Studying the Gluonic GFFs at JLab

RICHARD TYSON



J/ψ Near-Threshold Photoproduction

- Employ an electron beam for either:
 - Quasi-real photoproduction where a virtual photon mediates the interaction at Q^2 close to 0.
 - Real photoproduction where a real photon produced via bremsstrahlung interacts with the target
- J/ ψ decays to a lepton pair, either e^+e^- or $\mu^+\mu^-$. Other decay channels are OZI suppressed.
- Produce J/ψ close to its 8.2 GeV photoproduction threshold.
- Aim to measure:
 - > Total cross section as a function of E_{γ} .
 - Differential cross section as a function of t.



J/ψ quasi-real photoproduction on a proton target

Probing the Mechanical Properties of the Nucleon

- The mechanical properties of the nucleon are encoded by Gravitational Form Factors (GFFs) [1] defined from the matrix elements of the energy-momentum tensor (EMT).
- Any spin-2 field couples to the EMT and gives rise to a force indistinguishable from gravity [2].
- The quark GFFs have already been investigated in the context of DVCS. This led to estimates of sheer and pressure distributions of quarks inside the proton [3,4].
- A two-gluon exchange forms a spin-2 coupling between J/ψ and the nucleon. This allows to estimate the gluonic GFFs [5-7].



Spin-2 fields in graviton-proton scattering and DVCS.

J/ψ and the GFFs

- VMD based models relate J/ψ photoproduction to the J/ψ-nucleon scattering. The J/ψ-nucleon scattering amplitude gives access to the EMT [5,9].
- In holographic QCD a higher dimensional duality relates spin-2 fields to gravity. J/ψ is produced by the exchange of gravitons (tensor 2++glueballs) and scalar (0++) glueballs [6,10].
- In the GPD framework, large skewness at threshold allows to relate the scattering amplitude to gluon GPDs. The GFFs are extracted from the first moments of the GPDs [7,11,12].



The gluon contribution to the pressure distribution inside the proton from a GPD based model fit to lattice and J/ψ photoproduction data [7].

JLab

- The Thomas Jefferson National Accelerator Facility (JLab) is located in Newport News, Virginia.
- The Continuous Electron Beam Accelerator Facility (CEBAF) produces a 12 GeV electron beam.
- Upcoming Solenoidal Large Intensity Device (SoLID) will be located in Hall A.
- The CEBAF Large Acceptance Spectrometer (CLAS12) is located in Hall B.
- The J/ ψ 007 Collaboration located in Hall C.
- The GLUonic Excitation Experiment (GlueX) is located in Hall D.



SoLID in Hall A

- SoLID is a future facility planned for the 2030s.
- High luminosity ~ 10^{37} cm⁻²s⁻¹.

6

- Large acceptance with full azimuthal coverage and polar angles from ~7 to ~24 degrees.
- Good momentum resolution ~2% from 0.8 to 7 GeV (for electrons).
- Planned J/ψ experiment will utilize a 30µA beam on 15 cm LH2 target to reach counts of ~800k J/ψ (in photoproduction).
- Increased statistics at large t will be crucial for the extraction of the gluonic GFFs [11,12].





The CLAS12 Detector

Beam energies up to 11 GeV are delivered to Hall B.

The Forward Detector has polar angle coverage of 5 to 35 degrees.

The Central Detector has polar angle coverage of 35 to 125 degrees.

Both have full azimuthal coverage.



J/ψ at CLAS12

- CLAS12 took data with a liquid hydrogen target and a liquid deuteron target.
 - Shown here is the differential cross section measured in:
 - ► $en_{bound} \rightarrow (e')e^+e^-n$
 - ▶ $ep_{bound} \rightarrow (e')e^+e^-p$
 - $\blacktriangleright ep \rightarrow (e')\mu^+\mu^-p$

 These measurements will soon be updated with an improvement in efficiency of ~25% per charged particles.



Differential Cross Section [AU]

Hall C

Beam energies up to 11 GeV are delivered to Hall C.

- Hall C employs a high precision spectrometer.
- This allows for precise measurements with low background at specific kinematic points.



J/ψ at Hall C

- The J/ψ 007 Collaboration has made high precision measurements of the differential cross section as a function of t [13].
- The main objective was to probe the mechanical properties of the nucleon.
- Measurements in the decay channel of J/ψ to two muons are underway.



as a function of t energy and theoretical predictions [13].

GlueX

Beam energies up to 12 GeV are delivered to Hall D.

- GlueX produces a bremsstrahlung photon beam using the diamond wafer.
- Electrons are tagged to determine the photon energy.
- Full azimuthal coverage with polar angular coverage from ~1 to 120 degrees.



J/ψ at GlueX

- GlueX has made high precision measurements of the J/ψ total cross section [14,15].
- The differential cross section as a function of t was also measured [14,15].



Measurements of the J/ ψ total cross section as a function of the photon beam energy [15].

J/w at GlueX

- The low t differential cross section is consistent with a dominant gluon exchange production mechanism.
- Structure in the total and differential cross sections are consistent with contributions from other production mechanisms.



Measurements of the J/ψ differential cross section as a function of –t in bins of photon beam energy [15].

Gluonic GFFs

The $A_g(t)$ and $D_g(t)$ GFFs were estimated at Hall-C from J/ ψ photoproduction [13].

Relate to:

- Momentum fraction of gluons in the nucleon.
- Shear forces and pressure distribution of the gluonic content of the nucleon.
- The discrepancy between GPD and holographic models is now resolved [11].



The $A_g(t)$ and $D_g(t)$ GFFs estimated using holographic QCD [10] (orange) and GPD [7] (green) models compared to lattice QCD predictions [16] (blue). $(k^2 \equiv |t|)$

Trace Anomaly Contribution to the Nucleon Mass

The nucleon mass can be decomposed into the contributions from the quark masses, the energy of quarks and gluons and the trace anomaly contribution [17].

Estimates of the magnitude of the trace anomaly contribution to the proton mass were obtained from GlueX and Hall C data [13,14,17].



Mass Radius of the Nucleon

A scalar gravitational form factor G(t) gives access to the mass radius of the nucleon [9].

$$\frac{d\sigma}{dt} = G(t)$$

Assuming a dipole form for G(t): $G(t) = \left(\frac{M_p}{(1 - \frac{t}{m_s^2})^2}\right)^2$

The mass radius r_m is calculated from the free parameter m_s : $\sqrt{12}\hbar c$

 $m_{\rm s}$

 r_m



J/ψ differential cross section as a function of -t. Data from the GlueX Collaboration [14], plot taken from [9].

Mass Radius of the Nucleon at JLab

The mass radius of the proton has been measured from GlueX and Hall C data [13,14].

A larger charge radius than mass radius suggests than the quark radius within the nucleon is larger than the gluon radius.

Deviations at lower photon energies might be indicative of a region where the assumption of two-gluon exchange dominance is invalid.



Two or Three-Gluon Exchange?

- Near threshold, the 3-gluon exchange's contribution to the cross section is expected to dominate that of the 2gluon exchange [8].
- GlueX has already identified a path forward based on an energy upgrade at JLab [18].



Measurements of the J/ψ total cross section as a function of the photon beam energy [14] compared to predictions of the 2- and 3-gluon exchange contribution [8].

Open-Charm Photoproduction?

- J/ ψ near-threshold photoproduction could be dominated by open charm production of $\Lambda^c \overline{D}^{(*)}$ [19].
- JPAC analysis of GlueX and Hall C data suggests a non-negligible contribution from open-charm intermediate states [20].
- Additional data, in particular at high t, would be helpful [19,20]. Measuring $\Lambda^c \overline{D}^{(*)}$ cross sections even more so.
- J/ψ production on the neutron could place further constraint on the open charm production mechanism.



Predictions for the total cross section due to the open charm production of J/ ψ p [19], which is consistent with the GlueX measurements [15] in red.

Key Takeaways

The study of the gluonic GFFs are an active area of research at JLab.

Understanding the J/ψ near-threshold production mechanism(s) will be key in establishing the validity of the estimates of the gluonic GFFs from J/ψ photoproduction.

Need additional data and complementary measurements of open charm channels.

References

[1] H. Pagels, Phys. Rev. 144 (1966).

[2] C. W. Misner, K.S. Thorne, J.A Wheeler, W. H. Freeman 1973 Box 18.1.

[3] V.D. Burkert, L. Elouadrhiri, F.X. Girod , *Nature* **557** 7705 (2018).

[4] V.D. Burkert, L. Elouadrhiri, F.Girod, arXiv:2104.02031 (2021).

[5] D. Kharzeev, H. Satz, A. Syamtomov, and G. Zinovev, Nucl. Phys. A 661 568 (1999).

[6] Y. Hatta, D.-L. Yang, Phys. Rev. D 98 074003 (2018).

[7] Y. Guo, X. Ji, Y. Liu, Phys. Rev. D 103, 096010 (2021).

[8] S.J. Brodsky, E. Chudakov, P. Hoyer, J.M. Laget, *Phys.Lett. B* 498 23 (2001)

[9] D. Kharzeev, Phys. Rev. D 104 054015 (2021).

[10] Kiminad A. Mamo and Ismail Zahed *Phys. Rev. D* **106**, 086004 (2022).

[11] Y. Guo, X. Ji, Y. Liu, J. Yang, Phys. Rev. D 108, 034003 (2023).

[12] Y. Guo, X. Ji, F. Yuan, Phys. Rev. D 109 014014 (2024).

[13] D. Duran, et al. $(J/\psi$ -007 Collaboration), Nature 615 (2023).

[14] A. Ali, et. al. (GlueX Collaboration), Phys. Rev. Lett. **123**, 072001 (2019).

[15] S. Adhikari et al. (GlueX Collaboration) *Phys. Rev. C* **108**, 025201 (2023).

[16] D. A. Pefkou, D. C. Hackett, P. E. Shanahan, *Phys. Rev. D* **105** 054509 (2022).

[17] R. Wang, X. Chen, J. Evslin, Eur. Phys. J. C 80 507 (2020).

[18] L. Pentchev, Presentation available online at https://indico.ectstar.eu/event/152/contributions/3133 /attachments/2001/2612/LPentchev_Jpsi_Trento.pdf.

[19] M.-L. Du, V. Baru, F.-K. Guo, C. Hanhart, U.-G. Meißner, A. Nefediev, I.Strakovsky, Eur. Phys. J. C 80 1053 (2020).

[20] D. Winney, et. al. (Joint Physics Analysis Center), Phys. Rev. D 108, 054018 (2023).

Backup Slides

J/ψ Photoproduction in VMD

- In the VMD picture J/ψ is produced on the proton in the following steps:
 - Photon γ fluctuates into a $c\bar{c}$ pair of size r_{\perp} at a distance l_c
 - $c\bar{c}$ pair scatters off the proton with impact parameter b
 - $c\bar{c}$ forms J/ ψ at some distance l_F after the scattering
 - Near the J/ ψ production threshold, the large mass of J/ ψ imposes a small transverse size r_{\perp} and impact parameter b and a large minimum momentum transfer t_{min} .
- In order to have elastic scattering at threshold, quarks in the nucleon must share the large momentum transfer and must therefore be in a compact Fock state.

$$\frac{d\sigma_{\gamma N \to VN}}{dt} = \kappa \frac{3\Gamma(V \to e^+e^-)}{\alpha_{em}m_V} \frac{d\sigma_{VN \to VN}}{dt}$$



Gluon Exchange

At high energies, leading twist (single gluon) exchange dominates due to a requirement for all involved quarks to be in area of $\frac{1}{m_c^2}$ [8].

At threshold, higher twist (2,3-gluon) exchange dominates due to the large momentum transfer [8].



