

Strong Interactions under Extreme Conditions: a Review

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HGS-HIRe for FAIR
Helmholtz Graduate School for Hadron and Ion Research

- ▶ Motivation to study QCD under extreme conditions
- ▶ What we know already and how we know it
- ▶ What we do and don't know: Equation of State
- ▶ What we do and don't know: Dynamics
- ▶ Areas of active research

QCD is theory to describe protons+neutrons and other *hadrons* in terms of their constituents, *quarks and gluons*.

We know *degrees of freedom* and the *Lagrangian*

$$\mathcal{L} = \frac{1}{2g_s^2} \text{Tr} G_{\mu\nu} G^{\mu\nu} + \sum_{i=uds c b t} \bar{\psi}_i (\not{D} + m_i) \psi_i$$

We know the (scale-dependent) parameters g_s^2 and $m_{uds c b t}$ accurately

We know the *bound-state spectrum* ($\pi, K, \eta, \rho, p, n, \Delta, \dots$) from both experiment and theory

We know the *high-energy behavior* (jets and their properties) from a mix of experiment and theory

Few-body scattering at intermediate energies
(Experiments: pp , $pD \rightarrow pn$, $p\pi$, pK , but not $\pi\pi$, KK , ...)

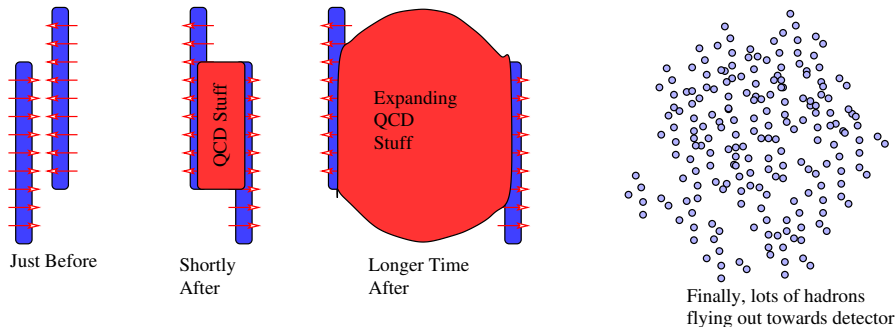
large coupling, difficult for first-principles tools like lattice QCD

Many-body physics and collective dynamics

- ▶ Equilibrium: Equation of state $P(T, \mu_B, \mu_I, \mu_S)$
- ▶ Nonequilibrium: dynamics of large assemblages of QCD matter with lots of energy
 - Heavy ion collisions
 - Neutron star mergers
 - Early Universe

I will concentrate on these questions

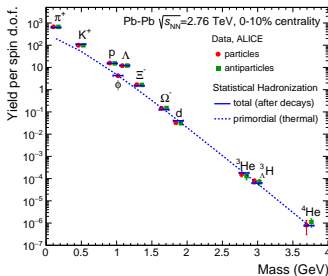
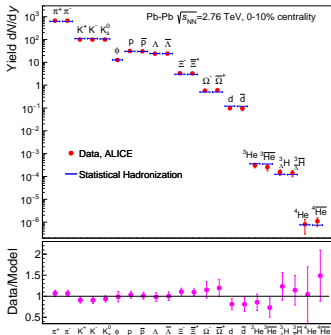
Collide (largest available) nuclei together at (largest available) high energy



Some complex dynamics lead to final-state expanding ball of hadrons.

If system in equilibrium at temp T , chem potential μ , volume V

- ▶ Thermal statistical distribution $1/(\exp((-μ_i + \sqrt{p^2 + m_i^2})/T) \mp 1)$
- ▶ Include unstable particles, but decay them (feeddown)
- ▶ Vary T, μ_B, V to fit particle yield data

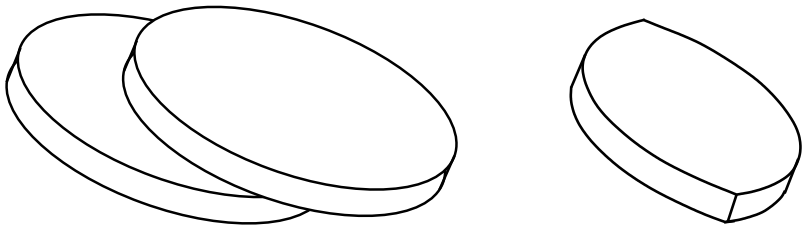


Andronic
Braun-Munzinger
Redlich
Stachel
arXiv:2101.05747

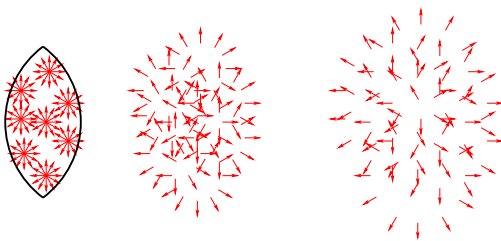
Works very well. (For pp or e^+e^- it does not.)

Nuclear collisions generally not perfectly central

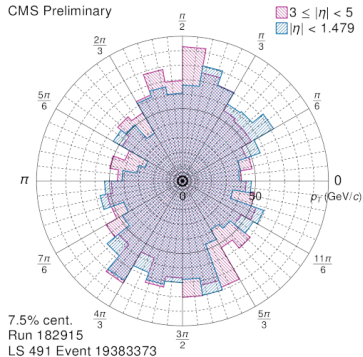
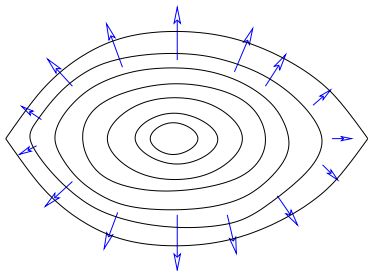
leading to



Without re-interactions: final state isotropic.

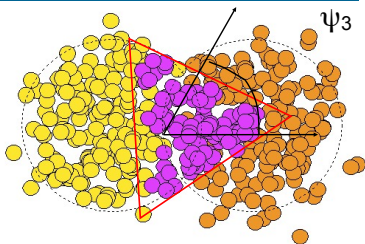


Strong reinteraction+fluid behavior:
Pressure contours imply system will “blow out” to the sides.
Coincides with appearance of real events (CMS):



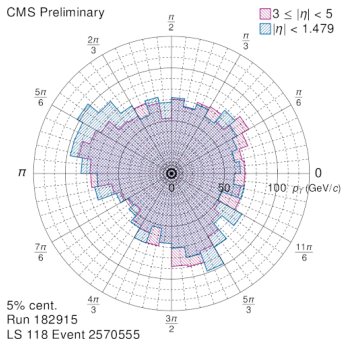
Geometry of collision:
Elongated (elliptical) region
Random fluctuations:
triangular/square/pentagon....

Alver+Roland [arxiv.org:1003.0194](https://arxiv.org/abs/1003.0194)

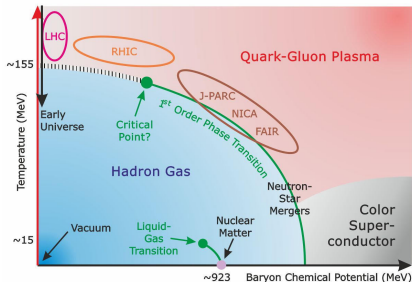


Some events clearly display this pattern.
Fluctuations in geometry are important.
system must be near-equilibrium to turn
the geometry into flow in this way.

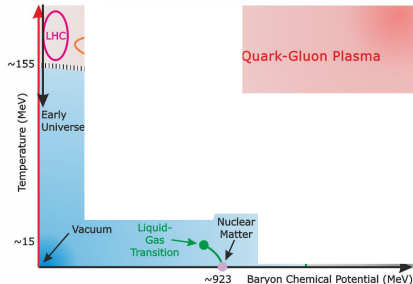
Makes sense to ask about thermodynamics



Thermodynamics means Phase Diagram + Equation of State.
So what do we expect – and what do we know for sure?

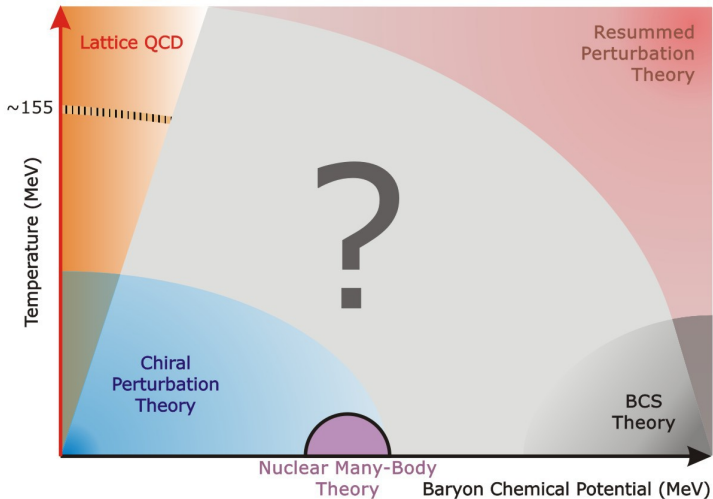


What we expect



What we know for sure

What first-principles tools do we have, and where do they work?



Lagrangian from second slide

$$\mathcal{L} = \frac{1}{2g_s^2} \text{Tr} G_{\mu\nu} G^{\mu\nu} + \sum_{i=udsctb} \bar{\psi}_i (\not{D} + m_i) \psi_i$$

Use this in the *Path Integral*

$$Z = e^{-iHt} = \int \mathcal{D}(G_\mu, \bar{\psi}, \psi) \exp \left(i \int d^4x \mathcal{L} \right)$$

Problems: Integral involves phases, phase cancellations

Stationary phases works *if* g^2 small.

But g^2 is scale-dependent, large at long distances

Perturbation theory expands about stationary phases –

works if all relevant scales are short distance (large T and/or μ)

For thermodynamics, sufficient to explore

$$Z_{\text{therm}} = e^{-H/T} = \int \mathcal{D}(G_\mu, \bar{\psi}, \psi) \exp \left(- \int_0^{1/T} dt \int d^3x \mathcal{H} \right)$$

Now integral converges without phases (if $\mu_B = 0$ anyways).

Thermodynamic information all still there

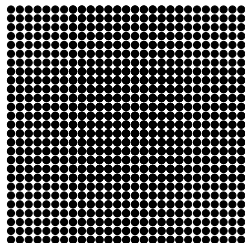
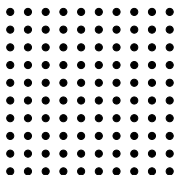
But information about *dynamics* is now all *indirect*.

But it's one integration variable

- ▶ per direction G_μ (4)
- ▶ per color component G_μ^a (4×8)
- ▶ per point in spacetime $G_\mu^a(x)$ ($4 \times 8 \times \infty^4$)

Replace continuous spacetime
with finite-volume lattice

$$\text{Answer} = \text{Correct} + \mathcal{O}(a^2) \\ + \mathcal{O}(\exp(-m_\pi L))$$



Lattice implementation preserves *gauge symmetry exactly*.

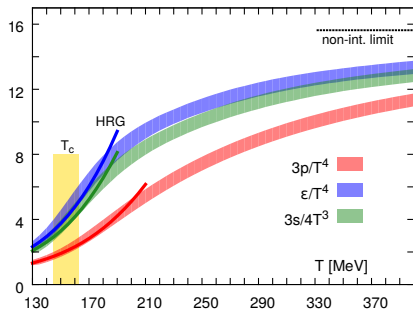
Fermions are trickier, but OK in $a \rightarrow 0$ limit.

Modern simulations: $\sim 1\%$ statistical + systematic errors
for spectrum, thermodynamics ($P, \varepsilon, s, \chi_B, \dots$) etc.

Dynamics (Scattering, transport, etc) **much** harder....

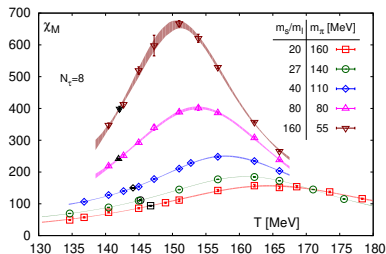
Old-days of quenched approx. or unphysically heavy quarks are over.

Smooth crossover from hadronic to quark-gluon behavior.
 Would be sharper if up, down quarks were lighter (right)



Equation of State

HotQCD arxiv:1407.6387



Chiral susceptibility

HotQCD arxiv:1903.04801

Suppose we want finite density? Need chemical potentials μ .
Lattice techniques depend on importance-sampling

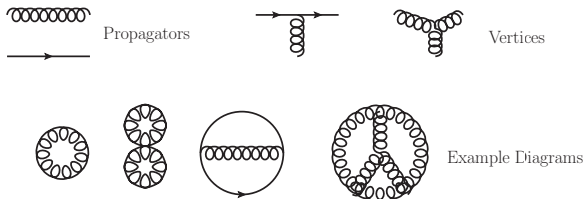
$$Z(\mu_q) = \int \mathcal{D}G_\mu \exp^{-\int d^4x \mathcal{H}_{\text{bos.}}} \prod_{q=udsc} \text{Det} \left(\not{D} + m_q + \mu_q \gamma^0 \right)$$

which requires \mathcal{H}_{bos} and Determinant to be real
Determinant generally complex for $\mu_i \neq 0$.

Perturb in small μ_q : results up to 6'th order.
Sufficient for a "slice" of μ values but not $\mu \gg T$.

Alternative techniques not yet quantitatively reliable at physical m_q .

Perturbation theory:
assume propagation
dominates interaction
do Loop Expansion



Convergence very poor without resummation techniques.

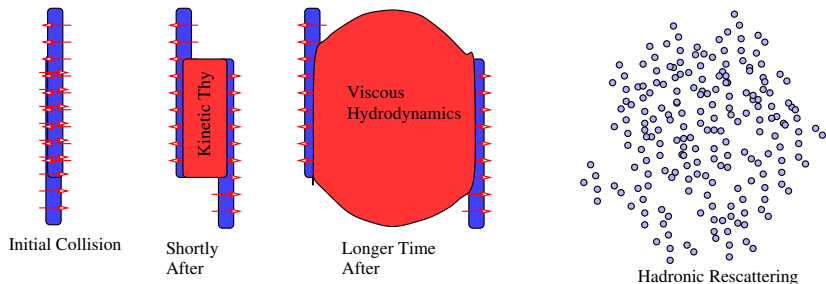
With resummations, work for $T > 350$ MeV.

“Low” T, n_B ($T < 100$ MeV, $\mu \lesssim m_p$): works in terms of (π, n, p)

Functional Renormalization Group: resummation technique which also allows bound states, condensation, collective dynamics ..

Also challenging to apply at $\mu_q \gg T$.

Total time scale ~ 20 fm long compared to inverse-energies $p_{\perp}^{-1} \sim 0.4$ fm



Series of “epochs” which may need different physics descriptions

- ▶ Initial: statistical Glauber, pQCD, Saturation/Colored Glass
- ▶ Early dynamics: kinetic theory
- ▶ Intermediate times/densities: relativistic viscous Hydro
- ▶ Freezeout and hadronic interactions

Local thermal equilibrium: *conserved charge densities*

$(\varepsilon, \vec{\pi}) = T^{\mu 0}$, n_B, ρ_e, n_s determine everything.

Stress tensor $T^{\mu\nu} = \varepsilon u^\mu u^\nu + P \eta^{\mu\nu}$, Currents $j_{B,e,s}^\mu = n_{B,e,s} u^\mu$

Their evolution determined by conservation equations

$$\partial_\mu T^{\mu\nu} = 0 \quad \partial_\mu j_{B,e,s}^\mu = 0$$

Together with equation of state $P = P(\varepsilon, n_{B,e,s})$, equations close.

Near equilibrium, derivative corrections

$$T^{\mu\nu} = T_{\text{eq}}^{\mu\nu} - \eta \left(\nabla^\mu u^\nu + \nabla^\nu u^\mu - \frac{2}{3} \Delta^{\mu\nu} \nabla \cdot u \right) - \zeta \Delta^{\mu\nu} \nabla \cdot u$$

Nice introduction: [Teaney arXiv:0905.2433](#); [Schäfer and Teaney arXiv:0904.3107](#)

Agnostic about Degrees of Freedom. Works for

- ▶ weakly-coupled quarks + gluons
- ▶ strongly-coupled “stuff” with no particle description
- ▶ dense gas of hadrons

provided that mean free path \ll system age.

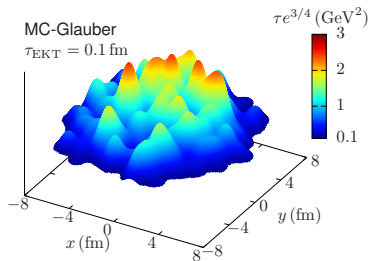
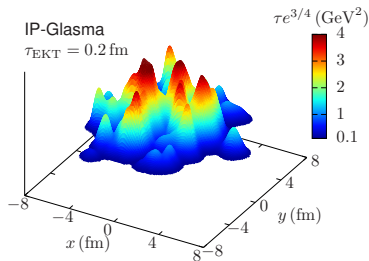
Large systems: this works. (Gold/Lead at RHIC/LHC)

Small systems: we'll see how/where it breaks down!

- ▶ Input required: Stress tensor from initial collision
- ▶ Output: energy+momentum density \Rightarrow temperature+flow velocity
- ▶ Output: local final anisotropy

Convert into hadrons and feed into final-state scattering model

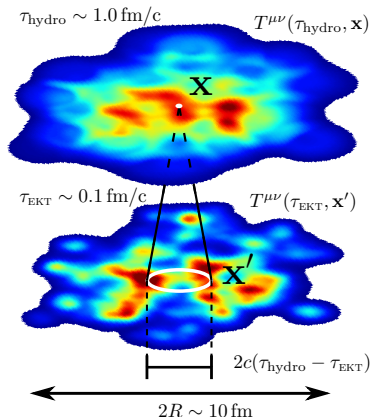
If Hydrodynamics applies at intermediate times,
All we need is the initial **Stress Tensor** $T^{\mu\nu}$ (early).



Different models on market, eg, IPGlasma (left), TRENTo (right)
Shown: 2D slice of initial energy density, to be fed into hydro

See eg Kurkela *et al* [arXiv:1805.00961](https://arxiv.org/abs/1805.00961)

Which model is “right” (or “best”)? Let the data decide.



Early dynamics propagate $T^{\mu\nu}$ a short distance, smoothing initial features...

Kurkela *et al* [arxiv:1805.01604](https://arxiv.org/abs/1805.01604)

Including this dynamics weakens dependence on “thermalization time” which may be ~ 1.5 Fermi/ c .

Old approach: end Hydrodynamics abruptly at “freezeout time”
Switch from perfectly-strong scattering to free streaming.

Modern approach: end Hydrodynamics when system is hadronic.
Switch to hadronic description but keep tracking hadronic scattering
At time of switch, hadronic description approximately hydrodynamic
Insensitive to exact time+details of switch

Each phase has unknowns. Roughly

- ▶ Initial/early: how much energy density, how clumpy?
- ▶ Hydrodynamics: Transport coefficients η, ζ, σ as functions of (T, n_B)
- ▶ Transition to hadrons: Transition temperature. How $T^{\mu\nu} - T_{\text{equil}}^{\mu\nu}$ shared between high/low momentum particles
- ▶ Final rescattering: details of hadronic interactions

Accommodated with *fitting parameters* which are varied to reproduce data

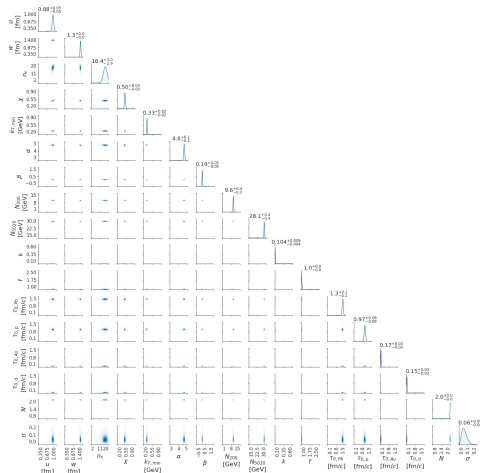
Old approach: experts on one aspect would vary fitting parameters in *their* stage of the process, holding other stages fixed.

Problem: changing, say, η/s changes best-fit initial energy density.

Choices in other stages influence best-fit in practitioner's stage.

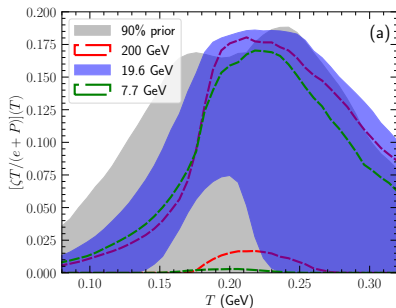
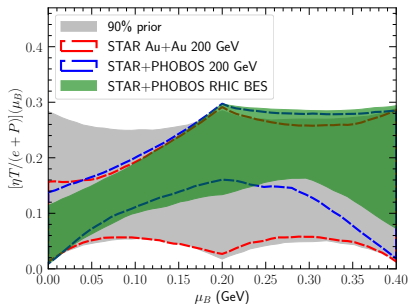
Collaboration with expertise in all phases **JETSCAPE**, or [arXiv:2306.08665](https://arxiv.org/abs/2306.08665)
 Code which interfaces output of one phase to input of the next
 Fit *all* data while varying *all* parameters in grand-fit Bayesian analysis

Parameter	Symbol	Range
Form width [fm]	u	0.35 – 1.0
Nucleon width [fm]	w	0.35 – 1.0
Constituent number	n_c	2.0 – 20.0
Structure	χ	0.2 – 0.9
Transverse mom. scale	$k_{T,\min}$	0.2 – 0.9
Shape parameter	α	3.0 – 5.0
Shape parameter	β	-0.5 – 1.5
Fireball norm. [GeV] ($\sqrt{s_{NN}} = 200$ GeV)	N_{200}	1.0 – 15.0
Fireball norm. [GeV] ($\sqrt{s_{NN}} = 5.02$ TeV)	N_{5020}	15.0 – 30.0
Fluctuation	k	0.1 – 0.6
Flatness	f	1.0 – 2.5
Hydrodyn. time [fm/c] (Pb-Pb 5.02 TeV)	$\tau_{0,Pb}$	0.1 – 1.5
Hydrodyn. time [fm/c] (p-Pb 5.02 TeV)	$\tau_{0,p}$	0.1 – 1.5
Hydrodyn. time [fm/c] (Au-Au 200 GeV)	$\tau_{0,Au}$	0.1 – 1.5
Hydrodyn. time [fm/c] (d-Au 200 GeV)	$\tau_{0,d}$	0.1 – 1.5
Overall scale	N_{scale}	0.8 – 2.0



Old goal: find η/s of Quark-Gluon Plasma.

New goal: find T, μ_B dependence of $\eta/s, \zeta/s$



Shen Schenke Zhao arXiv:2310.10787

Initial collision also produces

- ▶ Heavy quarks
- ▶ High-energy quarks and gluons \Rightarrow jets

Initial production (mostly) under control based on pp ...

Propagation through medium modifies energy distribution

Opportunity to learn medium properties / predict modification

Coupling large enough that many interactions nonperturbative

Makes quantitative first-principles predictions challenging

First-principles calculations already do a good job on

- ▶ Equation of state at $\mu_B/T < \pi$ from the lattice
- ▶ Hydrodynamic description of high-energy heavy-ion systems
- ▶ Extraction of η/s etc from data using Bayesian framework
- ▶ First-principles dynamics at extreme (unachievable) temperatures

We are making progress but there is much work to be done

- ▶ Initial conditions
- ▶ Early-time dynamics
- ▶ Hydrodynamics – or not? for small systems
- ▶ Hadronic reinteractions, late-stage evolution
- ▶ Hard probes – modification of heavy quarks and high-energy jets
- ▶ Equation of state at $T \sim 100$ MeV, $\mu_B \sim 1$ GeV
- ▶ Transport from the lattice
- ▶ Learning Equation of State from neutron stars + NS mergers

- ▶ Heavy ion collisions give a window into many-body Quantum Chromodynamics
- ▶ Thermodynamics at small μ_B : first principles lattice methods. Mature field, results well under control
- ▶ Thermodynamics at larger μ : theoretically much harder. Phase diagram partly conjecture
- ▶ Dynamics: Heavy Ion Collisions progress in stages: Initial conditions, early evolution, near-equilibrium/hydro, final-state hadrons
- ▶ Substantial progress in last 20+ years, with more to come!