Gamma} \overline{\beta}} \approx \frac{\beta + \overline{\beta}}{\beta - \overline{\beta}}$$

$$B' = \frac{\Gamma \beta + \overline{\Gamma} \overline{\beta}}{\Gamma \alpha - \overline{\Gamma} \overline{\alpha}} \approx \frac{\beta + \overline{\beta}}{\alpha - \overline{\alpha}}$$

Theoretical estimate $A \approx 5 \cdot 10^{-4}$ and

$$0.1O(B) \approx O(B') \approx 10O(A) \approx 100O(D)$$
(6)

(5)

500

Experimental limits: $A_{\Lambda} = 0.013 \pm 0.021$ (PDG average), $A_{\Xi} = (0.0 \pm 5.1 \pm 4.4) \cdot 10^{-4}$ (HYCP, 158M events, p Cu, 800 GeV). No experimental limits on B_{Ξ} and B'_{Ξ}



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PANDA Detector





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Conclusions

New for PANDA

- Very large increase in statistics for Λ and Σ hyperons and possibility to measure at higher beam momenta
- Opens up new channels: Ξ , Ω , Λ_c , ...
- Increase in statistics for CP violation measurements and access to the unmeasured (more sensitive parameters) B_{Ξ} and B'_{Ξ}

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Hyperon Properties

	Hyperon	Quarks	Mass [MeV/ c ²]	cτ [cm]	Main decay	B [%]	α _Y
Antihyperon Hyperon Physics for the	٨	uds	1116	8.0	рπ-	64	+0.64
PANDA Experiment	Σ+	uus	1189	2.4	рπ ⁰	52	-0.98
E. Thomé	Σ ⁰	uds	1193	2.2 · 10 ⁻⁹	Λγ	100	-
Outline	Σ-	dds	1197	2.4	$n\pi^{-}$	100	-0.07
Introduction to Hyperons	Ξ0	uss	1315	8.7	$\Lambda \pi^0$	99	-0.41
Existing Experimental	Ξ-	dss	1321	4.9	$\Lambda\pi^{-}$	100	-0.46
Data	Ω^{-}	SSS	1672	2.5	ΛK ⁻	68	-0.03
Simulations for PANDA							
$ \begin{array}{l} \bar{p}p \rightarrow \bar{\Lambda} \\ \bar{p}p \rightarrow \bar{\Xi}^+ \Xi^- \end{array} $	Λ_{c}^{+}	udc	2286	$6.0 \cdot 10^{-3}$	$\Lambda \pi^+$	1	-0.91(15)
$\bar{p}p \rightarrow \bar{\Omega}^+ \Omega^-$	Σ_{c}^{++}	uuc	2454		$\Lambda_{c}^{+} \pi^{+}$	100	
$\begin{array}{l} pp \rightarrow \Lambda_c \ \Lambda_c^{+} \\ \textbf{CP Violation} \end{array}$	Σ_{c}^{+}	udc	2453		$\Lambda_{c}^{+}\pi^{0}$	100	
Precession of the Polarisation Vector in	Σ _c	ddc	2454		$\Lambda_{c}^{+}\pi^{-}$	100	
the Magnetic Field	Ξ+	usc	2468	1.2 · 10 ^{- 2}	$\Xi^- \pi^+ \pi^+$	seen	
Conclusions	Ξ0	dsc	2471	2.9 · 10 ^{- 3}	$\Xi^- \pi^+$	seen	-0.6(4)
	Ω_c^0	SSC	2697	1.9 · 10 ^{- 3}	$\Omega^{-}\pi^{+}$	seen	



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 \begin{array}{l} \bar{\mathbf{p}}\mathbf{p} \rightarrow \bar{\mathbf{A}}\mathbf{A} \\ \bar{p}\mathbf{p} \rightarrow \bar{\Xi}^+ \Xi^- \\ \bar{p}\mathbf{p} \rightarrow \bar{\Omega}^+ \Omega^- \\ \bar{p}\mathbf{p} \rightarrow \bar{\Lambda}_c^- \Lambda_c^+ \\ \mathbf{CP} \text{ Violation} \end{array}
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Precession of the Polarisation Vector in the Magnetic Field

Conclusions

 $\bar{p}p \rightarrow \bar{\Lambda}\Lambda$

- Reconstruction efficiency between 14 and 24%, depending on beam momentum
- Between 100 and 1000 reconstructable events/s
- Polarisation and spin correlations reconstructable
- \bullet Most severe background problem is to separate Λ from Σ^0

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Precession of the Polarisation Vector in the Magnetic Field

Conclusions

 $\bar{\mathrm{pp}} \to \bar{\Xi}^+ \Xi^-$



- Only existing data from old bubble chamber experiments
- $\Xi^-/\bar{\Xi}^+$ bent by magnetic field, must be corrected for
- Differential cross section not known, isotropic used for simulations
- \bullet Four displaced vertices \rightarrow low background
- $\bullet~$ Daughter Λ also decays weakly $\rightarrow~$ more parameters reconstructable



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Conclusions

$$\bar{\mathrm{o}}\mathrm{p} \to \bar{\Xi}^+ \Xi^-$$

• Simulations at 4 GeV/c beam momentum

- \bullet Reconstruction efficiency $\sim 17\% \rightarrow \sim 30$ reconstructable events/s
- Background from $\bar{\rm p}{\rm p}\to\bar{\Sigma}^-(1385)\Sigma^+(1385)$ investigated. Background/Signal smaller than $2\cdot10^{-4}$



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Reconstruction Efficiency Depending on Production and Decay Angles







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 $\bar{p}p \rightarrow \Omega^+ \Omega^ \bar{p}p \rightarrow \bar{\Lambda}^-_c \Lambda^+_c$ CP Violation

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$\Xi^{-}/\overline{\Xi}^{+}$ Spin Variables

Spin variables can be calculated from averages of the reconstructed decay angles (\hat{k}) . The acceptance depends on the angle in the hyperon decay \Rightarrow need to use MC based acceptance functions

$$P_{y} = \frac{3}{\alpha} \sum_{n=1}^{N} \frac{k_{y,n}}{A(k_{y,n})}, \quad C_{ij} = \frac{9}{\alpha \bar{\alpha}} \sum_{n=1}^{N} \frac{\bar{k}_{i,n} k_{j,n}}{A(\bar{k}_{i,n}, k_{j,n})}$$
(7)

Reconstructed spin variables from 350000 reconstructed events. $(P = \sin 2\Theta, C_{ij} = \sin \Theta \text{ as input})$



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Conclusions

$\bar{\rm p}{\rm p}\to\bar{\Omega}^+\Omega^-$

- Differential cross section not known, isotropic used for simulations
- Cross section estimate ~ 2 nb, branching ratio for $\Omega \to \Lambda {\rm K}$ $\sim 68\%$
- Simulations at 12 GeV/c beam momentum
- Reconstruction efficiency $\sim 30\% \rightarrow \sim 80$ reconstructable events/hour
- Acceptance for all production angles of $\bar{\Omega}^+$





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Density Matrix of $\boldsymbol{\Omega}$

spin 3/2 particle \rightarrow 15 polarisation parameters *r* instead of only the vector polarisation as for spin 1/2 Parity invariance \rightarrow 8 are equal to zero

$$\rho = \frac{1}{4} \begin{bmatrix} 1 + \sqrt{3}r_0^2 & i\frac{3}{\sqrt{5}}r_{-1}^1 - \sqrt{3}r_1^2 & \sqrt{3}r_2^2 - i\sqrt{3}r_{-2}^3 & -i\sqrt{6}r_{-3}^3 \\ -i\frac{3}{\sqrt{5}}r_{-1}^1 - \sqrt{3}r_1^2 & 1 - \sqrt{3}r_0^2 & i2\sqrt{\frac{3}{5}}r_{-1}^1 + i3\sqrt{\frac{2}{5}}r_{-1}^3 & \sqrt{3}r_2^2 + i\sqrt{3}r_{-2}^3 \\ \sqrt{3}r_2^2 + i\sqrt{3}r_{-2}^3 & -i2\sqrt{\frac{3}{5}}r_{-1}^1 - i3\sqrt{\frac{2}{5}}r_{-1}^3 & 1 - \sqrt{3}r_0^2 & i\frac{3}{\sqrt{5}}r_{-1}^1 + \sqrt{3}r_1^2 \\ i\sqrt{6}r_{-3}^3 & \sqrt{3}r_2^2 - i\sqrt{3}r_{-2}^3 & -i\frac{3}{\sqrt{5}}r_{-1}^1 + \sqrt{3}r_1^2 & 1 + \sqrt{3}r_0^2 \end{bmatrix}$$
(8)

Degree of polarisation = $\sqrt{\sum r^2}$ Compare spin 1/2

$$\rho = \frac{1}{2} (\mathcal{I} + \tilde{P} \cdot \tilde{\sigma}) = \frac{1}{2} \begin{bmatrix} 1 + P_z & P_X + iP_y \\ P_X - iP_y & 1 - P_z \end{bmatrix} \rightarrow \frac{1}{2} \begin{bmatrix} 1 & iP_y \\ -iP_y & 1 \end{bmatrix}$$
(9)

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Polarisation of $\boldsymbol{\Omega}$

Angular distribution of $\Omega \rightarrow \Lambda K$ ($\alpha \approx 0$)

$$(\Theta, \phi) = \frac{1}{4\pi} \left(1 - \frac{\sqrt{3}}{2} (3\cos^2 \Theta - 1)r_0^2 - \frac{3}{2} (\sin^2 \Theta \cos 2\phi r_1^2 - \sin 2\Theta \cos \phi r_2^2) \right)$$
(10)

Angular distribution of $\Lambda \to \mathrm{p}\pi$

$$I(\Theta,\phi) = \frac{1}{4\pi} \left(1 + \alpha_{\Lambda} \sin \Theta (16\sqrt{15}r_{-1}^{1} + 15\sqrt{10}r_{-1}^{3}) (\beta_{1}\cos\phi + \gamma_{1}\sin\phi) \right)$$
(11)



Reconstructed r_0^2 from 30000 events. $r_0^2 = \frac{1}{\sqrt{3}} \sin 2\Theta$ as input

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Conclusions

 $\bar{p}p \rightarrow \bar{\Lambda}_c^- \Lambda_c^+$

- \bullet Treshold momentum 10.2 GeV/c, simulations at 12 GeV/c
- Differential cross section not known, isotropic used for simulations
- Acceptance for all production angles of $\bar{\Lambda}_c^-$
- Cross section \sim 0.1 μ b, branching ratio for $\Lambda_c
 ightarrow \mathrm{p}\pi \sim 1\%$
- Reconstruction efficiency $\sim 35\% \rightarrow \sim 25$ reconstructable events/day
- \bullet Background from Σ resonances, not investigated yet
- Large asymmetry parameter $\alpha = -0.91$, polarisation reasonably well reconstructed from a few thousand events (3500)



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 $\bar{p}p \rightarrow \Lambda\Lambda$ $\bar{p}p \rightarrow \bar{\Xi}^+\Xi^-$ $\bar{p}p \rightarrow \bar{\Omega}^+\Omega^-$ $\bar{p}p \rightarrow \bar{\Lambda}_c^-\Lambda_c^+$ CD Violation

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Other Charmed Hyperons

- Ξ_c^0 and Ω_c^0 in principle also accessible
- $\bullet\,$ Treshold momentum 12.0 and 14.6 GeV/c
- Assumption: Cross sections reduced by factors $\alpha_s^2=1/25$ and $\alpha_s^4=1/625$

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- Branching ratios not known
- Ξ_c^0 maybe possible, Ω_c^0 probably not



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CP Violation

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CP Violation

CP violation parmeters can be measured by detecting if the decay particles go up or down with respect to the hyperon production plane

$$p = \frac{N_{+-} - N_{-+}}{N} = \frac{\beta + \bar{\beta}}{4}$$
(12)

Experimental considerations:

- Uncertainty $\propto P \sqrt{rac{\mathrm{d}\sigma}{\mathrm{d}\Omega}}$, choose right angular interval
- Only use events where the hyperons decay in the vacuum of the beam pipe to avoid depolarisation in material
- Use events close to beam axis to make acceptance symmetric





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$$\frac{\mathrm{d}\bar{\Sigma}}{\mathrm{d}\tau} = -\left[\frac{1}{\gamma+1}\left(\bar{u}\times\frac{\mathrm{d}\bar{u}}{\mathrm{d}\tau}\right) + \frac{ge}{2m}\bar{B}_{(0)}\right]\times\bar{\Sigma}$$
(13)

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Uncharged hyperon: construct the coordinate system \hat{B}_0 , $\hat{B}_0 \times \hat{\Sigma}$, $(\hat{B}_0 \times \hat{\Sigma}) \times \hat{B}_0$.



The precession angle of the Λ polarisation vector at 4 GeV/c.



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Effect on the Reconstruction of Polarisation



The difference in calculated polarisation using the Λ rest system at production and a rest system which has been rotated in the same way as the polarisation vector. Spin correlations are not affected.



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- High statistics for the lighter hyperons, low background in general
- Many observables for Ξ⁻ and Ω⁻ can be measured for the first time: differential cross section, polarisation, spin correlations, CP violation parameters, ...
- Charmed hyperons: Λ_c^+ is possible to study, maybe Ξ_c^0 and probably not Ω_c^0
- A small effect from the strong magnetic field affecting the hyperon spins, for the uncharged hyperons this can be corrected for

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