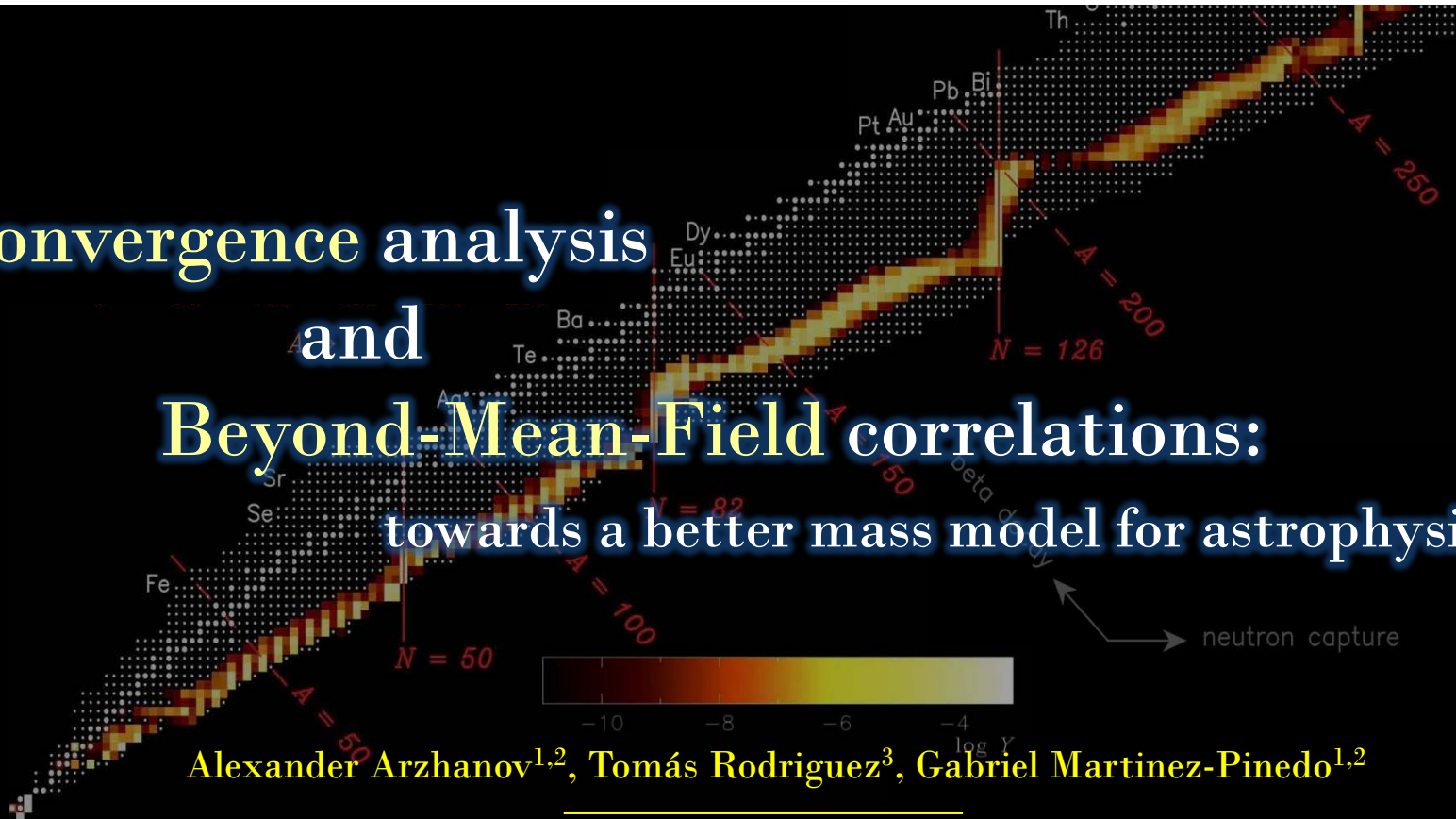




TECHNISCHE
UNIVERSITÄT
DARMSTADT

Convergence analysis and Beyond-Mean-Field correlations: towards a better mass model for astrophysics

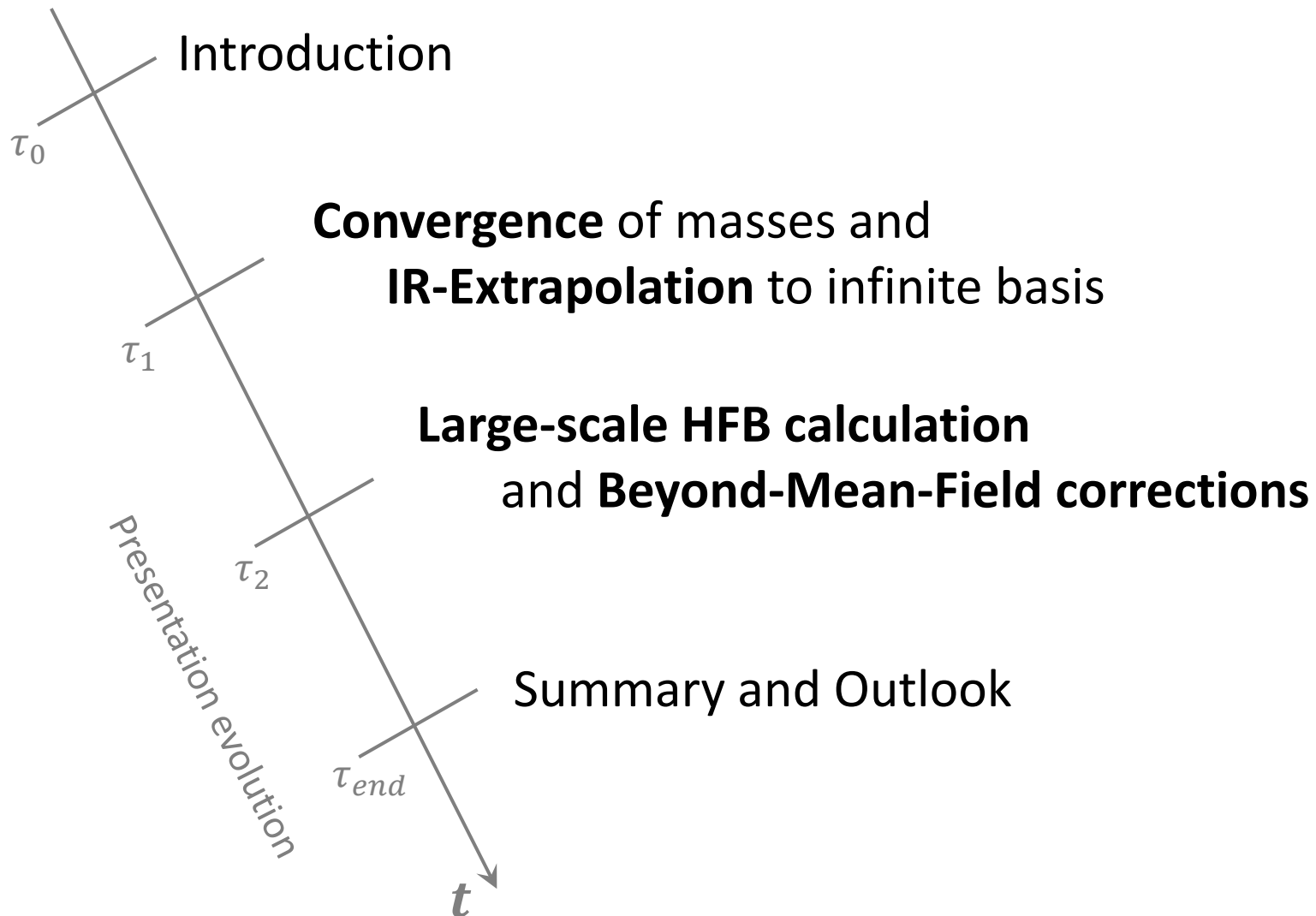


Alexander Arzhanov^{1,2}, Tomás Rodríguez³, Gabriel Martínez-Pinedo^{1,2}

¹Technische Universität Darmstadt, Germany

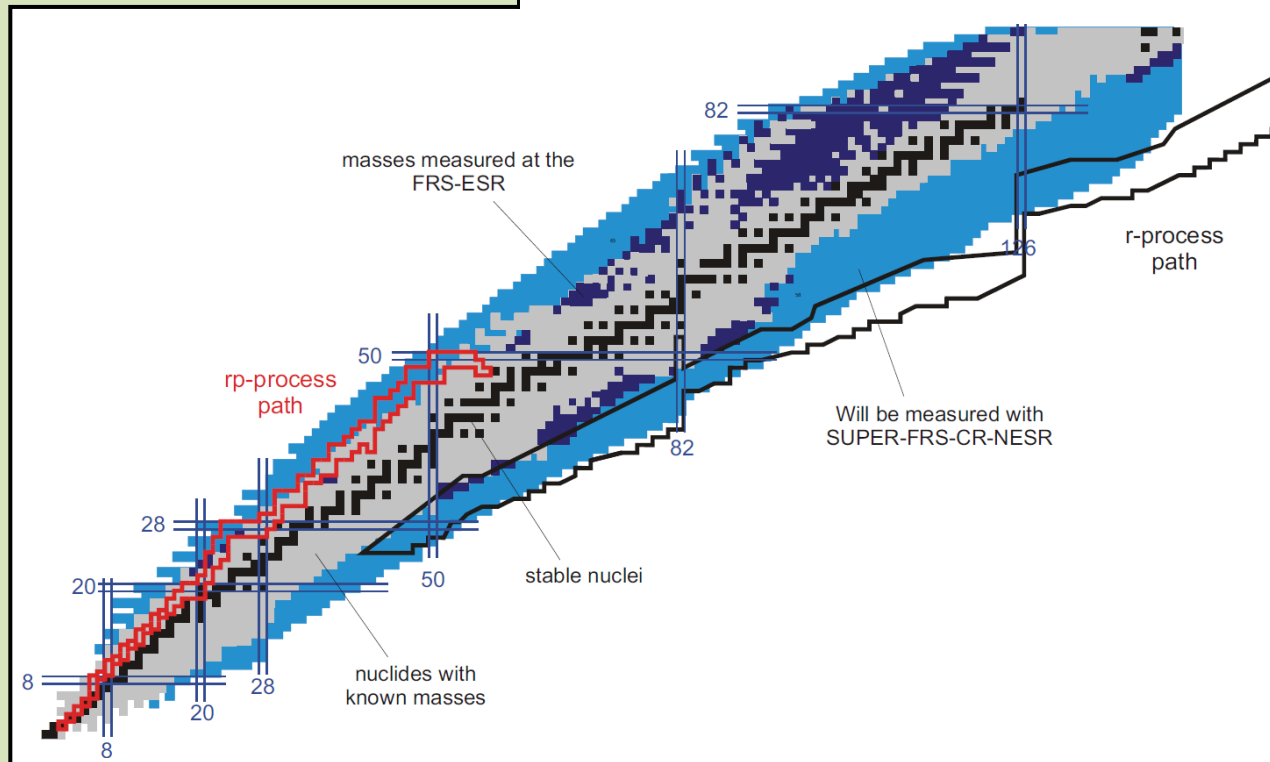
²GSI Helmholtzzentrum für Schwerionenforschung, Germany

³Departamento de Física Teórica, Universidad Autónoma de Madrid, Spain



- Astrophysical models require as an input thousands of **nuclear masses** → beyond experimental reach

most relevant input are the extracted
→ **neutron-separation energies S_n**
→ **beta-decay energies Q_β**
which determine thresholds of all nuclear reactions



- Astrophysical models require as an input thousands of **nuclear masses** beyond experimental reach
- Need accurate predictions from theoretical **global mass models** →

→ Modern mass market:

Finite Range Droplet Model (FRDM)

rms = 0.57 MeV

**Extended Thomas-Fermi
+ Strutinsky Integral (ETFSI)**

0.69 MeV

Duflo-Zuker (DZ)

0.36 MeV

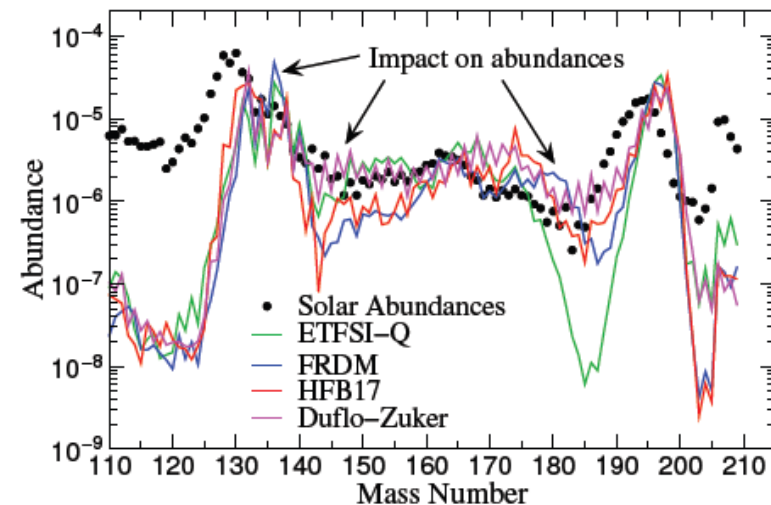
Weizsäcker-Skyrme (WS)

0.298 MeV

Hartree-Fock-Bogolyubov (HFB)

(Skyrme) 0.51 MeV

(Gogny) 0.798 MeV



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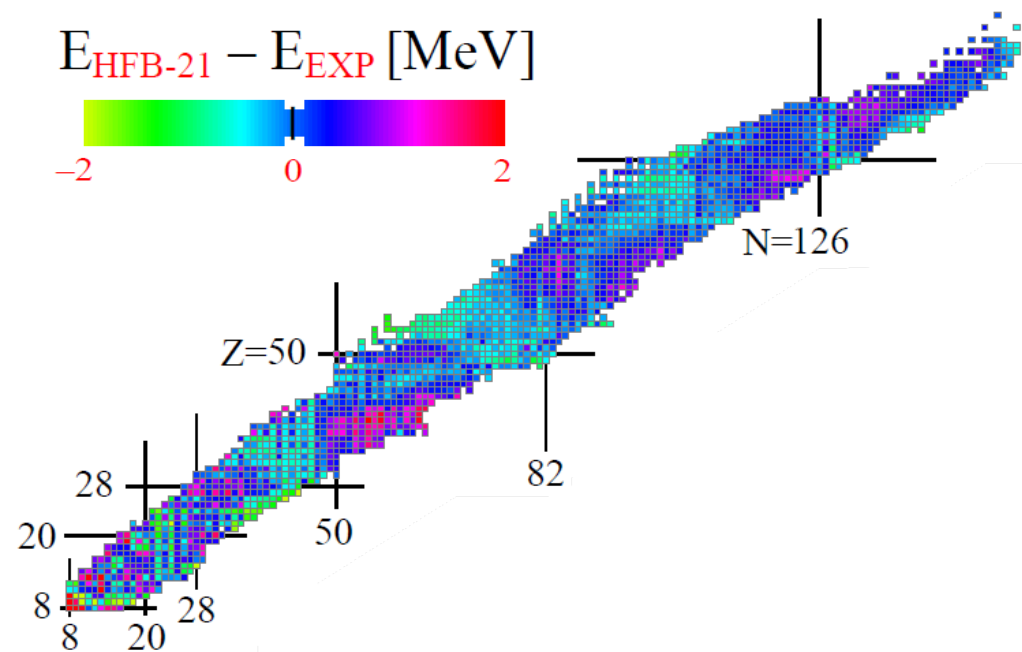
self-consistent mean field models

based on **Energy Density Functionals:**

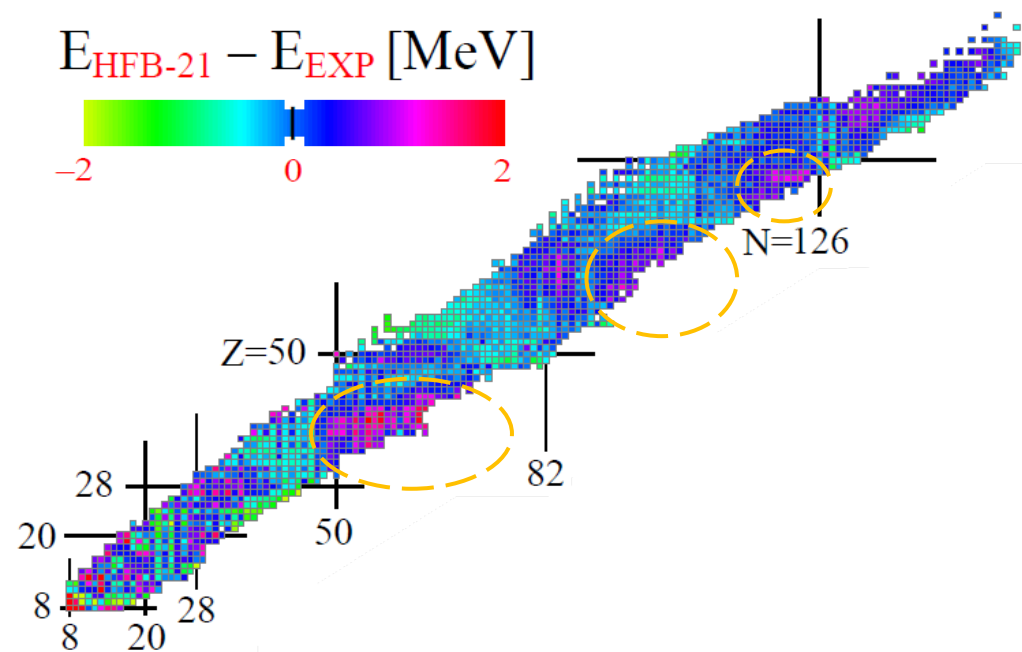
- **Skyrme HFB-*** (Goriely S. et al., *PRC88*, 2013)
- **Gogny D1S/D1M** (Goriely S. et al., *PRL102*, 2009)
- **UNEDF** (Erler J. et al., *Nature 486*, 2012)

more microscopic
↓

- Astrophysical models require as an input thousands of **nuclear masses** beyond experimental reach
- Need accurate predictions from theoretical **global mass models**
- All models have similar **rms**, but more fundamental **HFB models** should provide greater confidence in describing unknown isotopes
- **Still some problems with HFB models:**
 - I) Issue of **convergence** due to truncated model space
 - II) Missing some physics without **Beyond-Mean-Field** correlations
 - III) **Odd-mass nuclei** are not treated on the same footing



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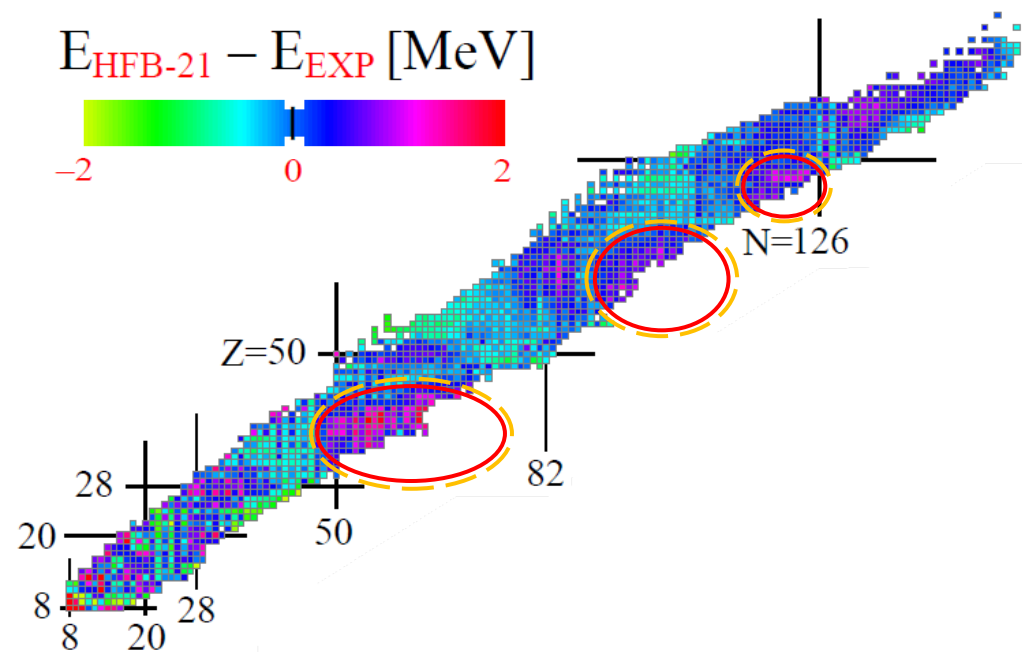
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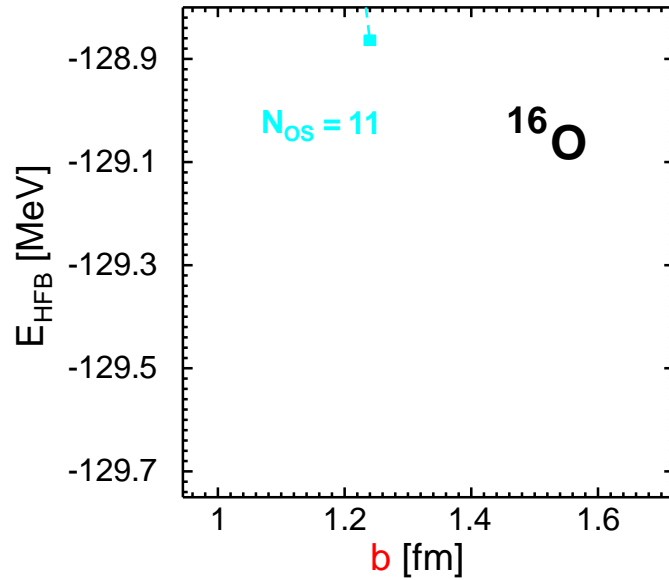
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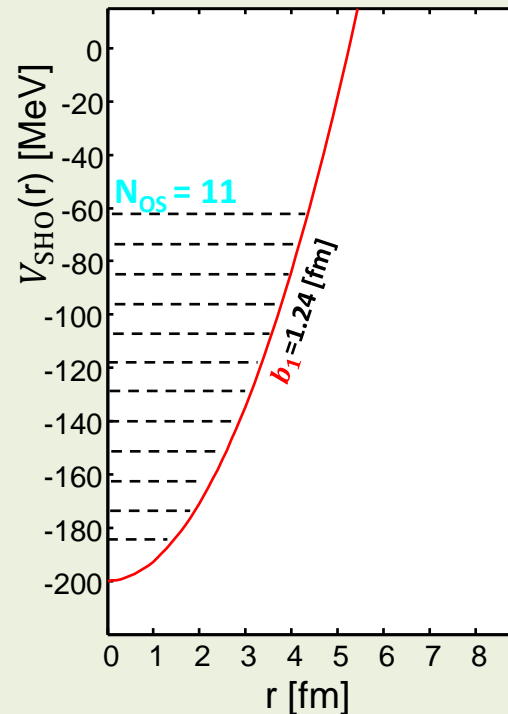


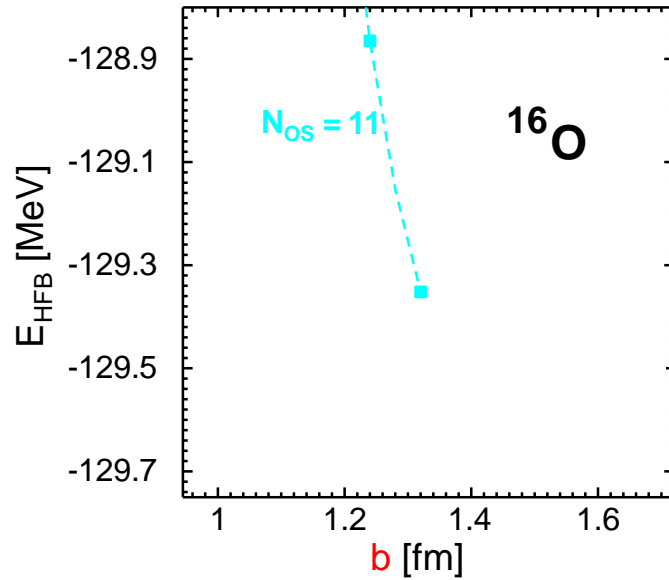
Convergence of masses and IR-Extrapolation to infinite basis



Convergence in finite oscillator space

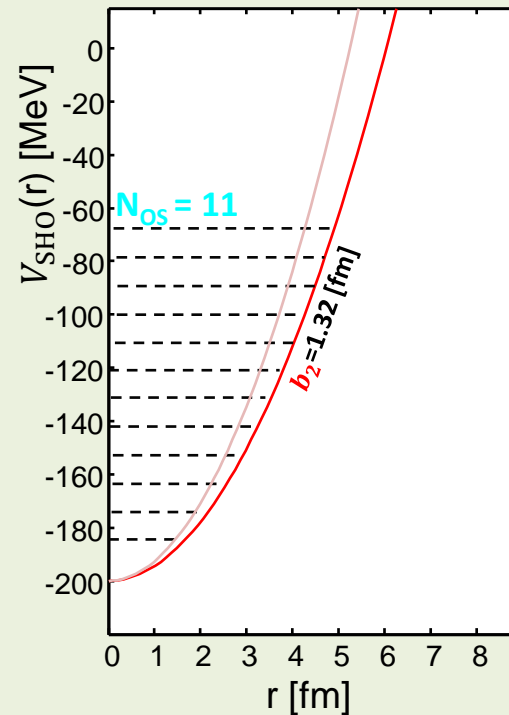
- Calculations are usually performed in finite spherical harmonic oscillator (SHO) basis with two parameters that define it:
 - N_{os} - number of major oscillator shells
 - b - length of SHO wavefunctions

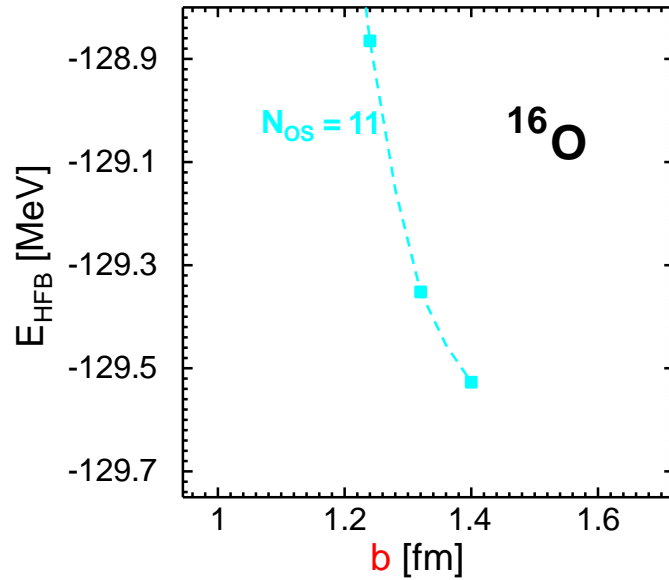




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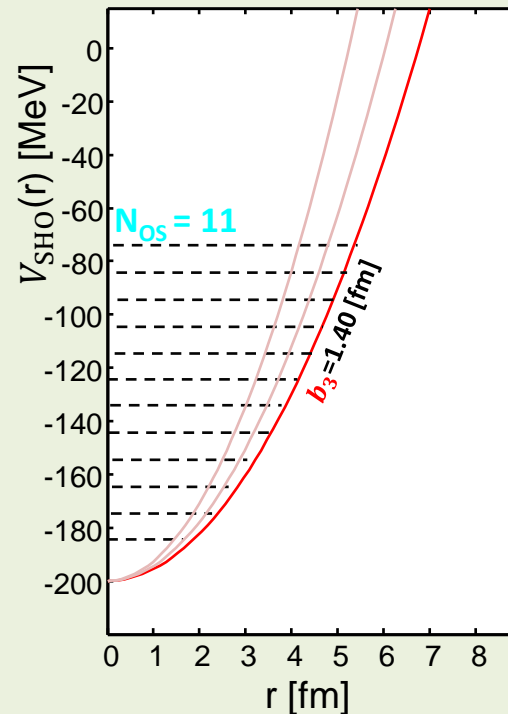
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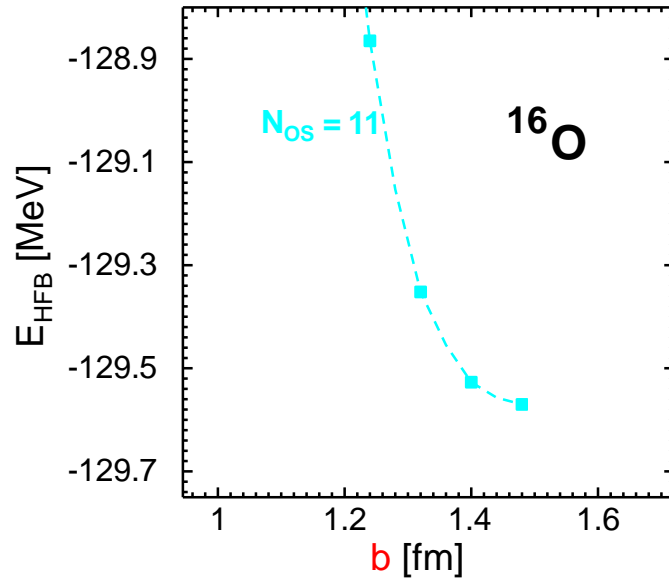




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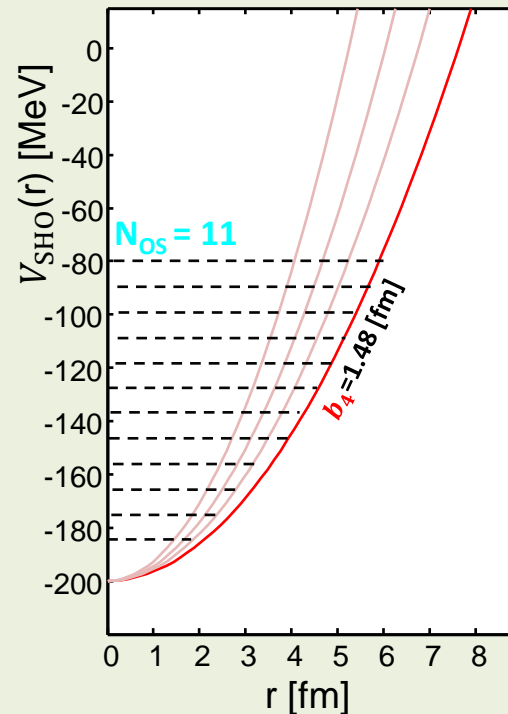
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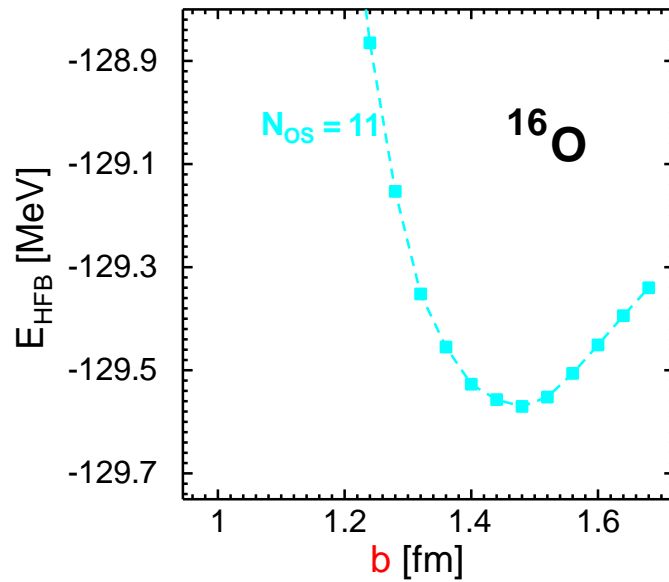




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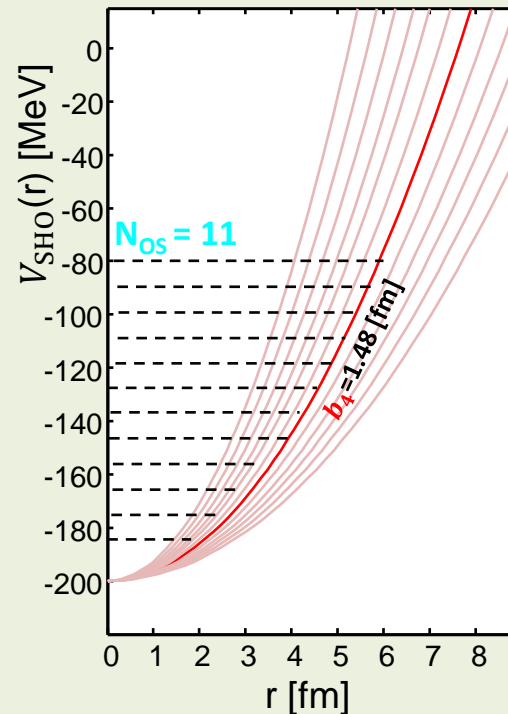
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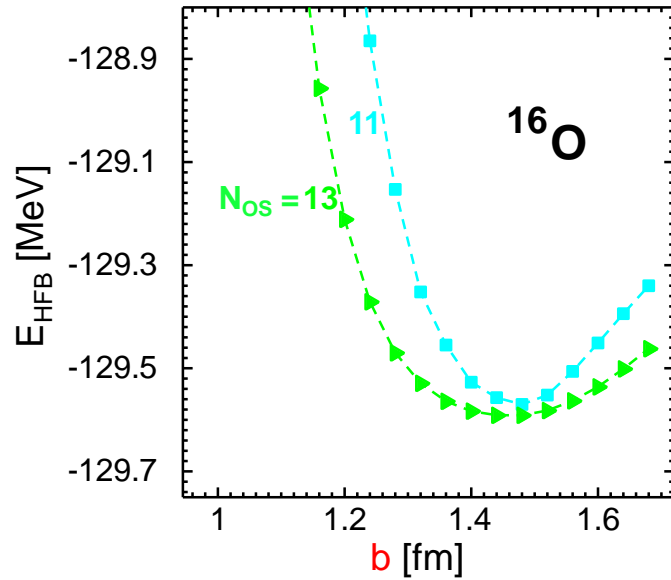




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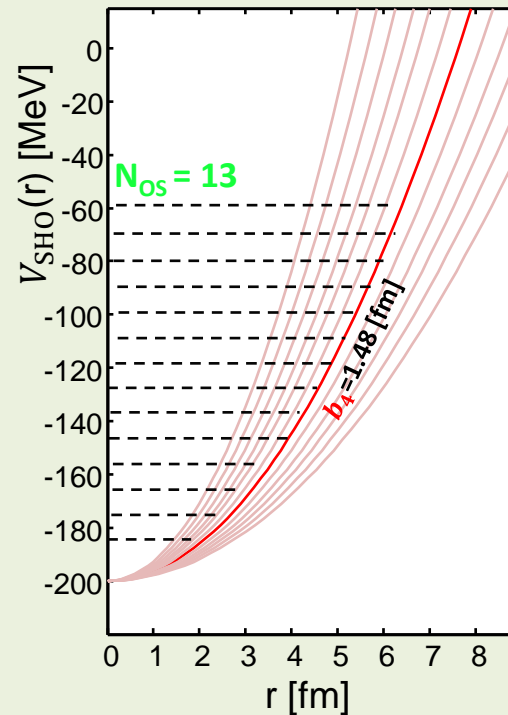


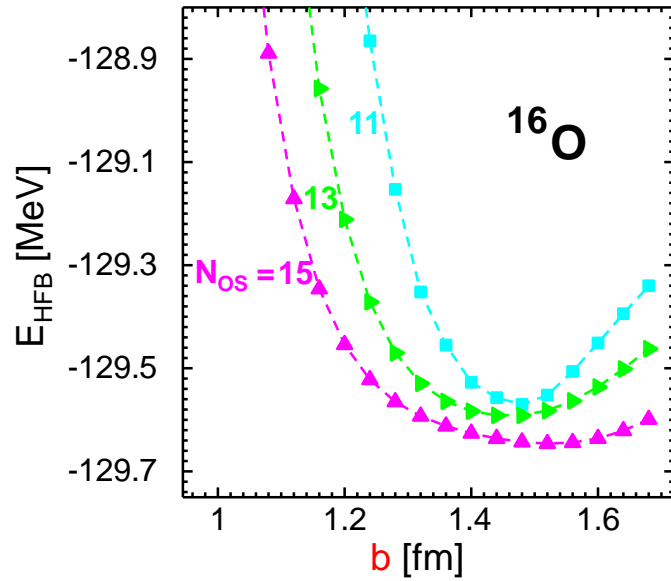


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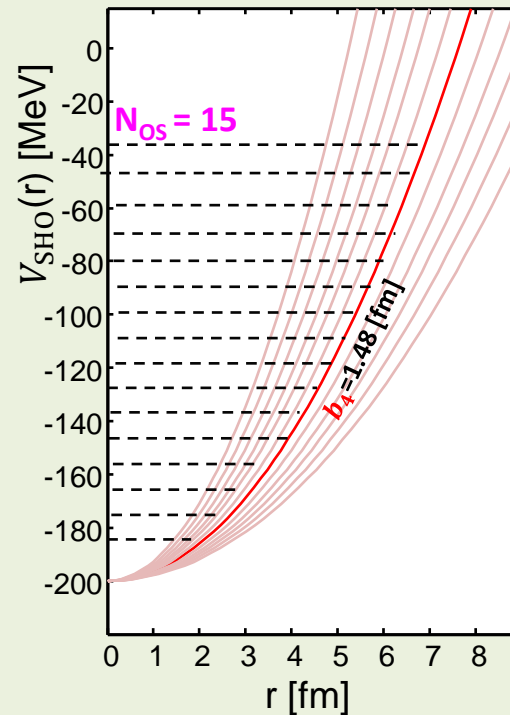
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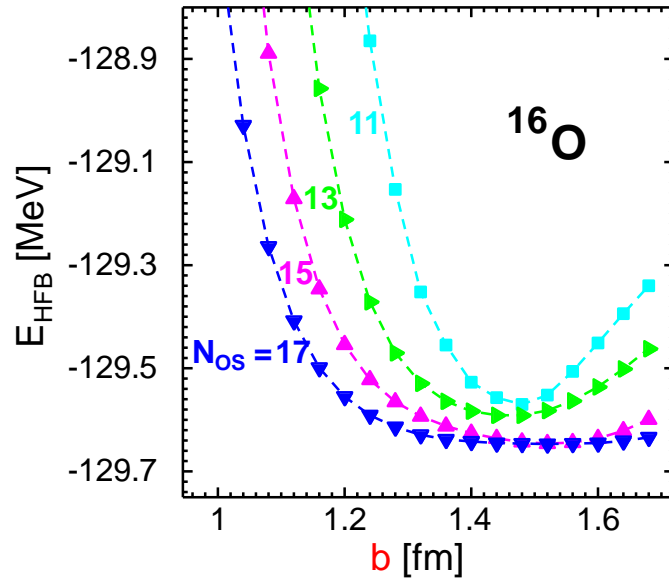




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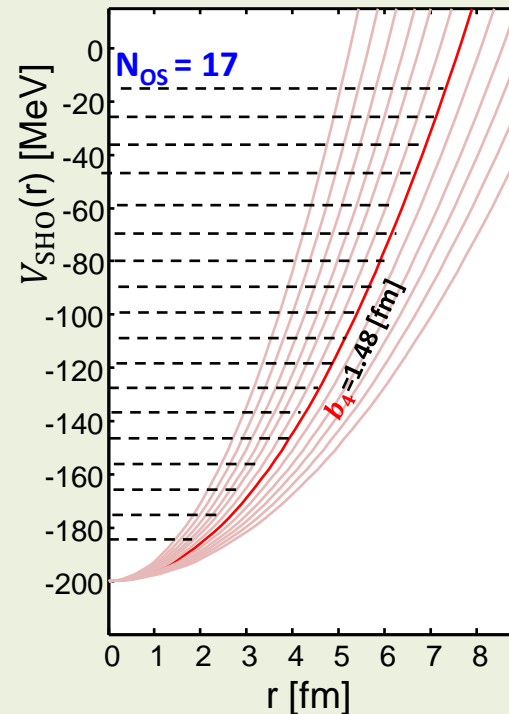


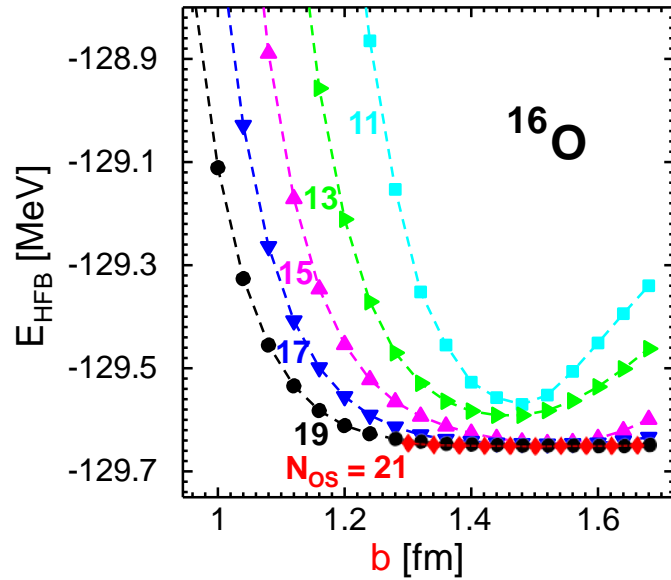
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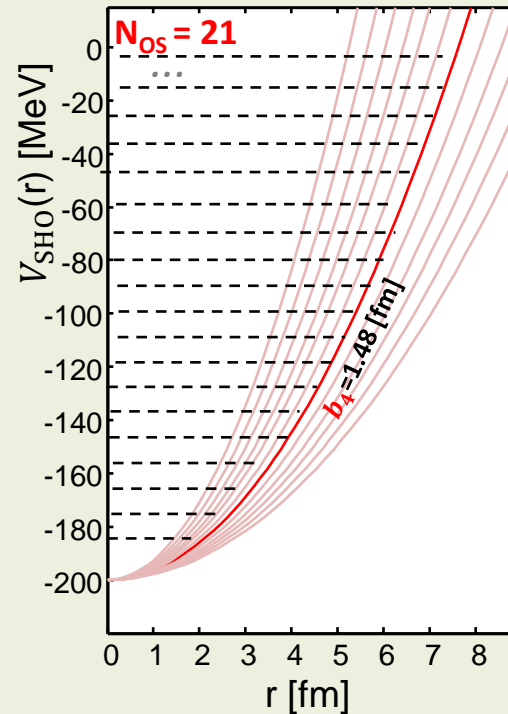


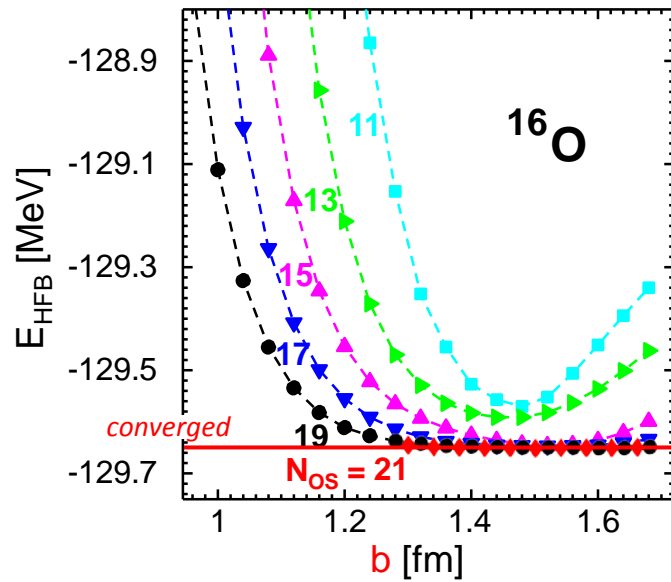
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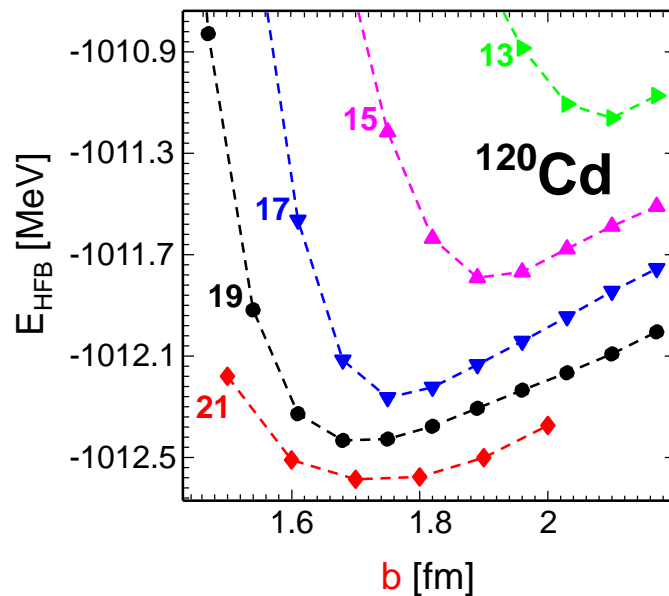
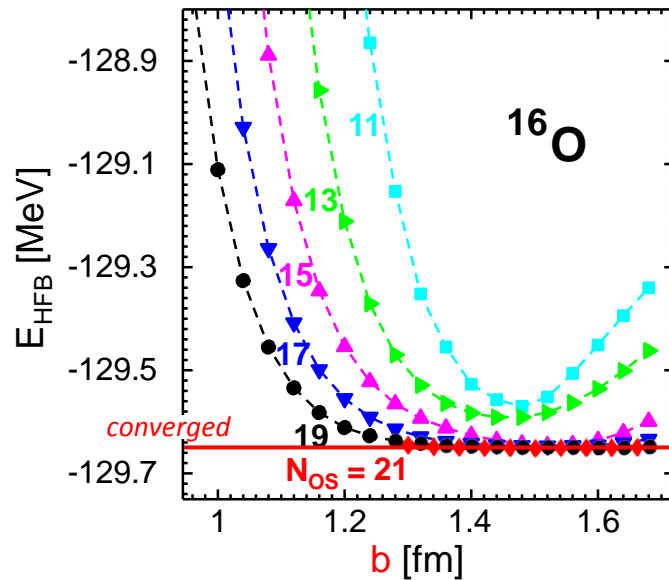
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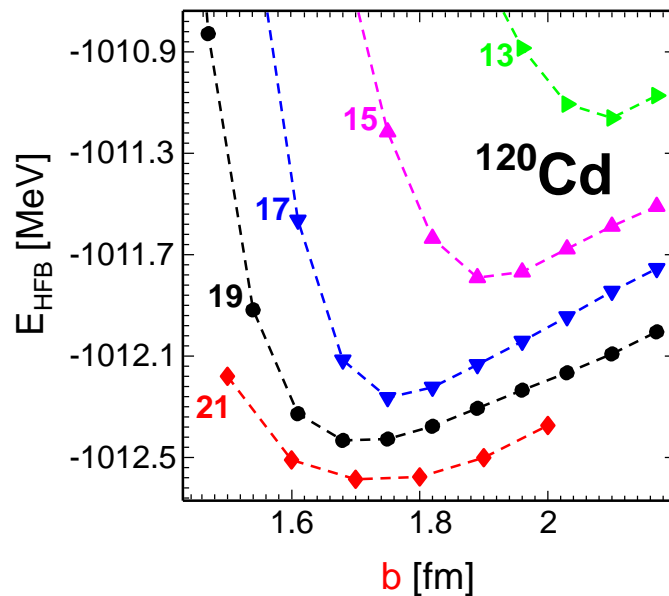
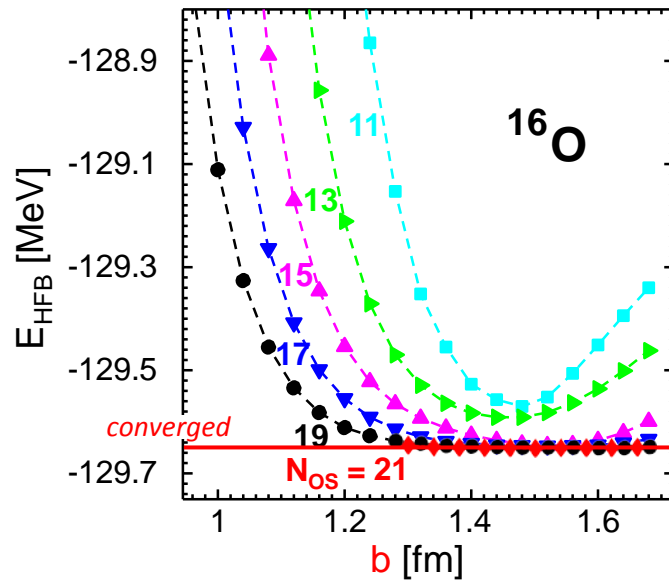
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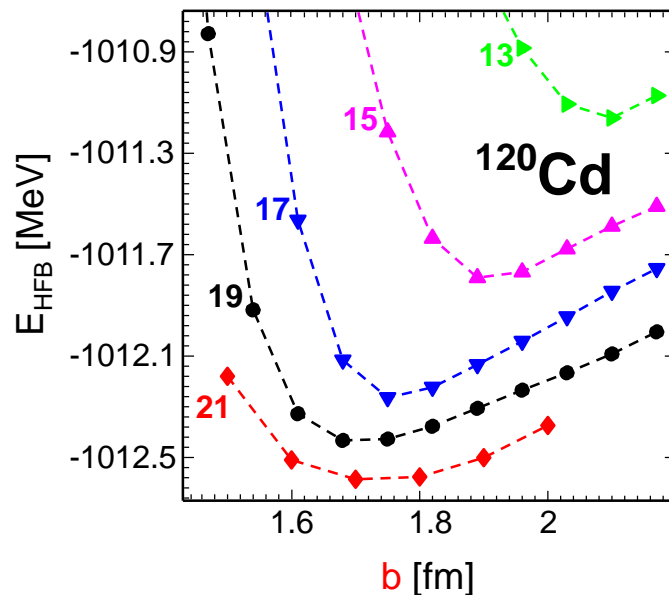
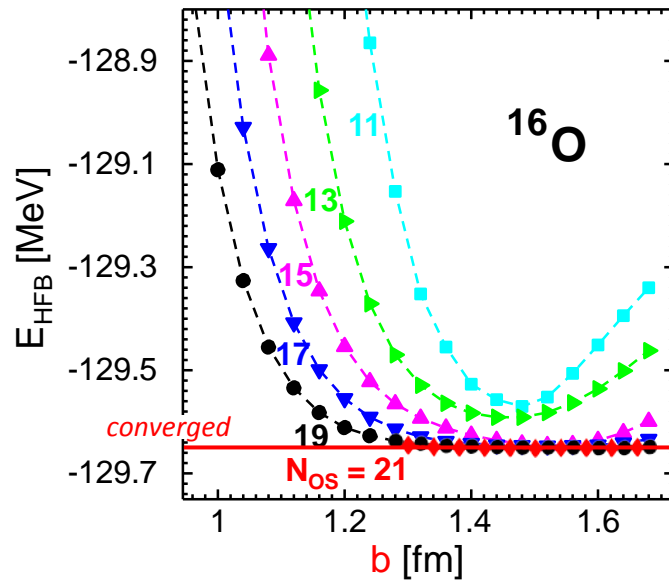
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- Recently new **IR-extrapolation** scheme with firm theoretical background developed

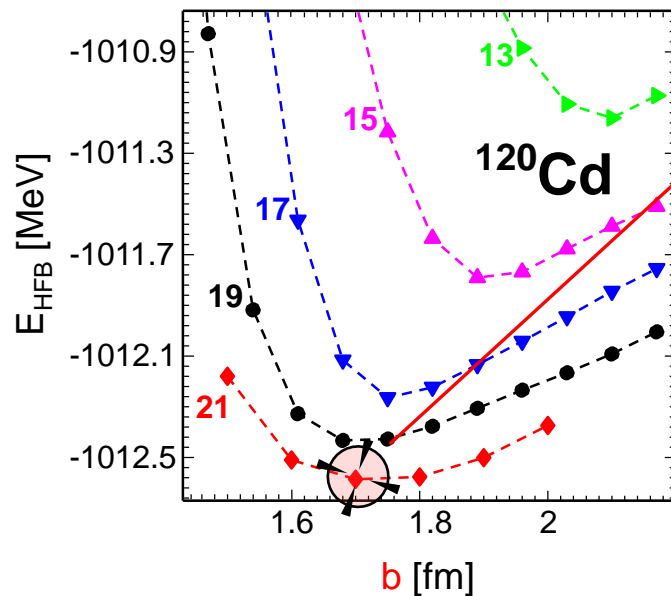
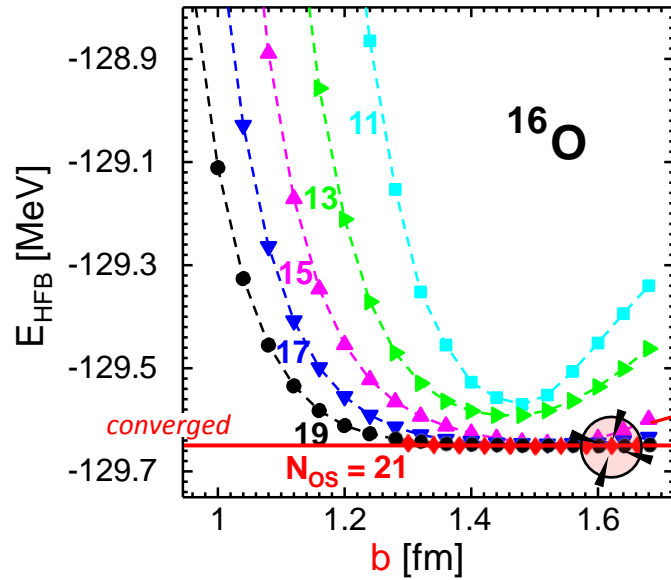
Furnstahl R.J., Hagen G., Papenbrock T., PRC86, 031301 (2012)

More S.N. et al., PRC87, 044326 (2013)

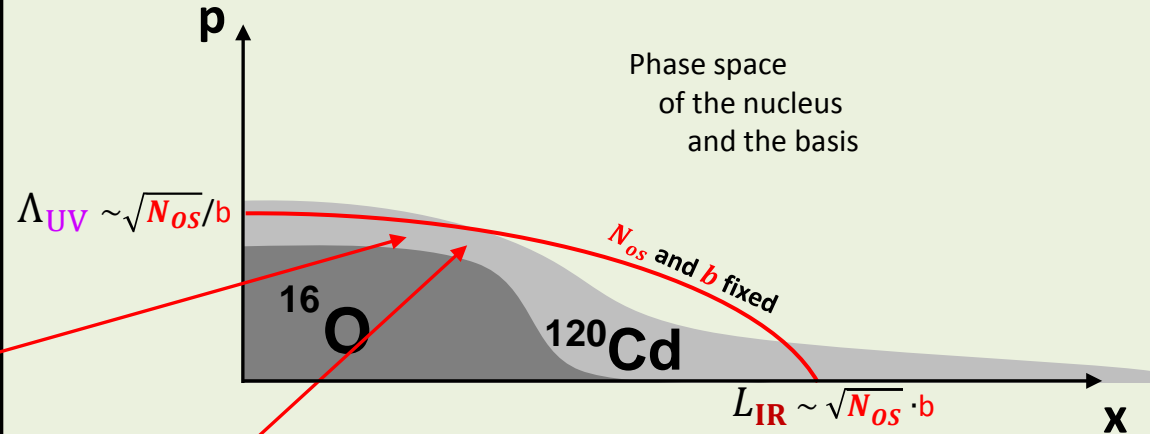
Furnstahl R.J., More S.N., Papenbrock T., PRC89, 044301 (2014)

Furnstahl R.J. et al, arXiv:1408.0252 (2014)

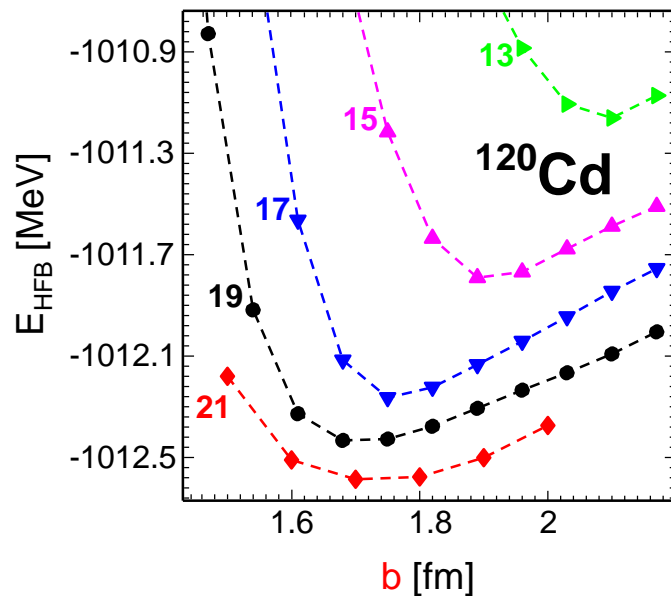
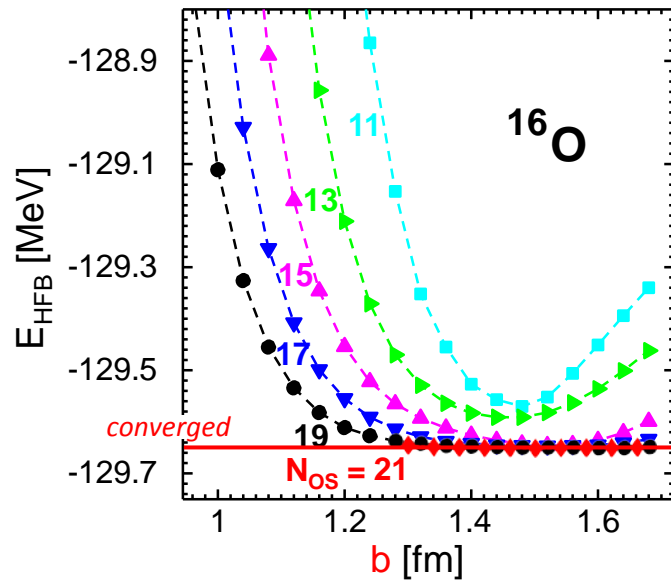
However, have not yet been systematically tested on whole isotopic chains.



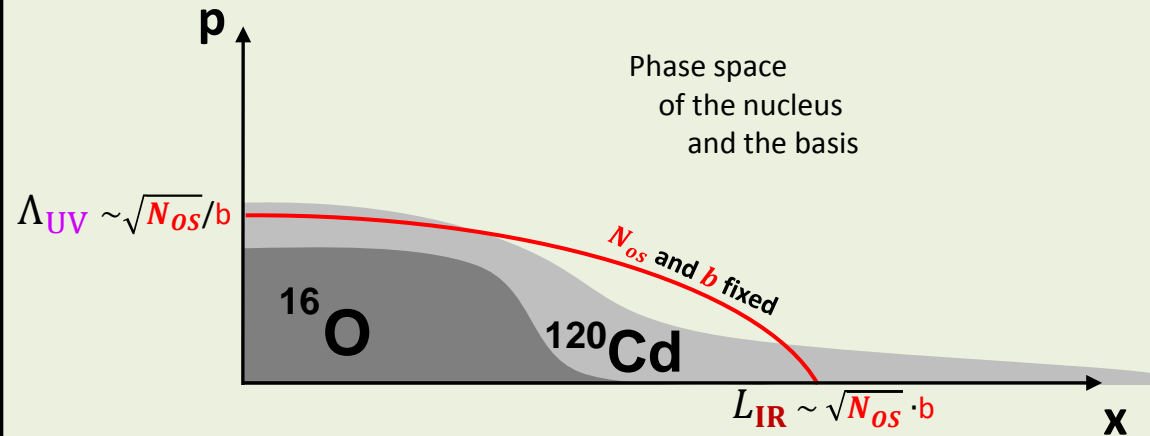
IR-Extrapolation to infinite basis



- Truncating working basis, we effectively impose
 - > **x**: a hard-wall L_{IR} cutoff
 - > **p**: analogous sharp Λ_{UV} cutoff



IR-Extrapolation to infinite basis



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- **Full Convergence** - when nucleus "fits" into SHO basis:

IR convergence:

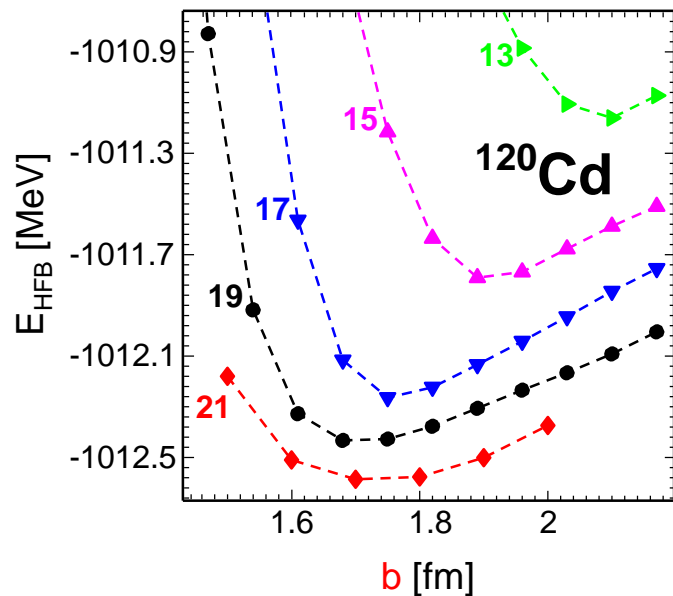
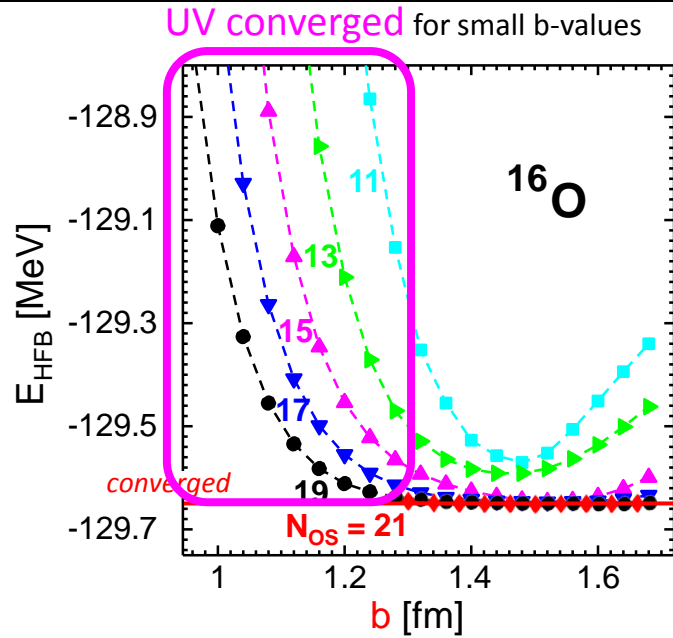
Spatial extent of the nucleus

$$r < L_{IR}$$

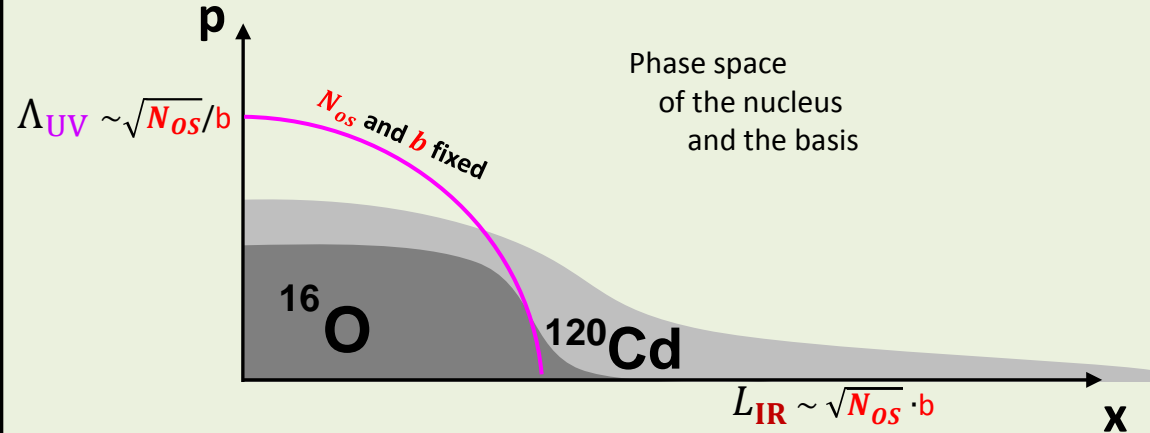
UV convergence:

Largest mom. scale of interaction

$$\lambda < \Lambda_{UV}$$



IR-Extrapolation to infinite basis



- Truncating working basis, we effectively impose
 - > x : a hard-wall L_{IR} cutoff
 - > p : analogous sharp Λ_{UV} cutoff

- IR-Extrapolation** - binding energy correction in the limit of UV converged results!

\times IR convergence:

exponential $r < L_{\text{IR}}$

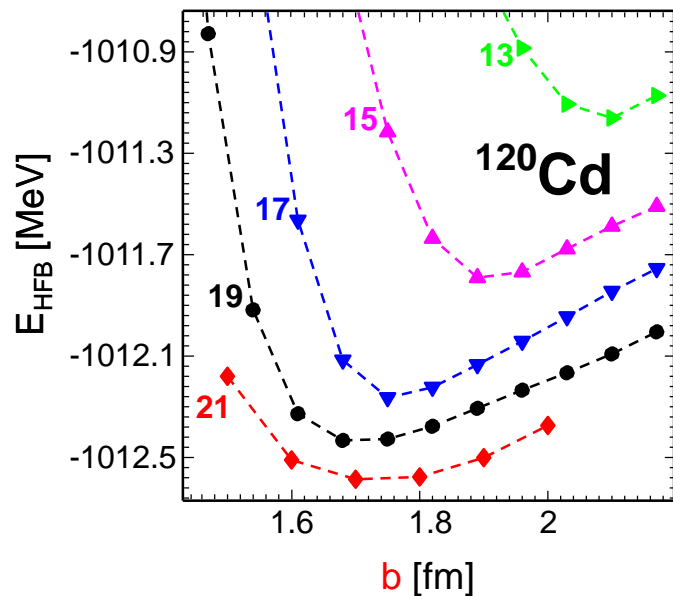
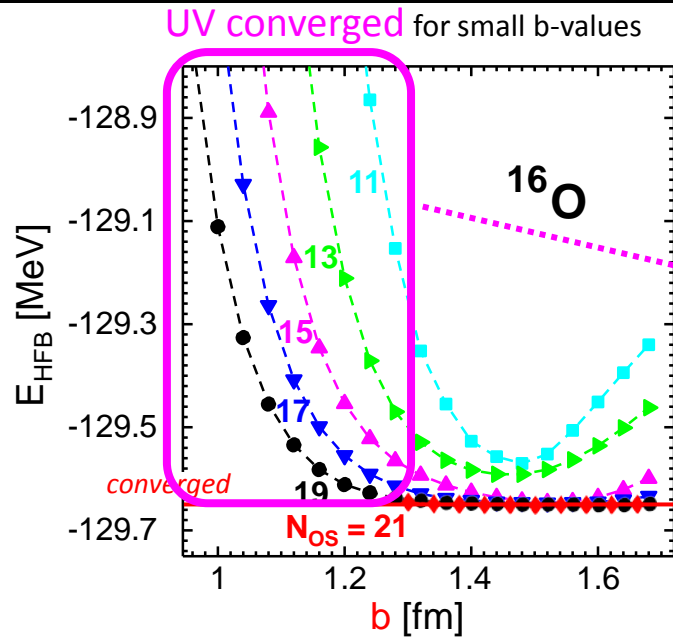
\checkmark UV convergence:

Gaussian Largest mom. scale of interaction $\lambda < \Lambda_{\text{UV}}$

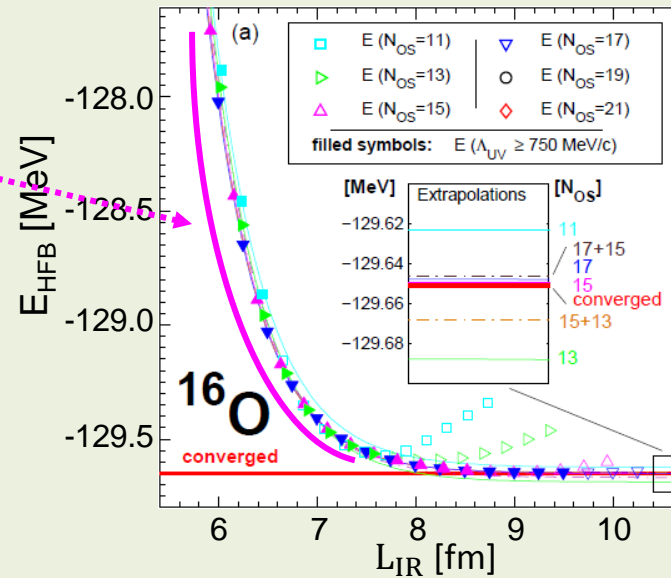
$$E_{\text{HFB}}(L_{\text{IR}}) = a_0 e^{-2k_\infty L_{\text{IR}}} + E_\infty$$

where a_0 , k_∞ and E_∞ are fit constants, and

$L_{\text{IR}} =$ depends on basis and nucleus



IR-Extrapolation to infinite basis



- **IR-Extrapolation** - binding energy correction in the limit of UV converged results!

✗ **IR convergence:** Spatial extent of the nucleus $r < L_{\text{IR}}$

✓ **UV convergence:** Largest mom. scale of interaction $\lambda < \Lambda_{\text{UV}}$

$$E_{\text{HFB}}(L_{\text{IR}}) = a_0 e^{-2k_\infty L_{\text{IR}}} + E_\infty$$

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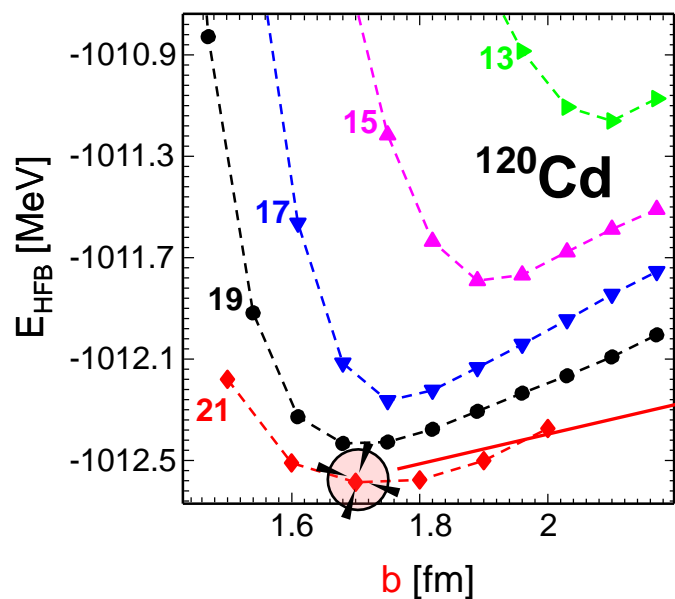
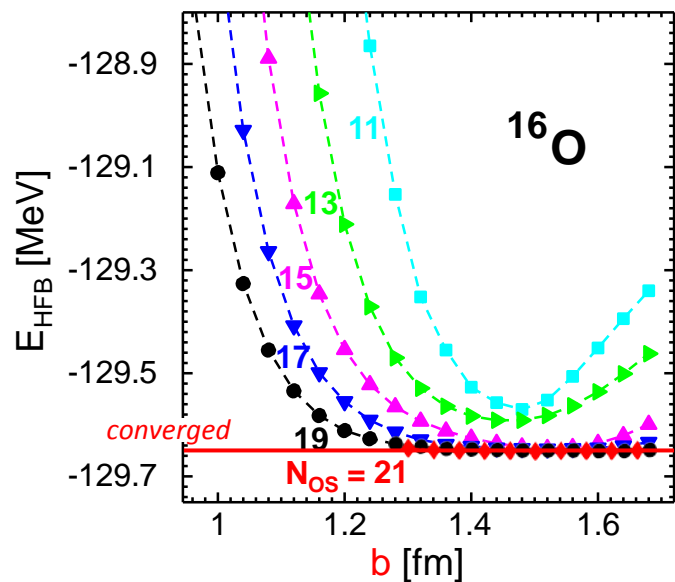
Convergence analysis

Convergence analysis
and Beyond-Mean-Field correlations

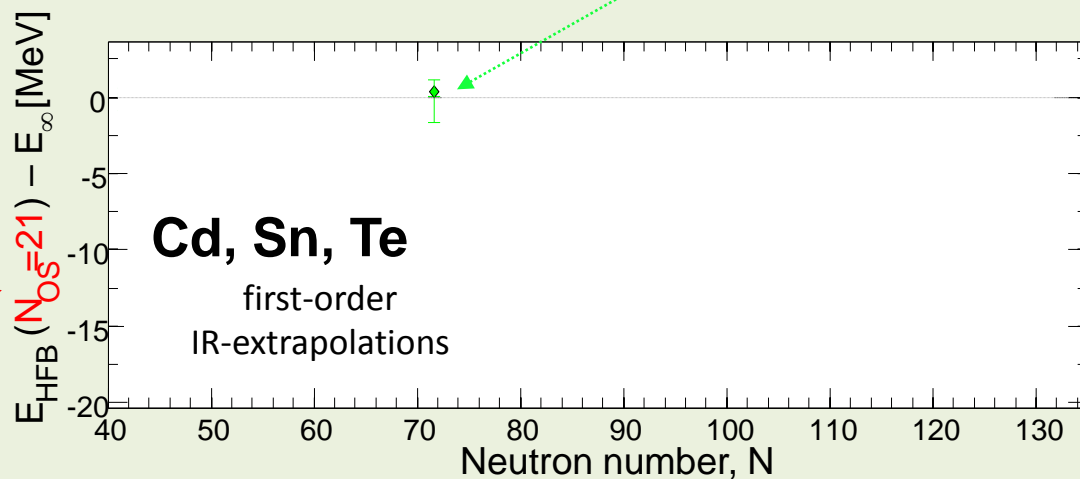
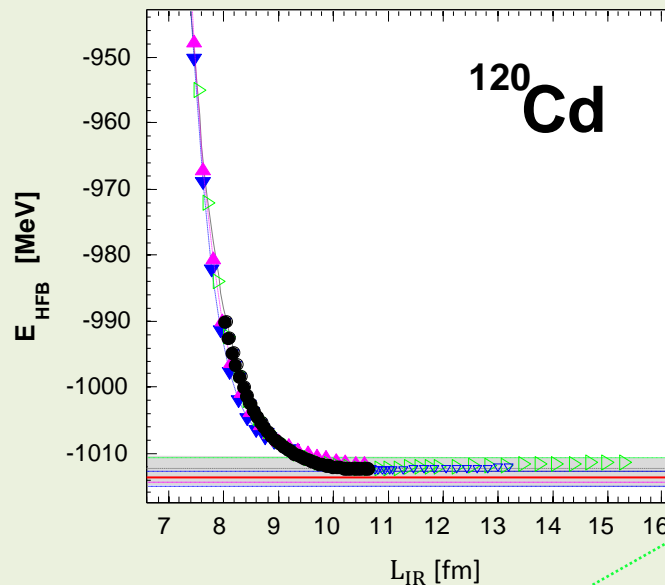
NAVI Physics Days 2015-02-26

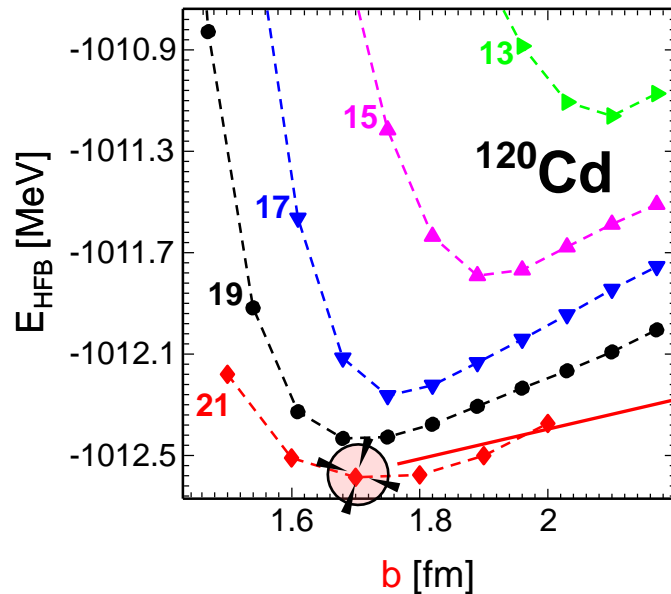
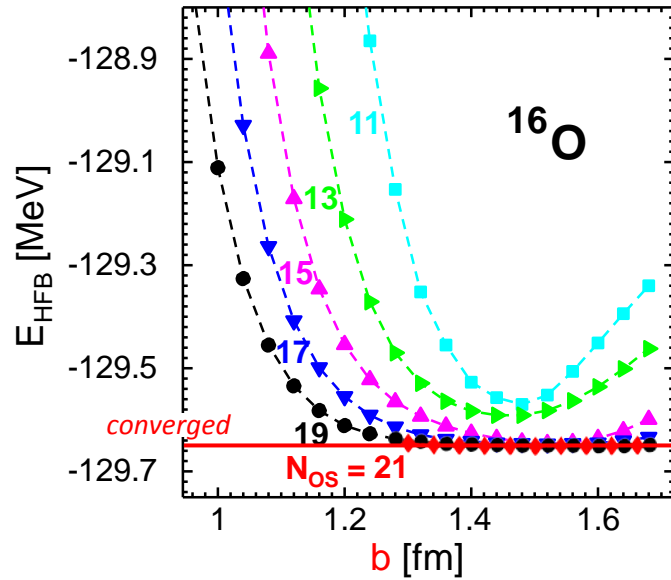
Alexander Arzhanov

$\tau = 27$



IR-Extrapolation to infinite basis

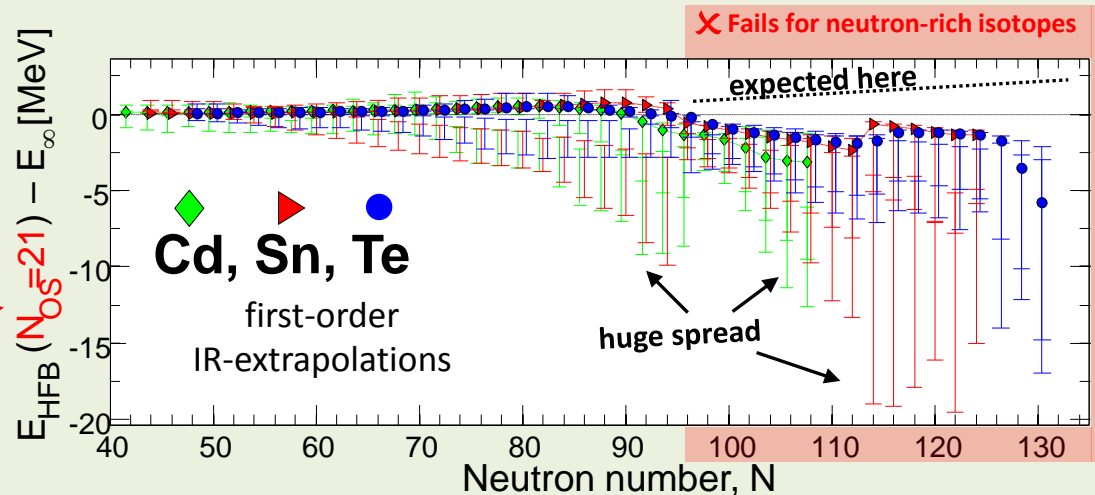




IR-Extrapolation to infinite basis

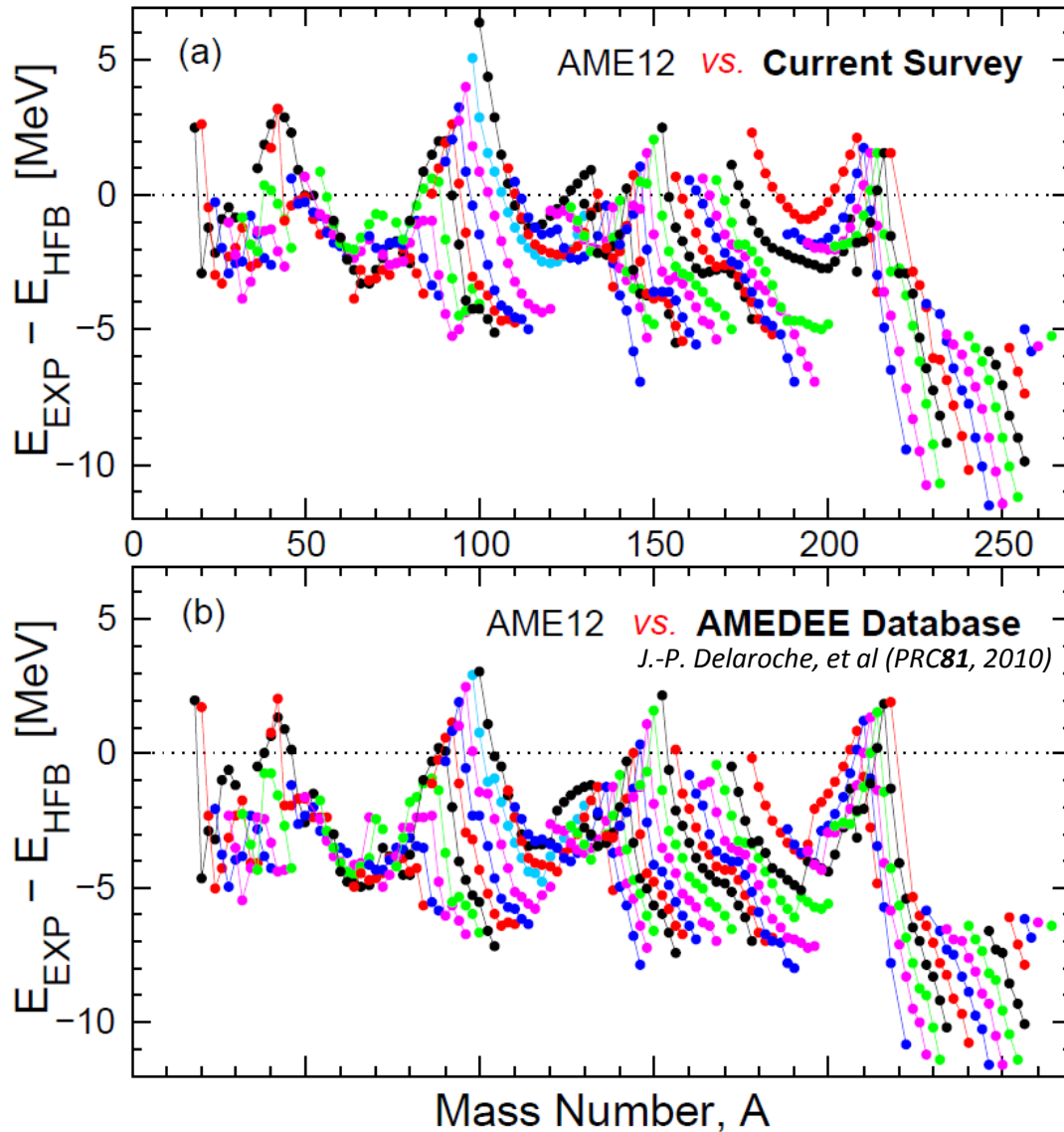
⇒ **Second-order IR-corrections for loosely bound nuclei?**
Preliminary checks unsatisfactory.
Still not adapted for atomic nuclei.

⇒ At present we do not have a reliable and universal extrapolation method for binding energies to the limit of an infinite basis for HFB-based models!

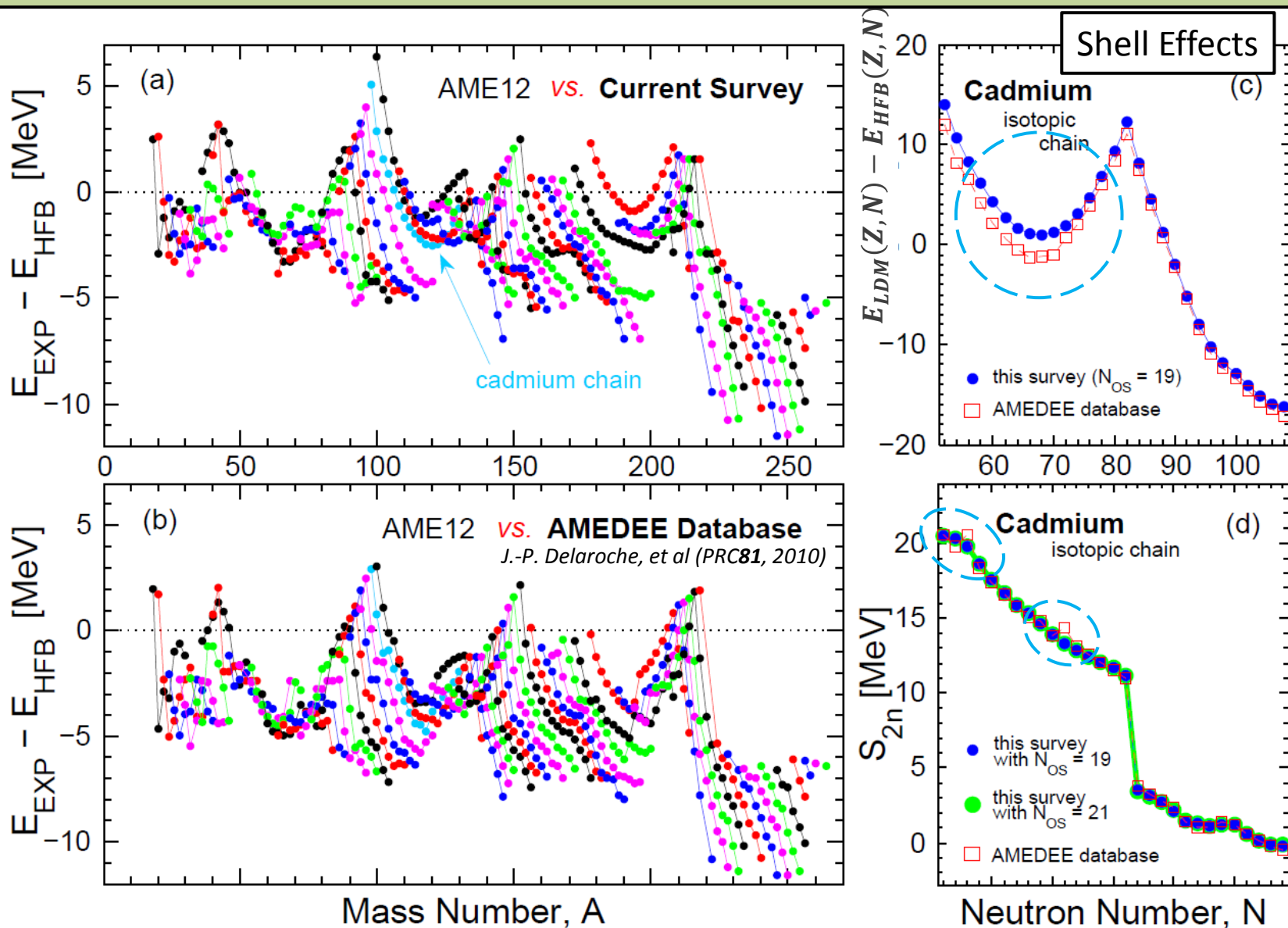


Large-scale HFB calculation and Beyond-Mean-Field corrections

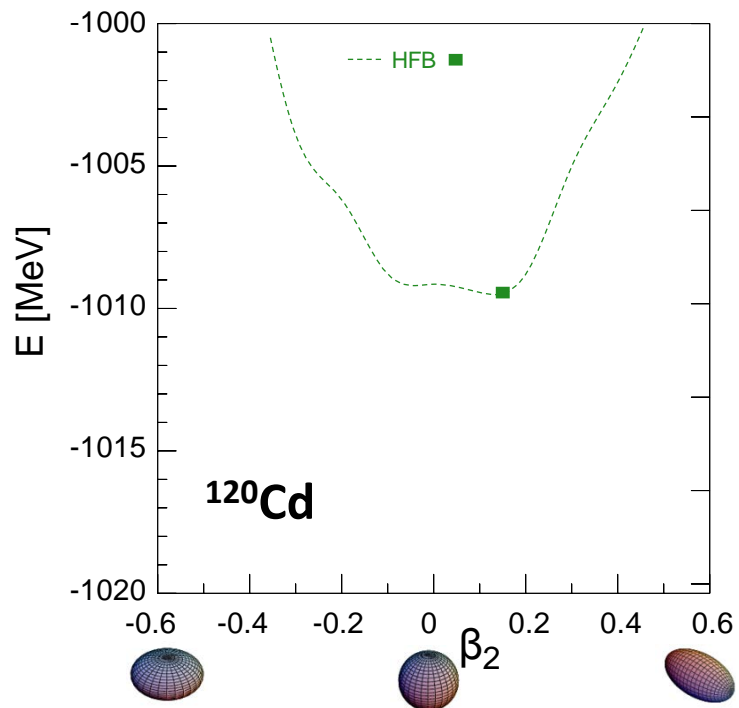
global mass surveys for *axially deformed Mean Field HFB-D1S* calculation for *even-even nuclei*



global mass surveys for axially deformed Mean Field HFB-D1S calculation for even-even nuclei



- **Mean Field approach of HFB:**
 - no symmetry conservations
 - no configuration mixing



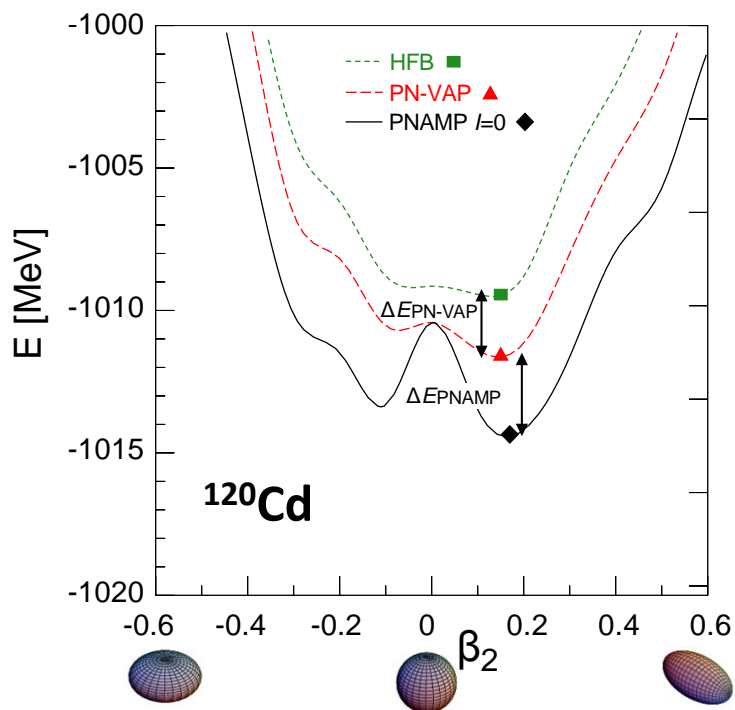
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Digging Beyond the Mean Field

Symmetry restoration by

- **Variation After Particle Number Projection (PN-VAP):**
 $\Delta E_{\text{PN-VAP}} \sim 2.3 \text{ MeV}$
- **Particle Number and J = 0 Angular Momentum Projection (PNAMP):**
 $\Delta E_{\text{PNAMP}} \sim 2.7 \text{ MeV}$



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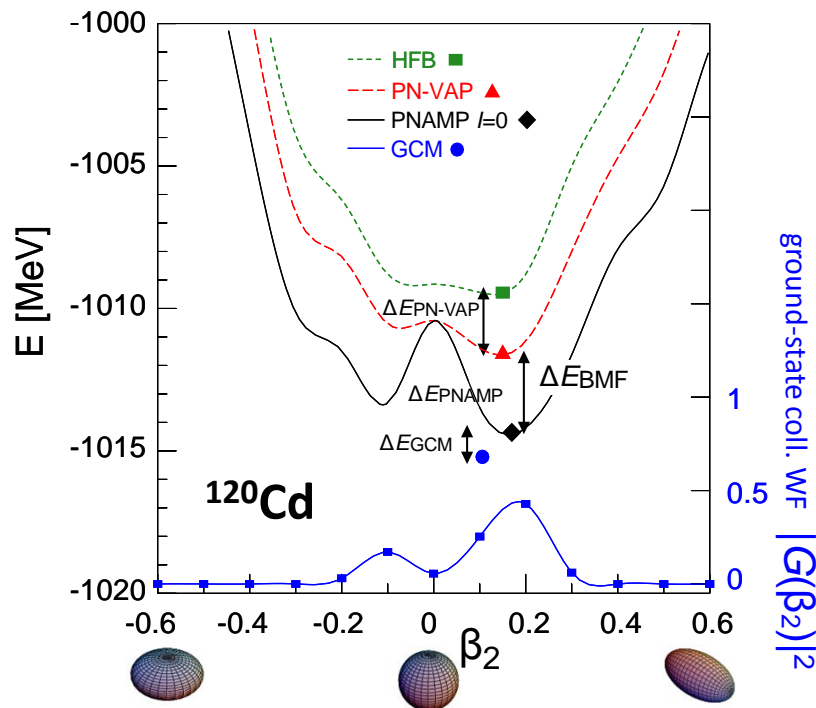
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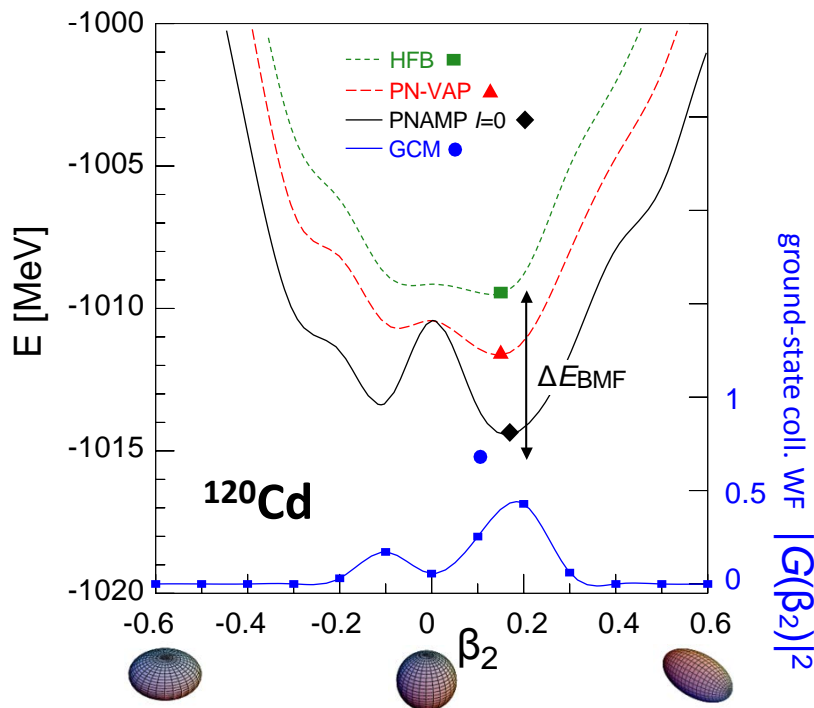
- Exact implementation of **Generator Coordinate Method (GCM):**
 $\Delta E_{\text{GCM}} \sim 0.8 \text{ MeV}$



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Configuration mixing by

- Exact implementation of **Generator Coordinate Method (GCM):**
 $\Delta E_{\text{GCM}} \sim 0.8 \text{ MeV}$

Total Energy with BMF correlations

$$E_{\text{GCM}} = E_{\text{HFB}}(N_{\text{OS}} = 19) - \Delta E_{\text{BMF}}$$

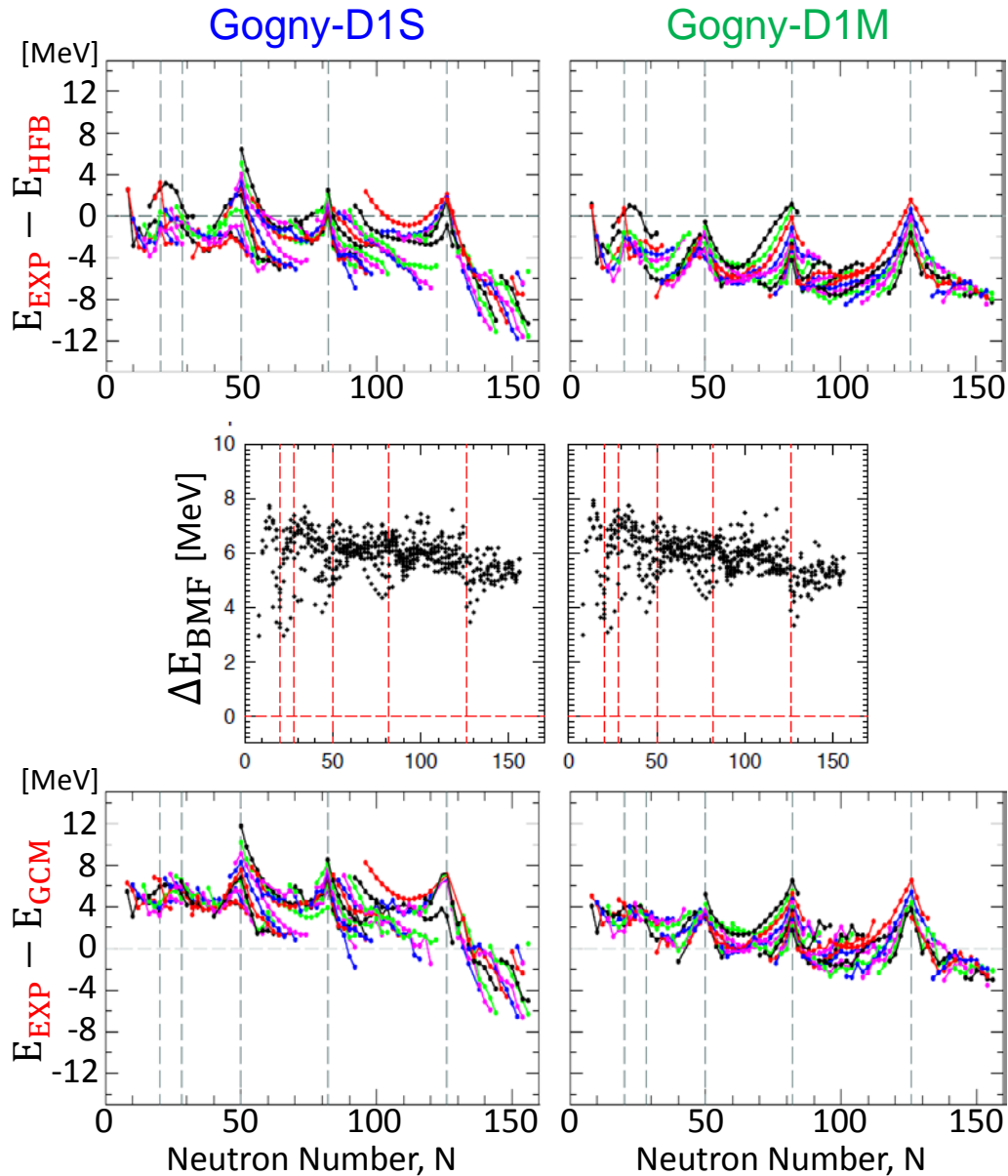
where the BMF correlations are calculated as

$$\Delta E_{\text{BMF}} = E_{\text{HFB}}(N_{\text{OS}} = 11) - E_{\text{BMF}}(N_{\text{OS}} = 11)$$

because of heavy computational burden

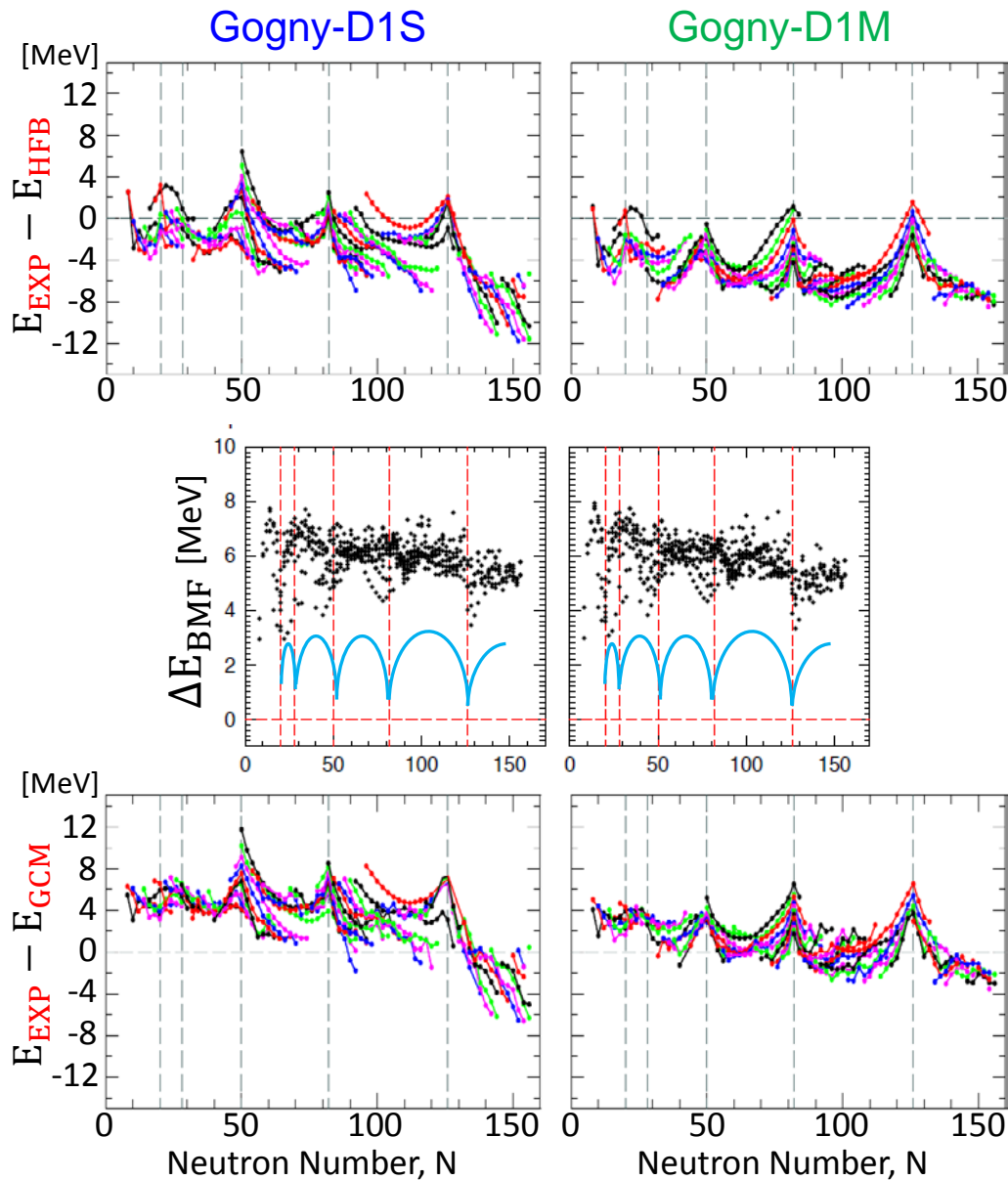
$$t_{\text{BMF}}(N_{\text{OS}} = 11) \approx 60\text{h} \quad t_{\text{BMF}}(N_{\text{OS}} = 19) > 1000\text{h}$$

BMF correlations



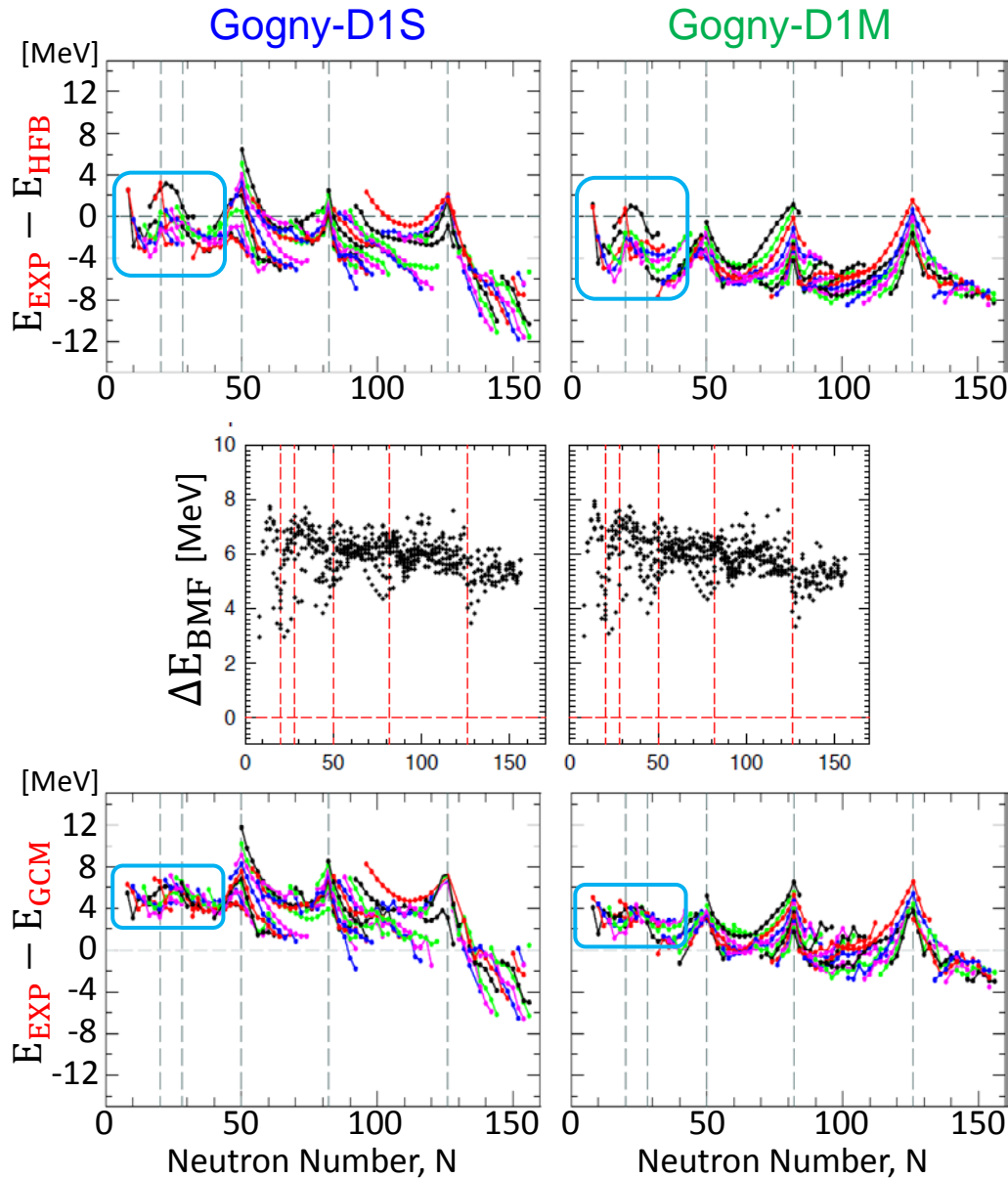
- Similar behavior of ΔE_{BMF} – corrections for both Gogny functionals **D1S** and **D1M** with $\Delta E_{\text{BMF}} \sim 5.8 \text{ MeV}$

BMF correlations



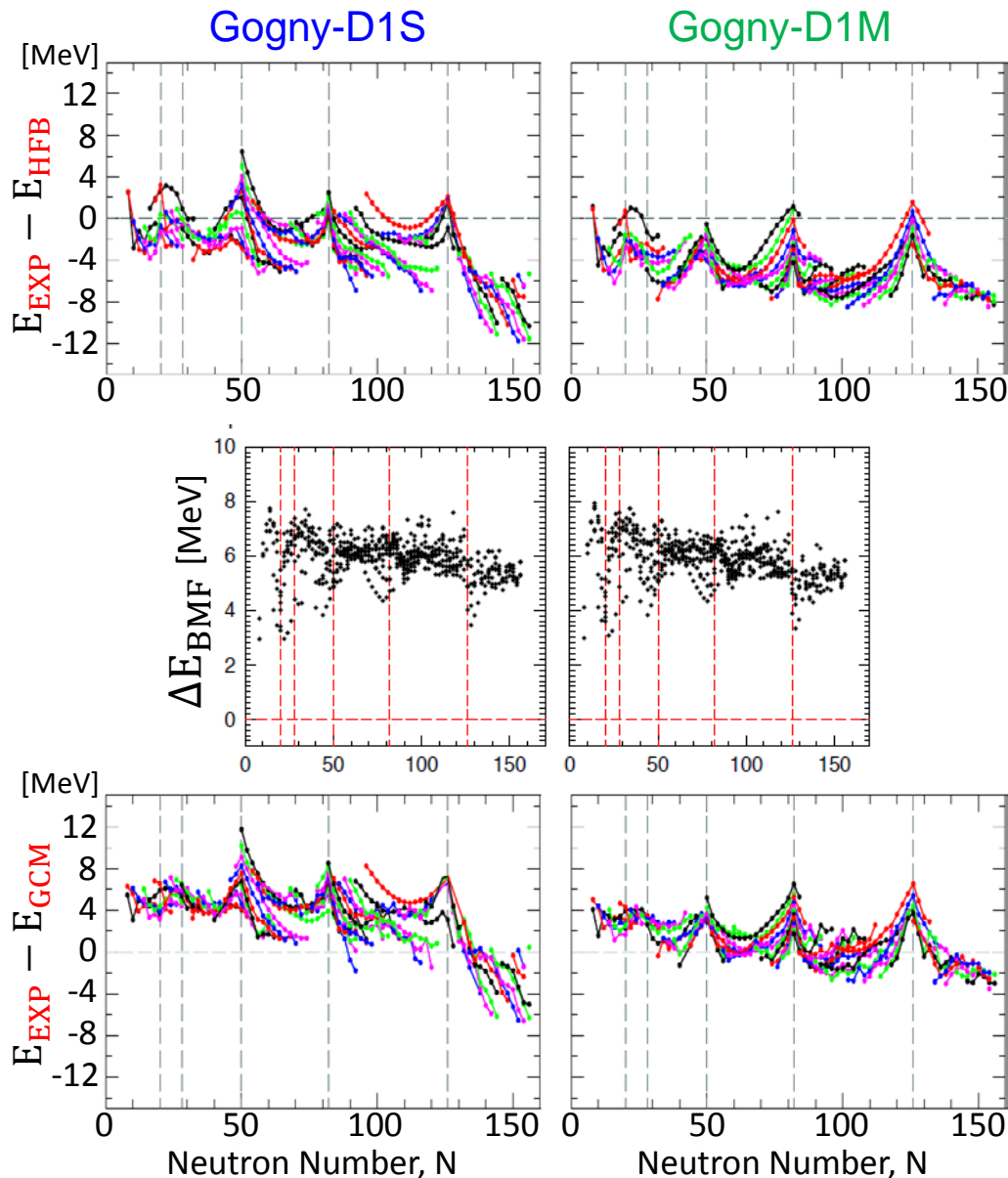
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- Inverse parabolic ΔE_{BMF} –corrections between shell closures tend to reduce the peaks at magic numbers *slightly*
- ... but *strong* Shell Effects are *not* washed out by BMF corrections

BMF correlations



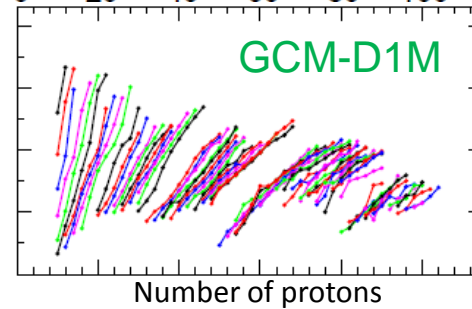
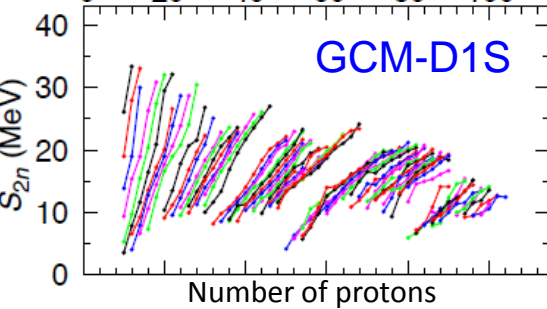
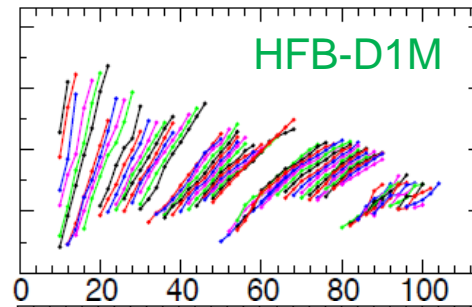
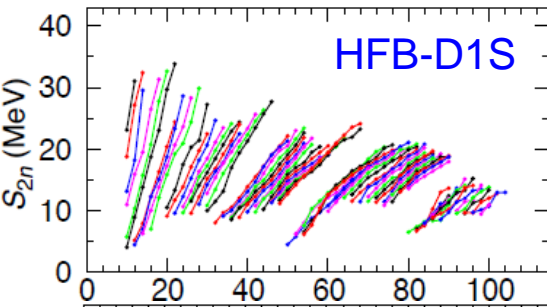
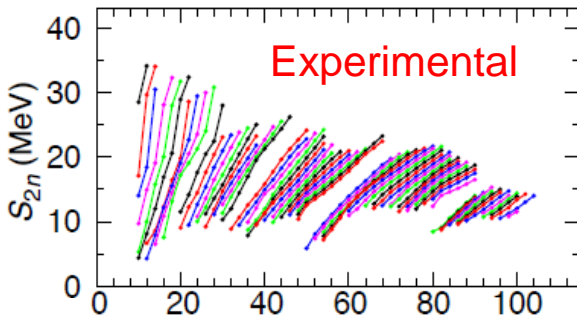
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BMF correlations



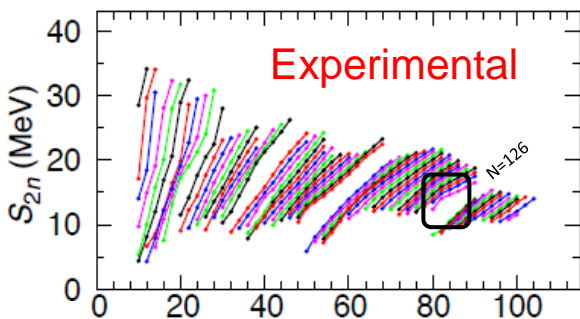
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- ... but *strong* Shell Effects are *not* washed out by BMF corrections
- Spread light nuclei ($N = 10 - 40$) is significantly reduced when BMF correlations are taken into account
- Overbinding for both **D1S** and **D1M** can be solved by re-fitting the EDF functional
- ... but it is still an open question whether re-fitting EDF functional with these and other BMF effects *self-consistently* can flatten the curves ?

BMF correlations



- Experimental S_{2n} are much smoother than both HFB and GCM results:
 - Convergence problem?
 - Missing triaxiality, octupolarity, etc.?

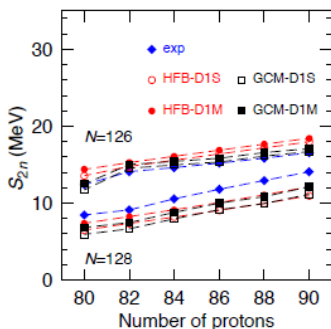
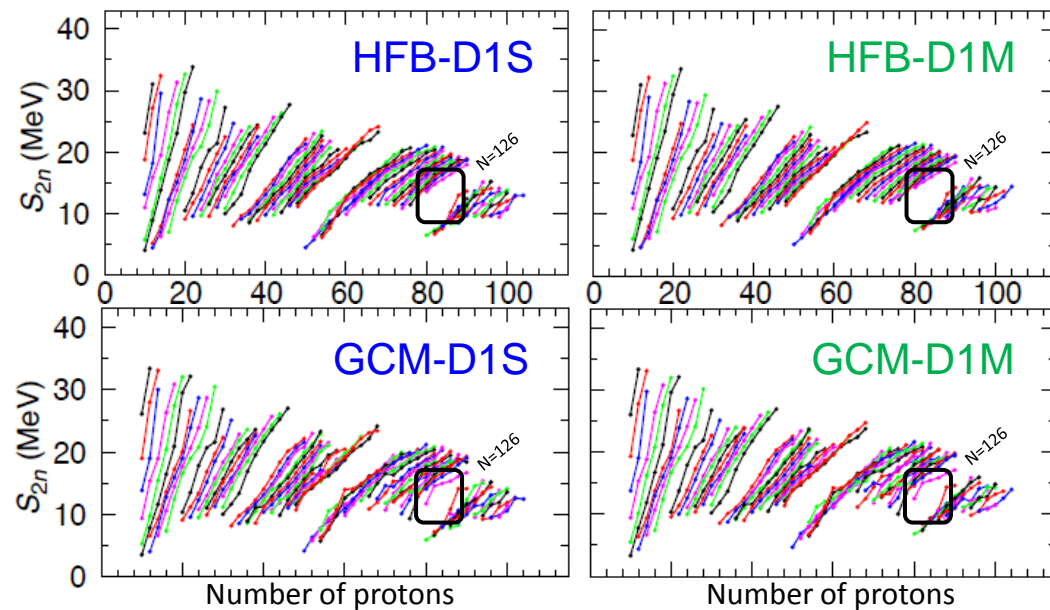
BMF correlations



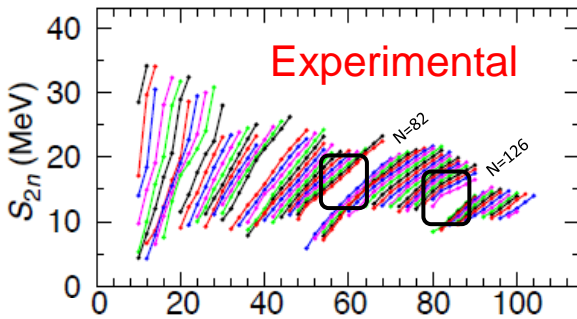
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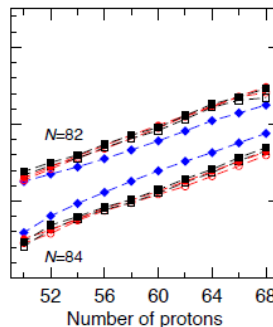
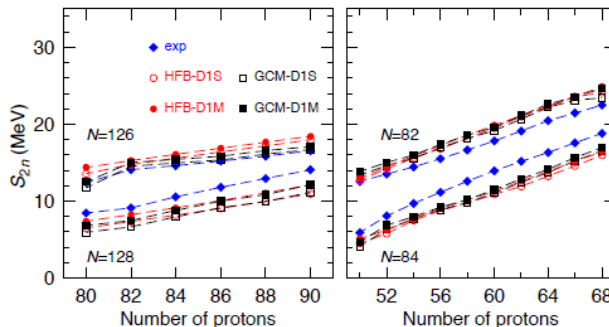
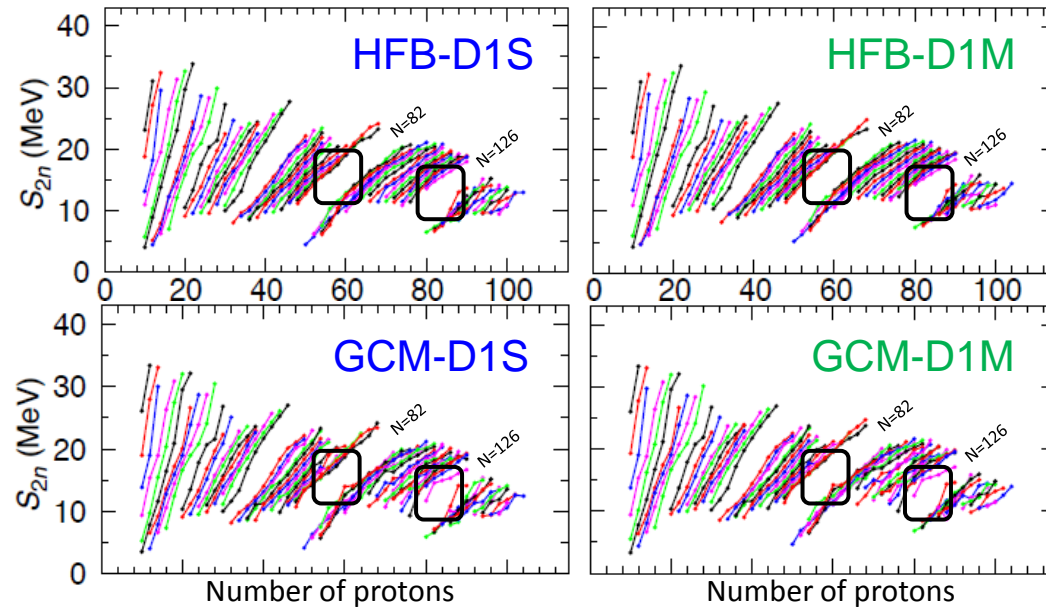


BMF correlations

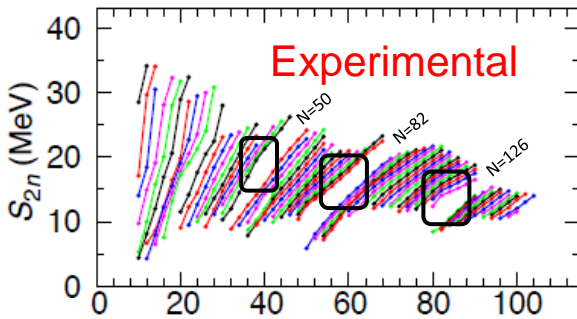


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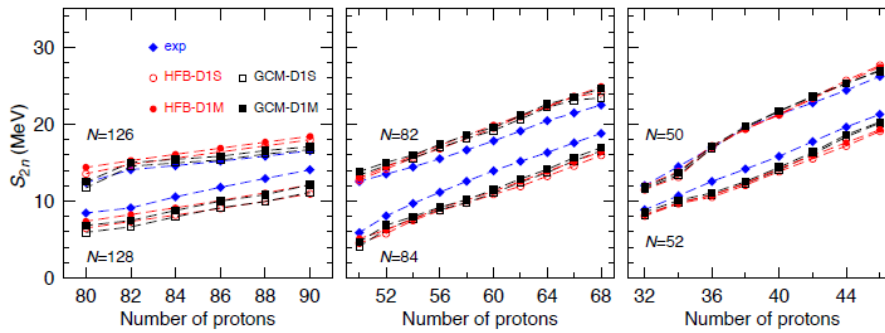
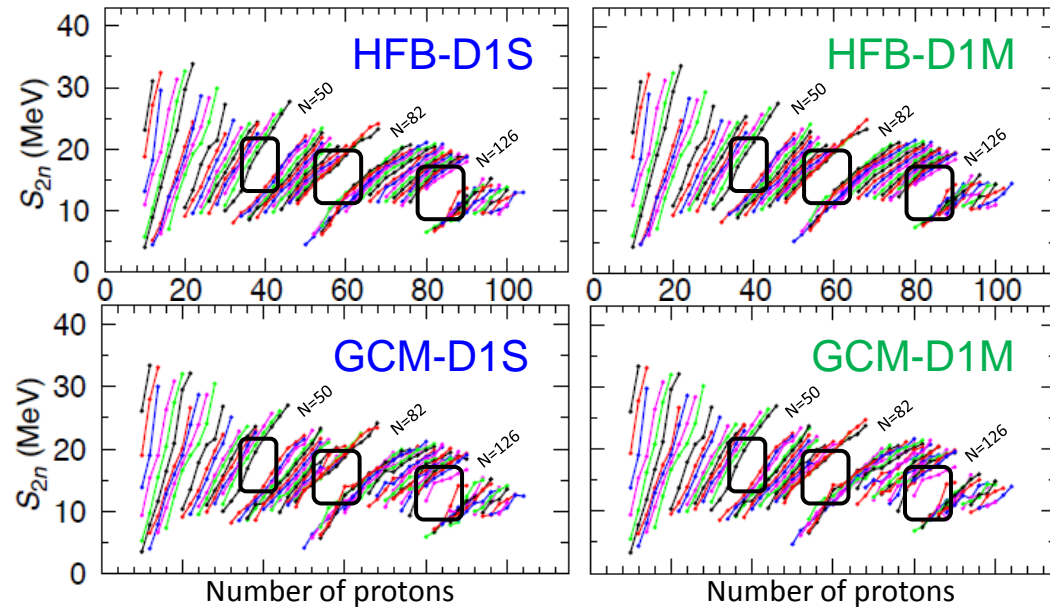


BMF correlations

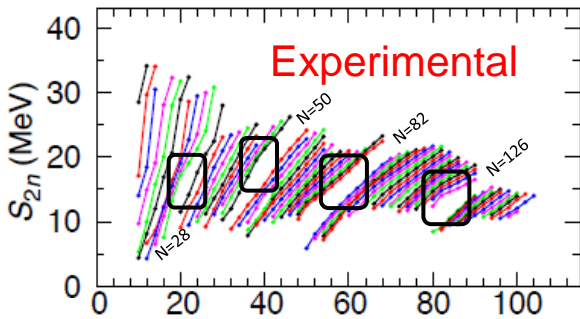


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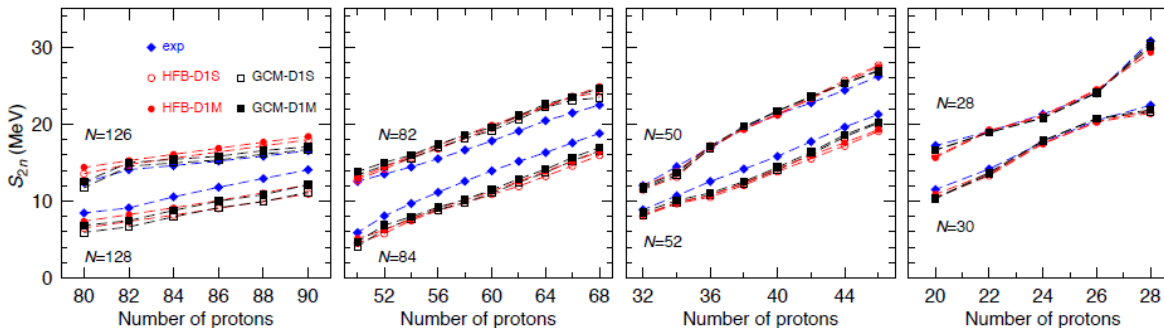
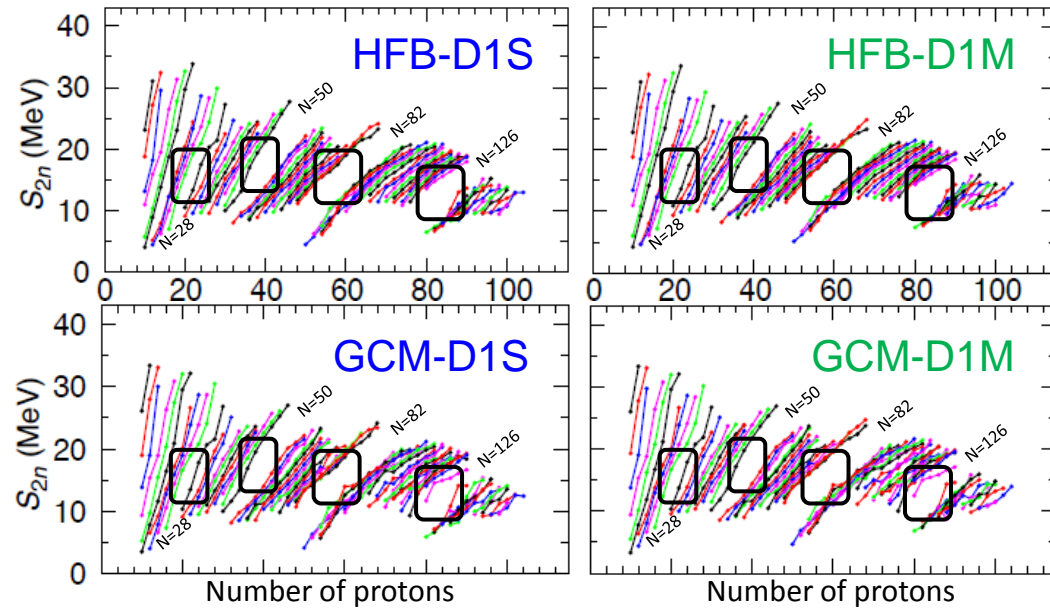


BMF correlations

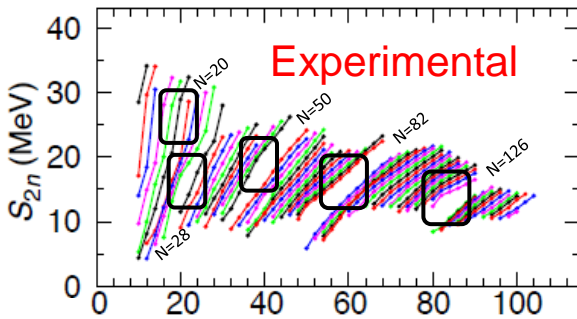


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BMF correlations

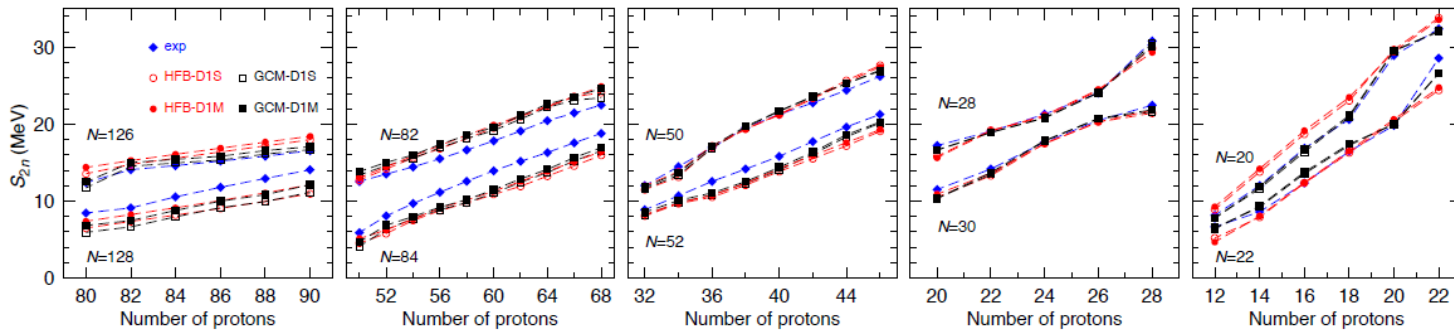
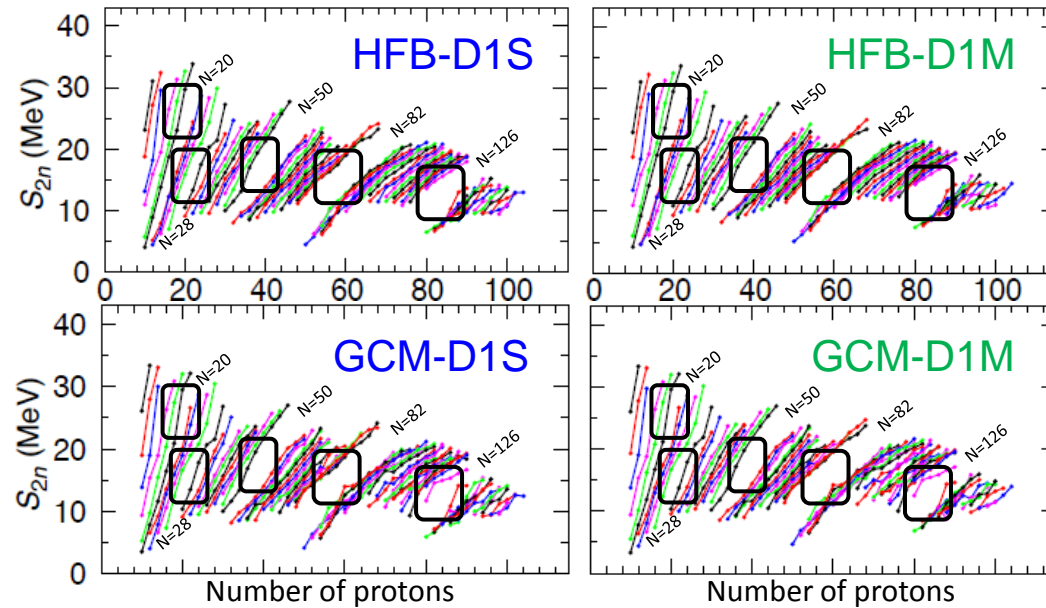


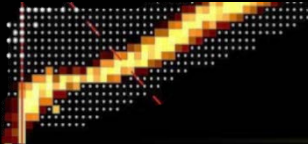
- Experimental S_{2n} are much smoother than both HFB and GCM results:
 - Convergence problem?
 - Missing triaxiality, octupolarity, etc.?

- BMF corrections tend to reduce the exaggerated shell gaps of HFB
- ... but the reduction is not as strong as in the similar surveys with BMF methods

Bender M, et al., *PRC78*, 054312, 2008 [GOA]
Delaroche J-P, et al., *PRC81*, 014303, 2010 [5DCH]

- Could be related to approximated BMF techniques used there (GOA, 5DCH)



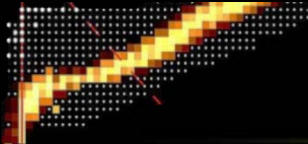


Summary and Outlook

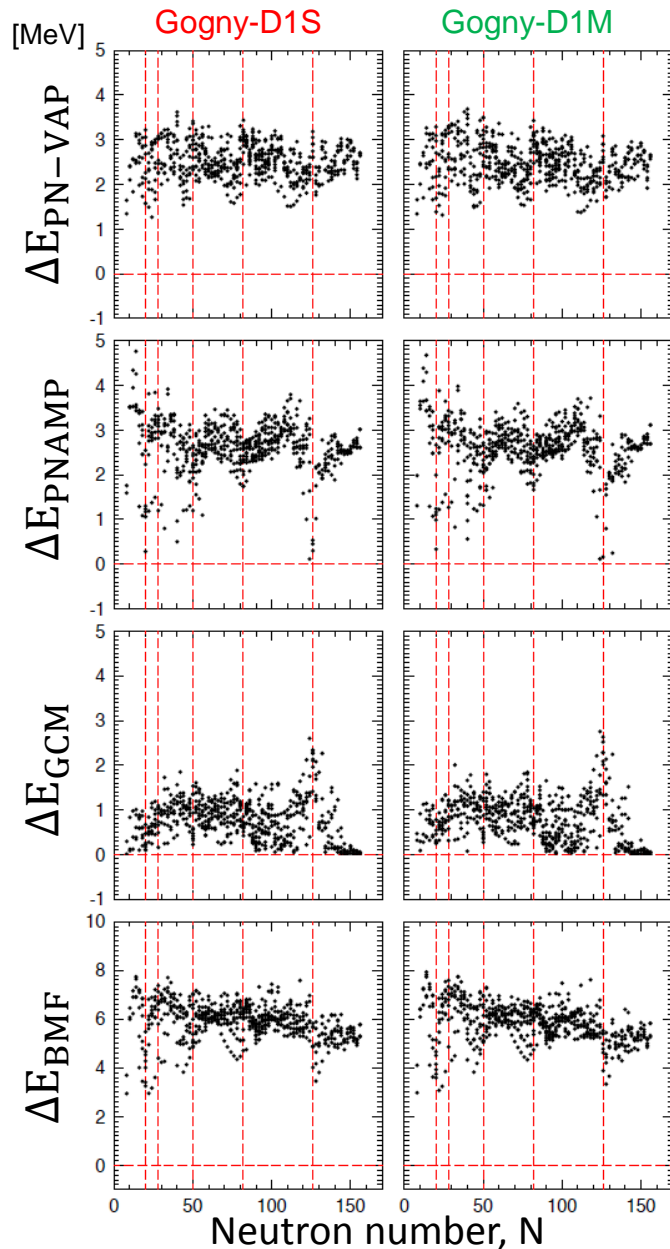
Summary and Outlook

- Despite that this global BMF-calculation with much improved convergence and GCM treatment is still far from precision level of other sophisticated mass formulas, this is the right step towards the microscopic global nuclear structure model that is reliably applicable to neutron-rich r-process nuclei.
- Additional degrees of freedom (*e.g.* triaxiality, particle-vibration coupling, octupole deformations) must be included explicitly to improve description of both spectral and ground state energies.
- Further investigation of odd-nuclei approximation techniques, or implementation of explicit time-reversal breaking is needed.
- Particular attention must be paid to the convergence properties of the harmonic oscillator working basis.
- Finally, a significant improvement is to be made from a new EDF parametrization tuned to include the relevant BMF effects.

τ_{END}



Additional Slides



Symmetry restoration by

- **Variation After Particle Number Projection (PN-VAP):**

$$\Delta E_{\text{PN-VAP}} \sim 2.3 \text{ MeV}$$

- **Particle Number and Angular Momentum Projection (PNAMP):**

$$\Delta E_{\text{PNAMP}} \sim 2.7 \text{ MeV}$$

Configuration mixing by

- Exact implementation of

Generator Coordinate Method (GCM):

$$\Delta E_{\text{GCM}} \sim 0.8 \text{ MeV}$$

Total Energy with BMF correlations

$$E = E_{\text{HFB}}(N_{\text{OS}} = 19) - \Delta E_{\text{BMF}}$$

where the BMF correlations are calculated as

$$\Delta E_{\text{BMF}} = E_{\text{HFB}}(N_{\text{OS}} = 11) - E_{\text{BMF}}(N_{\text{OS}} = 11)$$

because

$$t_{\text{BMF}}(N_{\text{OS}} = 11) \approx 60\text{h}$$

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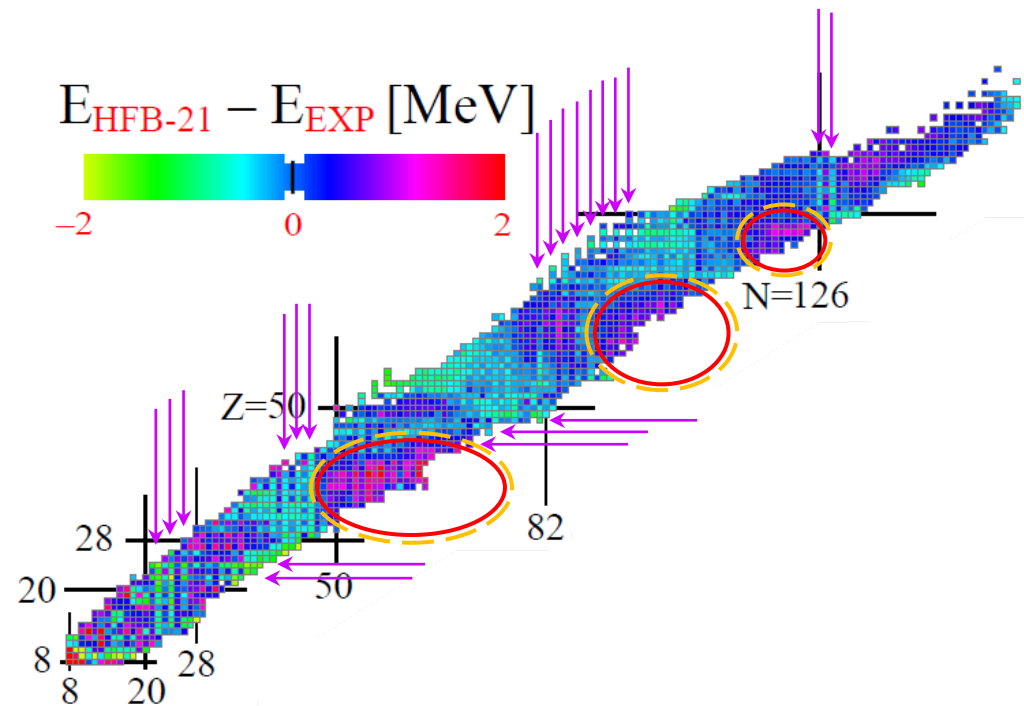
- Astrophysical models require as an input thousands of **nuclear masses** beyond experimental reach
- Need accurate predictions from theoretical **global mass models**
- All models have similar **rms**, but more fundamental **HFB models** *should provide* greater confidence in describing unknown isotopes

- **Still some problems with HFB models:**

I) **Lack of convergence** due to truncated model space

II) Missing some physics without **Beyond-Mean-Field** correlations

III) **Odd-mass nuclei** are not treated on the same footing



Outlook

- ▶ Construct a **complete mass table** by including odd-mass nuclei
- ▶ **Explore more degrees of freedom**
(triaxiality, particle-vibration coupling, octupole deformations, *etc.*)
- ▶ **New energy density functional parametrization**
adjusted to the Beyond-Mean-Field effects