

Causal Charge Diffusion in Relativistic Heavy-Ion Collisions

Christopher J. Plumberg
in collaboration with Joseph Kapusta

University of Minnesota

XQCD 2018



XQCD 2018, Frankfurt, DE
May 22, 2018



Quantifying the strength of Diffusion:



*Quantifying the
strength of Diffusion:*



*Quantifying the
strength of Diffusion:*



*Quantifying the
strength of Diffusion:*



Quantifying the strength of Diffusion:



Cups of coffee:



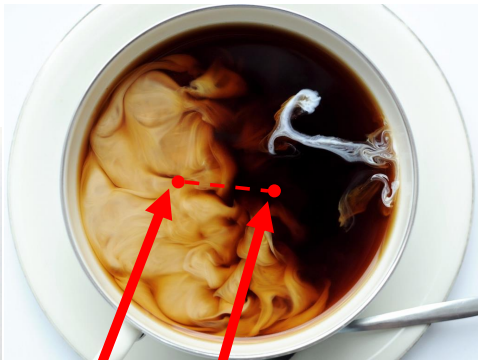
$$\langle \delta n_{\text{milk}}(x) \delta n_{\text{coffee}}(x') \rangle_{\text{lots of cups}}$$

Quantifying the strength of Diffusion:



Cups of coffee:

Heavy-ion collisions:



$$\langle \delta n_{\text{milk}}(x) \delta n_{\text{coffee}}(x') \rangle_{\text{lots of cups}}$$

$$\langle \delta n_{Q_a}(x) \delta n_{Q_b}(x') \rangle_{\text{lots of collisions}}$$

Quantifying the strength of Diffusion:

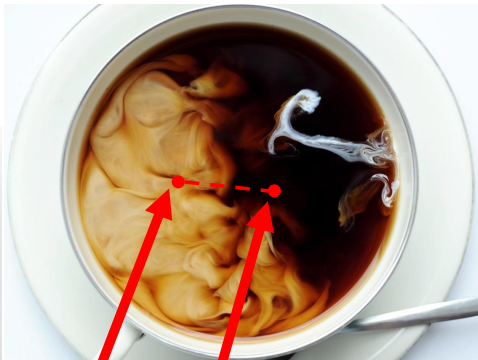


Cups of coffee:

$$\langle \delta n_{\text{milk}}(x) \delta n_{\text{coffee}}(x') \rangle_{\text{lots of cups}}$$

Heavy-ion collisions:

$$\langle \delta n_{Q_a}(x) \delta n_{Q_b}(x') \rangle_{\text{lots of collisions}}$$



$$B_{h+h-}(y_1 - y_2) \equiv \left\langle \frac{dN}{dy} \right\rangle_{\text{ev}}^{-1} \left\langle \delta \left(\frac{dN}{dy_1} \right) \delta \left(\frac{dN}{dy_2} \right) \right\rangle_{\text{ev}}$$

Charge Diffusion in Fluctuating Hydrodynamics

- Identify charge of interest: Q (electric charge)
- Construct 4-current J_Q^μ :

$$J_Q^\mu = \underbrace{n_Q u^\mu}_{\text{ideal}} + \underbrace{\Delta J_Q^\mu}_{\text{viscous}} + \underbrace{I^\mu}_{\text{stochastic}}$$

- Require J_Q to be conserved: $\partial_\mu J_Q^\mu = 0$ (EOM)
- Self-consistency of hydro $\implies \langle I^\mu I^\nu \rangle \propto \sigma_Q$ (FDT)
- Use solution to EOM + FDT to construct interesting quantities, e.g., $\langle \delta n_Q(x) \delta n_Q(x') \rangle$

$$\Delta J_Q^\mu = \sigma_Q T \Delta^\mu \left(\frac{\mu_Q}{T} \right),$$
$$I^\mu \equiv s(\tau) f(\xi, \tau) \left(\sinh \xi, \vec{0}_\perp, \cosh \xi \right)$$

with

$$\tau \equiv \sqrt{t^2 - z^2} \text{ and } \xi = \frac{1}{2} \ln \left| \frac{t+z}{t-z} \right|.$$

Steps to solve $\partial_\mu J_Q^\mu = 0$:

- Conservation of J_Q^μ yields Langevin-type equation of motion for $\delta\tilde{n}(k, \tau)$:

$$\partial_\mu J_Q^\mu = 0 \quad \Longrightarrow \quad \frac{\partial}{\partial \tau}(\tau \delta\tilde{n}) + \frac{D_Q k^2}{\tau^2}(\tau \delta\tilde{n}) = -iks \tilde{f}$$

Steps to solve $\partial_\mu J_Q^\mu = 0$:

- Conservation of J_Q^μ yields Langevin-type equation of motion for $\delta\tilde{n}(k, \tau)$:

$$\partial_\mu J_Q^\mu = 0 \quad \Longrightarrow \quad \frac{\partial}{\partial \tau}(\tau \delta\tilde{n}) + \frac{D_Q k^2}{\tau^2}(\tau \delta\tilde{n}) = -iks \tilde{f}$$

- Homogeneous equation (with $\tilde{f} = 0$) describes relaxation of single fluctuation mode $\delta\tilde{n}(k, \tau)$
→ solution defines the *response function* $\tilde{G}(k; \tau, \tau')$

Steps to solve $\partial_\mu J_Q^\mu = 0$:

- Conservation of J_Q^μ yields Langevin-type equation of motion for $\delta\tilde{n}(k, \tau)$:

$$\partial_\mu J_Q^\mu = 0 \quad \Longrightarrow \quad \frac{\partial}{\partial \tau}(\tau \delta\tilde{n}) + \frac{D_Q k^2}{\tau^2}(\tau \delta\tilde{n}) = -iks \tilde{f}$$

- Homogeneous equation (with $\tilde{f} = 0$) describes relaxation of single fluctuation mode $\delta\tilde{n}(k, \tau)$
→ solution defines the *response function* $\tilde{G}(k; \tau, \tau')$
- Solve inhomogeneous equation (with $\tilde{f} \neq 0$) with:

$$\delta\tilde{n}(k, \tau) = -\frac{1}{\tau} \int_{\tau_0}^{\tau} d\tau' s(\tau') \tilde{G}(k; \tau, \tau') \tilde{f}(k, \tau') .$$

Use 2-point functions to study fluctuations systematically:

$$\begin{aligned}\langle \delta \tilde{n}(k_1, \tau_1) \delta \tilde{n}(k_2, \tau_2) \rangle &= \frac{1}{\tau_1 \tau_2} \int_{\tau_0}^{\tau_1} d\tau'_1 s(\tau'_1) \int_{\tau_0}^{\tau_2} d\tau'_2 s(\tau'_2) \\ &\times \tilde{G}(k_1; \tau_1, \tau'_1) \tilde{G}(k_2; \tau_2, \tau'_2) \\ &\times \langle \tilde{f}(k_1, \tau'_1) \tilde{f}(k_2, \tau'_2) \rangle\end{aligned}$$

Use 2-point functions to study fluctuations systematically:

$$\begin{aligned}\langle \delta \tilde{n}(k_1, \tau_1) \delta \tilde{n}(k_2, \tau_2) \rangle &= \frac{1}{\tau_1 \tau_2} \int_{\tau_0}^{\tau_1} d\tau'_1 s(\tau'_1) \int_{\tau_0}^{\tau_2} d\tau'_2 s(\tau'_2) \\ &\times \tilde{G}(k_1; \tau_1, \tau'_1) \tilde{G}(k_2; \tau_2, \tau'_2) \\ &\times \langle \tilde{f}(k_1, \tau'_1) \tilde{f}(k_2, \tau'_2) \rangle\end{aligned}$$

Problem:

$$\langle \tilde{f}(k_1, \tau'_1) \tilde{f}(k_2, \tau'_2) \rangle = ?$$

Interlude:

White noise vs. **Colored** noise

Noise in a static medium

- White noise: $\langle f(x_1)f(x_2) \rangle = N\delta^4(x_1 - x_2)$
 - Pro: simple to implement
 - Con: acausal signal propagation ($v_Q^2 \rightarrow \infty$)

¹Aarts *et al.*, JHEP **1502**, 186 (2015)

Noise in a static medium

- White noise: $\langle f(x_1)f(x_2) \rangle = N\delta^4(x_1 - x_2)$
 - Pro: simple to implement
 - Con: acausal signal propagation ($v_Q^2 \rightarrow \infty$)
- Colored noise: $\langle f(x_1)f(x_2) \rangle = \frac{N}{2\tau_Q} e^{-|t_1-t_2|/\tau_Q} \delta^3(\vec{x}_1 - \vec{x}_2)$
 - Pro: causal ($v_Q^2 = D_Q/\tau_Q$)

$$D_Q \sim \tau_Q \sim \frac{1}{2\pi T} \implies v_Q^2 \sim 1^1$$

$$\Delta J^\mu \rightarrow \sigma_Q T \Delta^\mu (1 + \tau_Q u \cdot \partial)^{-1} \left(\frac{\mu_Q}{T} \right)$$

- Con: more complicated implementation

¹Aarts *et al.*, JHEP **1502**, 186 (2015)

Noise in a static medium

- White noise: $\langle f(x_1)f(x_2) \rangle = N\delta^4(x_1 - x_2)$
 - Pro: simple to implement
 - Con: acausal signal propagation ($v_Q^2 \rightarrow \infty$)
- Colored noise: $\langle f(x_1)f(x_2) \rangle = \frac{N}{2\tau_Q} e^{-|t_1-t_2|/\tau_Q} \delta^3(\vec{x}_1 - \vec{x}_2)$
 - Pro: causal ($v_Q^2 = D_Q/\tau_Q$)

$$D_Q \sim \tau_Q \sim \frac{1}{2\pi T} \implies v_Q^2 \sim 1^1$$

$$\Delta J^\mu \rightarrow \sigma_Q T \Delta^\mu (1 + \tau_Q u \cdot \partial)^{-1} \left(\frac{\mu_Q}{T} \right)$$

- Con: more complicated implementation

Complications:

- Colored two-point function is *non-local* on scales $\sim \mathcal{O}(\tau_Q)$
- $\partial \cdot u \sim \frac{1}{\tau}$, $\tau \sim \mathcal{O}(1 \text{ fm}/c)$
- Static approach breaks down at $\tau \sim \tau_Q$!

¹Aarts *et al.*, JHEP **1502**, 186 (2015)

Noise in a dynamical medium

Solution: rewrite two-point functions *locally*

■ White noise:

$$\langle \tilde{f}(k_1, \tau_1) \tilde{f}(k_2, \tau_2) \rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2)$$

■ Colored noise:

$$\begin{aligned} & \left\langle (1 + \tau_Q \partial / \partial \tau_1) \tilde{f}(k_1, \tau_1) \right. \\ & \times \left. (1 + \tau_Q \partial / \partial \tau_2) \tilde{f}(k_2, \tau_2) \right\rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2) \end{aligned}$$

Noise in a dynamical medium

Solution: rewrite two-point functions *locally*

■ White noise:

$$\langle \tilde{f}(k_1, \tau_1) \tilde{f}(k_2, \tau_2) \rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2)$$

■ Colored noise:

$$\begin{aligned} & \left\langle (1 + \tau_Q \partial / \partial \tau_1) \tilde{f}(k_1, \tau_1) \right. \\ & \left. \times (1 + \tau_Q \partial / \partial \tau_2) \tilde{f}(k_2, \tau_2) \right\rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2) \end{aligned}$$

Checklist:

✓ Solve $\partial_\mu J_Q^\mu = 0$ with $\tau_Q \neq 0$ for $\delta \tilde{n}(k, \tau)$

Noise in a dynamical medium

Solution: rewrite two-point functions *locally*

■ White noise:

$$\langle \tilde{f}(k_1, \tau_1) \tilde{f}(k_2, \tau_2) \rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2)$$

■ Colored noise:

$$\begin{aligned} & \left\langle (1 + \tau_Q \partial / \partial \tau_1) \tilde{f}(k_1, \tau_1) \right. \\ & \left. \times (1 + \tau_Q \partial / \partial \tau_2) \tilde{f}(k_2, \tau_2) \right\rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2) \end{aligned}$$

Checklist:

- ✓ Solve $\partial_\mu J_Q^\mu = 0$ with $\tau_Q \neq 0$ for $\delta \tilde{n}(k, \tau)$
- ✓ Fix $\langle \tilde{f}(k_1, \tau_1) \tilde{f}(k_2, \tau_2) \rangle$

Noise in a dynamical medium

Solution: rewrite two-point functions *locally*

■ White noise:

$$\langle \tilde{f}(k_1, \tau_1) \tilde{f}(k_2, \tau_2) \rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2)$$

■ Colored noise:

$$\begin{aligned} & \left\langle (1 + \tau_Q \partial / \partial \tau_1) \tilde{f}(k_1, \tau_1) \right. \\ & \left. \times (1 + \tau_Q \partial / \partial \tau_2) \tilde{f}(k_2, \tau_2) \right\rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2) \end{aligned}$$

Checklist:

- ✓ Solve $\partial_\mu J_Q^\mu = 0$ with $\tau_Q \neq 0$ for $\delta \tilde{n}(k, \tau)$
- ✓ Fix $\langle \tilde{f}(k_1, \tau_1) \tilde{f}(k_2, \tau_2) \rangle$
- ✓ Get $\langle \delta \tilde{n}(k_1, \tau_1) \delta \tilde{n}(k_2, \tau_2) \rangle$

Noise in a dynamical medium

Solution: rewrite two-point functions *locally*

■ White noise:

$$\langle \tilde{f}(k_1, \tau_1) \tilde{f}(k_2, \tau_2) \rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2)$$

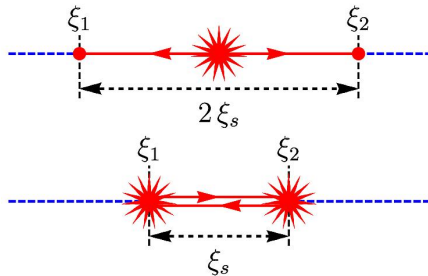
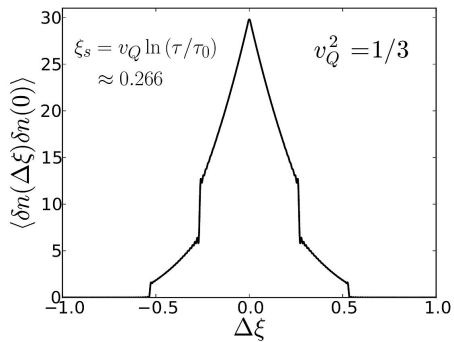
■ Colored noise:

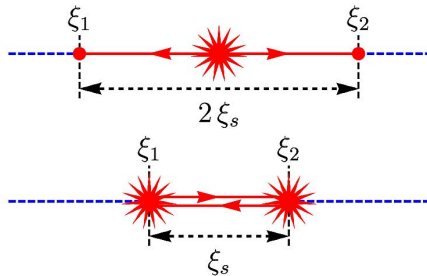
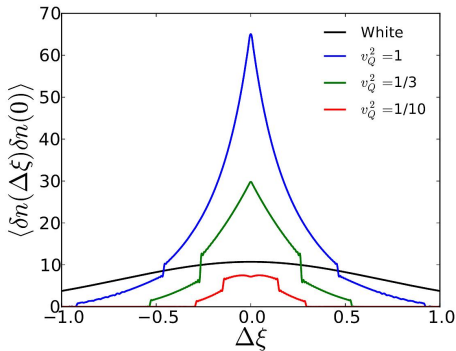
$$\begin{aligned} & \left\langle (1 + \tau_Q \partial / \partial \tau_1) \tilde{f}(k_1, \tau_1) \right. \\ & \left. \times (1 + \tau_Q \partial / \partial \tau_2) \tilde{f}(k_2, \tau_2) \right\rangle \propto \delta(k_1 + k_2) \delta(\tau_1 - \tau_2) \end{aligned}$$

Checklist:

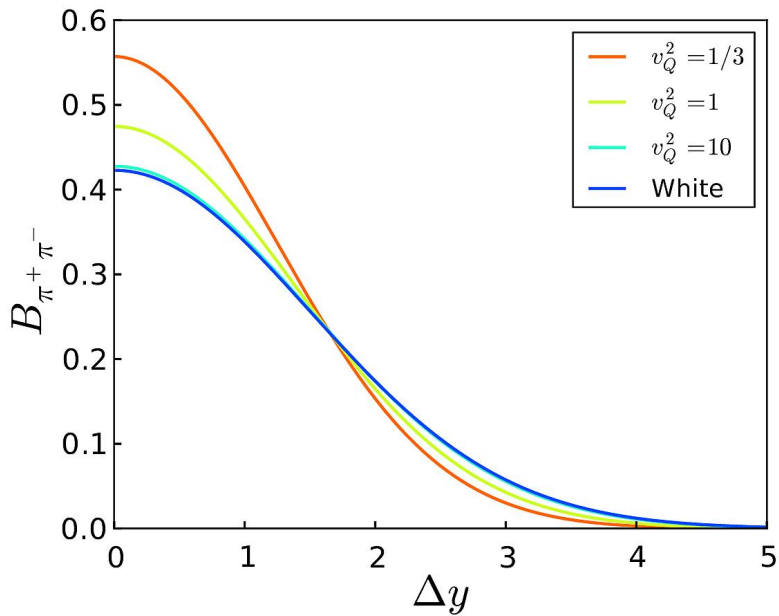
- ✓ Solve $\partial_\mu J_Q^\mu = 0$ with $\tau_Q \neq 0$ for $\delta \tilde{n}(k, \tau)$
- ✓ Fix $\langle \tilde{f}(k_1, \tau_1) \tilde{f}(k_2, \tau_2) \rangle$
- ✓ Get $\langle \delta \tilde{n}(k_1, \tau_1) \delta \tilde{n}(k_2, \tau_2) \rangle$
- ✓ Get $B_{h+h-}(y_1 - y_2)$

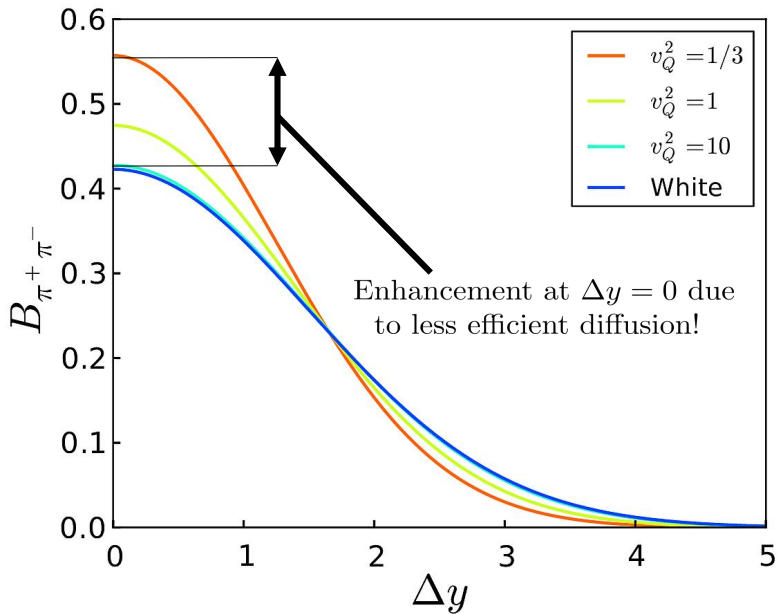
Results





- Both sets of wavefronts reflect essential aspects of causal signal propagation
- Wavefronts travel farther with larger v_Q^2
- *No* wavefronts for $v_Q^2 \rightarrow \infty$!





Summary:

- Diffusion is an essential aspect of heavy-ion collisions
- Hydrodynamic fluctuations offer a natural framework for modeling diffusion
- White noise leads to violations of relativistic causality
- Colored noise consistently imposes causality on diffusive mechanisms

Summary:

- Diffusion is an essential aspect of heavy-ion collisions
- Hydrodynamic fluctuations offer a natural framework for modeling diffusion
- White noise leads to violations of relativistic causality
- Colored noise consistently imposes causality on diffusive mechanisms

Future/ongoing work:

- More sophisticated models of colored noise (e.g., correlations in space and time)
- Match onto holographic dispersion relations
- Incorporate into 3+1D hydrodynamic simulations
- Subtract/project out self-correlations
- ...

Summary:

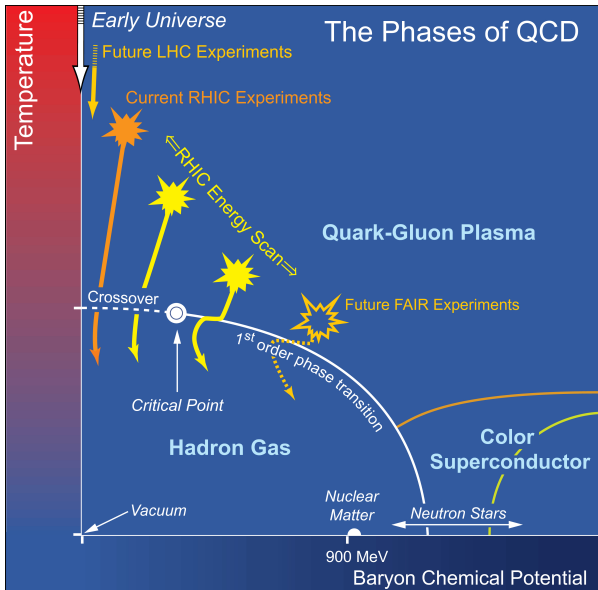
- Diffusion is an essential aspect of heavy-ion collisions
- Hydrodynamic fluctuations offer a natural framework for modeling diffusion
- White noise leads to violations of relativistic causality
- Colored noise consistently imposes causality on diffusive mechanisms

Future/ongoing work:

- More sophisticated models of colored noise (e.g., correlations in space and time)
- Match onto holographic dispersion relations
- Incorporate into 3+1D hydrodynamic simulations
- Subtract/project out self-correlations
- ...

Thanks for your attention!

Backup slides



Steps to solve $\partial_\mu J_Q^\mu = 0$ (with $\tau_Q \neq 0$):

$$\partial_\mu J_Q^\mu = 0 \quad \Longrightarrow$$

$$\begin{aligned} & \frac{\partial^2}{\partial \tau^2}(\tau \delta \tilde{n}) + \left[\frac{1}{\tau_Q} - \frac{\partial}{\partial \tau} \ln \left(\frac{\chi_Q T D_Q}{\tau} \right) \right] \frac{\partial}{\partial \tau}(\tau \delta \tilde{n}) + \frac{D_Q k^2}{\tau_Q \tau^2}(\tau \delta \tilde{n}) \\ &= -iks \left[\frac{\partial \tilde{f}}{\partial \tau} + \left(\frac{1}{\tau_Q} - \frac{1}{\tau} - \frac{\partial}{\partial \tau} \ln \left(\frac{\chi_Q T D_Q}{\tau} \right) \right) \tilde{f} \right]. \end{aligned}$$

T is the temperature, D_Q is the electric charge diffusion coefficient, σ_Q is the electric charge conductivity, and $\chi_Q = \sigma_Q/D_Q$ is the electric charge susceptibility. Finally, the quantity k is Fourier-conjugate to the spatial rapidity ξ ; for any quantity X , we define

$$X(\xi, \tau) \equiv \int_{-\infty}^{\infty} \frac{dk}{2\pi} e^{ik\xi} \tilde{X}(k, \tau)$$

Subtracting self-correlations

White noise density correlator:

$$\langle \delta n(\xi_1, \tau_f) \delta n(\xi_2, \tau_f) \rangle = \frac{\chi_Q(\tau_f) T_f}{A \tau_f} \left[\delta(\xi_1 - \xi_2) - \frac{1}{\sqrt{\pi w^2}} e^{-(\xi_1 - \xi_2)^2 / w^2} \right]$$

- First term: “self-correlations”

- Represent trivial correlations of a particle with itself
- Not measured experimentally

- Second term: diffusive correlations

- Represent physical, non-trivial correlations of distinct particles
- Are actually what we care about

→ Self-correlations need to be subtracted out to compare with experiment!

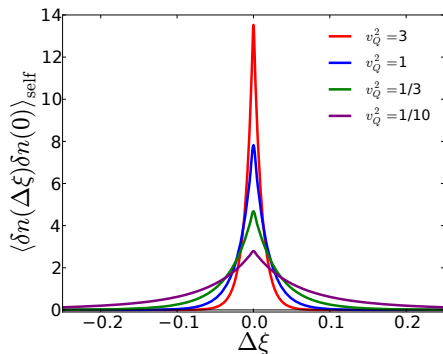
→ Not so hard to do for white noise...

→ ...but highly non-trivial for colored noise!

Subtracting self-correlations

Colored noise density self-correlations ($v_Q \gg 1$):

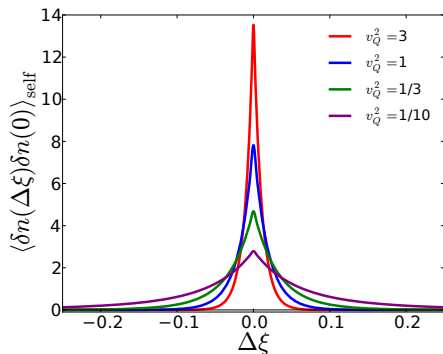
$$\langle \delta n(\xi_1, \tau_f) \delta n(\xi_2, \tau_f) \rangle_{\text{self}} \approx \frac{\chi_Q(\tau_f) T_f}{A \tau_f} \frac{v_Q \tau_f}{2 D_Q} \exp\left(-\frac{v_Q \tau_f}{D_Q} |\xi_1 - \xi_2|\right)$$



Subtracting self-correlations

Colored noise density self-correlations ($v_Q \gg 1$):

$$\langle \delta n(\xi_1, \tau_f) \delta n(\xi_2, \tau_f) \rangle_{\text{self}} \approx \frac{\chi_Q(\tau_f) T_f}{A \tau_f} \frac{v_Q \tau_f}{2 D_Q} \exp\left(-\frac{v_Q \tau_f}{D_Q} |\xi_1 - \xi_2|\right)$$

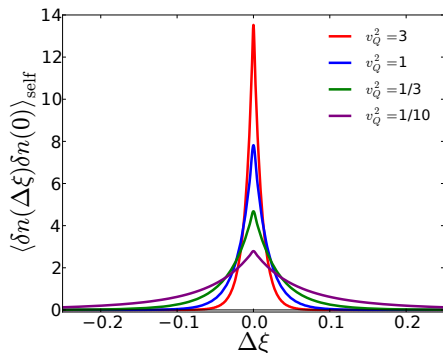


- “Adiabatic limit” ($v_Q \gg 1$)
reduces to exponential form on
quasi-static background

Subtracting self-correlations

Colored noise density self-correlations ($v_Q \gg 1$):

$$\langle \delta n(\xi_1, \tau_f) \delta n(\xi_2, \tau_f) \rangle_{\text{self}} \approx \frac{\chi_Q(\tau_f) T_f}{A \tau_f} \frac{v_Q \tau_f}{2 D_Q} \exp\left(-\frac{v_Q \tau_f}{D_Q} |\xi_1 - \xi_2|\right)$$



- “Adiabatic limit” ($v_Q \gg 1$)
reduces to exponential form on
quasi-static background
- “Instantaneous limit” ($v_Q \ll 1$)
just takes all correlations to
zero

See these references for more detail:

- Ling, Springer, and Stephanov [PRC **89**, 064901 (2014)]
- Kapusta and CP [PRC **97**, 014906 (2018)]

Holographic considerations

Key idea: matching colored noise to holographic dispersion relations yields estimates for D_Q , τ_Q

Gurtin-Pipkin noise:

$$\begin{aligned}\frac{\partial}{\partial t} - D_Q \nabla^2 + \tau_1 \frac{\partial^2}{\partial t^2} + \tau_2^2 \frac{\partial^3}{\partial t^3} - \tau_3 D_Q \frac{\partial}{\partial t} \nabla^2 &= 0 \\ \Rightarrow \tau_2^2 \omega^3 + i\tau_1 \omega^2 - (1 + \tau_3 D_Q k^2) \omega - iD_Q k^2 &= 0\end{aligned}$$

Holography^{2,3} yields Kaluza-Klein-type tower of poles in holographic dispersion relation:

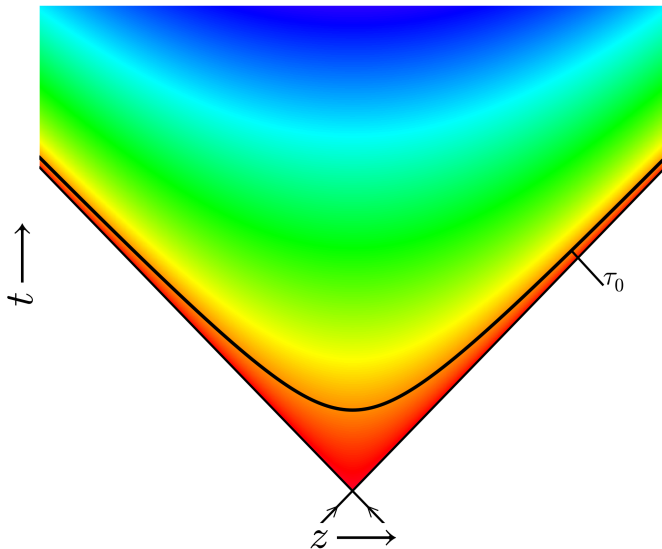
$$\omega(k=0) = (\pm n - in)2\pi T, \quad n = 0, 1, 2, \dots$$

Match GP noise onto three lowest frequency poles

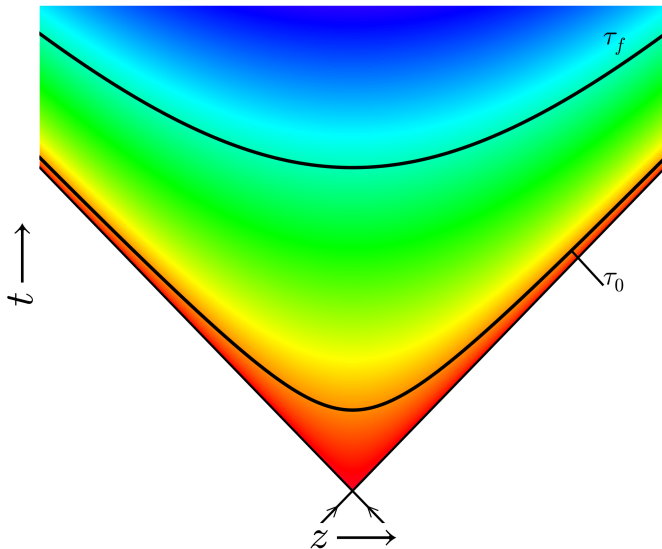
$$\Rightarrow D_Q = \tau_Q = \tau_1 = \frac{1}{2\pi T}, \quad \tau_2 = \tau_1/\sqrt{2}, \quad \tau_3 = \tau_1/2$$

²Nunez and Starinets [PRD **67**, 124013 (2003)]

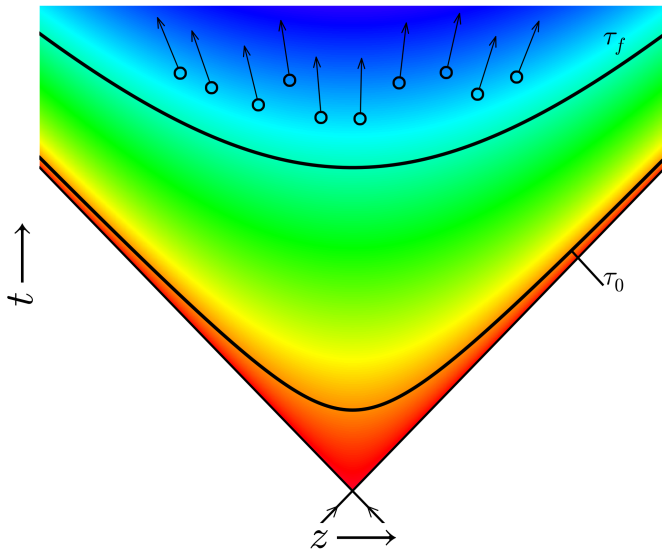
³PolICASTRO, Son and Starinets [JHEP **0209**, 043 (2002)]



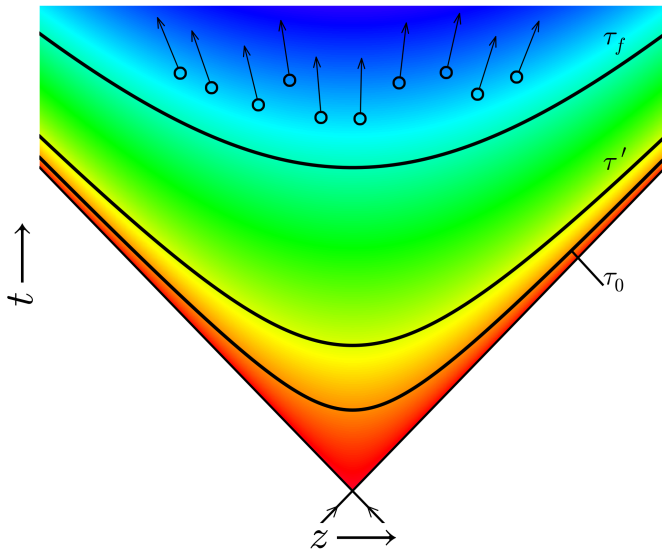
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} N(\tau') \tilde{G}_X(k_1; \tau_f, \tau') \tilde{G}_Y(k_2; \tau_f, \tau')$$



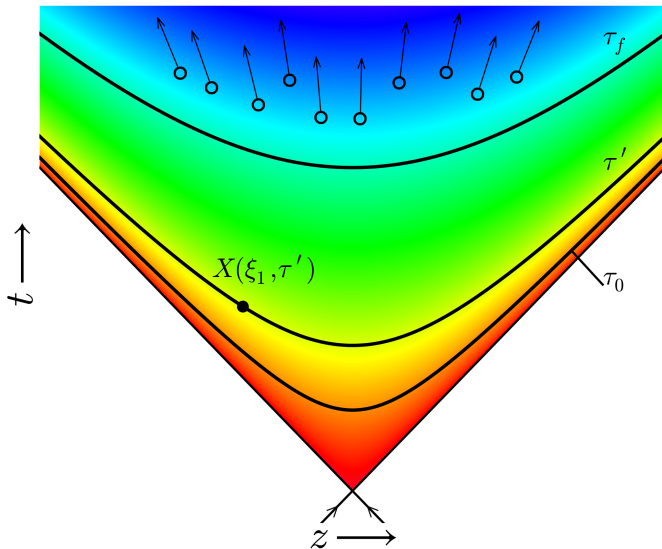
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} N(\tau') \tilde{G}_X(k_1; \tau_f, \tau') \tilde{G}_Y(k_2; \tau_f, \tau')$$



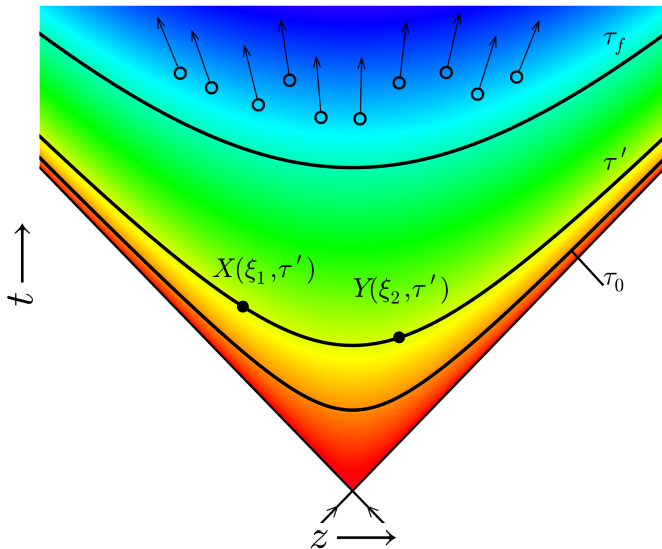
$$\langle \tilde{\mathbf{X}}(\mathbf{k}_1, \tau_f) \tilde{\mathbf{Y}}(\mathbf{k}_2, \tau_f) \rangle = 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} N(\tau') \tilde{G}_X(k_1; \tau_f, \tau') \tilde{G}_Y(k_2; \tau_f, \tau')$$



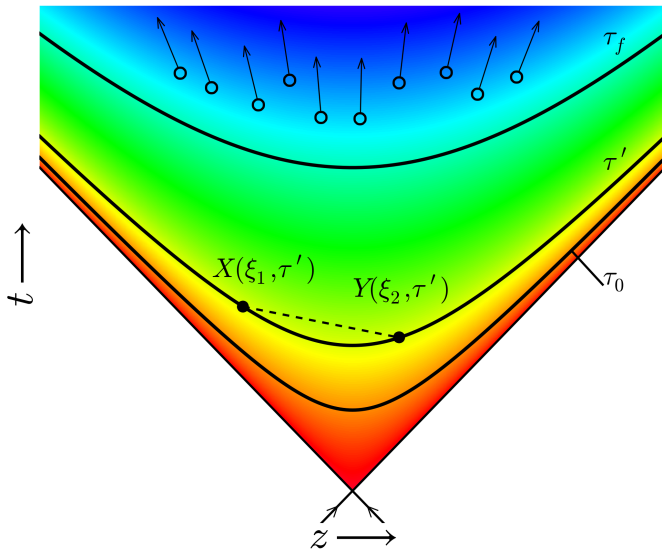
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} N(\tau') \tilde{G}_X(k_1; \tau_f, \tau') \tilde{G}_Y(k_2; \tau_f, \tau')$$



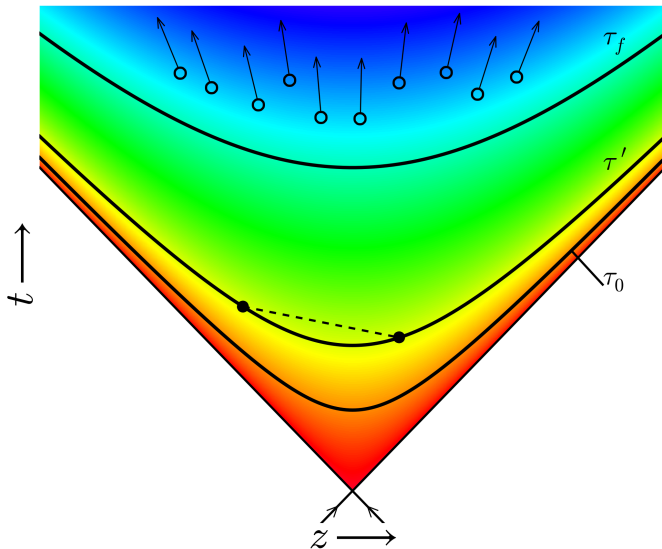
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} N(\tau') \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \tilde{G}_Y(k_2; \tau_f, \tau')$$



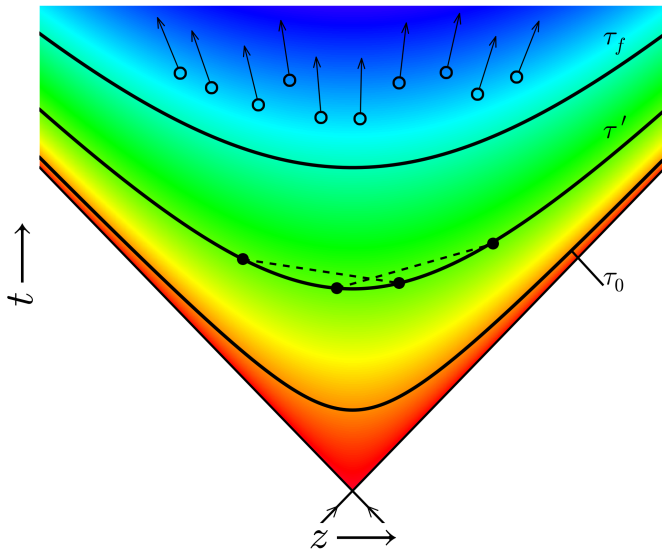
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} N(\tau') \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau')$$



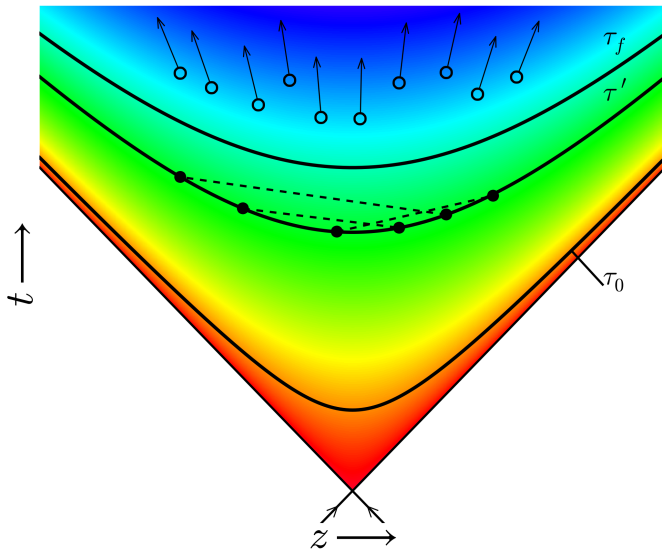
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \mathbf{N}(\tau') \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau')$$



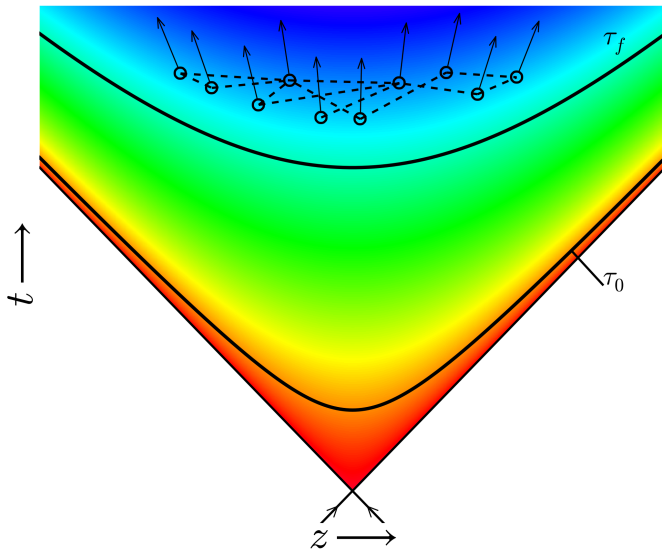
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \mathbf{N}(\tau') \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau')$$



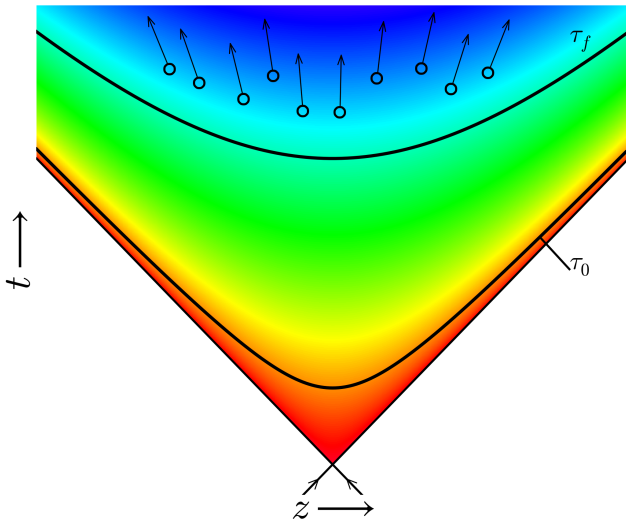
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \mathbf{N}(\tau') \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau')$$



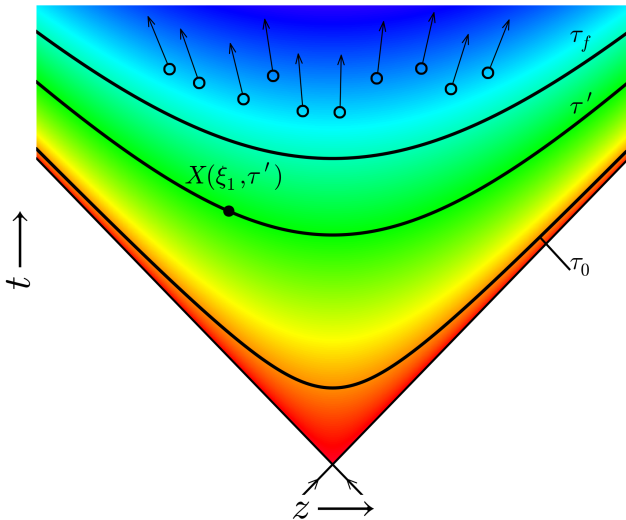
$$\langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle = 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \mathbf{N}(\tau') \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau')$$



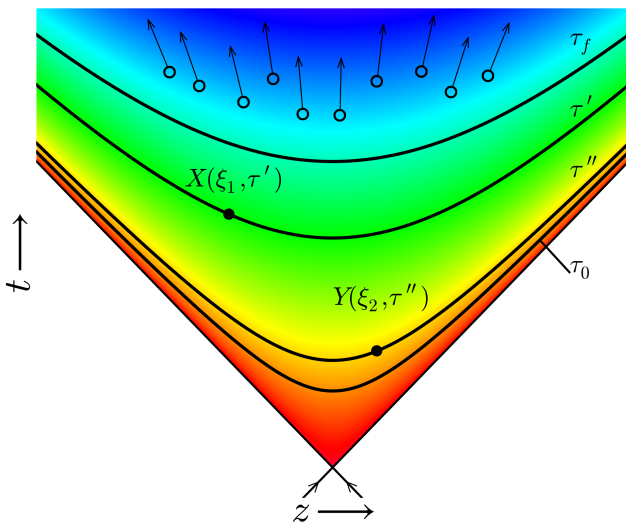
$$\langle \tilde{\mathbf{X}}(\mathbf{k}_1, \tau_f) \tilde{\mathbf{Y}}(\mathbf{k}_2, \tau_f) \rangle = 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \mathbf{N}(\tau') \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau')$$



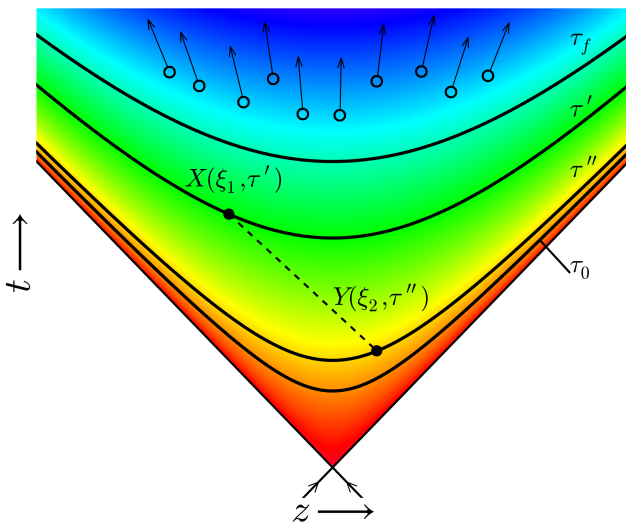
$$\begin{aligned}
 \langle \tilde{\mathbf{X}}(\mathbf{k}_1, \tau_f) \tilde{\mathbf{Y}}(\mathbf{k}_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{G}_X(k_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{G}_Y(k_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



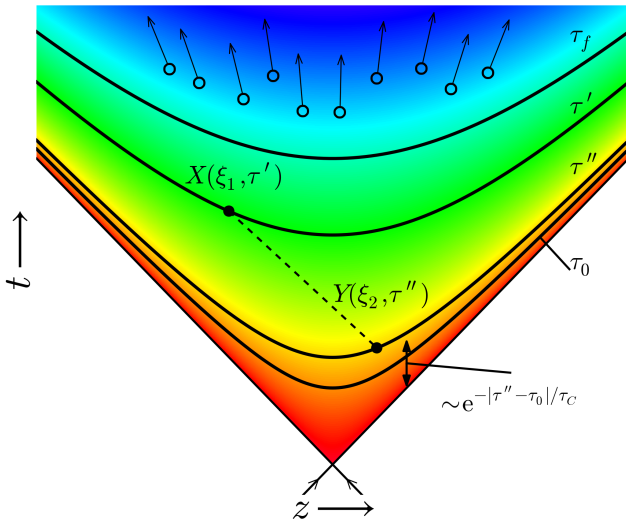
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{G}_Y(k_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



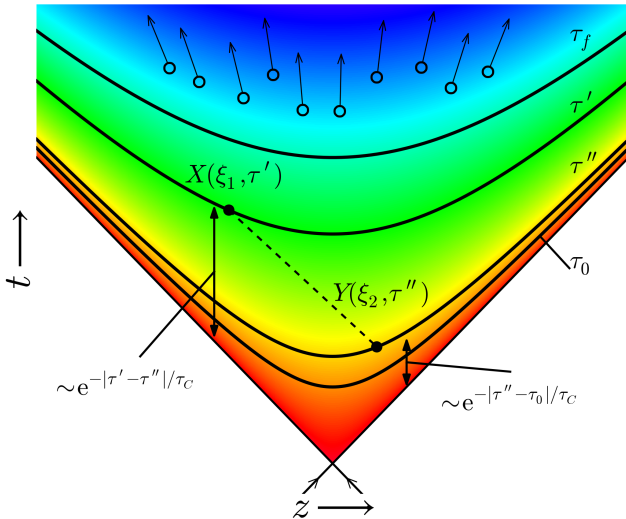
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



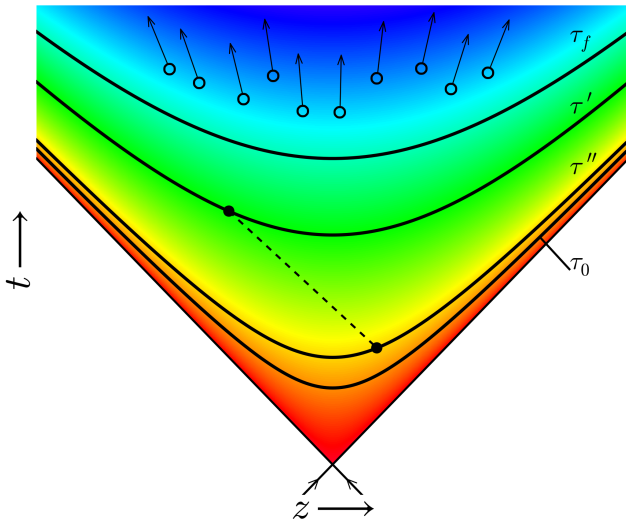
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



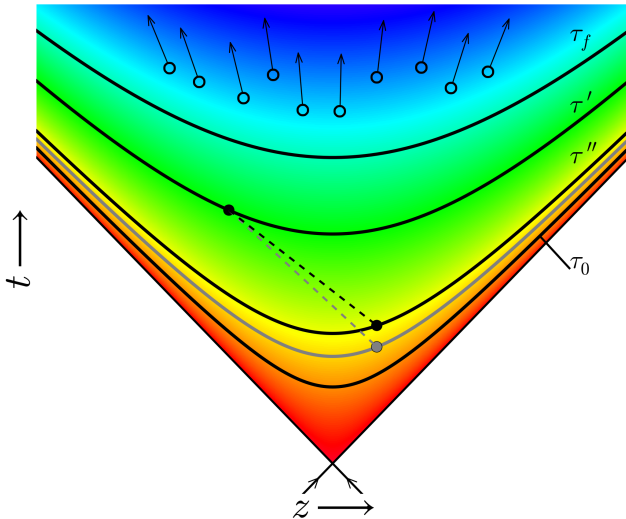
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} \mathbf{N}(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



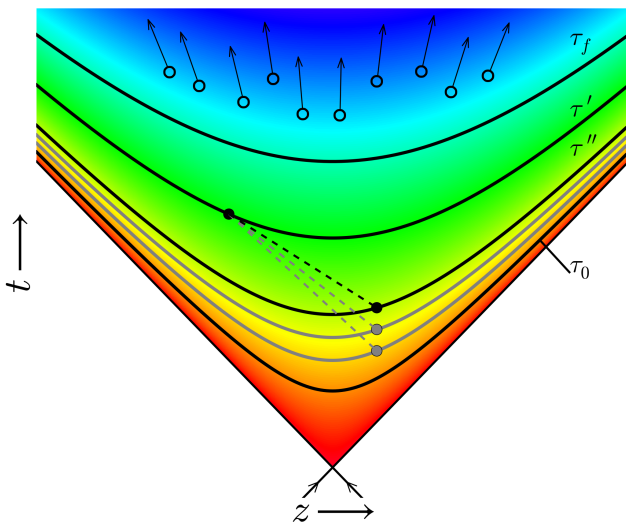
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



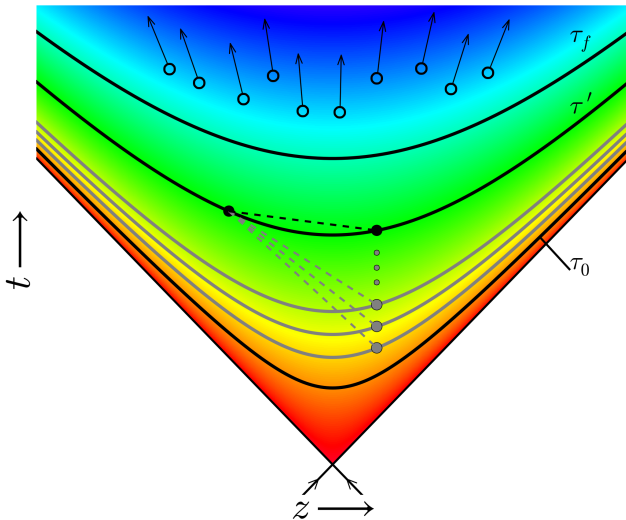
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



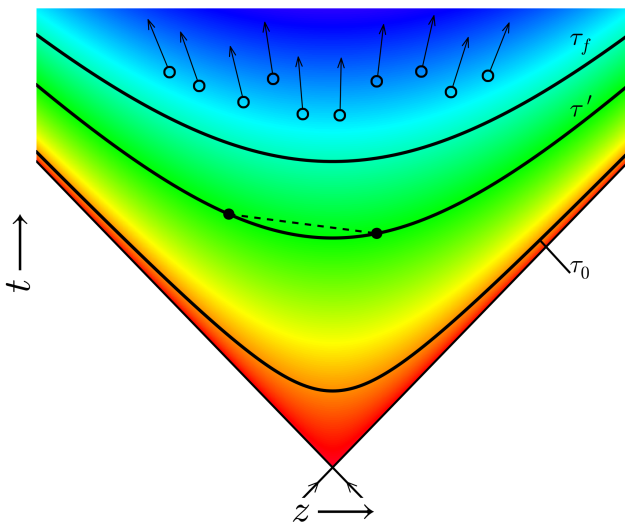
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi \delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



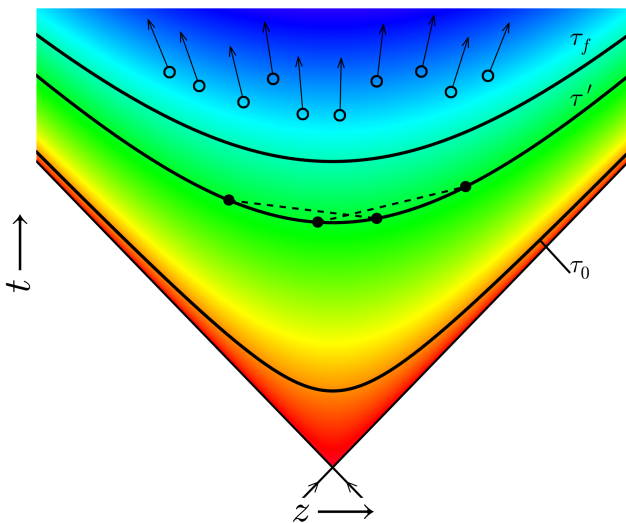
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



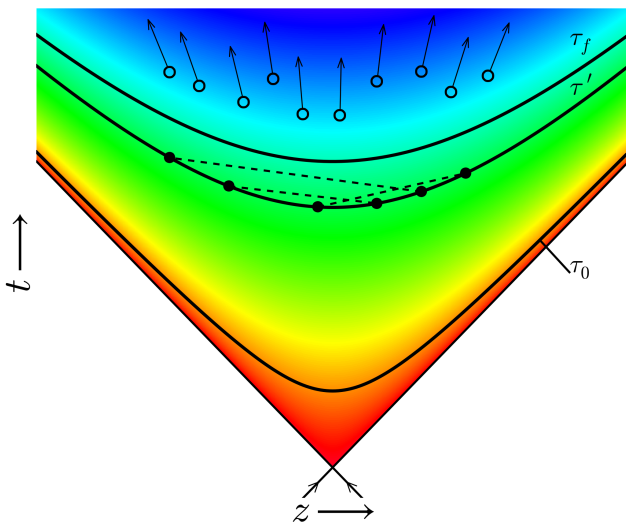
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



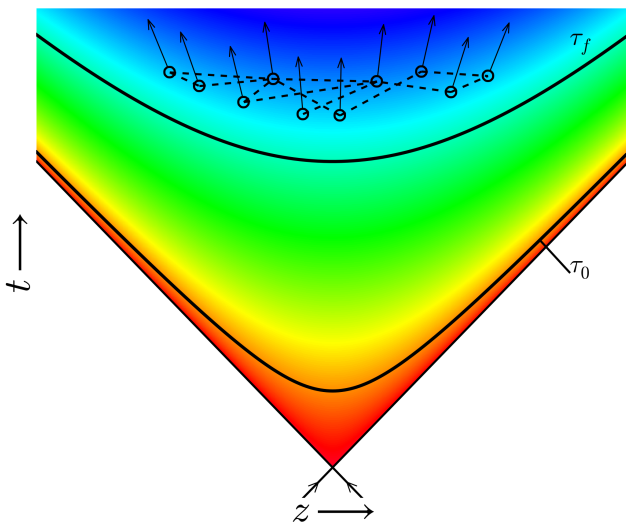
$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



$$\begin{aligned}
 \langle \tilde{X}(k_1, \tau_f) \tilde{Y}(k_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$



$$\begin{aligned}
 \langle \tilde{\mathbf{X}}(\mathbf{k}_1, \tau_f) \tilde{\mathbf{Y}}(\mathbf{k}_2, \tau_f) \rangle &= 4\pi\delta(k_1 + k_2) \int_{\tau_0}^{\tau_f} \frac{d\tau'}{\tau'} \tilde{\mathbf{G}}_{\mathbf{X}}(\mathbf{k}_1; \tau_f, \tau') \int_{\tau_0}^{\tau_f} \frac{d\tau''}{\tau''} \tilde{\mathbf{G}}_{\mathbf{Y}}(\mathbf{k}_2; \tau_f, \tau'') \\
 &\times \left(\frac{1}{2\tau_C} e^{-\frac{|\tau' - \tau''|}{\tau_C}} \right) \int_0^\infty \frac{d\tau}{\tau_C} e^{-\frac{\tau}{\tau_C}} N(\min(\tau', \tau'') - \tau)
 \end{aligned}$$