

# Chiral symmetry restoration by parity doubling and the structure of neutron stars

Michał Marczenko\*, David Blaschke, Krzysztof Redlich, Chihiro Sasaki  
Institute of Theoretical Physics, University of Wrocław, Poland



Uniwersytet  
Wrocławski

## Introduction

We discuss the implications of the chiral symmetry restoration by parity doubling of baryons on the mass-radius relation for compact stars. Most important for this study are the nucleon (neutron, proton) and N(1535) states. We also consider the threshold for the direct URCA process for which a new relationship is given. We show that the existence of high-mass stars might not necessarily signal the deconfinement of quarks.

## SU(2) Hybrid QMN model

- Thermodynamic potential in the mean-field approximation

$$\Omega = V_\sigma + V_\omega + V_\rho + V_b + \sum_x \Omega_x, \quad x = p_\pm, n_\pm, u, d$$

$$\Omega_x = \gamma_x \int \frac{d^3p}{(2\pi)^3} T [\ln(1 - n_x) + \ln(1 - \bar{n}_x)]$$

- Hadrons

- coupled to  $\sigma$ ,  $\pi$ ,  $\omega$ ,  $\rho$  and  $b$
- parity doubling structure ( $m_{p_\pm} = m_{n_\pm} \equiv m_\pm$ )

$$m_\pm = \frac{1}{2} \left[ \sqrt{(g_1 + g_2)^2 \sigma^2 + 4m_0^2} \mp (g_1 - g_2)\sigma \right] \xrightarrow{\sigma \rightarrow 0} m_0$$

- Quarks

- coupled to  $\sigma$  and  $b$
- Standard linear sigma for quarks ( $m_u = m_d \equiv m_q$ )

$$m_q = g_\sigma \sigma \xrightarrow{\sigma \rightarrow 0} 0$$

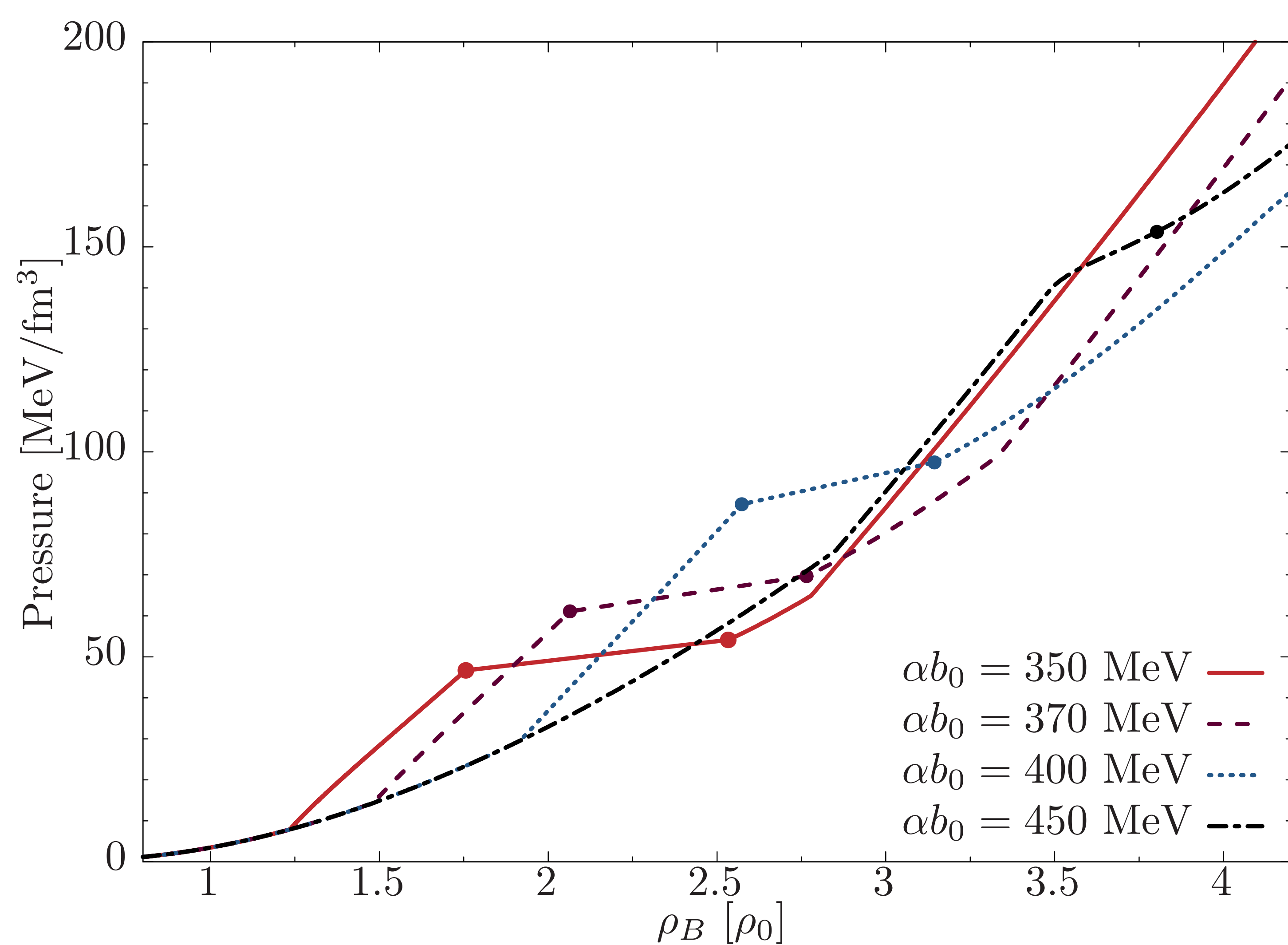
## Statistical (de)confinement

- Modification of the Fermi-Dirac distribution functions

$$n_q = \theta(\mathbf{p}^2 - b^2) f_q \quad n_\pm = \theta(\alpha^2 b^2 - \mathbf{p}^2) f_\pm$$

- $\alpha$  - an external parameter to be studied here
- Uncertainty principle  $\rightarrow$  size of a hadron  $\sim 1/b$
- $b$  is promoted to a field generated by a scalar potential  $V_b \sim -b^2 + b^4$
- $b$  is not an order parameter (it is not connected to any fundamental QCD symmetry)!

## Equation of State under the NS conditions



Equation of state for different values of the  $\alpha$  parameter under the NS conditions. Circles indicate the onset and restoration of the chiral phase transition.

## References

- [1] S. Benic, I. Mishustin and C. Sasaki, Phys. Rev. D **91**, 125034 (2015)
- [2] M. Marczenko and C. Sasaki, Phys. Rev. D **97**, 036011 (2018)
- [3] M. Marczenko, D. Blaschke, K. Redlich and C. Sasaki, arXiv:1805.06886

## Direct URCA threshold

- direct URCA process

$$n_+ \rightarrow p_+ + e + \bar{\nu}_e$$

- Chiral symmetry broken

$$\frac{1}{1 + (1 + \sqrt[3]{Y_e})^3} \approx 11 - 15\%$$

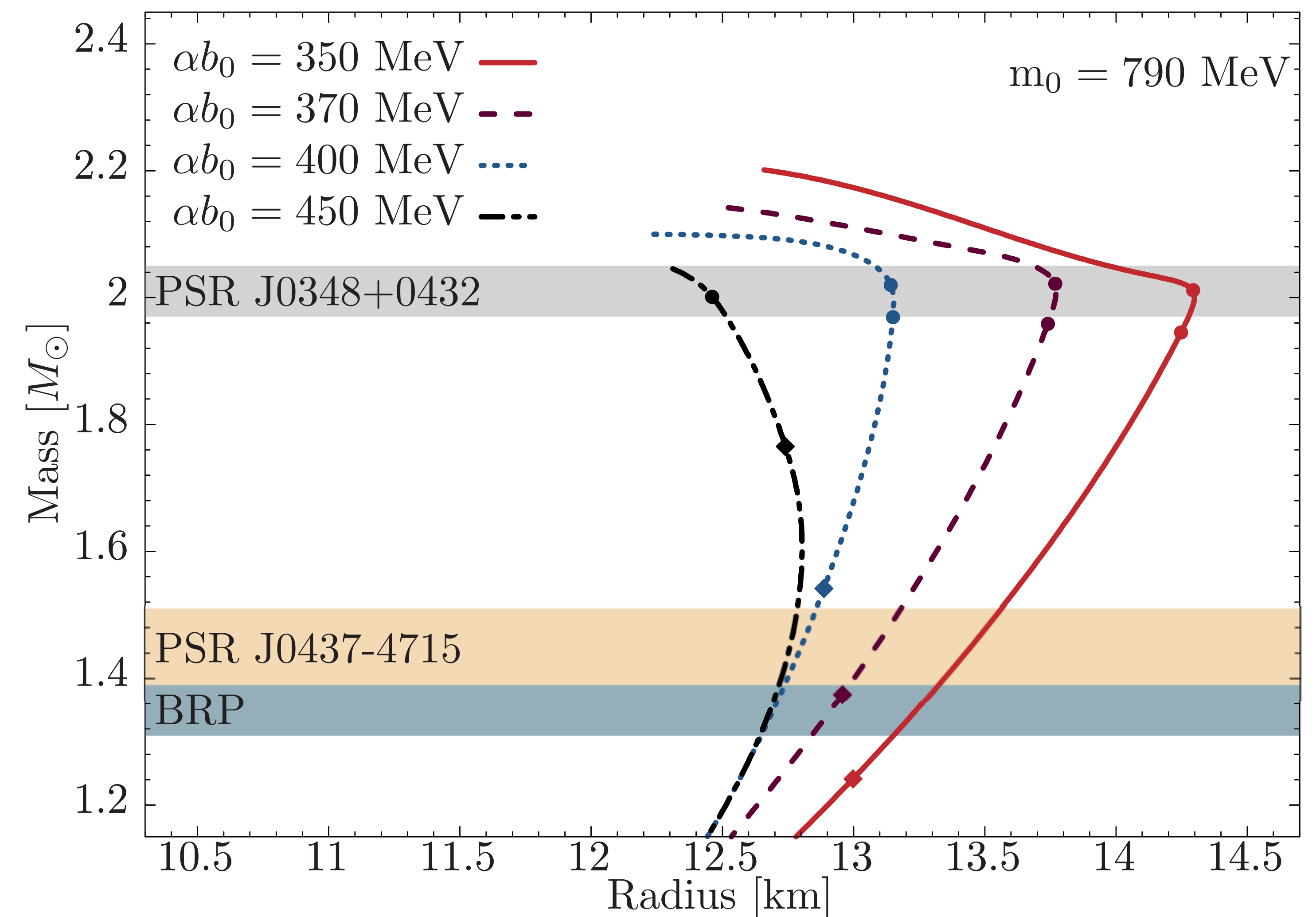
- Proton fraction

$$Y_{p_+} = \frac{\rho_{p_+}}{\rho_{p_+} + \rho_{n_+} + \rho_{p_-} + \rho_{n_-}}$$

- Chiral symmetry restored

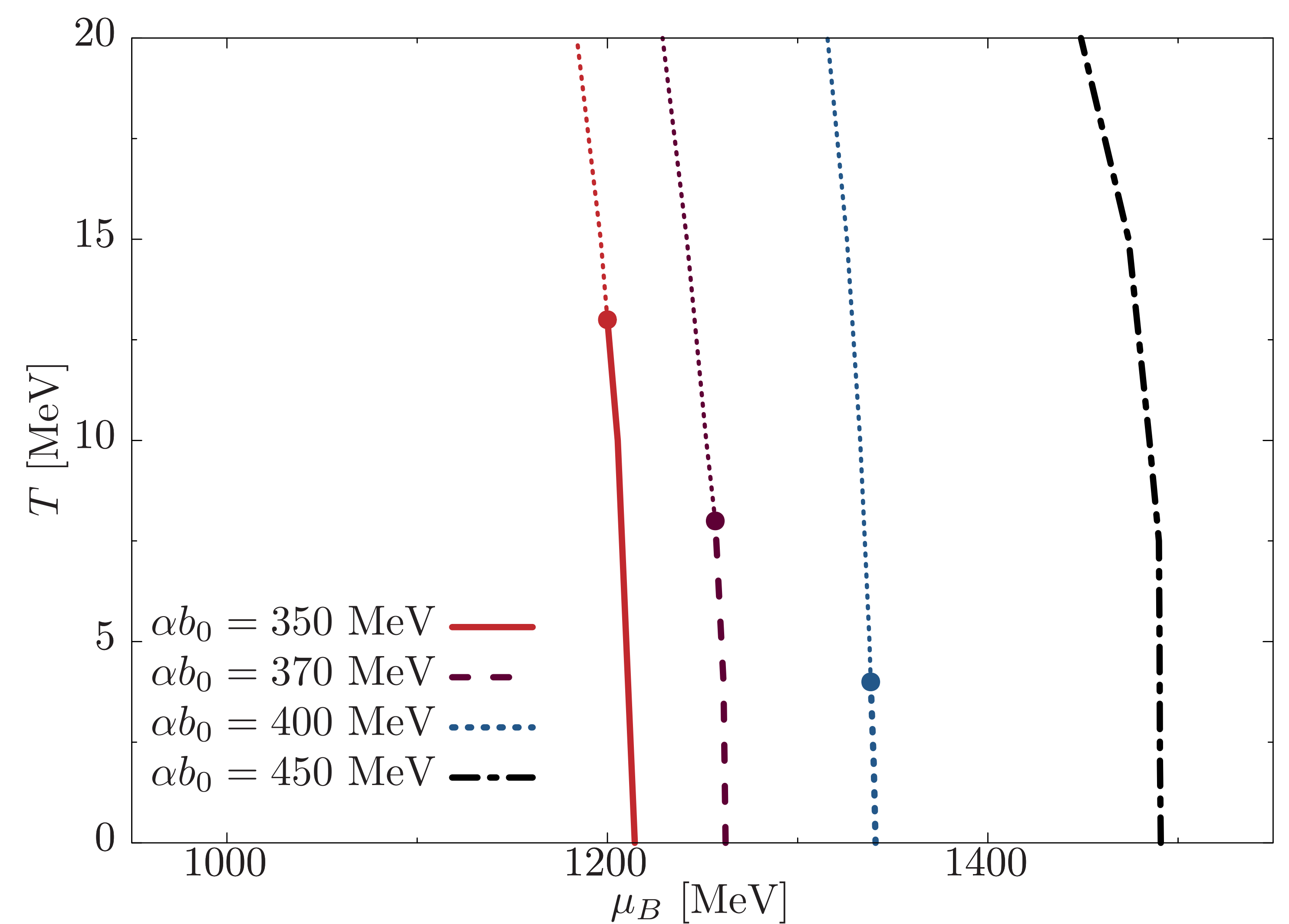
$$\frac{1}{1 + (1 + \sqrt[3]{2Y_e})^3} \approx 8 - 11\%$$

## Neutron Star Structure



Mass-radius relations. Circles indicate the onset and restoration of the chiral phase transition, diamonds indicate the direct URCA threshold.

## QCD phase diagram for isospin-symmetric matter



Predicted low-temperature QCD phase diagram for isospin-symmetric matter. Shown are the chiral phase transitions for different values of  $\alpha$ . The transitions from hadrons to quarks happen at higher baryon chemical potentials and are not shown here.

## Summary

- High-mass stars & phase transition  $\rightarrow$  not necessarily signal for deconfinement
- Parity doubling  $\rightarrow$  lower estimate for direct URCA threshold in the chirally restored phase.