



Understanding the QCD spectrum: progress and prospects from Lattice QCD

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Colloquium @ GSI 10th May 2016

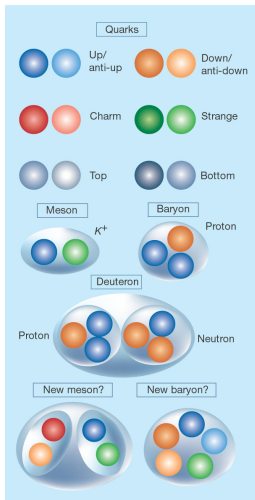
PLAN

- **The QCD spectrum**
 - Quark models and QCD.
 - New discoveries and further puzzles.
- **A consumers guide to Lattice QCD**
 - compromises and consequences
- **Discussion and selected results (mostly charm/charmonium)**
 - parallel tracks for progress
 - old challenges and new results
 - new challenges and exploratory results
 - **precision spectroscopy of single hadron states including excited and exotic states**
 - **spectroscopy of scattering states - progress and challenges**
- **Summary**

WHY HADRON SPECTROSCOPY?

- Many recently discovered hadrons have unexpected properties.
- Understand the hadron spectra to separate EW physics from strong-interaction effects
- Techniques for non-perturbative physics useful for physics at LHC energies.
- Understanding EW symmetry breaking may require nonperturbative techniques at TeV scales, similar to spectroscopy at GeV.
- Better techniques may help understand the nature of masses and transitions

OBJECTS OF INTEREST



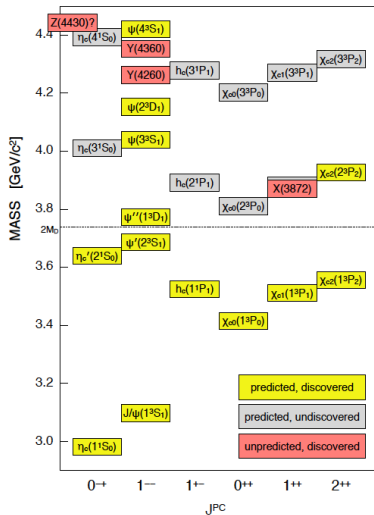
- Built from fundamental objects: quarks and gluons
- Fields of Lagrangian in colorless combinations: confinement

quark model object	structure
meson	$3 \otimes \bar{3} = 1 \oplus 8$
baryon	$3 \otimes 3 \otimes 3 = 1 \oplus 8 \oplus 8 \oplus 10$
hybrid	$\bar{3} \otimes 8 \otimes 3 = 1 \oplus 8 \oplus 8 \oplus 8 \oplus 10 \oplus 10$
glueball	$8 \otimes 8 = 1 \oplus 8 \oplus 8 \oplus 10 \oplus 10$
⋮	⋮

- **This is a model.** QCD does not always respect this constituent picture! There can be strong mixing.

HADRON STATES: QUESTIONS AND PUZZLES NOT RESOLVED BY MODELS

- States classified by $J^{P(C)}$ multiplets (representations of the Poincaré symmetry).
- In quark models, mesons with $P = (-1)^J$ and $CP = -1$ forbidden. **Some J^{PC} combinations don't appear: 0^{+-} , 0^{-+} , 1^{-+} , 2^{+-} , ...** These **exotics** (not just a $q\bar{q}$ pair) allowed in **QCD**.
- Many more baryon states predicted than observed - the missing resonance problem.
- Where are the other states **QCD** allows - hybrids, glueballs, ... ?



from D. Bettoni CIPANP2015

WHY LATTICE QCD ?

- A systematically-improvable non-perturbative formulation of QCD
 - Well-defined theory with the lattice a UV regulator
- Arbitrary precision is in principle possible
 - of course algorithmic and field-theoretic “wrinkles” can make this challenging!
- Starts from first principles - i.e. from the QCD Lagrangian
 - inputs are quark mass(es) and the coupling - can explore mass dependence and coupling dependence but getting to physical values can be hard!

A LATTICE QCD PRIMER

Start from the QCD Lagrangian:

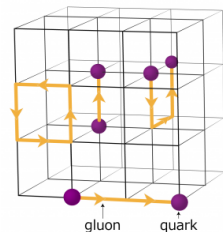
$$\mathcal{L} = \bar{\Psi} (i\gamma^\mu D_\mu - m) \Psi - \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a$$

Gluon fields on links of a hypercube;

Quark fields on sites: approaches to fermion discretisation -

Wilson, Staggered, Overlap;

Derivatives → finite differences.



Solve the QCD path integral on a finite lattice with spacing $a \neq 0$ estimated stochastically by Monte Carlo. Can only be done effectively in a Euclidean space-time metric (no useful importance sampling weight for the theory in Minkowski space).

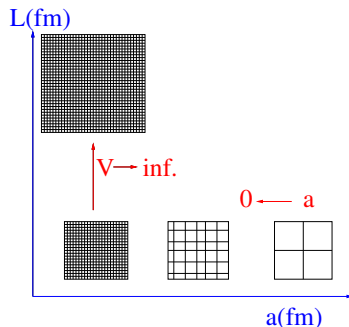
Observables determined from (Euclidean) path integrals of the QCD action

$$\langle O \rangle = 1/Z \int DUD\bar{\Psi}D\Psi O[U, \bar{\Psi}, \Psi] e^{-S[U, \bar{\Psi}, \Psi]}$$

Compromises and the Consequences

1. Working in a finite box at finite grid spacing

- Identify a “scaling window” where physics doesn't change with a or V . Recover continuum QCD by extrapolation.



A costly procedure but a regular feature in lattice calculations now

2. Simulating at physical quark masses

- Computational cost grows rapidly with decreasing quark mass $\rightarrow m_q = m_{u,d}$ costly. Care needed vis location of decay thresholds and identification of resonances.
- c-quark can be handled relativistically. b-quark with: NRQCD, FNAL etc.

Better algorithms for physical light quarks and/or chiral extrapolation. Relativistic m_b in reach

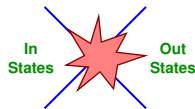
2. Breaking symmetry



- Lorentz symmetry broken at $a \neq 0$ so $SO(4)$ rotation group broken to discrete rotation group of a hypercube.

Classify states by irreps of O_h and relate by subduction to J values of O_3 . Lots of degeneracies in subduction for $J \geq 2$ and physical near-degeneracies. Complicates spin identification.

Spin identification at finite lattice spacing: 0707.4162, 1204.5425



3. Working in Euclidean time.

- Scattering matrix elements not directly accessible from Euclidean QFT [*Maiani-Testa theorem*]. Scattering matrix elements: asymptotic $|\text{in}\rangle, |\text{out}\rangle$ states:
 $\langle \text{out} | e^{i\hat{H}t} | \text{in} \rangle \rightarrow \langle \text{out} | e^{-\hat{H}t} | \text{in} \rangle$. Euclidean metric: project onto ground state. Analytic continuation of numerical correlators an ill-posed problem.

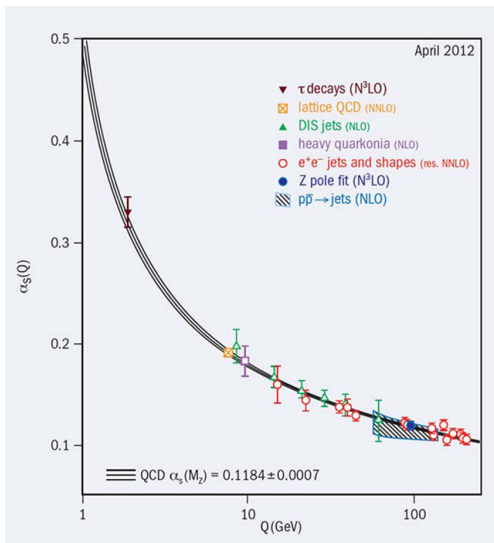
Lüscher and generalisations of: method for indirect access.

4. Quenching

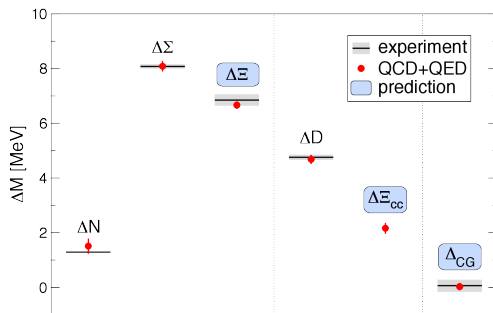
No longer an issue: Simulations done with $N_f = 2, 2 + 1, 2 + 1 + 1$.

Validation:
can we reproduce known results and make verified predictions?

VALIDATION

The running coupling, α_s 

Baryon electromagnetic mass splittings

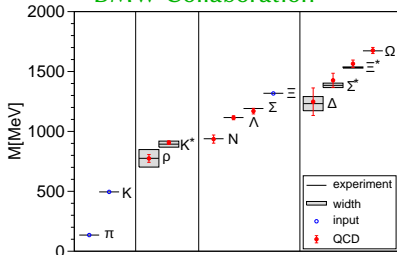


QED + QCD

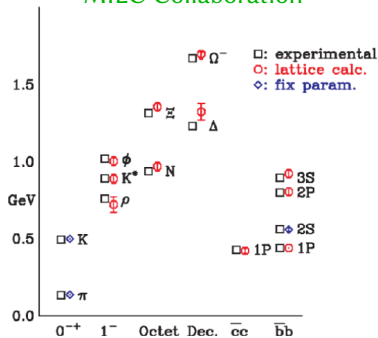
BMW Collab. Science 347 (2015) 1452

CONVERGENCE THROUGH UNIVERSALITY

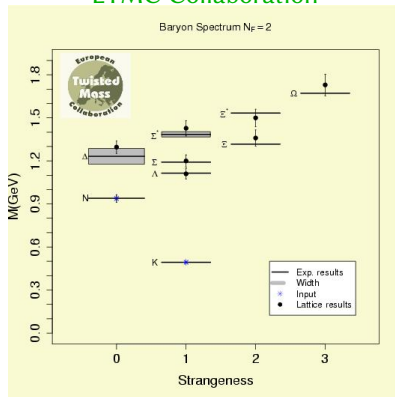
BMW Collaboration



MILC Collaboration



ETMC Collaboration



BMW: SW-Wilson [Science 322:1224-1227,2008.]

ETMC: Twisted Mass [arXiv:0910.2419,0803.3190]

MILC: Staggered [arXiv:0903.3598]

TWO STRATEGIES FOR PROGRESS

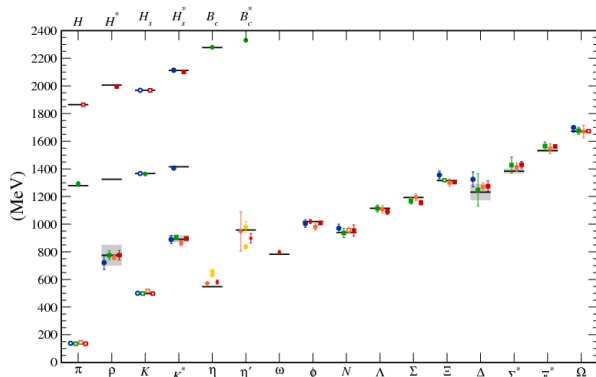
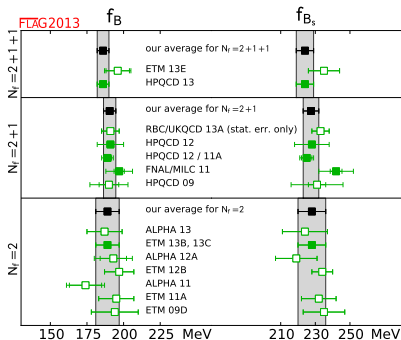
Gold-plated quantities

- e.g. single hadron states, or decays below thresholds
- phenomenologically relevant
- incremental progress
- robust/well-tested methods
- careful error budgeting

New directions

- new ideas - theoretical and algorithmic that open new avenues
- recent examples are scattering states, $g-2$, ...
- also improves gold-plated
- pioneering, error budgets not yet “robust”

STRATEGIES FOR PROGRESS: GOLD PLATED QUANTITIES - A SELECTION



FLAG 2013 itpwiki.unibe.ch/flag/

A. Kronfeld, *Ann.Rev.Nucl.Part.Sci.* 62 (2012)

- Stable single-hadron states, below thresholds
- Including continuum extrapolation, realistic quark masses, renormalisation etc

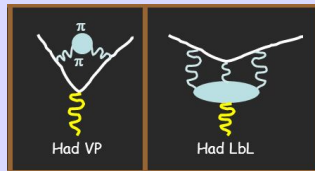
STRATEGIES FOR PROGRESS: NEW DIRECTIONS - A SELECTION

New ideas in hadron spectroscopy

- **Distillation** for quark propagation enabled isoscalars, precision spectroscopy, efficient calculation and motivated ...
- **Scattering and Coupled channels** new theoretical ideas to tackle scattering states and study (X,Y,Z) , resonance parameters in eg $\pi K, \pi \eta$...

New ideas for $g-2$

- Dominant uncertainty is in hadronic contributions - HVP and HLbL



- lots more!

Lattice Hadron Spectroscopy

precision & pioneering results

- (i) Precision spectroscopy of single-hadron states
- (ii) Exploratory studies of “exotic” and scattering states

A RECIPE FOR (MESON) SPECTROSCOPY

- Construct a basis of local and non-local operators $\bar{\Psi}(x)\Gamma D_i D_j \dots \Psi(x)$ from **distilled fields - the key enabling idea!** [PRD80 (2009) 054506].
- Build a correlation matrix of two-point functions

$$C_{ij} = \langle 0 | \mathcal{O}_i \mathcal{O}_j^\dagger | 0 \rangle = \sum_n \frac{Z_i^n Z_j^{n\dagger}}{2E_n} e^{-E_n t}$$

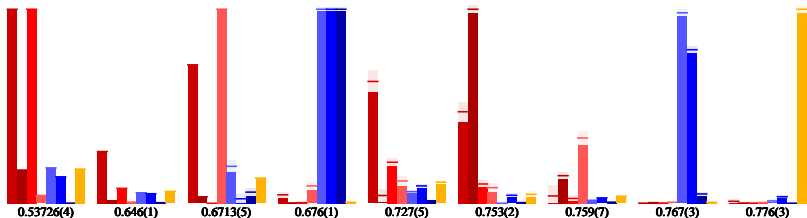
- Ground state mass from fits to $e^{-E_n t}$
 - Beyond ground state: Solve generalised eigenvalue problem
- $$C_{ij}(t) v_j^{(n)} = \lambda^{(n)}(t) C_{ij}(t_0) v_j^{(n)}$$

- eigenvalues: $\lambda^{(n)}(t) \sim e^{-E_n t} [1 + O(e^{-\Delta E t})]$ - principal correlator
- eigenvectors: related to overlaps $Z_i^{(n)} = \sqrt{2E_n} e^{E_n t_0/2} v_j^{(n)\dagger} C_{ji}(t_0)$

- operators of definite J^{PC} constructed in step 1 are subduced into the relevant irrep
- a subduced irrep carries a “memory” of continuum spin J from which it was subduced - it **overlaps** predominantly with states of this J .

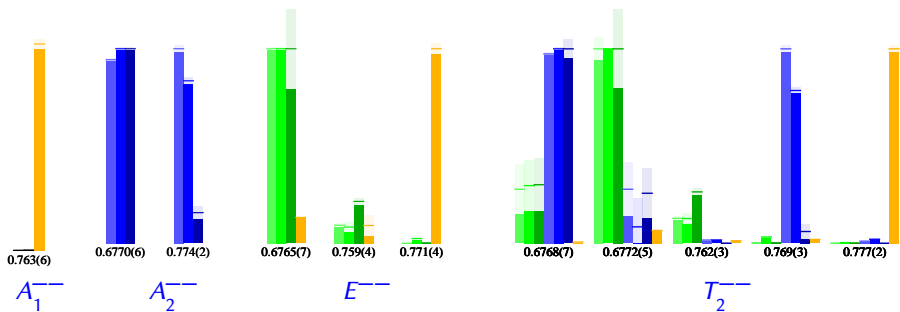
J	0	1	2	3	4
A_1	1	0	0	0	1
A_2	0	0	0	1	0
E	0	0	1	0	1
T_1	0	1	0	1	1
T_2	0	0	1	1	1

- Using $Z = \langle 0|\Phi|k\rangle$, helps to identify continuum spins
- For high spins, can look for agreement between irreps
- Data below for T_1^- irrep, colour-coding is **Spin 1**, **Spin 3** and **Spin 4**.



... THE REST OF THE SPIN-4 STATE

- All polarisations of the spin-4 state are seen
- Spin labelling: **Spin 2**, **Spin 3** and **Spin 4**.

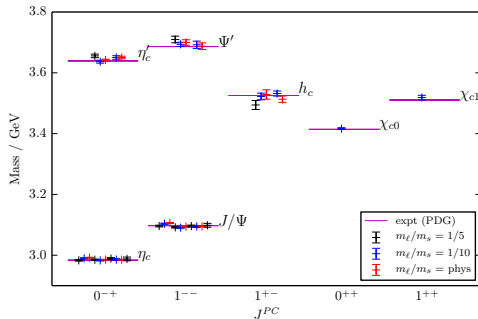


**Precision Spectroscopy:
states below strong decay thresholds**

SINGLE HADRON STATES: BELOW THRESHOLD - “GOLD-PLATED”

- Methods: tested, validated.
- High statistics and improved actions for precise results.
- Different actions in agreement.
- Simulation at m_q^{phys} or extrapolation
 $m_q \rightarrow m_q^{\text{phys}}$.
- Discretisation errors $\mathcal{O}(am_c)$ and $\mathcal{O}(am_b)$ under control,

Charmonium, HPQCD 1411.1318



Continuum limit, physical quark masses

No disconnected diagrams in $c\bar{c}$ spectrum: OZI suppressed - assumed to be small
 \Rightarrow mixing with lighter states not included

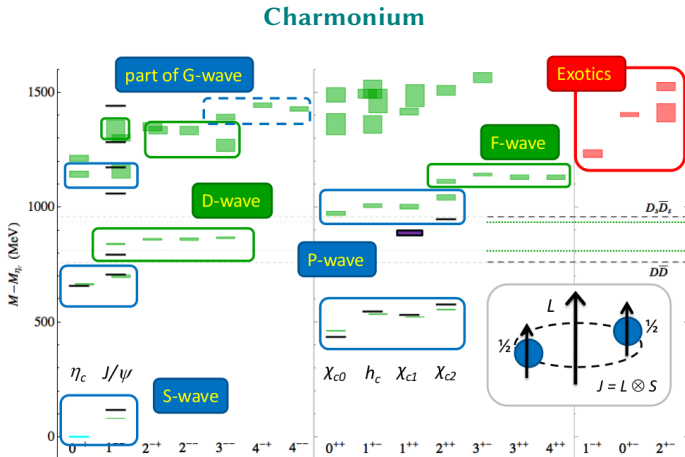
**Precision Spectroscopy:
single hadron states near/above thresholds**

SINGLE HADRON STATES: ABOVE THRESHOLD

Precision calculation of high spin ($J \geq 2$) and exotic states is relatively new

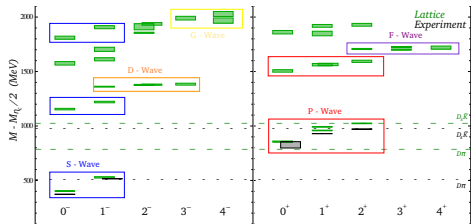
Caveat Emptor

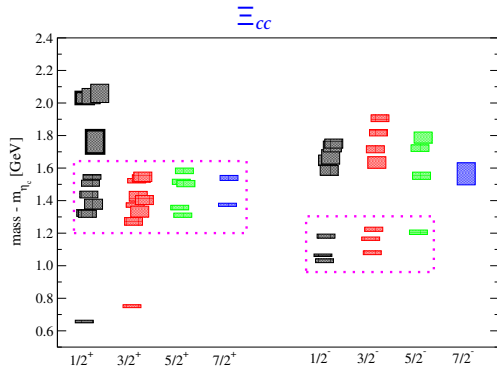
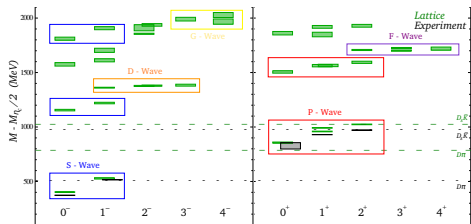
- Only single-hadron operators
- Physics of multi-hadron states appears to need relevant operators
- No continuum extrapolation
- Relatively heavy pions ← already changing



from HSC 2012

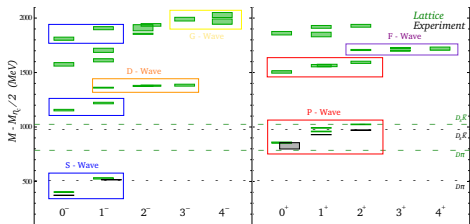
→ Expect improvements now methods established



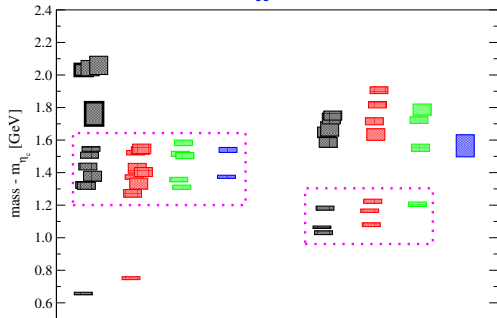


D Meson Spectrum - By J^P

[arXiv:1301.7670]



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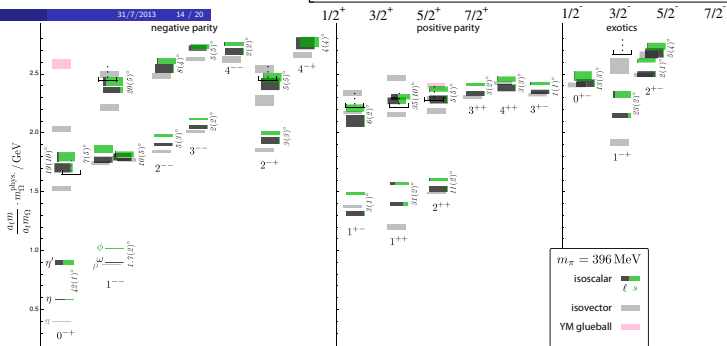
Graham Moir (TCD)

Charm Physics

31/7/2013 14 / 20

HadSpec
results

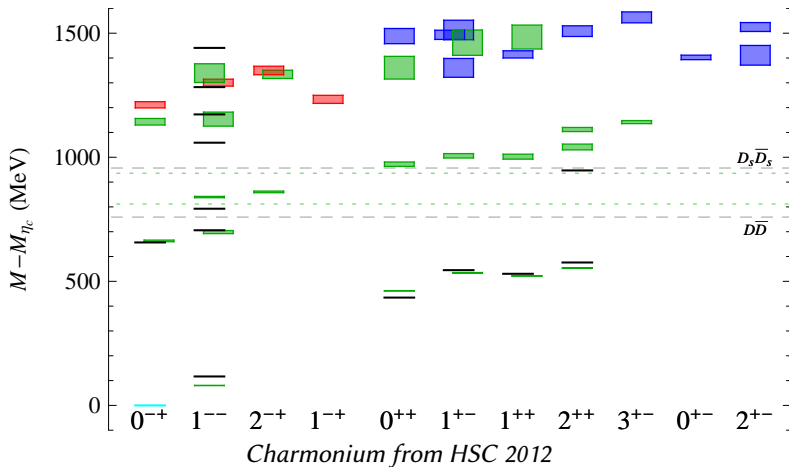
light mesons



$m_\pi = 396$ MeV

isovector ℓs
isovector
YM glueball

HYBRIDS



Lightest hybrid supermultiplet and excited hybrid supermultiplet same pattern and scale as in open charm and light [\[HadSpec:1106.5515\]](#) sectors.



EXPLORATORY STUDIES OF SCATTERING STATES

Characterised by

- **New methods (developed/applied in last 5 years)**
 - **algorithmic:** distillation allows access to all elements of propagators *and* construction of sophisticated basis of operators.
 - **theoretical:** spin-identification; construction of multi-hadron operators and mesons in flight; scattering below inelastic thresholds; coupled-channels (new in '14).
- **Generally high statistics, improved actions etc - results can be very precise.**
- **Systematic errors not all controlled in exploratory studies: e.g. no continuum extrapolation, relatively heavy pions ...**

Rapid progress in the last 5 years!

SCATTERING IN A EUCLIDEAN THEORY

Lose direct access to scattering in (Euclidean) lattice calculations

Lüscher found a way to extract scattering information in the elastic region from LQCD .

[NPB354, 531-578 (1991)]

- related **lattice energy levels in a finite volume** to a decomposition of the scattering amplitude in **partial waves in infinite volume**

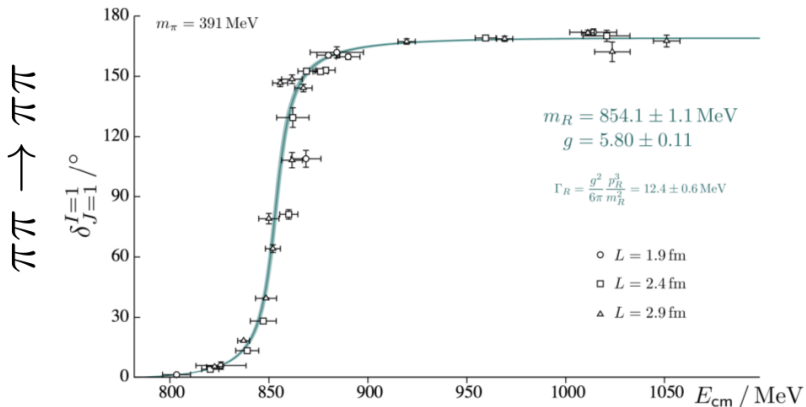
$$\det \left[\cot \delta(E_n^*) + \cot \phi(E_n, \vec{P}, L) \right] = 0$$

and $\cot \phi$ a known function (containing a generalised zeta function).

- The idea dates from the quenched era. To use it in a dynamical simulation need energy levels at extraordinary precision. This is why it has taken a while ...

USING LÜSCHER'S IDEA

Now in use to determine resonance parameters



$$m_\pi = 391 \text{ MeV}$$

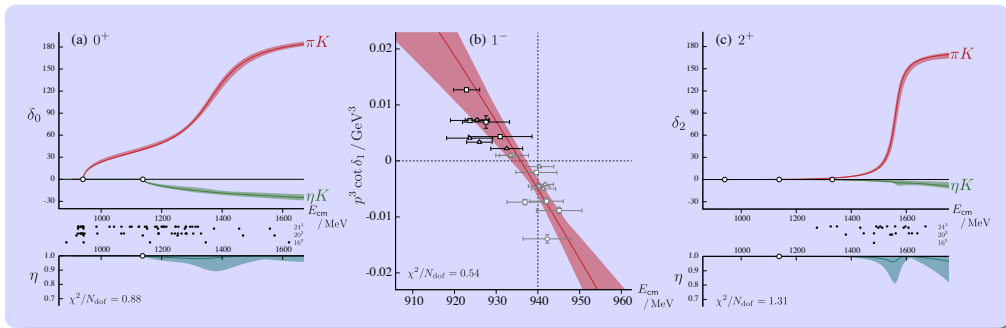
from Dudek, Edwards, Thomas in *Phys.Rev. D87* (2013) 034505

Many talks at Lattice 2015 & 2016

EVEN MORE RECENT PROGRESS

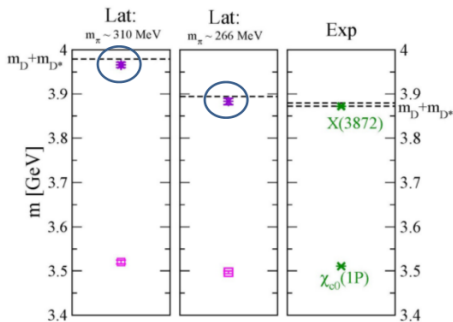
- Generalised for: moving frames; non-identical particles; multiple two-particle channels, particles with spin, by many authors.
- The precision and robustness of some numerical implementations is now very impressive. *[See talks at Lattice 2015 & 2016]*
- First coupled-channel resonance in a lattice calculation

$\pi K \rightarrow \eta K$ by D. Wilson et al 1406.4158 and 1507.02599



X(3872) - A FIRST LOOK

Prelovsek & Leskovec 1307.5172



ground state: $\chi_{c1}(1P)$

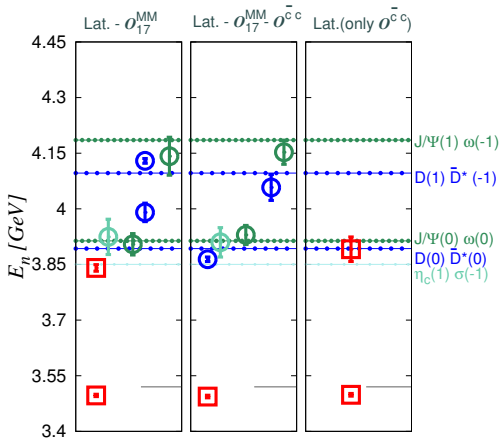
$D\bar{D}^*$ scattering mx: pole just below thr.

Threshold $\sim m_{u,d}$ and m_c discretisation?

Also results from Lee et al 1411.1389

State is within 1MeV of $D^0\bar{D}^{*0}$ and 8MeV of D^+D^* thresholds: isospin breaking effects important?

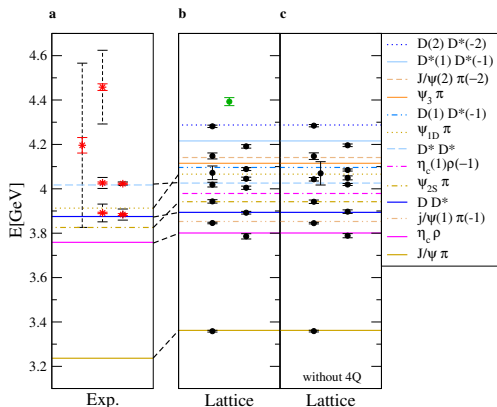
Padmanath, Lang, Prelovsek 1503.03257



X(3872) not found if $c\bar{c}$ not in basis.

Z_c^+ - FIRST LOOK ON THE LATTICE

Prelovsek, Lang, Leskovec, Mohler: 1405.7615



- 13 expected 2-meson e' states found (black)
- no additional state below 4.2 GeV
- no Z_c^+ candidate below 4.2 GeV

Similar conclusion from [Lee et al \[1411.1389\]](#) and [Chen et al \[1403.1318\]](#)

Why no eigenstate for Z_c ? Is Z_c^+ a coupled channel effect? What can other groups say? Work needed!

Z_c^+

An “exotic” hadron i.e. does not fit in the quark model picture.

There are a number of exploratory calculations on the lattice.

Challenges:

- The Z_c^+ (and most of the XYZ states) lies above several thresholds and so decay to several two-meson final states
- requires a coupled-channel analysis for a rigorous treatment
- on a lattice the number of relevant coupled-channels is large for high energies.

State of the art in coupled-channel analysis:

- Lüscher: $K\pi$, $K\eta$ [HSC 2014,2015]
- HALQCD: Z_c [preliminary results]

MANY OTHER STATES BEING INVESTIGATED

Tetraquarks:

- Double charm tetraquarks ($J^P = 1^+, I = 0$) by HALQCD [PLB712 (2012)]
 - attractive potential, no bound tetraquark state
- Charm tetraquarks: variational method with DD^* , D^*D^* and tetraquark operators finds no candidate.

Y(4140)

- Ozaki and Sasaki [1211.5512] - no resonant Y(4140) structure found
- Padmanath, Lang, Prelovsek [1503.03257] considered operators: $c\bar{c}$, $(\bar{c}s)(\bar{s}c)$, $(\bar{c}c)(\bar{s}s)$, $[\bar{c}\bar{s}][cs]$ in $J^P = 1^+$. Expected 2-particle states found and χ_{c1} , $X(3872)$ **not** Y(4140).

⋮

See Prelovsek @ Charm2015 for more



FORM FACTORS AND RESONANCES: NEW IN 2016

a_0 resonance in $\pi\eta, K\bar{K}$

- First calculation of coupled channel meson-meson scattering in $I=1$, G-parity negative sector from lattice QCD
- $a_0(980)$ -like resonance strongly coupled to both $\pi\eta$ and $K\bar{K}$ identified as a pole in the complex energy plane. [Hadron Spectrum Collaboration, Dudek et al, 1602.05122](#)

Transition amplitudes

- $\pi\pi \rightarrow \pi\gamma^*$ and the resonant $\rho \rightarrow \pi\gamma^*$ transition. [Hadron Spectrum Collaboration 1604.03530](#)
- First lattice calculation of the form factor of an unstable hadron, albeit at unphysical pion mass (400MeV).
- Extensions to nucleon resonances, heavy flavour decays and an extension to the coupled channel case, will allow calculations of radiative transitions with exotic hybrid mesons.



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Expect many new results at Lattice 2016!



CHALLENGES

There have been many successes in lattice spectroscopy in the last 5 years including

- “Gold-plated” quantities increasingly well-determined
- New ideas have led to rapid progress in spectroscopy - precision excited and exotic states and scattering analyses

Many challenges remain

- ◆ Improving existing calculations - understanding the effect of lighter light quarks on thresholds etc, simulations at multiple and larger volumes
- ◆ Handling the large number of coupled-channels that emerge on larger volumes
- ◆◆ A general framework for coupled channels for scattering involving more than 2 hadrons. Some progress [M. Hansen @ Lattice 2015]
- Haven’t discussed the many other open problems including finite density, BSM, ... !



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Thanks for listening!

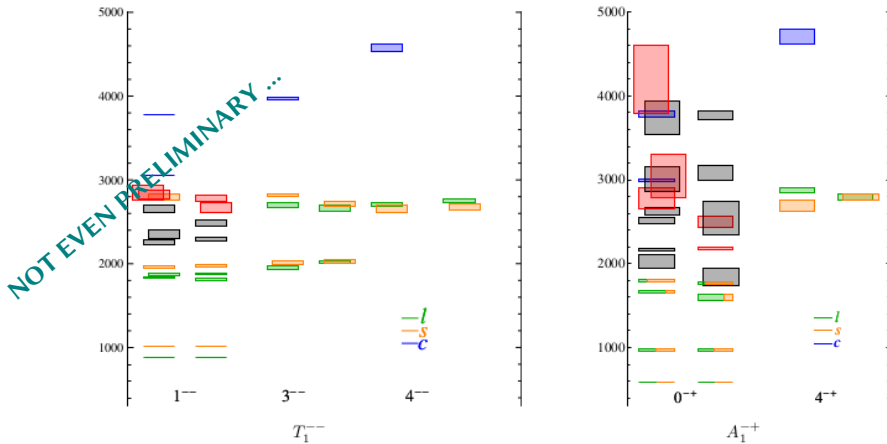
Backup Slides

LAST COMMENT ON SINGLE-HADRON SPECTRUM

Disconnected diagrams a remaining uncertainty in most $c\bar{c}$ calculations.
Distillation allows precision determination. BUT it's a can of worms!

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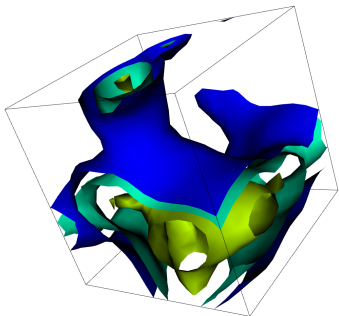
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$$\Delta(1^{--}) = -17(16)\text{MeV}$$

from HadSpec

WHAT'S THE PLOT?



Distillation

- A new approach to quark propagation by redefining smearing as a projection operator
- Basis vectors of the distillation operator (lattice laplacian) look like confining blobs